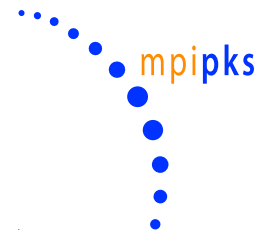


INTERACTION-INDUCED
EDGE EFFECTS
IN 1D OPEN-BOUNDARY LATTICES

Masud Haque

Max-Planck Institute for Physics of
Complex Systems (MPI-PKS)

Dresden, Germany

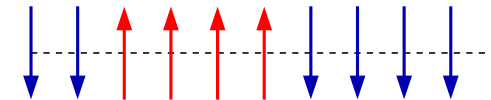
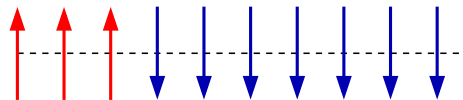
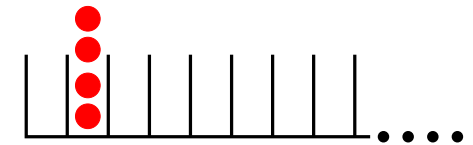
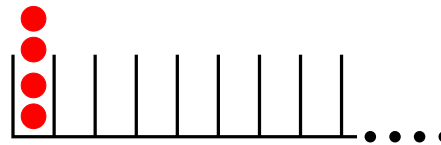


EDGE-LOCALIZATION IN 1D LATTICE MODELS

Bose-Hubbard chain

spinless fermion model

XXZ chain



PHYSICS:

Eigenstates far from ground state

Far-from-equilibrium dynamics

Intricate structures in spectrum (**FRACTAL**)

QUANTUM CONTROL:

Locking and **release** of magnetization/state

Designing a **quantum switch**

FOR DETAILS

R. A. Pinto, M. Haque, and S. Flach; [Phys. Rev. A **79**, 052118 \(2009\)](#).

Edge-localized states in quantum one-dimensional lattices.

M. Haque, [Phys. Rev. A **82**, 012108 \(2010\)](#).

Self-similar spectral structures and edge-locking hierarchy in open-boundary spin chains.

NICE PEOPLE THANK THEIR COLLABORATORS



Sergej Flach

MPI-PKS Dresden



Ricardo Pinto

MPI-PKS Dresden



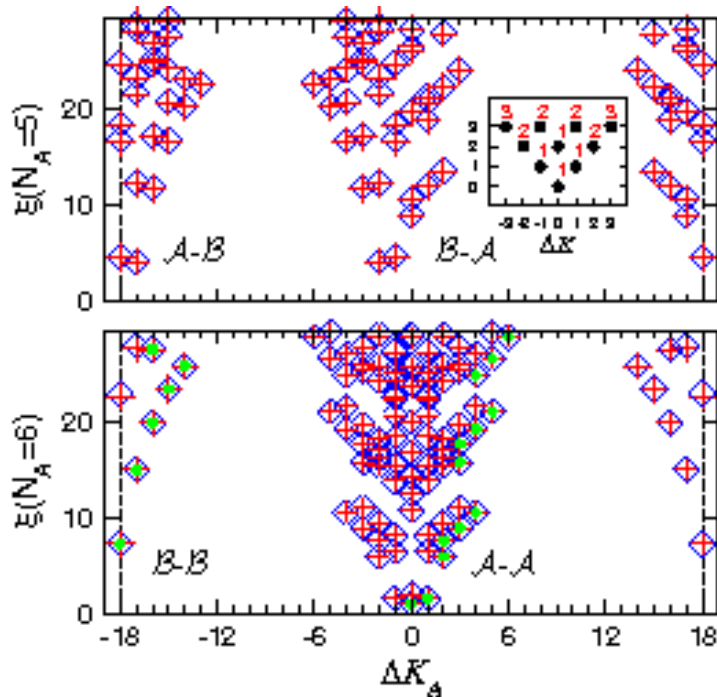
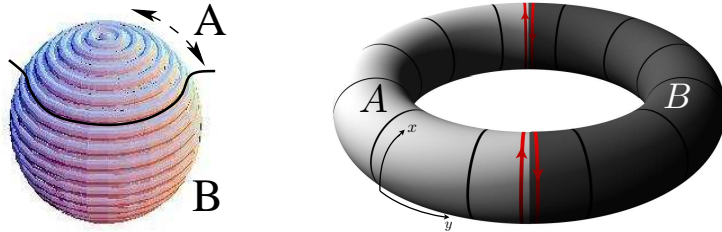
Riverside

BEFORE I START...

AN ADVERTISEMENT

(on unrelated topic)

ENTANGLEMENT & TOPOLOGICAL ORDER IN FQH STATES



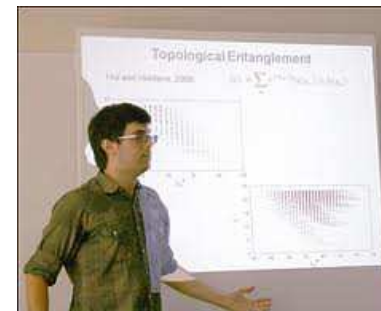
Haque, Zozulya, Schoutens; P.R.L. 2007
 Zozulya, Haque, Schoutens, Rezayi; P.R.B 2007
 Zozulya, Haque, Regnault; P.R.B 2009
 Läuchli, Bergholtz, Suorsa, Haque; P.R.L. 2010
 Läuchli, Bergholtz, Haque; N.J.P. 2010



Talk by Bergholtz next week.

Related series of work:
 Bernevig,
 Regnault, Haldane,

Talk by Papic last week

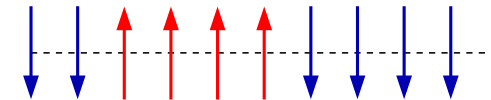
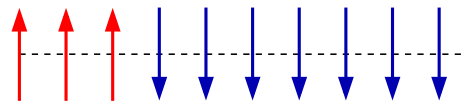
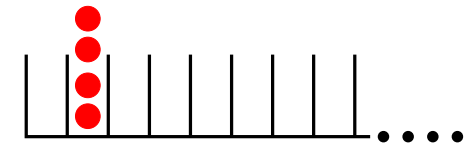
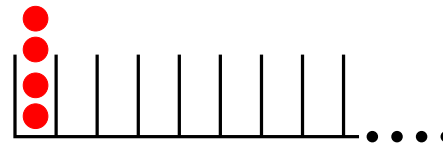


EDGE-LOCALIZATION IN 1D LATTICE MODELS

Bose-Hubbard chain

spinless fermion model

XXZ chain



PHYSICS:

Eigenstates far from ground state

Far-from-equilibrium dynamics

Intricate structures in spectrum (**FRACTAL**)

QUANTUM CONTROL:

Locking and release of magnetization/state

Designing a quantum switch

HAMILTONIANS & SMALL PARAMETERS

Hamiltonians:

$$H_{\text{Bose.Hubbard}} = -t \sum (a_j^\dagger a_{j+1} + a_{j+1}^\dagger a_j) + \frac{U}{2} \sum a_j^\dagger a_j^\dagger a_j a_j$$

$$H_{\text{sp.ferm.}} = -t \sum (c_j^\dagger c_{j+1} + c_{j+1}^\dagger c_j) + V \sum c_j^\dagger c_{j+1}^\dagger c_{j+1} c_j$$

$$H_{\text{XXZ}} = J_x \sum \left[S_j^x S_{j+1}^x + S_j^y S_{j+1}^y + \Delta S_j^z S_{j+1}^z \right]$$

Small Parameters:

$$t/U$$

$$t/V$$

$$1/\Delta$$

I take these Hamiltonians **seriously**

(not only low-energy sector)

START WITH SOME GUESSING GAMES

1D Bose-Hubbard model in an OPEN chain (has edges)

$$\hat{H} = -t \sum_{j=1}^{L-1} (a_j^\dagger a_{j+1} + a_{j+1}^\dagger a_j) + \frac{U}{2} \sum_{j=1}^L a_j^\dagger a_j^\dagger a_j a_j$$

I'm interested in large U/t .

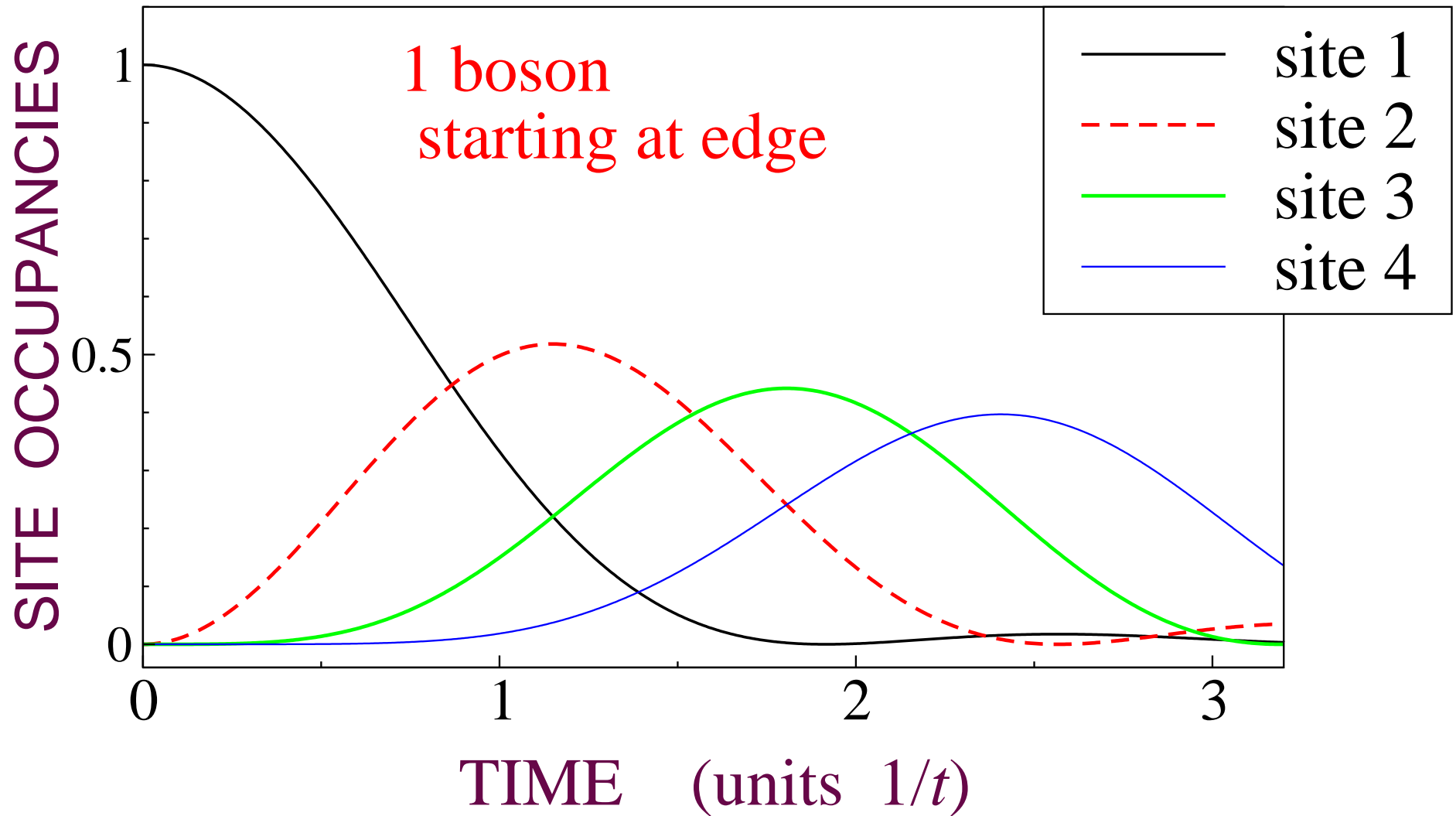


1 0 0 0 0 0 0 0

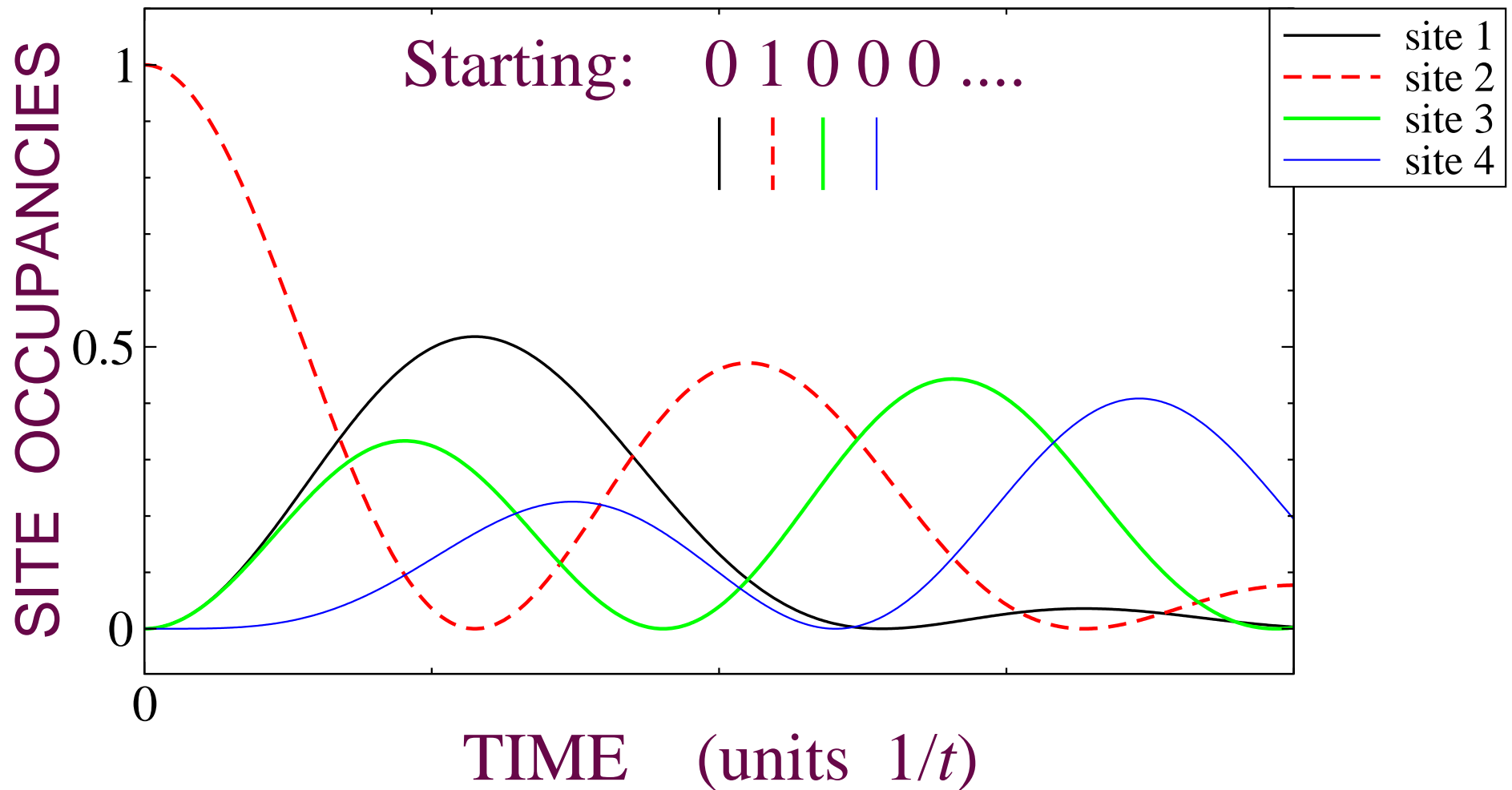
How does this evolve?

