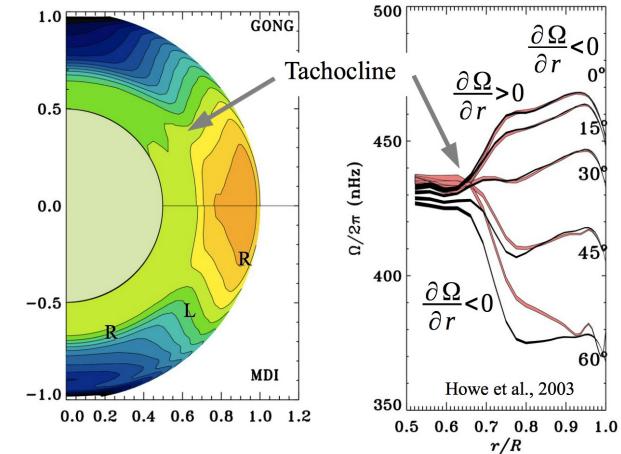


Solar differential rotation hints to obtain a near- surface shear layer



Gustavo Guerrero
Solar Physics (HEPL)
Stanford University

P. Smolarkiewicz,
A. Kosovichev, N. Mansour

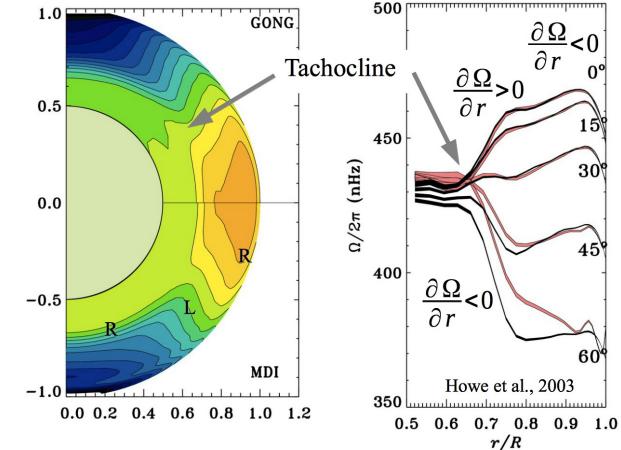


