

Black Branes as Piezoelectrics

Jakob Gath

Niels Bohr Institute

October 4, 2012

[based on hep-th 1209.2127 with Jay Armas and Niels A. Obers]

The Holographic Way:
String Theory, Gauge Theory and Black Holes

Finding ways to describe and understand gravity

- Fluid/Gravity (inspired by AdS/CFT)
- Solid/Gravity (novel)

SETTING

This can be done by studying the dynamics and properties of
higher dimensional black branes

TOOL

Blackfold Approach

- Long wavelength effective theory.

[Empanan, Harmark, Niarchos, Obers, 0910.1601]

- Directly derivable from the Einstein equations.

[Camps, Empanan, 1201.3506]

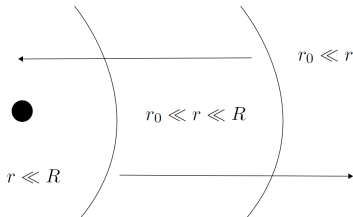
Applicability

- Applicable when one has separation of scales.
- Applicable when the fine-structure corrections dominates over backreaction corrections.
- Formalism can work for various backgrounds.

THE BLACKFOLD APPROACH

The effective dynamics of a black hole

Separation of scales: $r_0 \ll R$.



'Near zone' $r \ll R$

$$ds^2 = ds_{\text{sch.}}^2 + \mathcal{O}\left(\frac{r}{R}\right)$$

'Far zone' $r \gg r_0$

$$ds^2 = ds_{\text{background}}^2 + \mathcal{O}\left(\frac{r_0}{r}\right)$$

'Matching zone'

$$r_0 \ll r \ll R$$

First step of matched asymptotic expansion.

THE BLACKFOLD APPROACH

With black p -branes the near zone is described by

$$ds^2 = \left(\eta_{ab} + \frac{r_0^n}{r^n} u_a u_b \right) d\sigma^a d\sigma^b + \frac{dr^2}{1 - \frac{r_0^n}{r^n}} + r^2 d\Omega_{(n+1)}^2 .$$

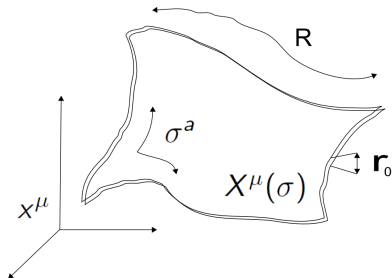
In $D = p + n + 3$ spacetime dimensions.

The thickness, velocity parameters, and embedding functions are promoted to collective fields over the worldvolume,

$$r_0 \rightarrow r_0(\sigma), \quad u_a \rightarrow u_a(\sigma), \quad X^\mu \rightarrow X^\mu(\sigma),$$

with worldvolume coordinates σ^a .

THE BLACKFOLD APPROACH



A blackfold is a *black* p -brane with thickness r_0 wrapped over a submanifold \mathcal{W}_{p+1} with characteristic length scale $R \gg r_0$ in the ambient space-time.

$$ds^2 = \left(\gamma_{ab}(X^\mu(\sigma)) + \frac{r_0^n(\sigma)}{r^n} u_a(\sigma) u_b(\sigma) \right) d\sigma^a d\sigma^b + \frac{dr^2}{1 - \frac{r_0^n(\sigma)}{r^n}} + r^2 d\Omega_{(n+1)}^2 + \dots$$

THE BLACKFOLD APPROACH

The black brane is in the far region ascribed an effective stress-energy tensor which can be computed using the quasilocal Brown-York formalism on a surface at large r .

For the black p -brane the effective stress tensor is,

$$T^{ab} = \frac{\Omega_{(n+1)}}{16\pi G} r_0^n(\sigma) \left(n u^a(\sigma) u^b(\sigma) - \gamma^{ab}(\sigma) \right) + \dots,$$

which at leading order is of the perfect fluid form,

$$T^{ab} = (\rho + P) u^a u^b + P \eta^{ab}.$$

THE BLACKFOLD APPROACH

The equation of motion of the blackfold

$$\nabla_\mu T^{\mu\nu} = 0,$$

Projecting along parallel and orthogonal directions this leads to

$$\mathcal{D}_a T^{ab} = 0, \quad \textit{intrinsic equations} \ (\mathcal{D}_\tau m = 0)$$

$$T^{ab} K_{ab}{}^\rho = 0, \quad \textit{extrinsic equations} \ (ma^\mu = 0)$$

When the brane is carrying charge

$$\mathcal{D}_a J^a = 0, \quad \textit{intrinsic}$$

Extrinsic Curvature: $K_{ab}{}^\rho = \perp^\rho{}_\lambda (\partial_a \partial_b X^\lambda + \Gamma_{\mu\nu}^\lambda \partial_a X^\mu \partial_b X^\nu)$

THE BLACKFOLD APPROACH

Past successes of the blackfold approach:

- Uncovering new horizon topologies for black holes.
[Empanan, Harmark, Niarchos, Obers, 0912.2352]
- Predicting the onset of GL-instabilities of black branes.
[Camps, Empanan, Haddad, 1003.3636]
- New method for thermal probe branes.