At timescales $\sim \hbar/t$

ONE BOSON STARTING AT SITE 1



1 BOSON STARTING AT SITE 2 (NEXT-TO-EDGE)



NEXT: TWO BOSONS



2 0 0 0 0 0 0

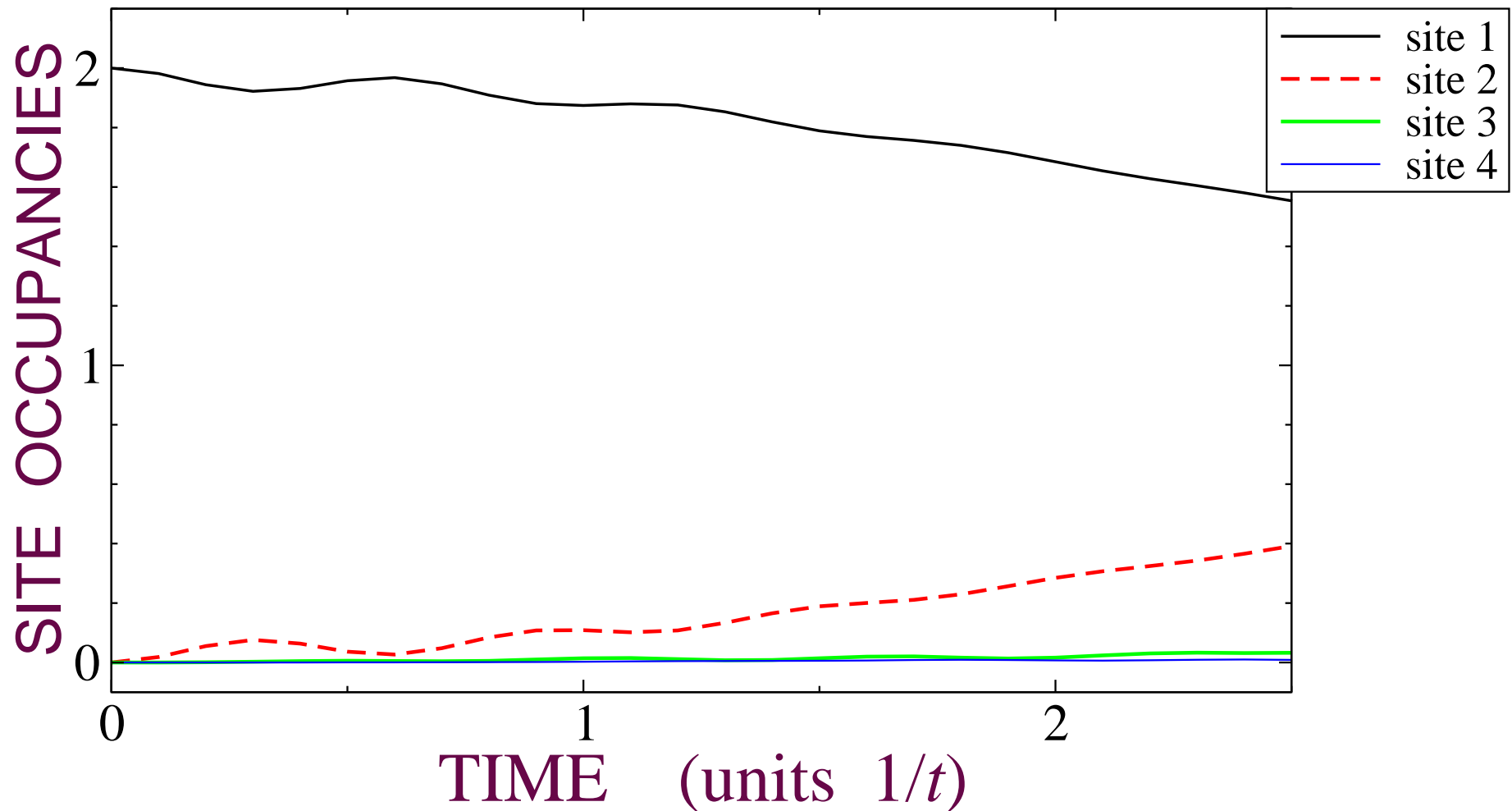
How does this evolve?

At timescales $\sim 1/t \sim 1$

At timescales $\sim 1/(t^2/U) \sim U$

TWO BOSONS AT EDGE: TIMESCALES $\sim 1/t \sim 1$

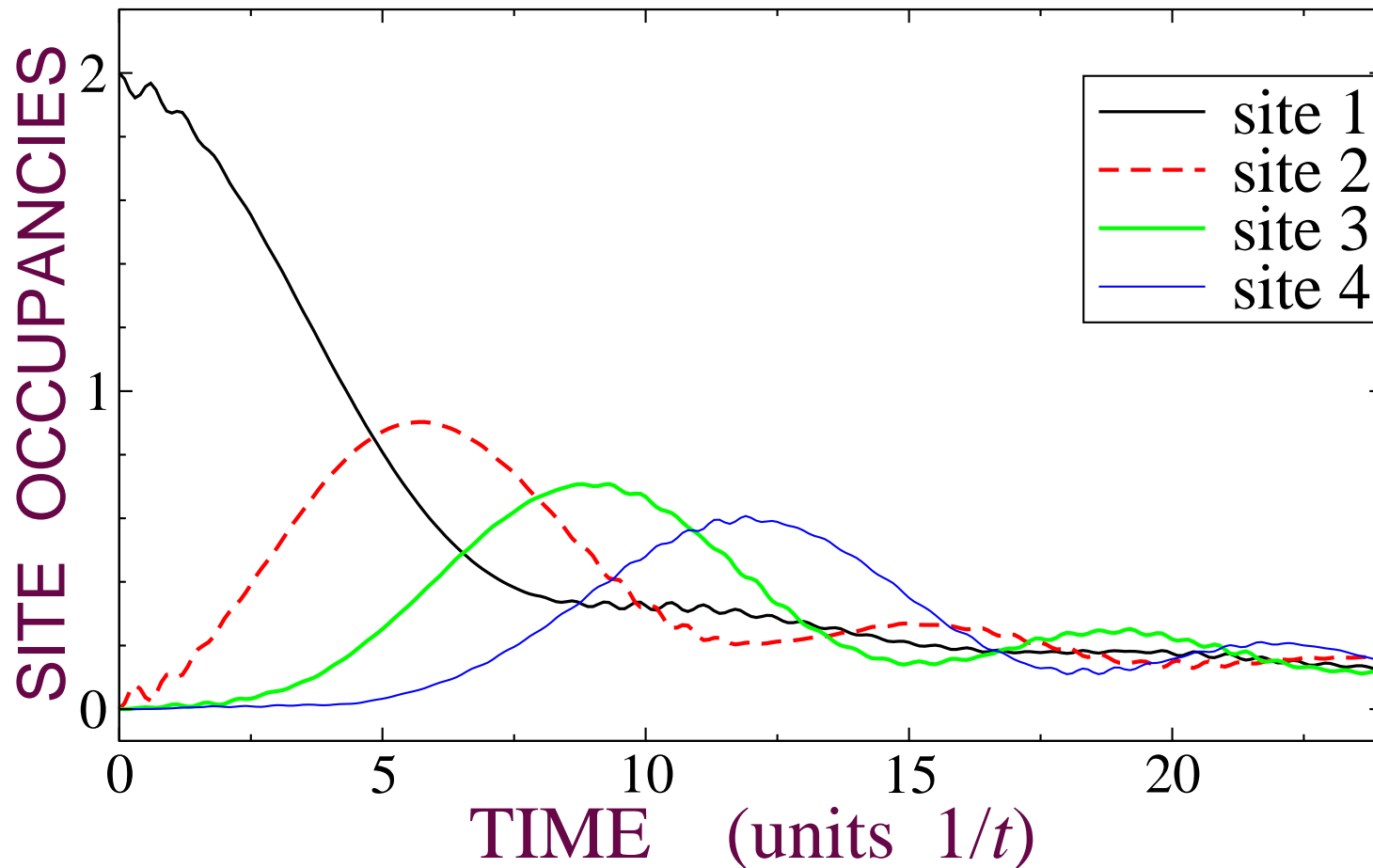
$U = 10$ Starting: 20000 ...



2 BOSONS AT EDGE: TIMESCALES $\sim 1/(t^2/U) \sim U$

$U = \pm 10$

Starting: 2 0 0 0 ...



LARGE U ENCOURAGES CORRELATED PAIR MOTION

Single particle hopping timescale $\sim 1/t \sim 1$

Pair hopping time scale $\sim 1/\left(\frac{t^2}{U}\right) \sim U$

“Repulsively bound pairs”

Triplet hopping time scale $\sim 1/\left(\frac{t^3}{U^2}\right) \sim U^2$

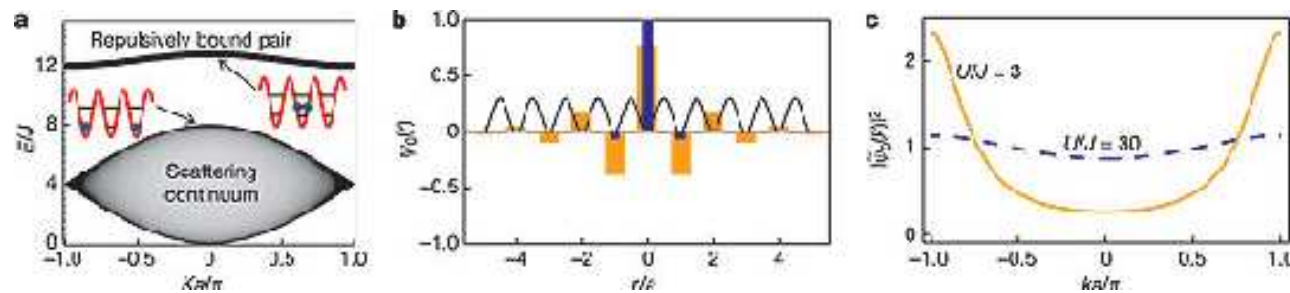
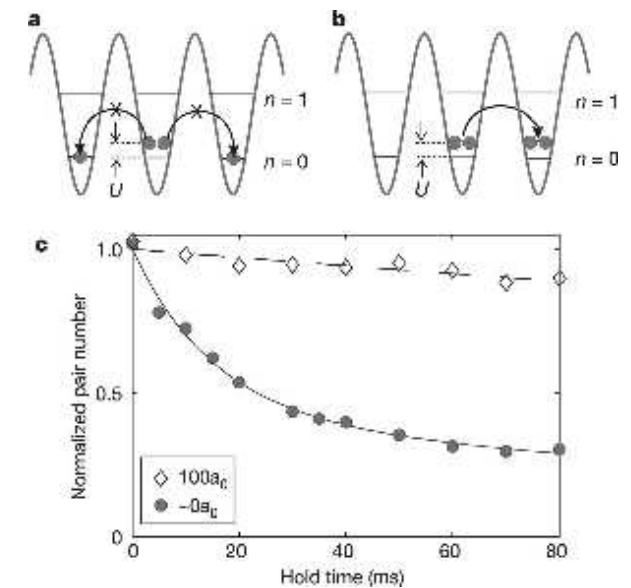
REPULSIVELY BOUND PAIRS

Vol 441 | 15 June 2006 | doi:10.1038/nature04918 nature

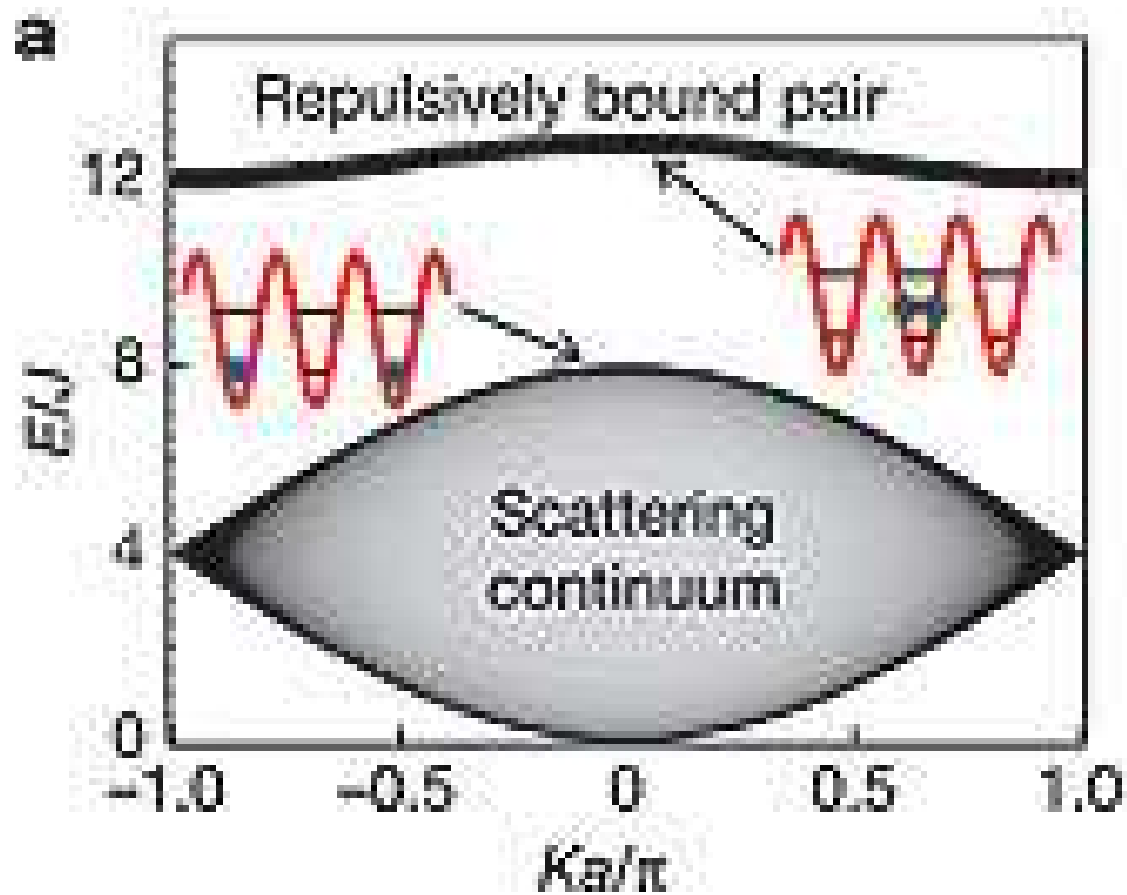
LETTERS

Repulsively bound atom pairs in an optical lattice

K. Winkler¹, G. Thalhammer¹, F. Lang¹, R. Grimm^{1,3}, J. Hecker Denschlag¹, A. J. Daley^{2,3}, A. Kantian^{2,3}, H. P. Büchler^{2,3} & P. Zoller^{2,3}



