

Evidence of anisotropic Kondo coupling in nanostructured devices

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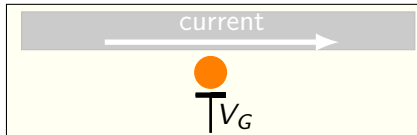
- Elementary geometries

- Single electron transistor



- V_G controls dot occupation n_d

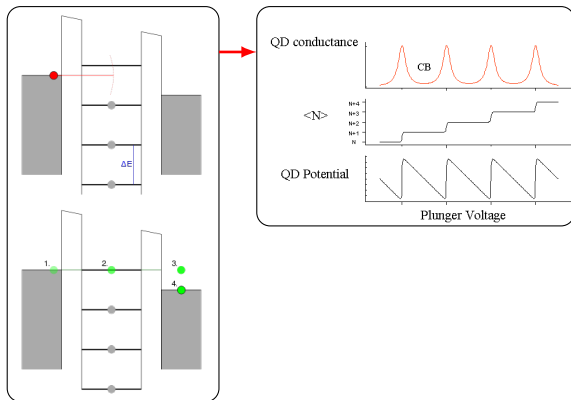
- Side-coupled device



Coulomb blockade



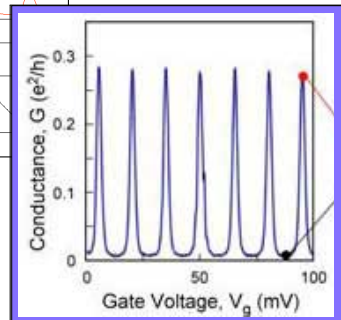
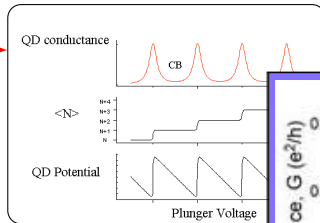
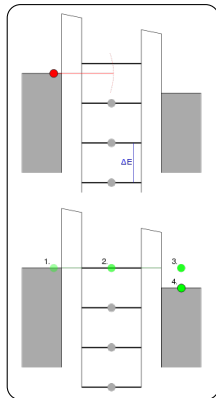
- Conduction only when n_{dot} is a half-integer



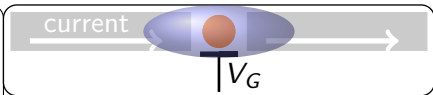
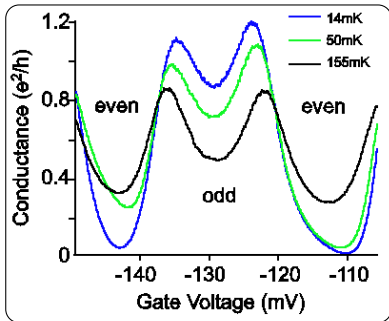
Coulomb blockade



- Conduction only when n_{dot} is a half-integer

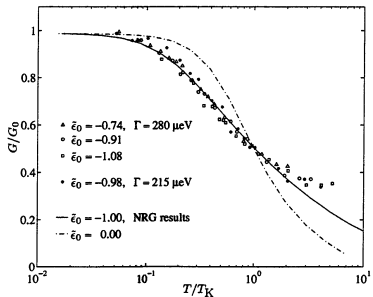


- Coulomb blockade can be lifted at very low T



Goldhaber-Gordon et al.,
PRL **81**, 5225 (1998)

- Odd n_{dot}
 - Dot has magnetic moment
 - Conduction spins screen moment
 - Screening allows conduction



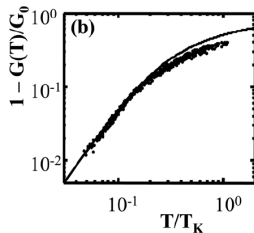
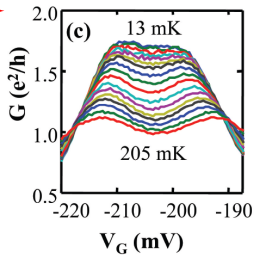
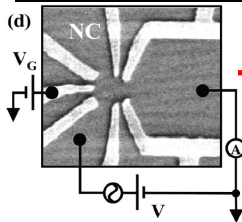
Outline of the talk

- Introduction
- More recent experimental results
- Theory
- Theory vs. experiment and a puzzle
- Conclusions

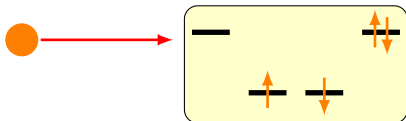
Ten years after 1998



Grobis et al, PRL **100**, 246601 (2008).



