

Harnessing ultraintense x-rays: from atomic response to applications



Linda Young Nobel Symposium 158 Sigtunahöjden, Sigtuna, Sweden 14 Jun 2015



Compare ultraintense optical and x-ray sources



X-ray sources: accelerator-based vs laser-based HHG



XFELs 10⁸ "brighter" than HHG sources

From: Miao, Ishikawa, Robinson, Murnane, Science 348, 530 (2015)

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Outline

- Birth of world's first hard x-ray FEL LCLS
- Non-resonant high intensity x-ray phenomena LCLS Experiment 1: Oct 1 - 6, 2009
 L. Young *et al.*, Nature 466, 56 (2010)
- Resonant high intensity x-ray processes
 LCLS Experiment 5: Oct 1 6, 2009
 E.P. Kanter *et al.*, PRL 107, 233001 (2011)
- Towards single particle imaging



Linac Coherent Light Source at SLAC

X-FEL based on last 1-km of existing 3-km linac

Proposed by C. Pellegrini in 1992

Injector (35^o) at 2-km point

Existing 1/3 Linac (1 km)

New e⁻ Transfer Line (340 m)

Far Experiment

Hall

X-ray Transport Line (200 m)

1.5-15 Å

(14-4.3 GeV)

Undulator (130 m)

Near Experiment Hall

-



NATIONAL ACCELERATOR LABORATOR

April 10, 2009: LCLS lases at 1.5 Angstroms



In my life before SLAC, I had the privilege to participate (in various capacities) in the design, construction and commissioning of two linacs, two synchrotrons, four storage rings and three FELs (free-electron lasers). Now I have had the privilege to be in SLAC's Main Control Center on April 10, when the Linac Coherent Light Source became a 1.5 Ångstrom laser. I don't expect I will ever, as long as I live, see such a beautiful, smooth turn-on of any light source. With each undulator placed on the beam path, the FEL power increased by a factor of about 2.3; two hours into the first attempt at lasing, the pinpoint of FEL light from twelve undulators was nearly 2,000-fold more intense than plain old undulator radiation. The team called it quits at 11:30 p.m. that night. When they returned at 8:00 a.m. the next morning, the FEL light came back as soon as the shutter was opened. – John Galayda

21:33:52 • 10 mm

21:03:03





Yes I Do Smile on Occasion





LCLS saturation at 1.5 Å



• Saturation after ~65 meters of undulator!

Paul Emma PAC 2009 proceedings



Science Drivers for LCLS



AMO: Atomic Molecular and Optical
SXR: Soft X-ray Materials Science
XPP: X-ray Pump-Probe
XCS: X-ray Correlation Spectroscopy
CXI: Coherent X-ray Imaging
MEC: Materials in Extreme Conditions

AMO

- Understand and control x-ray atom/molecule interactions at ultrahigh x-ray intensity as a foundation for other applications.
- Provide diagnostics of the LCLS radiation



X-ray Diffraction Pattern

AMO questions at the ultraintense x-ray frontier

- fundamental nature of x-ray damage at high intensity

 Coulomb explosion
 electronic damage
 behavior at 10²² W/cm² - 1Å
- nonlinear x-ray processes
 role of coherence
- quantum control of inner-shell processes



Neutze, Wouts, van der Spoel, Weckert, Hajdu Nature 406, 752 (2000)

LCLS Experiment 1 - Oct 1, 2009

Nature of the electronic response to

10⁵ x-rays/Å² 80 - 340 fs 800 - 2000 eV

 $\sim 10^{18} \, W/cm^2$

Original single molecule imaging parameters, Neutze et al. Nature (2000) $3 \times 10^{12} x$ -rays/(100 nm)² = $3 \times 10^{6} x$ -rays/Å² 10 fs ~ $10^{22} W/cm^{2}$

Our approach to understanding ultraintense x-ray interactions

Start with a well-characterized target





 Probe changes in interaction from outer- to inner-shell between 800-2000 eV **Guided by theory**

Theory: Rohringer & Santra, PRA 76, 033416 (2007)



Three target energies: 800 eV, 1050 eV, 2000 eV

Valence ionization, core ionization and Auger decay



Sequential single photon processes dominate the interaction

How does one arrive at a particular charge state?