“BANDS” IN ENERGY SPECTRUM, 2 BOSONS

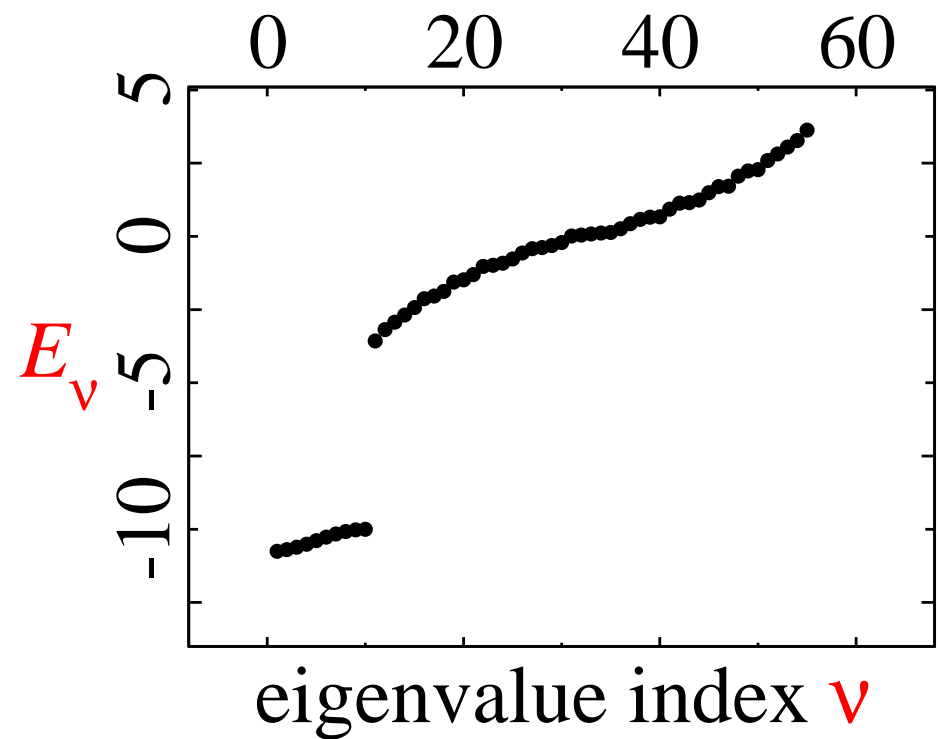
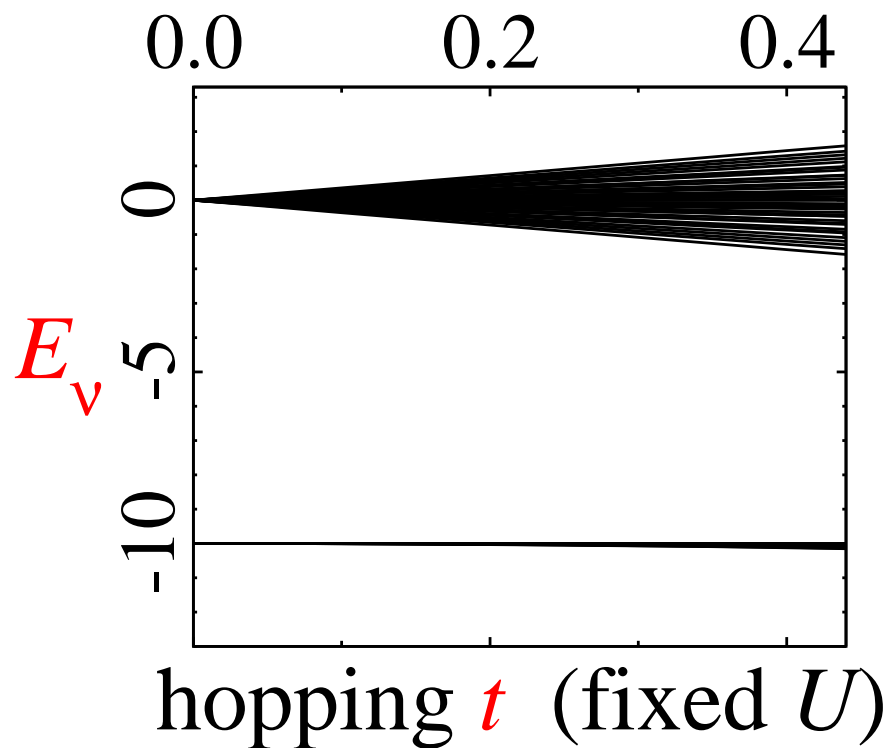


Pairs cannot break
without losing energy,
 \Rightarrow without energy
relaxation mechanism.

(translation-invariant picture)

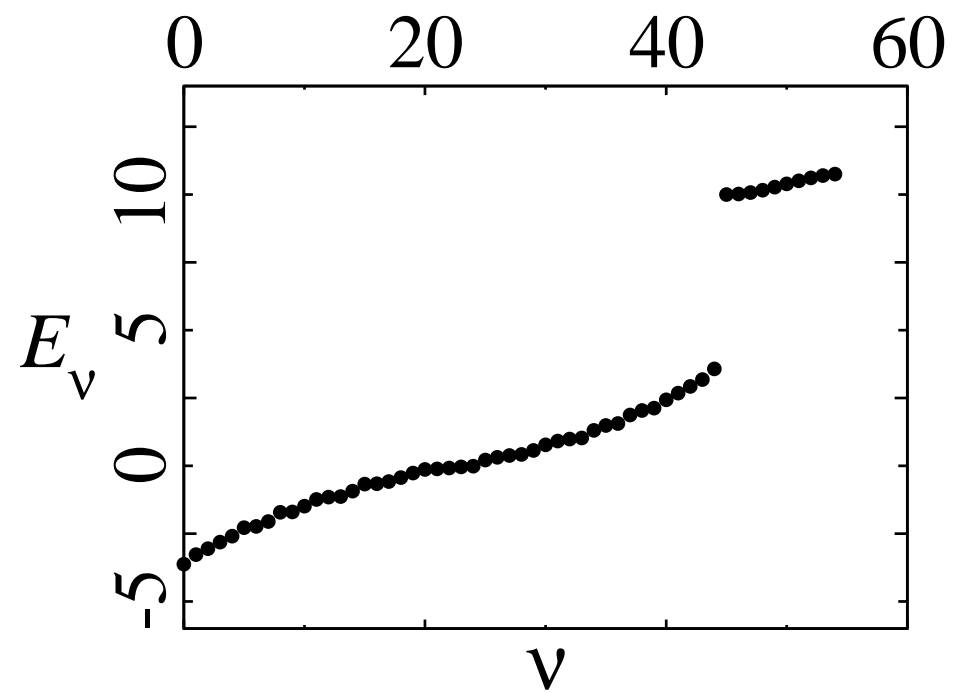
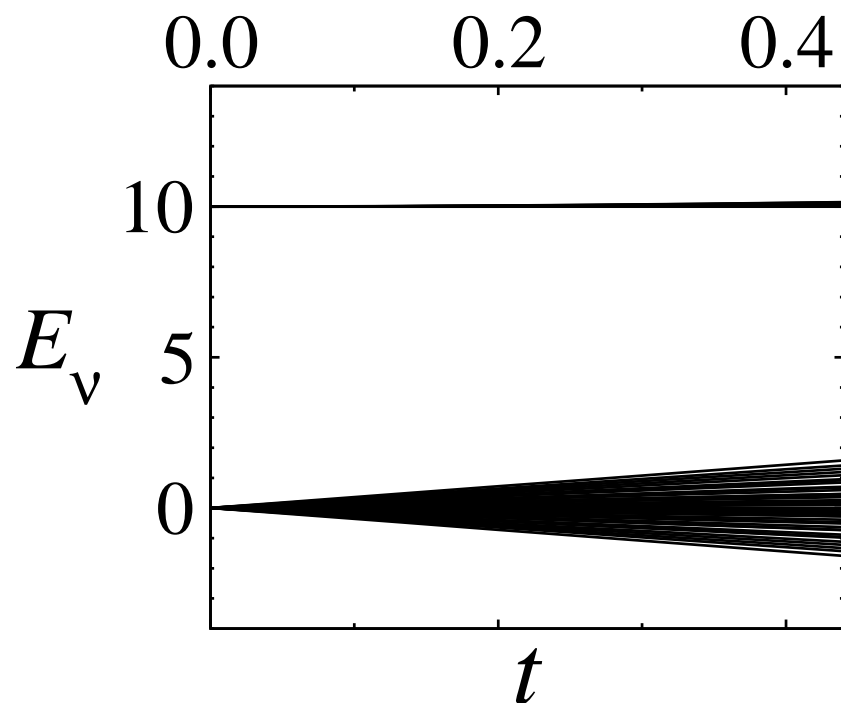
2-BOSON SPECTRUM, BANDS

2 Bosons in 10-site open chain. Negative U !! $U = -10$



2-BOSON SPECTRUM, BANDS, POSITIVE U

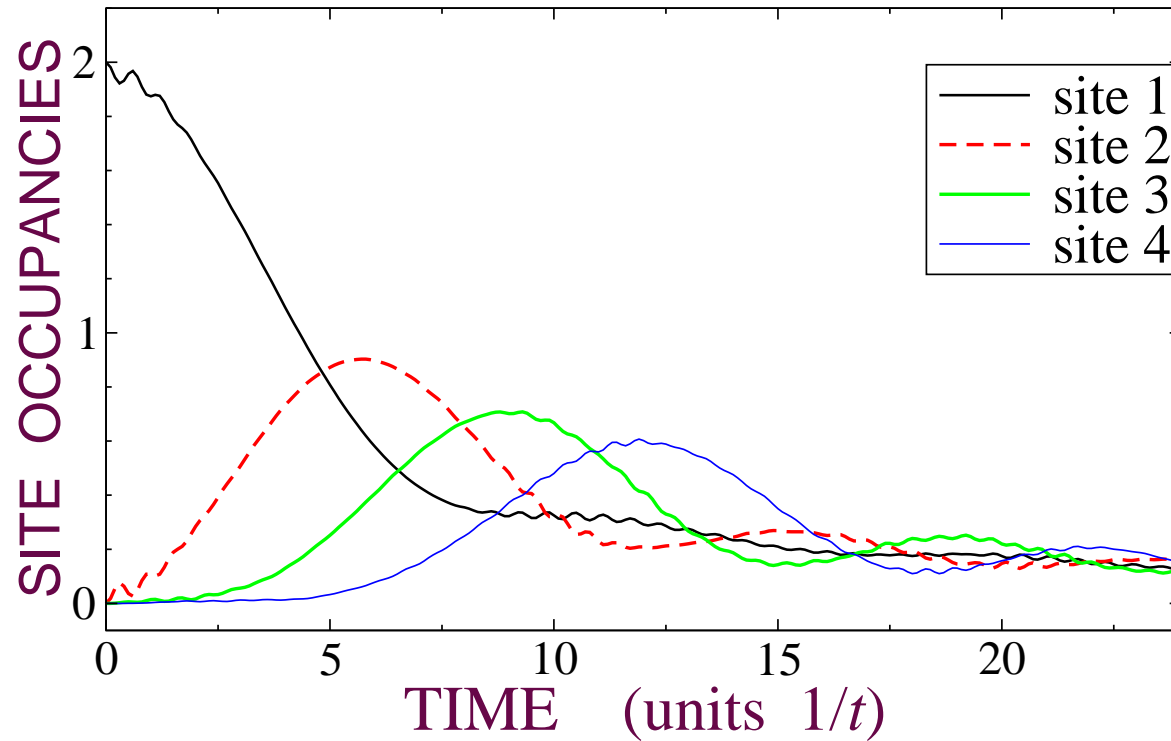
2 Bosons in 10-site open chain. $U = +10$



TWO BOSONS

$U = \pm 10$

Starting: 20000 ...



Long time-scale \rightarrow
hopping mostly within
bound-pair band.

High-frequency
oscillations \rightarrow
inter-band processes.

LET'S MOVE ON: THREE BOSONS



3 0 0 0 0 0 0

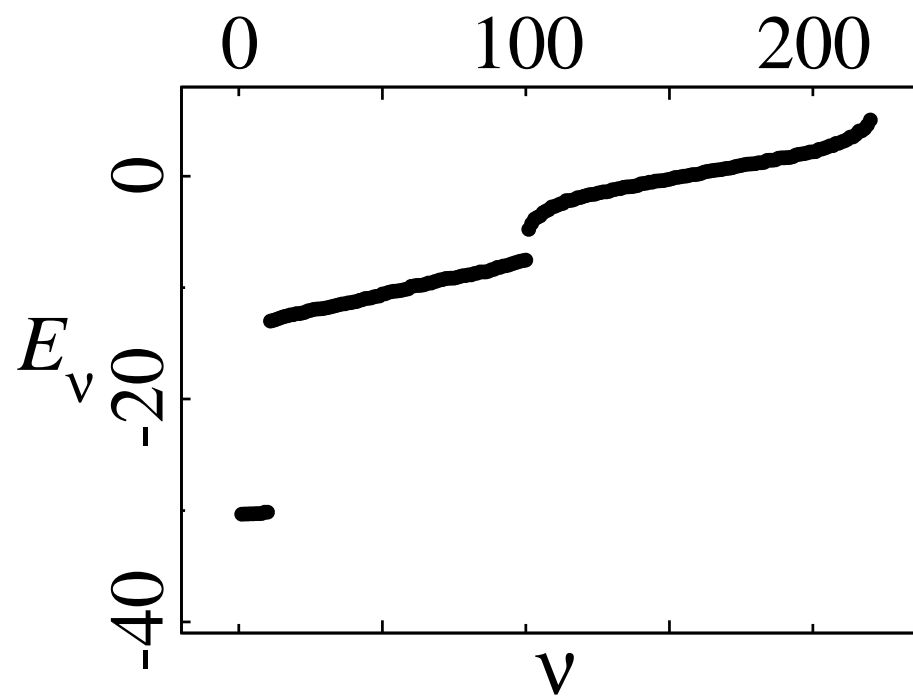
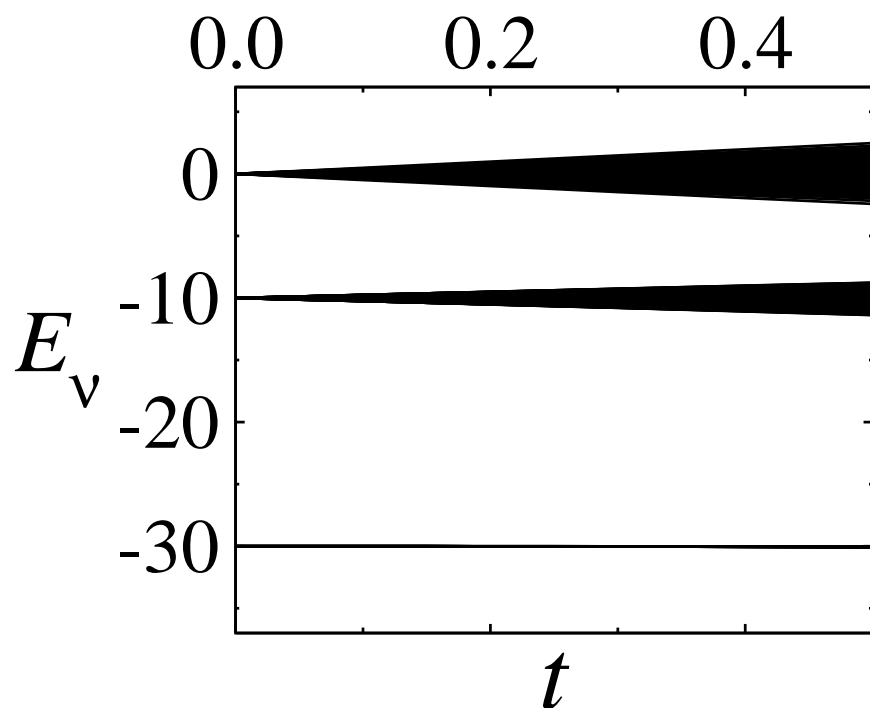
How does this evolve?