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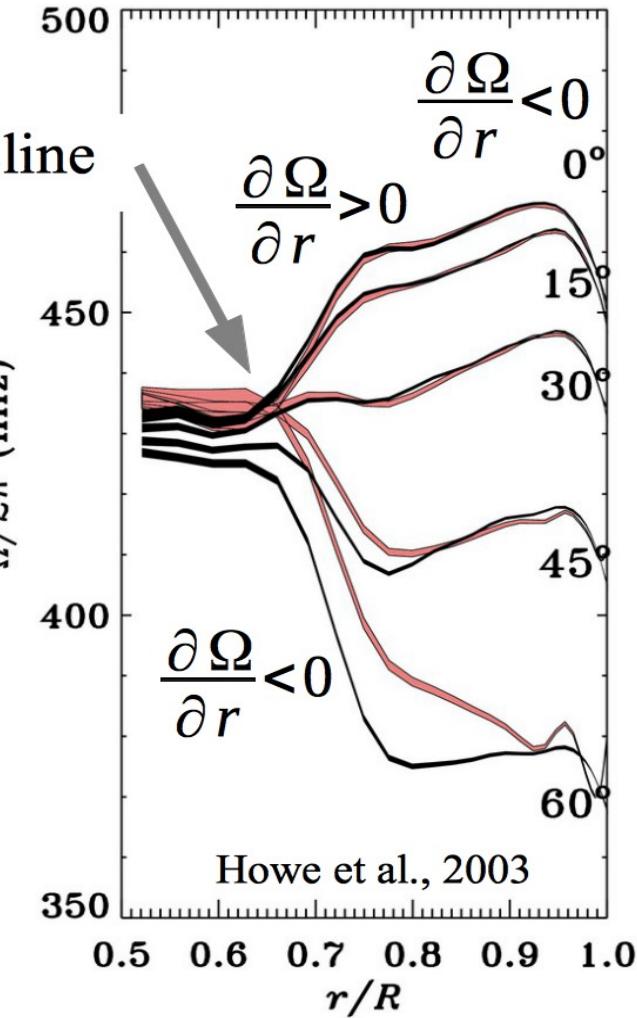
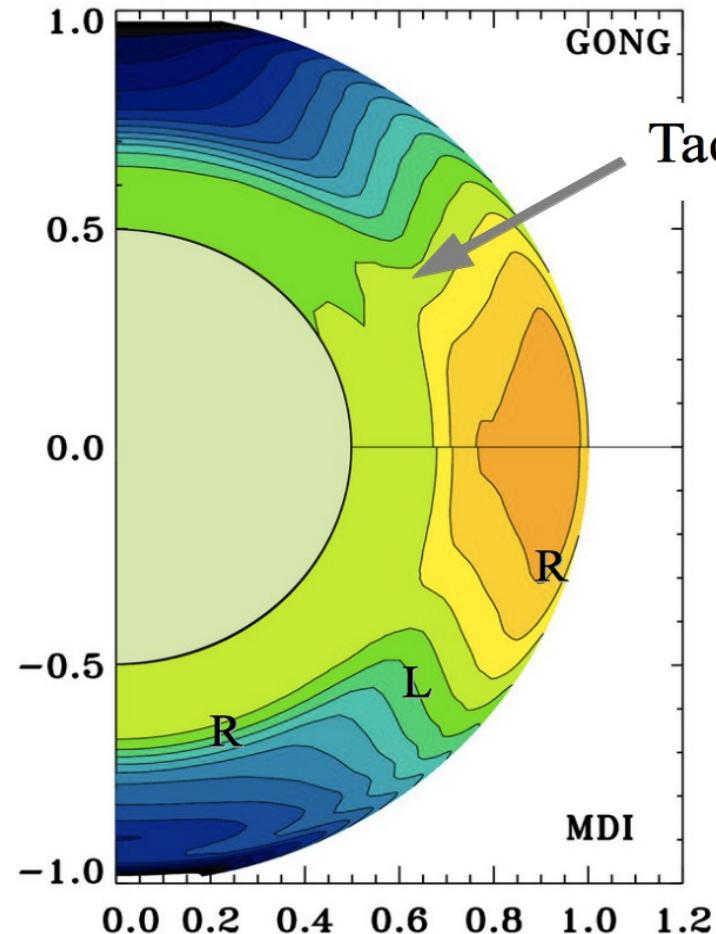


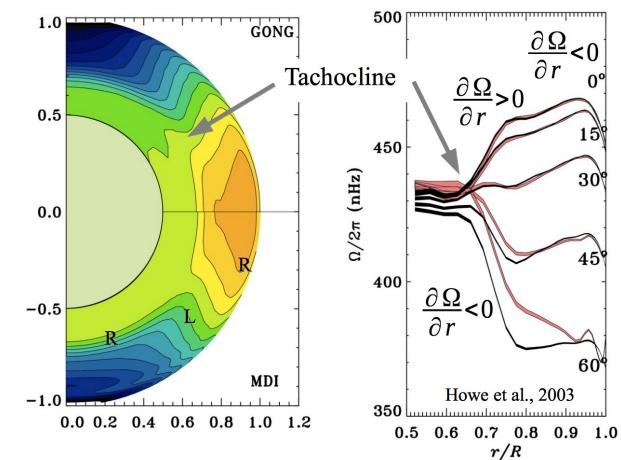
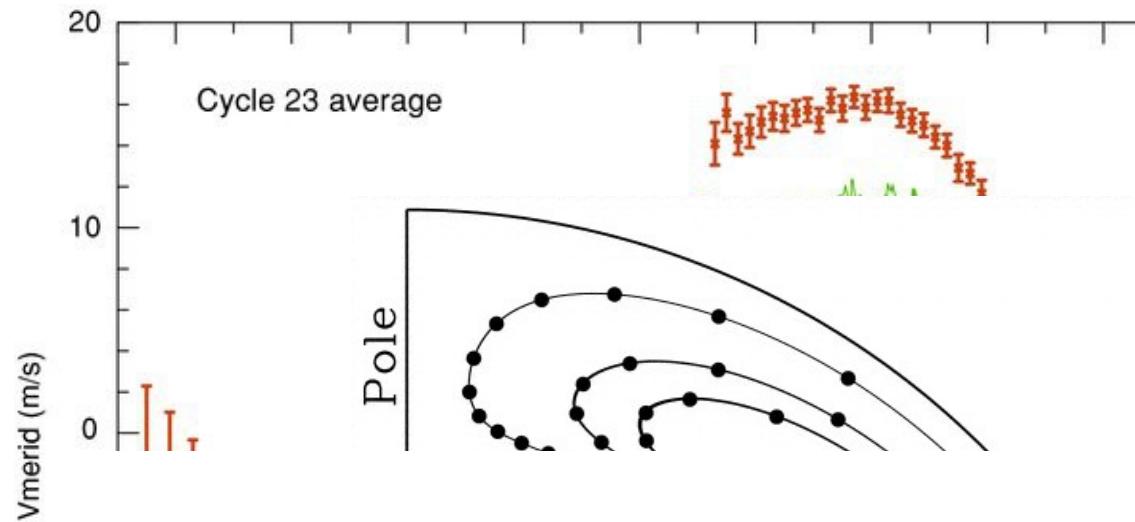
Outline

