

Nobel Symposium on Free Electron Laser Research
14th-18th June, 2015
Sigtunahöjden, Sweden

SACLA and FERMI:
**New opportunities for atomic,
molecular and cluster science**

Kiyoshi Ueda
Tohoku University, Sendai 980-8577, Japan

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SACLA



Masaki

FERMI



Carlo

*Coherent control via phase
controlled multicolor FEL!*

*We appreciate your support for our
SLiT-J (3GeV ring+FEL) project!*

EUV-X FELs in the world



LCLS in operation since 2009



SCSS test accelerator in operation since 2008 ; closed down in 2013

SACLA in operation since March 2012

SCSS+ will start operation in 2016



FERMI starts operation in December 2012!

Swiss FEL (2017), Korean FEL, Shanghai FEL, etc., are coming!

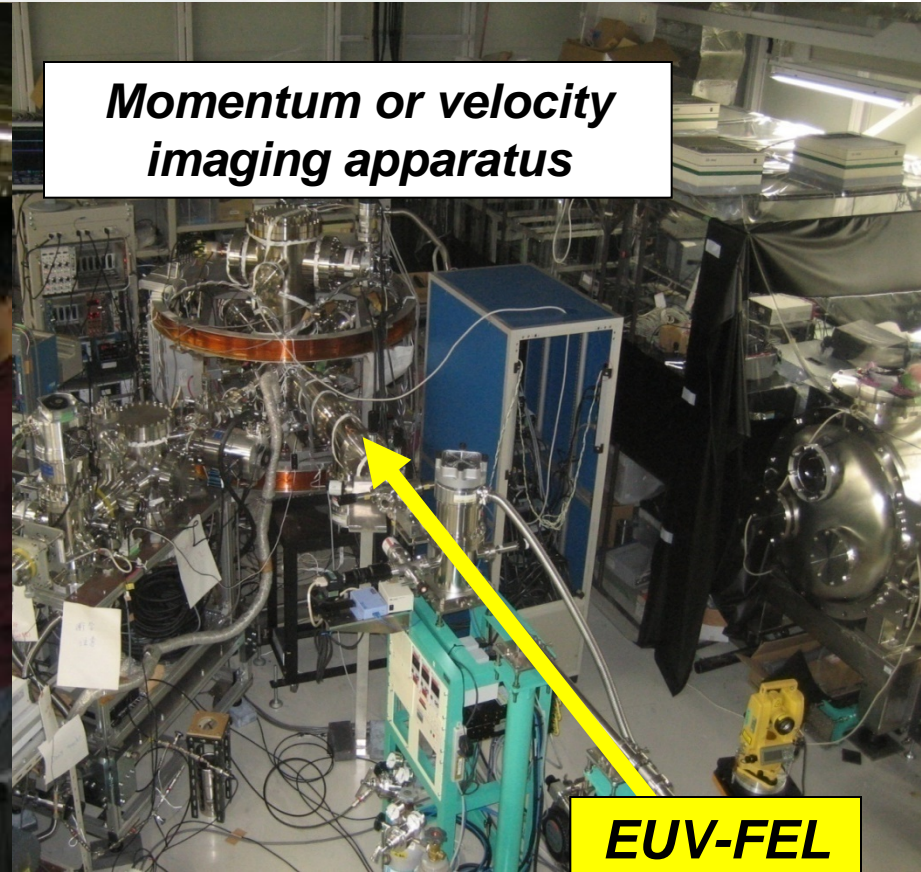


*FLASH in operation since 2005
European XFEL will start operation in 2017*

Outline

- From SCSS (SASE EUVFEL) to FERMI (seeded EUVFEL)
 - Interatomic Colombic decay (ICD) vs nanoplasma formation
 - ICD cascades, inelastic scattering of ICD electrons in neon clusters
 - Two-photon excitation of ICD states and the time-resolved study in neon dimers
 - Coherent control for photoionization of the neon atom by phase-controlled two-color (w-2w) pulses
- to SACLA (SASE hard X-ray FEL)
 - Deep inner-shell multi-photon ionization of Ar and Xe atoms
 - Photoion-photoion coincidence imaging following deep inner-shell multi-photon ionization of CH₃I and 5I-uracil
 - Electron spectroscopy on cold nanoplasma formation from argon, krypton and xenon clusters
 - Single-shot imaging of xenon nano-clusters
 - IR-probe experiment of XFEL-ignited nanoplasma dynamics

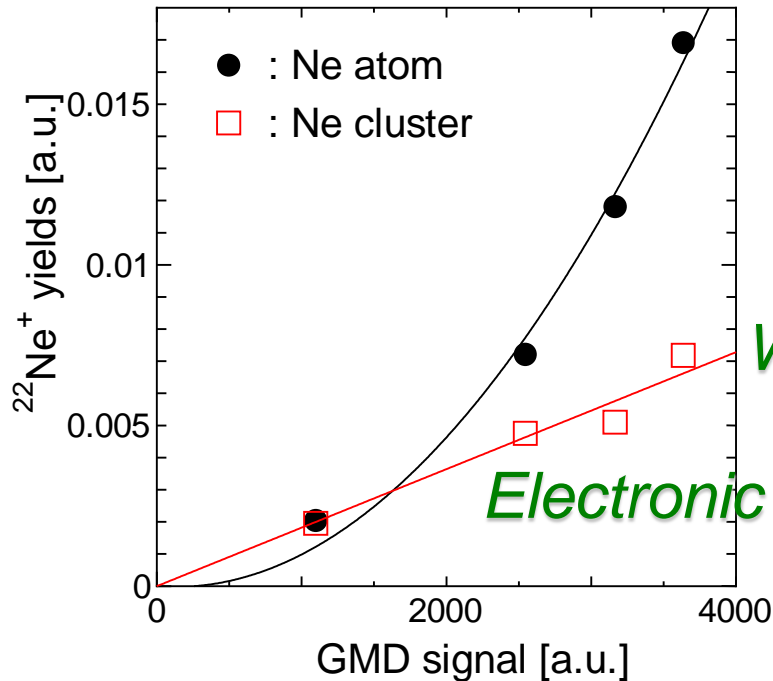
SCSS test accelerator : EUV-FEL (20-24 eV)



EUV-FEL

Multiple ionization of rare gas clusters: with M. Yao's group
EUV-pump EUV-probe: with J. Ullrich's group
Electron spectroscopy with VMI: with M. Vrakking's group

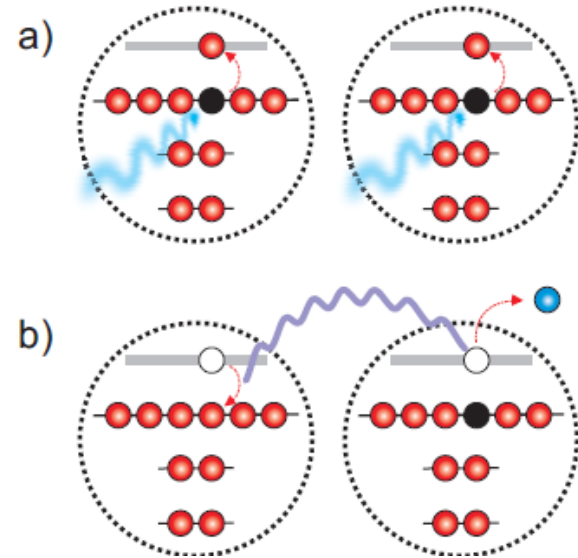
Linear FEL power dependence of energetic Ne⁺ yields



WHY?

Electronic decay ?

Interatomic Coulombic decay
in multiply excited clusters

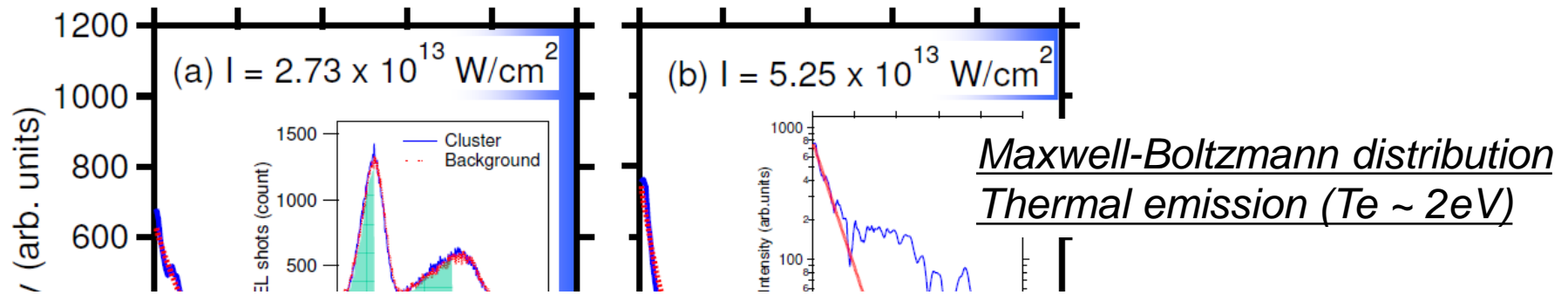


Kuleff et al. PRL 105, 043004 (2010).

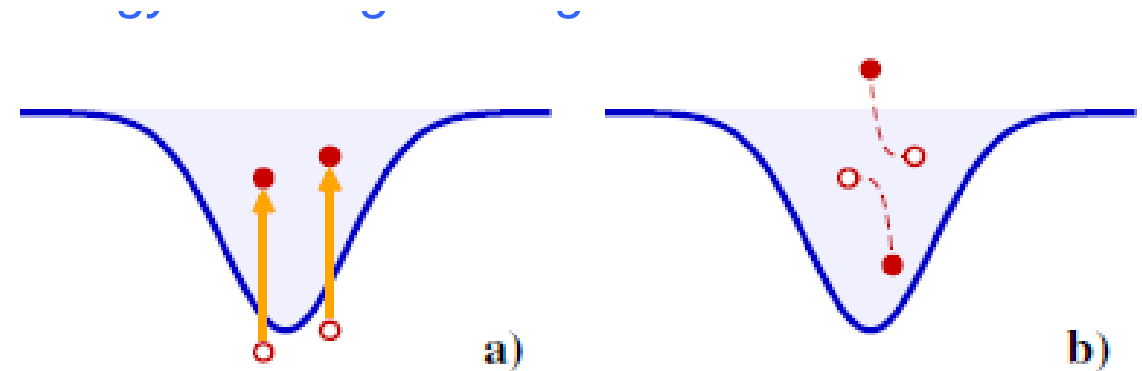
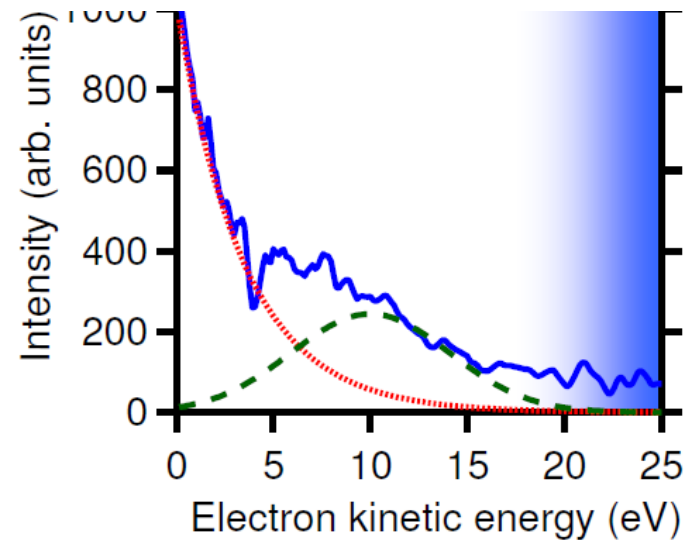
We found linear FEL-power dependence of energetic Ne⁺ yields ejected from clusters ($\langle N \rangle = 1000$), in sharp contrast with quadratic FEL-power dependence of Ne⁺ yields from the atomic beam.

K. Nagaya, A. Sugishima, H. Iwayama, H. Murakami, M. Yao, H. Fukuzawa, X.-J. Liu, K. Motomura, K. Ueda, N. Saito, L. Foucar, A. Rudenko, M. Kurka, K.-U. Kuehnel, J. Ullrich, A. Czasch, R. Doerner, R. Feifel, M. Nagasono, A. Higashiya, M. Yabashi, T. Ishikawa, T. Togashi, H. Kimura, and H. Ohashi, J. Phys. B: At. Mol. Opt. Phys. 46, 164023 (2013)

Thermal emission from Ne clusters



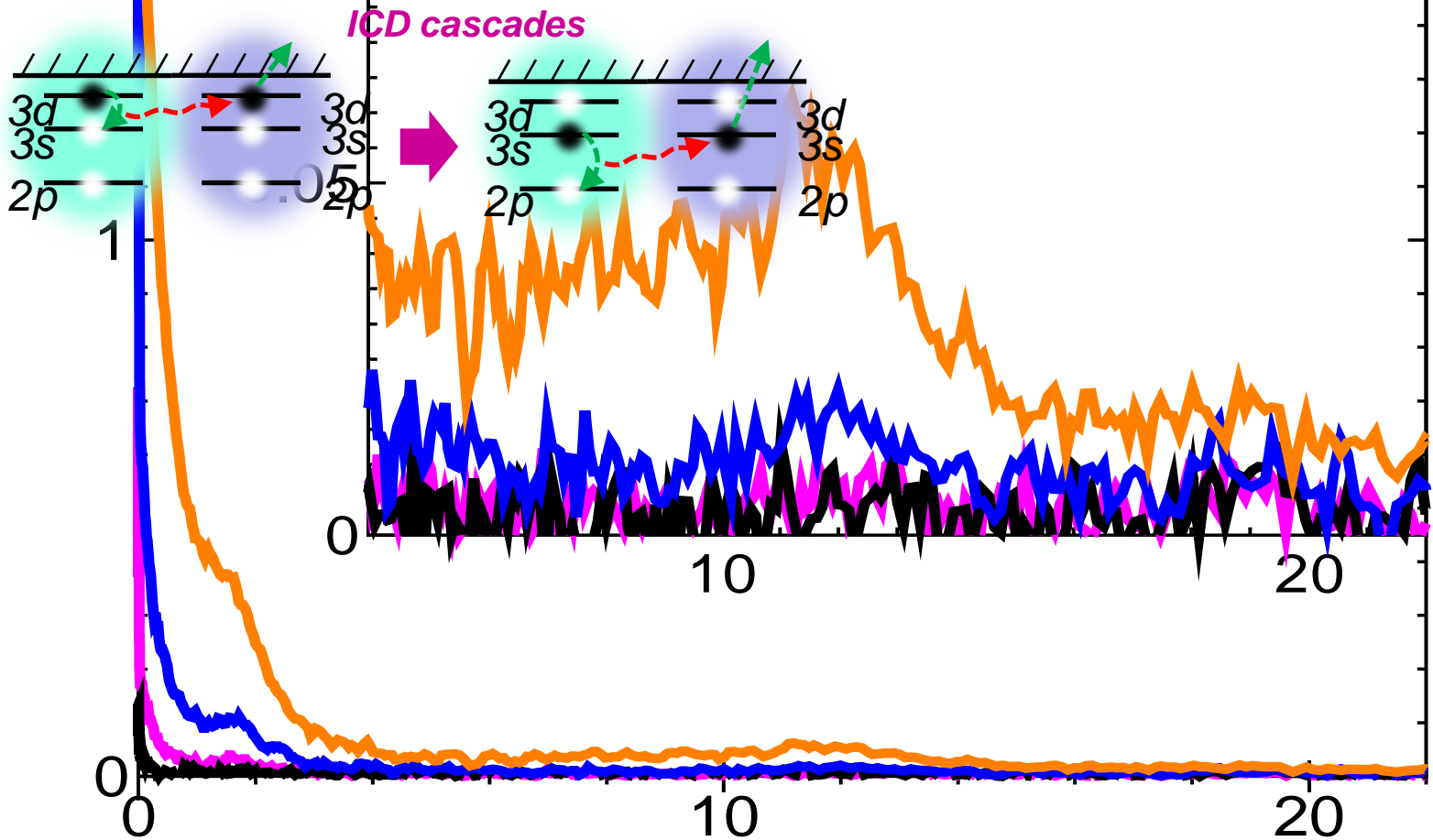
*S. Yase, K. Nagaya, Y. Mizoguchi, M. Yao, H. Fukuzawa, K. Motomura, A. Yamada, Ri Ma, K. Ueda, N. Saito, M. Nagasono, T. Togashi, K. Tono, M. Yabashi, T. Ishikawa, H. Ohashi, and Y. Senba
Phys. Rev. A 88, 043203 (2013).*



*collisions between "free" electrons
Bostedt et al. NJP 12, 083004 (2010).*

Evidence of ICD cascades

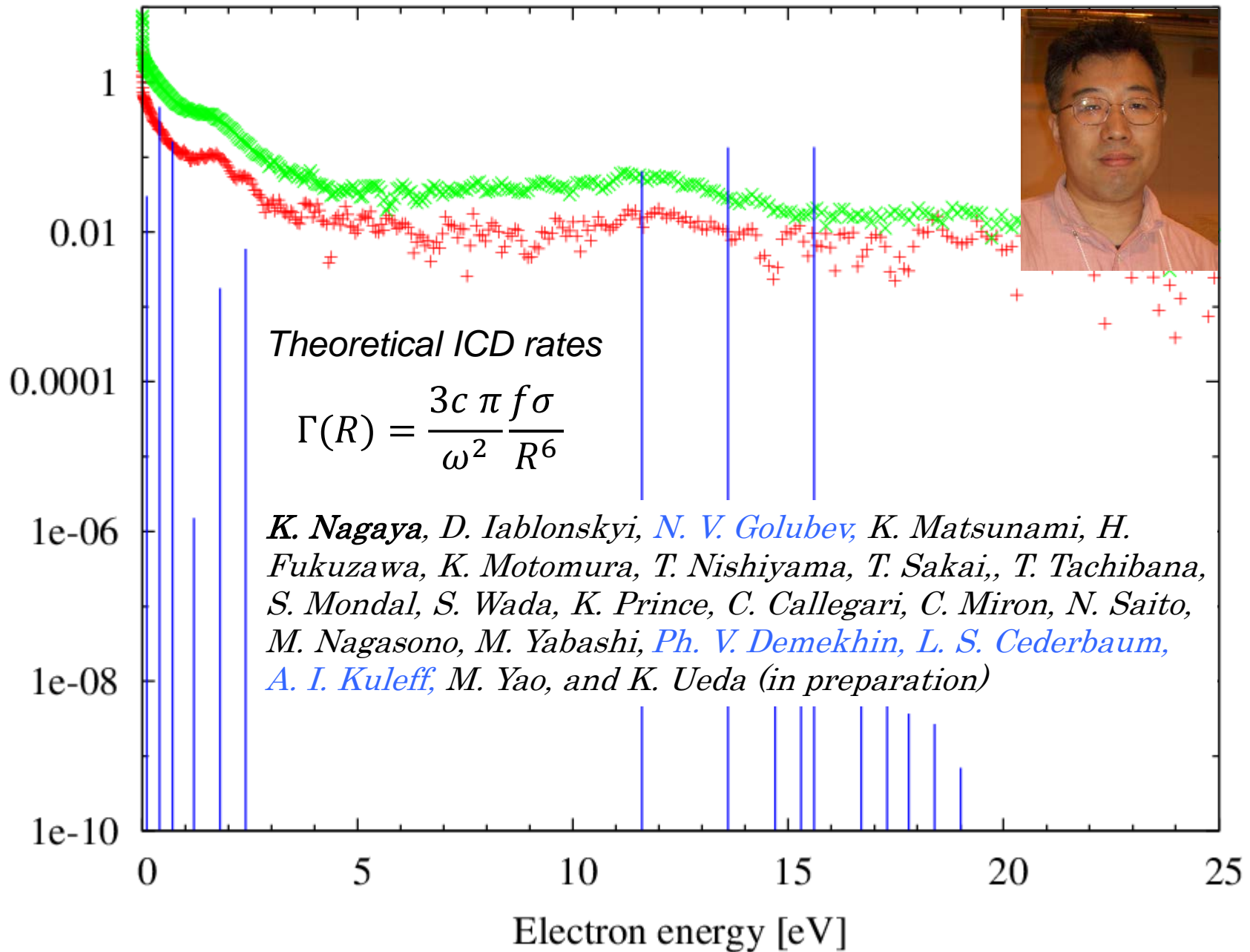
$3d + 3d \rightarrow 3p + e(\text{ICD})$ $3d + 3d \rightarrow 3s + e(\text{ICD})$ $3s + 3s \rightarrow 2p + e(\text{ICD})$ $3d + 3d \rightarrow 2p + e(\text{ICD})$



ICD cascades: experiment and theory



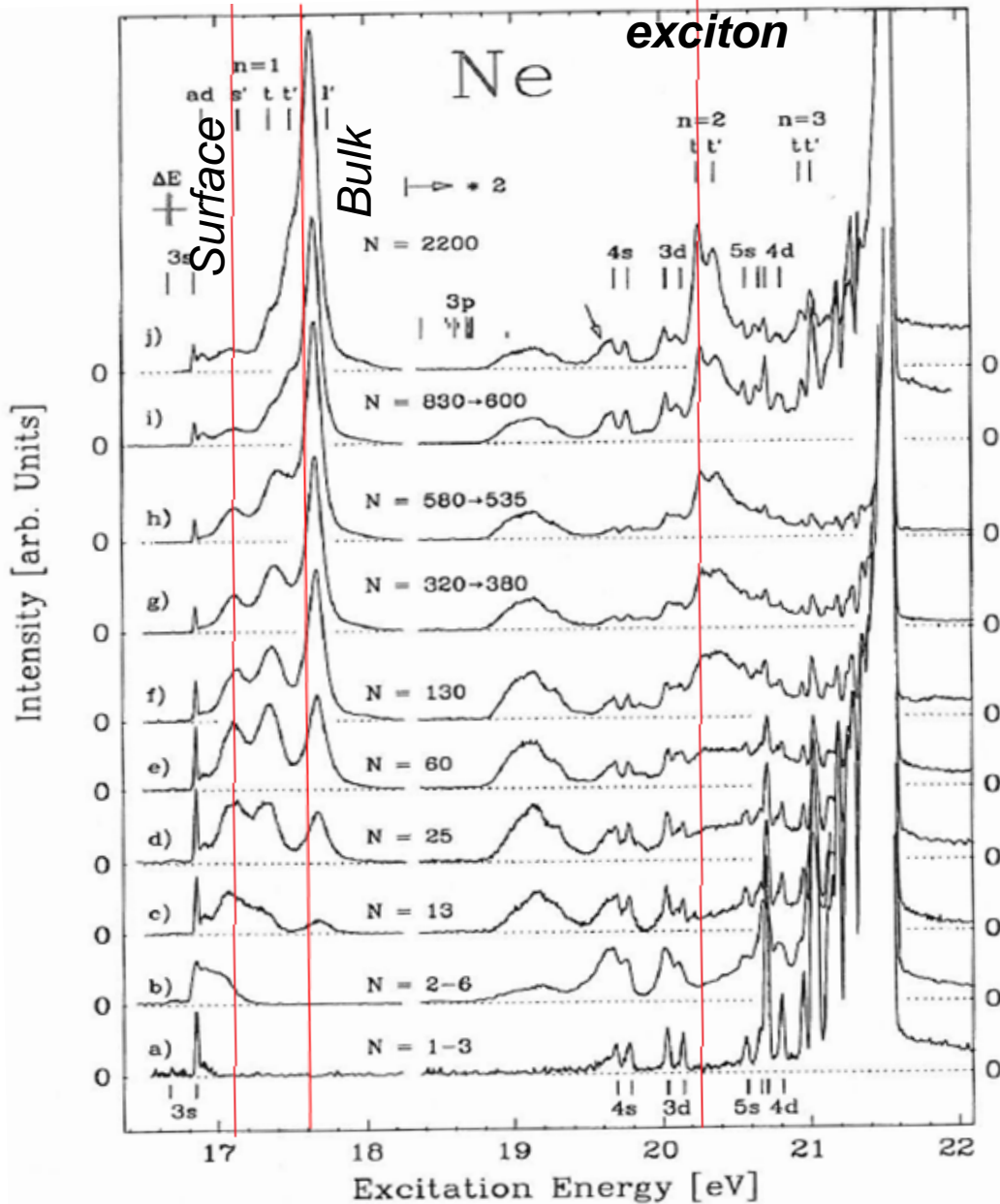
Intensity [arb. units]



Absorption spectra of neon clusters

Frenkel exciton

**Wannier
exciton**



M. Joppien, Ph.D. thesis
Universität Hamburg (1994)

FERMI@ELETTRA: seeded FEL source

Operation for users started in December 2012



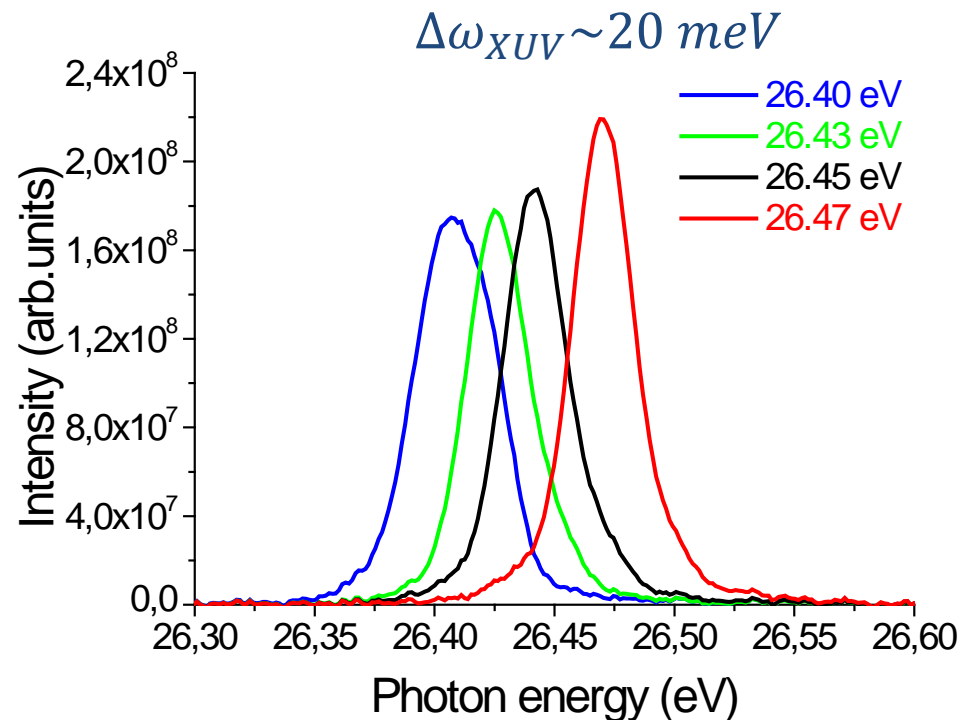
Photon energy range:
19-62 eV

Pulse duration: < 100 fs

Pulse Energy: < 100 μ J



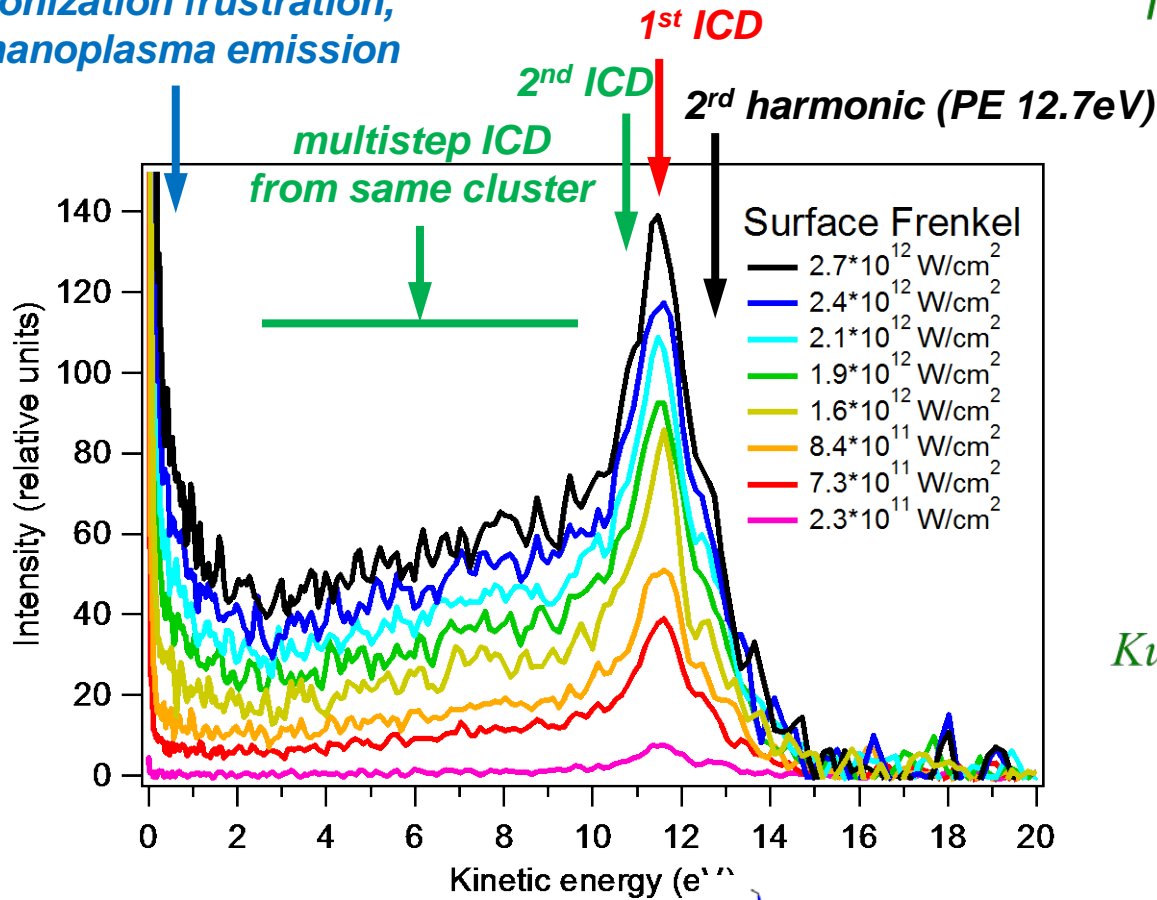
Photon energy tunability



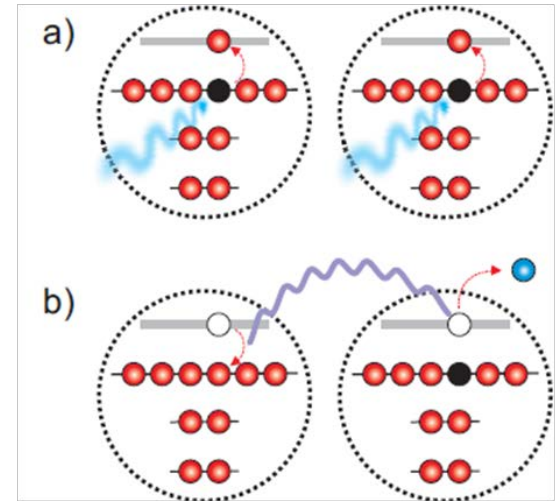
Ideal source for investigating nonlinear resonant effects!

Surface Frenkel exciton ($2p^6 \rightarrow 2p^5 3s$ @17.12eV)

ionization frustration,
nanoplasma emission



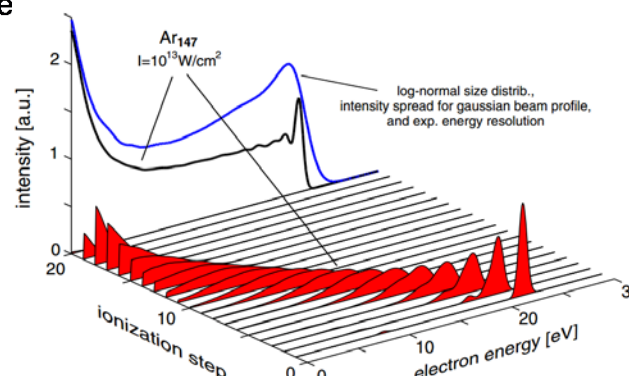
Interatomic Coulombic decay
in multiply excited clusters



Kuleff et al. PRL 105, 043004 (2010).

Electrons are bounded to
the individual atomic ions.

Multistep direct ionization of Ar_{147} clusters



C. Bostedt et al., Phys. Rev. Lett. 100, 133401 (2008)

Bulk Frenkel exciton ($2p^6 \rightarrow 2p^5 3s$ @17.65eV)

*ionization frustration,
nanoplasma emission*

ICD:



Internal inelastic scattering ICD:



(knock-up)

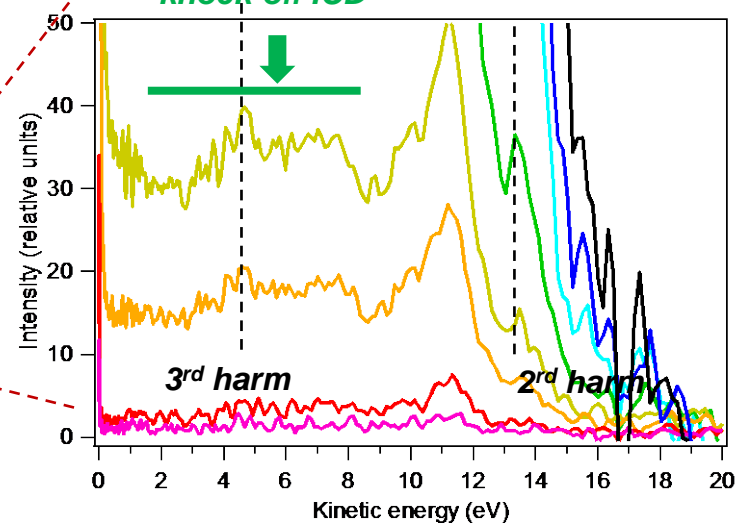
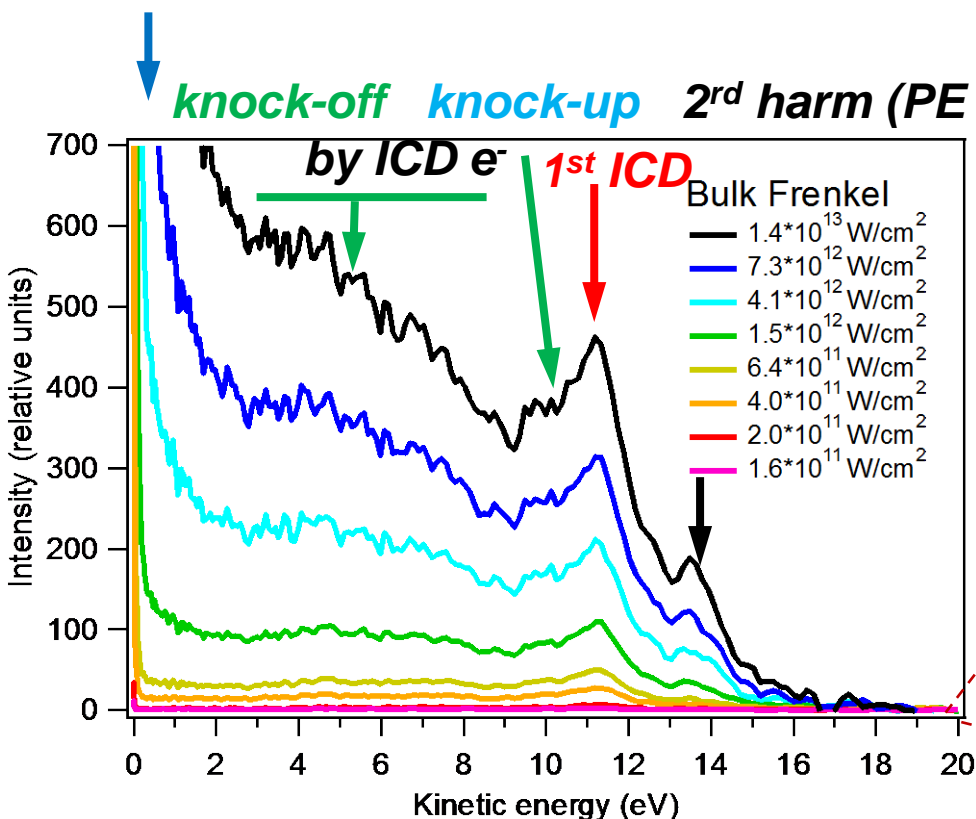


(knock-off)

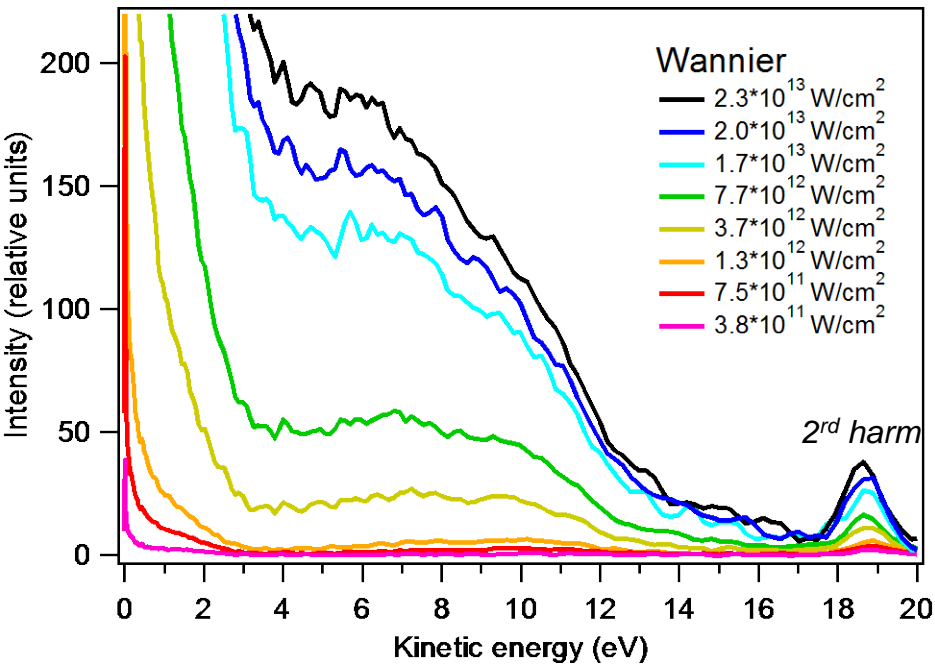
1st ICD

knock-up ICD

knock-off ICD



Wannier exciton ($2p^6 \rightarrow 2p^5 3d$ @20.26eV)



2-body ICD:

$3d+3d \rightarrow 2p+e$ – quenched!