At timescales $\sim 1/t$

At timescales
 $\sim 1/(t^3/U^2) \sim U^2/t^3$

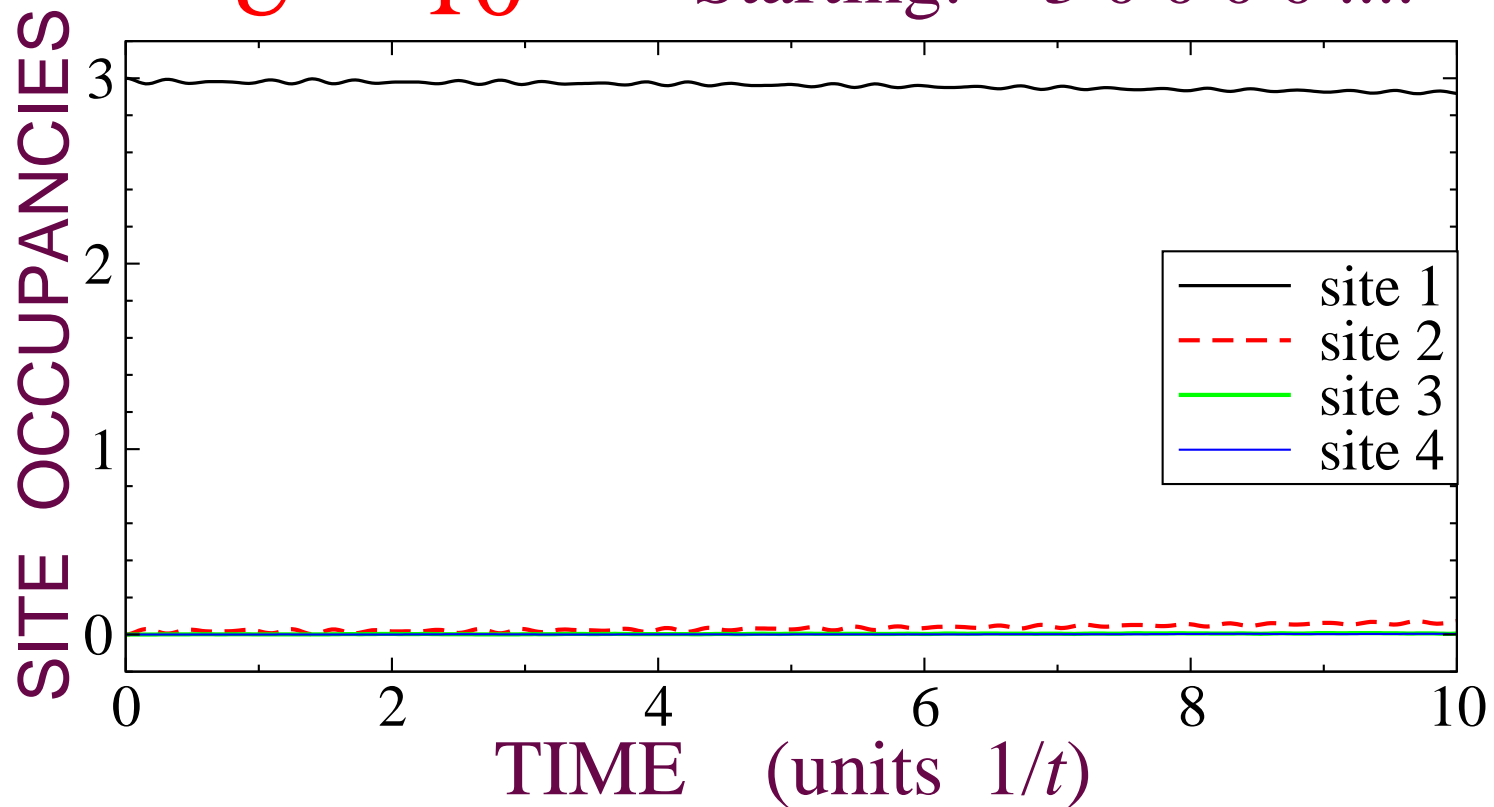
3-BOSON SPECTRUM, BANDS

3 Bosons in 10-site open chain. Negative U ; $U = -10$



THREE BOSONS AT EDGE: TIMESCALES $\sim 1/t$

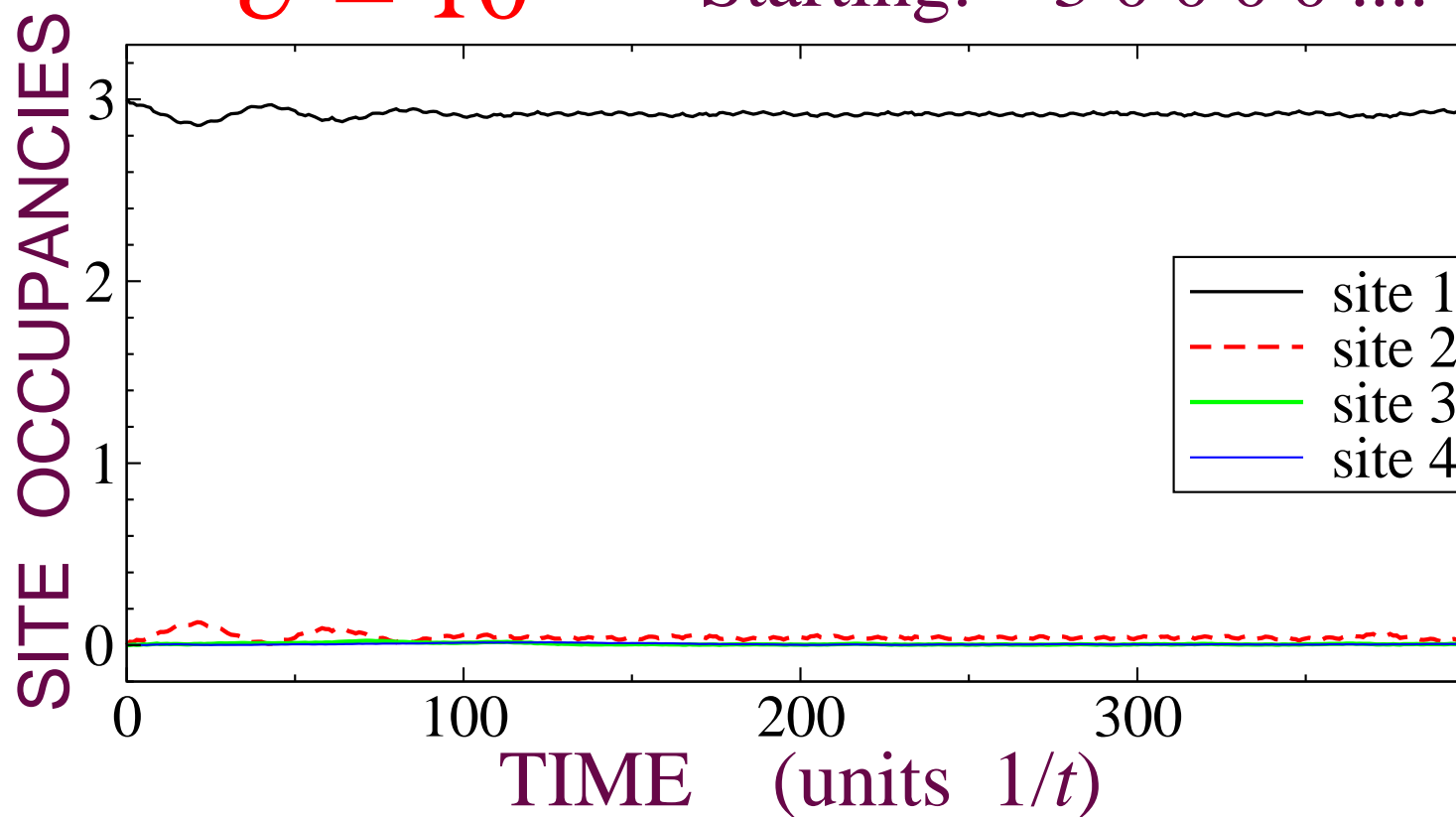
$U = 10$ Starting: 3 0 0 0 0 ...



No big surprise.

THREE BOSONS AT EDGE: TIMESCALES $\sim U^2$

$U = 10$ Starting: 3 0 0 0 ...

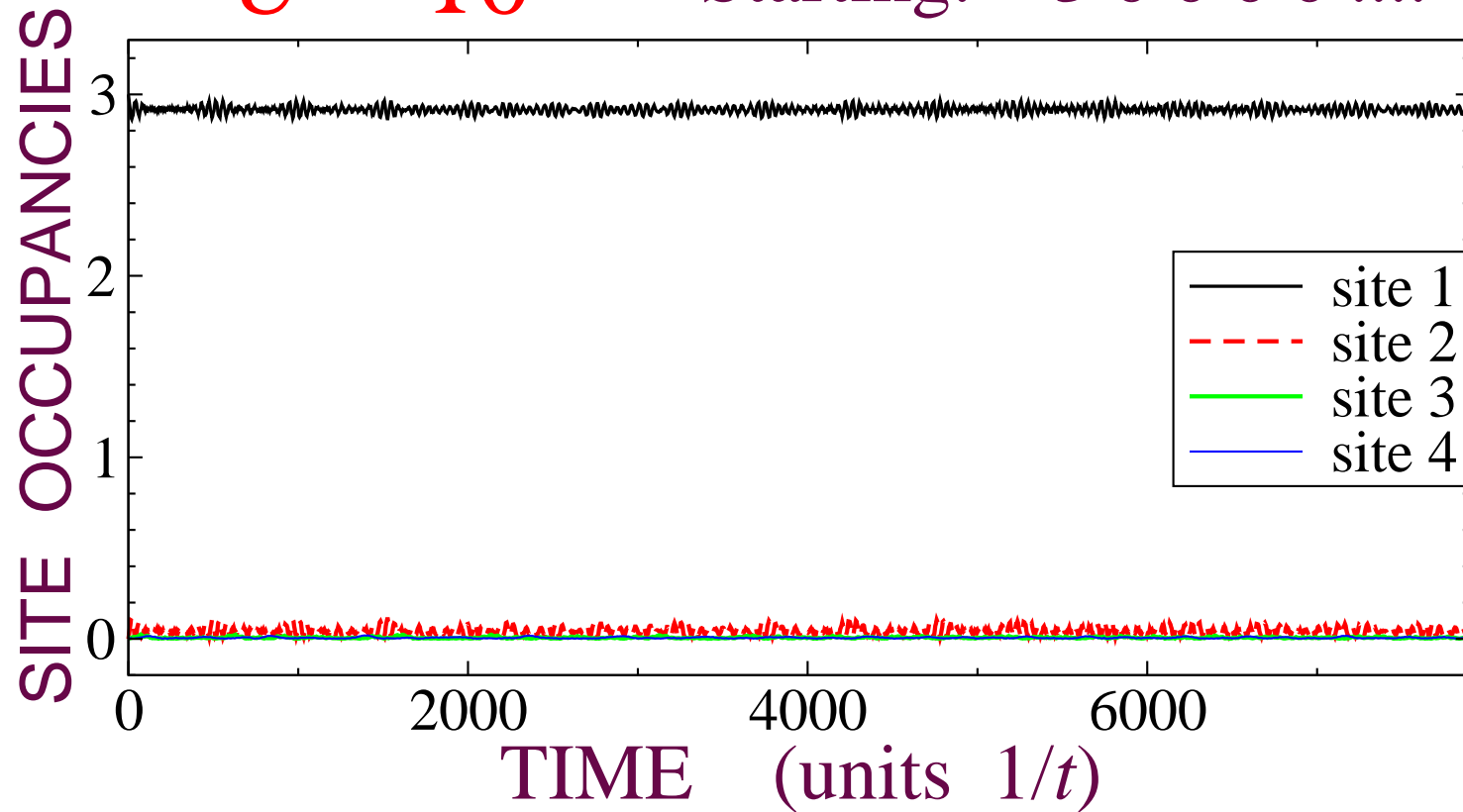


CONFUSED: TRY TIMESCALES $\gg \sim U^2$

TRYING TIMESCALES $\gg \sim U^2$

$U = 10$

Starting: 3 0 0 0 0 ...



? ? ? ? ? ? ? ? ?

You should be surprised

WE'VE FOUND A **STABLE** STATE

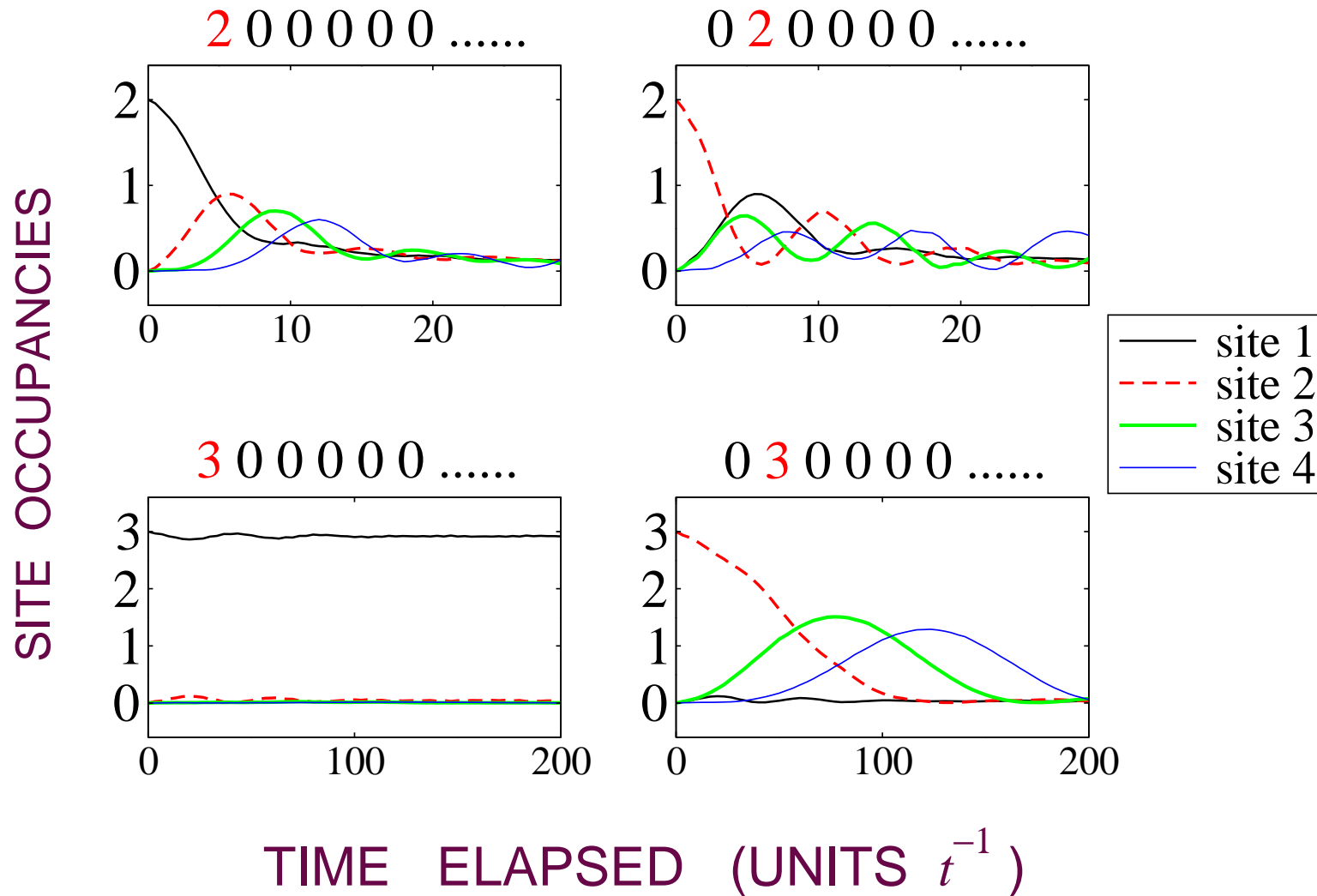


3 0 0 0 0 0

For $n \geq 3$ bosons, **edge states** are stable.

Stable should mean “close” to an eigenstate?

TIME EVOLUTION SUMMARY



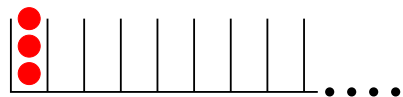
EDGE-LOCALIZED CONFIGURATIONS



NOT
STABLE



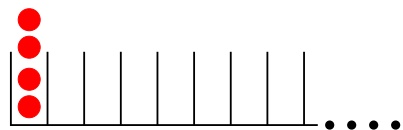
NOT
STABLE



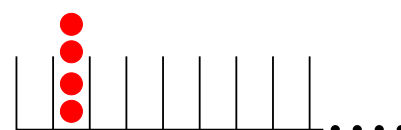
STABLE



NOT
STABLE



STABLE

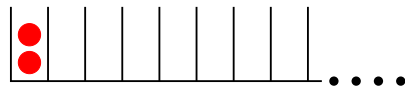


NOT
STABLE

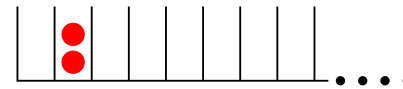
..... beginning of a 'hierarchy'

'Stable' means almost an eigenstate at large U/t .

