Interacting fermions in flux lattices: Ferromagnetism and nematic ordering

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see also Phys. Rev. Lett. 109, 265301 (2012)

Motivation

Itinerant ferromagnetism with cold fermions is (so far) hampered by losses



Novel effects of interacting fermions in optical flux lattices

- Momentum dependent interactions
- Topological flat bands (2D)



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rapid decay into bound pairs is observed over times on the order of 10 h/E_F, preventing the study of equilibrium phases of strongly repulsive fermions. Our work suggests that a Fermi gas with strong short-range repulsive interactions does not undergo a ferromagnetic phase transition."

Sanner et al, PRL 2012

Novel effects of interacting fermions in optical flux lattices

 Momentum dependent interactions
 Topological flat bands (2D)



- Flux lattice: Example spin 1/2
- Toy model: Consider spin texture in 2D



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- Flux lattice: Example spin 1/2
- Toy model: Consider spin texture

x-y components of Bloch vector





Wraps 1/2 of Bloch sphere

• Flux lattice: Lattice potential with periodic texture with net positive flux



Gauge potential $\mathbf{A} = i \langle \psi | \nabla_{\mathbf{r}} \psi \rangle$ Flux density $n_{\phi} = (\nabla \times \mathbf{A}) / (2\pi)$ $n_{\phi} = \frac{-1}{4\pi} \mathbf{n} \cdot (\partial_x \mathbf{n} \times \partial_y \mathbf{n})$

 $4 \times (1/2) = 2$ Flux quanta per unit cell

• Flux lattice: Lattice potential with periodic texture with net positive flux



Spin dep. potential

$$H = \frac{p^2}{2m} + \sigma \cdot \mathbf{B}(\mathbf{r})/2$$

Adiabatic approx.

$$\psi(\mathbf{r}) = \begin{pmatrix} e^{-i\phi}\cos(\theta/2) \\ \sin(\theta/2) \end{pmatrix}$$

 $\mathbf{n}(\mathbf{r}) = \psi^{\dagger}(\mathbf{r})\sigma\psi(\mathbf{r}) \mathbf{n} \mid| -\mathbf{B}$

• Flux lattice: Lattice potential with periodic texture with net positive flux



Spin dep. potential p^2

$$H = \frac{P}{2m} + \sigma \cdot \mathbf{B}(\mathbf{r})/2$$

state dep. light shift $\sigma \cdot \mathbf{B}(\mathbf{r}) = \begin{pmatrix} B_z & B_x + iB_y \\ B_x - iB_y & -B_z \end{pmatrix}$

Raman dressing



Cooper and Dalibard, EPL 95, 66004 (2011)

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Two photon Raman-scheme for
e.g. <sup>171</sup>Yb with J=1/2
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3 Parameters to play with:

1) Lattice depth V_0/E_R 2) Polarization angle θ 3) Raman coupling ϵ



Cooper and Dalibard, EPL 95, 66004 (2011)

Two photon Raman-scheme for e.g. ¹⁷¹Yb with J=1/2







Cooper and Dalibard, EPL 95, 66004 (2011)

Two photon Raman-scheme for e.g. ¹⁷¹Yb with J=1/2

Fermi surface KK'(weak lattice)





Interactions

Project interactions onto lowest band

$$H = \sum_{\mathbf{k}} c_{\mathbf{k}}^{\dagger} c_{\mathbf{k}} \epsilon_{\mathbf{k}} + \sum_{\mathbf{k}_{1} \mathbf{k}_{2} \mathbf{k}_{3} \mathbf{k}_{4}} V_{\mathbf{k}_{1} \mathbf{k}_{1} \mathbf{k}_{3} \mathbf{k}_{4}} c_{\mathbf{k}_{1}}^{\dagger} c_{\mathbf{k}_{2}}^{\dagger} c_{\mathbf{k}_{3}} c_{\mathbf{k}_{4}}$$
$$V_{\mathbf{k}_{1} \mathbf{k}_{2} \mathbf{k}_{3} \mathbf{k}_{4}} = g_{2D} \int d^{2} \mathbf{r} \sum_{\sigma} \phi_{\mathbf{k}_{1}\sigma}^{*}(\mathbf{r}) \phi_{\mathbf{k}_{2}\bar{\sigma}}^{*}(\mathbf{r}) \phi_{\mathbf{k}_{3}\bar{\sigma}}(\mathbf{r}) \phi_{\mathbf{k}_{4}\sigma}(\mathbf{r})$$

Interacting spinless fermions! (starting from contact s-wave int.)

Cooper, Dalibard EPL (2011); Williams et al., Science (2012); Cui PRA (2012) ...

Interactions

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$$V_{\mathbf{k}_{1} \mathbf{k}_{2} \mathbf{k}_{3} \mathbf{k}_{4}} = g_{2D} \int d^{2} \mathbf{r} \sum_{\sigma} \phi_{\mathbf{k}_{1}\sigma}^{*}(\mathbf{r}) \phi_{\mathbf{k}_{2}\overline{\sigma}}^{*}(\mathbf{r}) \phi_{\mathbf{k}_{3}\overline{\sigma}}(\mathbf{r}) \phi_{\mathbf{k}_{4}\sigma}(\mathbf{r})$$

• Hartree-Fock decoupling

$$\begin{split} E[\{n_{\mathbf{k}}\}] &= \sum_{\mathbf{k}} \epsilon_{\mathbf{k}} n_{\mathbf{k}} + \frac{1}{2} \sum_{\mathbf{k}\mathbf{k}'} V_{\mathbf{k}\mathbf{k}'} n_{\mathbf{k}} n_{\mathbf{k}'} \\ V_{\mathbf{k}\mathbf{k}'} &= V_{\mathbf{k}\mathbf{k}'\mathbf{k}'\mathbf{k}} - V_{\mathbf{k}\mathbf{k}'\mathbf{k}\mathbf{k}'} \\ \mathbf{Hartree} \quad \mathbf{Fock} \end{split}$$

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Hartree-Fock phase diagram



Hartree-Fock phase diagram



Up to 4 times weaker bare interaction needed for Stoner.

Huge increase in lifetime at Stoner transition (~100)

Phasediagram: finite T 0 (a) 0.14 N Temperature/Bandwidth ky 0.12 0.10

2.0

$$\begin{split} \tilde{g}_{2D} &\sim 1.9 \\ V_0 &\sim 2E_R \\ (W + \Delta)/W &\sim 4 \end{split}$$

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1.5

1.0

0.08

0.06

0.04

0.02

Nematic order of fermions: Ubiquitous in correlated fermion systems



See review by Fradkin, Kivelson, Lawler, Eisenstein, McKenzie, Ann. Rev. Cond. Mat. 1, 153 (2010)

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Origin of nematic order: Anisotropic interaction between identical fermions



Connection to models with Ising nematic order

Mean-field model for Ising nematic (2D square)



Yamase, Oganesyan, Metzner, PRB 72, 35114 (2005)

Connection to models with Ising nematic order

Mean-field model for Ising nematic (2D square)



Conclusion

- Optical flux lattices offer route to itinerant ferromagnetism at weaker bare interaction strength
- Coupling between (pseudo) spin and lattice gives rise to lsing-nematic symmetry breaking
- Open questions: Competition of nematic order/magnetism with other correlated states (FQH)