1. Large-scale flows in the solar interior
2. Mean-field models, DNS and LES
3. Anelastic simulations with EULAG
 - a) Angular momentum transport
 - b) Convergence
 - c) Near-surface shear layer
4. Conclusions

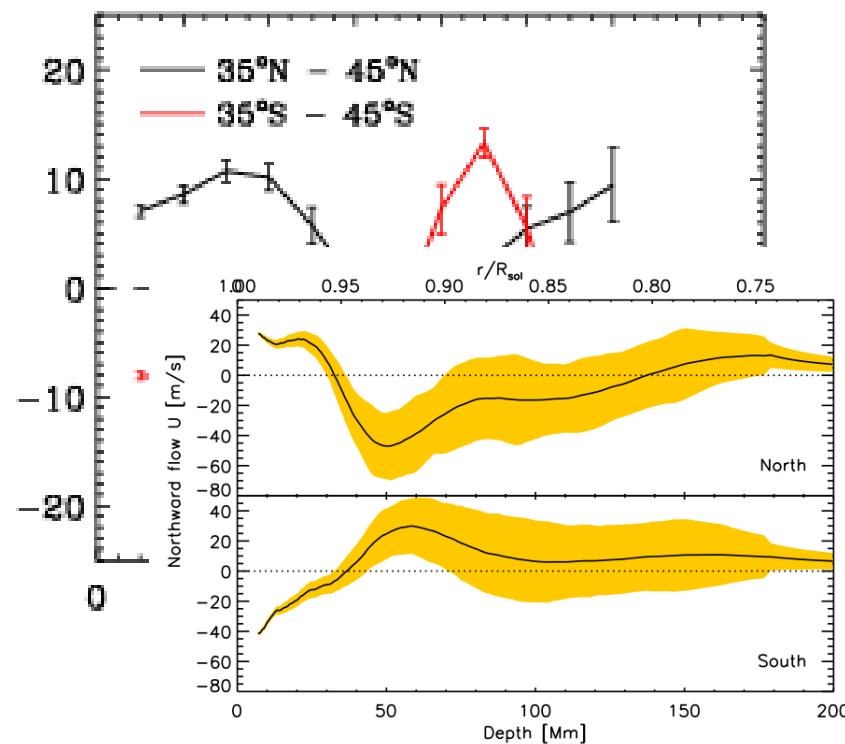
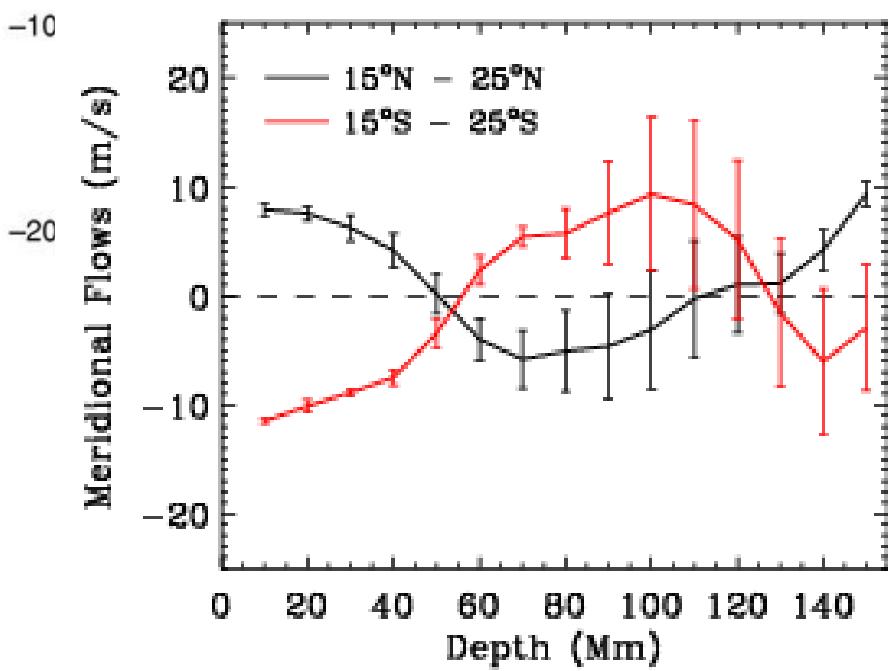