$3d+3d \rightarrow 3s+e$ (~1.3eV) – intra-

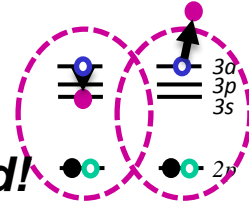
Rydberg ICD

$3s+3s \rightarrow 2p+e$ – cascade ICD

Internal inelastic scattering ICD:

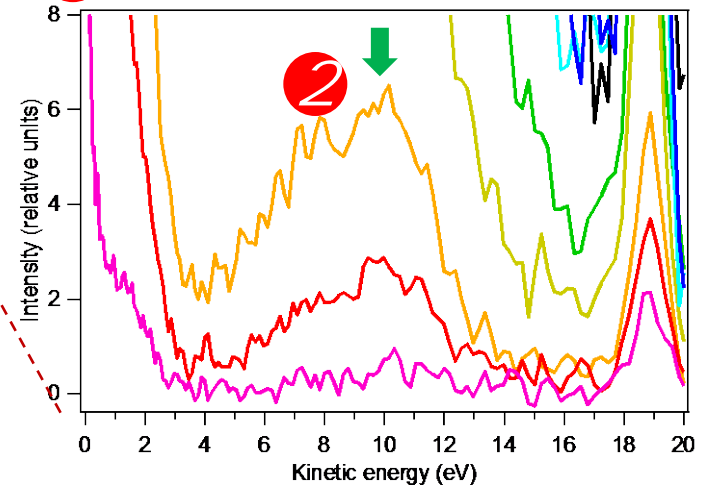
numerous knock-up and knock-off CD channels

intra-Rydberg ICD
(super Coster-Kronig)

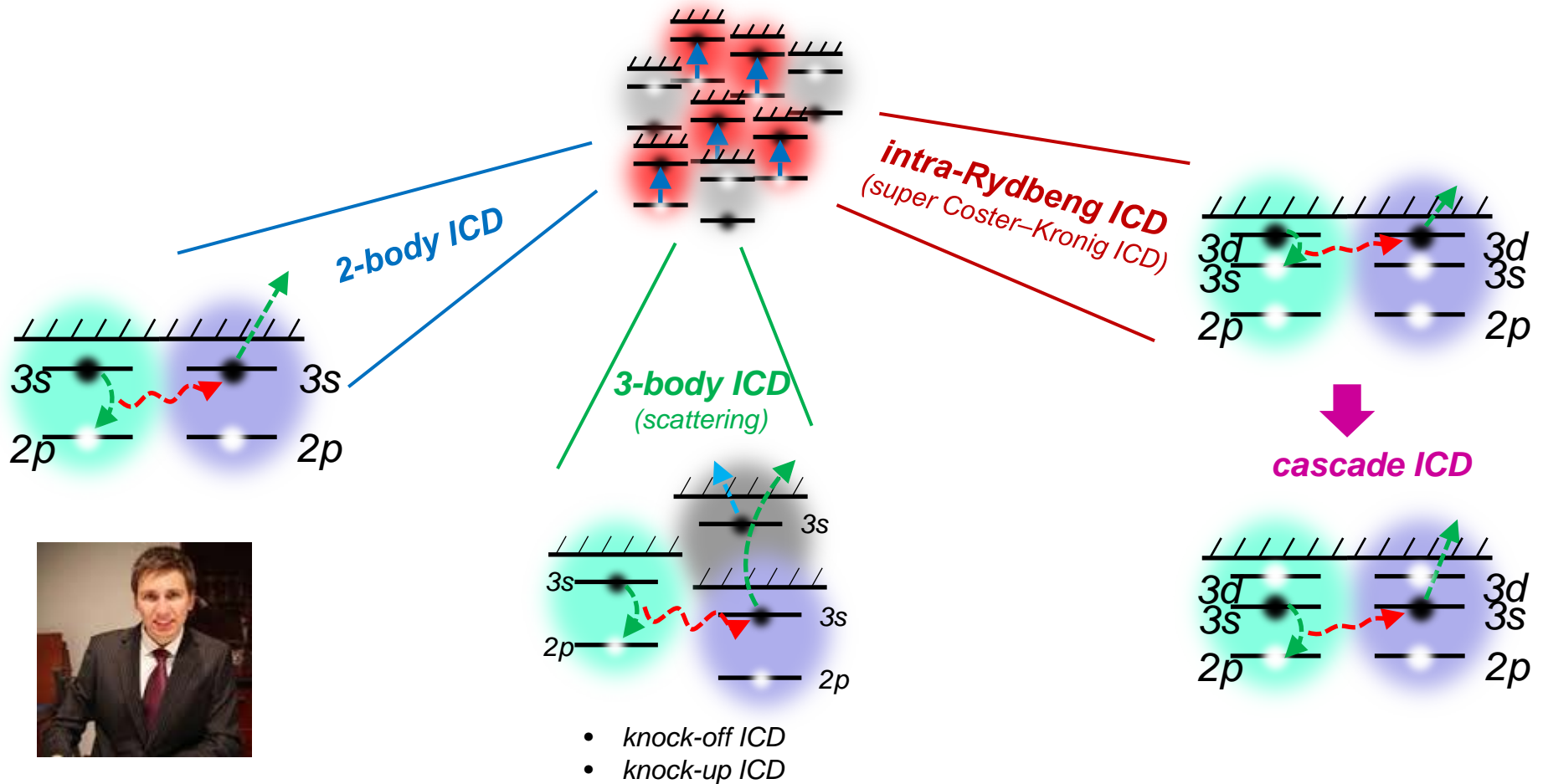


cascade 3s ICD

2nd harm



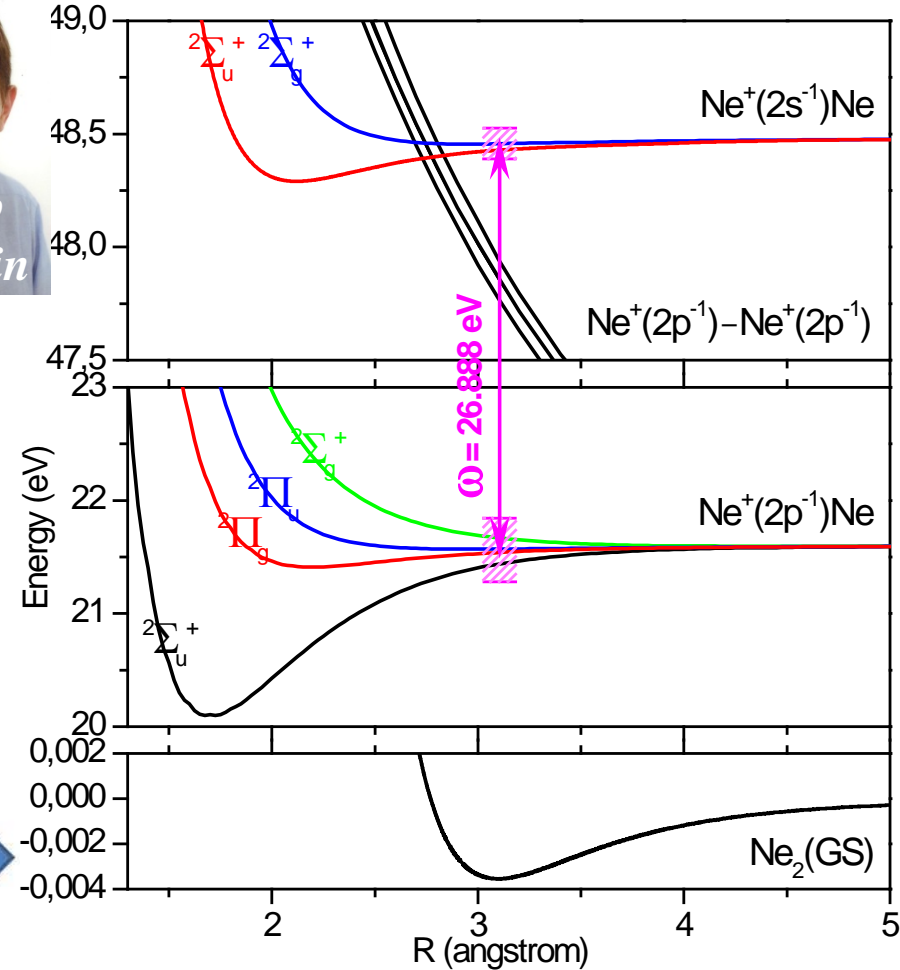
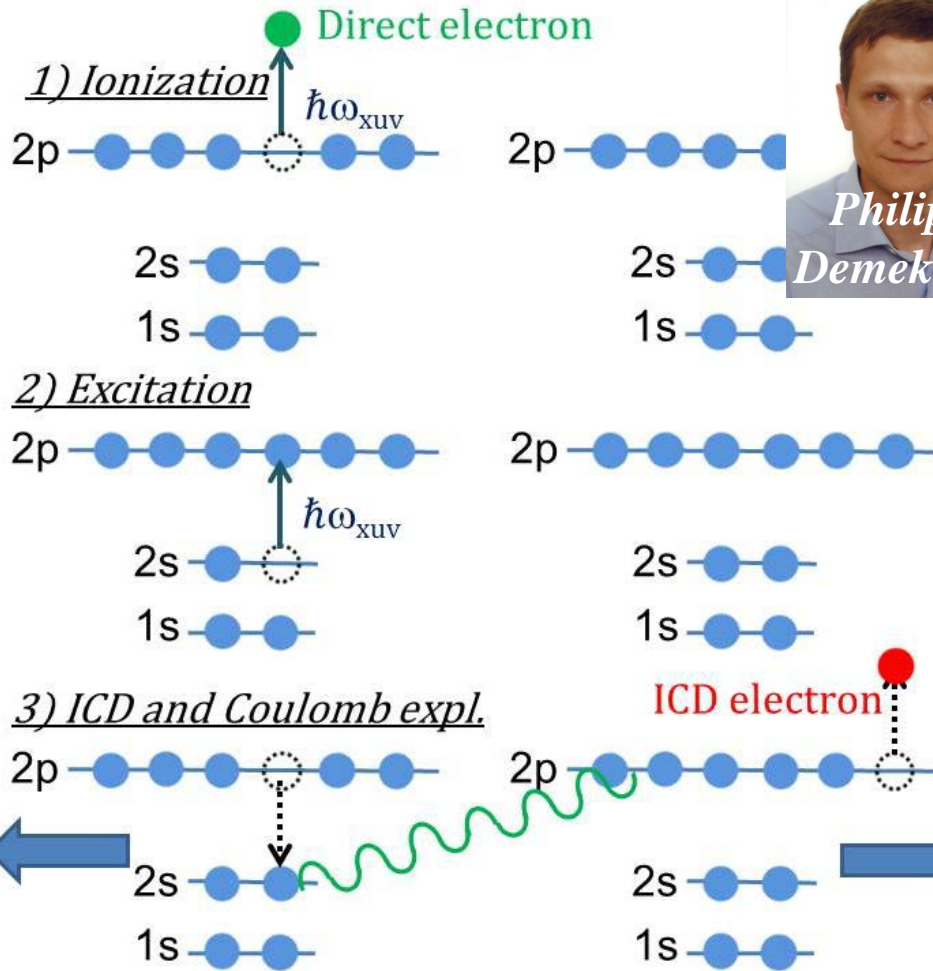
Summary of observed various ICDs



D. Iablonsky, P. Johnsson, T. Takanashi, H. Fukuzawa, K. Motomura, Y. Kumagai, S. Mondal, T. Tachibana, T. Nishiyama, K. Nagaya, P. Piseri, G. Sansone, A. Dubrouil, M. Reduzzi, P. Carpeggiani, C. Vozzi, M. Devetta, M. Negro, D. Faccialà, F. Calegari, A. Trabattoni, M. Castrovilli, Y. Ovcharenko, T. Moeller, M. Mudrich, F. Stienkemeier, M. Coreno, M. Alagia, B. Schütte, N. Berrah, C. Callegari, O. Plekan, P. Finetti, K. C. Prince, L. Giannessi, C. Spezzani, E. Ferrari, E. Allaria, G. Penco, C. Serpico, G. De Ninno, B. Diviacco, S. Di Mitri, M. Yao, and K. Ueda (in preparation)

Two-photon excitation of ICD in neon dimers

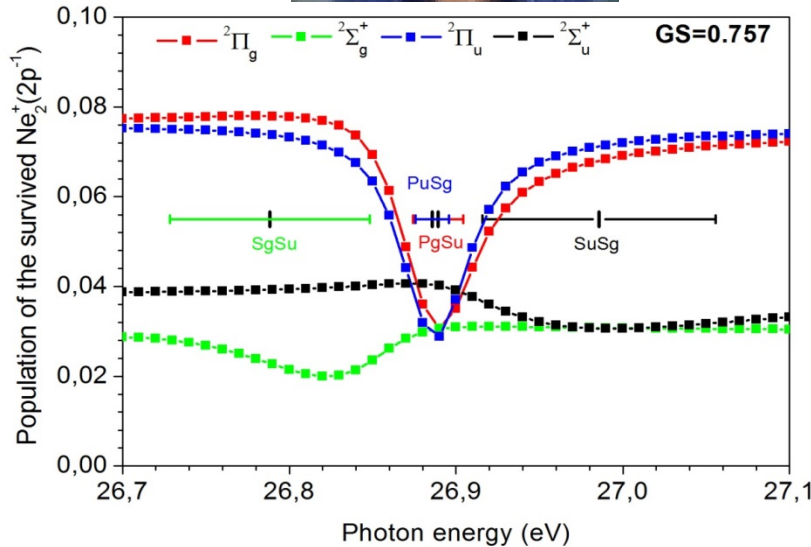
Theoretical prediction



$$E_{2s} - E_{2p} = \hbar\omega = 26.9 \text{ eV}$$

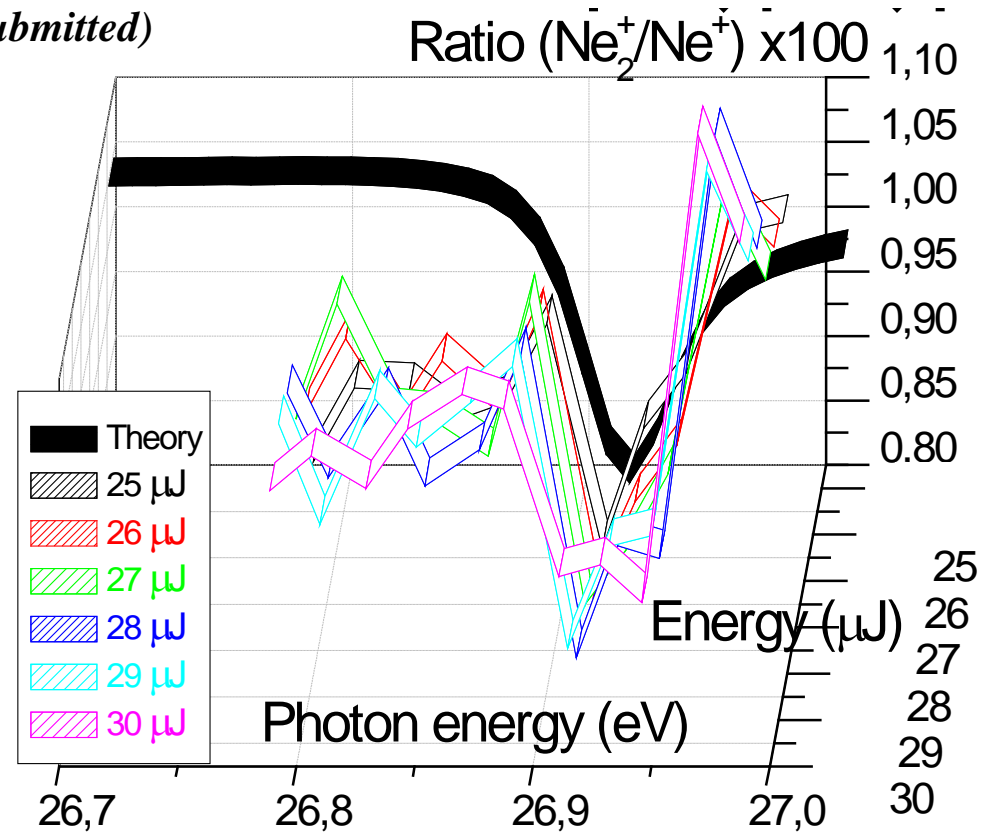
Two-photon-excitation-induced ICD in Ne₂

A Dubrouil, M Reduzzi, M Devetta, C Feng, J Hummert, P Finetti, O Plekan, C Grazioli, M Di Fraia, V Lyamayev, A La Forge, R Katzy, F Stienkemeier, Y Ovcharenko, M Coreno, N Berrah, K Motomura, S Mondal, K Ueda, K C Prince, C Callegari, *A I Kuleff, Ph V Demekhin, G Sansone*



$$E_{2s} - E_{2p} = \hbar\omega = 26.9 \text{ eV}$$

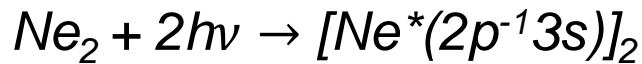
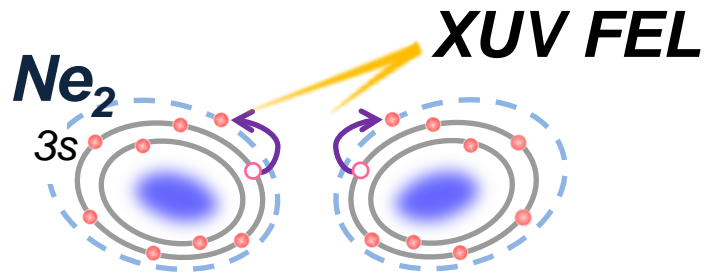
(submitted)



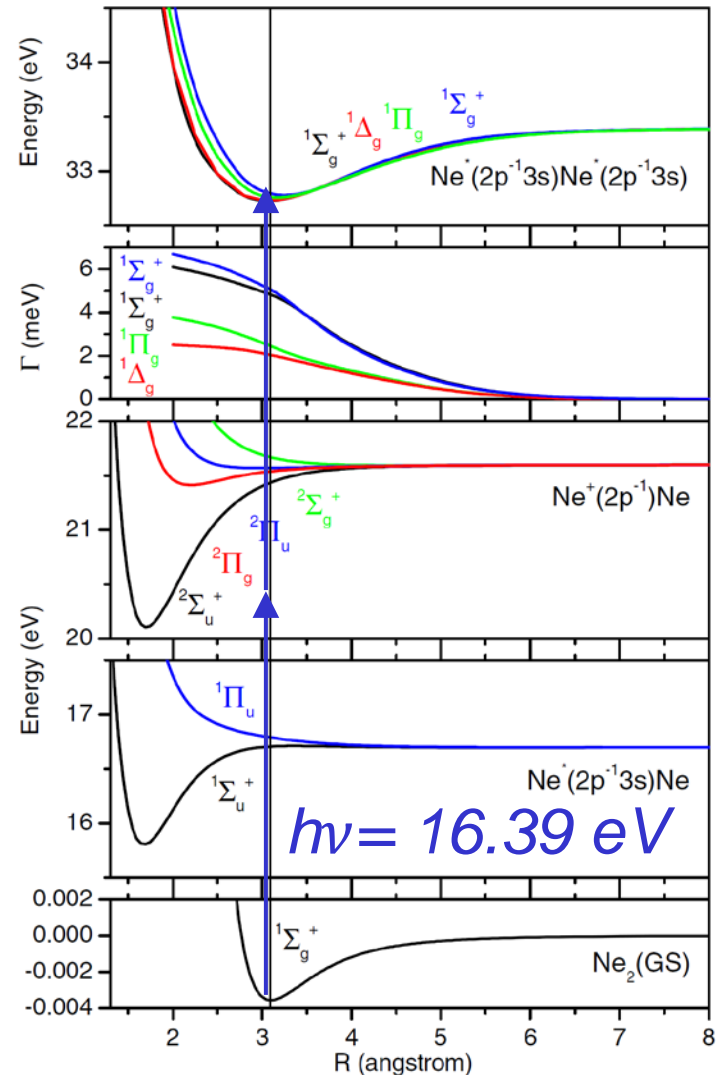
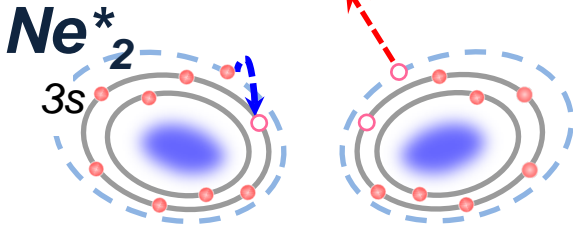
We have experimentally proved the depletion of ionized neon dimers Ne₂⁺ through the resonance as predicted by theory!

Two-photon two-site excitation of ICD

ICD processes induced by two-photon double excitation in Ne_2



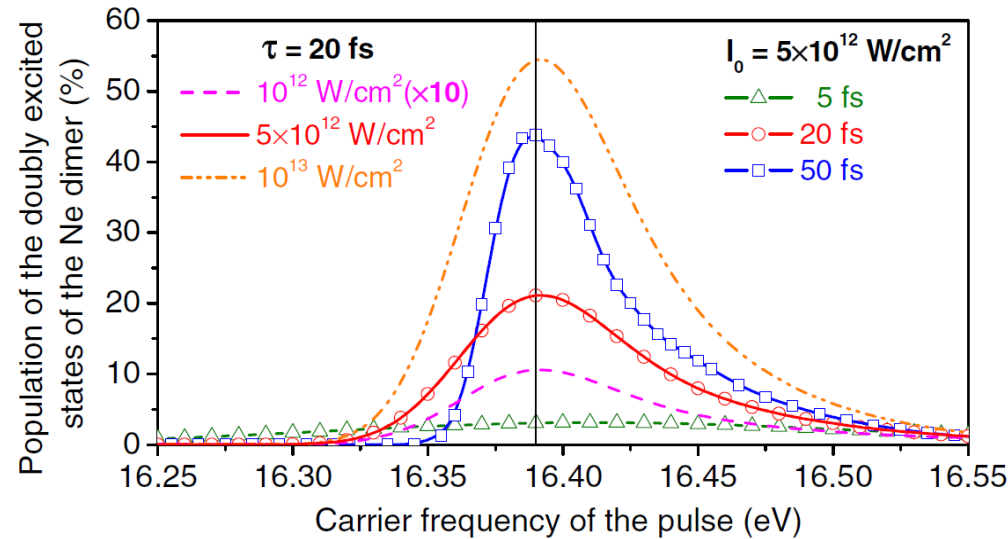
• • • two photon double excitation



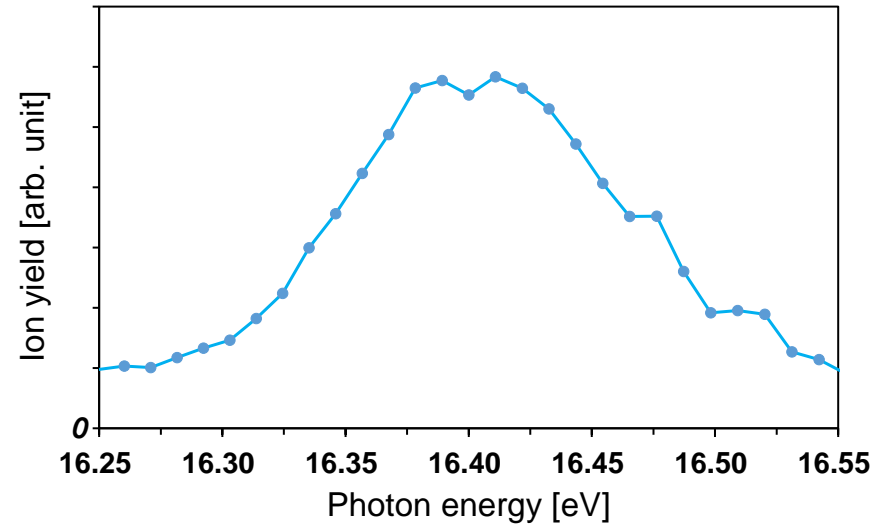
Resonant energy to $[\text{Ne}^*(2p^{-1}3s)]_2$: prediction and measurement

To observe the doubly excited states, we have measured Ne_2^+ yield as a function of FEL photon energy.

Theoretical curve



Ne_2^+ yield – photon energy plot



Two-photon absorption of Ne_2 leads to $[\text{Ne}^*(2p^{-1}3s)]_2$ that is subject to ICD.

Resonant energy: $h\nu = 16.39$ eV

Cf. Atomic resonant energy: 16.85 eV for $\text{Ne} \rightarrow \text{Ne}^*(2p^{-1}3s)$

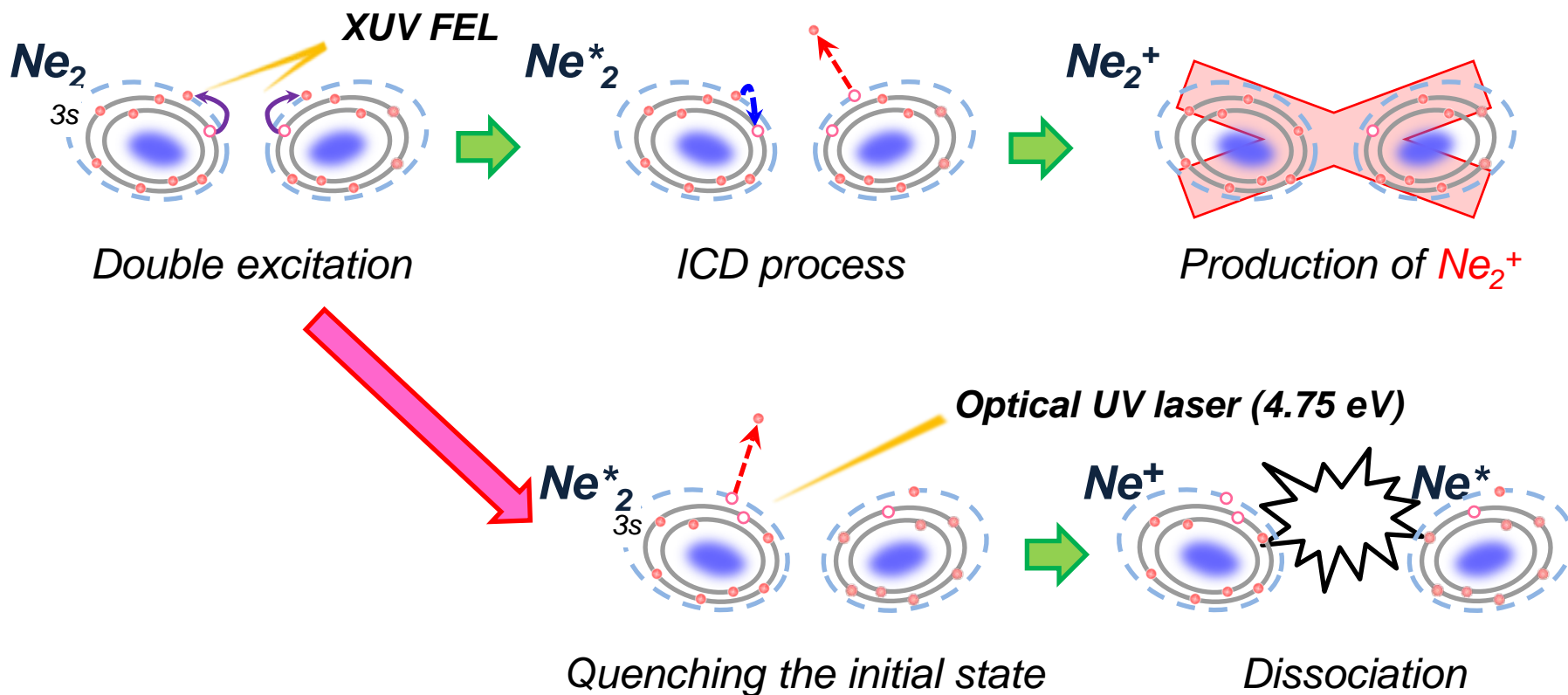
Doubly excited state $[\text{Ne}^*(2p^{-1}3s)]_2$ was observed!

Resonant FEL energy was 16.4 eV in agreement with the theoretical one!

Pump-probe scheme of the ICD process

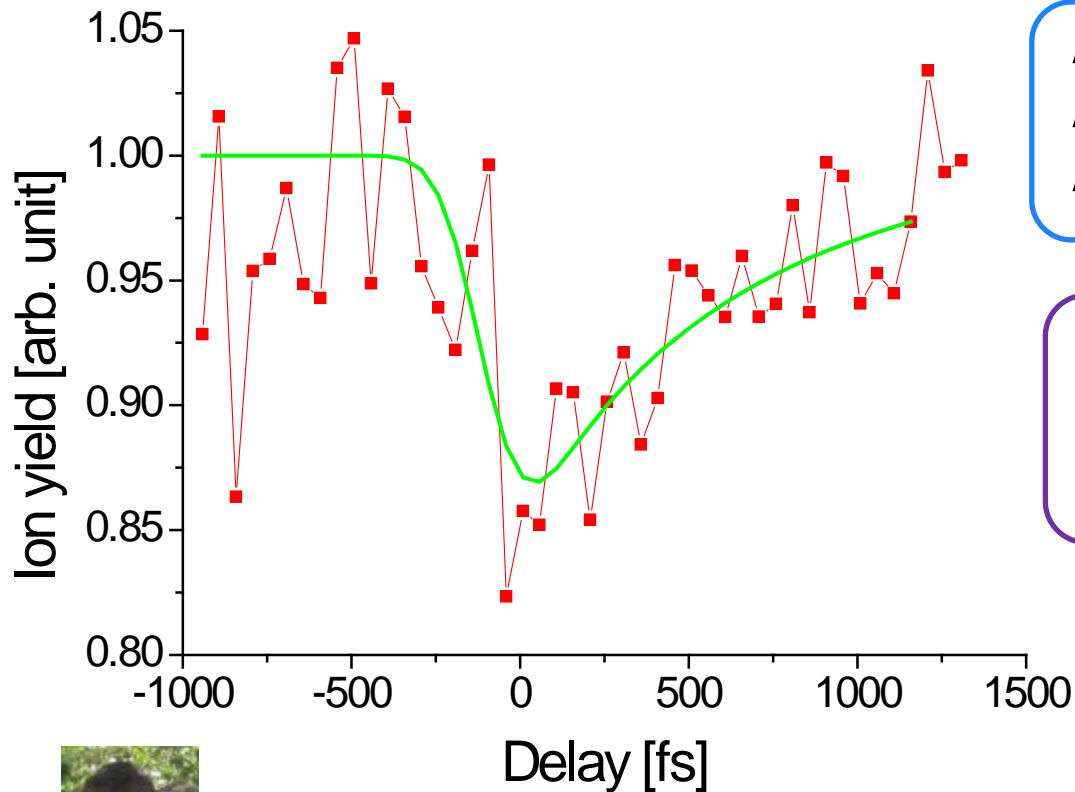
➤ The target process

* Quenching the initial state of the ICD process



→ Measuring the time constant of ICD process by scanning UV delay time

Time-resolved measurements of the ICD process



XUV-FEL pulse

- Photon energy: 16.40 eV
- Pulse duration (FWHM): ~ 70 fs
- Power density: $\sim 1.2 \times 10^{13}$ W/cm²

UV laser pulse

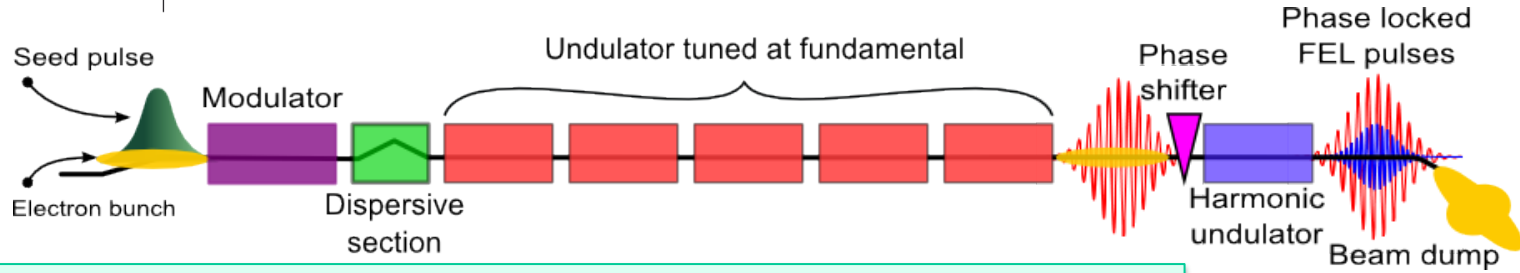
- Photon energy: 4.75 eV
- Pulse duration (FWHM): ~ 230 fs
- Power density: $\sim 0.8 \times 10^{13}$ W/cm²

ICD rates

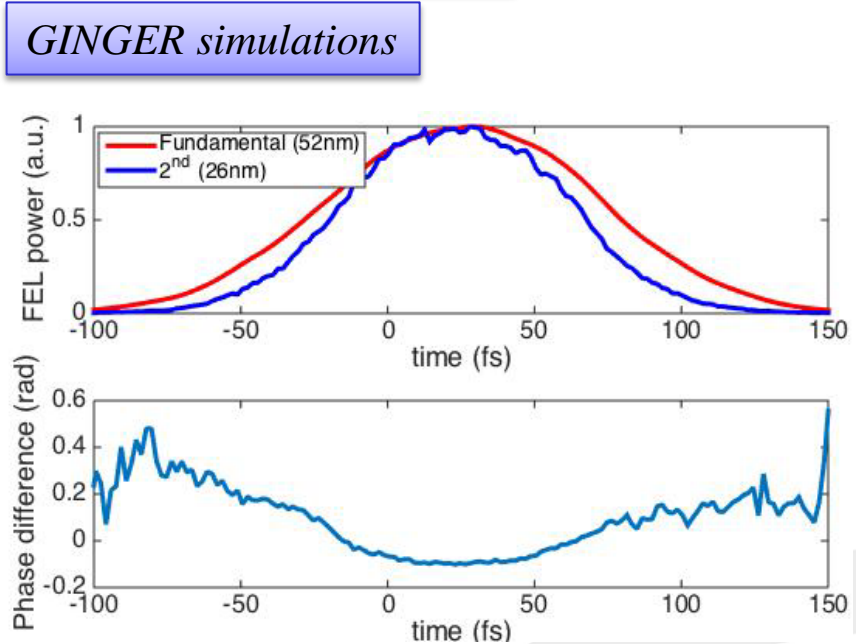
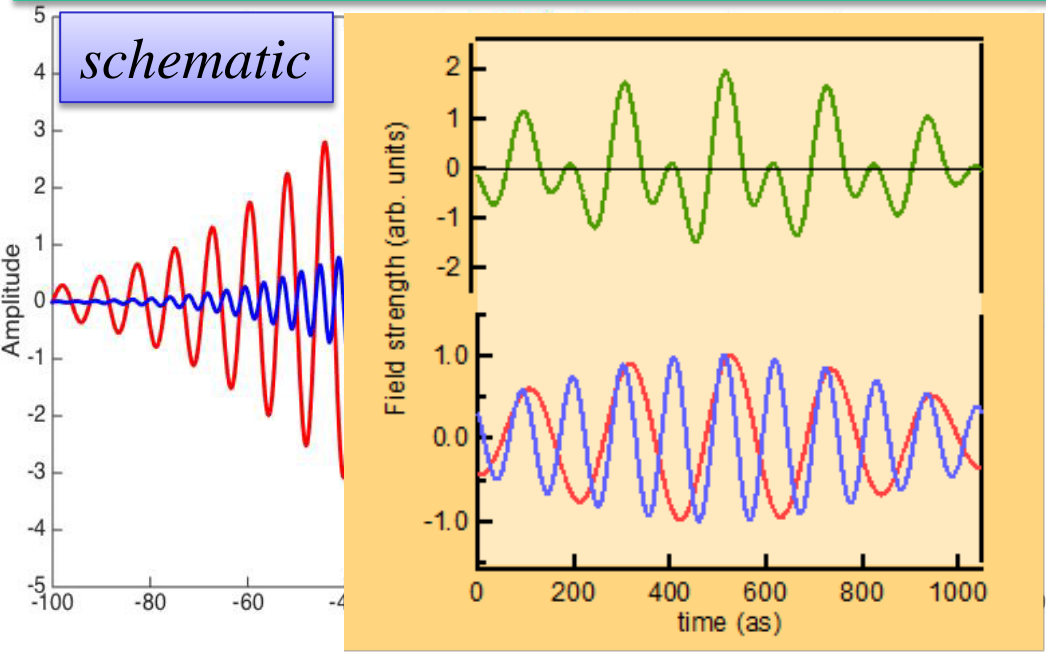
- $\tau_1 = \sim 100$ fs
- $\tau_2 = \sim 340$ fs (cf. 308 fs in theory)

T. Takanashi, H. Fukuzawa, K. Motomura, Y. Kumagai, S. Mondal, T. Tachibana, T. Nishiyama, K. Matsunami, K. Nagaya, P. Johnsson, P. Piseri, G. Sansone, A. Dubrouil, M. Reduzzi, P. Carpeggiani, C. Vozzi, M. Devetta, M. Negro, D. Faccialà, F. Calegari, A. Trabattoni, M. Castrovilli, Y. Ovcharenko, M. Mudrich, F. Stienkemeier, M. Coreno, M. Alagia, B. Schütte, N. Berrah, C. Callegari, O. Plekan, P. Finetti, K. C. Prince, L. Giannessi, C. Spezzani, E. Ferrari, E. Allaria, G. Penco, C. Serpico, G. De Ninno, B. Diviacco, S. Di Mitri, A. Kullef, N. Golubev, P. Demekhin, M. Yao, and K. Ueda (in preparation)

Phase-controlled harmonics!



Tuning the last undulator at one harmonic gives the possibility to control the phase between two pulses with different wavelength.



figures and data courtesy of Enrico Allaria

Atomic photoionization with phase control

1st+3rd harmonic:

Chen, Yin, Elliot, PRL **64**, 507 (1990)

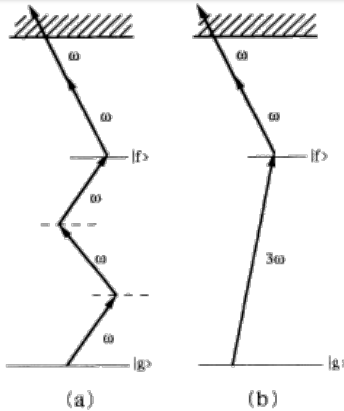


FIG. 1. The two processes which interfere in this observation. The transition $|g\rangle \rightarrow |f\rangle$ is three- and one-photon allowed, as shown in (a) and (b).

1st+2nd harmonic:

Yin et al., PRL **69**, 2352 (1992)

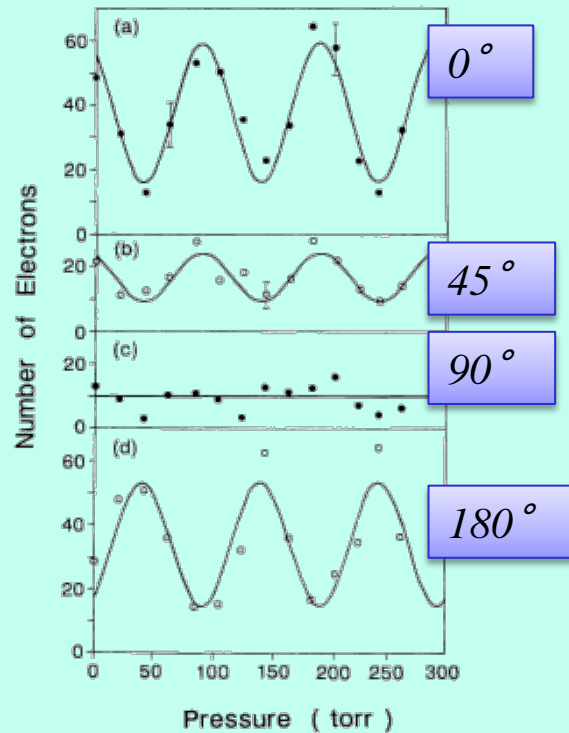


FIG. 3. Experimental data. The total electron count as a function of pressure of N_2 gas in the phase delay cell for the four detectors positioned at (a) 0° , (b) 45° , (c) 90° , and (d) 180° . The solid line is the result of a least-squares fit of a sinusoidally varying curve to the data.

