

Defect formation

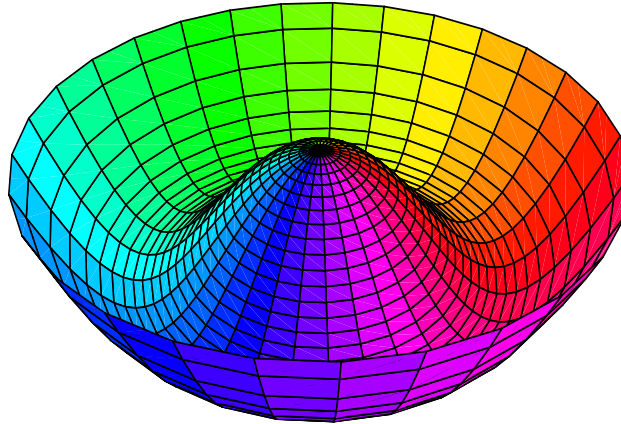
Arttu Rajantie

Topological Defects

- Localised, topologically stable objects:
Different vacuum in different directions
- Exist in systems with
 - spontaneously broken symmetry
 - topologically non-trivial vacuum manifold \mathcal{M} :
 - $\pi_0(\mathcal{M}) \neq 1$: Domain walls
 - $\pi_1(\mathcal{M}) \neq 1$: Vortices/strings
 - $\pi_2(\mathcal{M}) \neq 1$: Monopoles
 - $\pi_3(\mathcal{M}) \neq 1$: Textures



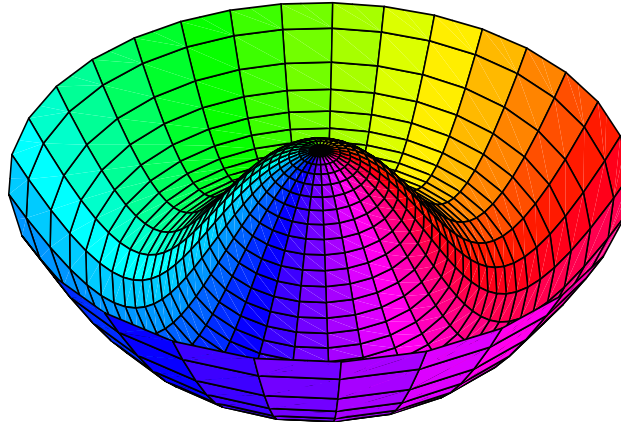
Strings



$$\mathcal{L} = \partial_\mu \phi^* \partial^\mu \phi - m^2 \phi^* \phi - \lambda (\phi^* \phi)^2$$

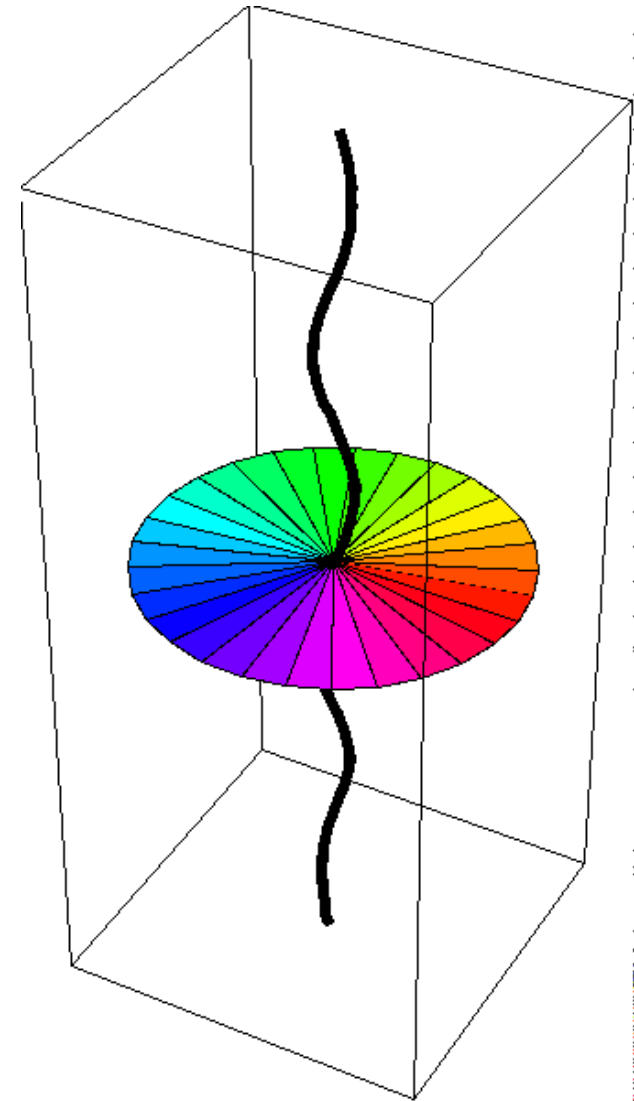
- Continuous U(1) symmetry $\phi \rightarrow e^{i\alpha} \phi$ in
- When $m^2 < 0$, spontaneously broken:
Circle of degenerate vacua $\langle \phi \rangle = v e^{i\theta}$

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- Continuous U(1) symmetry $\phi \rightarrow e^{i\alpha} \phi$ in
- When $m^2 < 0$, spontaneously broken:
Circle of degenerate vacua $\langle \phi \rangle = v e^{i\theta}$
- Nielsen-Olesen string:
 - Different vacuum in different directions
 - Phase angle θ changes by 2π around the string



Cosmic Strings (Kibble 1976)

- Stable, massive, line-like objects
- Tension μ = energy/length
- Formed in the early universe?
- Predicted by many GUTs:

$$G\mu \sim (M_{\text{GUT}}/M_{\text{Pl}})^2 \sim 10^{-6}$$



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- Tension μ = energy/length
- Formed in the early universe?
- Predicted by many GUTs:
- Cosmic superstrings (Copeland et al 2004):

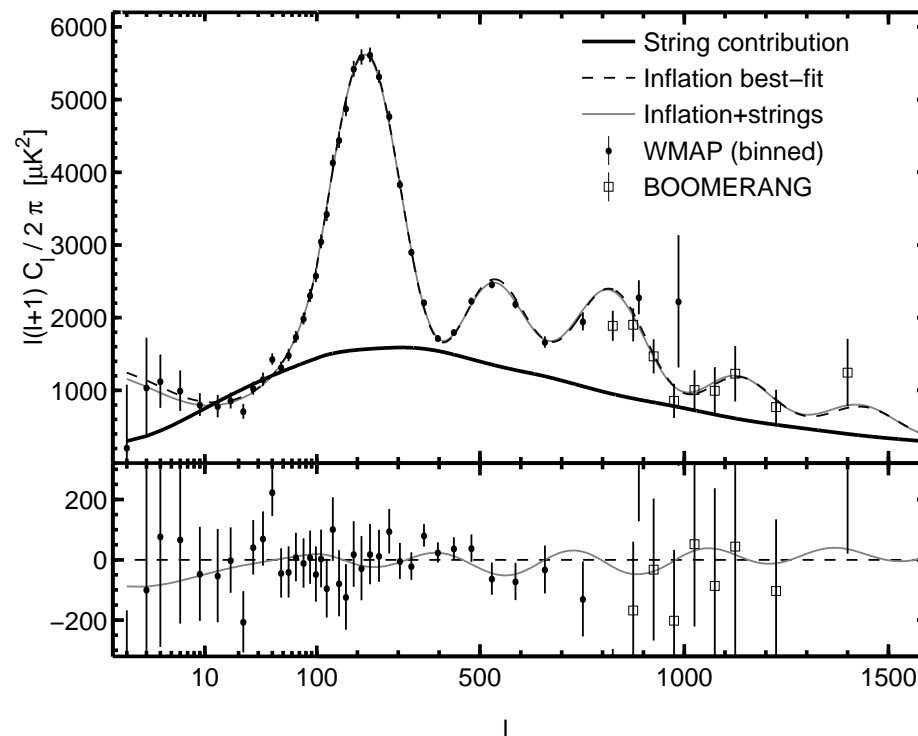
$$G\mu \sim (M_{\text{GUT}}/M_{\text{Pl}})^2 \sim 10^{-6}$$

$$G\mu \sim 10^{-12} \dots 10^{-6}$$



Observing Cosmic Strings?

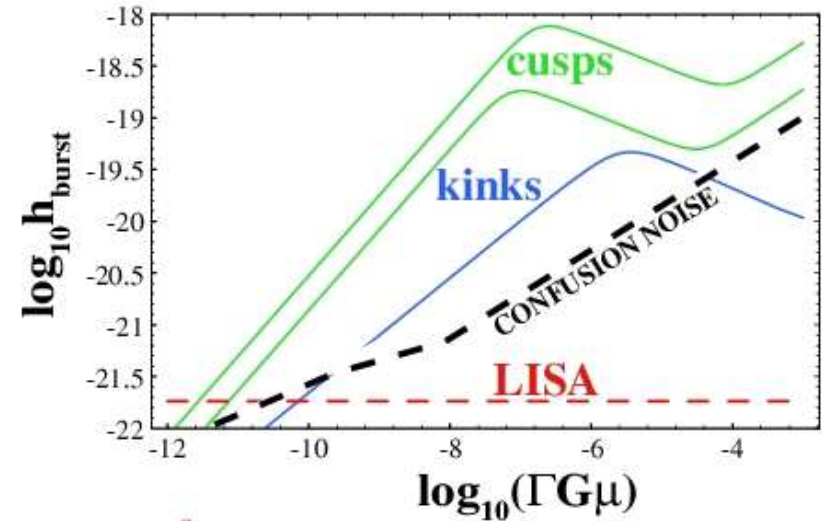
- Observational constraints:
 - Cosmic microwave background
 - Temperature $G\mu \lesssim 10^{-6}$
 - Polarisation $\Rightarrow G\mu \lesssim 10^{-7}$?
(Bevis et al 2007)