- Hollow atoms produced at high x-ray intensity
- Electron spectroscopy can define the mechanism



High field physics chamber



Day 1 - two interesting observations

Single ~100 fs pulse at 2000 eV fully strips neon

6-photon, 10-electron process



 Shorter pulses with equal pulse energy & fluence suppress absorption & damage.

Theory can model ultraintense x-ray-induced electronic damage in neon



Theory

- Intensity averaged
- Fluence determined by experiment

Consistent with "measured" pulse energy and focus.

Sang-Kil Son, Robin Santra – refined calcs include shakeoff – G. Doumy et al, PRL 2011)

Atoms become transparent at high x-ray intensity !



- x-ray absorption is due to the presence of 1s electrons
- high x-ray intensities eject both 1s electrons rendering the atom transiently transparent
- slowing atomic clocks create transparency at surprisingly long timescales

Electron spectrometers track ionization mechanism



"Slow" 1s photoelectrons along x-ray polarization axis

"Fast" valence photoelectrons and Augers along polarization axis

Clean hollow atom signature double-core-hole Auger $\theta = 90^{\circ}$

Hollow atom production: deliberate, huge and an a an indicator of x-ray pulse duration



Hollow atom yield
@ LCLS ~10%
@ synchrotron ~0.3%
due to electron correlation

1050 eV, nominal electron bunch duration ~80 fs



Absorption vs scattering: normal and hollow atoms



	photo elastic	$rac{\sigma_{Compton}}{\sigma_{elastic}}$
2 keV	360	0.05
8 keV 8 keV hollow	20	0.60

Impact of hollow atom formation on coherent x-ray scattering Sang-Kil Son, L Young, R Santra Phys. Rev A. **83**, 033402 (2011)

Prescient prediction

"Furthermore, during short intense pulses numerous Kholes may be present at any one time, reducing the photoelectric cross-sections of atoms in which they were produced and thus lowering the total number of primary ionization events in the sample (see trend in Fig. 1). This effect makes the system radiation hardened to photo-ionization during very short exposures, and is more pronounced at higher radiation intensities."

Neutze, Wouts, van der Spoel, Weckert, Hajdu Nature (2000)



NEWS & VIEWS

ATOMIC PHYSICS

X-ray laser peels and cores atoms

Justin Wark

The world's first kiloelectronvolt X-ray laser produces such a high flux of photons that atoms can be 'cored'. In other words, the light source can knock out both the electrons of an atom's innermost shell.



Femtosecond electronic response of atoms to ultra-intense X-rays Nature 466, 56 (2010).

L. Young¹, E. P. Kanter¹, B. Krässig¹, Y. Li¹, A. M. March¹, S. T. Pratt¹, R. Santra^{1,2}, S. H. Southworth¹, N. Rohringer³, L. F. DiMauro⁴, G. Doumy⁴, C. A. Roedig⁴, N. Berrah⁵, L. Fang⁵, M. Hoener^{5,6}, P. H. Bucksbaum⁷, J. P. Cryan⁷, S. Ghimire⁷, J. M. Glownia⁷, D. A. Reis⁷, J. D. Bozek⁸, C. Bostedt⁸ & M. Messerschmidt⁸



Summary: non-resonant ultraintense x-ray interactions

- Ultraintense x-ray interactions multiphoton processes rule!
 - establish sequential single photon absorption as dominant ionization mechanism fully stripped neon: six-photon, ten-electron ($\sim 10^{12}/\mu m^2$)
 - multiple photon absorption probability high when fluence > 1/ σ
 - controlled electron stripping (outer v inner shells)
- X-ray induced transparency a general phenomena
 - transient x-ray transparency caused by ejection of inner-shell electrons
 - induced transparency = frustrated absorption = core-level bleaching molecules: Hoener *et al.*, PRL **104**, 253002 (2010) clusters: Schorb *et al.*, PRL **108**, 233401 (2012) solids: Yoneda *et al.*, Nat. Comm. (2014), Rackstraw *et al.*, PRL (2015)
 - implications for imaging: σ_{scatt}/σ_{abs} is increased
- Femtosecond time-scale atomic processes provide FEL diagnostics

LCLS Experiment 5

Resonant x-ray processes at high intensity



Can we control inner-shell electron dynamics?

"Rabi flopping" may inhibit Auger decay & x-ray damage.





But LCLS linewidth ~ 8 eV!

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Rabi-flopping on 1s - 2p resonance more feasible



Observe Auger yield when x-rays scanned over 1s - 2p resonance. Observe broadening at resonance to indicate Rabi flopping Theory: Rohringer & Santra PRA (2008).