HIERARCHY OF EDGE-LOCKED STATES



NOT
STABLE



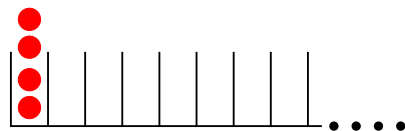
NOT
STABLE



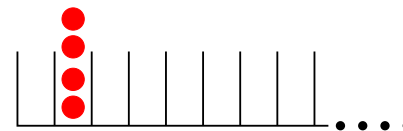
STABLE



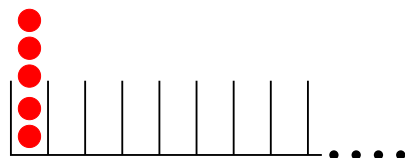
NOT
STABLE



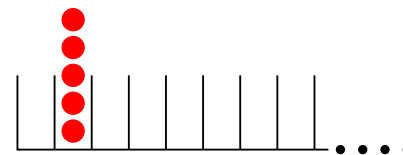
STABLE



NOT
STABLE

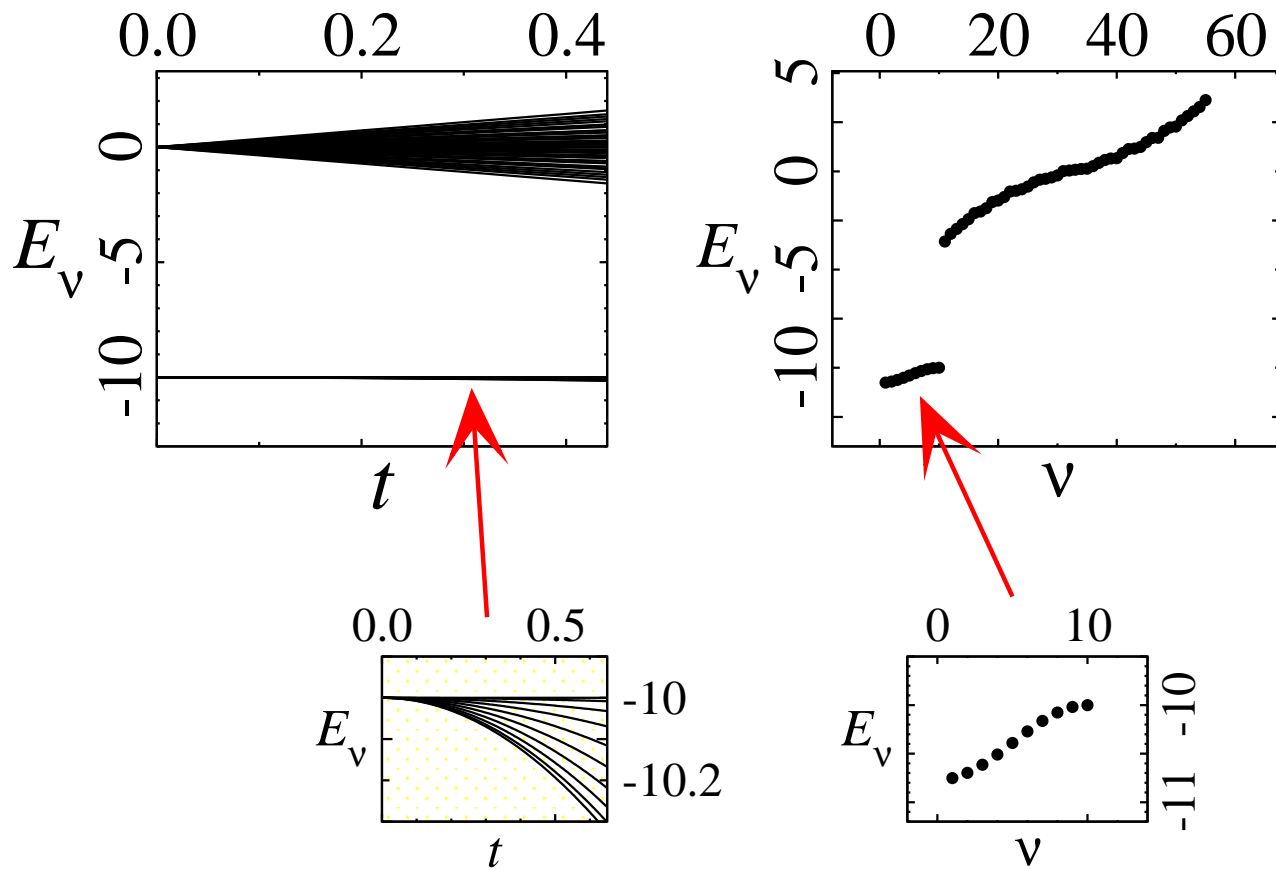


STABLE



STABLE

STRUCTURE OF 'BOUND' BAND: TWO BOSONS



Linear combinations of

$|20000\dots000\rangle$

$|02000\dots000\rangle$

$|00200\dots000\rangle$

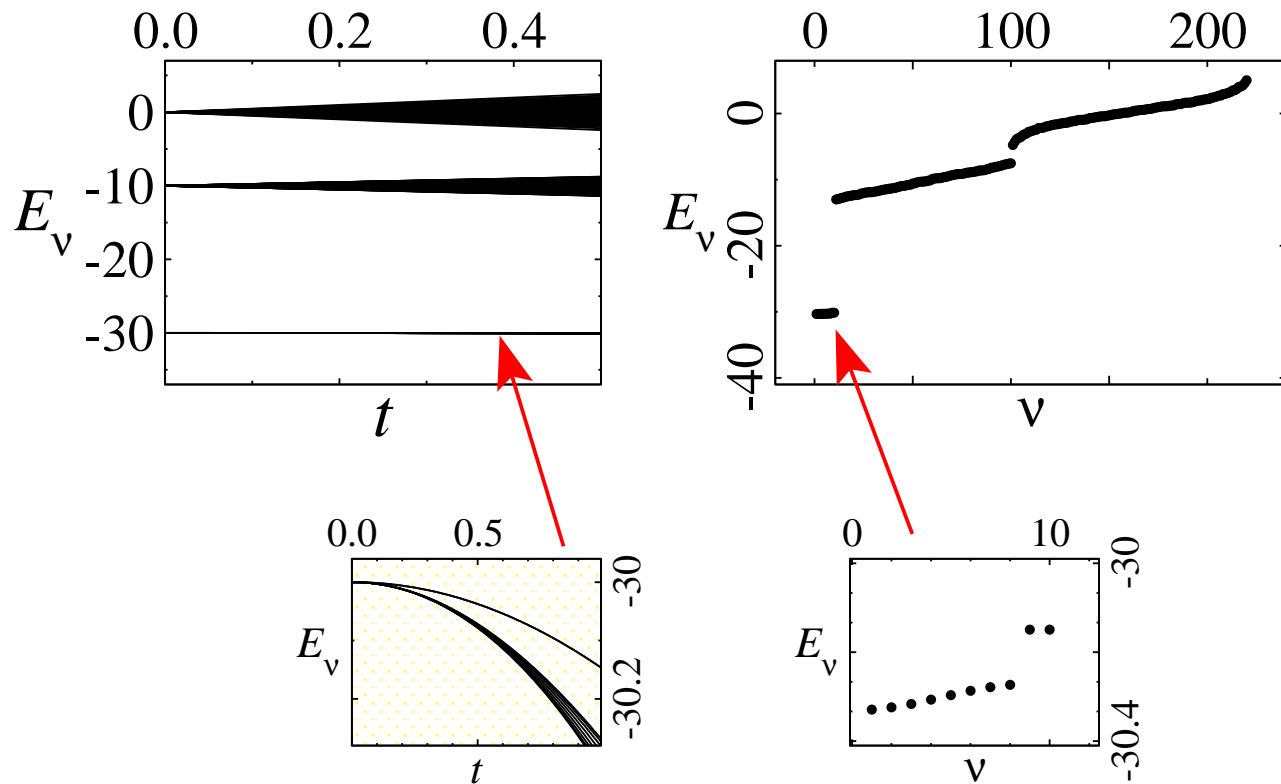
$|00020\dots000\rangle$

..

..

$|0000\dots002\rangle$

'BOUND' BAND: THREE BOSONS



Linear combinations of

$|03000\dots000\rangle$

$|00300\dots000\rangle$

$|00030\dots000\rangle$

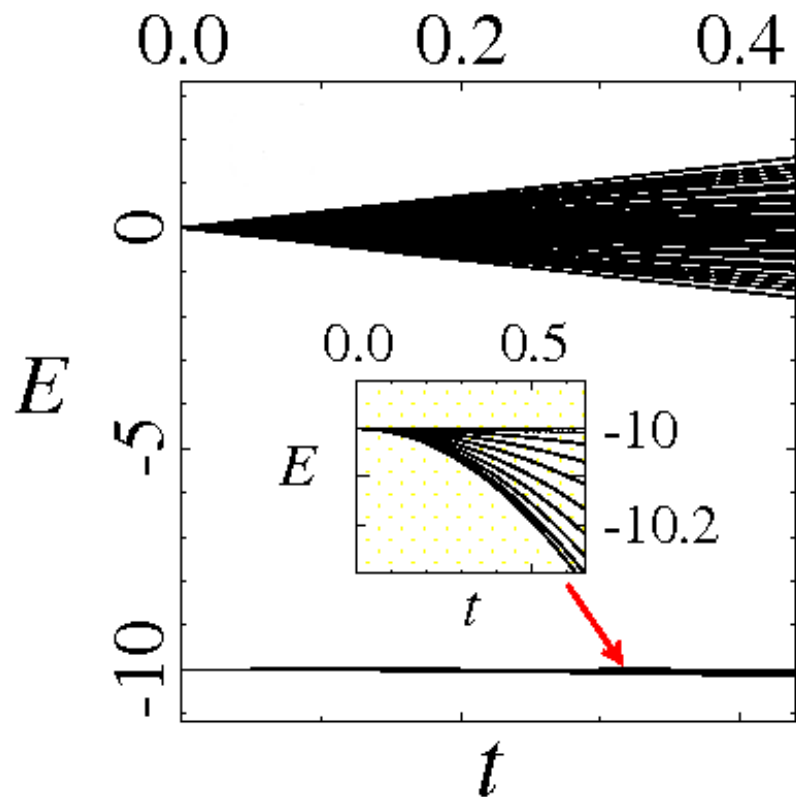
..

..

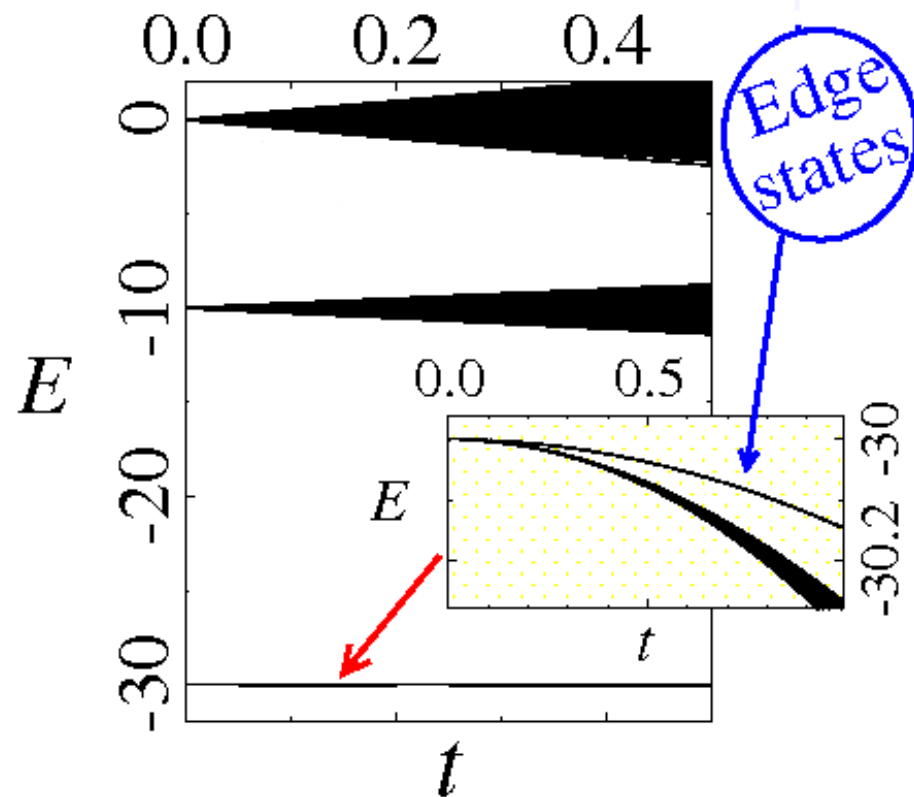
$|0000\dots030\rangle$

Separated out from the rest: $|30000\dots000\rangle$ and $|0000\dots003\rangle$.

2 bosons



3 bosons



$$|L\rangle = \begin{array}{c} \bullet \\ \bullet \\ \bullet \end{array} | \dots \rangle$$

$$|R\rangle = \dots | \begin{array}{c} \bullet \\ \bullet \\ \bullet \end{array} \rangle$$

Edge eigenstates: $|L\rangle + |R\rangle$ and $|L\rangle - |R\rangle$

TUNNEL TO OTHER EDGE?

$$|L\rangle = |3000\dots 00\rangle \quad \text{and} \quad |R\rangle = |00\dots 0003\rangle$$

Question: Why doesn't $|L\rangle$ tunnel to $|R\rangle$?

Answer: It will. After some astronomically long time.

$|L\rangle \leftrightarrow |R\rangle$ tunneling exponentially suppressed.

Splitting between $|L\rangle + |R\rangle$ and $|L\rangle - |R\rangle$ exponentially small.

SPECTRAL SEPARATION EXPLAINS
STABILITY OF EDGE STATES

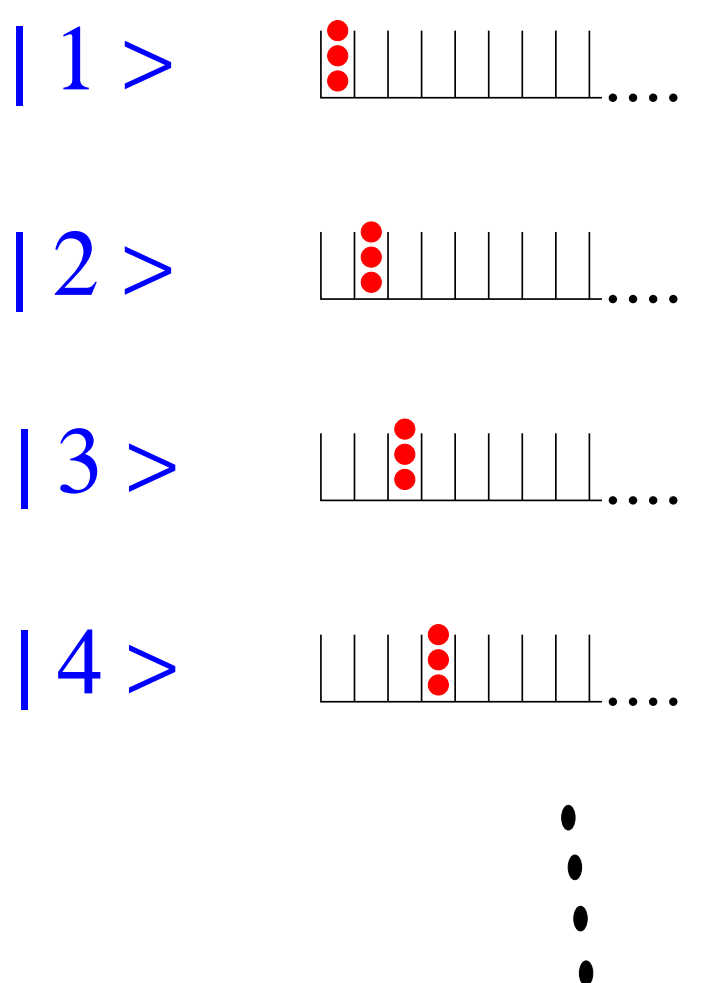
SPECTRAL SEPARATION EXPLAINS STABILITY OF EDGE STATES

Who ordered the spectral separations?