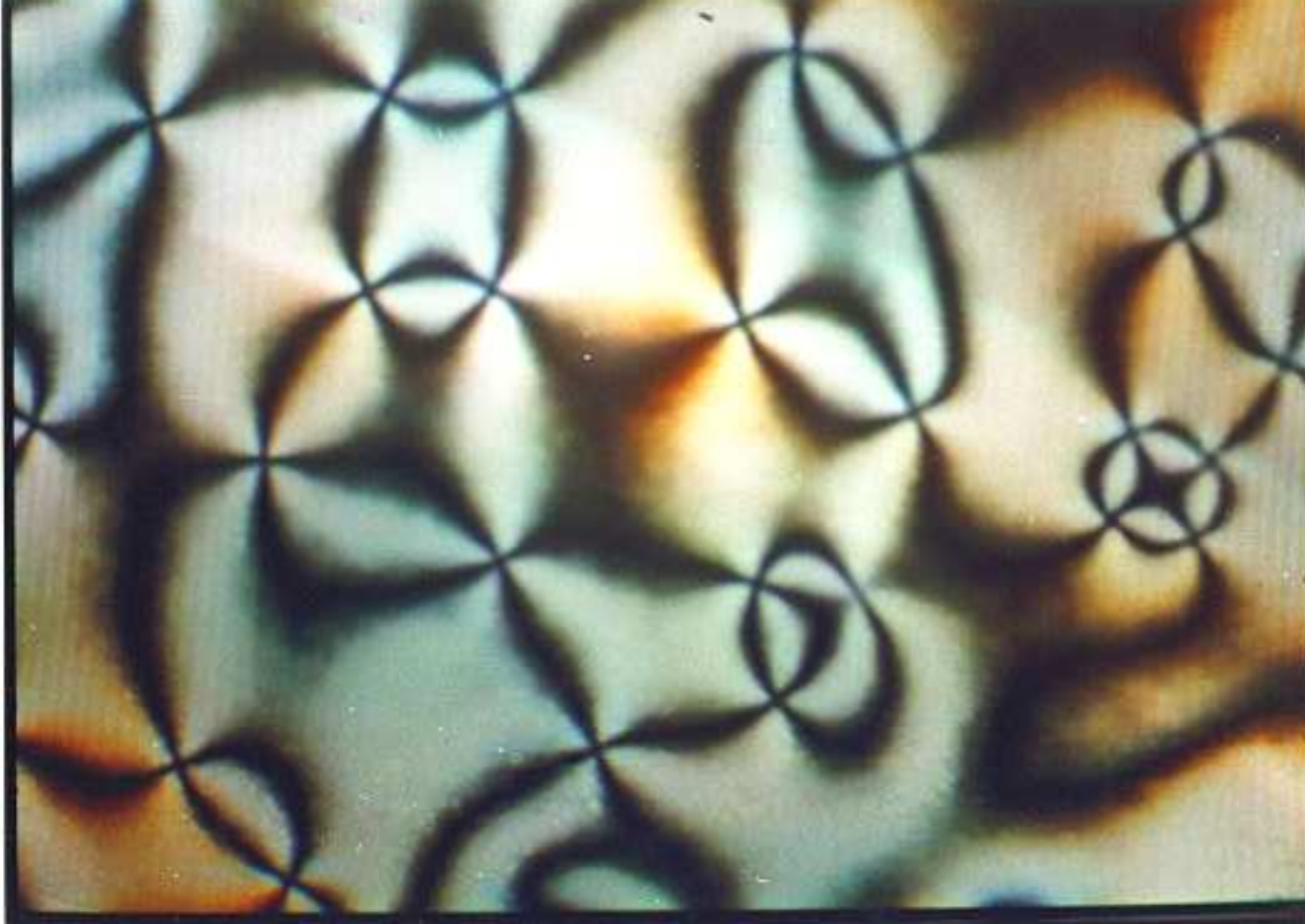
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Observing Cosmic Strings?

- Observational constraints:
 - Cosmic microwave background
 - Temperature $G\mu \lesssim 10^{-6}$
 - Polarisation $\Rightarrow G\mu \lesssim 10^{-7}$?
(Bevis et al 2007)
 - Gravitational waves
 - Pulsar timing $G\mu \lesssim 10^{-6}$ (Jenet et al 2006)
 - Gravity wave experiments $\Rightarrow G\mu \lesssim 10^{-10}$?
(Damour&Vilenkin 2001)

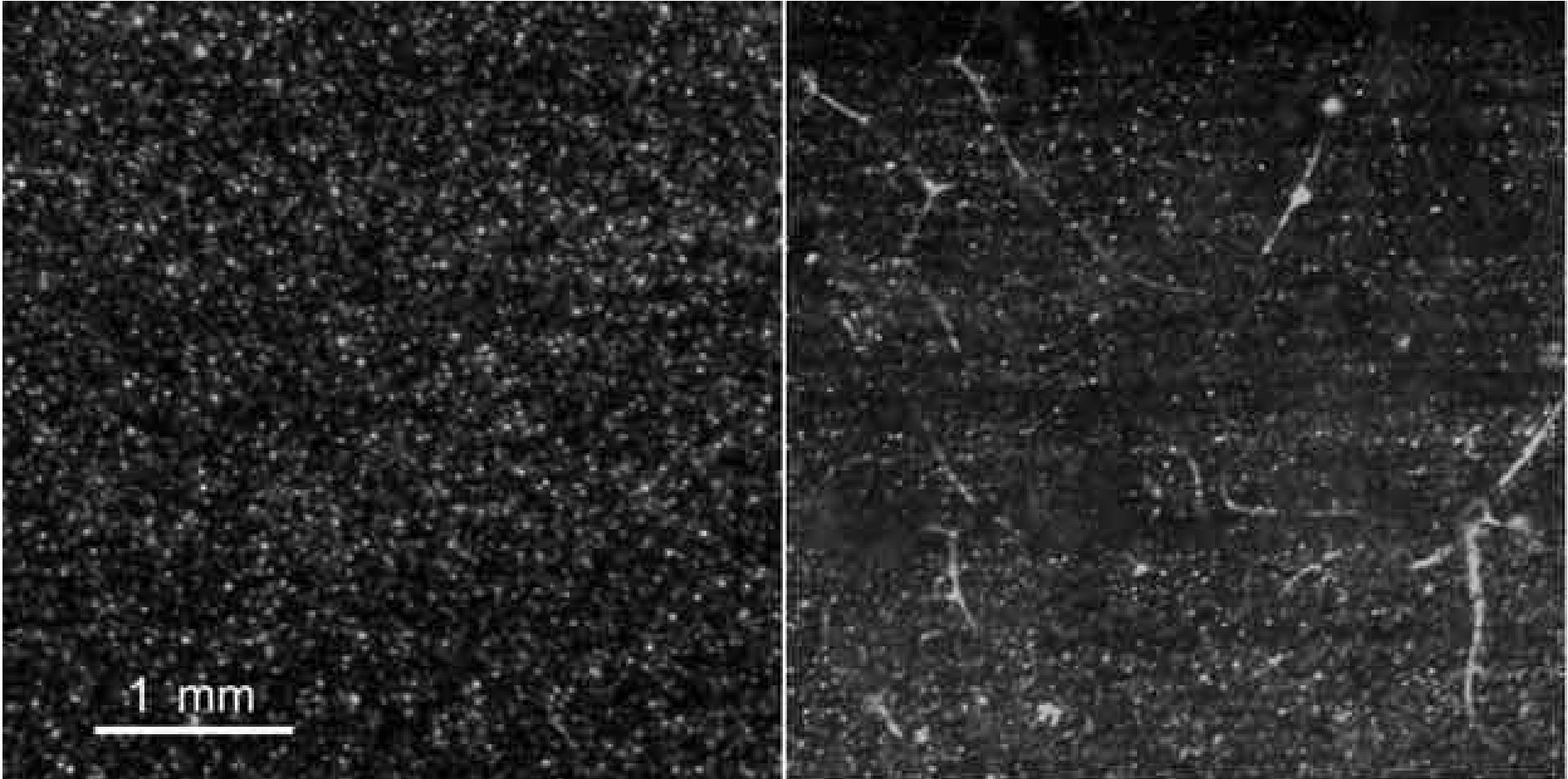


Liquid Crystal Vortex



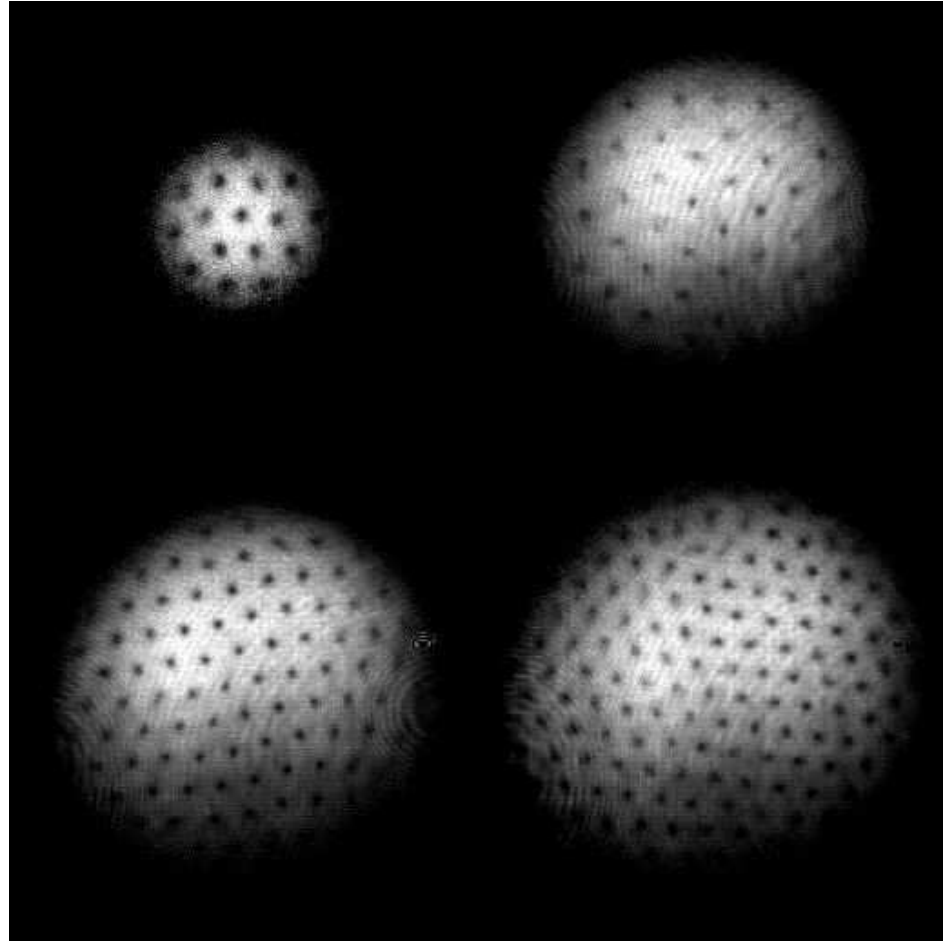
- Orientation of molecules (Digal et al 1999)