Our scheme

Ne 2p \rightarrow 4s for 1st harmonic

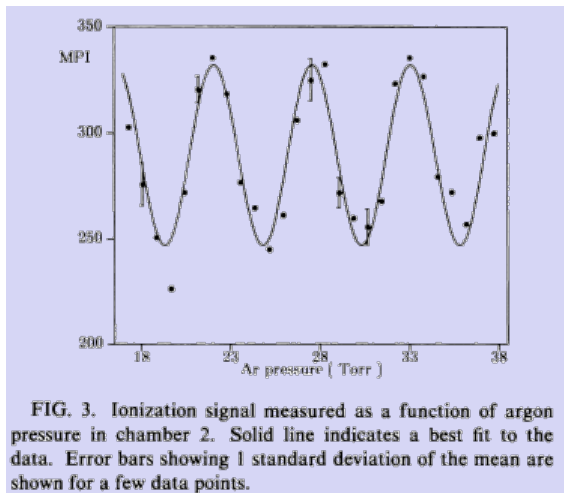
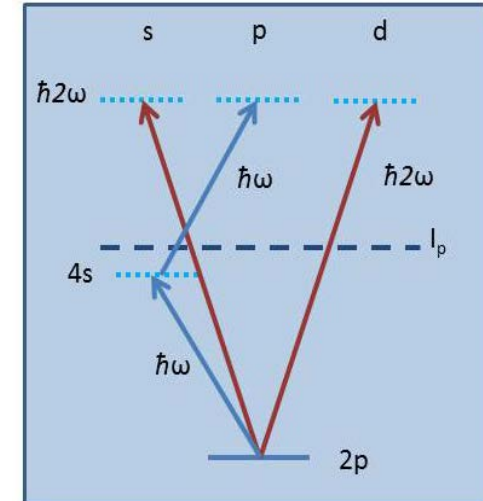
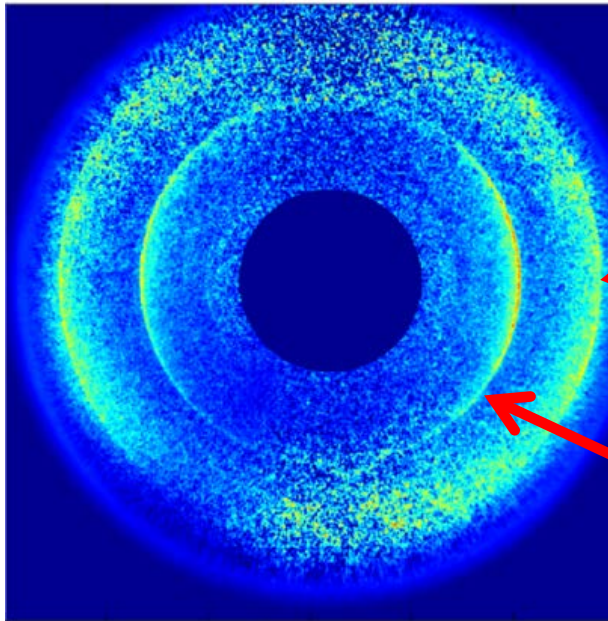


FIG. 3. Ionization signal measured as a function of argon pressure in chamber 2. Solid line indicates a best fit to the data. Error bars showing 1 standard deviation of the mean are shown for a few data points.

$$|M_1 + M_2(\phi)|^2 = |M_1|^2 + |M_2(\phi)|^2 + 2 \operatorname{Re}(M_1 M_2(\phi))$$

Two color coherent control at FERMI



First harmonic plus second harmonic

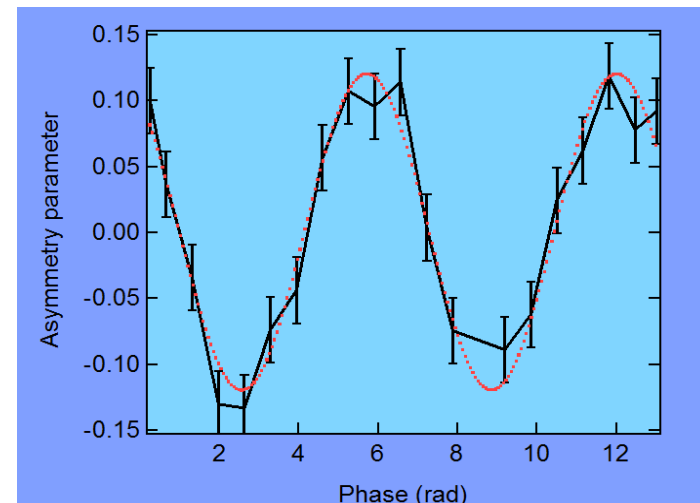
VMI image at 62.974 nm (Ne 2p⁵4s resonance).

Outer ring: a sharp line due to third harmonic radiation, plus a broad distribution due to fluorescence (the fluorescence excites photoelectrons which then impinge on the detector).

Inner sharp ring: electrons ionized by two photons.

The center part of the image has been blanked

K. C. Prince, E. Allaria, C. Callegari, R. Cucini, G. De Ninno, S. Di Mitri, B. Diviacco, E. Ferrari, P. Finetti, D. Gauthier, L. Giannessi, N. Mahne, G. Penco, O. Plekan, L. Raimondi, P. Rebernik, E. Roussel, C. Svetina, M. Trovò, M. Zangrando, G. Sansone, M. Reduzzi, P. Carpeggiani, A. Grum-Grzhimailo, E.V. Gryzlova, S.I. Strakhova, K. Bartschat, D. Iablonskyi, Y. Kumagai, T. Takanashi, K. Ueda, A. Fischer, F. Stienkemeier, E. Ovcharenko, T. Mazza, M. Meyer (submitted)



Asymmetry emerges from non-linear process
N. B. Baranova and B. Ya Zel'dovich, Sov. Phys. JETP 71, 1043 (1990)

FERMI can (potentially) generate up to 4 consecutive phase-locked harmonics

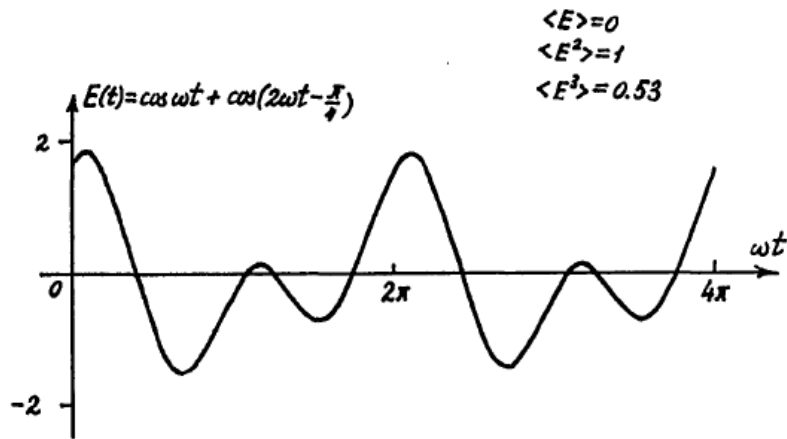
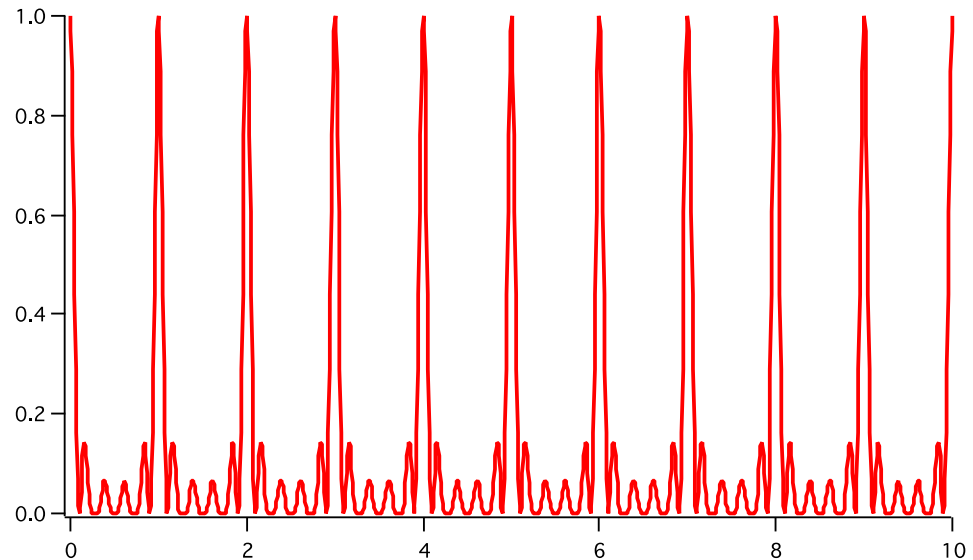
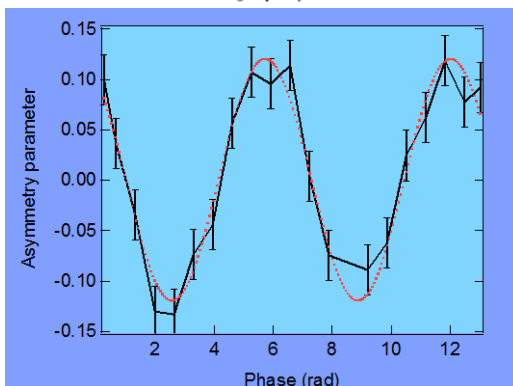


Fig. 1. Time dependence of the electrical field [Eq. (1)] at $E_1 = E_2 = 1$, $\Delta\varphi = \varphi_2 - 2\varphi_1 = 45^\circ$. The time-averaged field is equal to zero, $\langle E \rangle = 0$, but the asymmetry in the upper direction is evident and is characterized by $\langle E^3 \rangle = 0.53 > 0$.



e.g. interval ~ 200 as at 63 nm (20 eV)



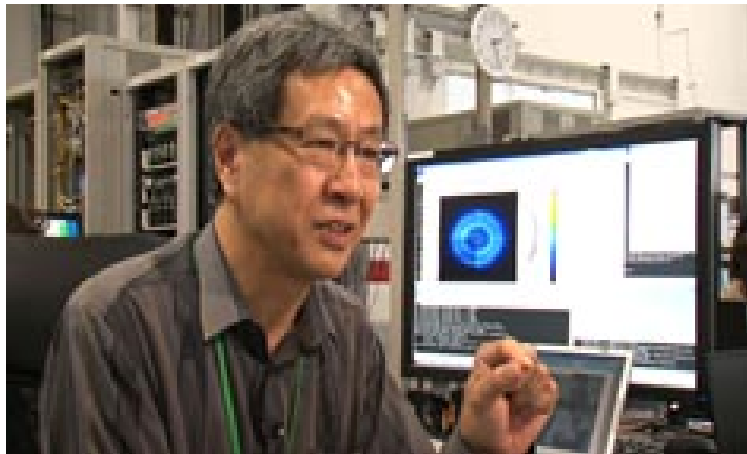
Summary

- ICD cascades and inelastic scattering of ICD electrons in neon clusters measured at SCSS and FERMI
- Time resolved study on two-photon excited ICD in neon clusters at FERMI
- Coherent control of photoionization with phase controlled harmonics at FERMI

Outlook

- Full characterization and control of photoionization and photo-dissociation with phase controlled harmonics at FERMI

SACLA XFEL (lased on 7 June 2011)



on air!



SACLA XFEL

Photon energy range: 4-20 keV

Photon numbers: $\sim 10^{11}$ photons/pulse (5-15 keV)

Repetition rate: 10~60 Hz

Pulse width ~ 10 fs

Focusing optics: $\sim 1 \mu\text{m}$ (1.5 m) \rightarrow 50 nm (0.5 m)

Commissioning beam time: Nov. 2011-Feb. 2012

7-11 Nov. 2011: Detector test (no real FEL beam...)

20-24 Feb. 2012: Serial femtosecond crystallography

User beam time started in March 2012

First two beam times in 2012:

Atoms, molecules and atomic clusters

Third beam time in 2013: pump-probe, failed

Fourth beam time in 2014:

Single-shot imaging of giant xenon clusters

Fifth and six beam time in 2014:

Pump-probe experiment on clusters

What should we do with SACLA?

Novel structure determination

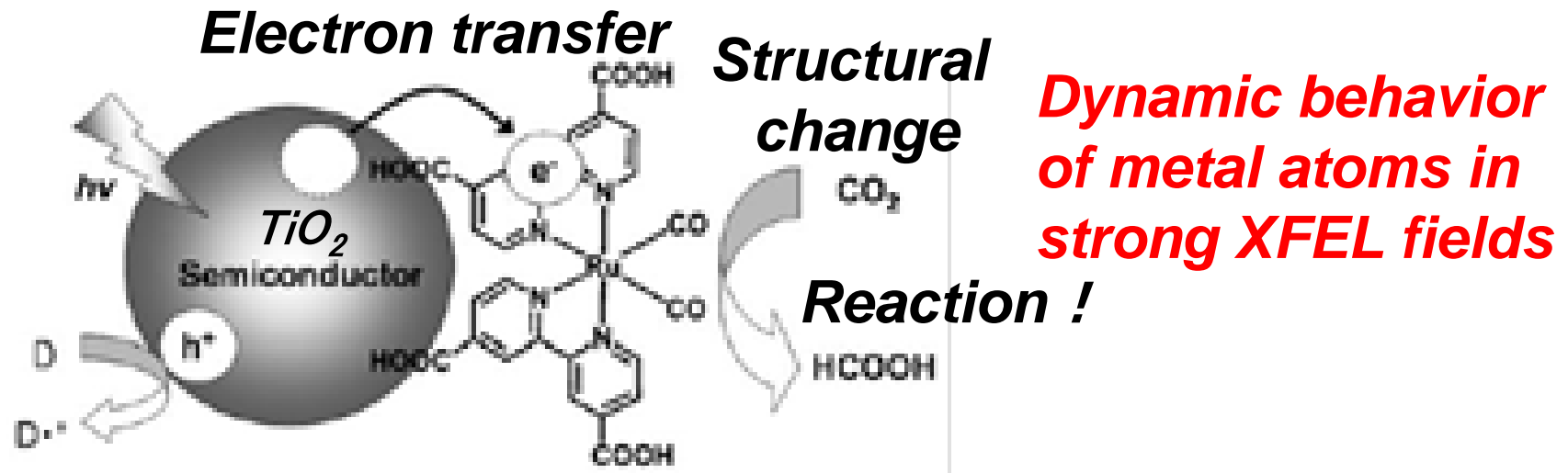
•phasing

•radiation damage

understanding dynamic behavior of heavy atoms!

Dynamic structure and light-induced reaction

—Femtosecond electronic and structure changes



Sato et al. Angew. Chem. 49, 5101 (2010) (Toyota)

An aerial photograph of a large scientific facility, likely a synchrotron or free-electron laser, featuring a prominent circular building and several long, rectangular structures. The facility is surrounded by lush green hills and a clear sky.

***Multi-photon, multiple ionization of
Atoms, Molecules and Rare-gas clusters***

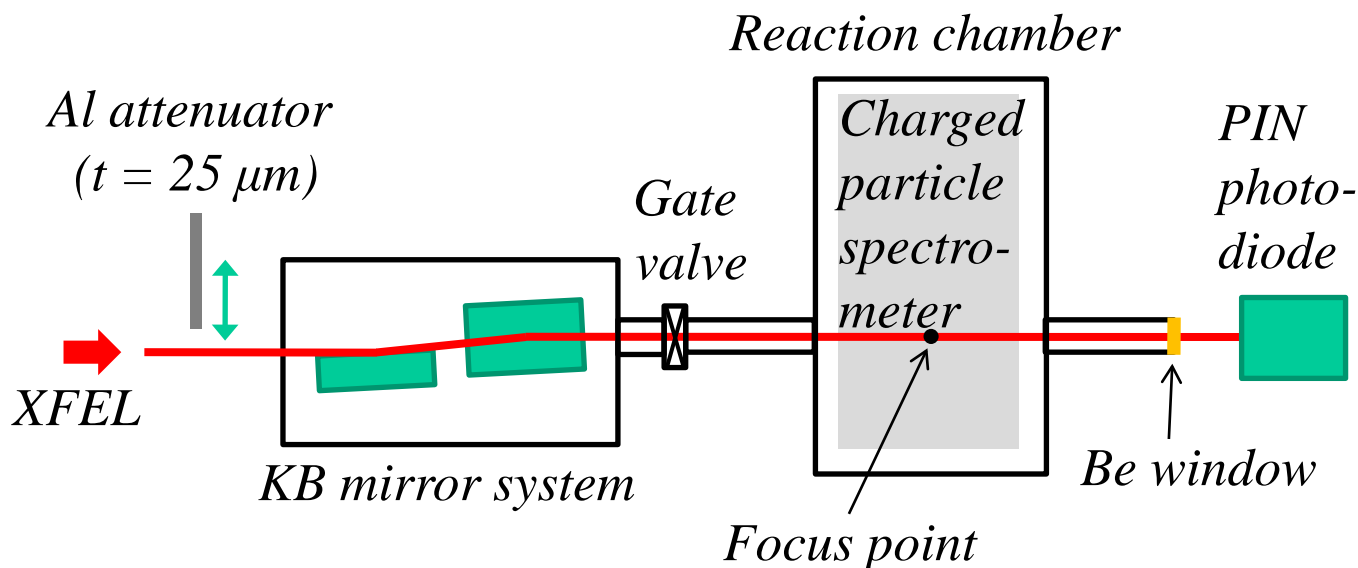
Single-shot imaging of Nano-clusters

***Pump-probe experiments for
Atoms, Molecules and Rare-gas clusters***

Aiming at probing

Ultrafast electron and structure dynamics

Experimental configuration @ SACLA BL3 EH3



XFEL pulses

Photon energy: 5 and 5.5 keV

(Wavelength: 0.25 and 0.22 nm)

Band width: ~ 60 eV (FWHM)

Repetition: 10-30 Hz

Pulse energy before KB mirror:

$\sim 240 \mu\text{J}$ ($\sim 3 \times 10^{11}$ photons) @5.5keV

Fluctuation of pulse energy:

$\pm 25\%$ (50% FWHM)

@Focus point

Focus size: $\sim 1.5 \mu\text{m}$ (FWHM)

Peak fluence:

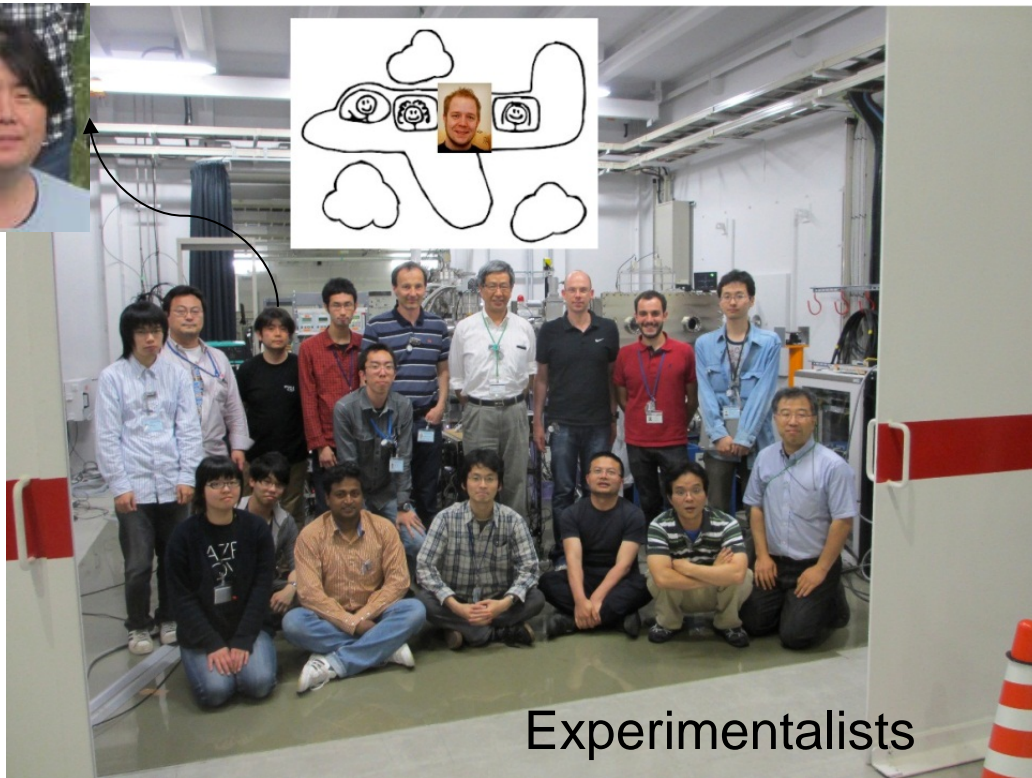
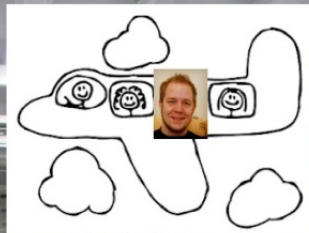
$\sim 47 \mu\text{J}/\mu\text{m}^2$ (atoms, clusters),

$\sim 26 \mu\text{J}/\mu\text{m}^2$ (molecules)

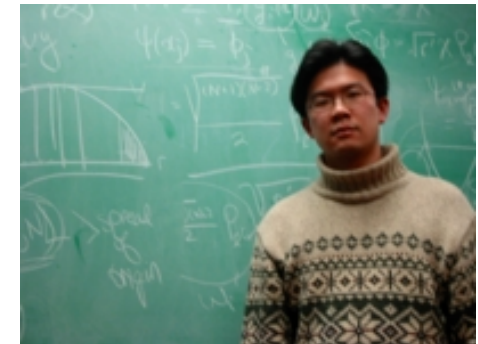
Sample gas was introduced as a pulsed super sonic gas jet to the focus point.

I. Deep inner-shell multiphoton absorption by intense x-ray free-electron laser pulses

H. Fukuzawa, S.-K. Son, K. Motomura, S. Mondal, K. Nagaya, S. Wada, X.-J. Liu, R. Feifel, T. Tachibana, Y. Ito, M. Kimura, T. Sakai, K. Matsunami, H. Hayashita, J. Kajikawa, P. Johnsson, M. Siano, E. Kukk, B. Rudek, B. Erk, L. Foucar, E. Robert, C. Miron, K. Tono, T. Togashi, Y. Inubushi, T. Sato, T. Katayama, T. Hatsui, T. Kameshima, M. Yabashi, M. Yao, R. Santra, and K. Ueda [PRL **110, 173005 (2013) & JPB **46**, 164024 (2013)]**



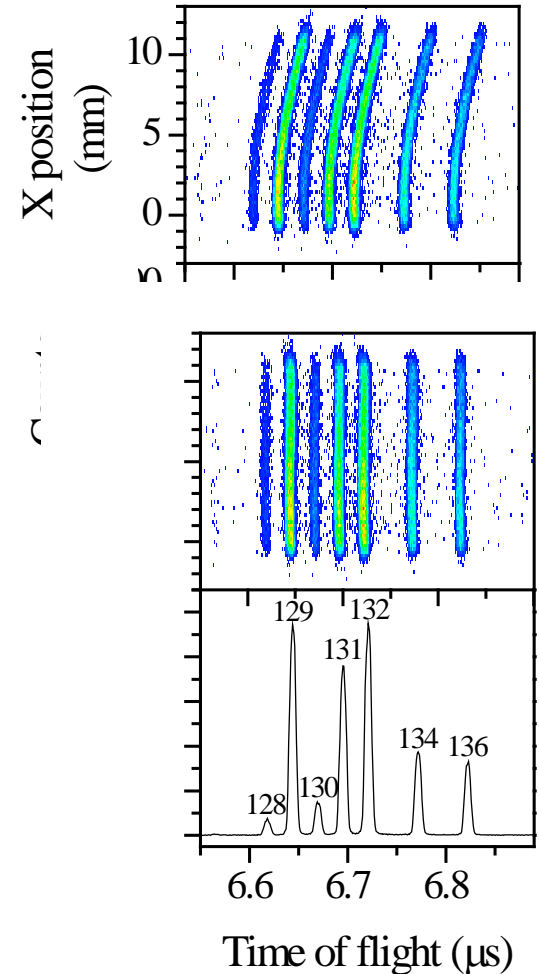
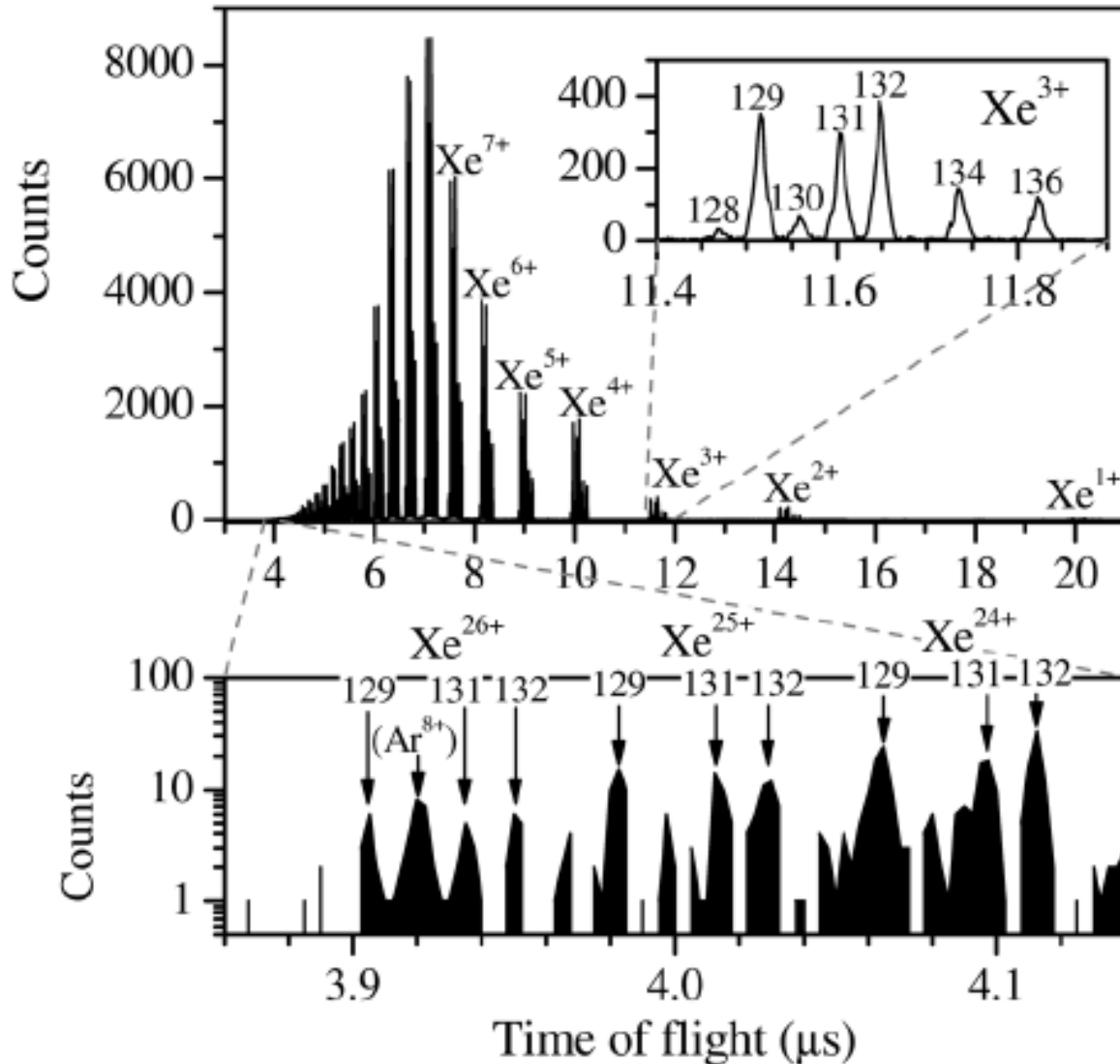
Experimentalists



Theorists

Time of Flight spectrum of xenon ions

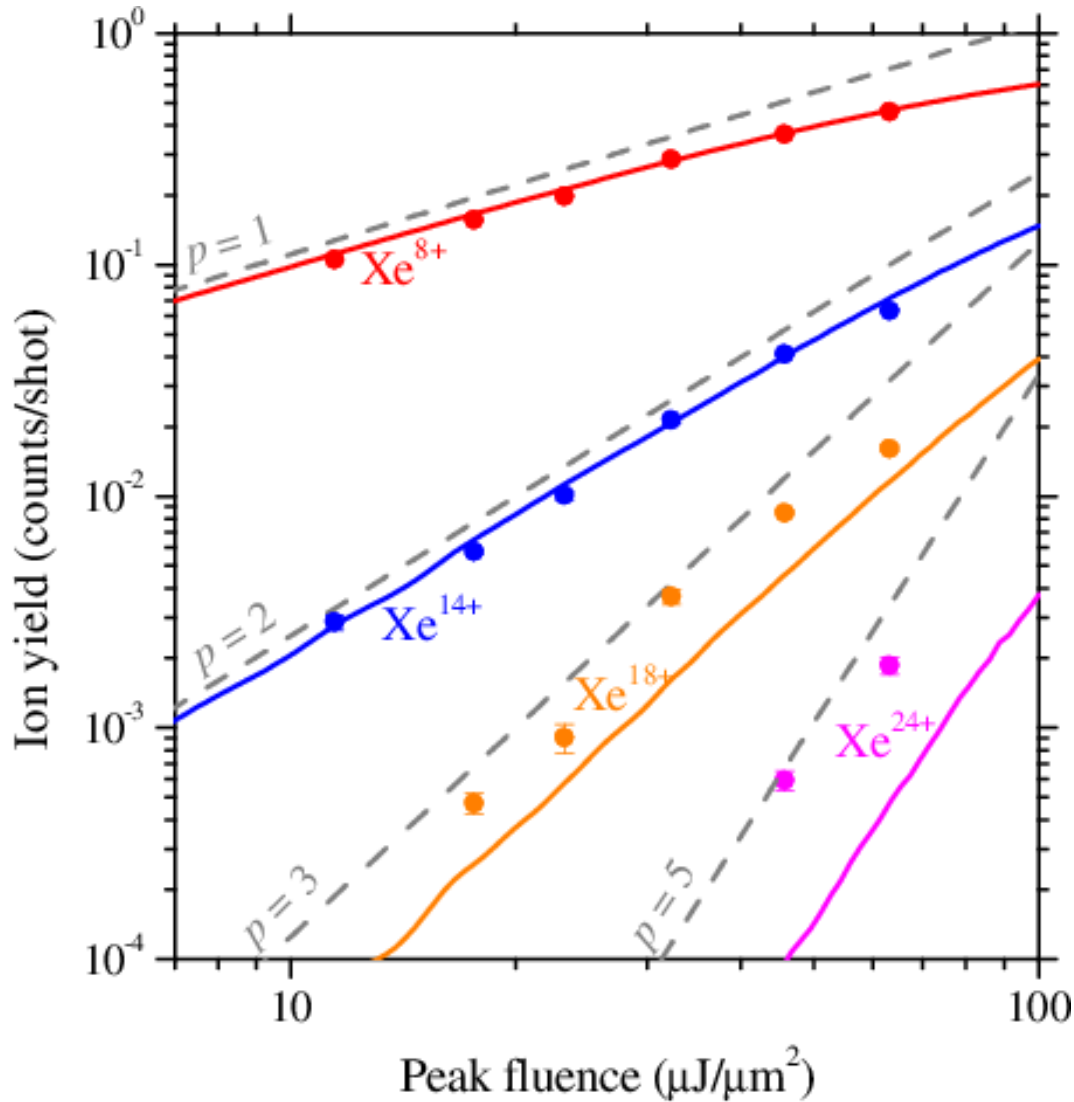
5.5 keV, 50 $\mu\text{J}/\mu\text{m}^2$ at SACLA



2D position resolved TOF improves the resolution!

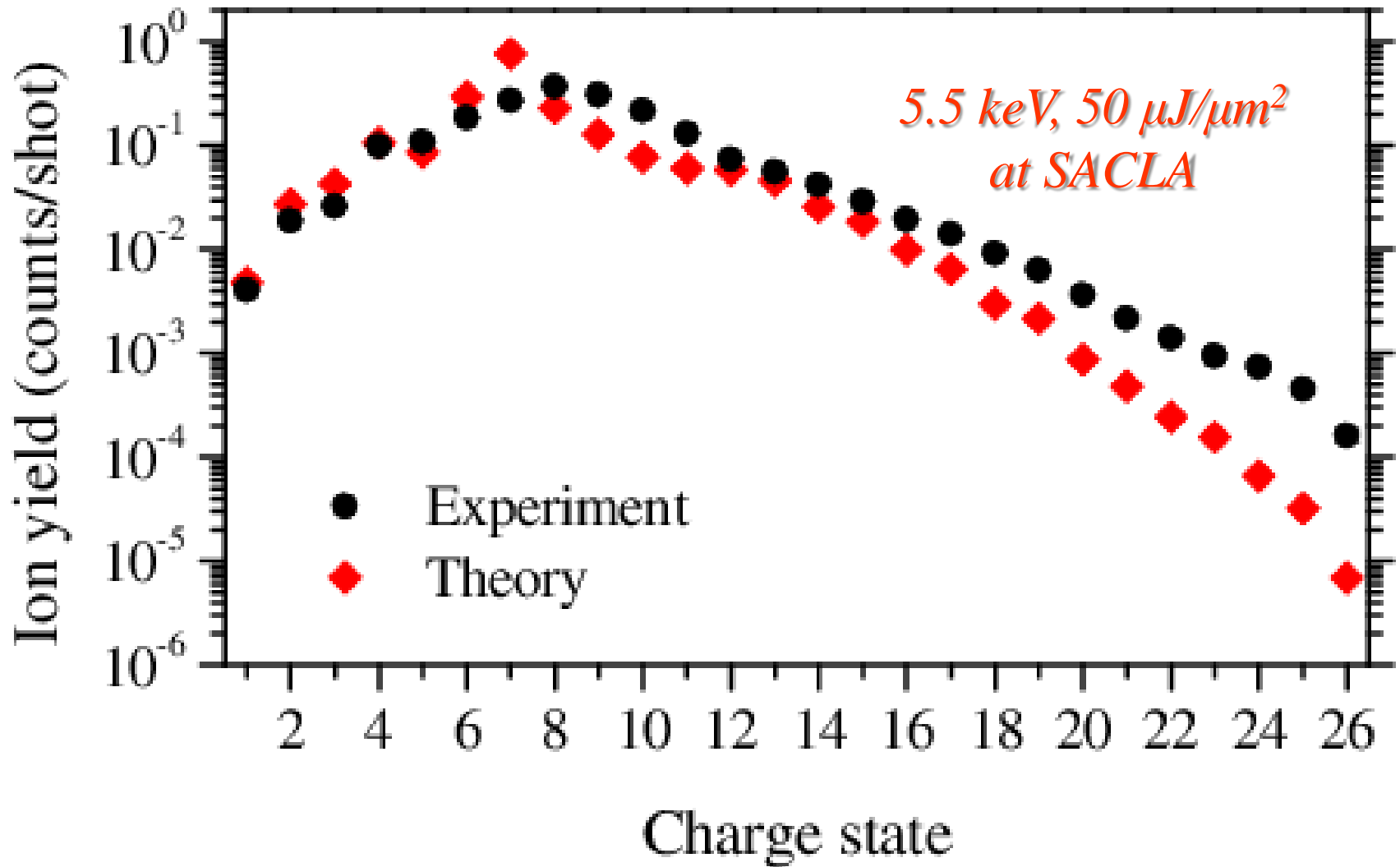
High charge states Xe^{n+} with n up to 26 are produced!

XFEL fluence dependence for Xe^{n+} yields



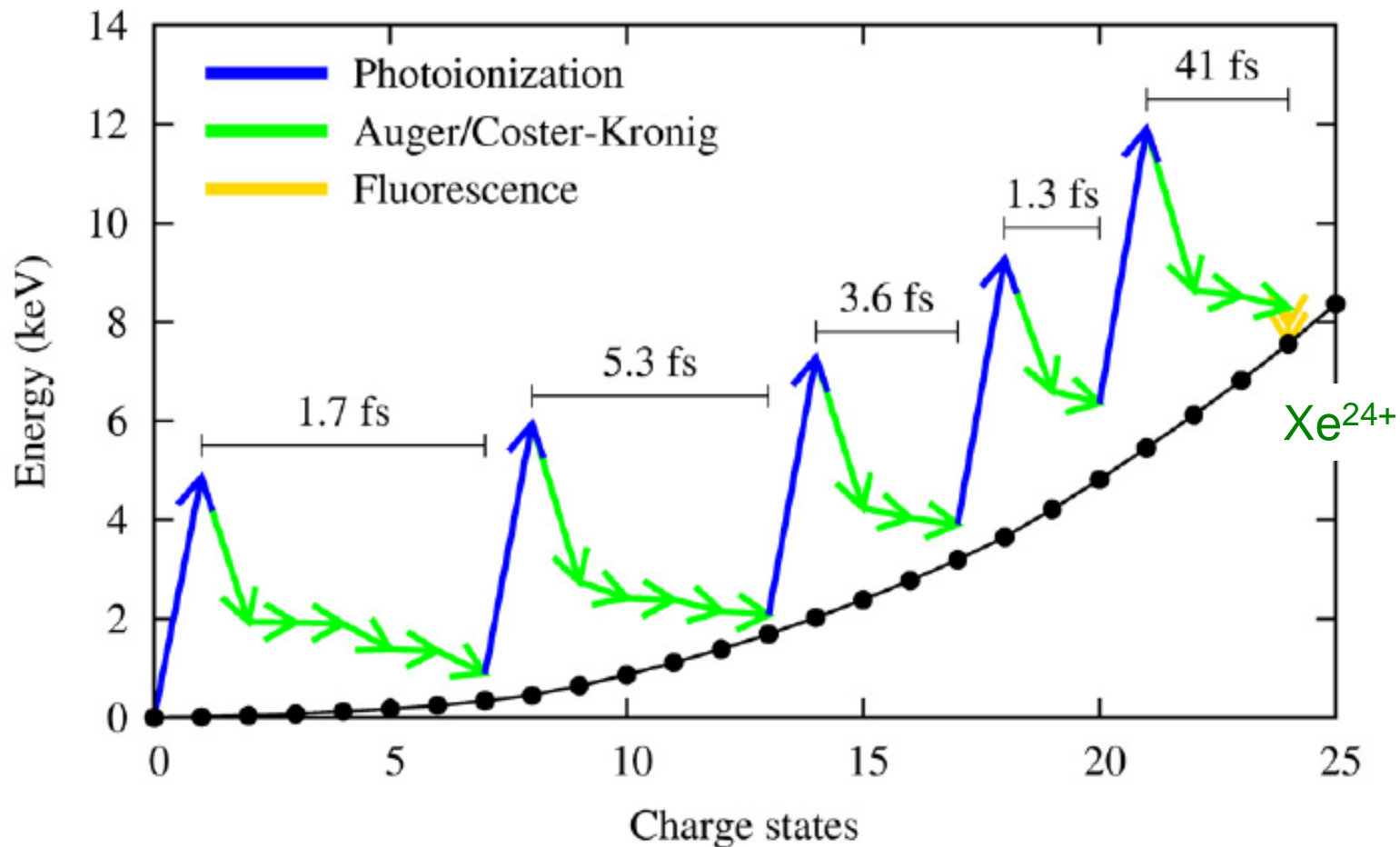
With help of ab initio calculations, we find that the observed high charge states ($n \geq 24$) are produced via five-photon absorption, evidencing the occurrence of multiphoton absorption involving deep inner shells.