Degenerate perturbation theory.

Competition between energy shifts at $\mathcal{O}(t^2)$ and manifold mixing at $\mathcal{O}(t^n)$.

DEGENERATE PERTURBATION THEORY



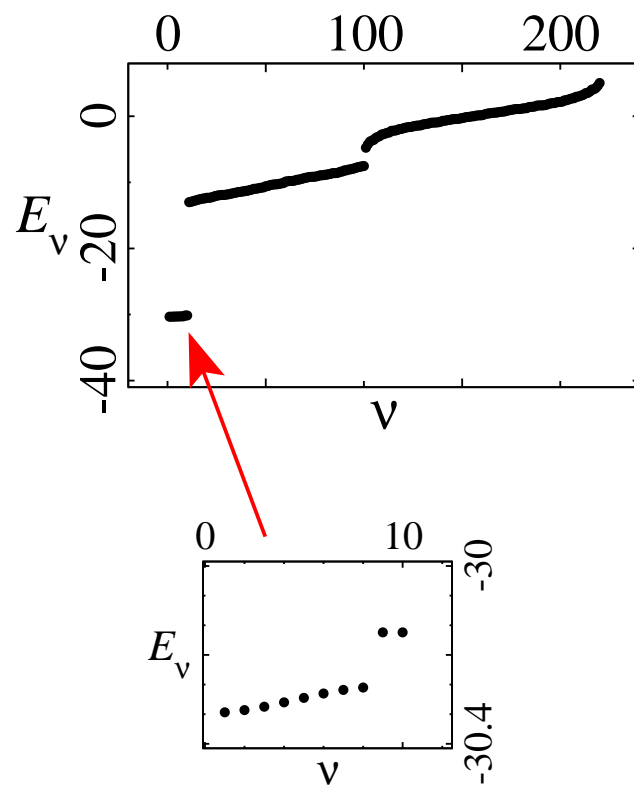
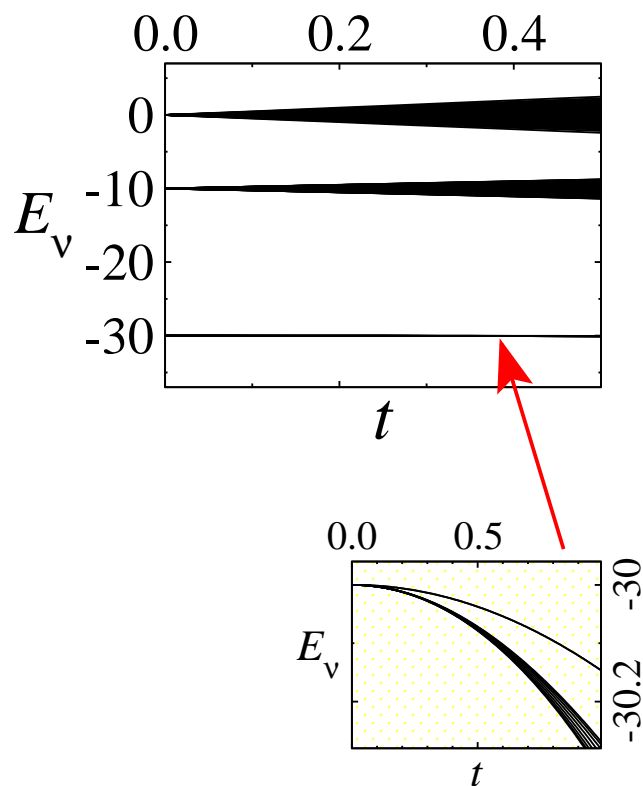
Degenerate manifold at $t/U = 0$.

States $|j\rangle$ and $|j + 1\rangle$ connect at $\mathcal{O}(t^n)$.

State $|1\rangle$ acquires different shift at $\mathcal{O}(t^2)$.

State $|2\rangle$ acquires different shift at $\mathcal{O}(t^4)$.

THREE BOSONS: $\mathcal{O}(t^2)$ VERSUS $\mathcal{O}(t^3)$



Linear combinations of

$|03000\dots000\rangle$

$|00300\dots000\rangle$

$|00030\dots000\rangle$

..

..

$|0000\dots0030\rangle$

Separated out from the rest: $|30000\dots000\rangle$ and $|0000\dots003\rangle$.

WHAT I'M MISSING....

There should be a
sum over histories
interpretation

3	0	0	0	0
0	3	0	0	0
0	0	3	0	0
2	1	0	0	0
2	0	1	0	0
1	2	0	0	0
1	0	2	0	0
1	1	1	0	0
1	1	0	1	0

SPINLESS FERMION MODEL: SIMILAR HIERARCHY

$$\hat{H} = -t \sum_{j=1}^{L-1} (c_j^\dagger c_{j+1} + c_{j+1}^\dagger c_j) + V \sum_{j=1}^{L-1} c_j^\dagger c_{j+1}^\dagger c_{j+1} c_j$$

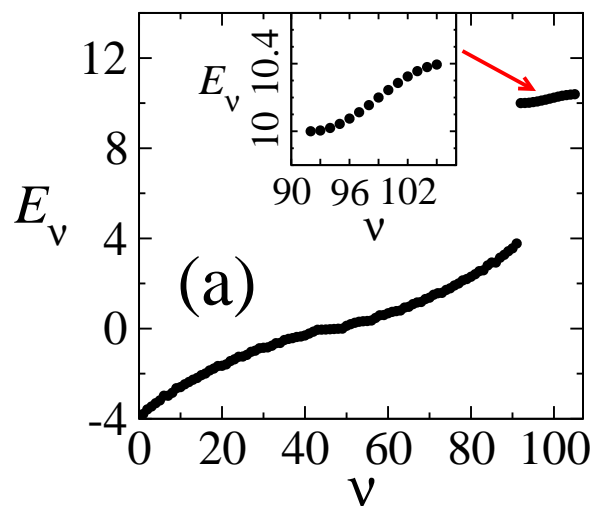
Sometimes called t - V model or Heisenberg-Ising model.

1	1	1	0	0	0	0	0	0	0	0
1	1	1	1	0	0	0	0	0	0	0
1	1	1	1	1	0	0	0	0	0	0
1	1	1	1	1	1	0	0	0	0	0
0	1	1	1	1	1	0	0	0	0	0
0	1	1	1	1	1	1	0	0	0	0

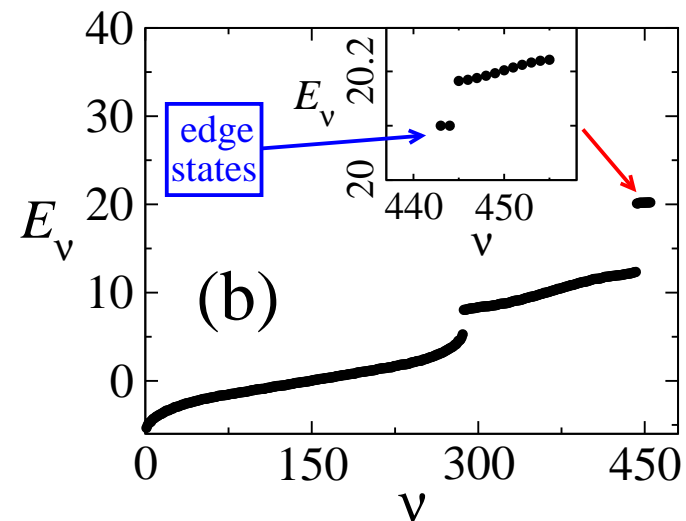
1	0	0	0	0	0	0	0	0	0	0
1	1	0	0	0	0	0	0	0	0	0
0	1	0	0	0	0	0	0	0	0	0
0	1	1	0	0	0	0	0	0	0	0
0	1	1	1	0	0	0	0	0	0	0
0	1	1	1	1	0	0	0	0	0	0

SPINLESS
FERMIONS:
SPECTRUM

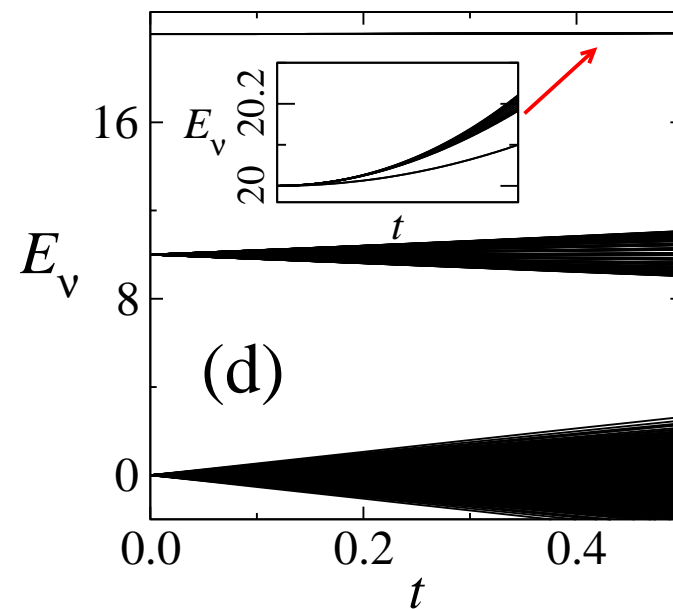
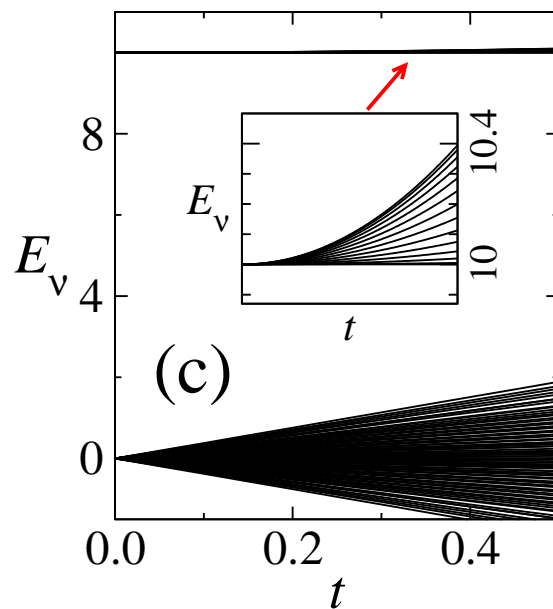
2 fermions