Superfluid Vortex



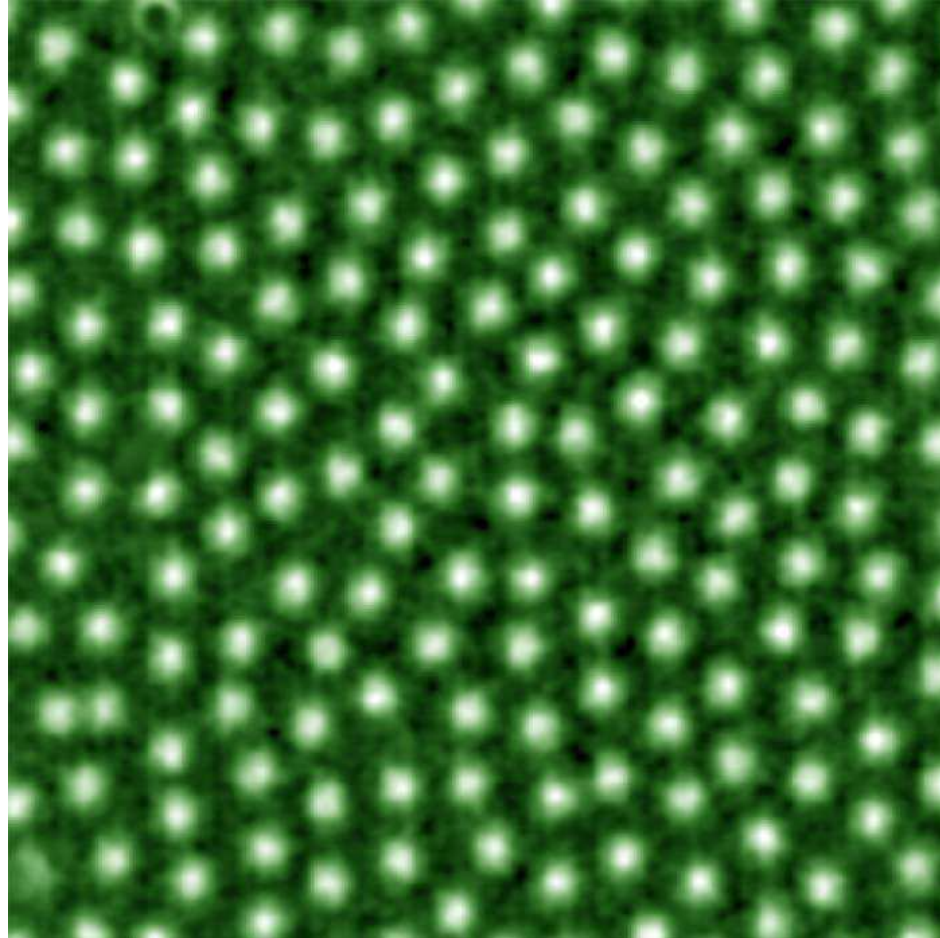
- Complex phase of atom wave function (Bewley et al 2006)

BEC Vortex



- Complex phase of atom wave function (Abo-Shaeer et al 2001)

Superconductor Vortex

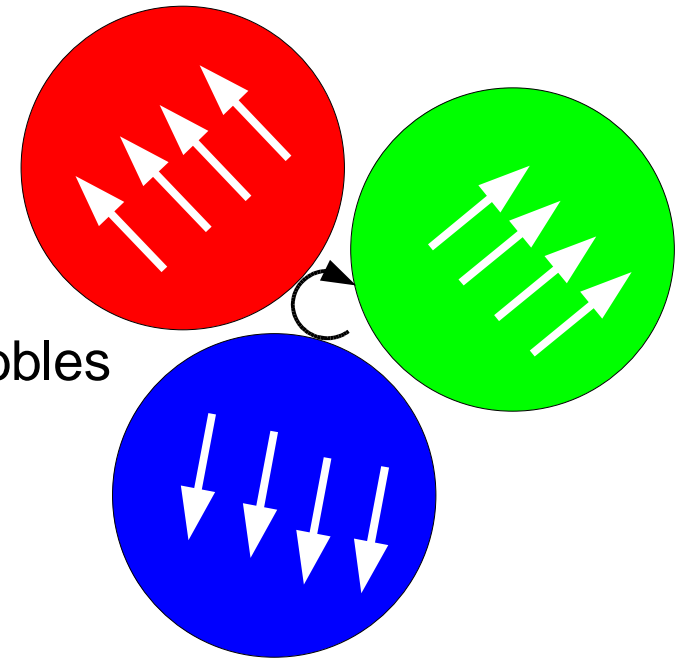


- Complex phase of Cooper pair wave function (Goa et al)

Kibble Mechanism in First Order Transitions (Kibble 1976)

- Spontaneous symmetry breaking:
Field has to “choose” a direction
- Typical bubble size R
 - Order parameter uncorrelated between bubbles
 - Finite probability to form a string
whenever three bubbles meet (Kibble 1976)
 - Number density per cross-sectional area:

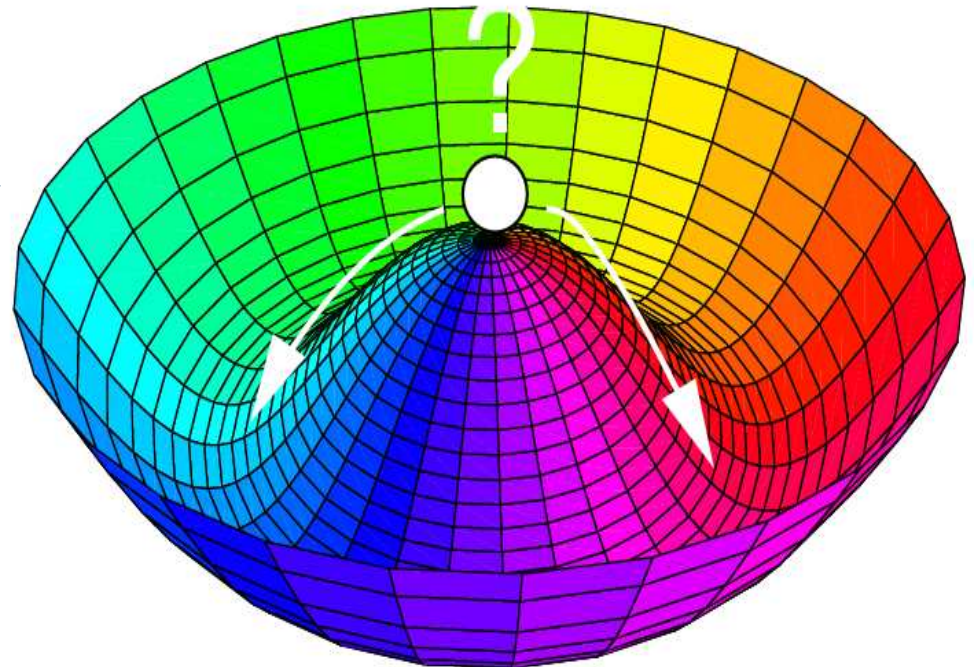
$$n \sim R^{-2}$$



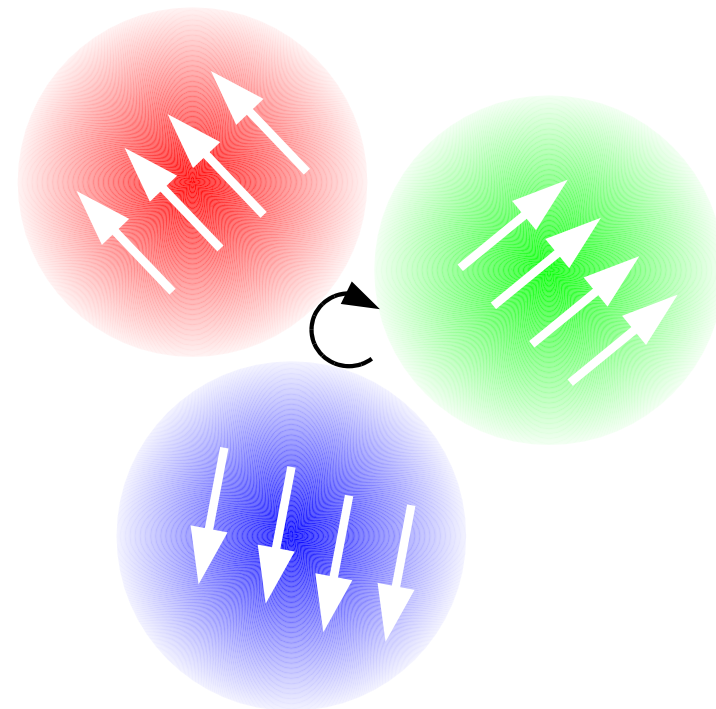
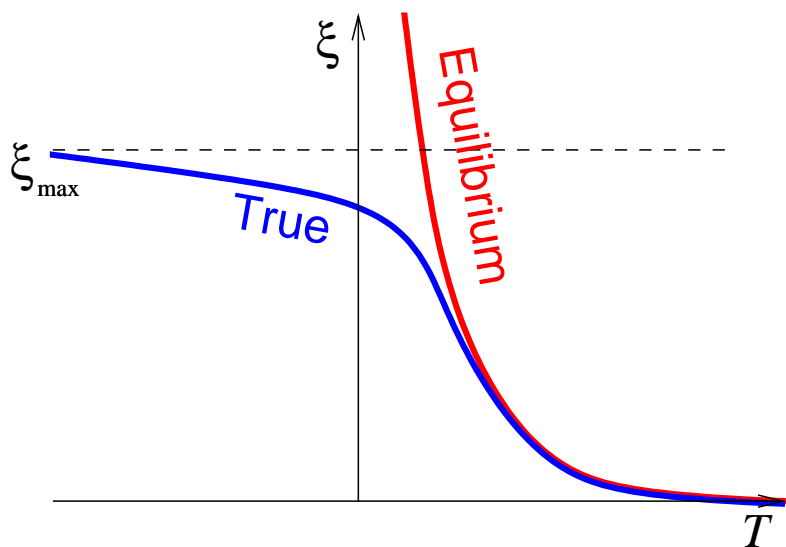
Second Order Phase Transition

$$V(\phi) = m^2 \phi^* \phi + \lambda (\phi^* \phi)^2$$

- Gradual change from $m^2 > 0$ to $m^2 < 0$:
Transition from symmetric to broken phase
- Caused by decreasing temperature,
end of inflation etc
- Choose a direction
- Uncorrelated at distances longer
than the correlation length $\xi < 1/H$
 \Rightarrow Kibble mechanism
- Number density:
 $n \sim \xi^{-2}$

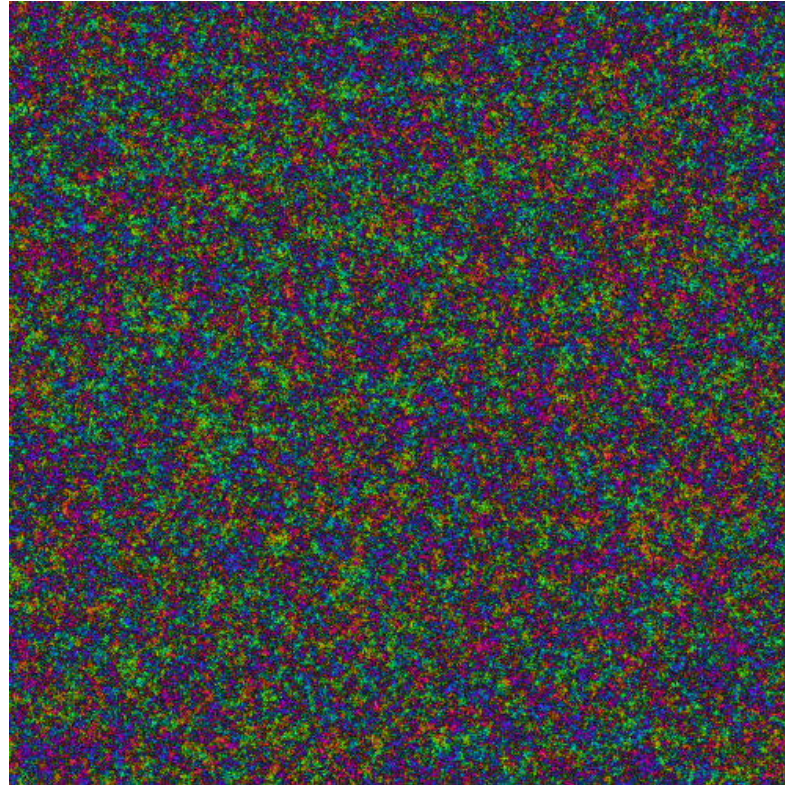


Kibble-Zurek Mechanism (Zurek 1985)