Xenon ion charge distributions (exper. vs theory)



A newly developed theoretical model shows good agreement with the experiment!

An exemplary pathway of multiphoton multiple ionization



A newly developed theoretical model elucidates the complex pathways of sequential electronic decay cascades accessible in heavy atoms, *revealing that L shell ionization and sequential electronic decay cycles are repeated multiple times within the XFEL pulse duration of ~ 10 fs.*

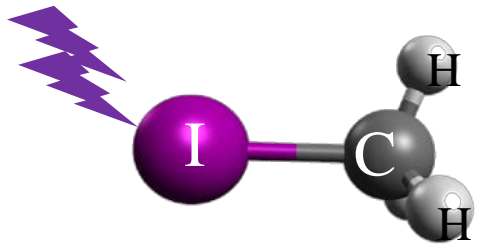
Fukuzawa, Son et al. PRL 110, 173005 (2013)

II. Charge transfer and molecular dissociation following deep inner-shell multi-photon multiple ionization of CH_3I and 5-Uracil molecules by intense x-ray free-electron laser pulses from SACLA

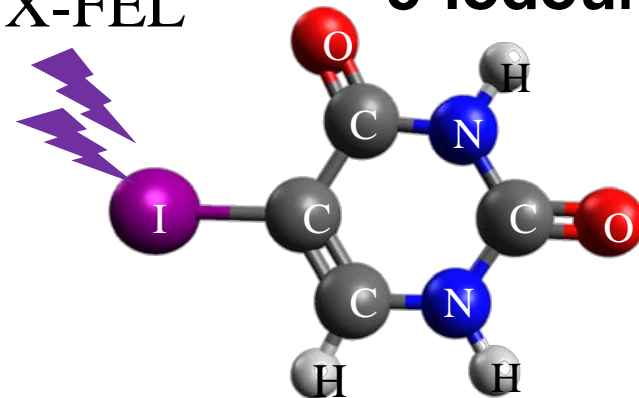


K. Motomura, E. Kukk, H. Fukuzawa, K. Nagaya, S. Omura, S. Wada, S. Mondal, T. Tachibana, Y. Ito, T. Sakai, K. Matsunami, A. Rudenko, C. Nicolas, X.-J. Liu, C. Miron, Y. Zhang, Y.H. Jiang, J. Chen, A. Milam, D. Kim, K. Tono, T. Hatsui, Y. Inubushi, M. Yabashi, H. Kono, M. Yao and K. Ueda (submitted).

X-FEL **iodomethane**



X-FEL **5-iodouracil**

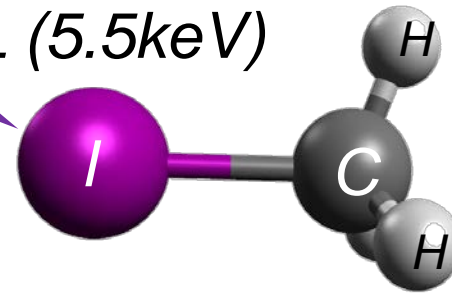


Ionization of iodomethane

We expect that the ionization proceeds with this sequence.

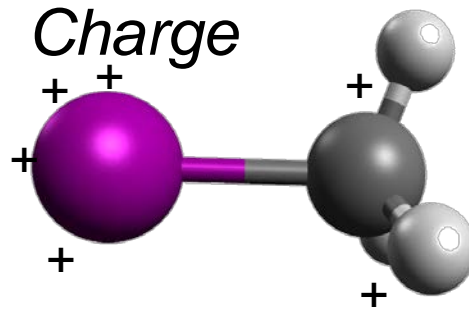
XFEL (5.5keV)

Selective ionization of iodine L-shell

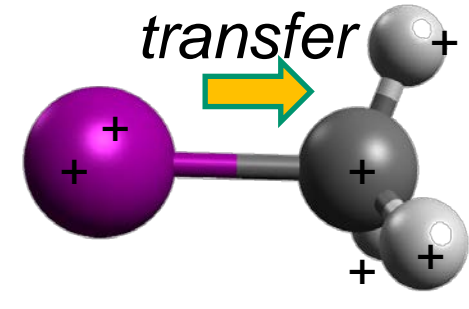


Next photon absorption

Production of highly charged iodine ion by Auger cascades



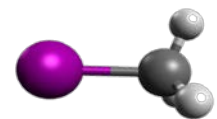
Charge transfer from iodine



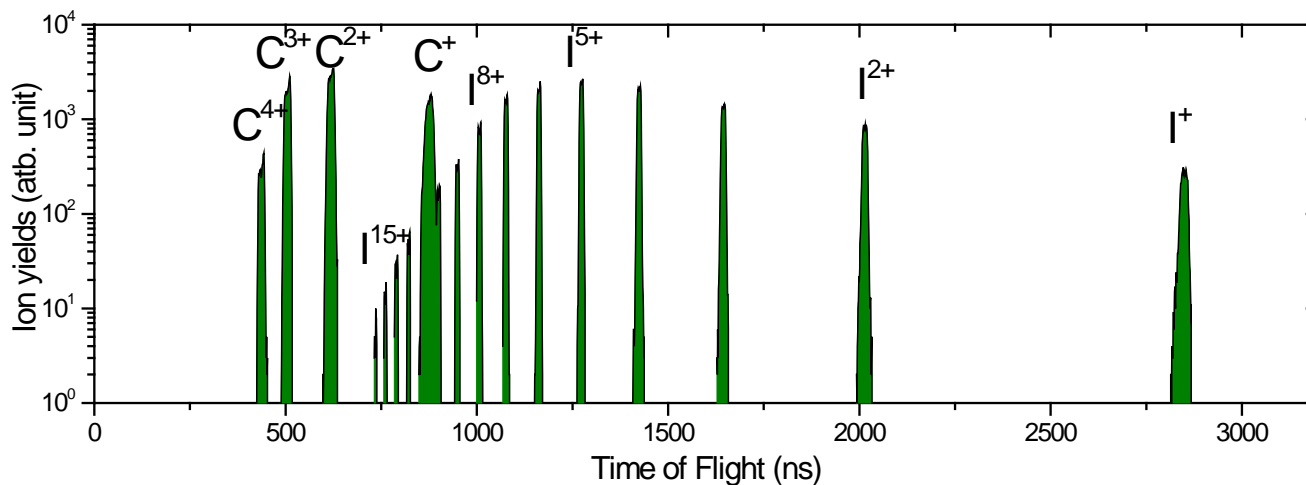
Coulomb explosion

We tried to understand these processes by measuring momentum of ion fragments.

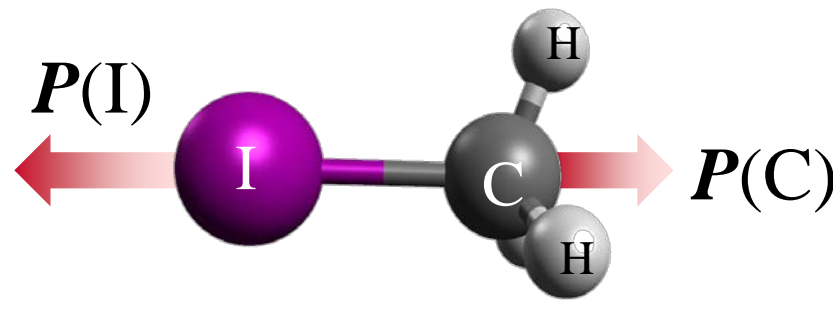
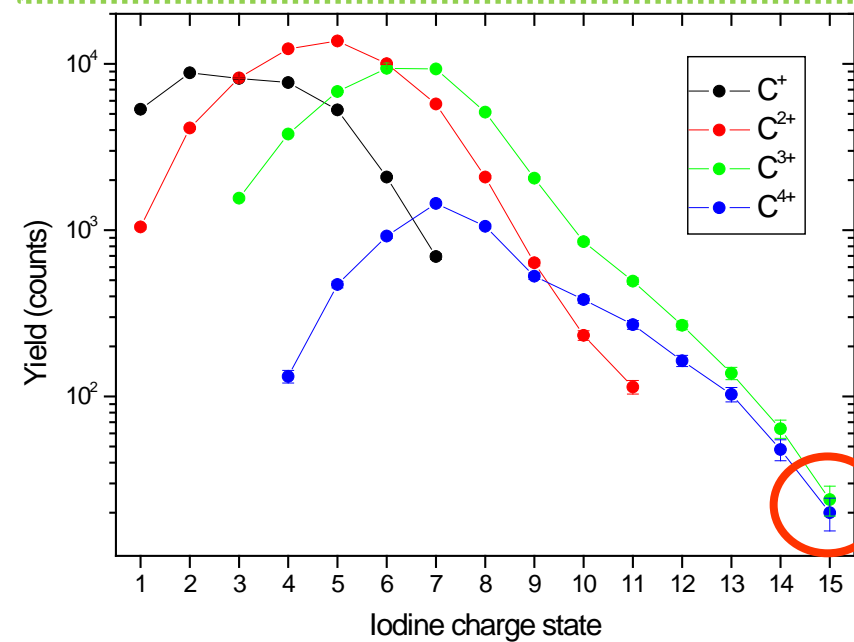




Iodomethane: Charge distribution

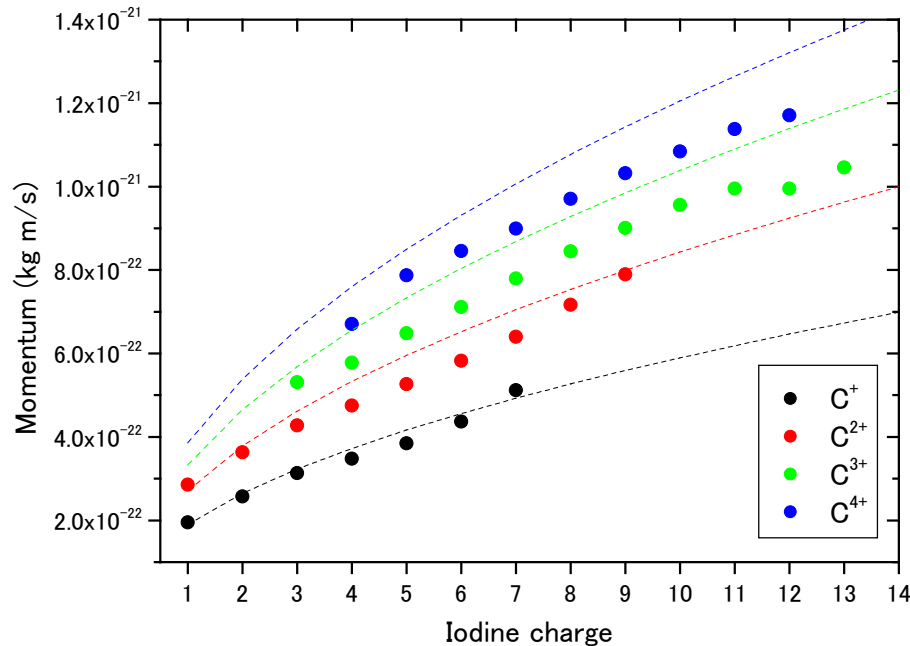


Momentum conservation: $P(H1)+P(H2)+P(H3)+P(C)+P(I) = 0$

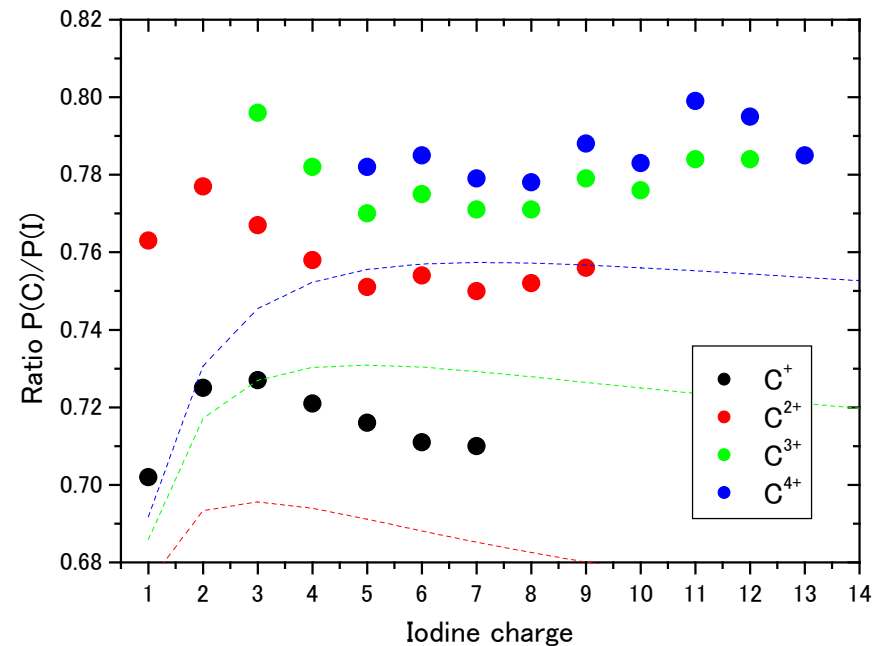


Charge state dependence of ion momentum

Charge state dependence for the momentum of carbon ions



Ratio of the momentum of carbon ions and iodine ions



Dashed line: Simulation with instant charge build up and transfer within the intact molecule; the charges are arranged before the Coulomb explosion starts.

The results of the simulation do not agree with the experimental results. Dissociation may compete with the charge buildup and transfer.

Charge build up and transfer model

We introduced two parameters “ τ ” and “ R ”.

XFEL ionizes iodine L-shell

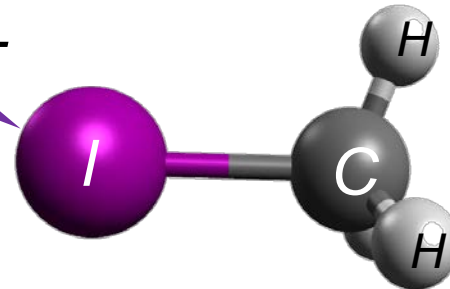
Charge build up in a molecule.

$$Q_{total}(t) = Q \left(1 - e^{-\frac{t}{\tau}} \right)$$

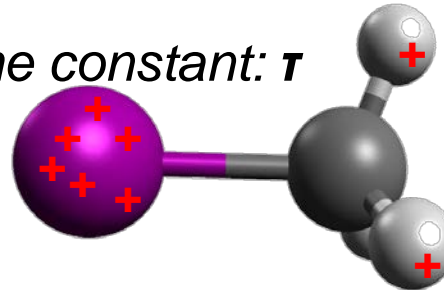
Charge transfer from iodine

$$\frac{dQ_{CH_3}(t)}{dt} = R \times Q_I(t)$$

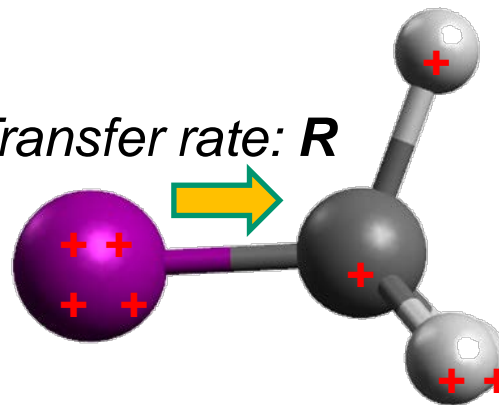
XFEL



Time constant: τ



Transfer rate: R



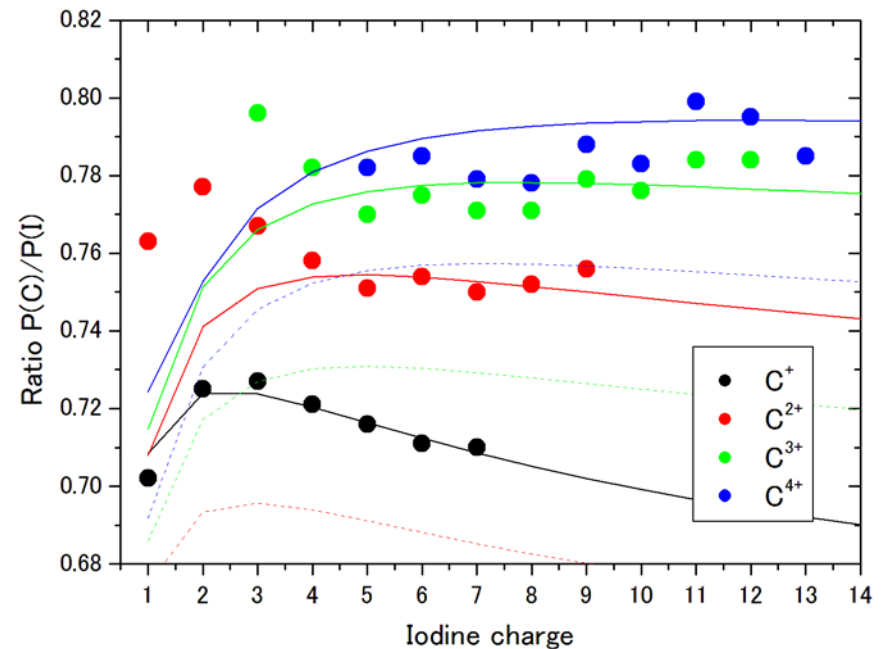
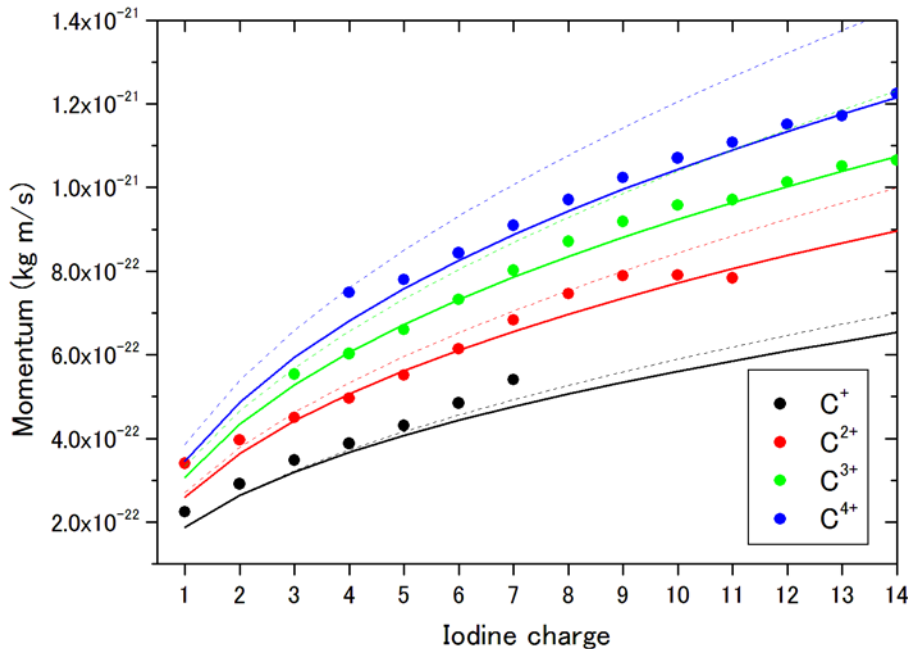
Next photon absorption

Coulomb explosion

Comparison with charge build up and transfer model

*Charge state dependence
for the momentum of carbon ions*

*Ratio of the momentum of
carbon ions and iodine ions*



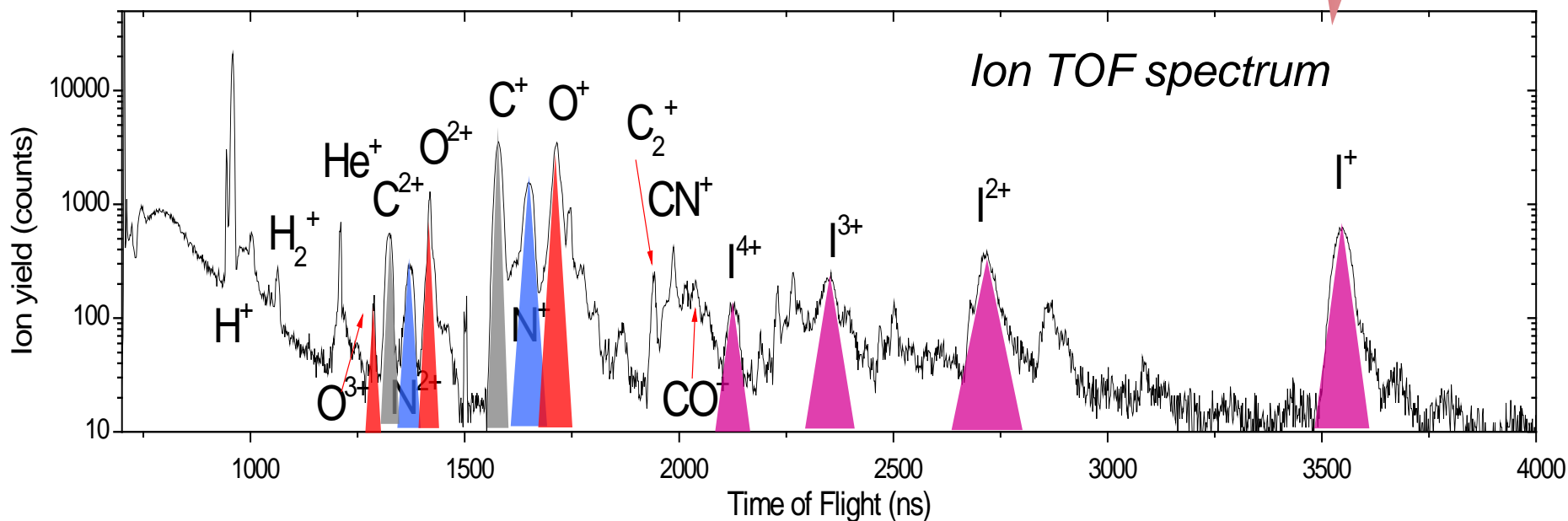
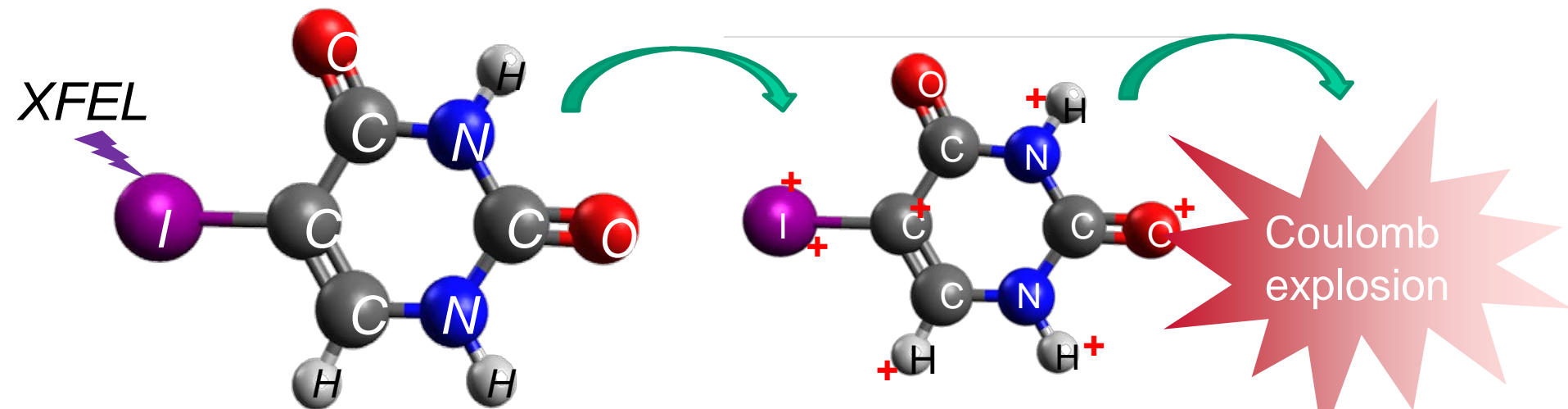
Solid line: Simulation with charge build up and transfer

Dashed line: Simulation with Instant charge buildup and transfer

Using the parameters $\tau = 9 \text{ fs}$ and $R = 0.37 \text{ fs}^{-1}$, the simulations agree with experimental results.

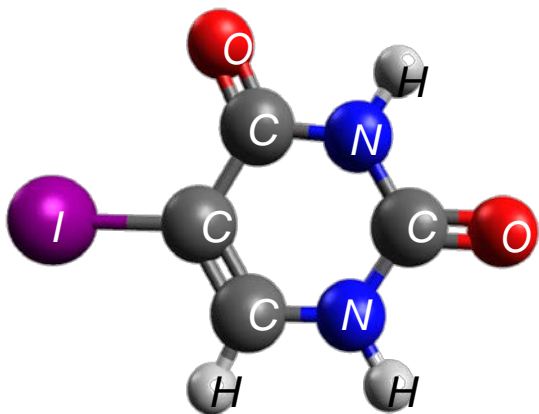
τ of 9 fs is consistent with the results of atomic xenon results and roughly the same as the XFEL pulse width ($\sim 10 \text{ fs}$).

5-iodouracil ($C_4H_3IN_2O_2$)



We observed iodine ions with the charge up to +4 and many more ions

Ion trajectory tracing by classical MD



Trajectories of emitted ion are traced by classical MD.

MD conditions

- Coulombic repulsive force between ions is considered.
- Time step is 0.2 fs, total MD time is 10 ps.
- Calculate 1,000 ensembles for each condition.

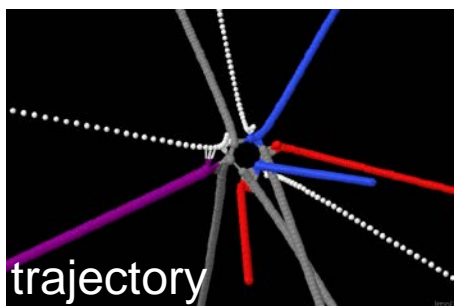
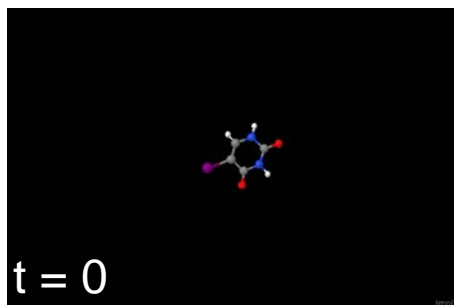
Charge distribution

- We adopt the charge distribution of CH_3I molecule for ion distribution of parent 5-I-Uracil molecule.
- **One charge is fixed to iodine** (absorber). Rest of charges are randomly distributed within constituent at $t=0$ of MD.
- Charges in ion are increased according to $(1 - \exp(-t/\tau))$ with **a charge build-up time constant τ** .

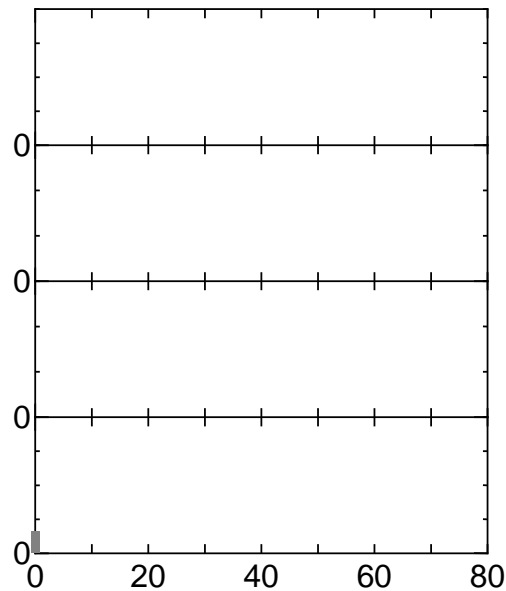
Initial structure of MD calculations

- Initial molecular structures of MD are prepared by B3LYP/3-21Gd level calculations by GAMESS package
- Thermally randomized structure and velocity distribution is estimated by GAMESS using Nose/Hoover method.

MD with $T = 300 \text{ K}$ and $t = 10 \text{ fs}$ reproduces experimental results well.



Influence of charge-build up time for KED in MD



Solid bars: experiment

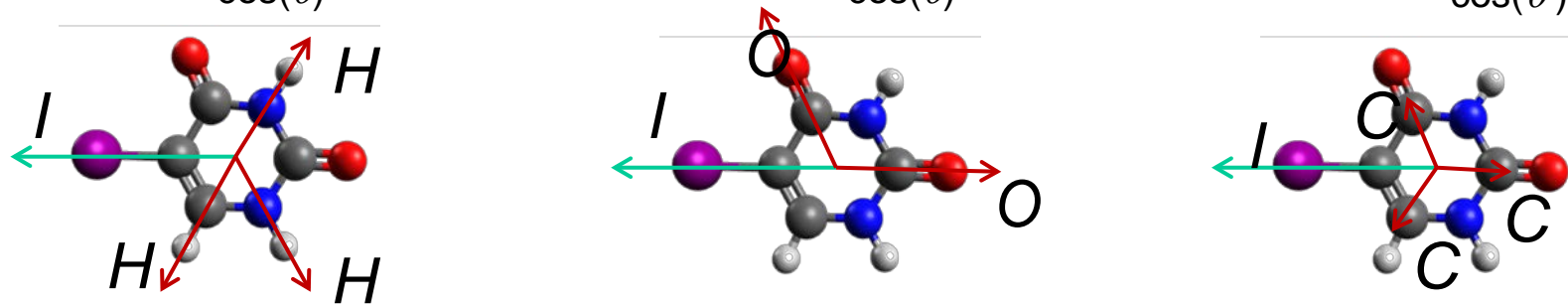
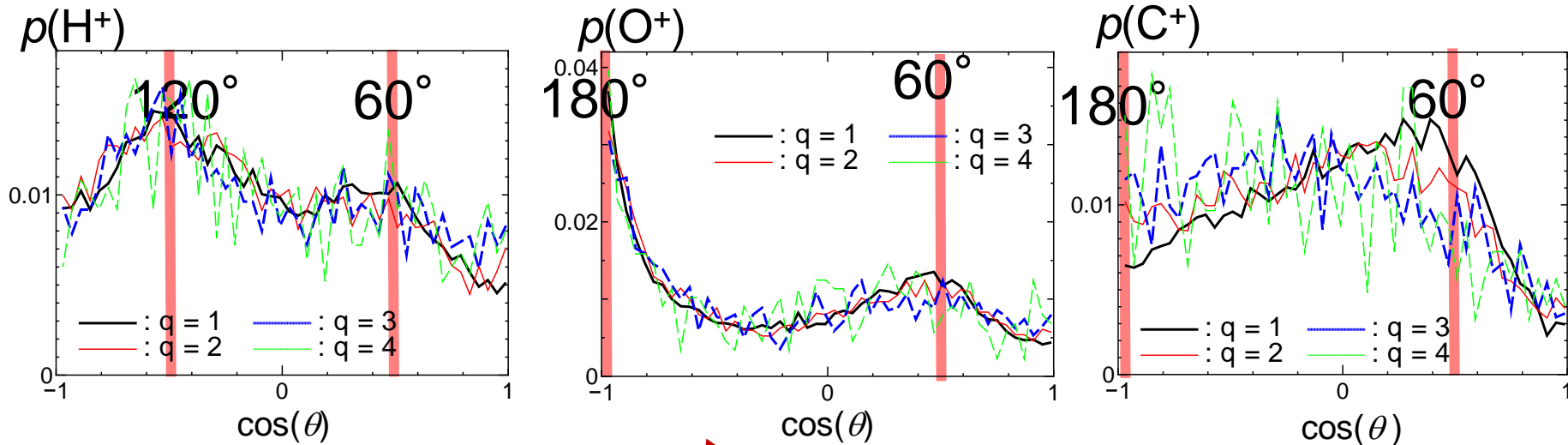
MD @ 300K

Dotted lines: $\tau = 0fs$

Solid lines: $\tau = 10fs$

2-body angular correlation

Black: I^+ , Red: I^2+ , Blue: I^3+ , Green: I^4+



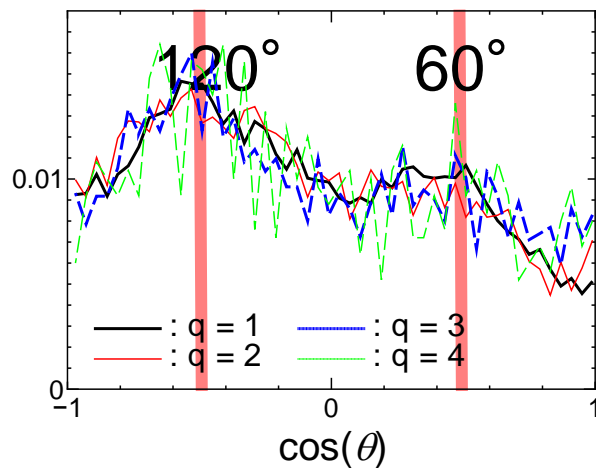
Red lines indicate angle between I direction and H, O, or C directions from the center of the aromatic ring in the neutral molecule.

Angular correlations of fragment ions (except carbon ions) reflect the shape of parent molecule.

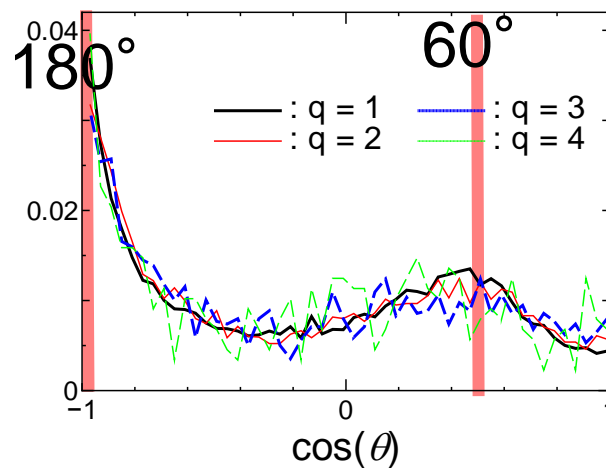
Comparison between Experiment and MD

2-body angular correlation

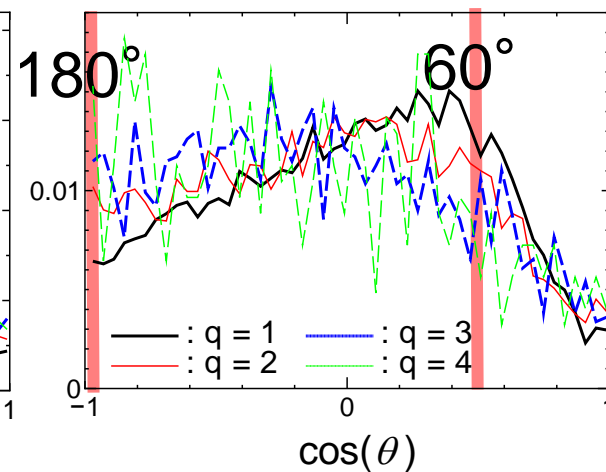
Exp. : $p(\text{H}^+)$



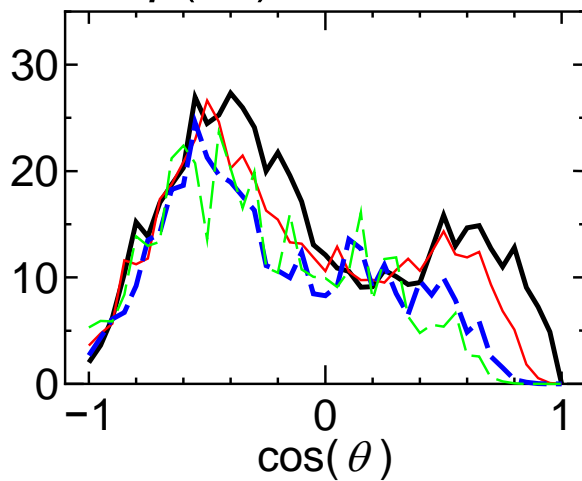
Exp. : $p(\text{O}^+)$



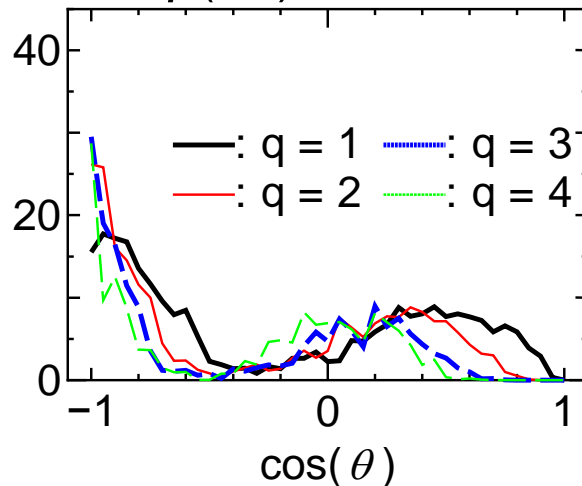
Exp. : $p(\text{C}^+)$



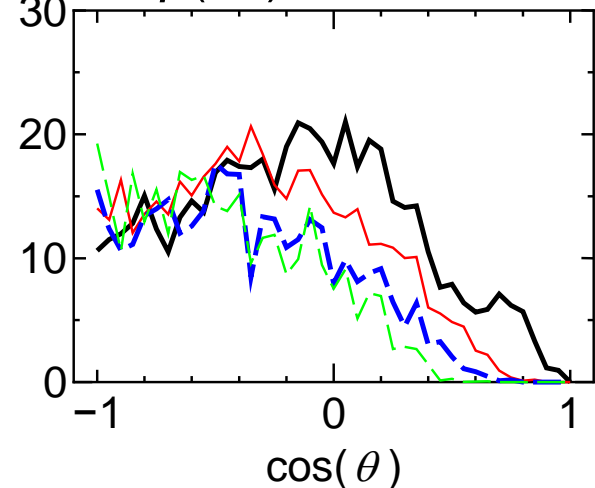
MD : $p(\text{H}^+)$



MD : $p(\text{O}^+)$



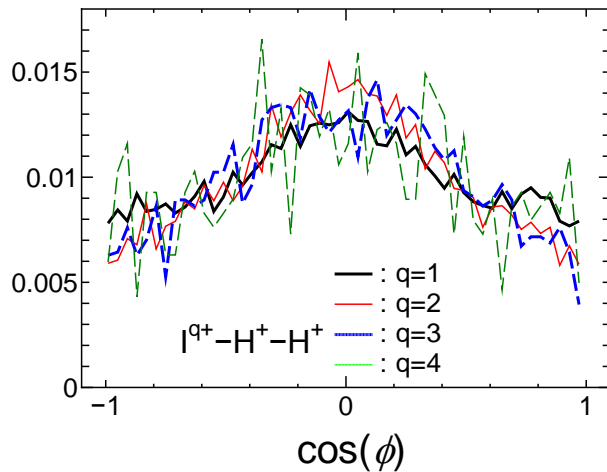
MD : $p(\text{C}^+)$



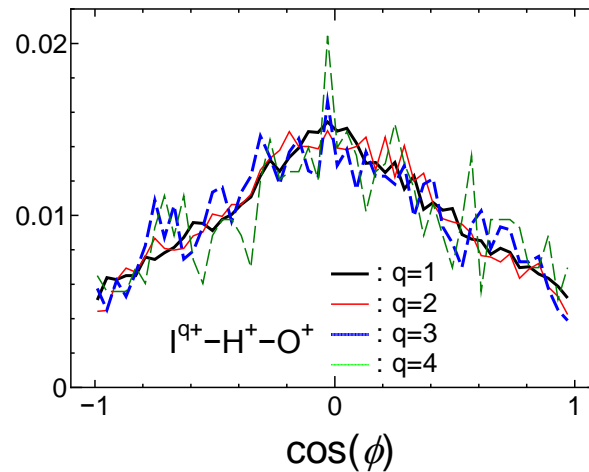
Black: I^+ , Red: I^{2+} , Blue: I^{3+} , Green: I^{4+}

3-body angular correlation

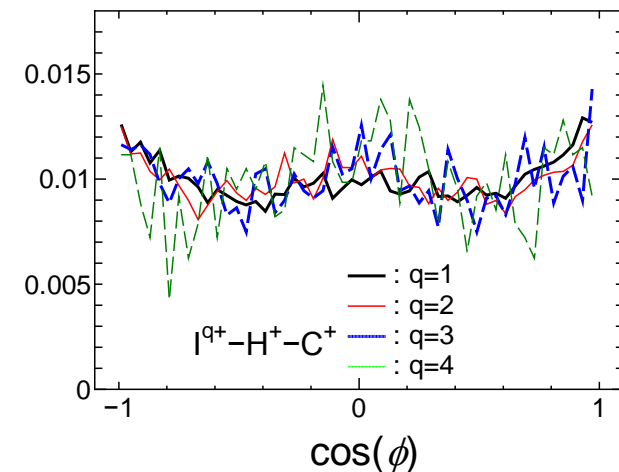
$I^+-H^+-H^+$



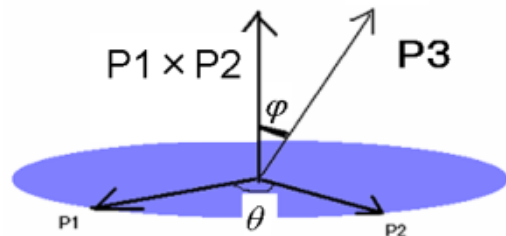
$I^+-H^+-O^+$



$I^+-H^+-C^+$



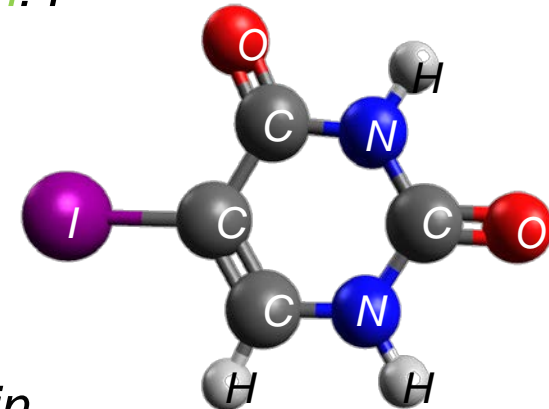
Black: I^+ , Red: I^{2+} , Blue: I^{3+} , Green: I^{4+}



$$\cos(\phi) = \frac{(\mathbf{P}_1 \times \mathbf{P}_2) \cdot \mathbf{P}_3}{(|\mathbf{P}_1| |\mathbf{P}_2| |\mathbf{P}_3| \sin(\theta))}$$

\mathbf{P}_1 , \mathbf{P}_2 and \mathbf{P}_3 are momentum of ions.