Calculated "Resonant Auger effect at high x-ray intensity"



-> Look for Auger line broadening on resonance

N. Rohringer & R. Santra, PRA 77, 053404 (2008)

Looking for Rabi flopping: unveiling and driving hidden resonances with LCLS pulses



 High fluence pulse alters target to reveal enormous "hidden" resonances ~1000x larger than background

 X-ray absorption spectrum changes rduring the fs duration pulse



Characterization of SASE x-ray pulse required

Shot-to-shot photon energy jitter

-Photoelectron energy spectra give x-ray energy centroid

-Jitter derived from shot-to-shot GeV electron beam energy measurements

 $E_x (eV) = 44.25 E_{e(}(GeV)^2$

Conditions	FWHM photon energy jitter (eV)	
40 pC (<10 fs) 850 eV 0.3 mJ 4500 A	4.25	
250 pC (100 fs) 787 eV 1.5 mJ 2500 A	4.79	
250 pC (100 fs) 769 eV 1.5 mJ 2500 A	5.24	



Intrinsic x-ray bandwidth

Electron kinetic energy (eV)

Conditions	Intrinsic x-ray pulse bandwidth (from 2s photopeak) (eV) (FWHM)	Intrinsic x-ray pulse bandwidth (from 2p photopeak) (eV) (FWHM)	Average bandwidth (eV) (FWHM)	%
40 pC (<10 fs) 850 eV 0.3 mJ 4500 A	4.3	4.5	4.4	0.5 %
250 pC (100 fs) 787 eV 1.5 mJ 2500 A	7.1	7.8	7.45	0.9%
250 pC (100 fs) 769 eV 1.5 mJ 2500 A	7.7	7.8	7.77	1%

Is the ¹D Auger line broadened on 1s-2p resonance?



E.P. Kanter *et al.*, PRL (2011)

Theory from N. Rohringer and R. Santra

SASE vs Gaussian pulse for Rabi flopping





Summary: resonant x-ray processes at high intensity

- First hint of Rabi cycling for inner-shell electrons: Ne 1s 2p resonance
- Need XFEL with improved longitudinal coherence SEEDING
 - Quantum control multidimensional spectroscopies
 - Single particle imaging (reduced radiation damage & increased x-ray intensity)
- "Hidden" resonances critical in atomic response to ultraintense x-rays
 - Enhanced two-photon absorption probability
 - Doumy *et al.*, PRL **106**, 083002 (2011)
 - Ionization beyond sequential single photon model
 - Schorb et al., PRL 108, 233401 (2012) Ar
 - Rudek *et al.*, Nat. Phot. **6**, 858 (2012) Xe
 - Rudek et al., Phys. Rev. A 87, 023413 (2013) Kr



Improving the X-ray Laser

Self-seeding techniques and their importance for XFELs

SASE pulses, baseline mode of operation: poor longitudinal coherence



Figure 5.2.4 Temporal (top) and spectral (bottom) structure for 12.4 keV XFEL radiation from SASE 1. Smooth lines indicate averaged profiles. Right side plots show enlarged view of the left plots. The magnetic undulator length is 130 m.

Source: The European XFEL TDR - DESY 2006-097 (2006)

$$\frac{\Delta\omega}{\omega}\sim 2\rho\sim 10^{-3}$$

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$$\left(\frac{\Delta\omega}{\omega}\right)_{spike} \sim \frac{1}{\sigma_T \omega} \sim 10^{-5}$$

- Hundreds of longitudinal modes
- A lot of room for improvement
- Self-seeding schemes answer the call for increasing longitudinal coherence

European

Hard x-ray self-seeding proposed 2010





Diamond C(004): 100 μ m λ = 0.15 nm, θ_B = 57°

arXiv:1008.3036v1 Geloni ³⁸

Hard x-ray self seeding realized Jan 2012 - P. Emma et al.