- Correlation length grows as $\xi \sim (T - T_c)^{-\nu}$
- Critical slowing down: $\tau \sim (T - T_c)^{-\mu}$
- ξ freezes out to $\hat{\xi} \propto \tau_Q^{\nu/(1+\mu)}$:
Sets the string number density

Simulation Tests



- Classical field theory simulations confirm predictions

Gauge Fields: Abelian Higgs Model

$$\epsilon = \frac{1}{2} \vec{B}^2 + |\vec{D}\phi|^2 + m^2 \phi^* \phi + \lambda (\phi^* \phi)^2$$

- $\vec{D} = \vec{\nabla} + ie\vec{A}$, $\vec{B} = \vec{\nabla} \times \vec{A}$
- Gauge symmetry $\phi(x) \rightarrow e^{i\alpha(x)}\phi(x)$, $\vec{A} \rightarrow \vec{A} - (1/e)\vec{\nabla}\alpha$
 - Phase angle not θ physical observable - Kibble argument?
- Higgs phase:
 - Minimize energy: $\vec{D}\phi = 0 \leftrightarrow e\vec{A} \approx \vec{\nabla}\theta$
 - Gauge field cancels gradient energy

Cold Electroweak Baryogenesis

- Electroweak theory: $\pi_3(\mathcal{M}) = \mathbb{Z}$
⇒ Textures
- Gauge field compensates: $N_{CS} = N_w$
 - Either N_{CS} or N_w changes, depending on texture size
 - Biased by CP violation
 - Anomaly $\Delta B = 3\Delta N_{CS}$ leads to baryon asymmetry (Turok&Zadrozny 1990)
- Can take place at the end of inflation: $T = 0$
(Krauss&Trodden; Copeland,Lyth,AR&Trodden; Tranberg&Smit)

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(Krauss&Trodden; Copeland,Lyth,AR&Trodden; Tranberg&Smit)
 - In thermal transition dominated by gauge dynamics:
Thermal sphalerons

Flux Quantisation

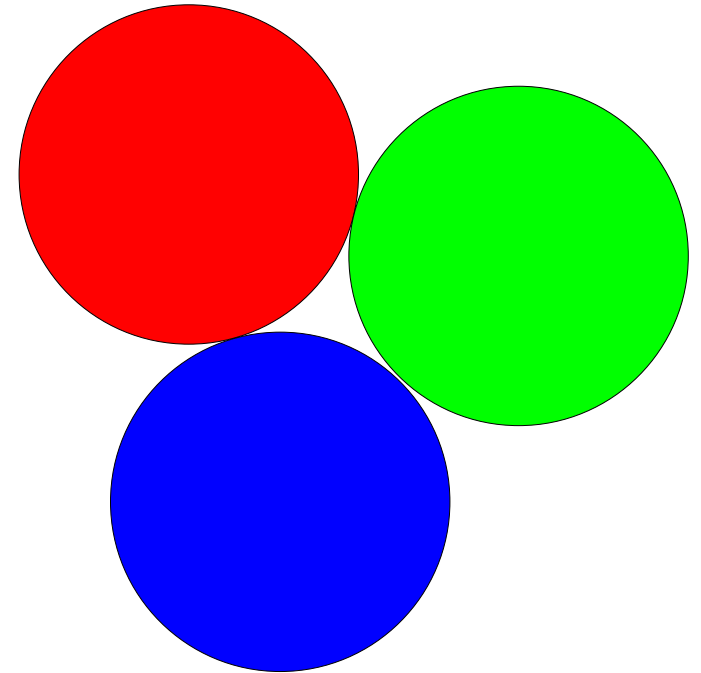
- Minimize energy: $\vec{D}\phi = 0 \leftrightarrow e\vec{A} \approx \vec{\nabla}\theta$
- No magnetic field in the Higgs phase:
 $\vec{\nabla} \times \vec{A} \approx (1/e)\vec{\nabla} \times \vec{\nabla}\theta = 0$
- Confined into strings

$$\Phi = \oint d\vec{x} \cdot \vec{A} = (1/e) \oint d\vec{x} \cdot \vec{\nabla}\theta = N_w \Phi_0$$

- Flux quantum $\Phi_0 = 2\pi/e$ per string

First Order Phase Transition

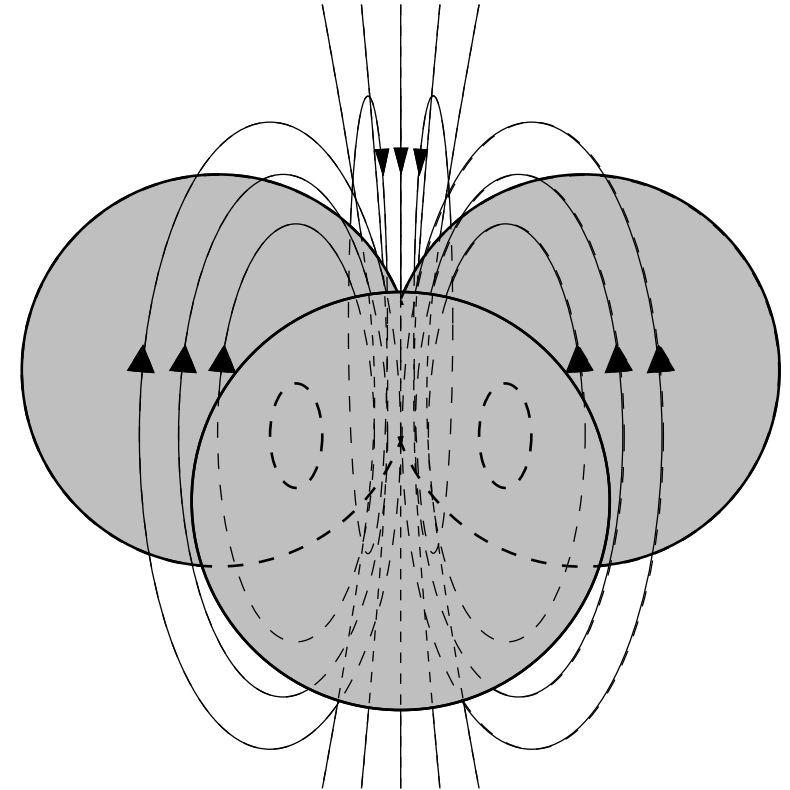
- Bubble radius R
- Hole of area $\sim R^2$ between bubbles



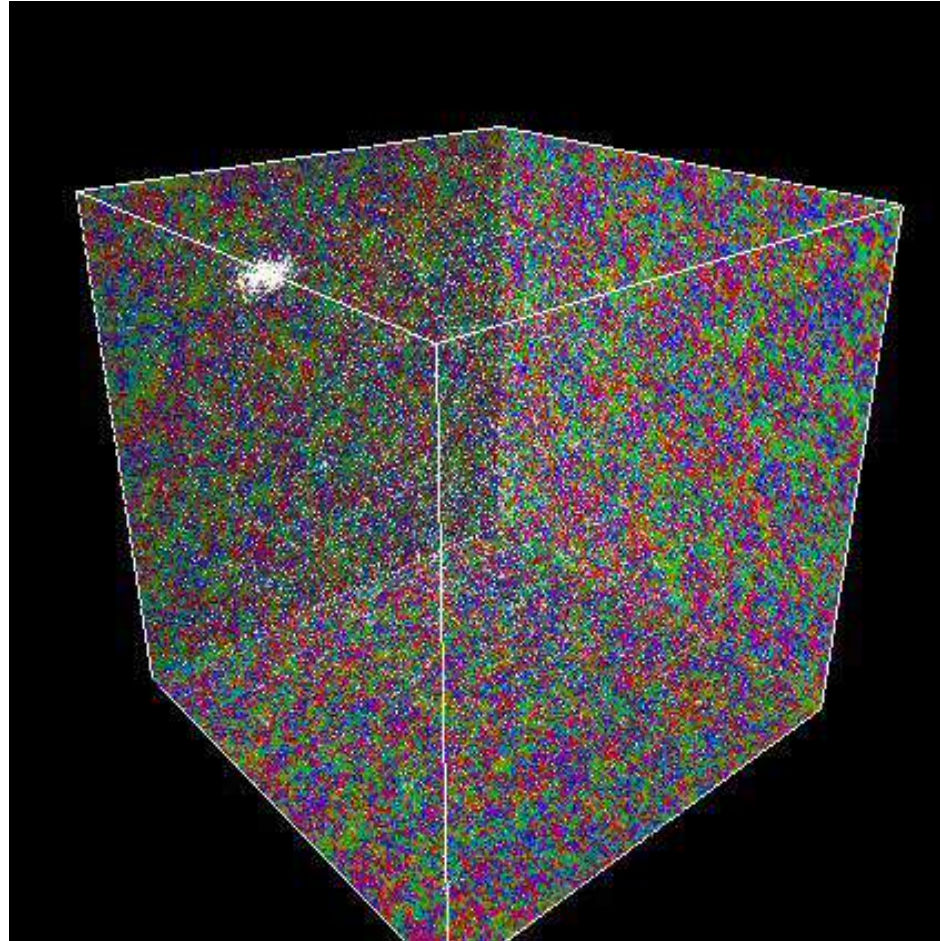
First Order Phase Transition

- Bubble radius R
- Hole of area $\sim R^2$ between bubbles
- Thermal flux through hole: $\Phi \approx \sqrt{RT}$
- Conservation of flux:
Must get confined in a string
with winding $N_w = \Phi/\Phi_0 \approx e\sqrt{RT}$