1. Large-scale motions in the solar interior





Courtesy J. Zhao



Mitra-Kraev & Thompson (2007)

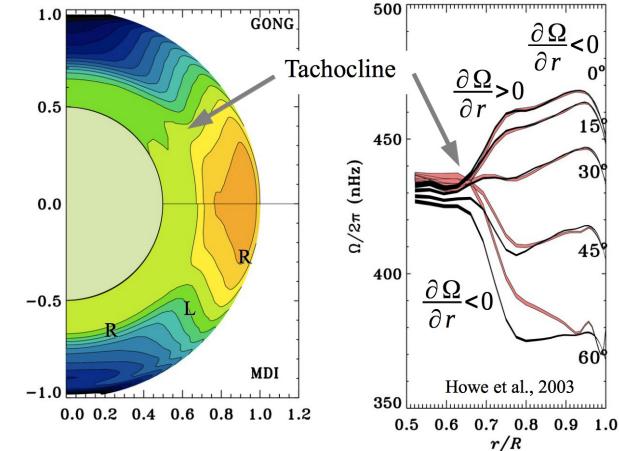
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2. mean-field models

DNS and LES

Large-scale motions as well as large-scale magnetic field are the result of collective effects of turbulence



$$\frac{\partial \overline{\mathbf{B}}}{\partial t} = \nabla \times (\overline{\mathbf{U}} \times \overline{\mathbf{B}} + \mathbf{E} + \eta \nabla \times \overline{\mathbf{B}})$$

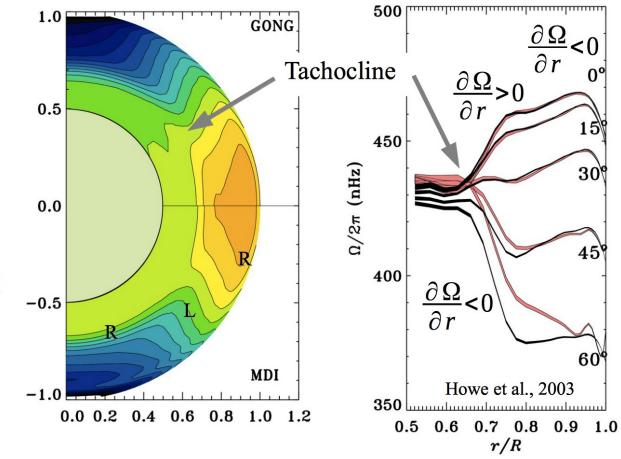
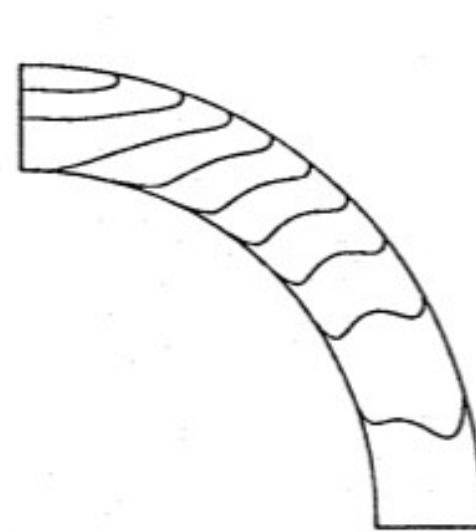
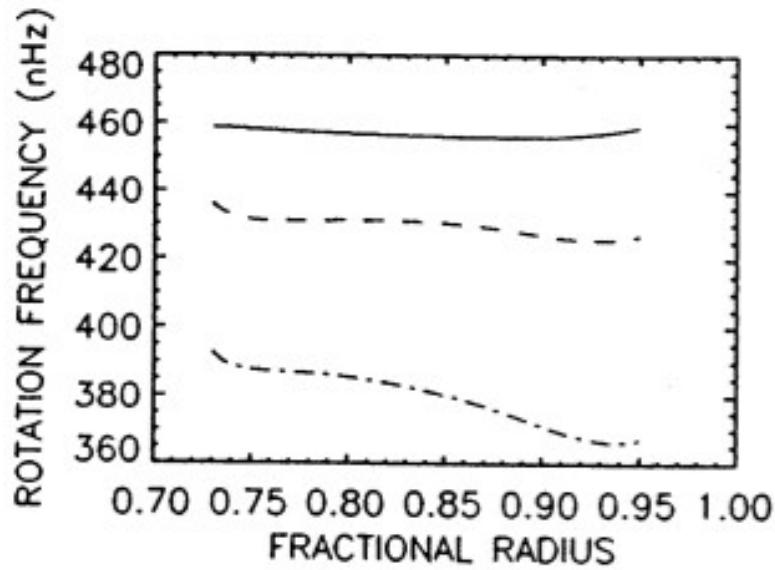
$$\mathcal{E} = \langle \mathbf{u}' \times \mathbf{b}' \rangle \quad \text{(Steenbeck et al. 66)}$$