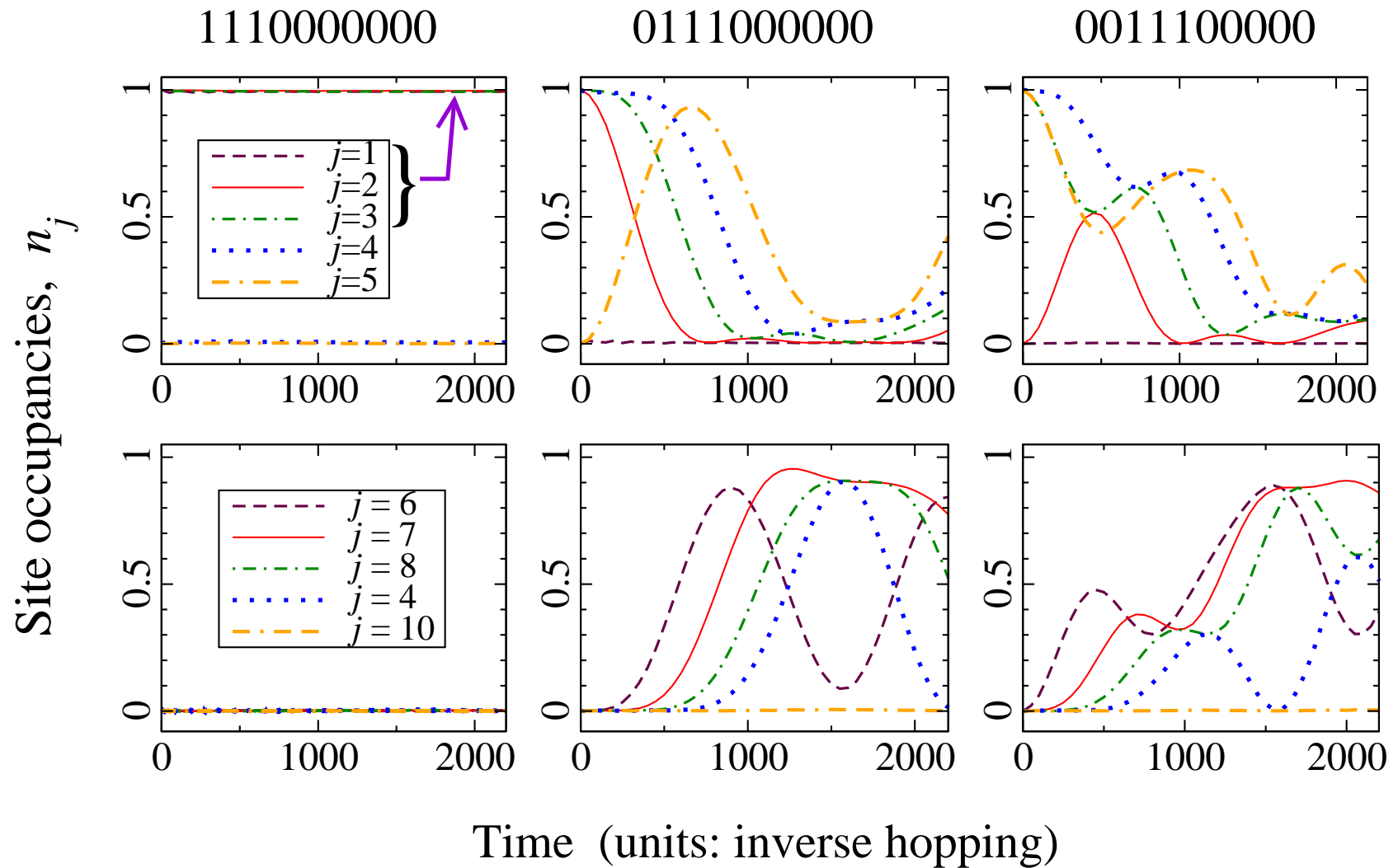
3 fermions



Two, three
fermions
in 15 sites



SPINLESS FERMIONS: DYNAMICS



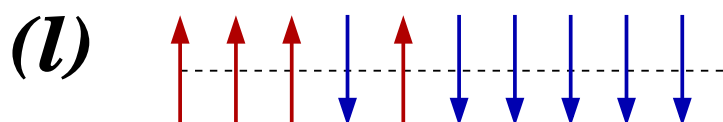
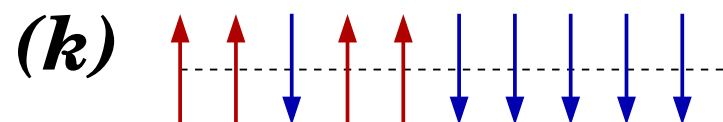
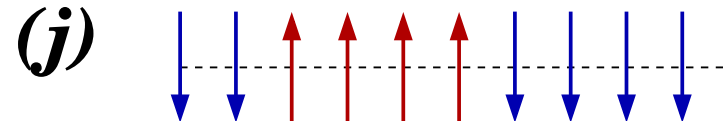
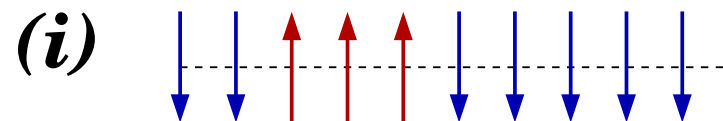
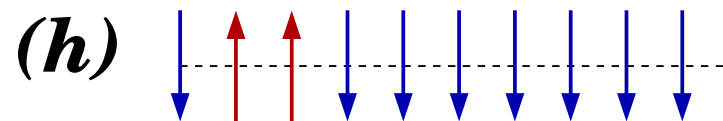
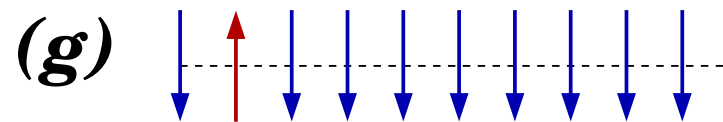
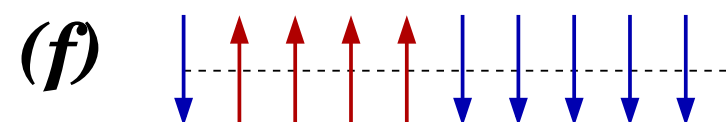
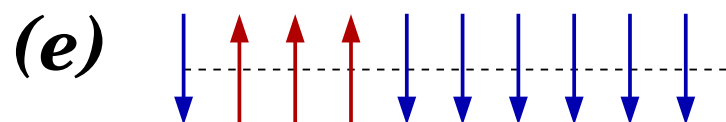
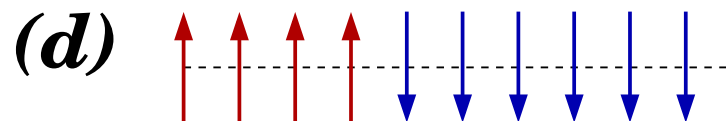
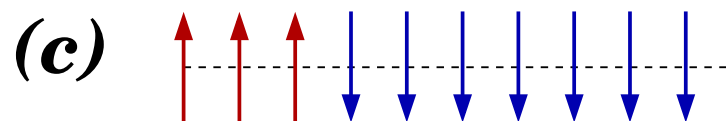
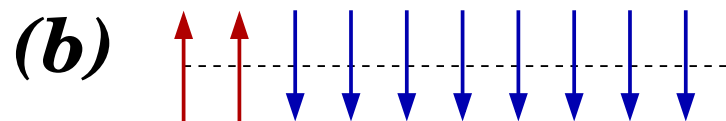
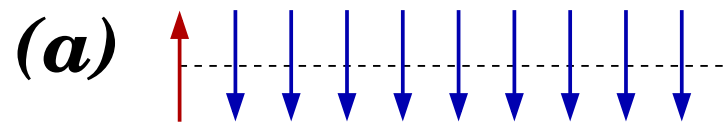
ANISOTROPIC HEISENBERG (XXZ) CHAIN

$$H = J_x \sum_{j=1}^{L-1} \left[S_j^x S_{j+1}^x + S_j^y S_{j+1}^y + \Delta S_j^z S_{j+1}^z \right]$$

Edge-locking hierarchy \rightarrow surprisingly different from t - V model.

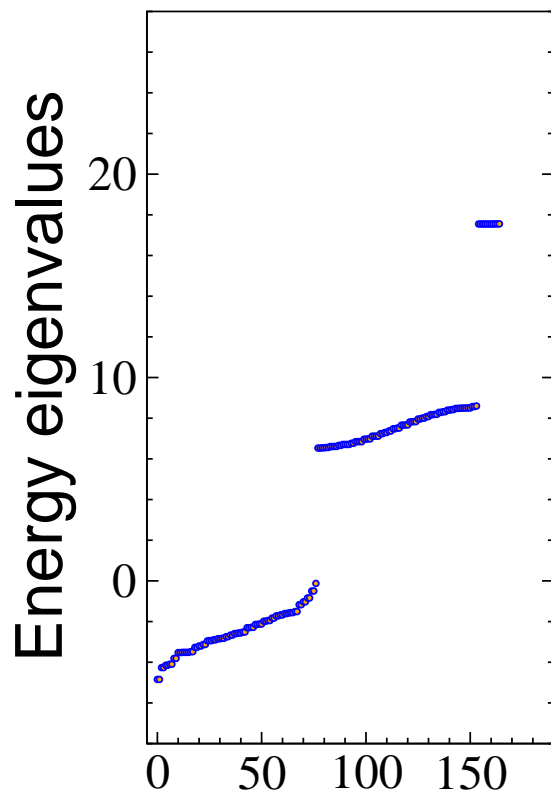
Physical t - V model has $V n_i n_{i+1}$, not $V(n_i - \frac{1}{2})(n_{i+1} - \frac{1}{2})$.

Physical t - V model does not have
empty-empty or empty-occupied energy.
(Only occupied-occupied energy.)

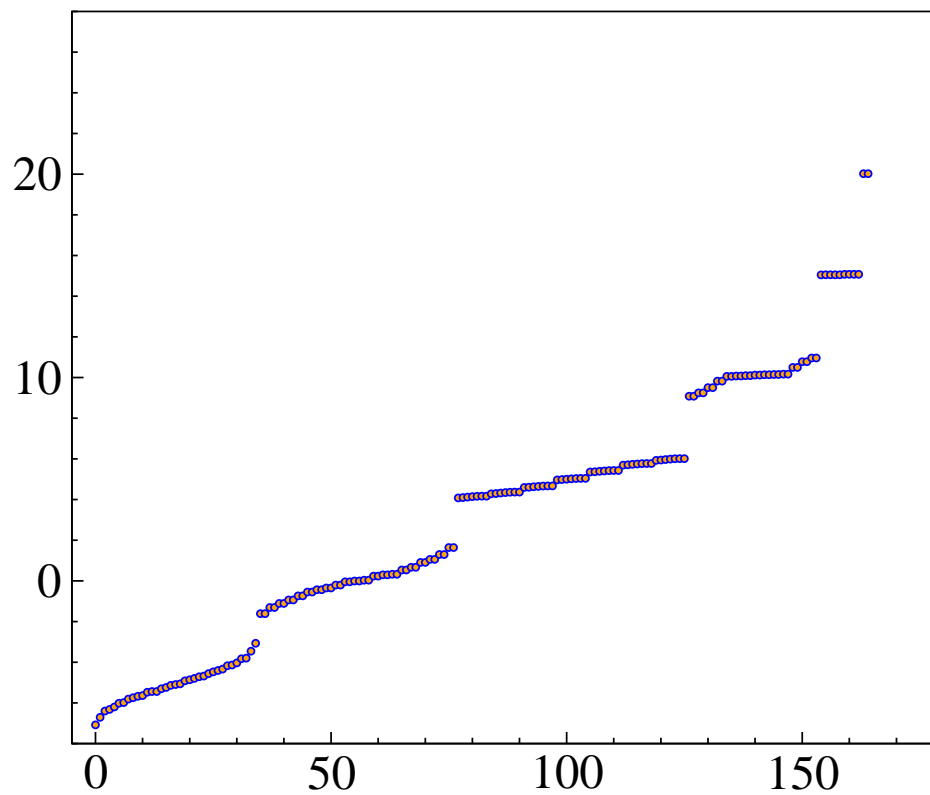


XXZ CHAIN: PERIODIC VERSUS OPEN SPECTRA

(a) Periodic chain



(b) Open chain



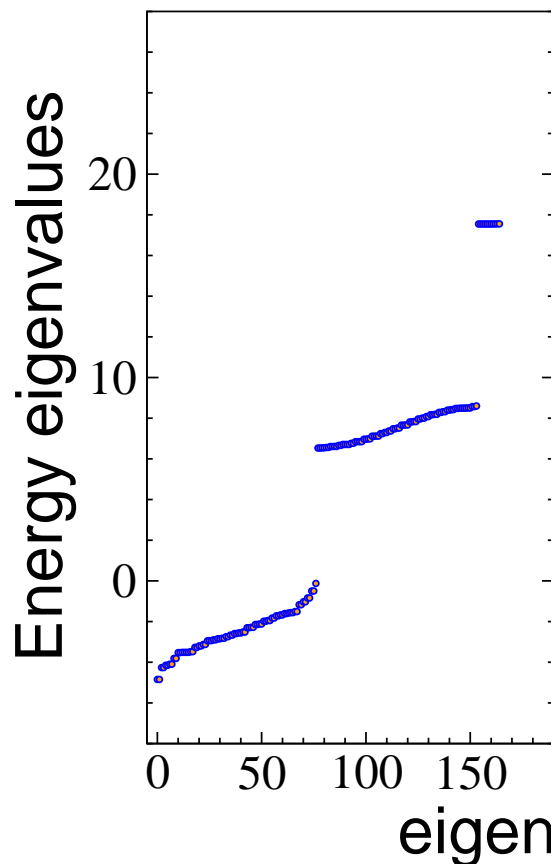
$$N_{\uparrow} = 3$$

11 sites

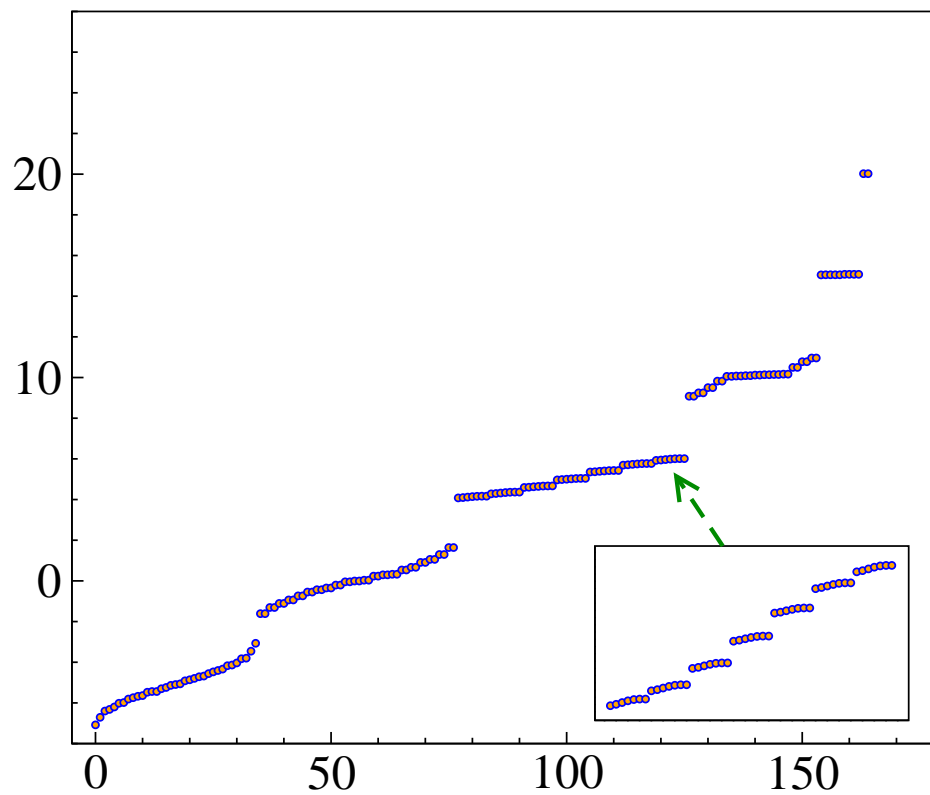
Many
extra
spectral
features
in open chain

XXZ CHAIN: PERIODIC VERSUS OPEN SPECTRA

(a) Periodic chain



(b) Open chain



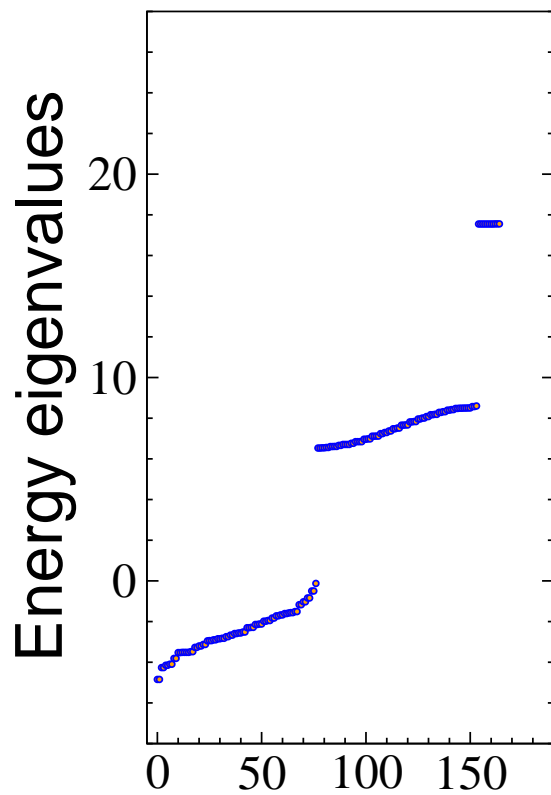
$$N_{\uparrow} = 3$$

11 sites

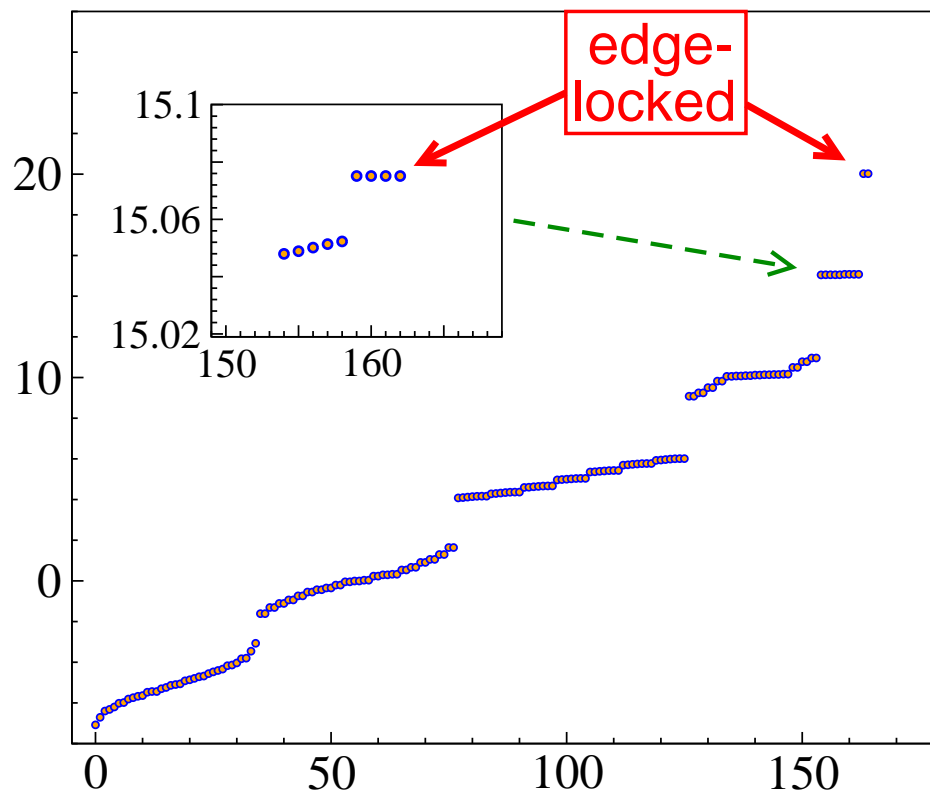
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XXZ CHAIN: PERIODIC VERSUS OPEN SPECTRA