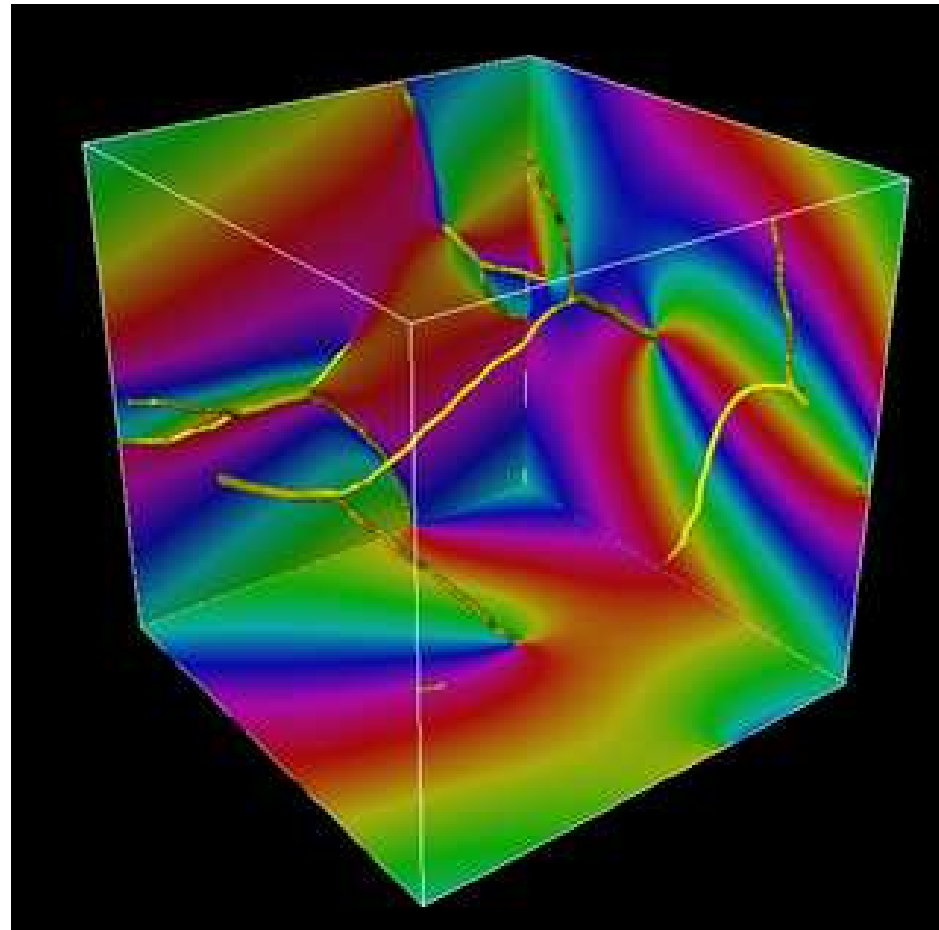
(Donaire&AR 2006)



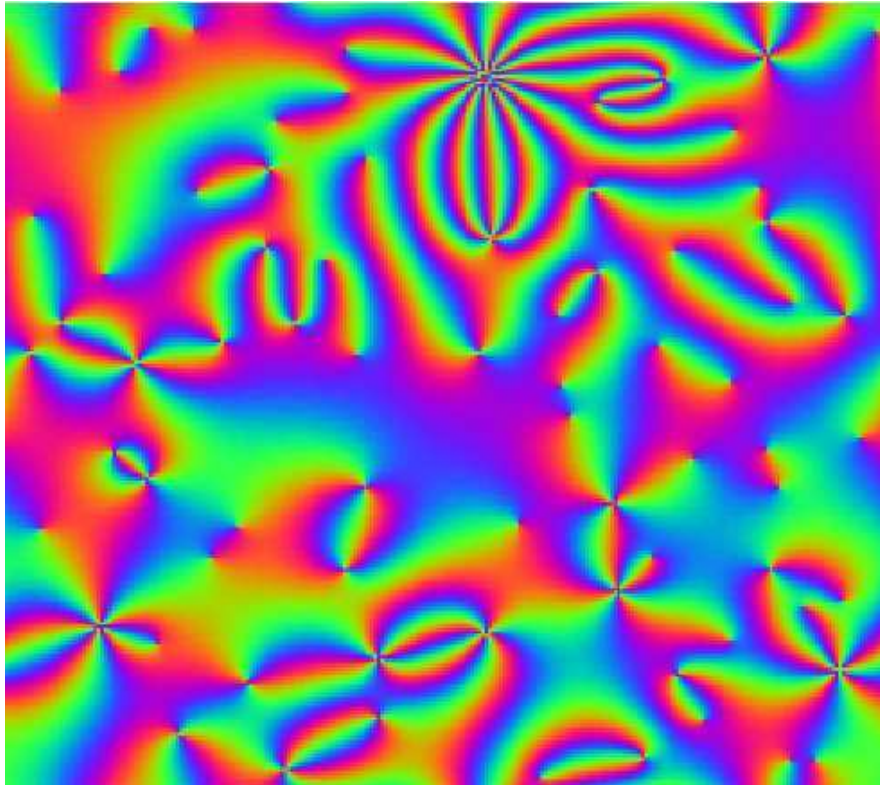
Simulation of a First Order Transition



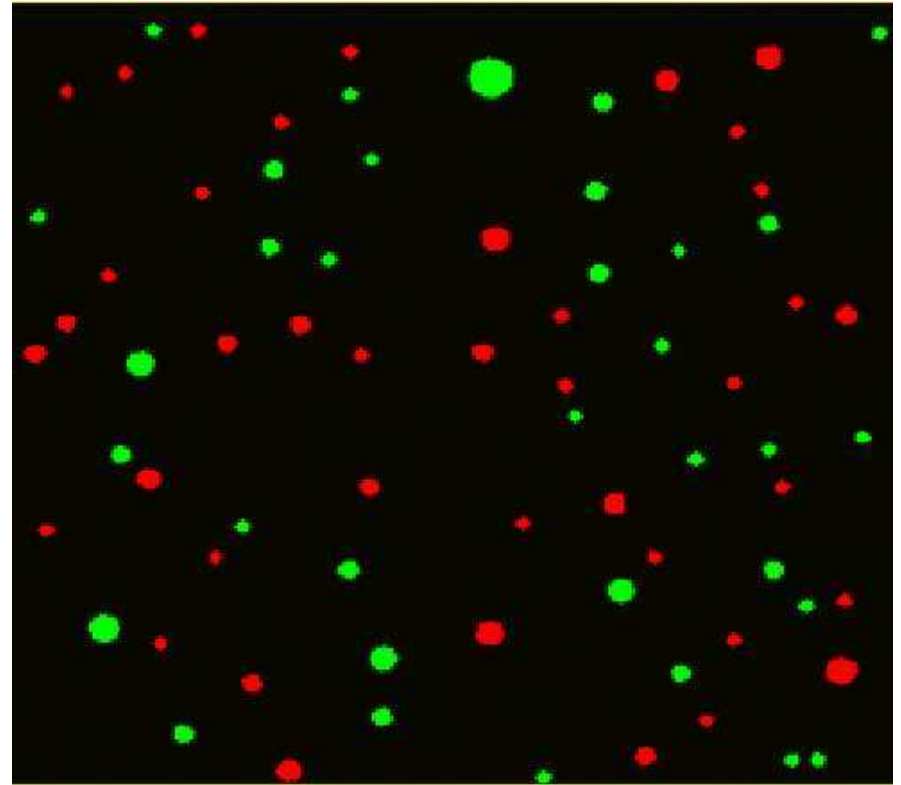
Simulation of a First Order Transition



First Order Phase Transition in 2D



Phase angle



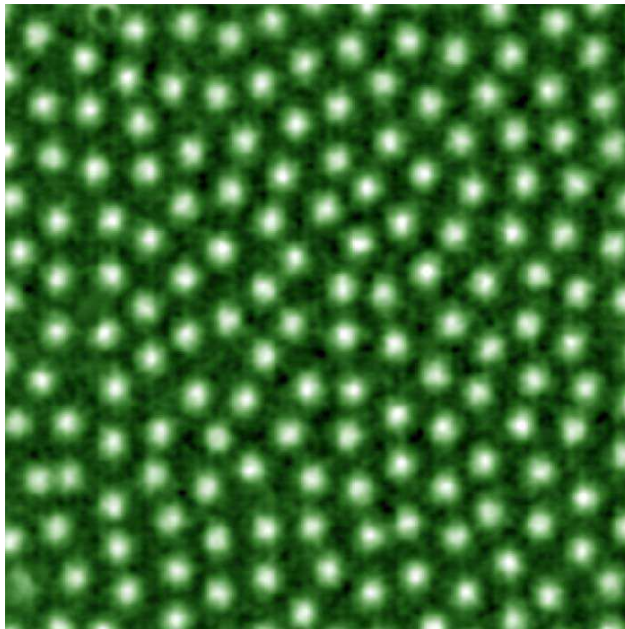
Magnetic field

Flux Trapping in a Continuous Transition

- Magnetic flux is conserved

$$\frac{d}{dt}\Phi = - \oint_{\partial C} d\vec{x} \cdot \vec{E}$$

- Uniform initial field: String number density $\vec{\rho} = (e/2\pi)\vec{B}$

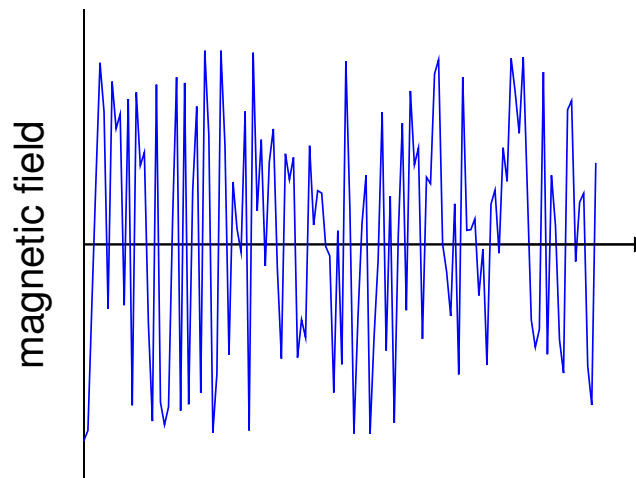


Flux Trapping in a Continuous Transition

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$$\frac{d}{dt}\Phi = - \oint_{\partial C} d\vec{x} \cdot \vec{E}$$

- Non-uniform field: Short-wavelength $\lambda \lesssim \lambda_c$ modes decay

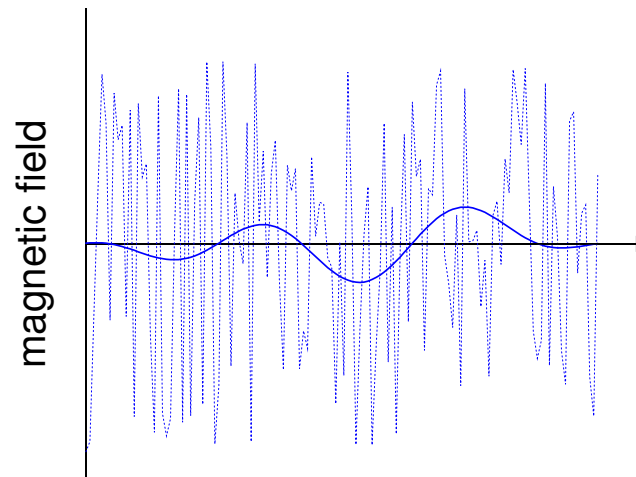


Flux Trapping in a Continuous Transition

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- Long-wavelength modes survive