Bandwidth <10⁻⁴ at 8-9 keV and tunable But ... did not achieve saturation and power jitter still present

Three lab collaboration:



J. Amann et al., Nature Photonics 6, 693 (2012)

On the road to a TW FEL: LCLS-TN-11-3

GENESIS simulation



W.M. Fawley¹, J. Frisch¹, Z. Huang¹, Y. Jiao¹, H.-D. Nuhn¹, C. Pellegrini^{1,2}, S. Reiche³, J. Wu^{1†}
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 ³Paul Scherrer Institute, Villigen PSI, 5232, Switzerland

Also arXiv Jun 2013: S. Serkez *et al.* 10 TW FEL, 10¹⁴ x-rays, 10 fs @ 3.5 keV

Towards single particle imaging

LETTER

doi:10.1038/nature09748

Single mimivirus particles intercepted and imaged with an X-ray laser Seibert et al., Nature 470, 78 (2011)





Mimivirus

-Largest known virus – 0.75 μm
-Does not crystallize
-Large for 3D cryoelectron microscopy
Single Shot Scattering Pattern 2D: 32 nm resolution
Set of 198 scattering patterns 3D: reconstruction to 120 nm



on Ekeberg et al., PRL (2015)

Viewpoint

X-Ray Imaging of a Single Virus in 3D

"And there are still open questions on the impact of electronic damage on x-ray scattering on femtosecond time scales: The above-mentioned work by Neutze et al. tracked the movements of the atomic nuclei of the biomolecule, showing they don't move on the fewfemtosecond timescale of an x-ray pulse. Electrons, however, move faster than nuclei. Since electrons are what scatters x rays, it is yet to be confirmed that fewfemtosecond pulses can probe an unperturbed electronic structure." - Keith Nugent



March 2, 2015

Beyond the sequential single photon ionization model



@ 480 eV

sequential single photon limit 10+ observe 13+ Schorb et al., PRL (2012)



Resonance-enabled x-ray multiple ionization



B. Rudek et al., Nat. Phot. (2012)

Tracking electronic configurations during XFEL pulse including resonances!

	# of ECs with no RE	# of ECs with RE
Ar	1.33×10^{3}	2.85×10^{13}
Kr	3.05×10^{5}	2.08×10^{19}
Xe	7.06×10^7	9.05×10^{22}



Ultraintense hard x-ray interactions: seeded v SASE

CXI endstation experiment Apr 2014: I >10²⁰ W/cm²

5.5 – 8.3 keV, 2 mJ, 30-40 fs "100 nm" focus 10⁵ x-rays/Å²

...but no shot-by-shot spectral monitor



Joined/Artem Rudenko, Daniel Rolles & team

Observations:

Ar¹⁸⁺: fully stripped
Kr³⁴⁺: two 1s electrons
Xe⁴⁸⁺: 6 electrons
Largely consistent with
sequential single photon
CH₃I: sum of charges 55+

Calculations by Phay Ho: Identified photon energies where seeded pulses reduce absorption & radiation damage relative to SASE pulses.

High Intensity X-ray Imaging of Nanosystems

- Intense x-ray pulses lead to efficient stripping of e⁻ from atoms/ions and significant displacement of e⁻, ions/atoms from equilibrium during the pulse – nanoplasma
- Our MC/MD model includes: Elastic Scattering, Photoionization, Auger, Fluorescence, Resonant excitation, Three-body recombination, Electron impact ionization, Lattice dynamics
- Investigate the impact of electronic damage on various observables (ion, photoelectron, Auger and fluorescence spectra, x-ray diffraction)
- Recorded diffraction patterns depend strongly on pulse parameters.





AMO "solutions" at the ultraintense x-ray frontier

- Fundamental nature of x-ray damage at high intensity

 Coulomb explosion
 electronic damage
 behavior at 10²² W/cm² – 1Å
 nanoplasma formation
- Expt'l AMO observables: Ion, photoelectron, Auger, fluorescence and ... x-ray diffraction pattern
- Computational studies on large scale systems

Phay Ho Chris Knight





Atomistic computational studies based upon Monte Carlo/Molecular Dynamics

New XFEL capabilities becoming available







International hard X-ray FELs here and on the horizon















Summary

- Our atomic physics expts have established fundamental understanding of the response of matter to ultraintense XFEL irradiation
 - sequential single photon ionization dominates
 - intensity-induced x-ray transparency (frustrated absorption)
 - intense x-rays can "control" inner-shell electron dynamics
 - resonances can be critical in XFEL interactions
- This fundamental understanding will aid in the quest for single molecule imaging and other applications, e.g. high energy density matter

- AMO methods (ion, electron, photon) in concert with computational studies promise fuller understanding of radiation damage in extended systems

 Future is bright with well-characterized ultraintense x-ray lasers, multiple pulse configurations ... with 1000x higher intensity.