$$\frac{D\overline{\mathbf{U}}}{Dt} = -\frac{\nabla \overline{p}}{\rho} - \nabla \cdot (\rho Q) + \rho \mathbf{g} + \nabla \cdot \boldsymbol{\pi}$$

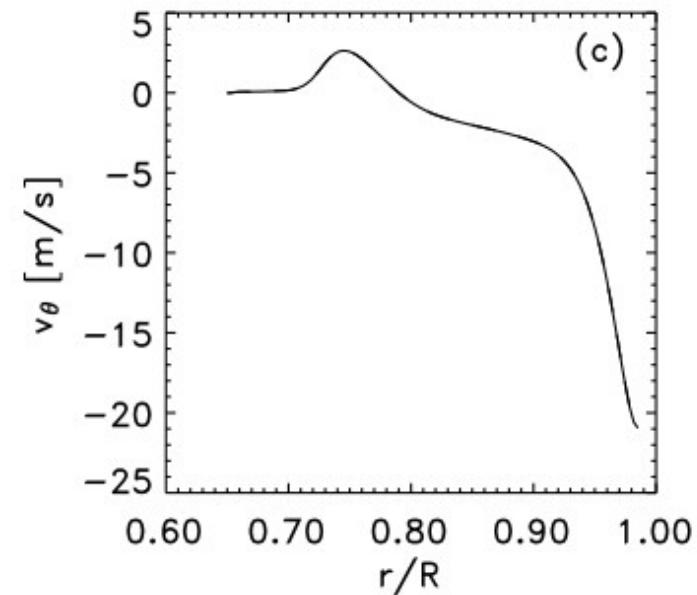
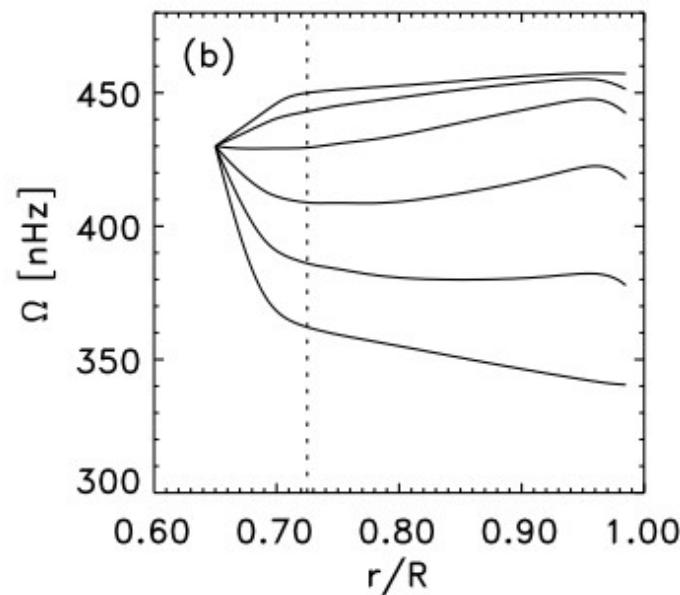
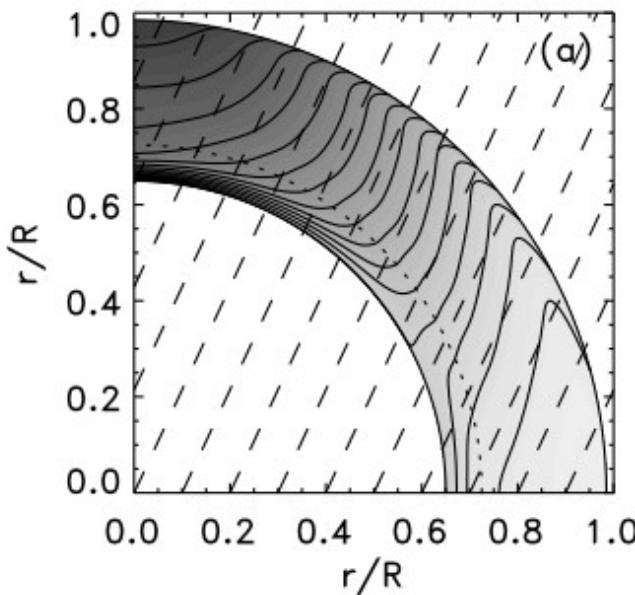
$$Q_{ij} = \langle u'_i u'_j \rangle \quad \text{(Kitchatinov & Rüdiger 93)}$$