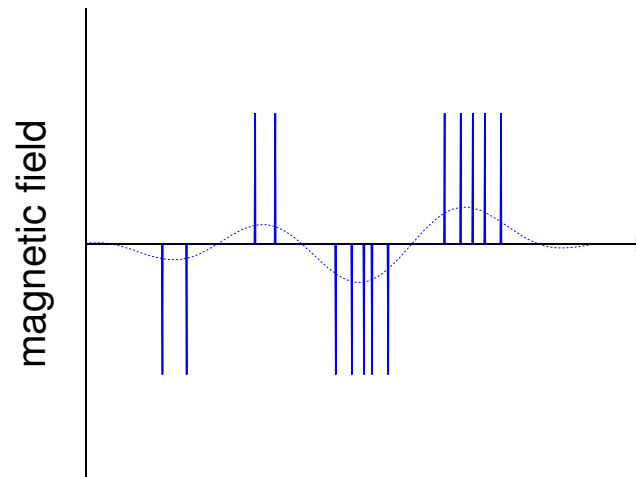


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- Long-wavelength modes form strings $\vec{\rho} = (e/2\pi)\vec{B}$ (Hindmarsh&AR 2000)
- Correlator at long distances $r \gg \lambda_c$:

$$\langle \rho_i(x)\rho_j(y) \rangle = (e/2\pi)^2 \langle B_i(x)B_j(y) \rangle$$

- Given by a scalar function $G(k)$:

$$\langle \rho_i(\vec{k})\rho_j(\vec{q}) \rangle = (2\pi)^3 \delta(\vec{k} + \vec{q}) \left(\delta_{ij} - \frac{k_i k_j}{k^2} \right) G(k)$$

Long-Range Correlations (AR 2009)

- Initial state:
 - Thermal ($T > 0$):

$$G(k) = \frac{e^2}{4\pi^2} G^B(k) \approx \frac{e^2}{4\pi^2} \frac{k}{2} \coth \frac{k}{2T}$$

$$\langle \rho_i(\vec{x}) \rho_j(\vec{y}) \rangle \sim -\frac{e^2 T}{|\vec{x} - \vec{y}|^3}$$

- Vacuum ($T = 0$):

$$G(k) = \frac{e^2}{4\pi^2} G^B(k) \approx \frac{e^2}{4\pi^2} \frac{k}{2}$$

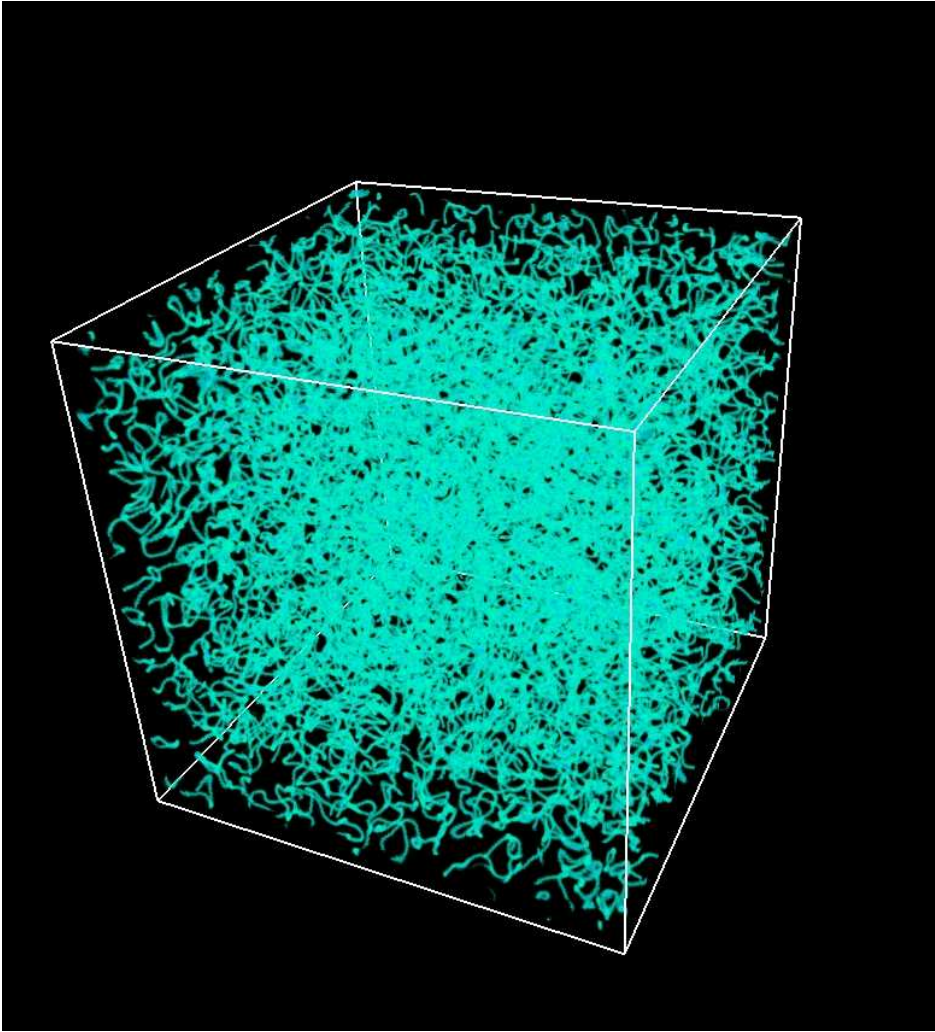
$$\langle \rho_i(\vec{x}) \rho_j(\vec{y}) \rangle \sim -\frac{e^2}{|\vec{x} - \vec{y}|^4}$$

3D Simulation

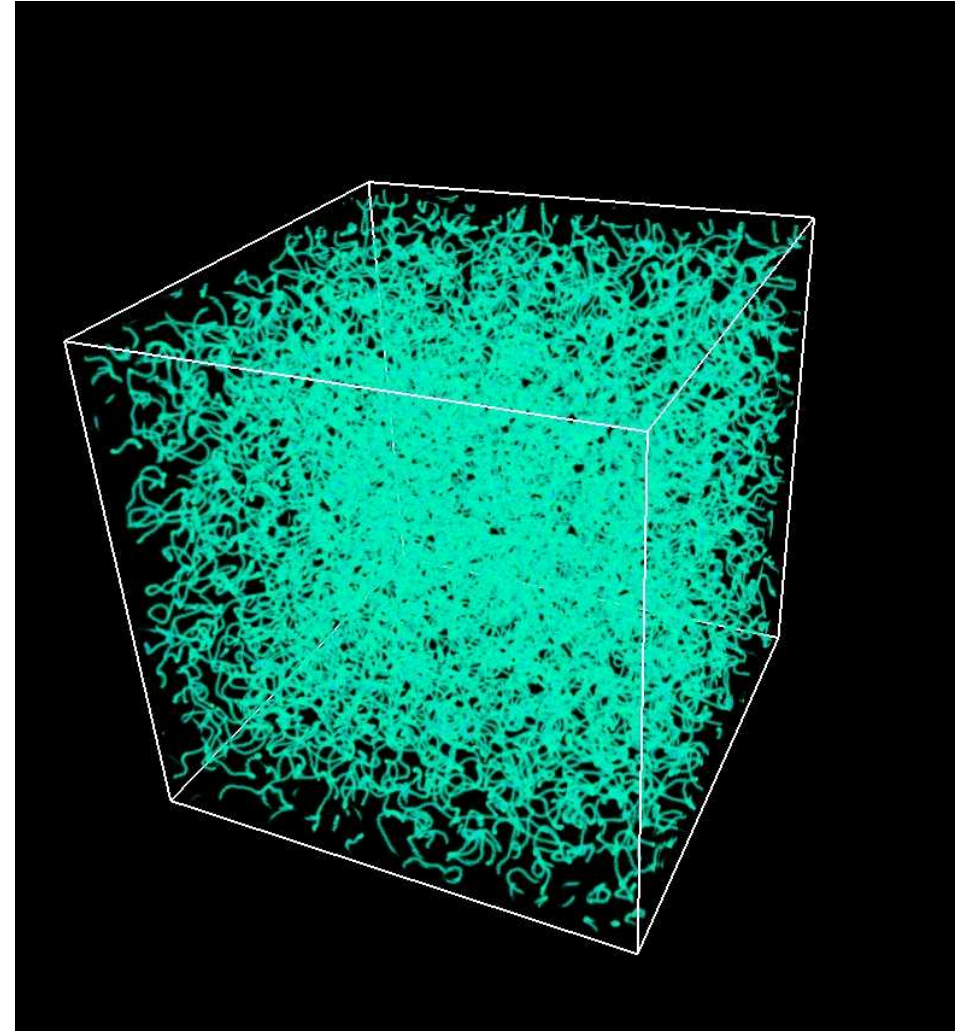
- 256x256x256 lattice, spacing $\delta x = 1$
- Thermal initial conditions with $T_{\text{ini}} = 0.5$
- Radiation dominated with $H_{\text{ini}} = 0.1$: Evolve until $a = 2$
- Scalar coupling $\lambda = 1$,
 m^2 such that transition takes place at $a = \sqrt{2}$

$$\begin{aligned}\partial_\tau^2 \tilde{\phi} &= \vec{D}^2 \tilde{\phi} + (m^2 a^2 + \partial_\tau^2 a/a) \tilde{\phi} - 2\lambda |\tilde{\phi}|^2 \tilde{\phi} \\ \partial_\tau \tilde{E}_i &= \partial_j F_{ij} + 2e \text{Im} \tilde{\phi}^* D_i \tilde{\phi}, \\ \partial_i \tilde{E}_i &= 2e \text{Im} \tilde{\phi}^* \partial_\tau \tilde{\phi}\end{aligned}$$