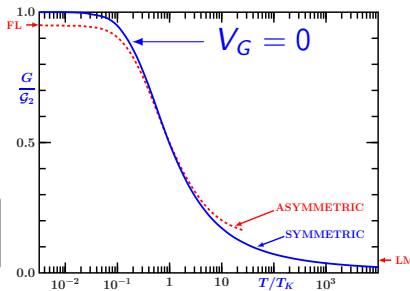
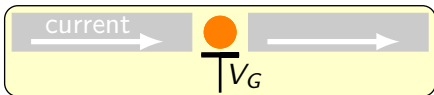
- Anderson Model



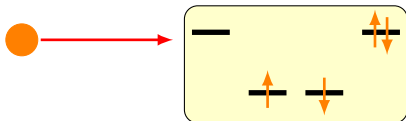
- $$H = \sum_k \epsilon_k n_k - \frac{U}{2} (n_{d\uparrow} - n_{d\downarrow})^2 + V_G n_d + V \sum_k (c_k^\dagger c_d + H. c.)$$

- $$G = \mathcal{G}_2 \mathcal{G}_S \left(\frac{T}{T_K} \right)$$

$$\mathcal{G}_2 \equiv \frac{2e^2}{h}$$

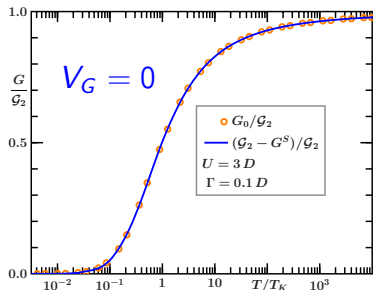
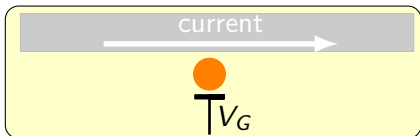


- Anderson Model



- $$H = \sum_k \epsilon_k n_k - \frac{U}{2} (n_{d\uparrow} - n_{d\downarrow})^2 + V_G n_d + V \sum_k (c_k^\dagger c_d + H. c.)$$

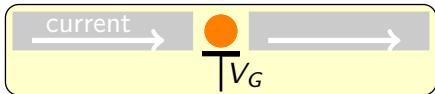
- $$G = \mathcal{G}_2 \left[1 - G_S \left(\frac{T}{T_K} \right) \right]$$



- Anderson model

- $$H = \sum_k \epsilon_k n_k - \frac{U}{2} (n_{d\uparrow} - n_{d\downarrow})^2 + V_g n_d + V \sum_k (c_k^\dagger c_d + H. c.)$$

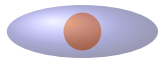
- $V_G \neq 0$



$$G = \alpha G_S \left(\frac{T}{T_K} \right) + \beta \left[1 - G_S \left(\frac{T}{T_K} \right) \right]$$

Seridonio et al.,
EPL **86**, 67006(2009).