[Grignani, Harmark, Marini, Obers, Orselli, 1012.1494], [Niarchos, Siampos, 1205.1535 & 1206.2935]

BLACK BRANES AS FLUIDS AND SOLIDS

Goal: Show that black branes have properties of fluids and solids
one has to do something with them

Two ways of perturbing the brane with derivative corrections

- Intrinsic perturbations along the worldvolume
- Extrinsic perturbations transverse to the worldvolume

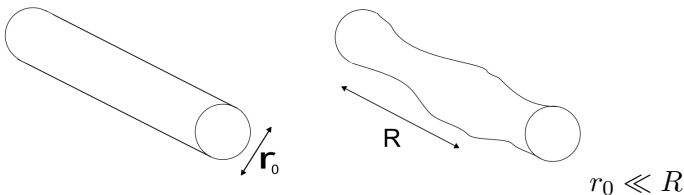
VISCOUS FLUIDS

First order derivative expansion

$$T_{ab} = \rho u_a u_b + P P_{ab} - \zeta \theta P_{ab} - 2\eta \sigma_{ab} + \mathcal{O}(\partial^2)$$

Viscosities

$$\text{shear: } \eta = \frac{s}{4\pi}, \quad \text{bulk: } \zeta = 2\eta \left(\frac{1}{p} - c_s^2 \right)$$



[Camps, Emparan, Haddad, 1003.3636]

Various connections between fluid and gravity have been found

- Fluid/AdS-gravity correspondence
- Membrane paradigm

The blackfold approach naturally introduces perturbations of extrinsic nature

SOLID/GRAVITY

First order Dirac-delta function series

$$T^{\mu\nu} \sim T_{(0)}^{\mu\nu} \delta(r) + T_{(1)\rho}^{\mu\nu} \partial_\rho \delta(r) + \dots$$

and similar for the electric current,

$$J^\mu \sim J_{(0)}^\mu \delta(r) + J_{(1)}^{\mu\rho} \partial_\rho \delta(r) + \dots$$

SOLID/GRAVITY

Working out the physical interpretation of the components within $T_{(1)}^{\mu\nu\rho}$ and $J_{(1)}^{\mu\rho}$ takes some work, one finds two interesting contributions,

- Intrinsic transverse angular momenta.
- Dipole moment of worldvolume stress-energy.

When a current is present

- Worldvolume electric dipole moment.

SOLID/GRAVITY

The dipole moments are developed as the response of bending which are probed by considering finite thickness corrections.

Stress-energy dipole contribution

$$T_{(1)}^{\mu\nu\rho} \sim u_a^\mu u_b^\nu d^{ab\rho}$$

with

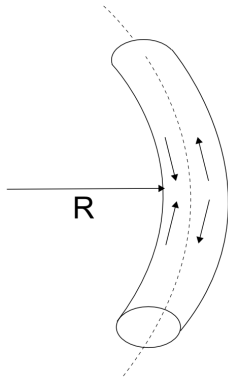
$$d^{ab\rho} = u_\mu^a u_\nu^b \perp^\rho{}_\lambda T_{(1)}^{\mu\nu\lambda}$$

and current dipole contribution

$$J_{(1)}^{\mu\rho} \sim u_a^\mu p^{a\rho}$$

with

$$p^{a\rho} = u_\mu^a \perp^\rho{}_\lambda J_{(1)}^{\mu\nu}$$



BLACK BRANES AS PIEZOELECTRICS

The dipole moment of worldvolume stress-energy $d^{ab\rho}$ and current $p^{a\rho}$ are not a priori constrained.

Under the expectation that bent black branes will behave like elastic solids and like piezoelectric material, one assume,

$$\begin{aligned}d^{ab\rho} &= \tilde{Y}^{abcd} K_{cd}{}^{\rho}, \\p^{a\rho} &= \tilde{\kappa}^{abc} K_{bc}{}^{\rho}\end{aligned}$$

with elastic moduli \tilde{Y}^{abcd} and piezoelectric moduli $\tilde{\kappa}^{abc}$.

MATCHED ASYMPTOTIC EXPANSION

Take the example of a black string

In order to construct *bent* string solutions one has to push the matched asymptotic expansion to next order in $1/R$.

- 0th (near/far); Boosted schwarzschild black string.
- 1st (far); Far field sourced by the effective stress tensor.
- 1st (near); Schwarzschild black string with $1/R$ corrections.

Boundary conditions are provided by matching with the far field and requiring regularity of the horizon.

Next: Read off the $1/R$ corrected stress tensor.

EXAMPLE

Black string carrying 0-brane charge in Einstein-Maxwell-Dilaton theory gives a realization of linear electroelasticity theory in gravitational physics.

$$\tilde{Y}^{ttzz} = -\frac{\Omega_{(n+1)}r_0^{n+2}}{16\pi G} \left(\frac{n^2(n+2)}{n+1} \sinh^2 \alpha + n^2 + 3n + 4 \right) \xi(n),$$

$$\tilde{Y}^{tzzz} = 0,$$

$$\tilde{Y}^{zzzz} = \frac{\Omega_{(n+1)}r_0^{n+2}}{16\pi G} (3n+4) \xi(n),$$

$$\tilde{\kappa}^{tzz} = -\frac{\Omega_{(n+1)}r_0^{n+2}}{8\pi G} (n+2) \xi(n),$$

$$\tilde{\kappa}^{zzz} = 0,$$

RESULT

The new response coefficients can be interpreted as black branes having piezoelectric behaviour.

This establish a realization between gravity and solids.

THANK YOU