$\cos(\phi) = 0$ means that three momentum vectors are in the same plane.

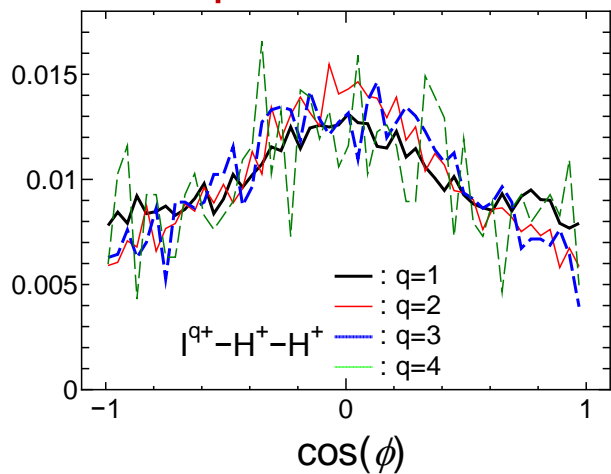


These results suggest that carbon ions are released to off-planar direction.

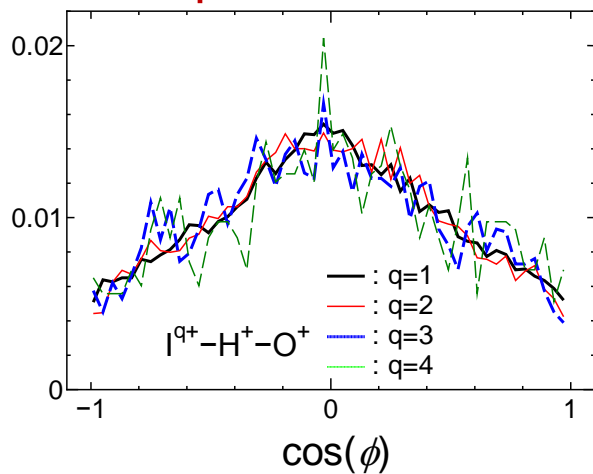
Comparison between Experiment and MD

3-body angular correlation

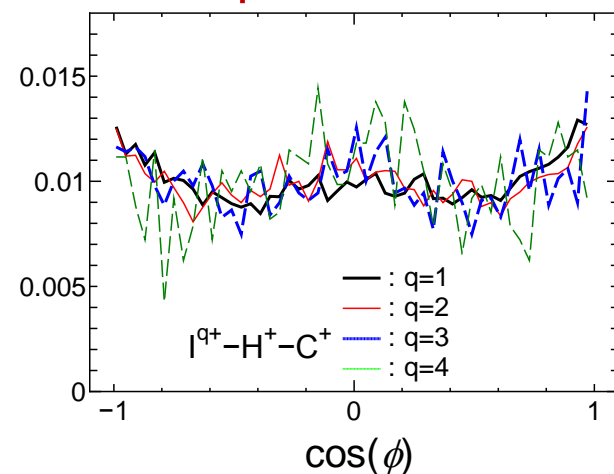
Exp. : $I^{q+}-H^+-H^+$



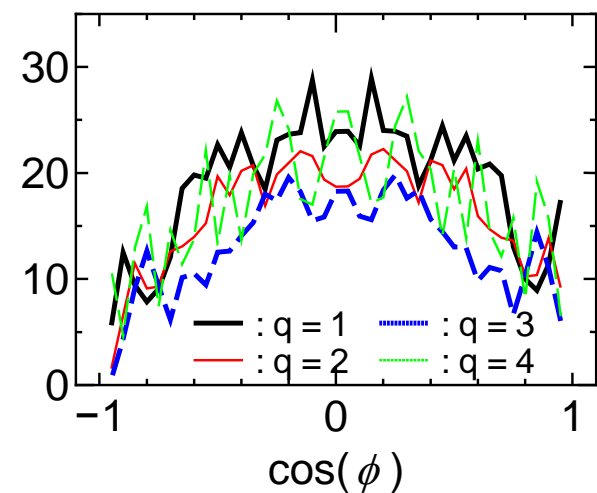
Exp. : $I^{q+}-H^+-O^+$



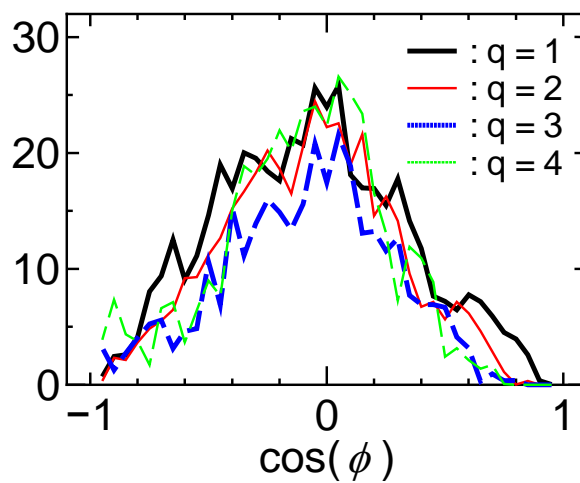
Exp. : $I^{q+}-H^+-C^+$



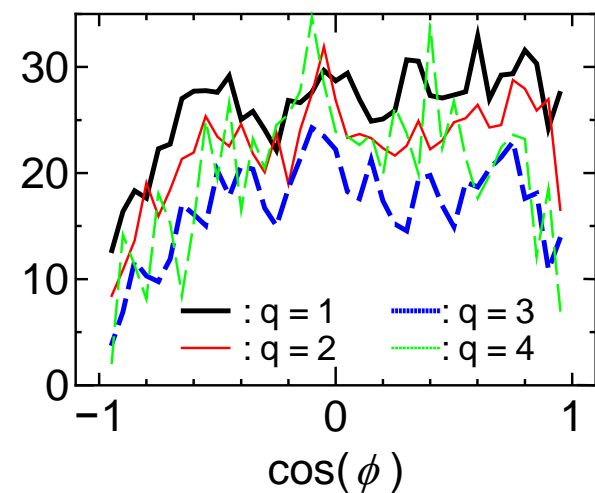
MD : $I^{q+}-H^+-H^+$



MD : $I^{q+}-H^+-O^+$



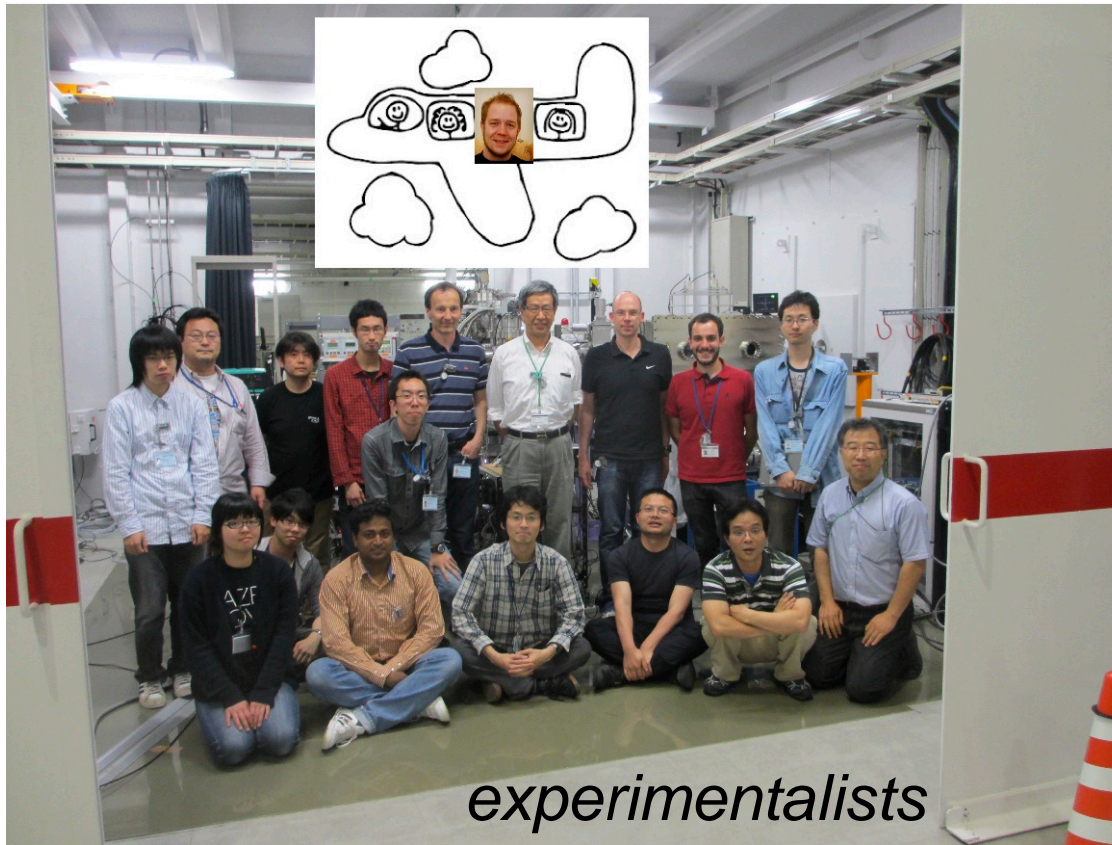
MD : $I^{q+}-H^+-C^+$



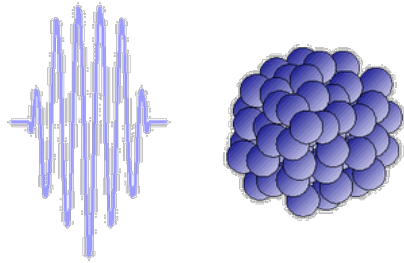
Black: I^+ , Red: I^2+ , Blue: I^3+ , Green: I^4+

III. Efficient Nanoplasma Formation from Argon Clusters Irradiated by the Hard X-ray Free Electron Laser

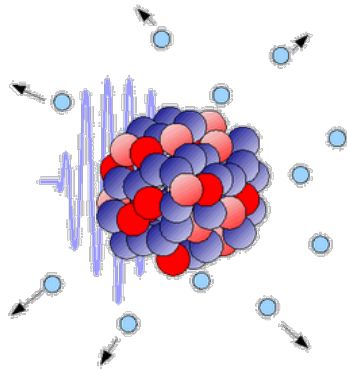
T. Tachibana, Z. Jurek, H. Fukuzawa, K. Motomura, K. Nagaya, S. Wada, P. Johnsson, M. Siano, S. Mondal, Y. Ito, M. Kimura, T. Sakai, K. Matsunami, H. Hayashita, J. Kajikawa, X.-J. Liu, E. Robert, C. Miron, R. Feifel, J. Marangos, K. Tono, T. Togashi, Y. Inubushi, T. Hatsui, M. Yabashi, B. Ziata, S. Son, M. Yao, R. Santra, and K. Ueda (*Scientific Reports; in press*).



Nanoplasma formation by intense laser irradiation

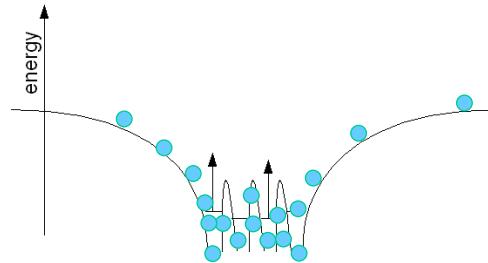
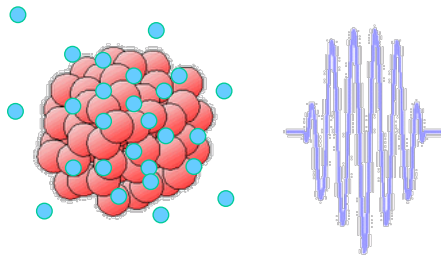


Laser irradiation into cluster



Many atoms in the cluster are ionized

Nanoplasma is formed when the electrons ejected from atoms trapped by the Coulomb potential of the multiply charged cluster ion.



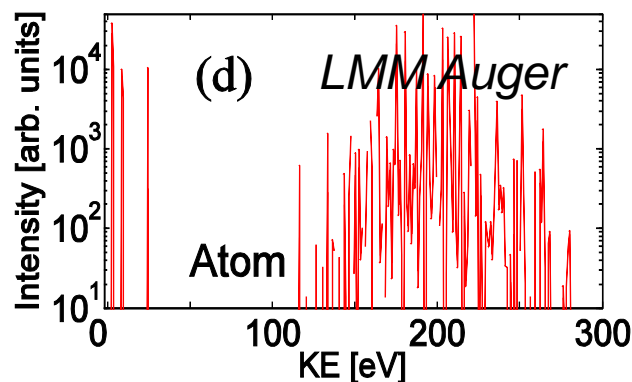
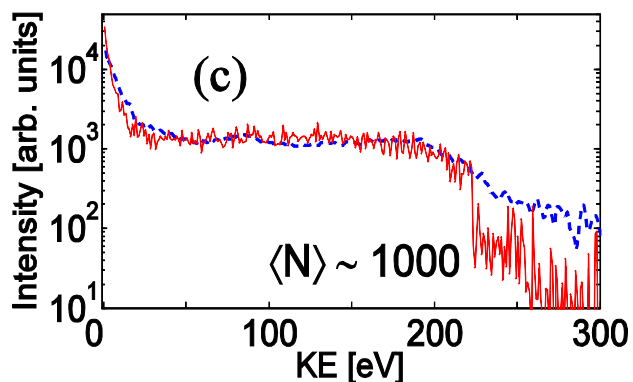
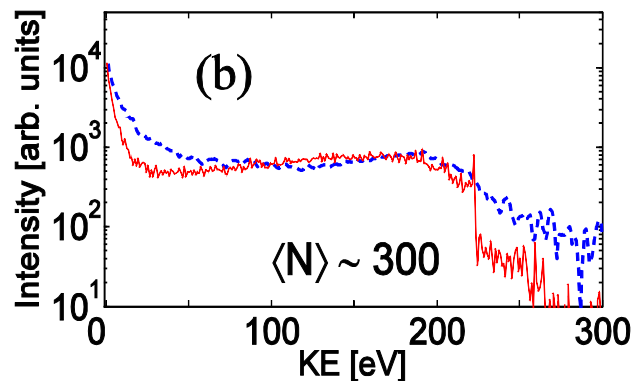
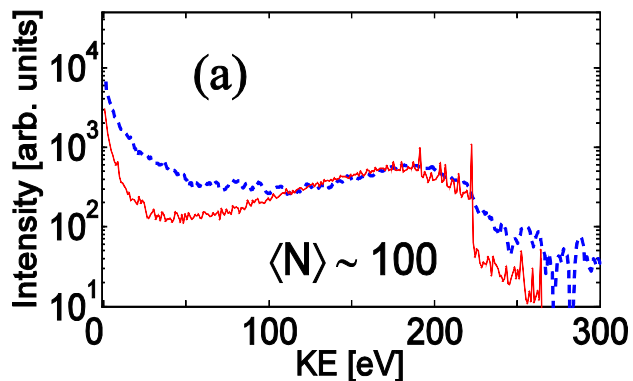
Is nanoplasma also formed by intense hard x-ray pulse irradiation?

How is nanoplasma formed?

Experimental & theoretical electron spectra of Ar clusters

Photon energy: 5 keV
Ar K edge: 3.2 keV

Dashed blue: Experiment,
Solid red: Theory

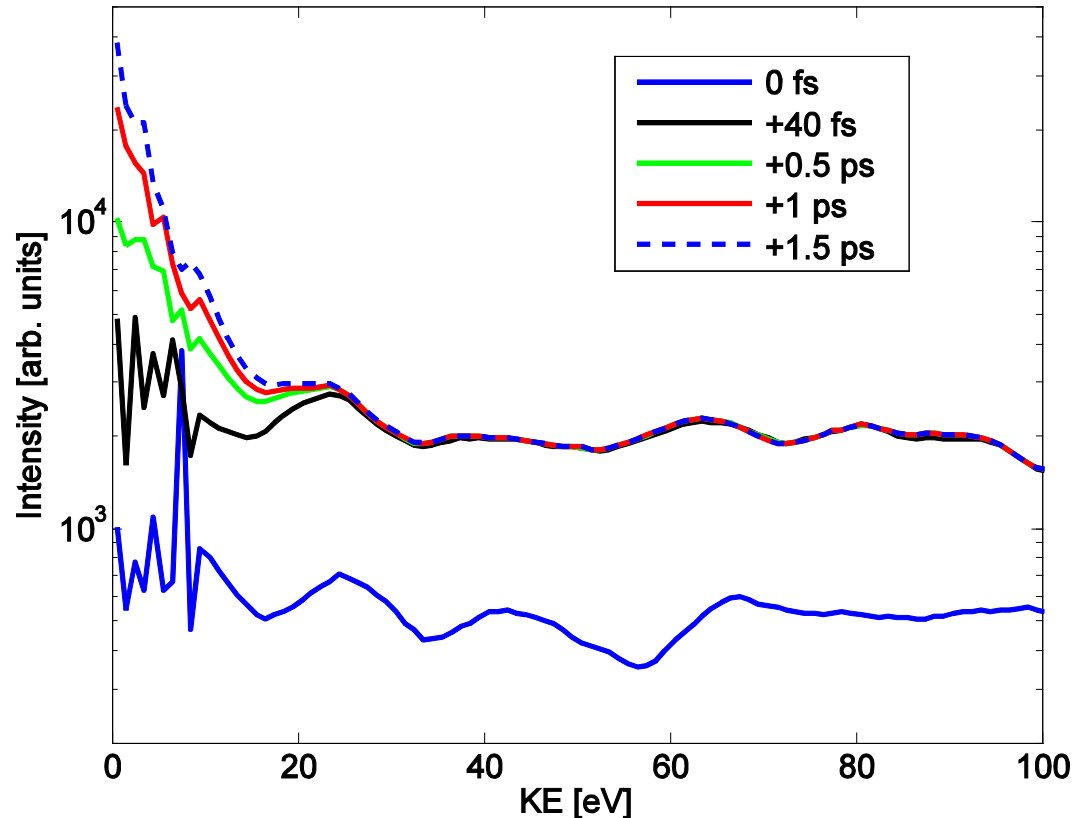


A plateau in the spectra is produced by the deceleration of the electrons. With the increase of the cluster size, a stronger potential builds up, decelerating the emitted electrons more.

The strong peak at zero kinetic energy is due to the thermal emission from nanoplasma.

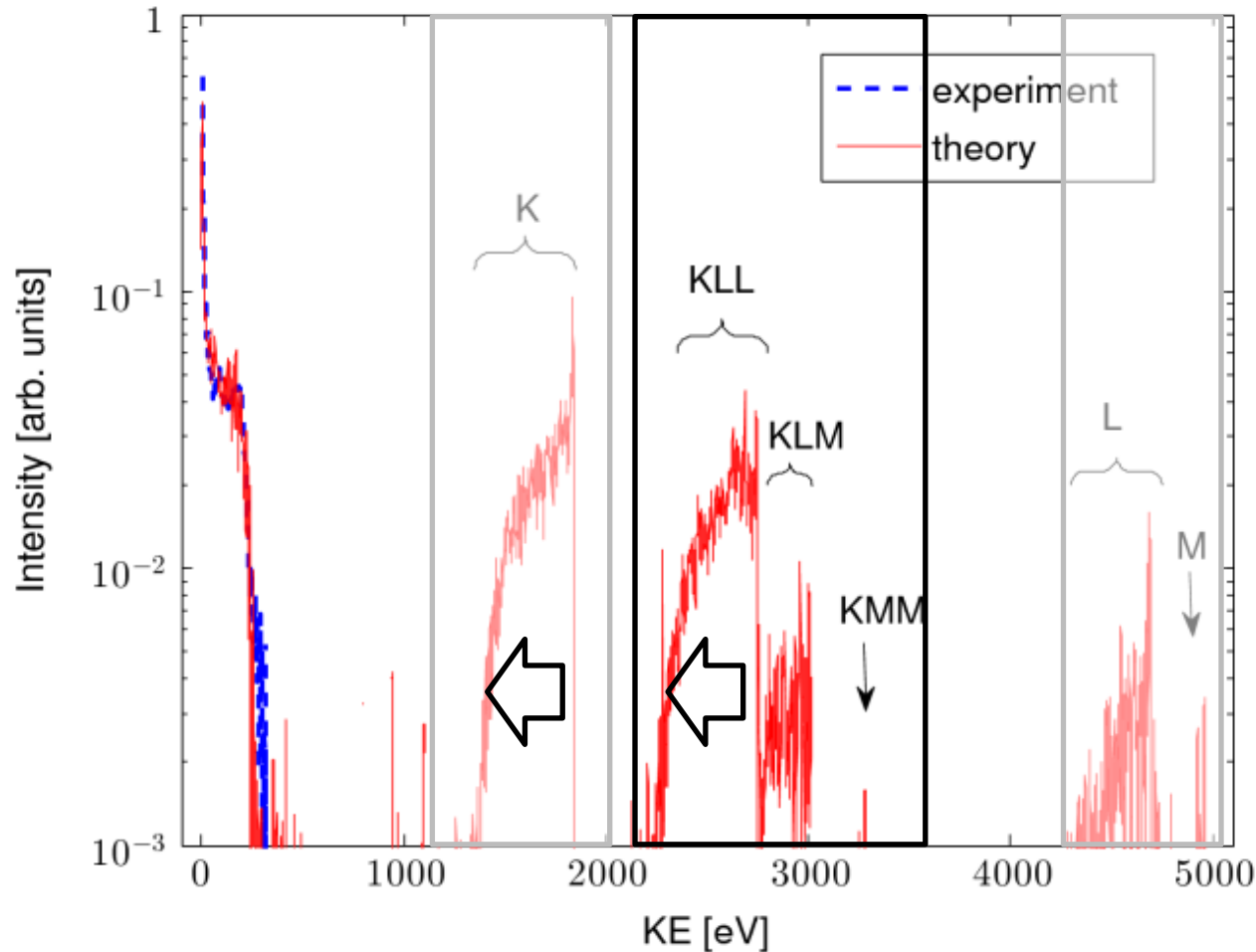
T. Tachibana, Z. Jurek, H. Fukuzawa et al., Scientific Reports (in press).

Time evolution for the theoretical electron spectrum



During the XFEL pulse, only the plateau is formed. The main peak at 0 eV develops after the XFEL pulse when the ionic system has started to expand and let some of trapped electrons escape from nanoplasma.

Electron spectra of Ar₁₀₀₀ in the whole region

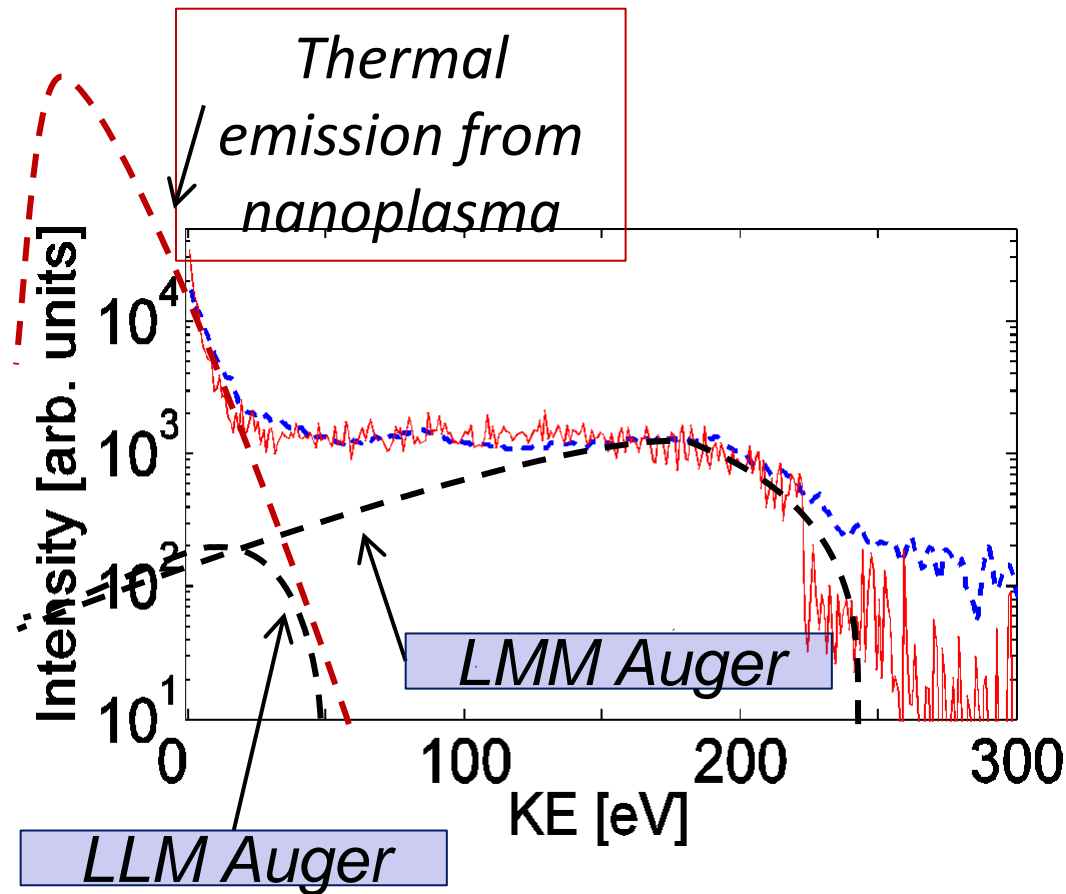


Photon energy : 5 keV
Ar K edge : 3.2 keV

Photo electron
Auger electron

Emitted electrons are decelerated ~500 eV

Origin of the slow electrons that can be trapped

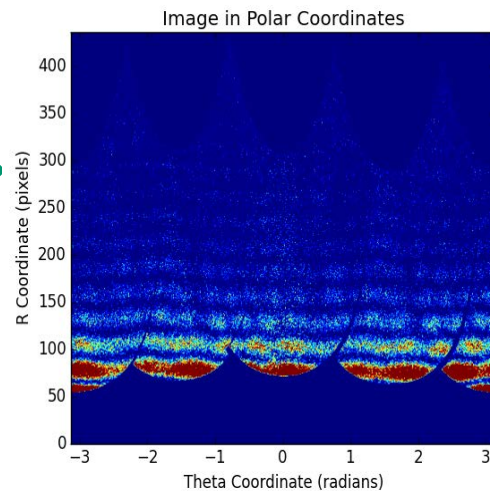
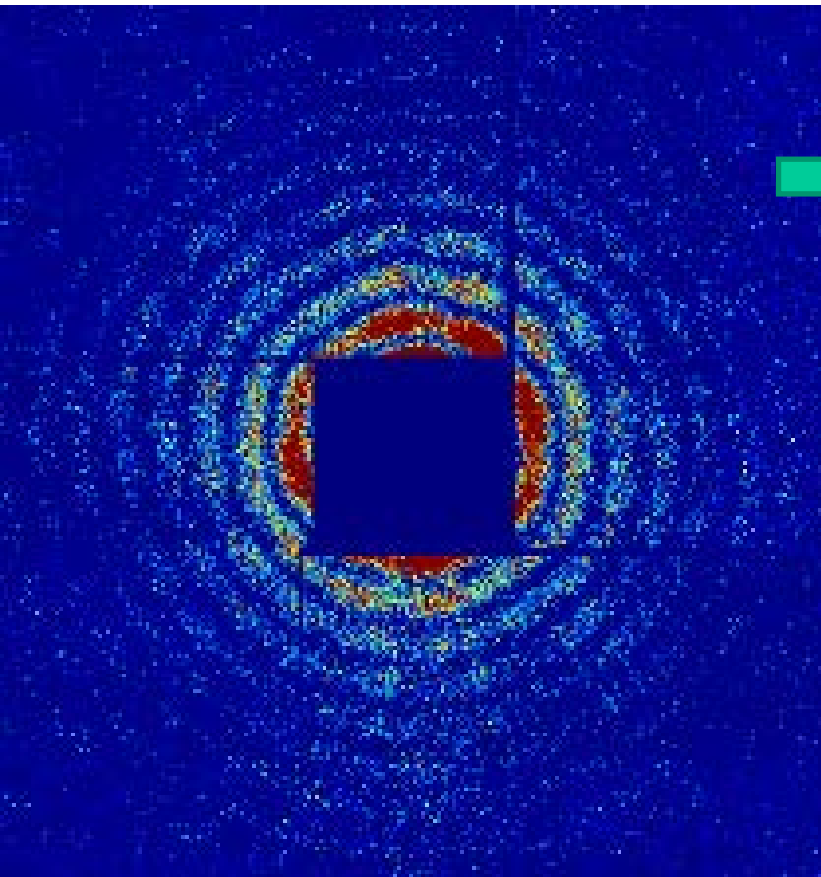


The majority of trapped electrons are created by impact ionizations caused by low-energy Auger and secondary electrons.

IV. Single-shot imaging of giant Xe clusters with X-ray free-electron laser (5.5 keV)



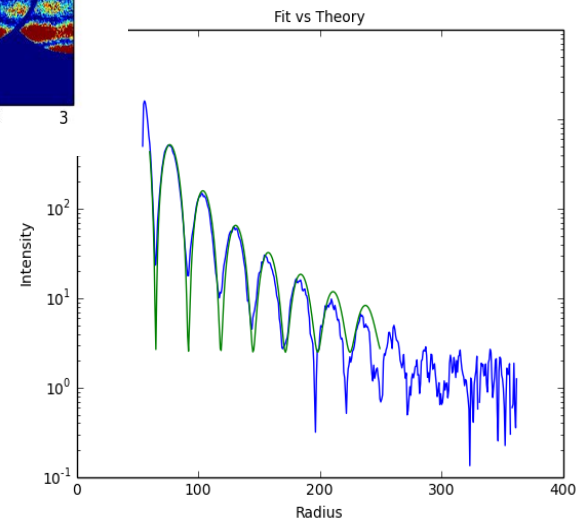
T. Nishiyama, C. Bostedt, K. Nagaya, K. R. Ferguson, C. Hutchison, H. Fukuzawa, K. Motomura, S. Wada, T. Sakai, K. Matsuami, T. Tachibana, Y. Ito, W. Q. Xu, S. Mondal, T. Umemoto, C. Nicolas, C. Miron, K. Kameshima, Y. Jochi, K. Tono, H. Hatsui, M. Yabashi, M. Yao, and K. Ueda (in preparation).