- $G(T \ll T_K) = \alpha = \kappa \mathcal{G}_2 \sin^2\left(\frac{\pi n_{LT}}{2}\right) \equiv G_{FL}$



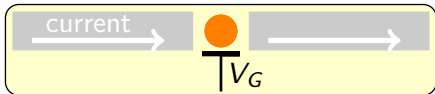
- $G(T \gg T_K) = \beta = \kappa \mathcal{G}_2 \sin^2\left(\frac{\pi n_{HT}}{2}\right) \equiv G_{LM}$



- Anderson model

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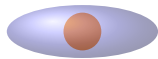
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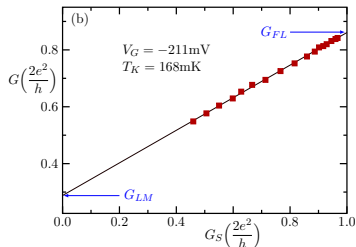
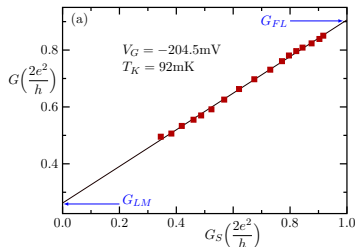
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Comparison with experiment

- $G\left(\frac{T}{T_K}\right)$ maps linearly onto $G_S\left(\frac{T}{T_K}\right)$

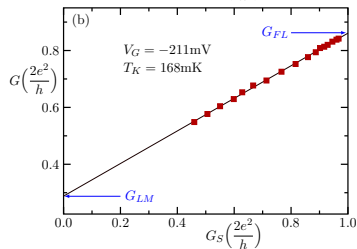
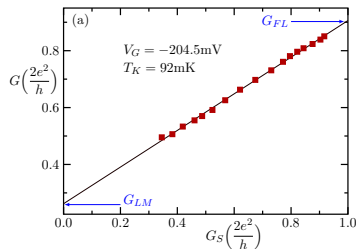
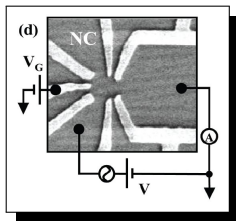
$$G\left(\frac{T}{T_K}\right) = G_S\left(\frac{T}{T_K}\right) G_{FL} + \left[1 - G_S\left(\frac{T}{T_K}\right)\right] G_{LM}$$



Comparison with experiment

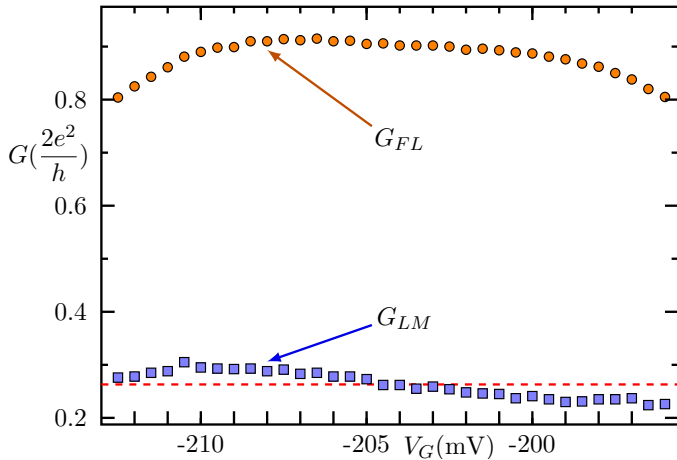
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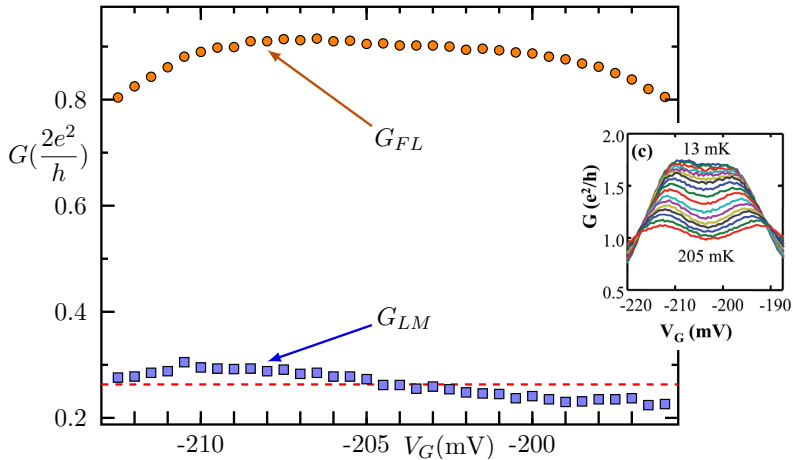
Comparison with experiment