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(Kitchatinov & Rudiger 95)



(Rempel 2005)

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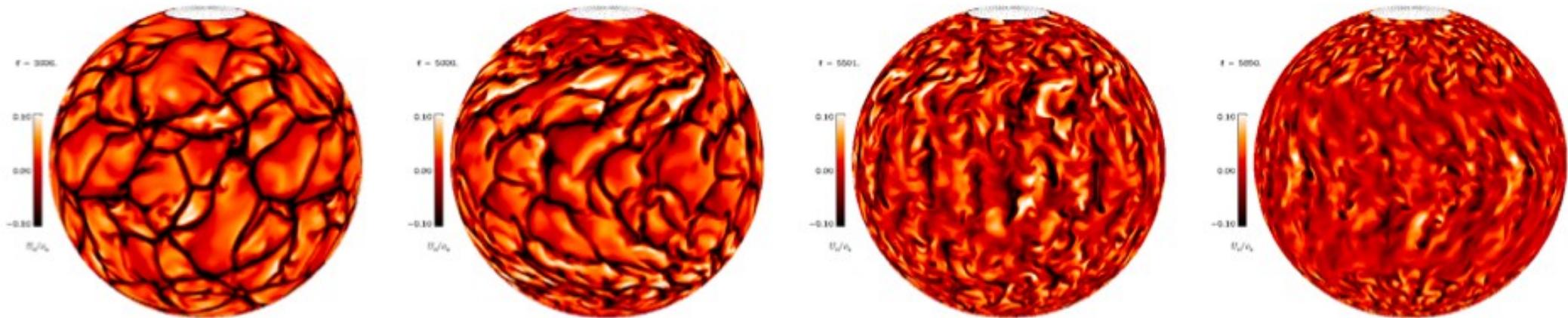
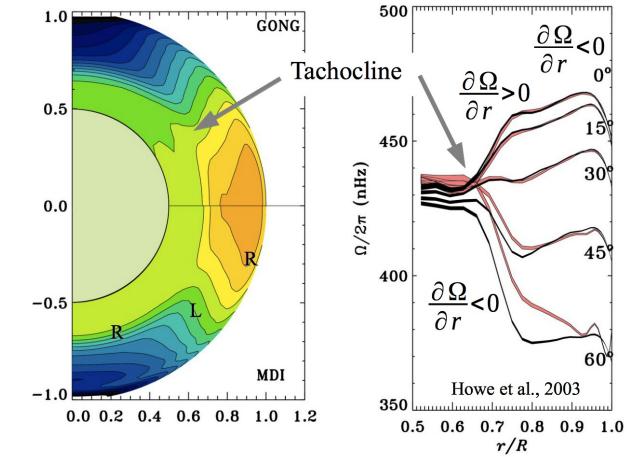


Direct Numerical Simulations (DNS)

Resolve scales down to the grid resolution. They could be global or local.

$$N^3 \geq \text{Re}^{9/4}$$

$$\text{Re} \simeq 10^{12} \dots 10^{13}$$



(Käpylä et al. 2011)

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Large eddy Simulations (LES)

Resolve scales down to the grid resolution,
the sub-grid scale (SGS) contribution is
considered via a turbulent model.