Diameter: 250nm
 $\sim 1.1 \times 10^8$ atoms

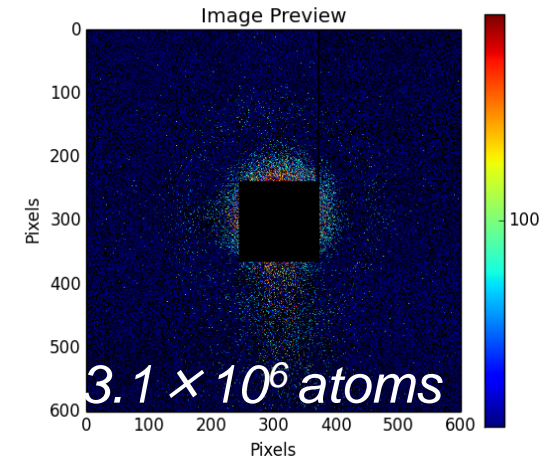
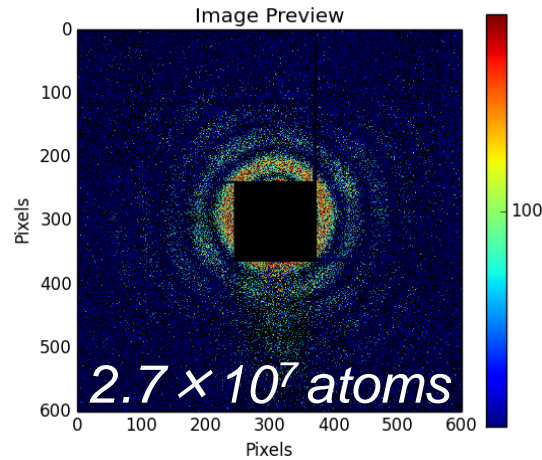
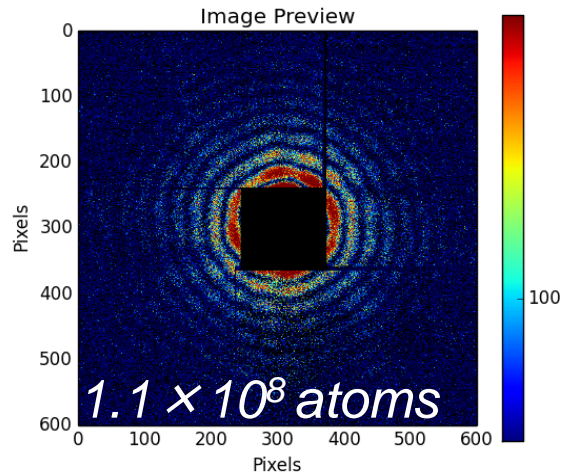


FEL fluence:
 $\sim 10 \mu\text{J}/\mu\text{m}^2$

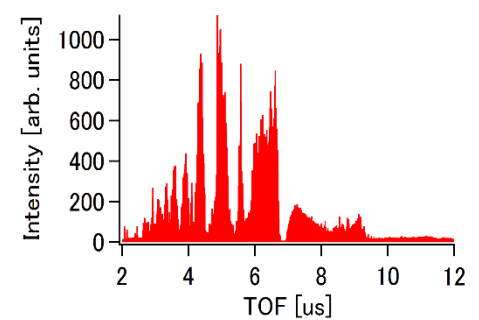
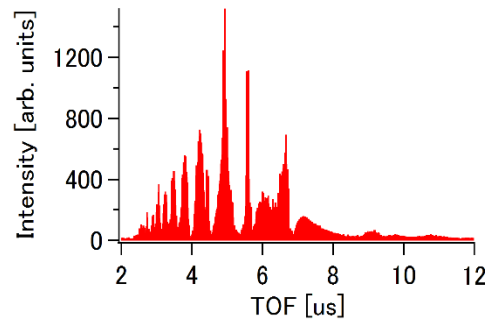
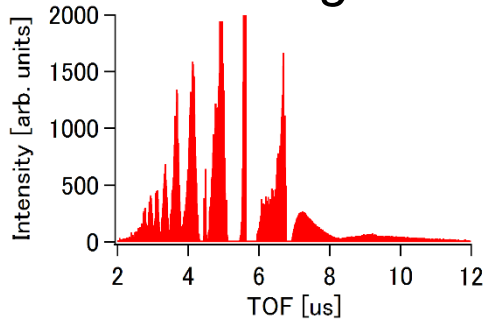


IV. Single-shot imaging of Xe clusters with XFEL

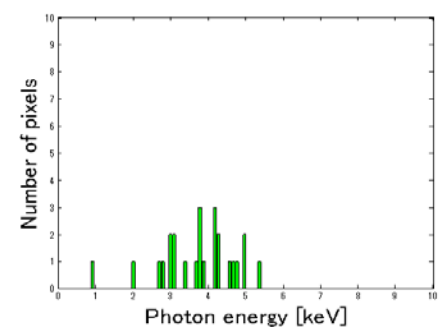
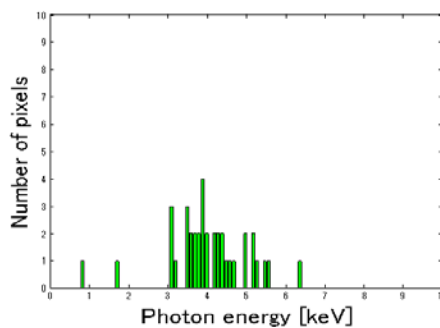
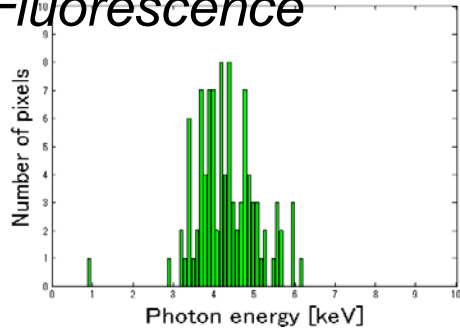
5.5 keV with fluence of $\sim 10 \mu\text{J}/\mu\text{m}^2$



Ion Time of Flights

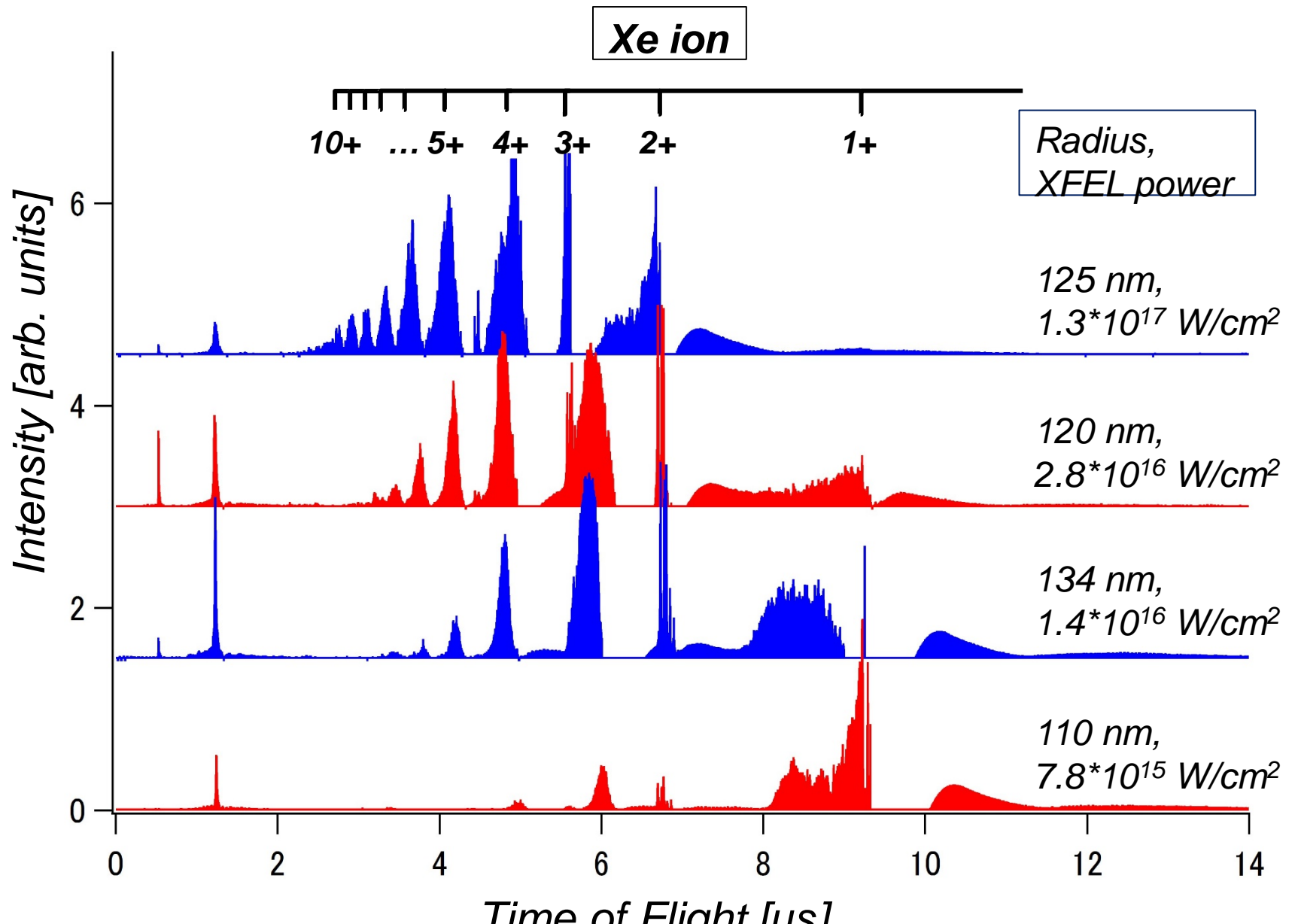


Fluorescence



IV. Single-shot imaging of Xe clusters with XFEL

5.5 keV with the cluster size of ~ 120 nm



V. Real-time study on the ultrafast plasmon resonance heating of nanoplasma produced by the XFEL irradiation to rare gas clusters at SACLA

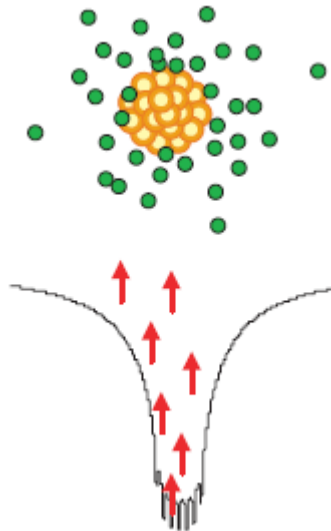
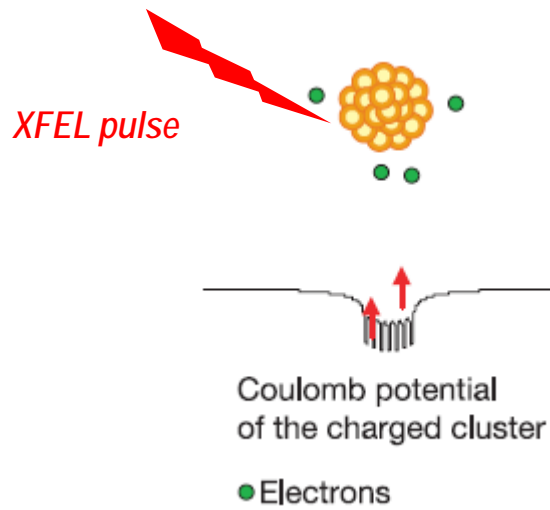


Y. Kumagai, W. Xu, Z. Jurik, H. Fukuzawa, K. Motomura, S. Mondal, T. Tachibana, Y. Ito, K. Nagaya, T. Sakai, K. Matsunami, T. Nishiyama, M. Yao, S. Wada, T. Umemoto, C. Nicolas, C. Miron, T. Togashi, K. Tono, Ogawa, S. Owada, M. Yabashi, B. Ziata, S. Son, R. Santra, and K. Ueda
(in preparation)

Dynamics of nanoplasma produced by XFEL

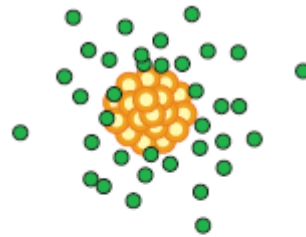
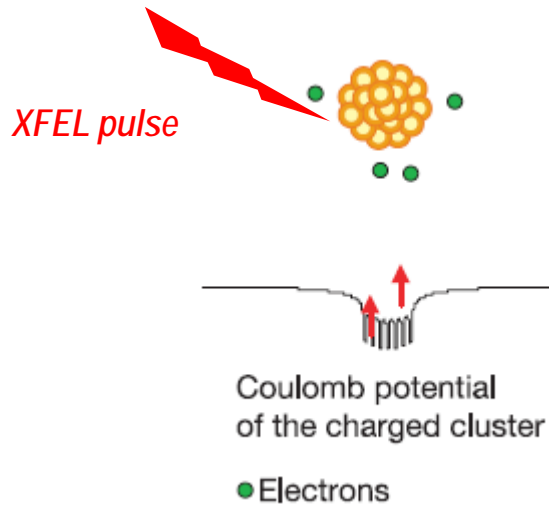
- **Photoionization**
Auger cascades

- **Nanoplasma formation**

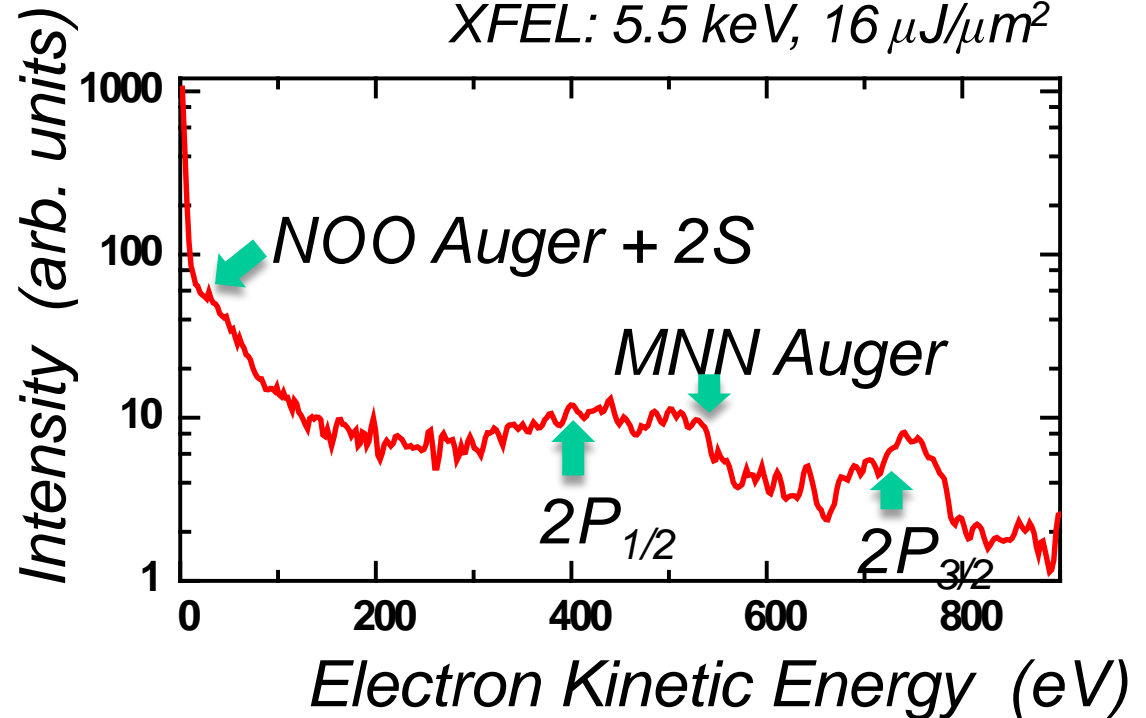


Dynamics of nanoplasma produced by XFEL

- Photoionization
- Nanoplasma formation
- Auger cascades

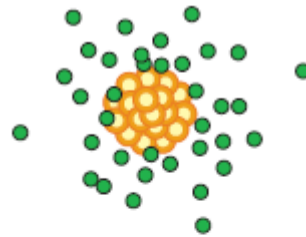
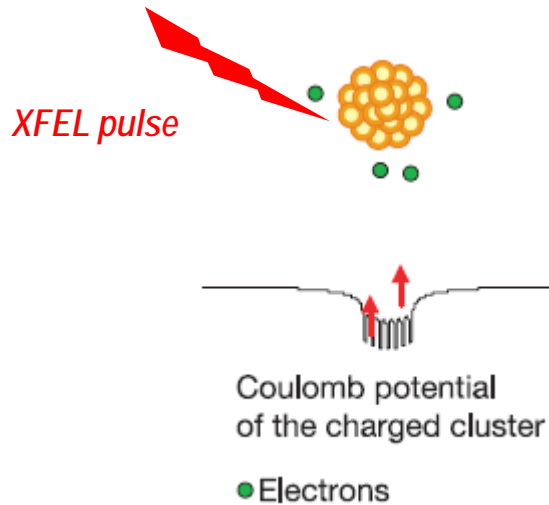


Xe clusters
 $\langle N \rangle \sim 5000$ @ SACLA
XFEL: 5.5 keV, $16 \mu\text{J}/\mu\text{m}^2$

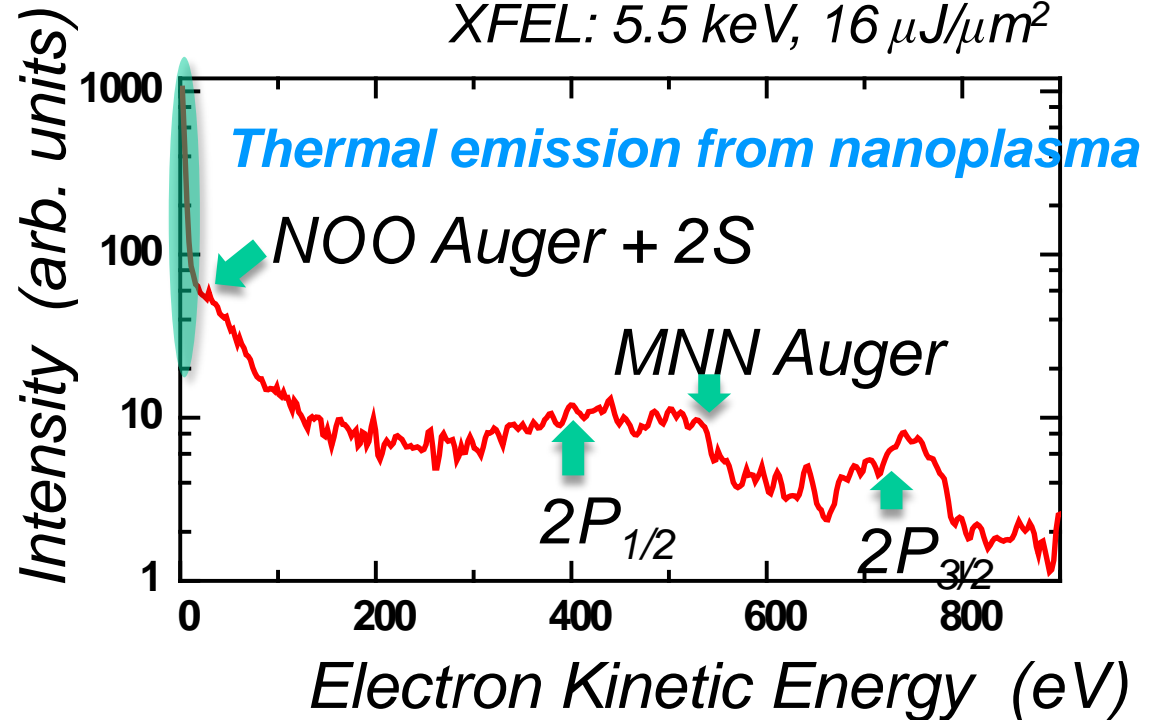


Dynamics of nanoplasma produced by XFEL

- **Photoionization**
- **Nanoplasma formation**
- **Auger cascades**



Xe clusters
 $\langle N \rangle \sim 5000$ @ SACLA
XFEL: 5.5 keV, $16 \mu\text{J}/\mu\text{m}^2$

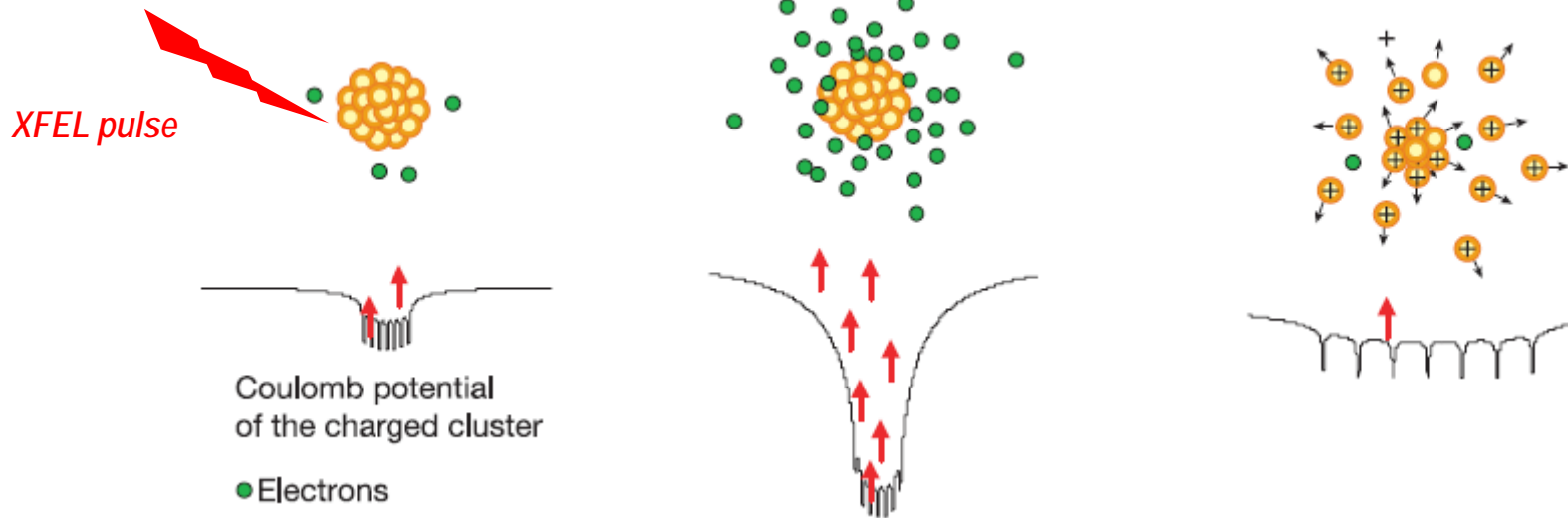


Dynamics of nanoplasma produced by XFEL

• **Photoionization**
Auger cascades

• **Nanoplasma formation**

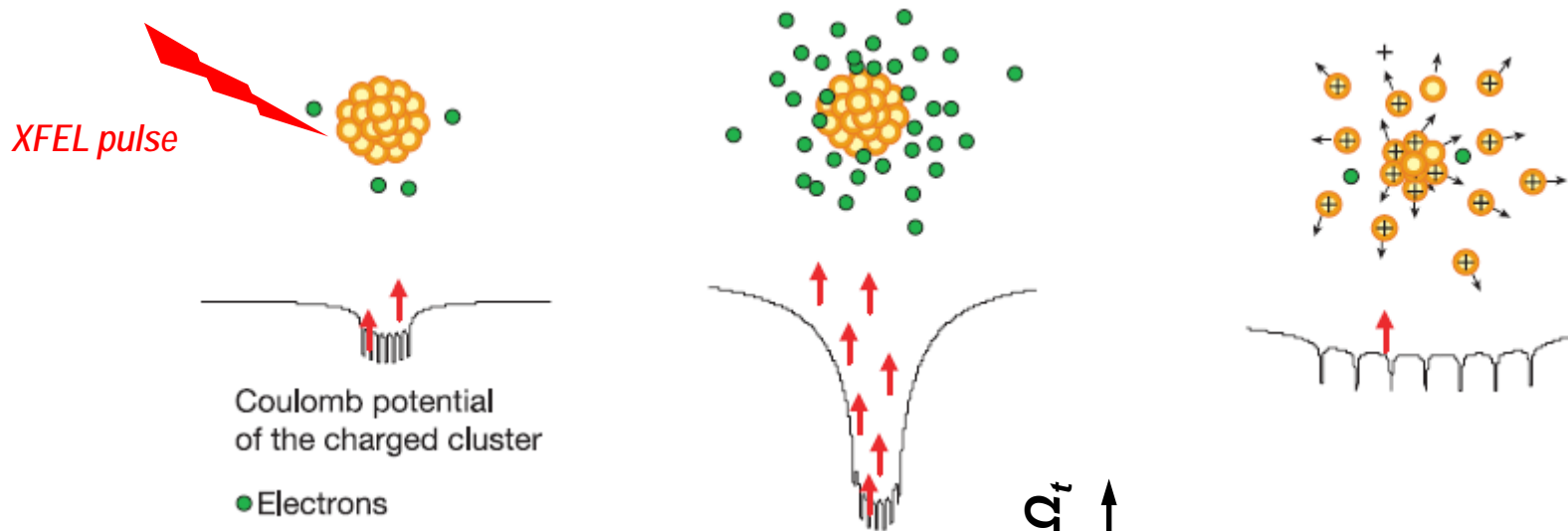
• **Coulomb explosion**



• • • *To probe nanoplasma formation and to investigate its dynamics we employ pump-probe technique*

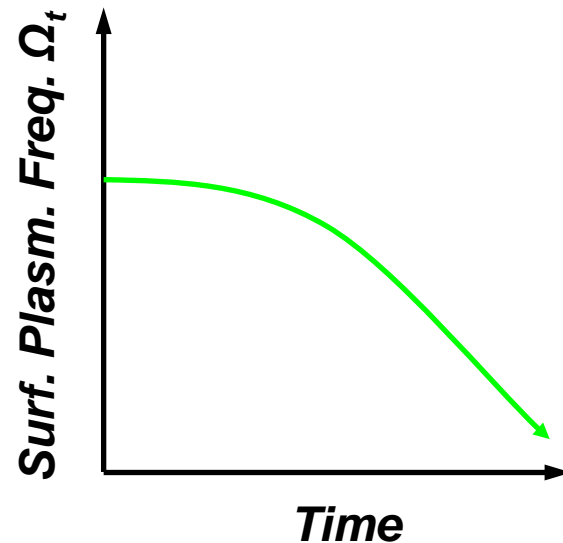
To search the plasmon resonance heating

- **Photoionization**
 - **Nanoplasma formation**
 - **Coulomb explosion**
- Auger cascades**



Surface-plasma frequency

$$\Omega_t = \sqrt{\frac{N_t Z_t}{R_t^3}} = \frac{\omega_{pl}}{\sqrt{3}}$$

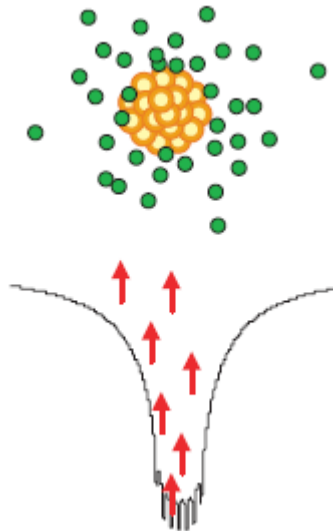
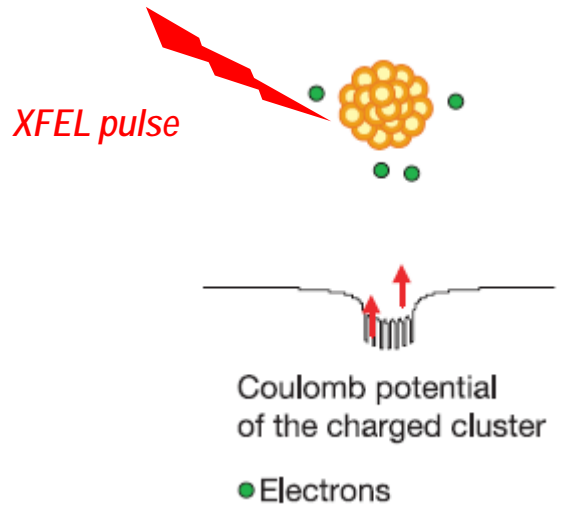


To search the plasmon resonance heating

**Photoionization
Auger cascades**

**Nanoplasma
formation**

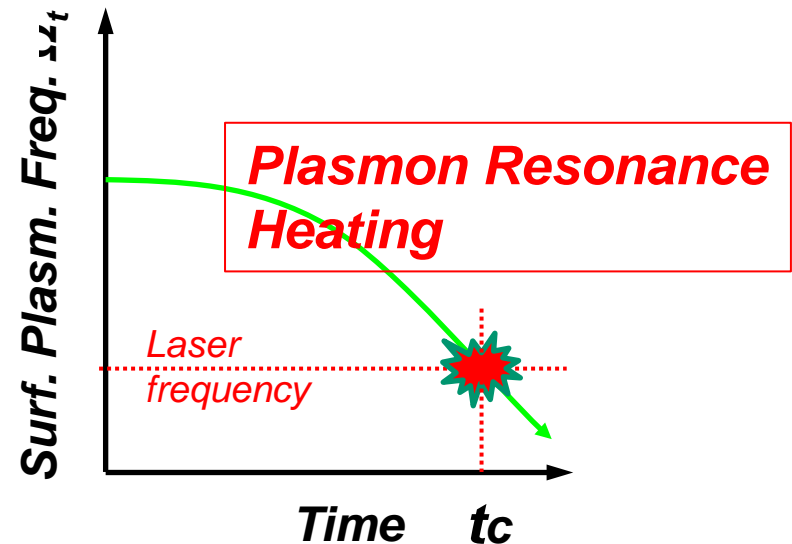
Coulomb explosion



$$\Omega_t \propto \sqrt{\frac{N_t Z_t}{R_t^3}} = \frac{\omega_{pl}}{\sqrt{3}}$$

Surface-plasma frequency

$$\Omega_t \propto \sqrt{\frac{N_t Z_t}{R_t^3}} = \frac{\omega_{pl}}{\sqrt{3}}$$



Experimental setups

■ XFEL @ SACLA

Photon energy: 5.5 keV

Spectral width: ~ 33 eV (FWHM)

Repetition rate: 30 Hz

Pulse duration: < 10 fs

Focused beam size: ~ 1 μm

Peak intensity: $\sim 3.28 / 3.38 \times 10^{16}$ W/cm²

■ NIR laser

Wavelength: 800 nm

Repetition rate : 30 Hz

Pulse duration : 80 fs

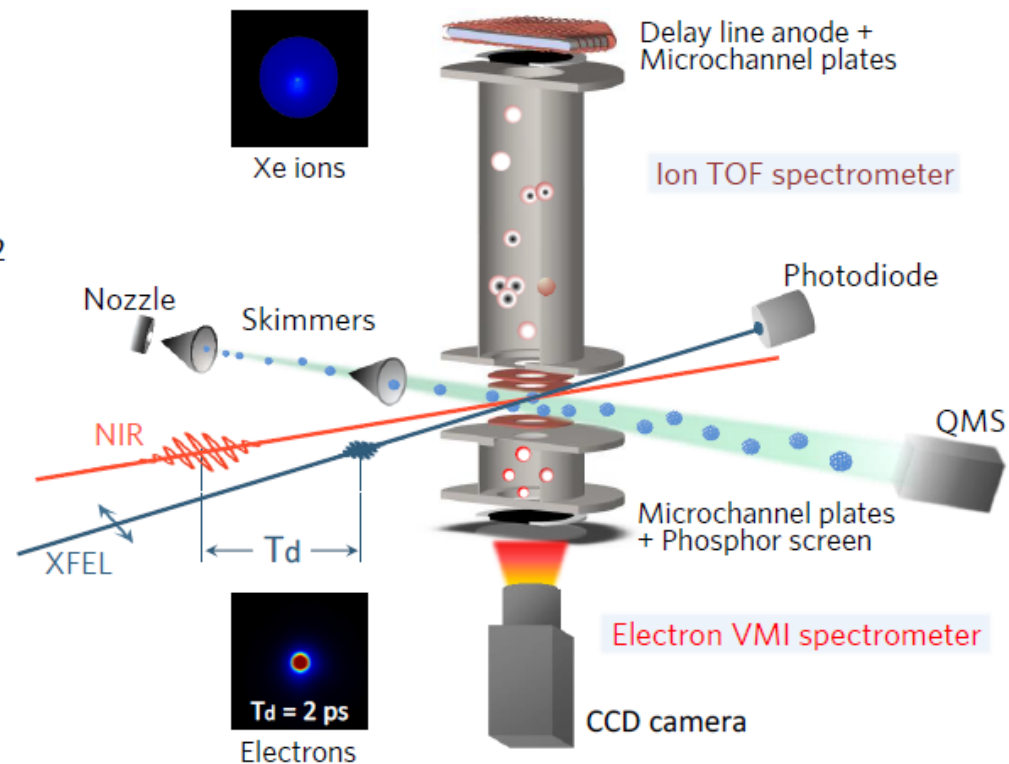
Focused beam size: ~ 200 μm

Intensity: $\sim 2.7 / 5.1 \times 10^{12}$ W/cm²

■ Clusters

Size of Xe clusters: ~ 5000 atoms

Nozzle diameter: ~ 250 μm



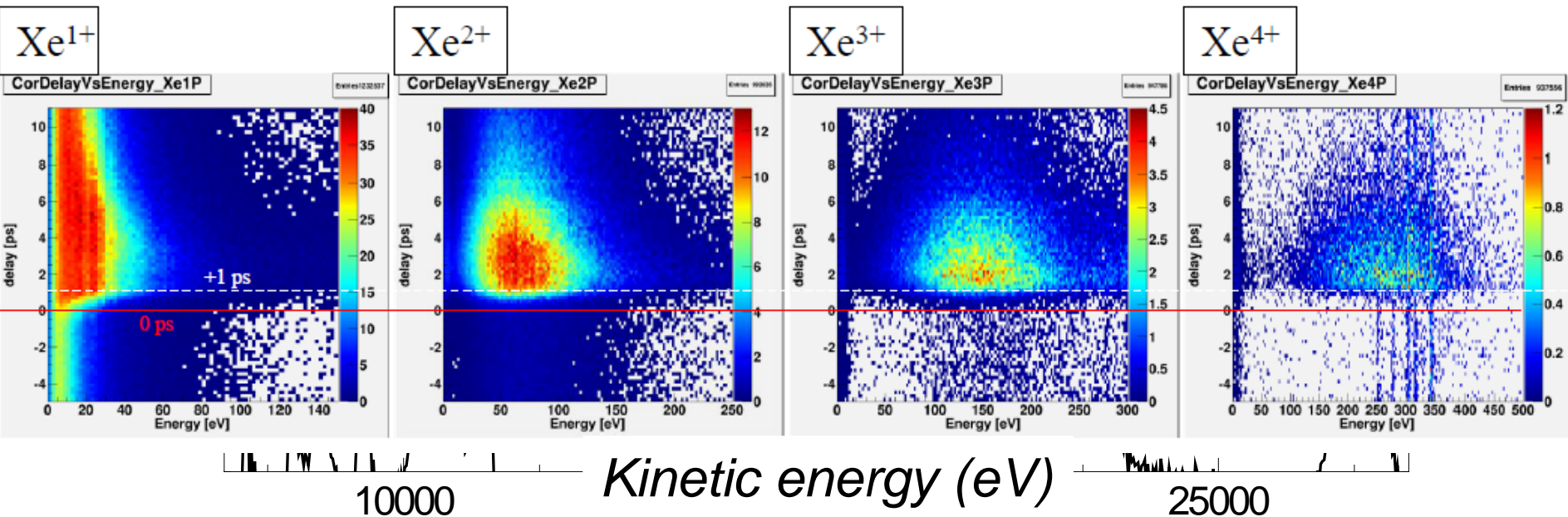
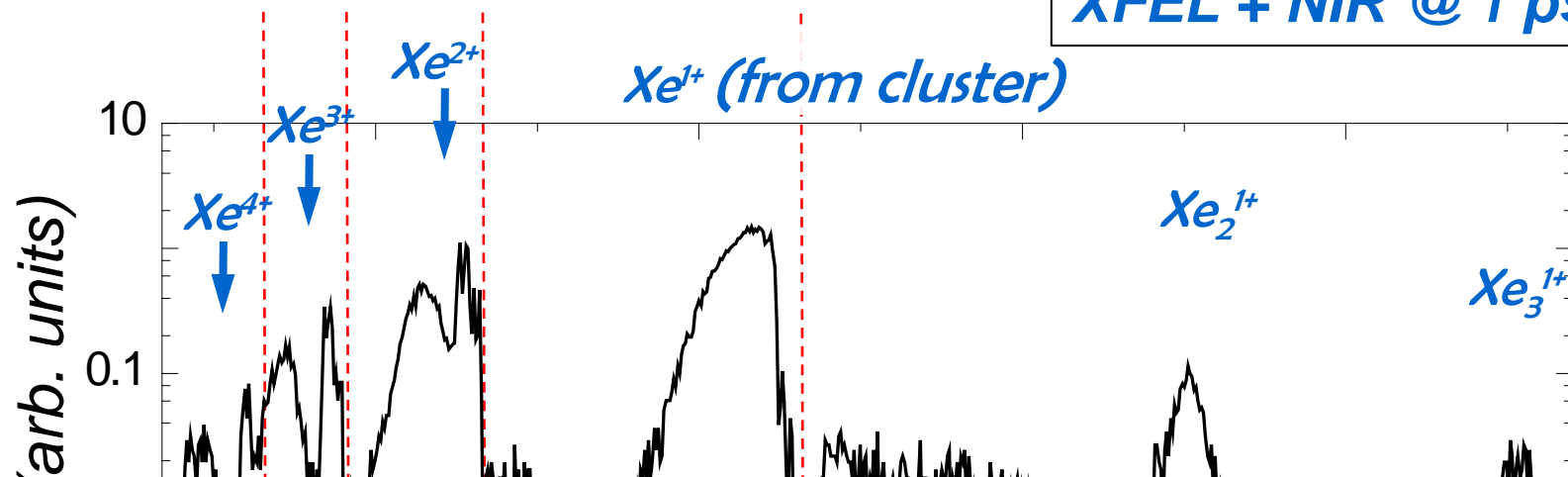
K. Tono *et al.*, *New J. Phys.* **15**, 083035 (2013);

A. T. J. B. Eppink and D. H. Parker, *Rev. Sci. Instrum.* **68**, 3477 (1997); K. Motomura *et al.*, *J. Phys. B.* **46**, 164024 (2003)

Time-resolved TOF spectra of Xe clusters

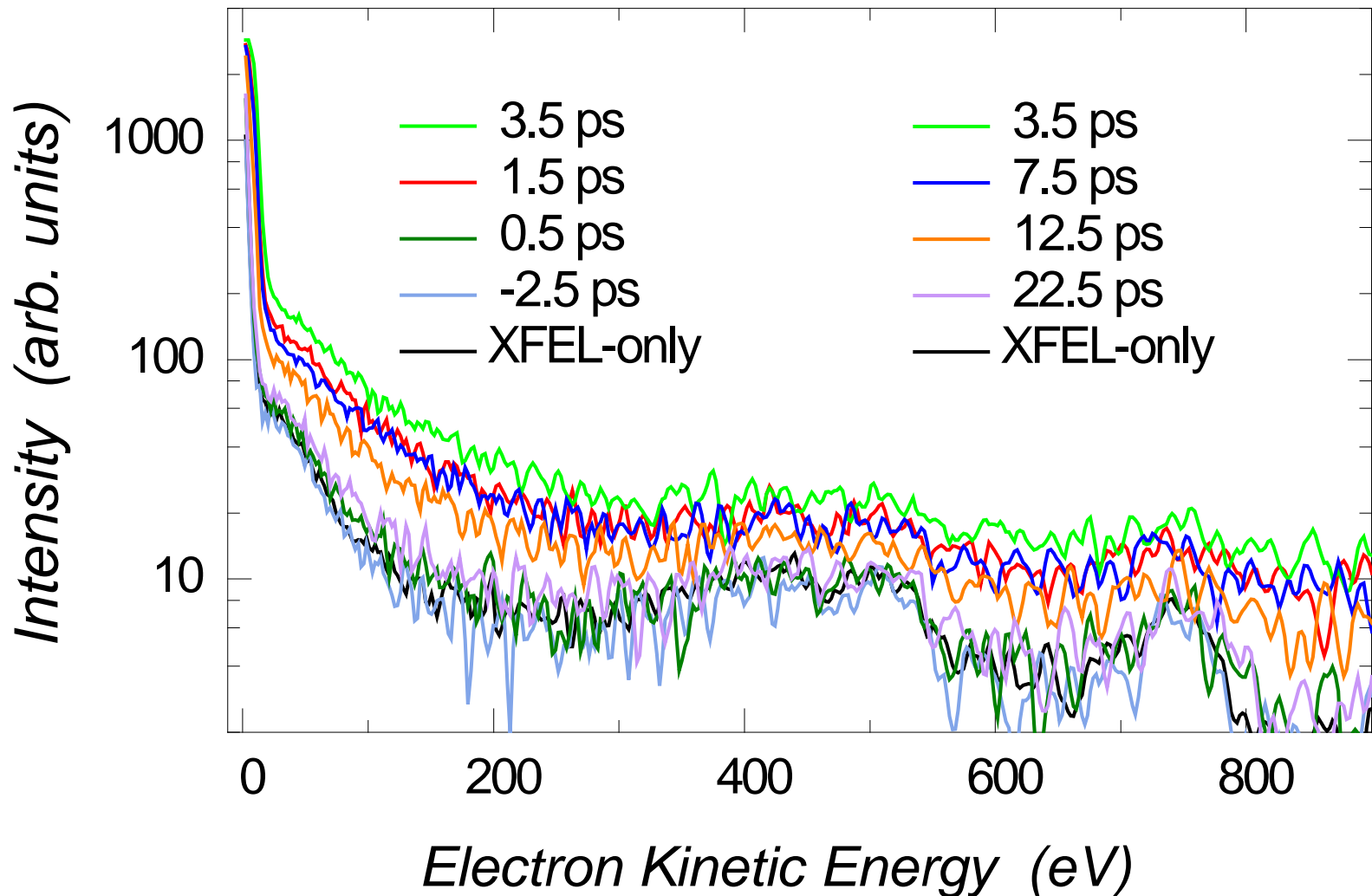
XFEL: $16 \mu\text{J}/\mu\text{m}^2$; NIR: $4.1 \text{ nJ}/\mu\text{m}^2$

XFEL + NIR @ 1 ps delay



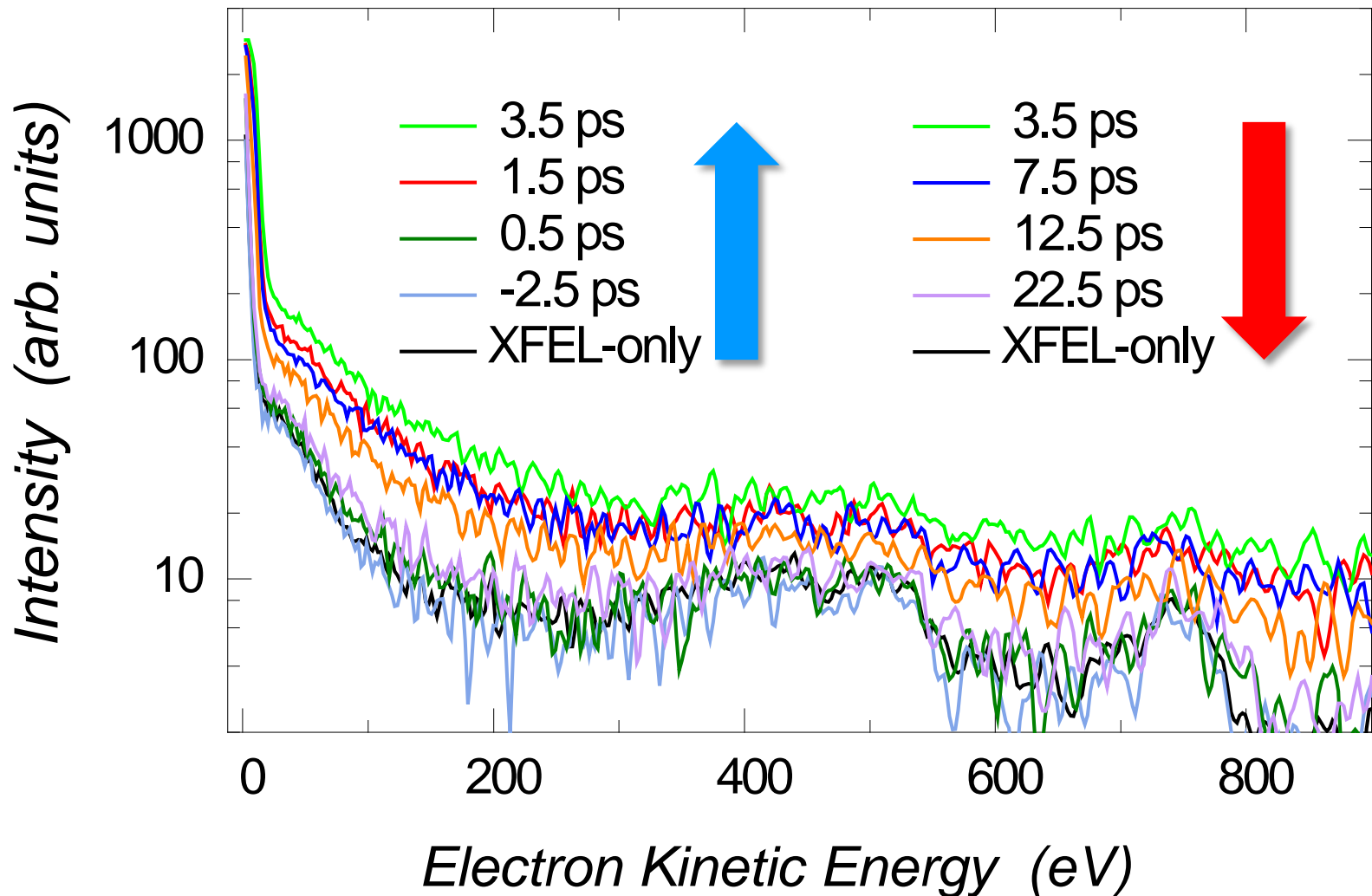
Plasmon heating in time-resolved PES

XFEL: 16 $\mu\text{J}/\mu\text{m}^2$; NIR: 4.1 $\text{nJ}/\mu\text{m}^2$



Plasmon heating in time-resolved PES

XFEL: 16 $\mu\text{J}/\mu\text{m}^2$; NIR: 4.1 $\text{nJ}/\mu\text{m}^2$



Summary

- Deep inner-shell multi-photon absorption of Ar and Xe atoms by **SACLA** XFEL pulses – **Electronic damage**
- Photoion-photoion coincidence imaging following deep inner-shell multi-photon absorption by **SACLA** XFEL pulses (CH₃I, 5I-uracil) – **Radiation damage in the atomic level**
- Electron spectroscopy of argon and xenon clusters heated by **SACLA** XFEL pulses – **Nanoplasma formation**
- Single-shot imaging of xenon nano-clusters – **Influence of the electronic damage to the imaging**
- IR-probe experiment of XFEL induced nanoplasma formation – **Nanoplasma dynamics**

What's next...