3D Simulation

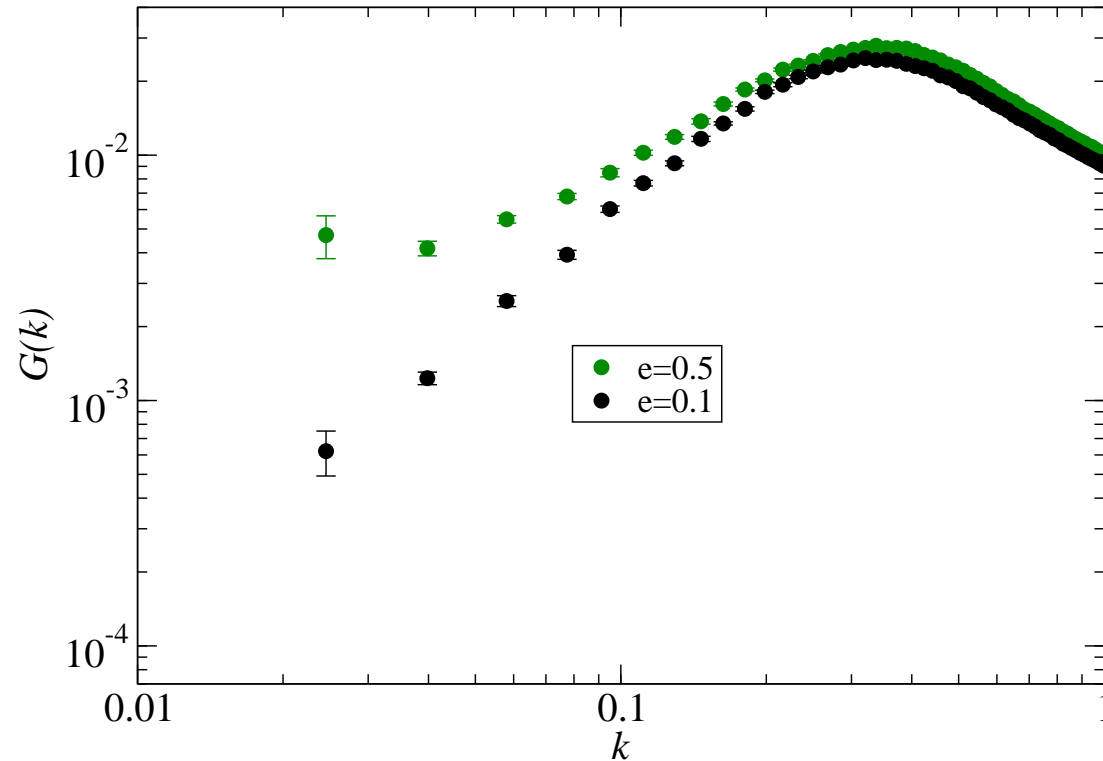


$e = 0.1$



$e = 0.5$

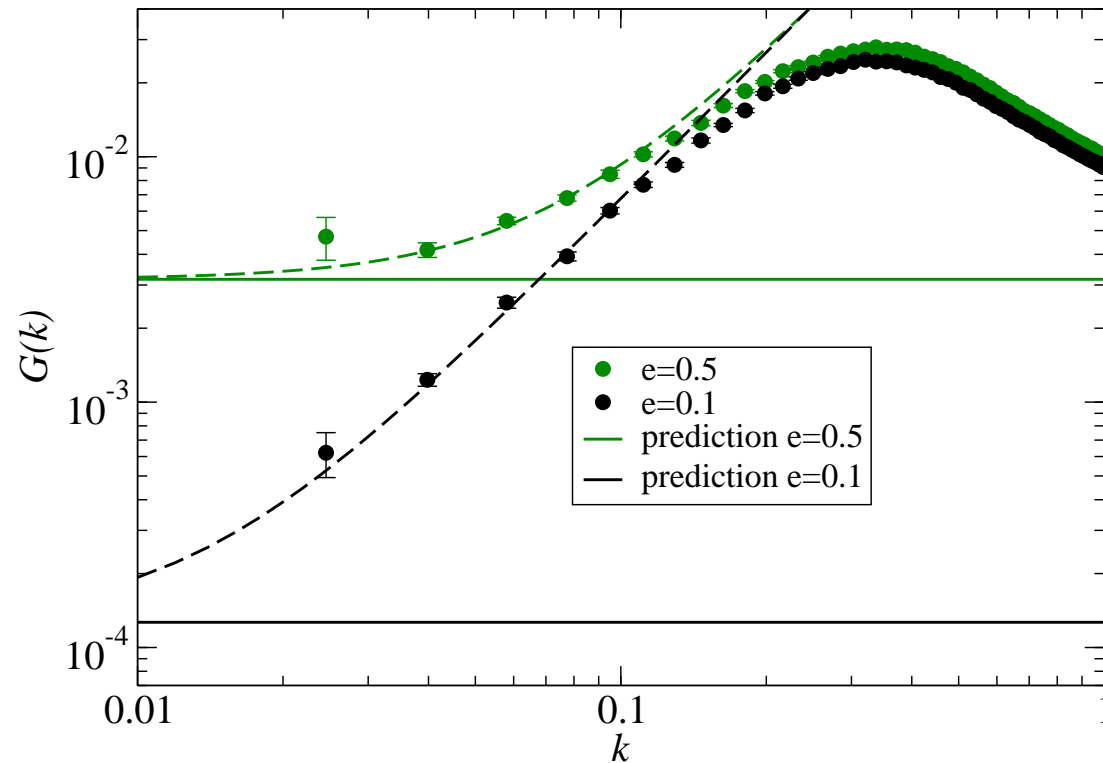
3D Simulation



- Winding number correlator:

$$\langle \rho_i(\vec{k}) \rho_j(\vec{q}) \rangle = (2\pi)^3 \delta(\vec{k} + \vec{q}) \left(\delta_{ij} - \frac{k_i k_j}{k^2} \right) G(k)$$

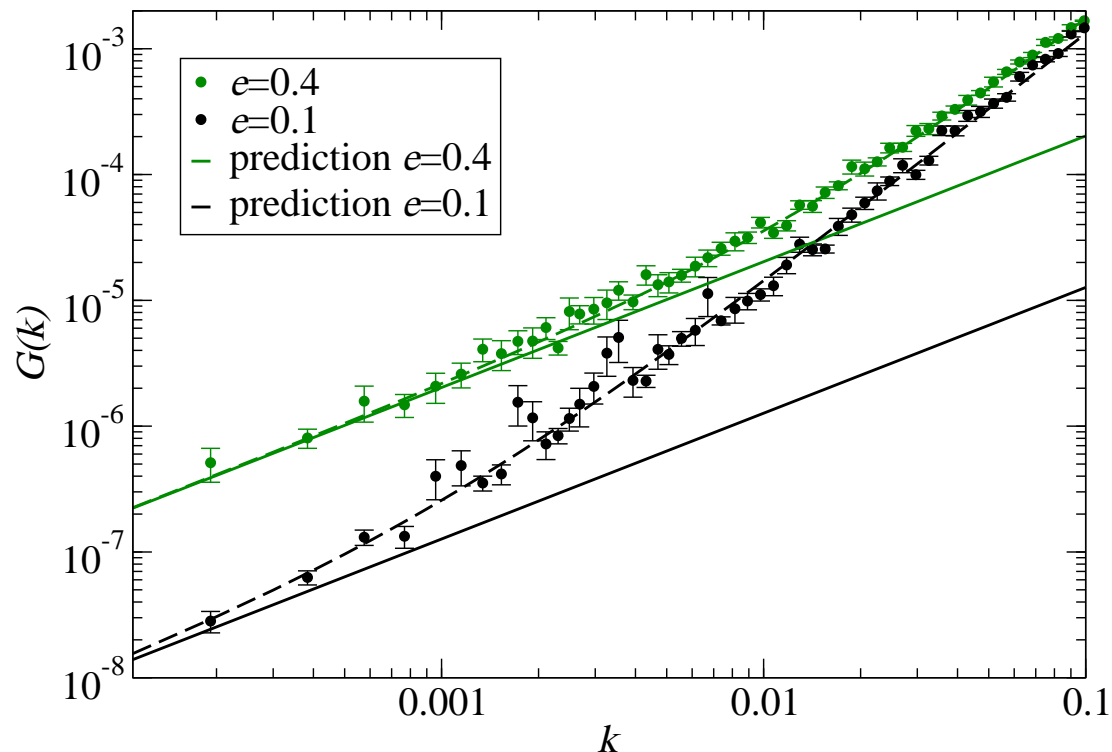
3D Simulation



- Flux trapping prediction:

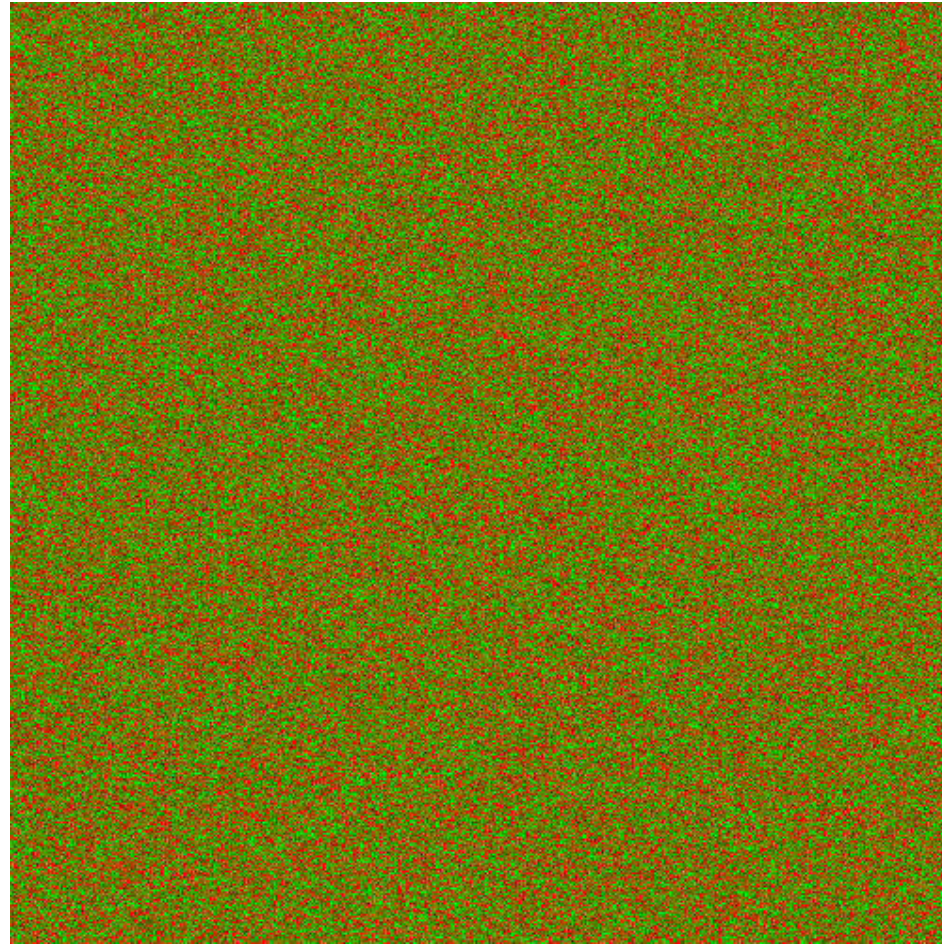
$$\langle \rho_i(\vec{k}) \rho_j(\vec{q}) \rangle = \left(\frac{e}{2\pi} \right)^2 \langle B_i(\vec{k}) B_j(\vec{q}) \rangle_{\text{ini}}$$

Zero-Temperature Phase Transition



- $128 \times 128 \times 32768$ lattice
- Classical field theory simulation:
Initial Gaussian fluctuations with the quantum vacuum two-point function