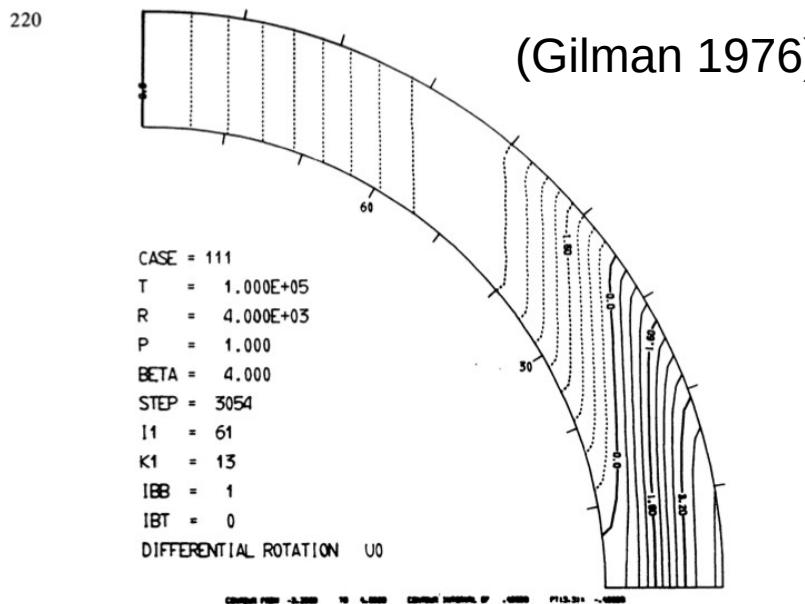
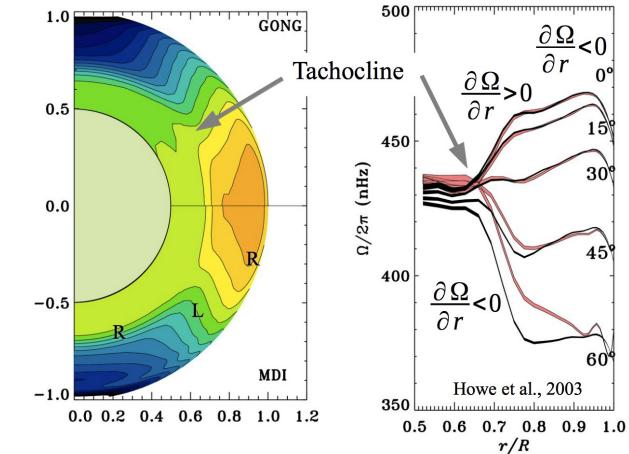
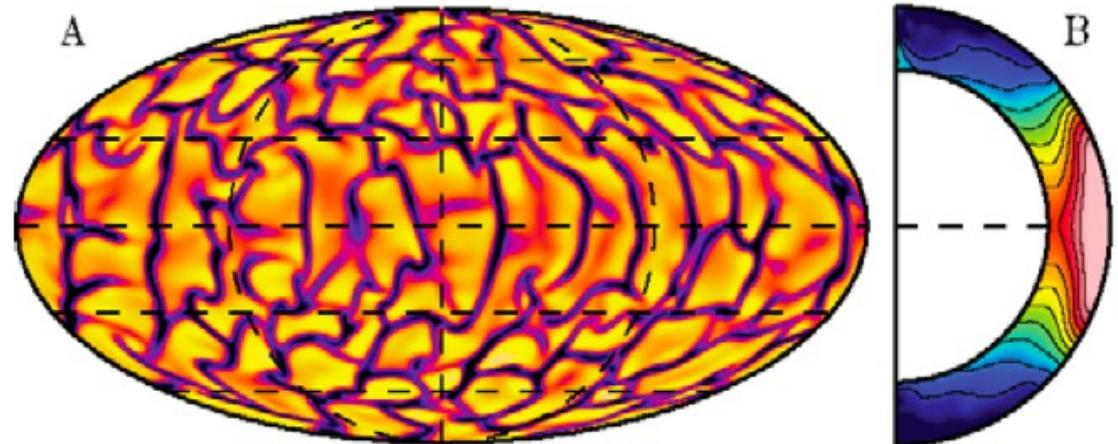


