

$$E_2 = J (\cos k_1 + \cos k_2 - 2\Delta)$$

$$\uparrow \quad p = k_1 + k_2$$

$$M=2$$

$$\Psi_2(j_1, j_2) = e^{ik_1 j_1 + ik_2 j_2} e^{-\frac{i}{2} \phi(k_1, k_2)}$$

$$- e^{ik_2 j_1 + ik_1 j_2} e^{+\frac{i}{2} \phi(k_1, k_2)}$$

Periodicity:

$$\Psi_2(j_2, j_1 + N) = \Psi_2(j_1, j_2) = \Psi_2(j_2 - N, j_1)$$

$$\Rightarrow e^{ik_1 N} = -e^{-i\phi(k_1, k_2)} \quad e^{ik_2 N} = -e^{+i\phi(k_1, k_2)}$$

Take logs:

$$Nk_1 + \phi(k_1, k_2) = 2\pi \tilde{I}_1$$

$$Nk_2 - \phi(k_1, k_2) = 2\pi \tilde{I}_2$$

$$\tilde{I} \in \mathbb{Z} + 1/2$$

General $M \leq N/2$ for any set $\{j_1, \dots, j_M\}$

$$\Psi_M(j_1, \dots, j_M) = \prod_{M \geq a > b \geq 1} \text{sgn}(j_a - j_b) \times$$

$$\sum_{\mathcal{P}_M} (-1)^{[\mathcal{P}]} e^{i \sum_{a=1}^M k_{p_a} j_a + \frac{i}{2} \sum_{M \geq a > b \geq 1} \text{sgn}(j_a - j_b) \phi(k_{p_a}, k_{p_b})}$$

This is the Bethe Ansatz

$$|\Psi_M\rangle = \sum_{\{j\}} \Psi_M(\{j\}) S_{j_1}^- \dots S_{j_M}^- |0\rangle$$

$$E_M = J \sum_{a=1}^M (\cos k_a - \Delta) \quad \mathcal{P} = \sum_a k_a$$

Bethe equations:

$$e^{ik_a N} = (-1)^{M-1} e^{-i \sum_{b=1}^M \phi(k_a, k_b)} \quad a=1, \dots, M$$

$$N k_a + \sum_{b=1}^M \phi(k_a, k_b) = 2\pi \tilde{I}_a$$

$$\tilde{I}_a \in \begin{cases} \mathbb{Z} + 1/2 & M \text{ even} \\ \mathbb{Z} & M \text{ odd} \end{cases}$$

XXX antiferromagnet ($J > 0$, $\Delta = 1$)

Parametrization: $k = \pi - 2 \arctan 2\lambda$

rapidities $\rightarrow \lambda = \frac{1}{2} \cot \frac{k}{2}$

Why? $\phi(k(\lambda_1), k(\lambda_2)) \rightarrow \Theta(\lambda_1 - \lambda_2)$

$$\Theta(\lambda) = 2 \arctan \lambda$$

$$\text{Bethe } \varrho_{\lambda}^{\pm} = \left[\frac{\lambda_a + i/2}{\lambda_a - i/2} \right]^N = \prod_{b \neq a}^M \frac{\lambda_a - \lambda_b + i}{\lambda_a - \lambda_b - i}$$

$$\text{or } \Theta_1(\lambda_a) - \frac{1}{N} \sum_{b=1}^M \Theta_2(\lambda_a - \lambda_b) = \frac{2\pi}{N} I_a$$

$$\Theta_n(\lambda) \equiv 2 \arctan \frac{2\lambda}{n} \quad I_a \in \begin{cases} \mathbb{Z} + \frac{1}{2} & N-M \text{ even} \\ \mathbb{Z} & \text{" odd} \end{cases}$$

$$E = J \sum_{a=1}^M \frac{-2}{4\lambda_a^2 + 1}$$

$$P = \sum_{a=1}^M \frac{1}{i} \ln \left[\frac{\lambda_a + i/2}{\lambda_a - i/2} \right] = \pi M - \frac{2\pi}{N} \sum_{a=1}^M I_a \pmod{2\pi}$$

Proper solutions to the Bethe equations: must have distinct rapidities.

(most cases: this translates to choosing different quantum numbers in Bethe equations)

Ground state: at fixed M, N even

quantum numbers $I_j^0 = -\frac{M+1}{2} + j \quad j=1, \dots, M$

$M=5^0$: $\circ \circ \quad \bullet \quad \bullet \quad \bullet \quad \bullet \quad \bullet \quad \bullet \quad \bullet \quad \bullet \quad \bullet$

-2 -1 0 1 2

Allowable quantum nos: careful (4 notes)

For real λ , I cannot be arbitrarily large.

$$\lim_{\lambda_a \rightarrow \infty} 2 \operatorname{atan} 2\lambda_a - \frac{1}{N} \sum_{b=1}^M 2 \operatorname{atan} (\lambda_a - \lambda_b)$$

$$= \pi - \frac{1}{N} \pi(M-1) = \frac{2\pi}{N} I_M^\infty$$

$$I_M^\infty = \frac{N-M+1}{2}$$

$$|\lambda_a| < \infty \text{ if } |I_a| \leq I_M^{\max} \equiv I_M^\infty - 1$$

$$I_{M=N}^\infty = \frac{N}{2} + \frac{1}{2}$$

Simple case: $M = N/2$. For GS:

$$\{I\} = \left\{ -\frac{M-1}{2}, -\frac{M-1}{2} + 1, \dots, \frac{M-1}{2} \right\} = \left\{ -\frac{N}{4} + \frac{1}{2}, \dots, \frac{N}{4} - \frac{1}{2} \right\}$$

$$I_{M=N}^{\max} = \frac{N}{2}$$