- UV/XFEL pump – XFEL probe for I-contained molecules
Intra-molecular charge transfer via ionic fragmentation
- X-ray imaging for UV/XFEL induced nanoplasma from giant clusters
Influence of nanoplasma formation to X-ray imaging....
- UV pump – XFEL probe for photocatalytic molecules
Intra-molecular charge transfer via TR X-ray spectroscopy
- Serial femtosecond X-ray crystallography
Phasing vs radiation damage

The end

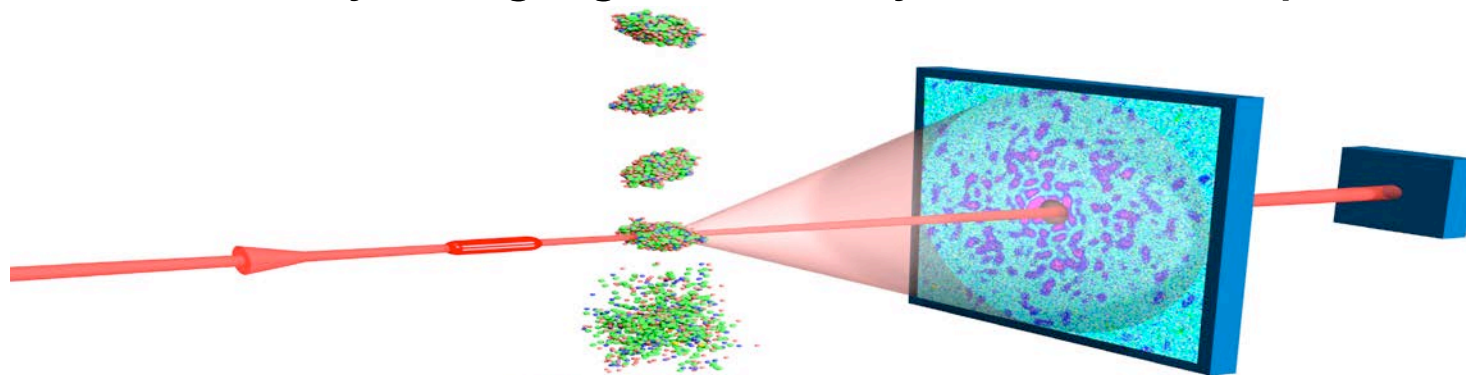


*Thank you very much for
your attention!*

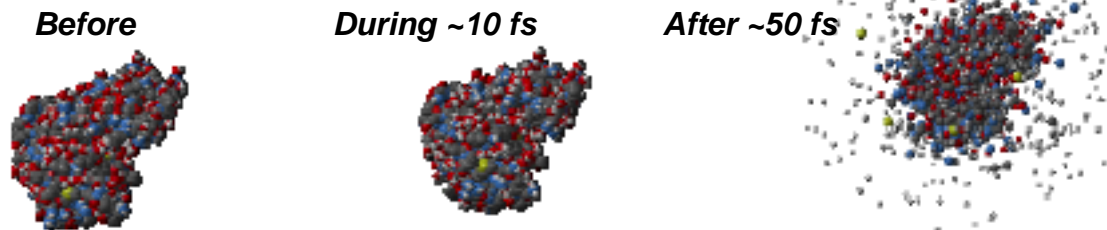
Characteristic properties of FEL pulses

*Coherent, intense, and ultra-short pulses
at short wavelengths (EUV to X-rays)*

Coherent X-ray imaging of non-crystallized samples



Gösta Huldt, Abraham Szöke, Janos Hajdu (J.Struct Biol, 2003 02-ERD-047)

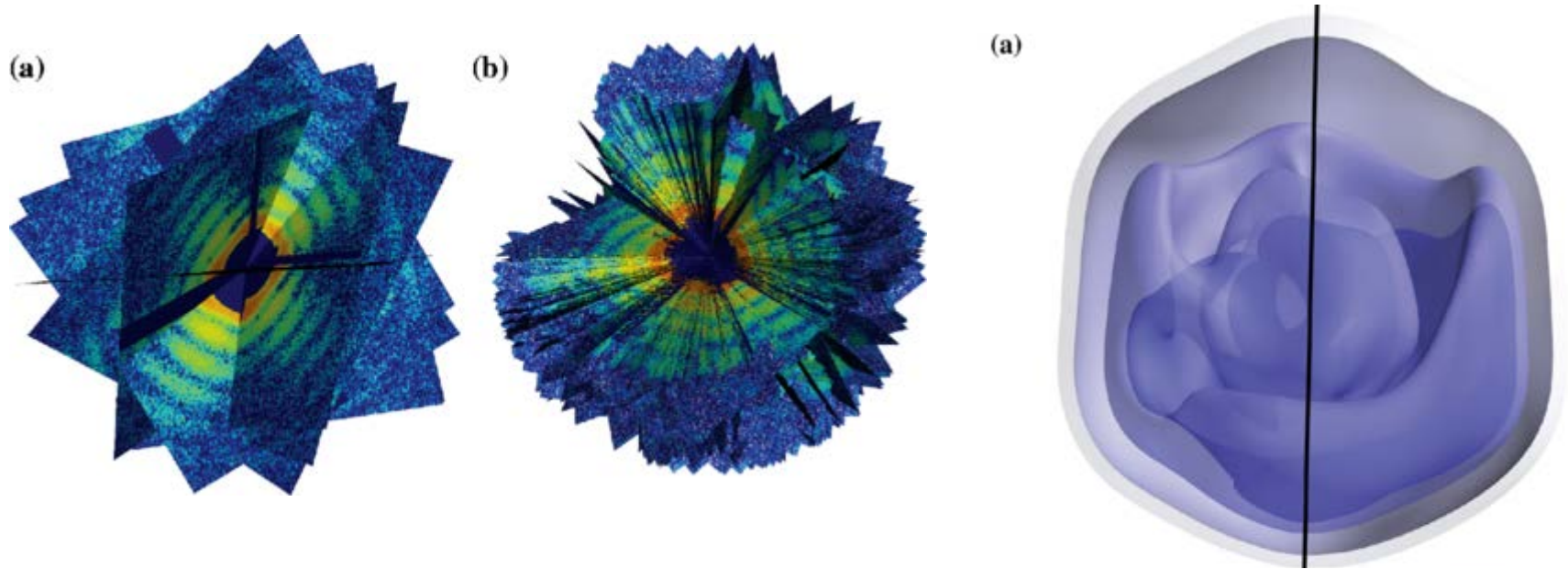


Neutze, Wouts, van der Spoel, Weckert, Hajdu Nature 406, 752 (2000)

Single Mimivirus Particles Intercepted and Imaged with an X-ray laser

Seibert et al. Nature 470, 78–81 (2011) only 2D.....

Coherent X-ray imaging of non-crystallized samples



Three-Dimensional Reconstruction of the Giant Mimivirus Particle
with an X-Ray Free-Electron Laser

Ekeberg *et al.* Phys. Rev. Lett. **114**, 098102 (2015)

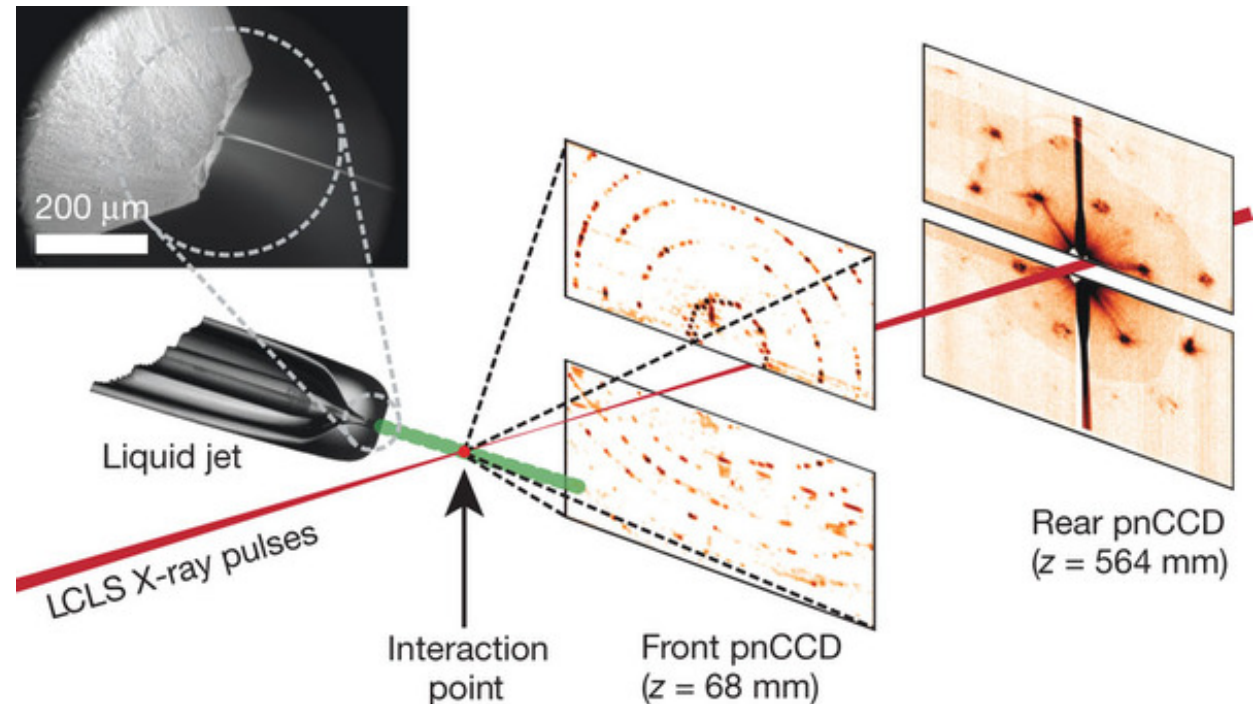
Can we get 3D image from a single shot data?

Characteristic properties of FEL pulses

Intense and ultra-short pulses at X-rays

Why X-rays? structure determination at atomic resolution

Femtosecond X-ray Protein Nano-crystallography, Chapman et al., *Nature* **470**, 73–77 (2011).



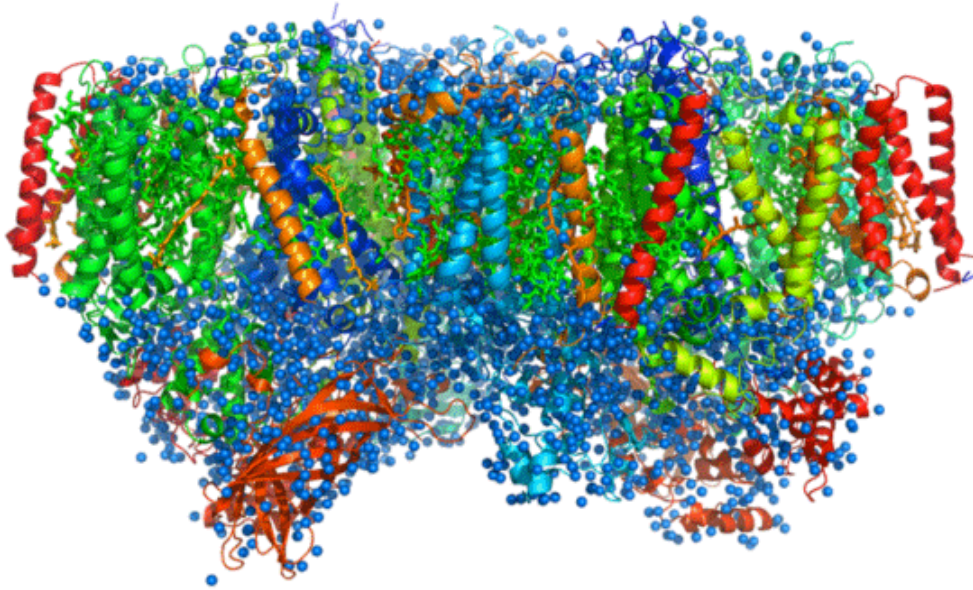
High-Resolution Protein Structure Determination by Serial Femtosecond Crystallography, Boutet et al. *Science* **337**, 362 (2012).

Natively Inhibited Trypanosoma brucei Cathepsin B Structure Determined by Using an X-ray Laser, Redecke et al. *Science* **339**, 227 (2013).

Photo-system II

Native structure of photosystem II at 1.95 Å resolution revealed by a femtosecond X-ray laser (SACLA)

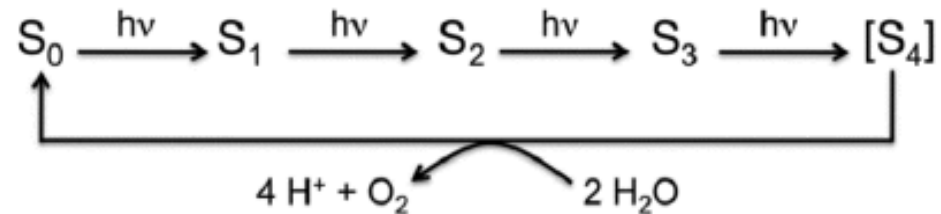
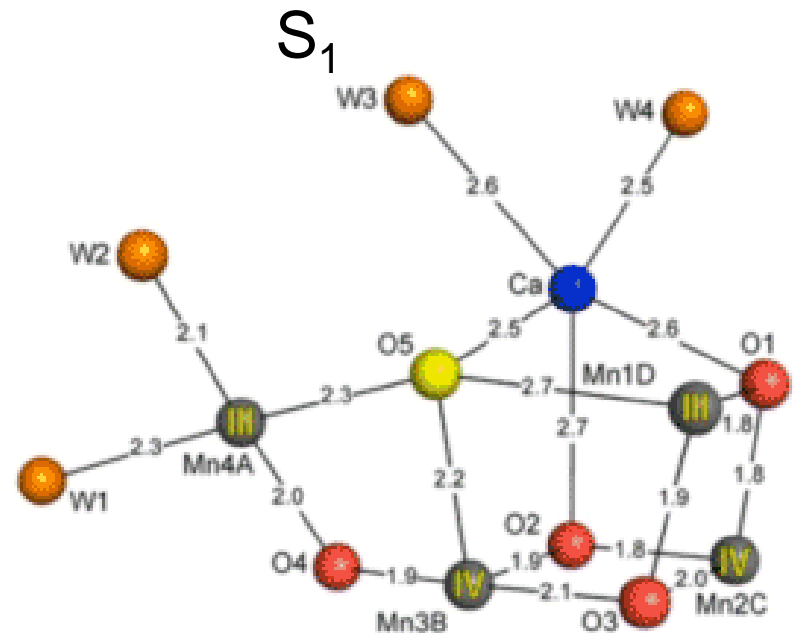
Suga et al, Nature (2014)



"Determination of damage-free crystal structure of an X-ray sensitive protein using an XFEL"

Nature Methods, 2014,
doi:10.1038/NMETH.2962

SR results (Nature 2011) had radiation damage...



Dynamic behavior of photo-system II

"The Mn_4Ca photosynthetic water-oxidation catalyst studied by simultaneous X-ray spectroscopy and crystallography using an X-ray free-electron laser"

Rosalie Tran et al

Phil. Trans. R. Soc. B 369 20130324
(2014) doi: 10.1098/rstb.2013.0324

"Taking snapshots of photosynthetic water oxidation using femtosecond X-ray diffraction and spectroscopy"

Jan Kern et al

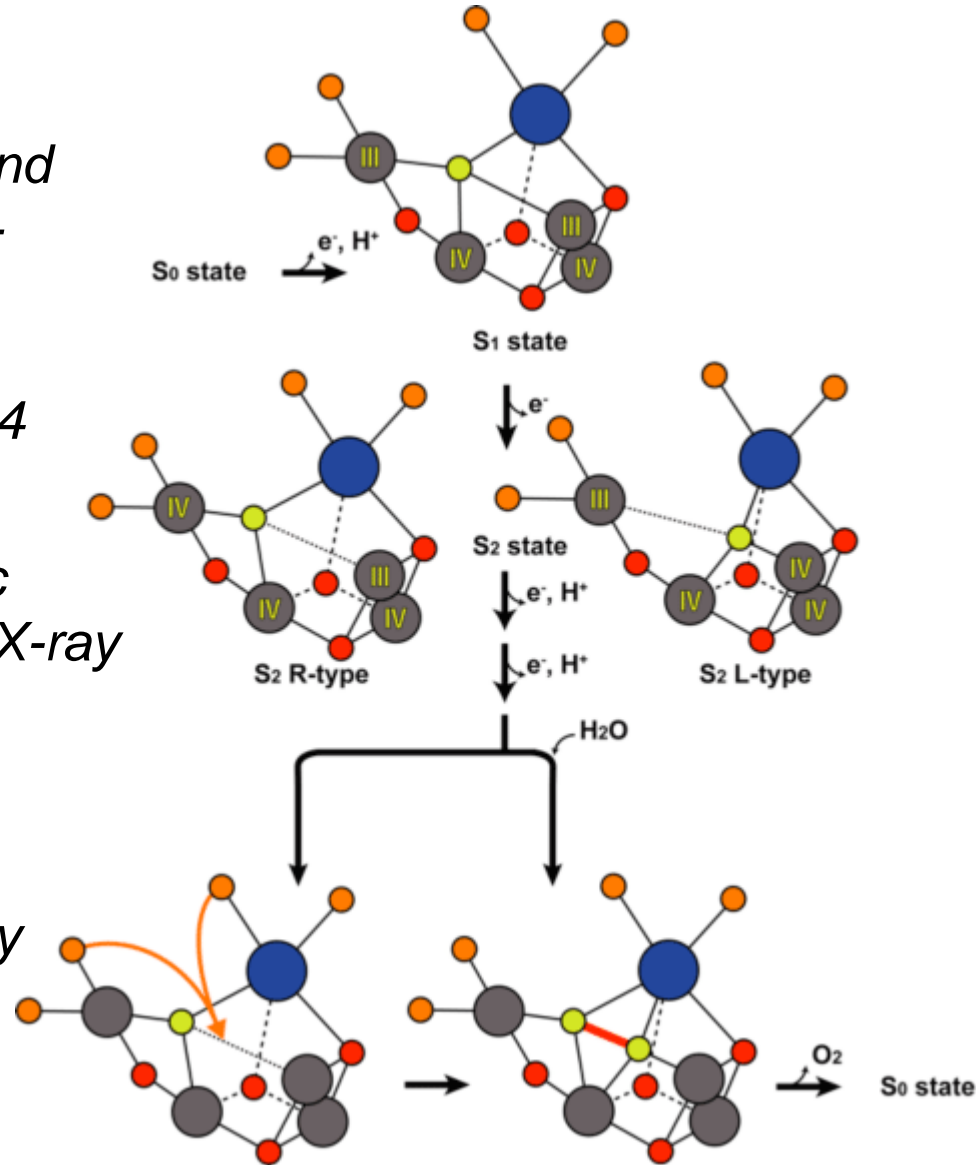
Nature Comm. 5. 4371 (2014)

"Serial time-resolved crystallography of photosystem II using a femtosecond X-ray laser"

C. Kupitz et al

Nature (2014)

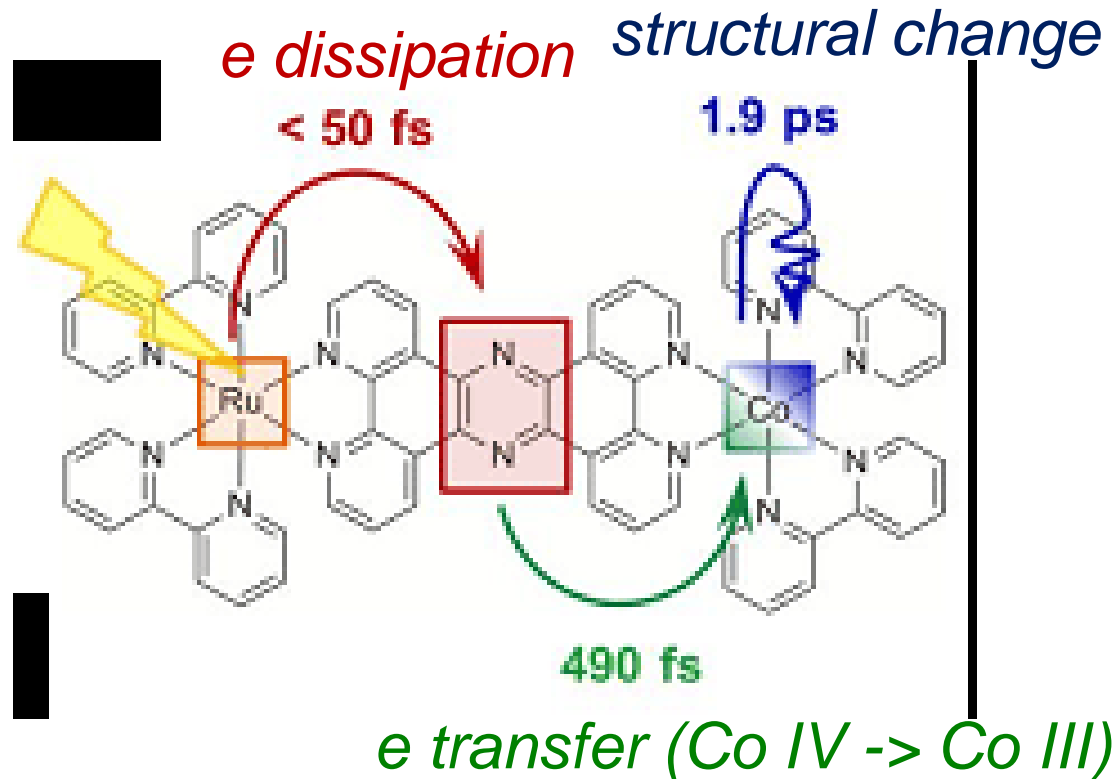
doi:10.1038/nature13453



Suga et al, Nature

Towards artificial photosynthesis

Visualizing the non-equilibrium dynamics of photo-induced intramolecular electron transfer with femtosecond X-ray Pulses (SACLA) Canton et al. *Nature Comm.* 6, 6359 (2015)



Serial femtosecond x-ray crystallography: phasing

Phasing of the serial femtosecond x-ray crystallography had been relying on molecular replacements...

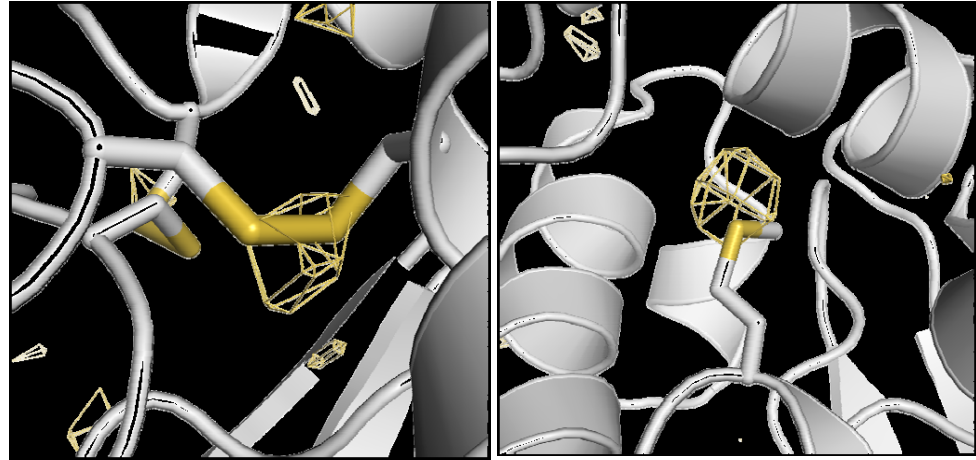
If the structure is completely unknown, phasing approaches make use of anomalous dispersion in the scattering signals from specific atoms.

Anomalous signal from S atoms in protein crystallographic data from an X-ray free-electron laser

T.R.M. Barends, L. Foucar, R.L. Shoeman, K. Ueda and I. Schlichting

*Acta Cryst. D **69**, 838-842 (2013).*

One of the first experiments at SACLA



*For phasing with a heavy atom, one has to take account of high x-ray intensity
Multi-wavelength anomalous diffraction at high X-ray intensity*

*S.-K. Son, H. N. Chapman, and R. Santra, Phys. Rev. Lett. **107**, 218102 (2011).*

However.....

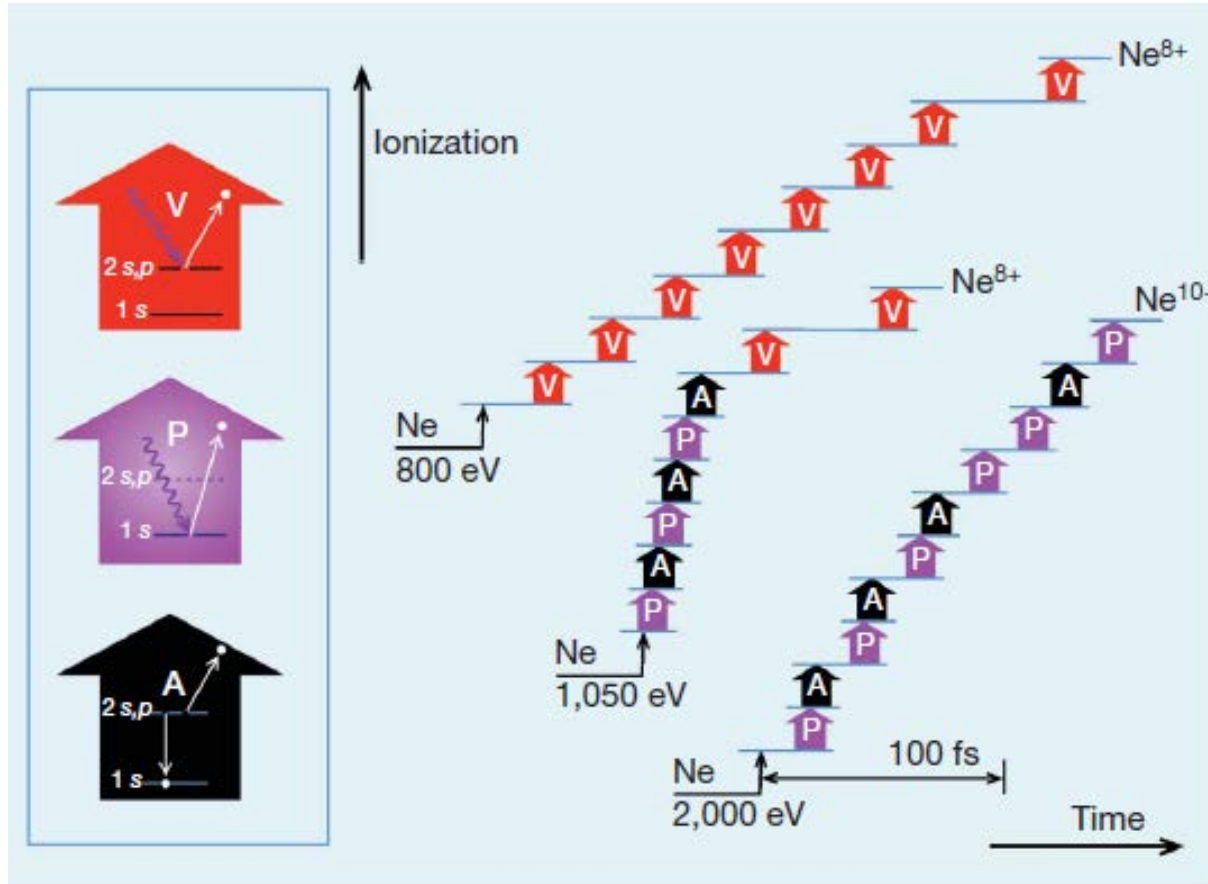
De novo protein crystal structure determination from X-ray free-electron laser data

*T.R.M. Barends et al., Nature **505**, 244 (2014).*

Conventional phasing method based on anomalous dispersion worked....

Characteristic properties of FEL pulses

Intense 10^{14} W/cm^2 (EUV) - 10^{20} W/cm^2 (X)



One LCLS pulse at 2 keV can remove all ten electrons from the neon atom.

The pulse is so intense that it causes electronic damage to the sample.

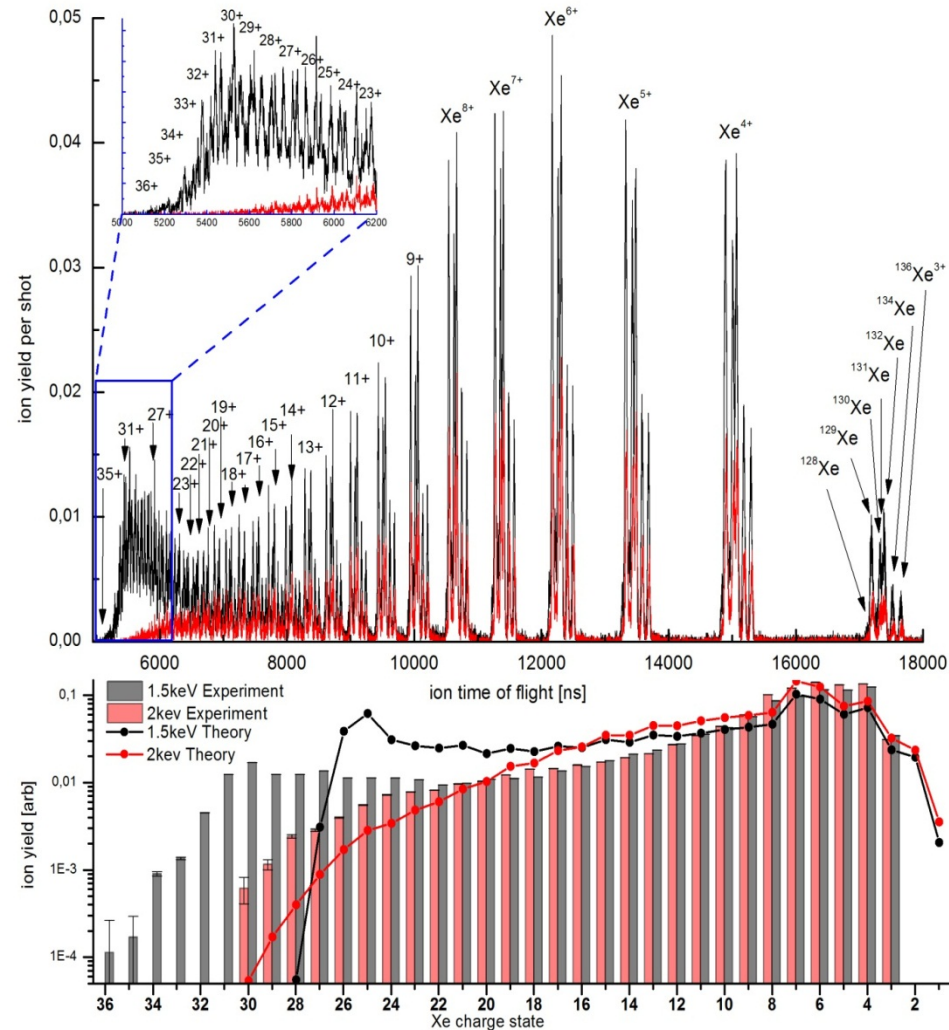
Femtosecond electronic response of atoms to ultra-intense x-rays

*L. Young et al., Nature **466**, 56 (2010).*

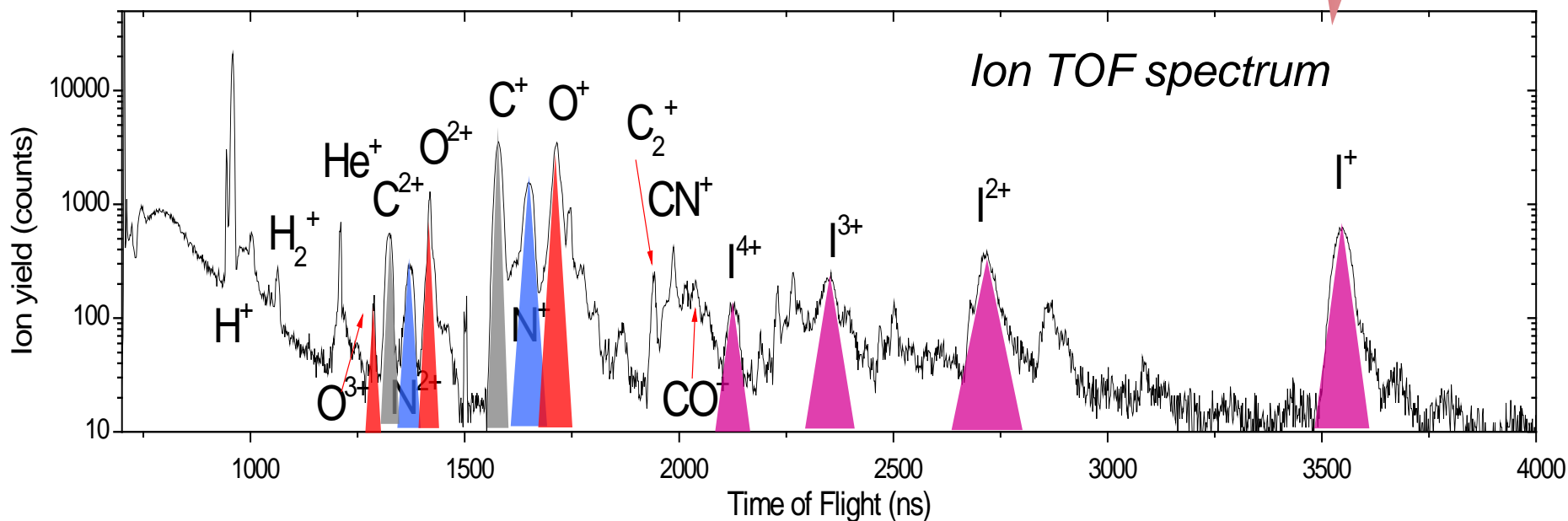
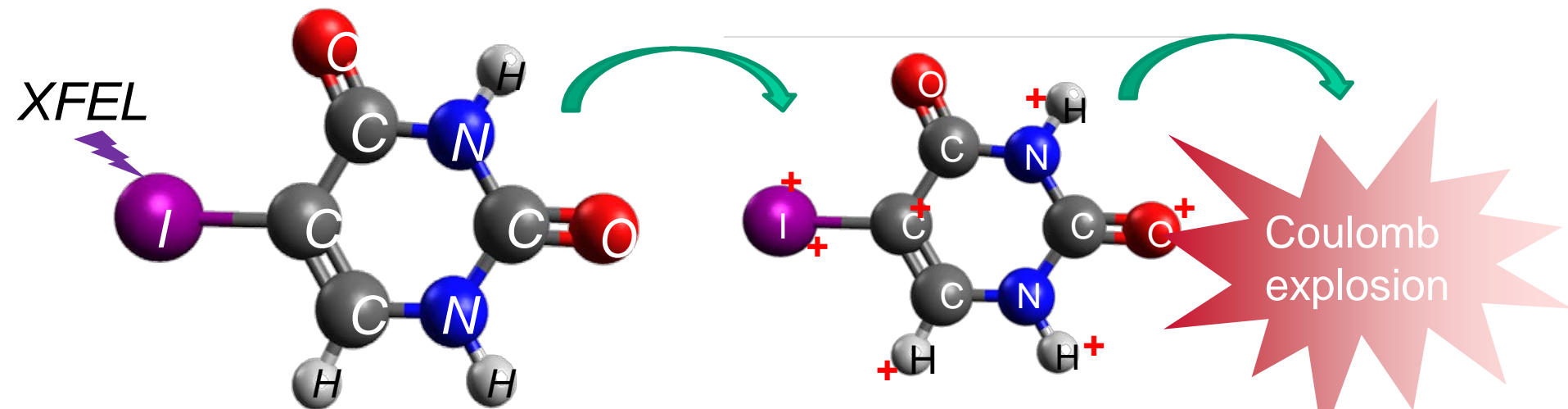
Non-linear X-ray atomic Physics

Ultra-Efficient Ionization of Heavy Atoms by Intense X-Rays

B. Rudek, S-K. Son, L. M. Foucar, S. W. Epp, B. Erk, R. Hartmann, M. Adolph, R. Andritschke, A. Aquila, N. Berrah, C. Bostedt, J. Bozek, N. Coppola, F. Filsinger, H. Gorke, T. Gorkhover, H. Graafsma, L. Gumprecht, A. Hartmann, G. Hauser, S. Herrmann, H. Hirsemann, P. Holl, A. Hömke, L. Journel, C. Kaiser, N. Kimmel, F. Krasniqi, K-U. Kühnel, M. Matysek, M. Messerschmidt, D. Miesner, T. Möller, R. Moshhammer, K. Nagaya, B. Nielsson, G. Potdevin, D. Pietschner, C. Reich, D. Rupp, R. Santra, G. Schaller, I. Schlichting, C. Schmidt, F. Schopper, S. Schorb, C-D. Schröter, J. Schulz, M. Simon, H. Soltau, L. Strüder, **K. Ueda**, G. Weidenspointner, J. Ullrich, A. Rudenko, and D. Rolles
Nature Photonics 6, 858 (2012).