- High (LM) and low (FL) temperatures



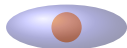
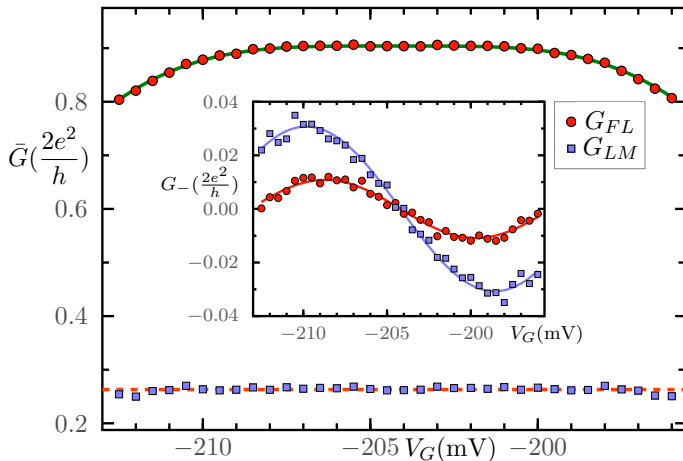
Comparison with experiment

- High (LM) and low (FL) temperatures



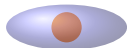
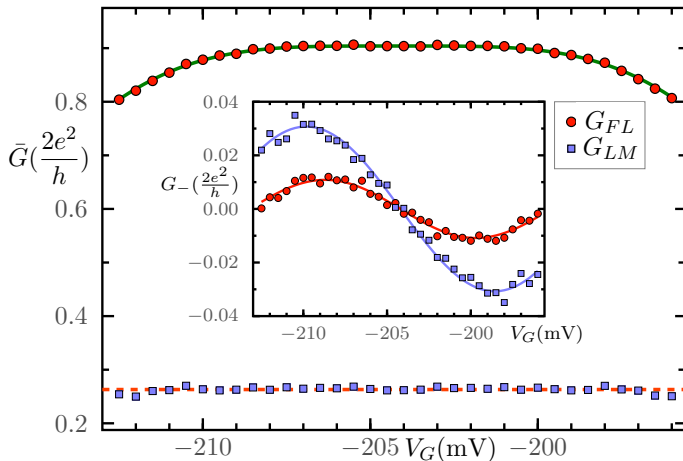
Comparison with experiment

- Symmetrized conductances at high (LM) and low T (FL)



Comparison with experiment

- Symmetrized conductances at high (LM) and low T (FL)

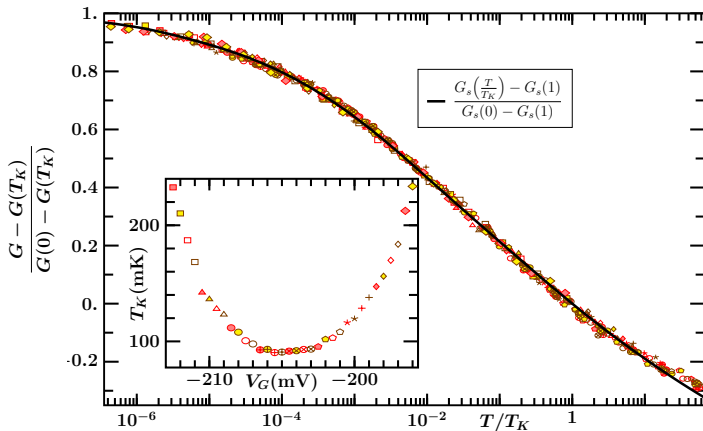


Anisotropy?



Comparison with experiment

- Results for 34 V_G 's



- Anderson model describes experimental data very well
 - Thanks to universality
 - There are perturbations outside the scope of the model
- Must allow for partial screening at high T
 - Anisotropic Kondo coupling?