(a) Periodic chain



(b) Open chain

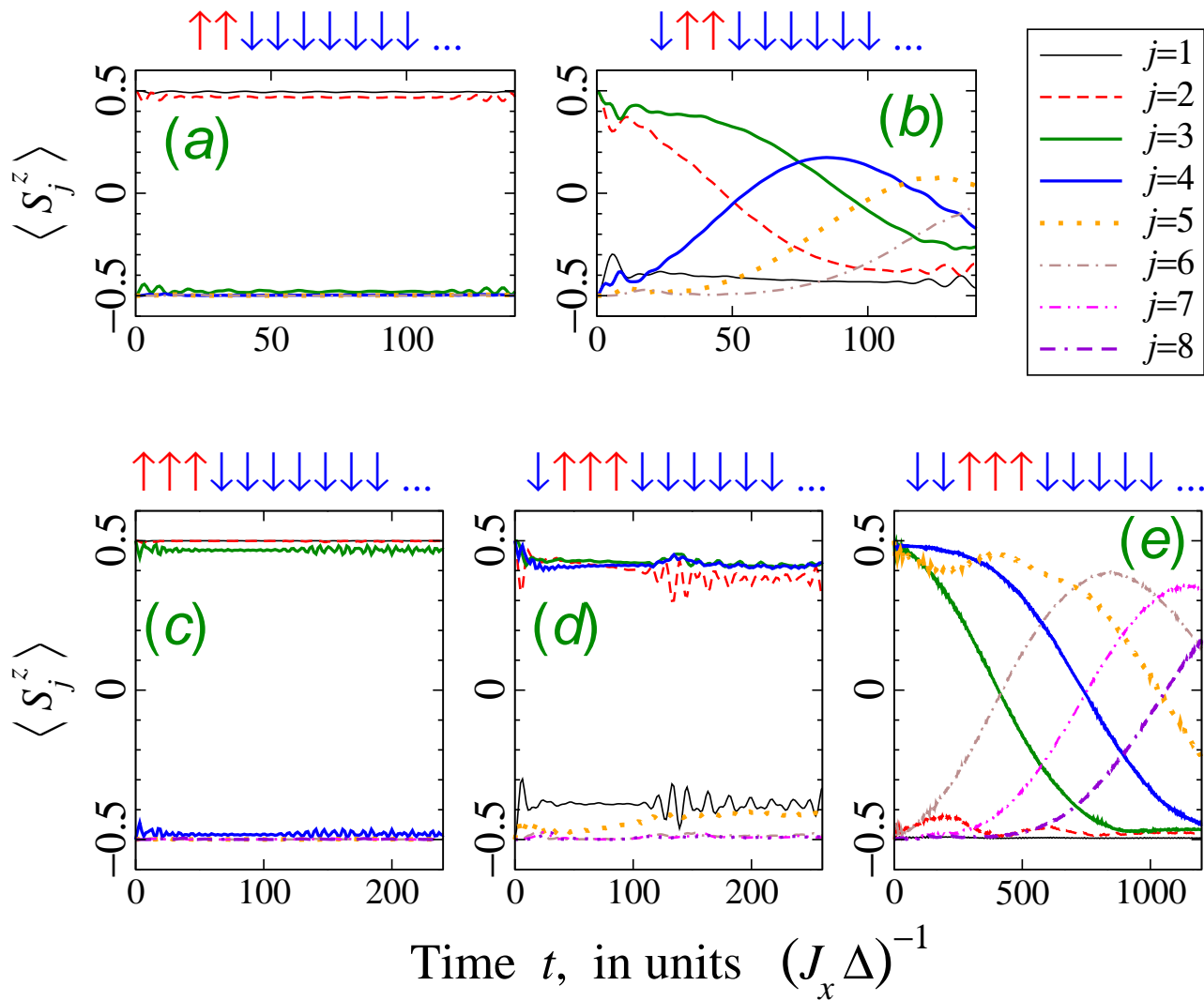


$$N_{\uparrow} = 3$$

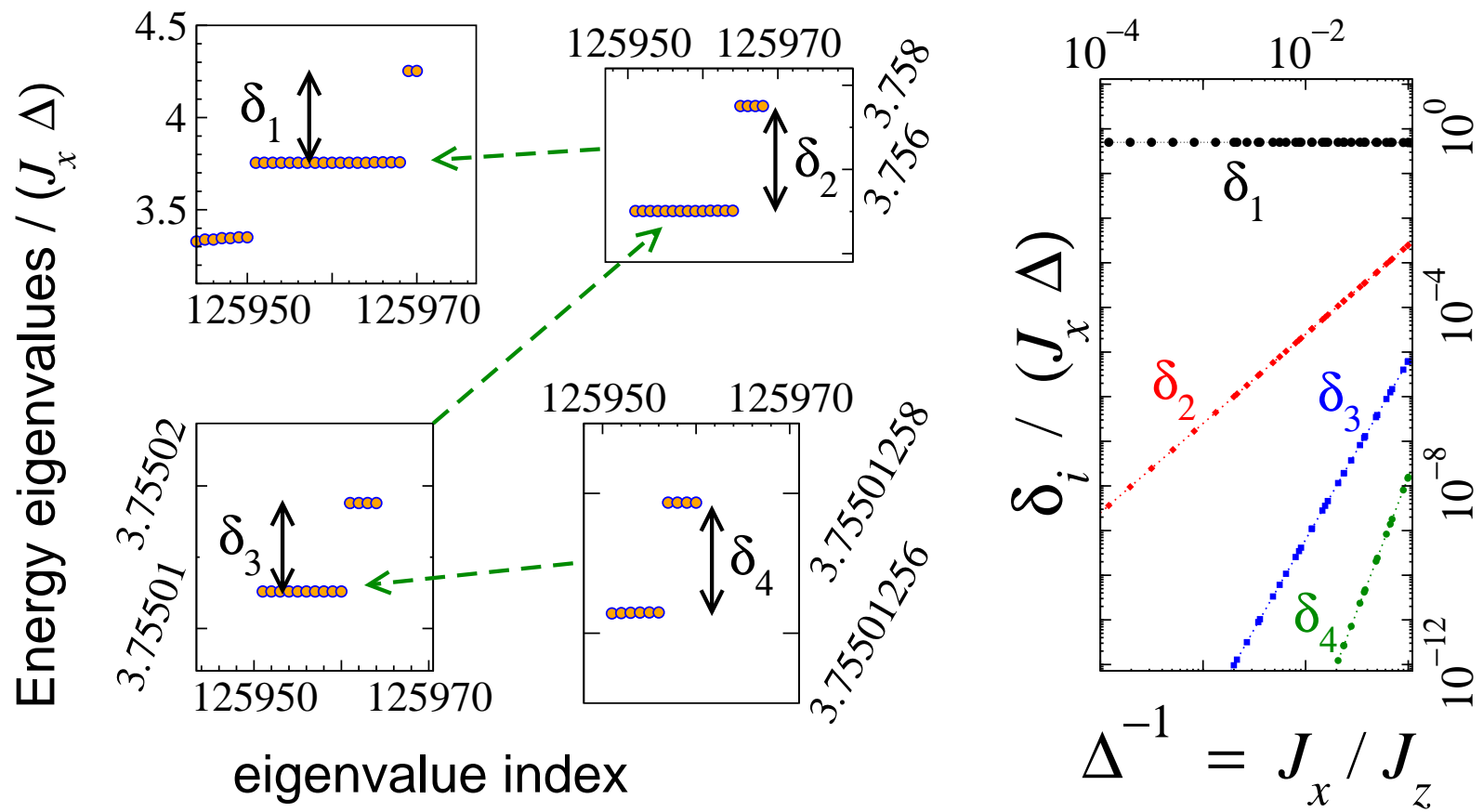
11 sites

Many
extra
spectral
features
in open chain

XXZ CHAIN: DYNAMICS



XXZ CHAIN: HIERARCHY



$N_{\uparrow} = 8$; 20 sites.

$\delta_1 \sim \Delta^0$

$\delta_2 \sim \Delta^{-2}$

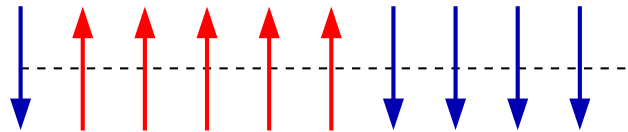
$\delta_3 \sim \Delta^{-4}$

HIERARCHY OF EDGE-LOCALIZATION

Energy spectrum contains structures at many different scales.

FRACTAL structure in spectrum

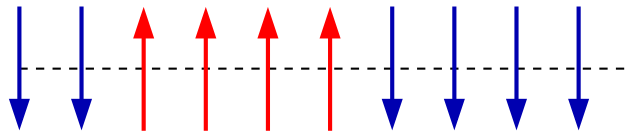
“QUANTUM CONTROL” OF MAGNETIZATION TRANSPORT



Single-site π -pulse



quantum switch
for unlocking magnetization.



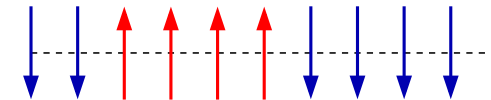
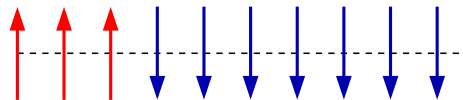
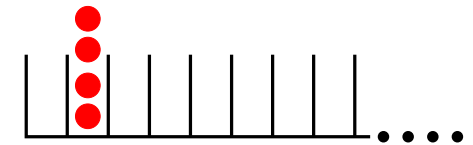
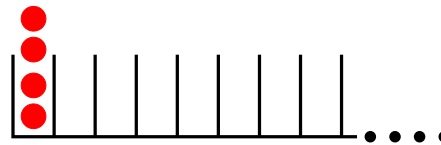
Many other control protocols....

EDGE-LOCALIZATION IN 1D LATTICE MODELS

Bose-Hubbard chain

spinless fermion model

XXZ chain



PHYSICS:

Eigenstates far from ground state

Far-from-equilibrium dynamics

Intricate structures in spectrum (**FRACTAL**)

QUANTUM CONTROL:

Locking and **release** of magnetization/state

Designing a **quantum switch**

EDGE LOCALIZATION: MORE BOSONS

For more bosons, a **hierarchy** of localization patterns.

$n \geq 5$ bosons \longrightarrow can also be bound in site 2

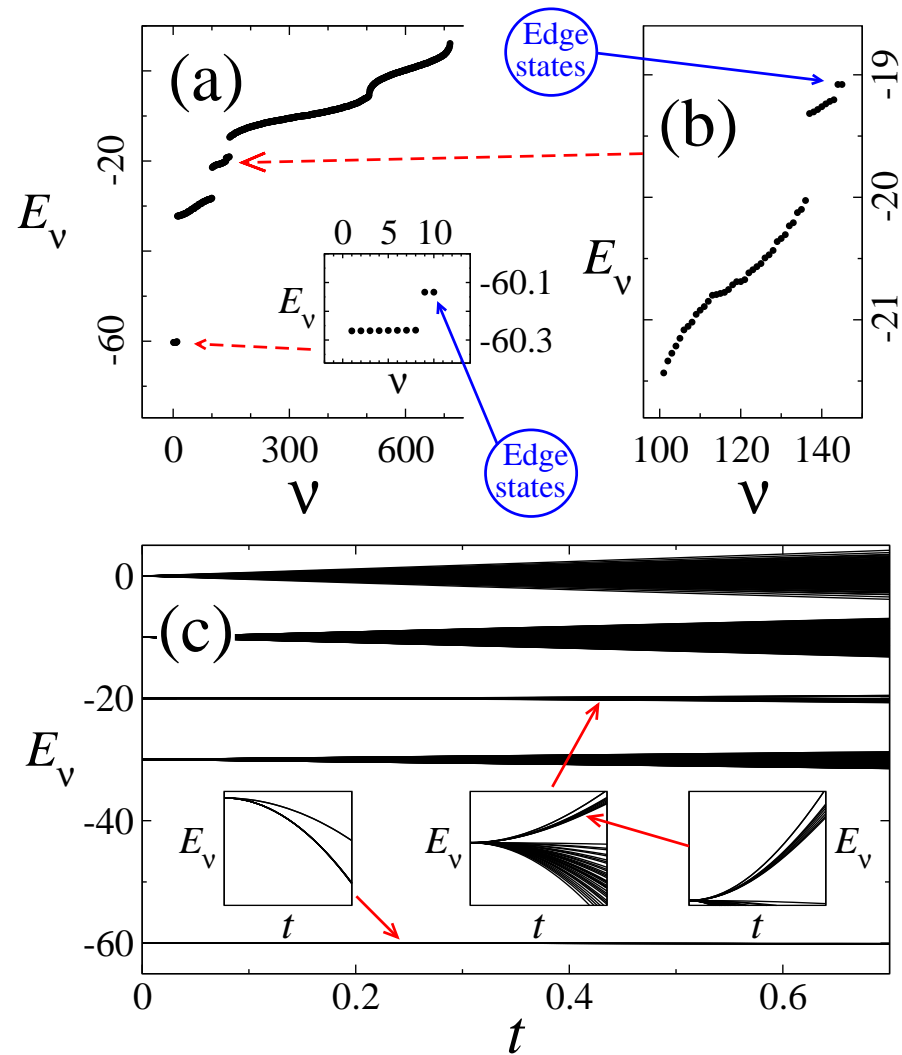
$n \geq 7$ bosons \longrightarrow can also be bound in site 3

... ..etc

Actually, several hierarchies, with other localization patterns:

2 2 0 0 0 0

EDGE-LOCALIZATION: SPECTRAL PICTURE (4 BOSONS)



ISN'T THIS JUST SELF-TRAPPING?

Question from
nonlinear dynamics
and/or BEC
community

ISN'T THIS JUST SELF-TRAPPING?

NO

AREN'T EDGE STATES UBIQUITOUS?

yes, but

No **single-particle** edge-localization in our tight-binding model

Our edge-locking is an **interaction** effect

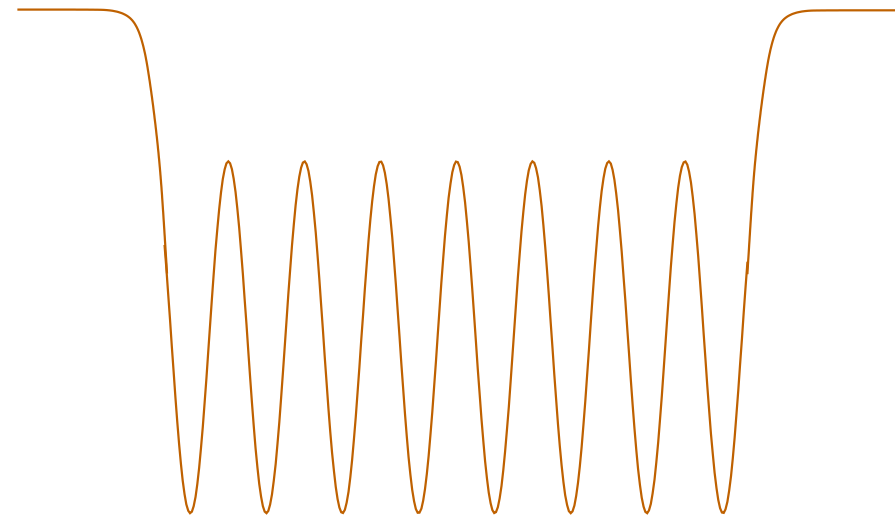
AREN'T EDGE STATES UBIQUITOUS?

yes, but

No single-particle edge-localization in our tight-binding model

Typical solid-state situation:

$$\hat{H} = -t \sum_{j=1}^{L-1} (c_j^\dagger c_{j+1} + c_{j+1}^\dagger c_j) - \lambda c_1^\dagger c_1 - \lambda c_L^\dagger c_L$$



EXPERIMENTAL REALIZATIONS

We would like

Clean and sharp edge,

Single-site addressability

Possible experimental settings:

Cavity polariton arrays

Josephson junction arrays

(ideal, but not yet realized)

1D optical lattices

Solid-state magnets with chain structures.

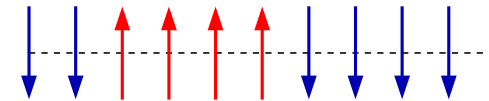
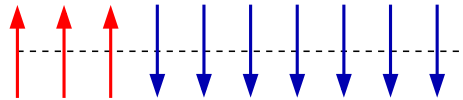
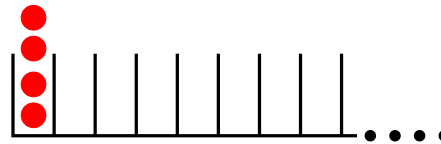
(not so ideal)

EDGE-LOCALIZATION IN 1D LATTICE MODELS

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