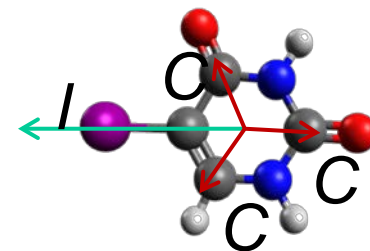
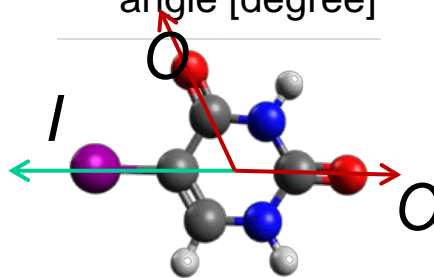
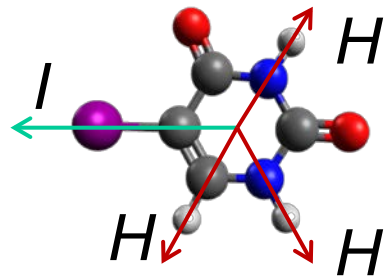
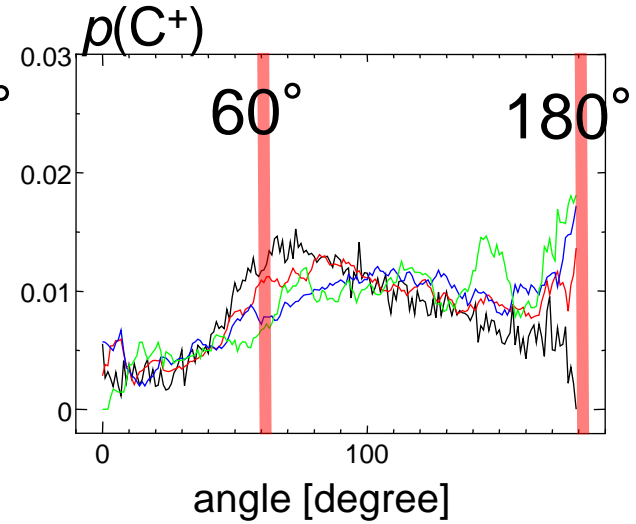
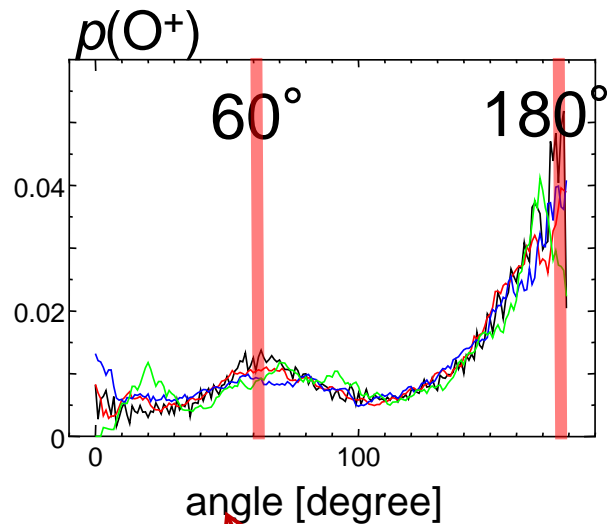
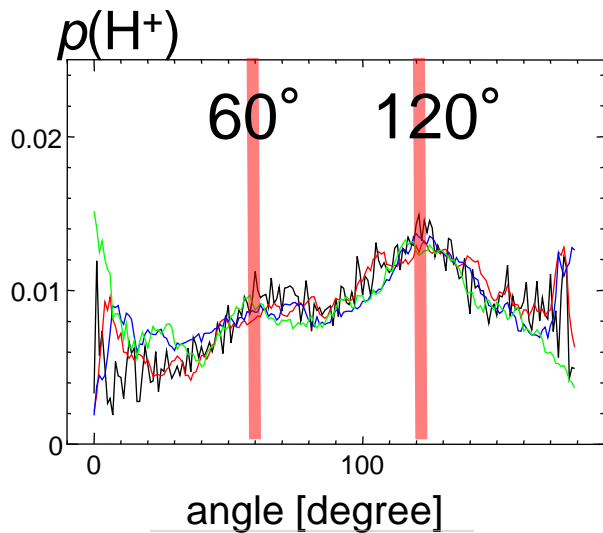
5-iodouracil ($C_4H_3IN_2O_2$)



We observed iodine ions with the charge up to +4 and many more ions

2-body angular correlation

Black: I^+ , Red: I^2+ , Blue: I^3+ , Green: I^4+

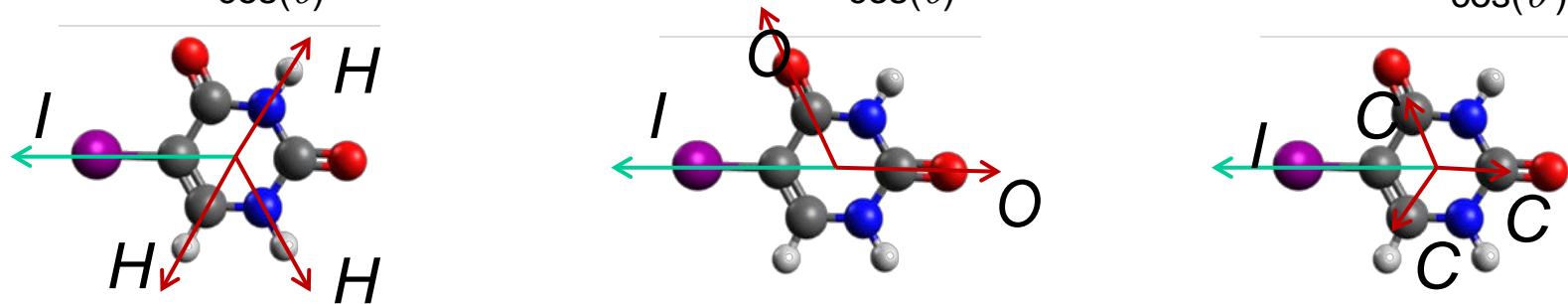
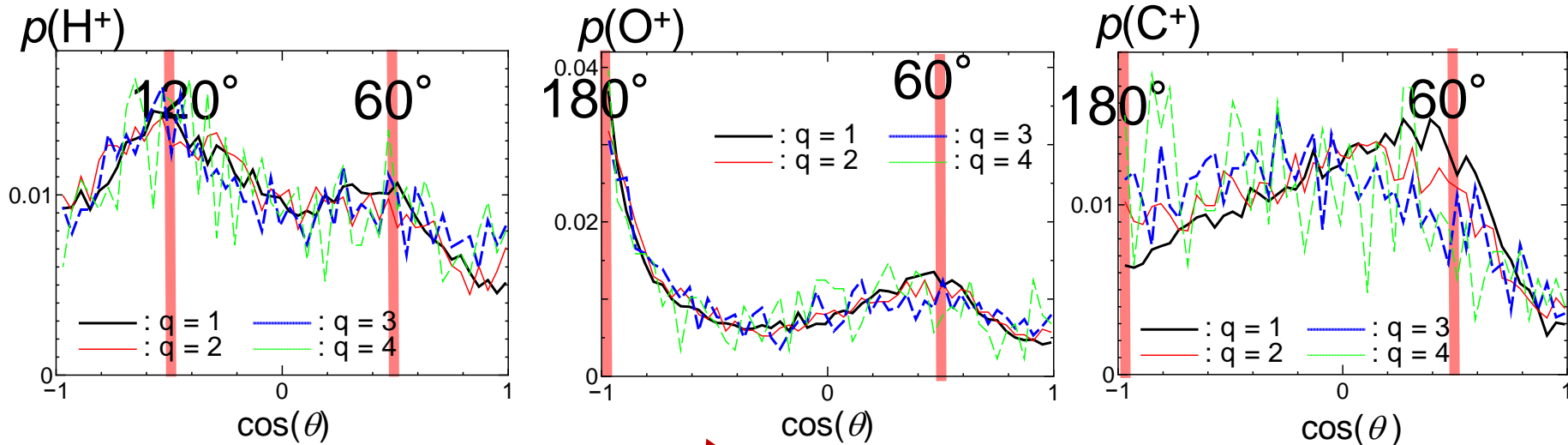


Red lines indicate angle between I direction and H , O , or C directions from the center of the aromatic ring in the neutral molecule.

Angular correlations of fragment ions (except carbon ions) reflect the shape of parent molecule.

2-body angular correlation

Black: I^+ , Red: I^2+ , Blue: I^3+ , Green: I^4+

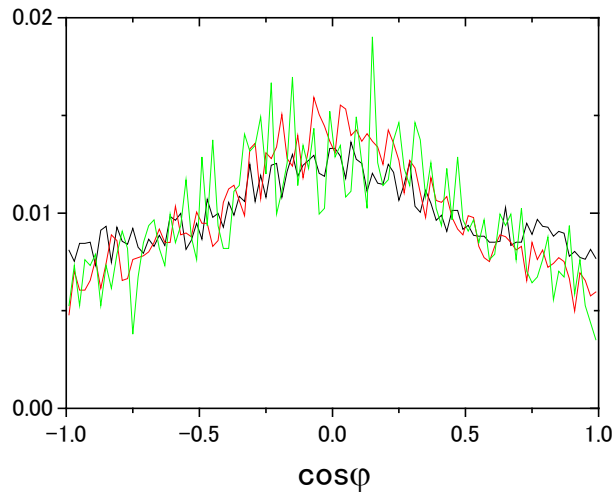


Red lines indicate angle between I direction and H , O , or C directions from the center of the aromatic ring in the neutral molecule.

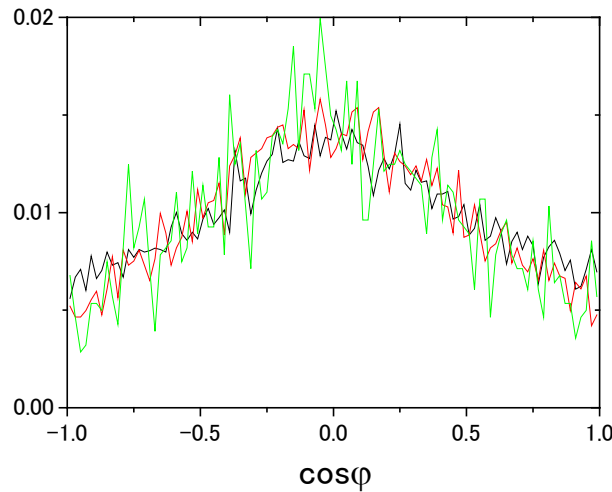
Angular correlations of fragment ions (except carbon ions) reflect the shape of parent molecule.

3-body angular correlation

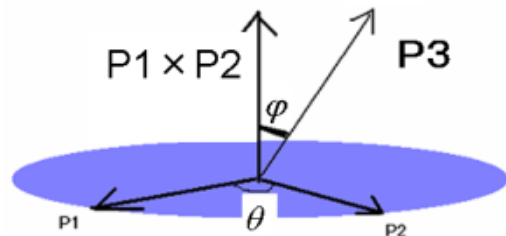
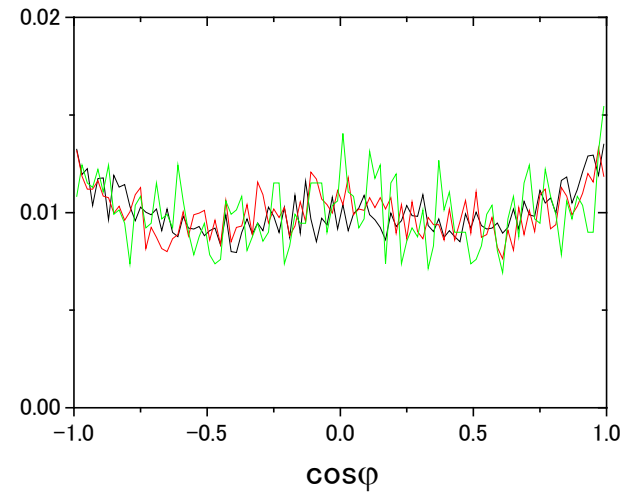
$I^+-H^+-H^+$



$I^+-H^+-O^+$



$I^+-H^+-C^+$

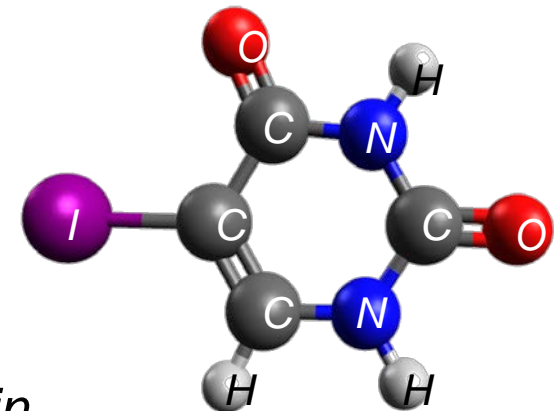


Black: I^+ , Red: I^{2+} , Green: I^{3+}

$$\cos(\varphi) = \frac{(\mathbf{P}_1 \times \mathbf{P}_2) \cdot \mathbf{P}_3}{(|\mathbf{P}_1||\mathbf{P}_2||\mathbf{P}_3|\sin(\theta))}.$$

\mathbf{P}_1 , \mathbf{P}_2 and \mathbf{P}_3 are momentum of ions.

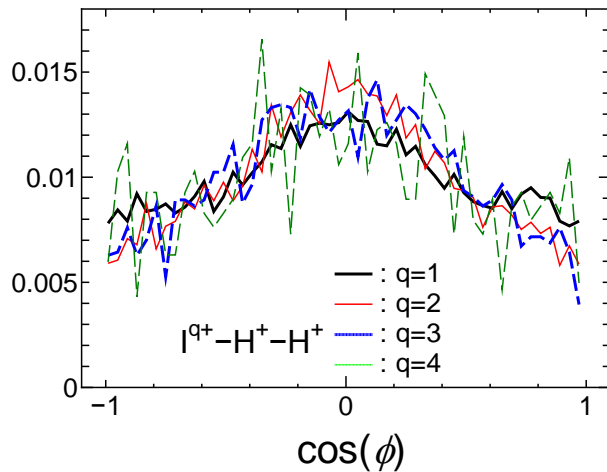
$\cos(\varphi) = 0$ means that three momentum vectors are in the same plane.



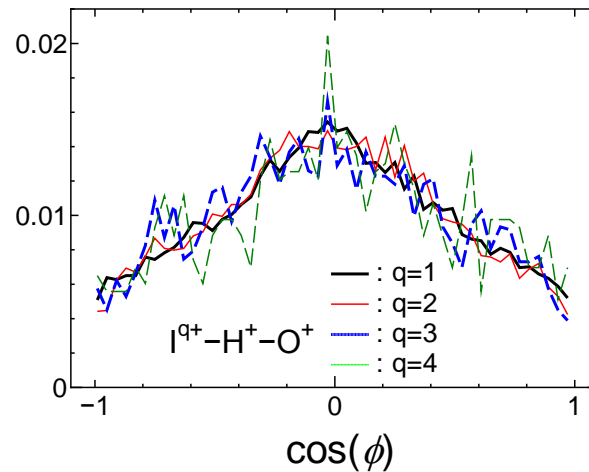
These results suggest that carbon ions are released to off-planar direction.

3-body angular correlation

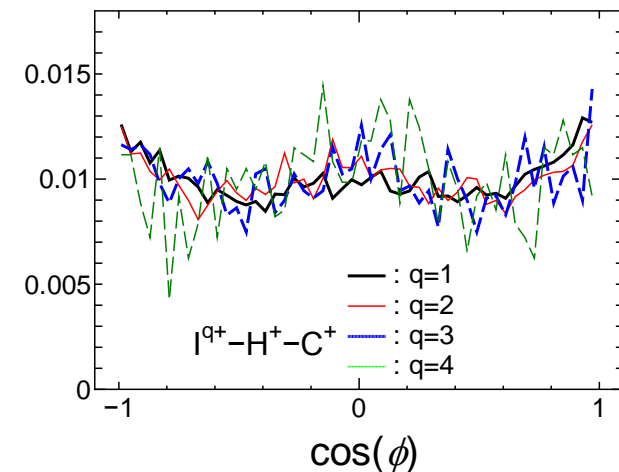
$I^+-H^+-H^+$



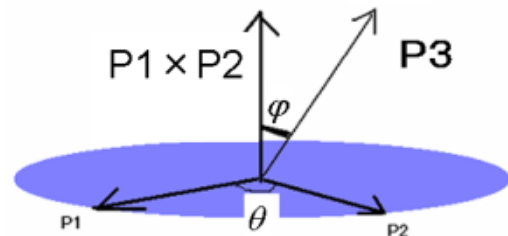
$I^+-H^+-O^+$



$I^+-H^+-C^+$



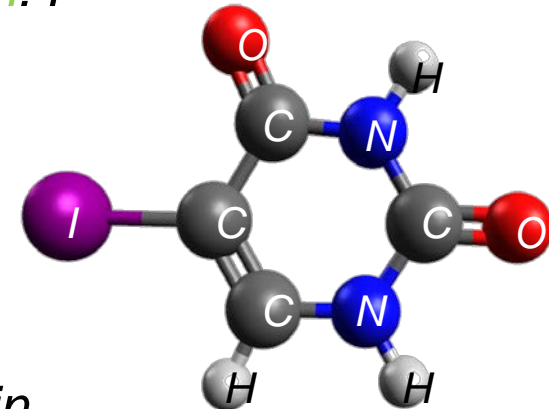
Black: I^+ , Red: I^{2+} , Blue: I^{3+} , Green: I^{4+}



$$\cos(\phi) = (\mathbf{P}_1 \times \mathbf{P}_2) \cdot \mathbf{P}_3 / (|\mathbf{P}_1| |\mathbf{P}_2| |\mathbf{P}_3| \sin(\theta)).$$

\mathbf{P}_1 , \mathbf{P}_2 and \mathbf{P}_3 are momentum of ions.

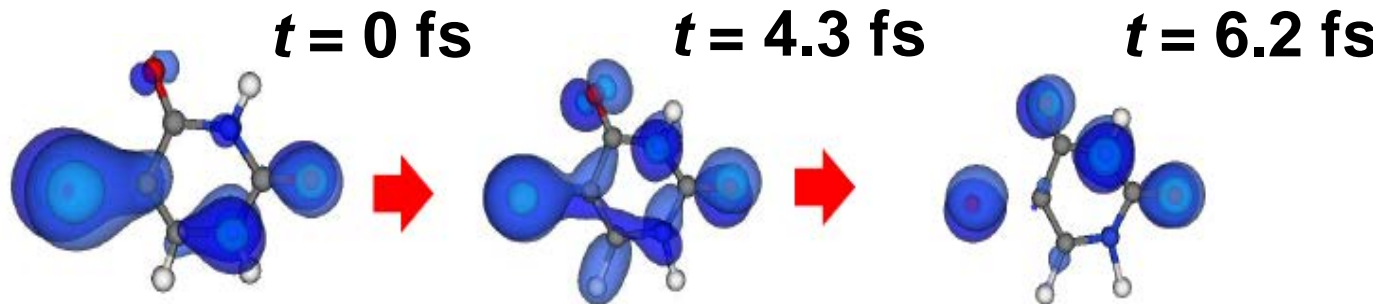
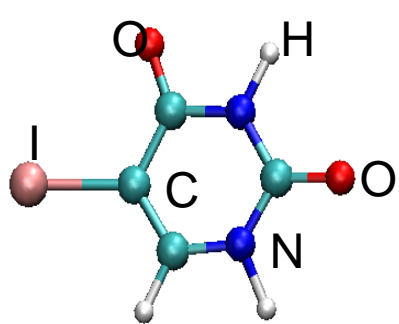
$\cos(\phi) = 0$ means that three momentum vectors are in the same plane.



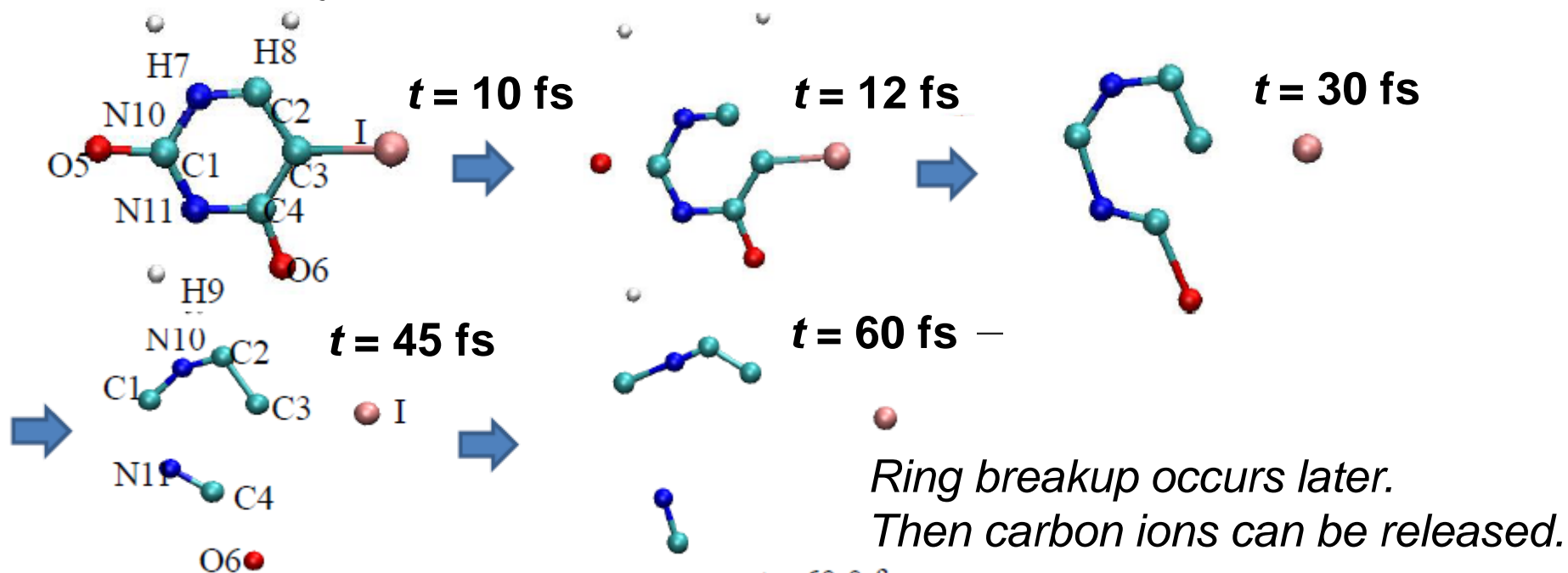
These results suggest that carbon ions are released to off-planar direction.

DFT simulations for 5-iodouracil

We found charge transfer (localization) is completed within 6 fs.



Protons are ejected within 10 fs.



These calculation results support well the experimental findings.

IR energy absorption of nanoplasma

The *energy absorbing rate per electron* via surface plasmon resonance can be described by the *Drude model*,

$$\frac{\partial E}{\partial t} = \frac{1}{n_e} E \cdot \frac{\partial \mathbf{D}}{\partial t} = 2U_p \frac{\Gamma}{\left(\frac{\omega_p^2}{3\omega^2} - 1\right)^2 + \Gamma^2/\omega^2}$$

Ponderomotive energy Damping rate
 Γ

ω_p : plasma frequency ω : NIR Laser frequency

$$\omega_p = \sqrt{\frac{e^2 n_e(t)}{m_e \epsilon_0}}$$

n_e : electron density $n_e(t) = 3ZN/4R(t)^3$

R : radius of nanoplasma $R(t) = vt + R_0$

R_0 : cluster size $R_0 = 4.4\text{nm}$

$$\Gamma(t) = \Gamma_0 \times [R_0/R(t)] \quad R(t) = R_{\text{res}} \quad @ \quad \omega_p = 3^{1/2}\omega$$

Fitting to the experimental data

$$\frac{\partial E^*}{\partial t} = \frac{4\pi I_{\text{NIR}}}{c} (R_0 + v_0 t_{\text{Res}})^3 \Gamma_0 \frac{\left[\frac{(R_0 + v_0 t)}{R_0} \right]^{-1}}{\left[\exp\left(-\frac{t-t_{\text{Res}}}{\tau_e}\right) \left(\frac{R_0 + v_0 t}{R_0 + v_0 t_{\text{Res}}}\right)^{-3} - 1 \right]^2 + \left[\Gamma_0 \left(\frac{R_0 + v_0 t}{R_0}\right)^{-1} / \omega \right]^2} \exp\left(-\frac{t-t_{\text{Res}}}{\tau_e}\right)$$

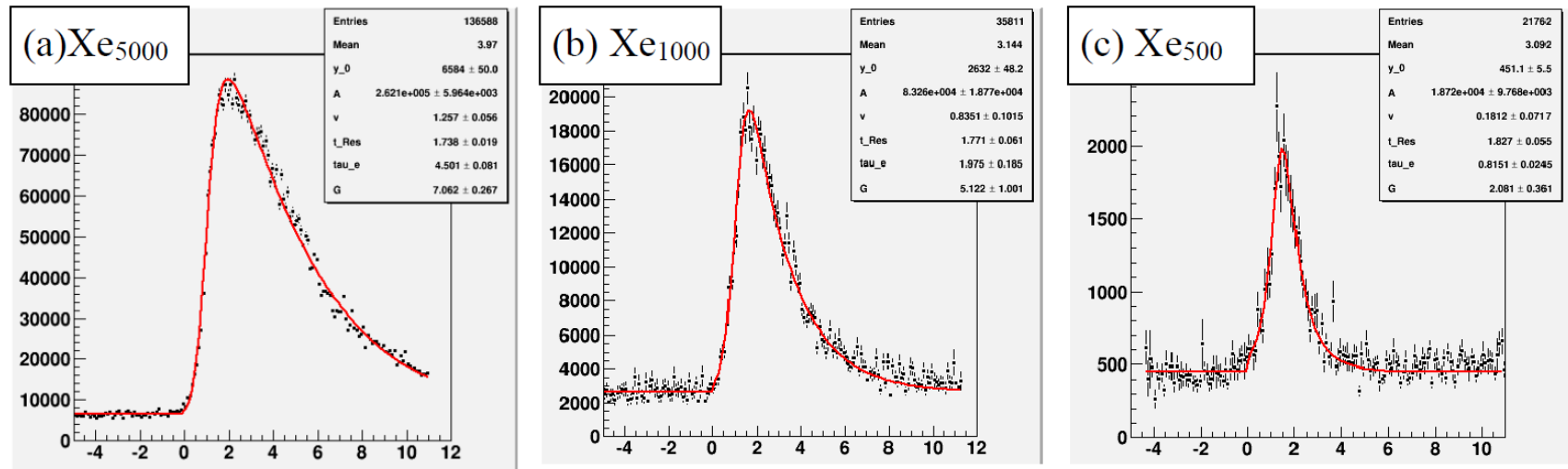
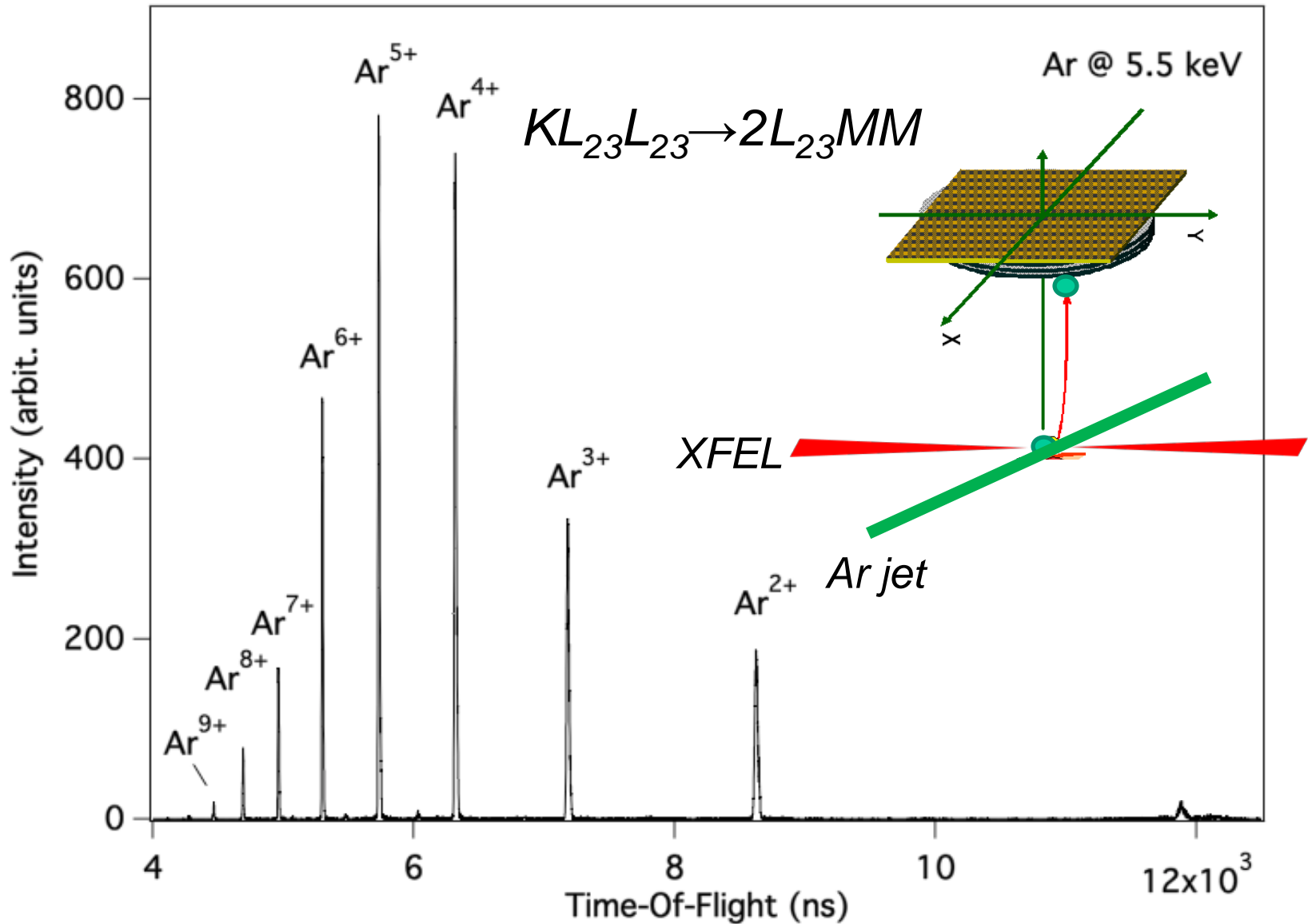


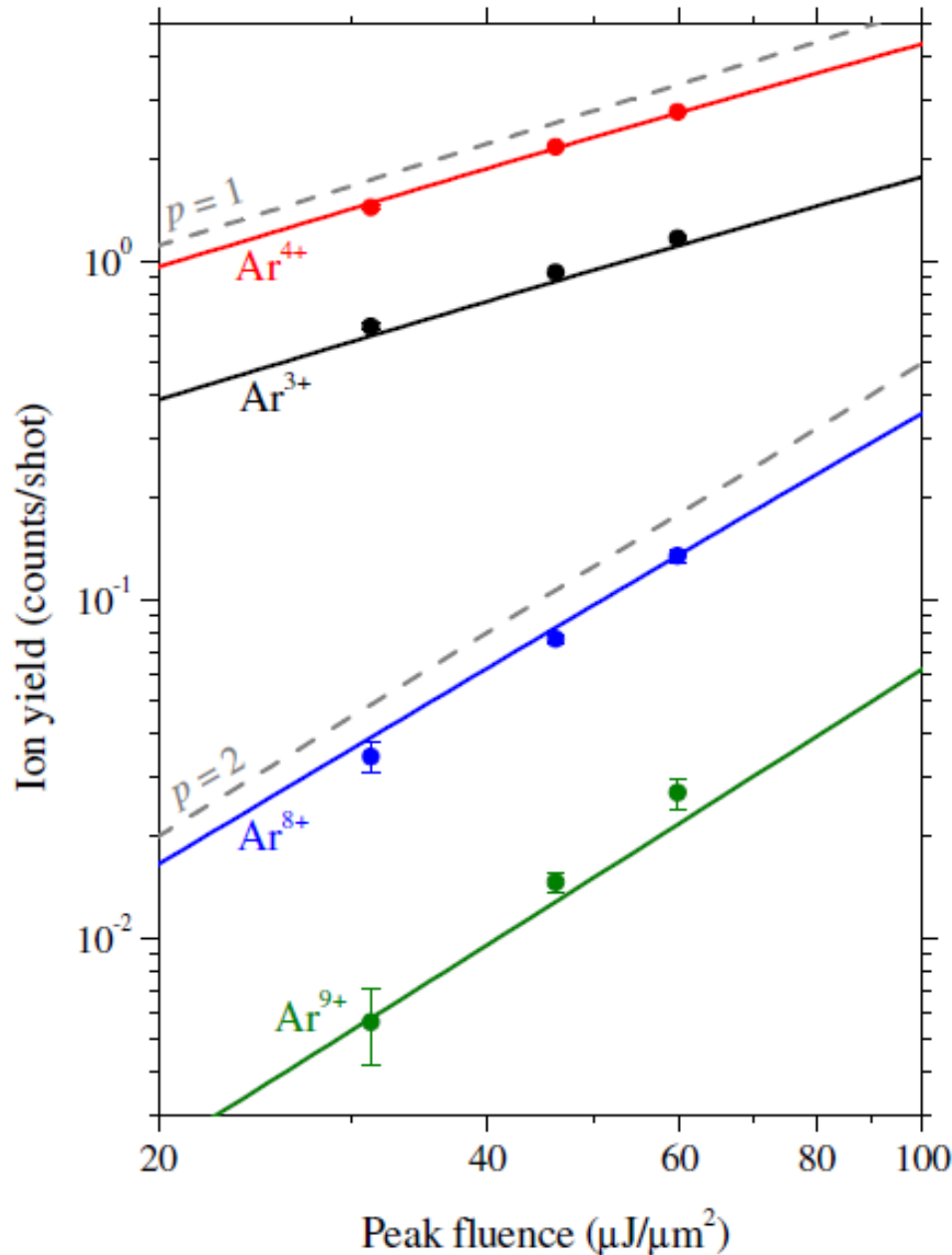
Fig3-1. Total energy absorption as a function of delay time. (a)Xe₅₀₀₀, (b) Xe₁₀₀₀, (c) Xe₅₀₀. Red curves show the fitting results by equation (4).

We can characterize XFEL-induced nanoplasma by investigating its expansion and recombination by IR probe !

Time of Flight spectrum of argon ions



XFEL fluence dependence for Ar^{n+} yields



*Ar^{4+} : single photon K-shell ionization
 $\rightarrow KL_{23}L_{23}$ Auger $\rightarrow 2L_{23}MM$ Auger*

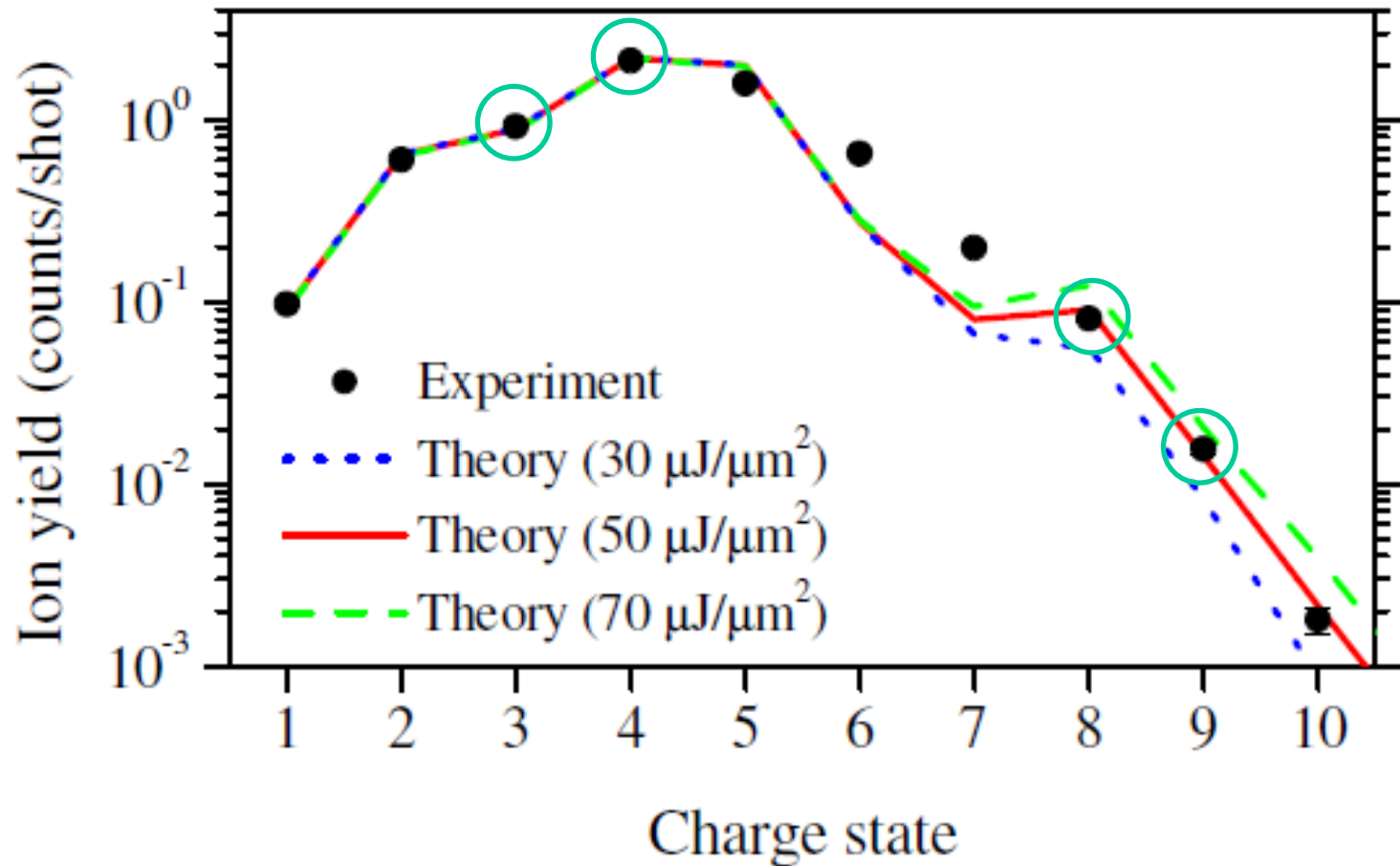
*Ar^{3+} : single photon K shell absorption
 $\rightarrow KL_{23}M$ Auger $\rightarrow L_{23}MM$ Auger*

Ar^{8+} , Ar^{9+} : sequential two photon K shell ionization

Bench mark ab initio calculation reproduces fluence dependence and relative ratios.

In the theory, the pulse shape of Gaussian of 30 fs (FWHM), and Gaussian focal shape of $1 \mu\text{m}$ (FWHM) $\times 1 \mu\text{m}$ (FWHM) are assumed.

Charge state distribution of Ar: experiment and theory



By comparison with theory, we obtained peak fluence of $50 \mu\text{J}/\mu\text{m}^2$ in the experiment!