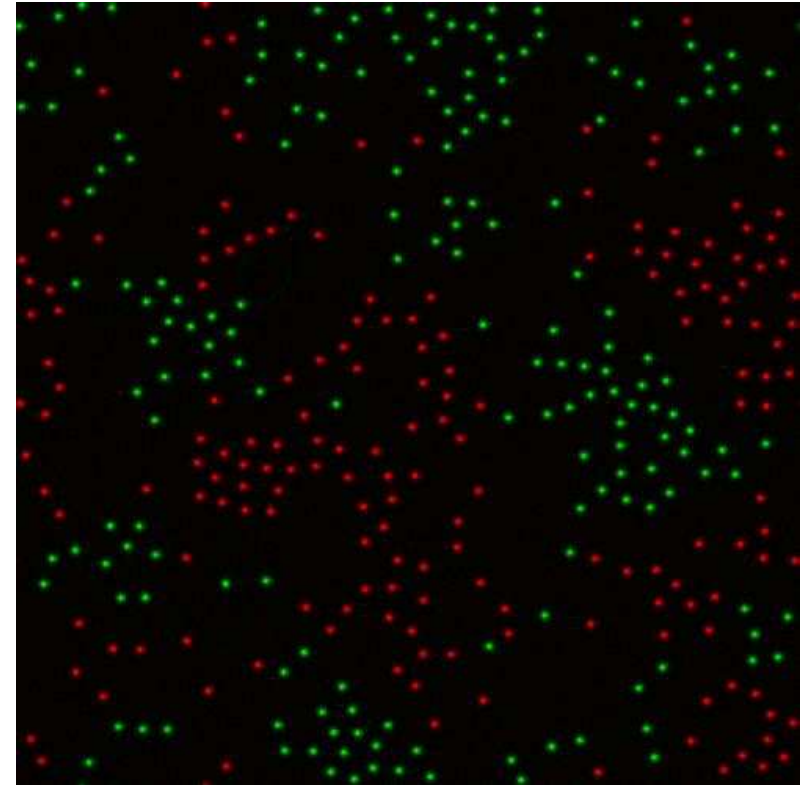
Short-Distance Effects



2D Simulation

Short-Distance Effects

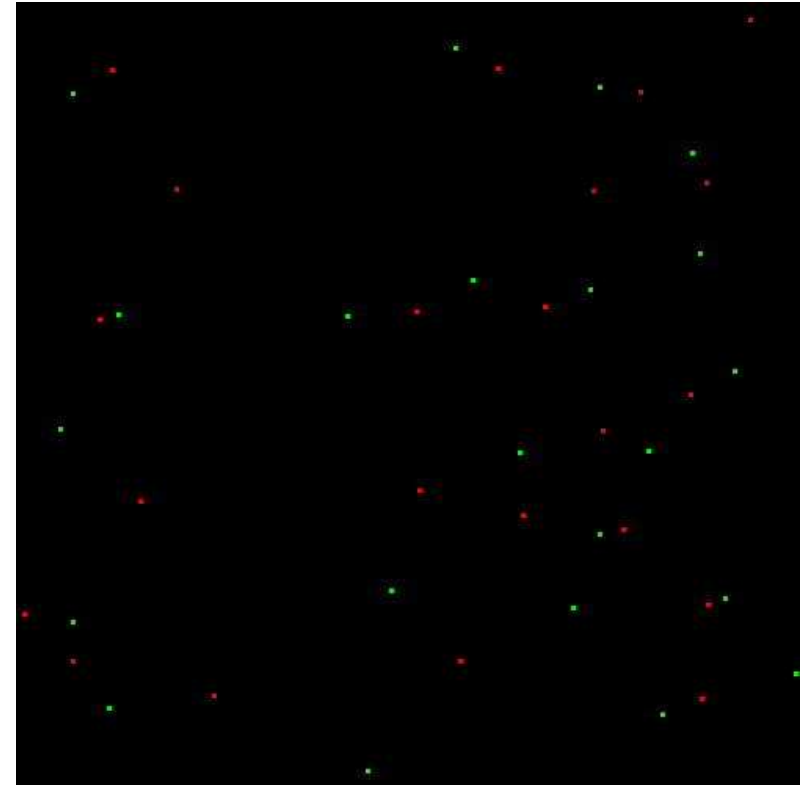
- Clusters of equal-sign vortices:
 - $N \approx eT^{1/2} \lambda_c^{2-D/2}$ per cluster
(Hindmarsh&AR, 2000)
 - Kibble mechanism: No clusters



gauge

Short-Distance Effects

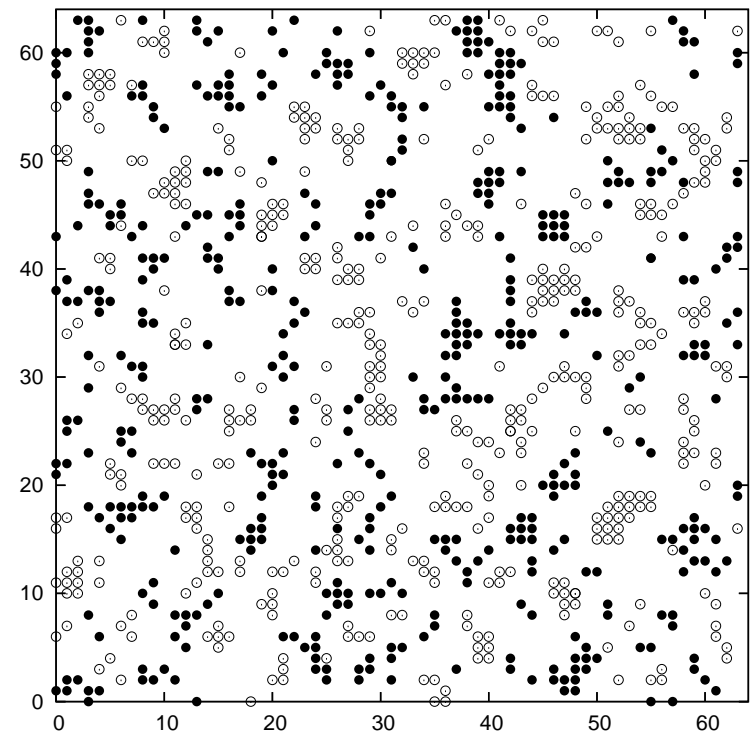
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global

Short-Distance Effects

- Clusters of equal-sign vortices:
 - $N \approx eT^{1/2} \lambda_c^{2-D/2}$ per cluster
(Hindmarsh&AR, 2000)
 - Kibble mechanism: No clusters
- Also in 3D



(Blanco-Pillado et al 2007)

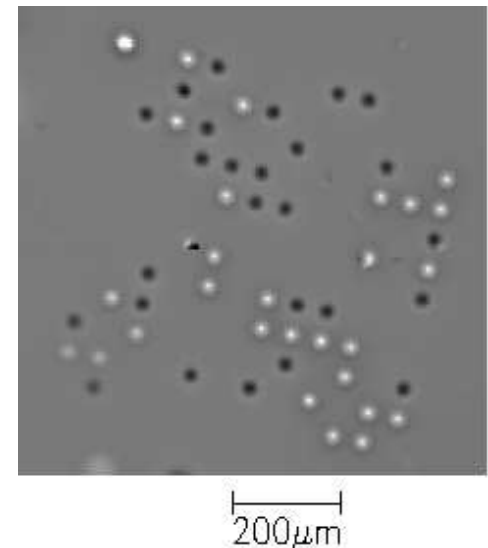
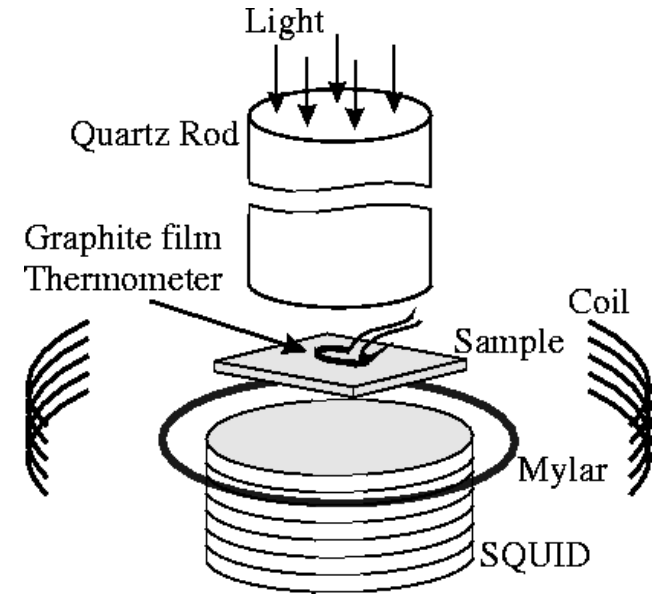
Cosmology

- Short-distance effects washed out quickly
 - Clustering probably not relevant
 - Thick strings ($N_w > 1$) can modify string evolution
- Long-range correlations potentially very important
 - Correlations on superhorizon scales:
Strings pile up as the horizon grows
 - Thermal initial state ruled out: $T_{\text{CMB}}/H_0 \sim 10^{14}$ infinite strings
 - Vacuum initial state: $M_{\text{Pl}}/g_*^{1/2} T_{\text{inf}} \sim 10$ infinite strings
 - with preferred direction
 - correlations from “nothing”
 - are classical arguments valid?



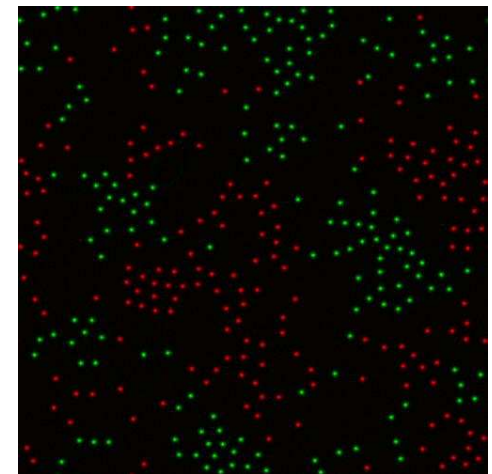
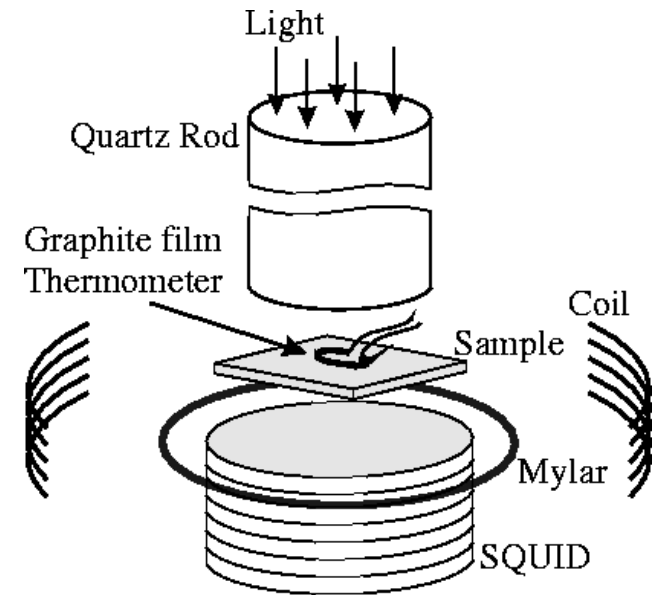
Superconductor Experiments

- Quench from $T > T_c$ to $T < T_c$
 - Fully quantum mechanical
 - Better test than any simulation
- Measurements of net flux
(Carmi et al, Monaco et al)
- Array of rings (Kirtley et al 2003)
 - Shows clustering
- Magneto-optical imaging:
 - Experiments being planned (Golubchik 2008)
 - Individual vortices: Correlations



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Conclusions

- Global symmetry: Kibble-Zurek mechanism
- Gauge symmetry: Kibble-Zurek + flux trapping
 - Thick strings from first order transitions
 - Dominant mechanism at long distances
 - Gauge field correlations survive in defect distribution:
Infinite-range correlations from “nothing”?
- Can be tested in superconductor experiments