Fig. 7. Computer produced meridional cross section of linear rotation velocity u_0 relative to uniformly rotating frame, for Taylor number $T = 10^5$, Rayleigh number $R = 4 \times 10^3$, Prandtl number $P = 1$. Units dimensionless, with velocity scaled by κ/d , in which κ is thermal diffusivity, d is convection zone depth.
Positive u_0 indicated by solid contours; negative u_0 by dashed contours.

$$\nu_{sgs} = \text{cnt} \gg \nu$$

- ASH code, e.g. Miesch et al. (2011)
 - Anelastic approximations, LES.

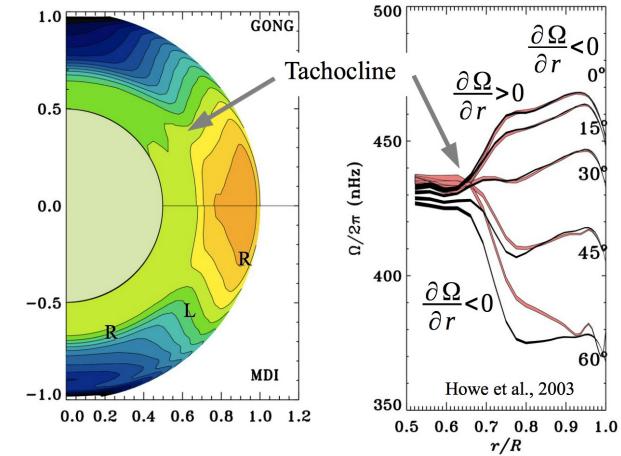
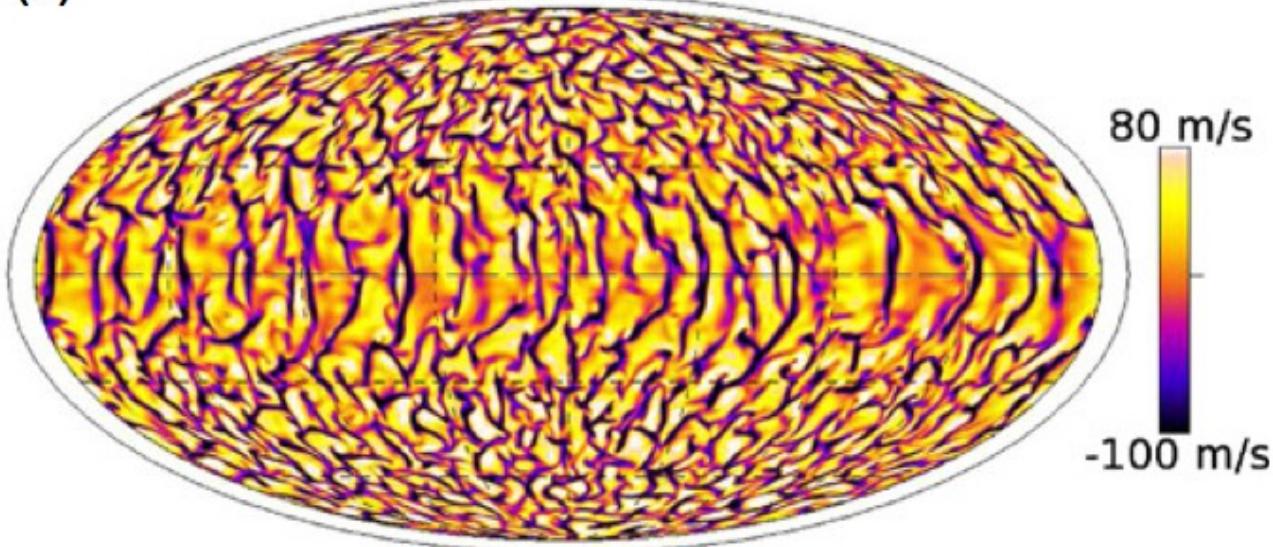


$$\nu_{SGS} \propto \rho^{-1}(\rho^{-1/2}) \gg \nu$$

(c)

D3b

(Nelson et al. 2013)



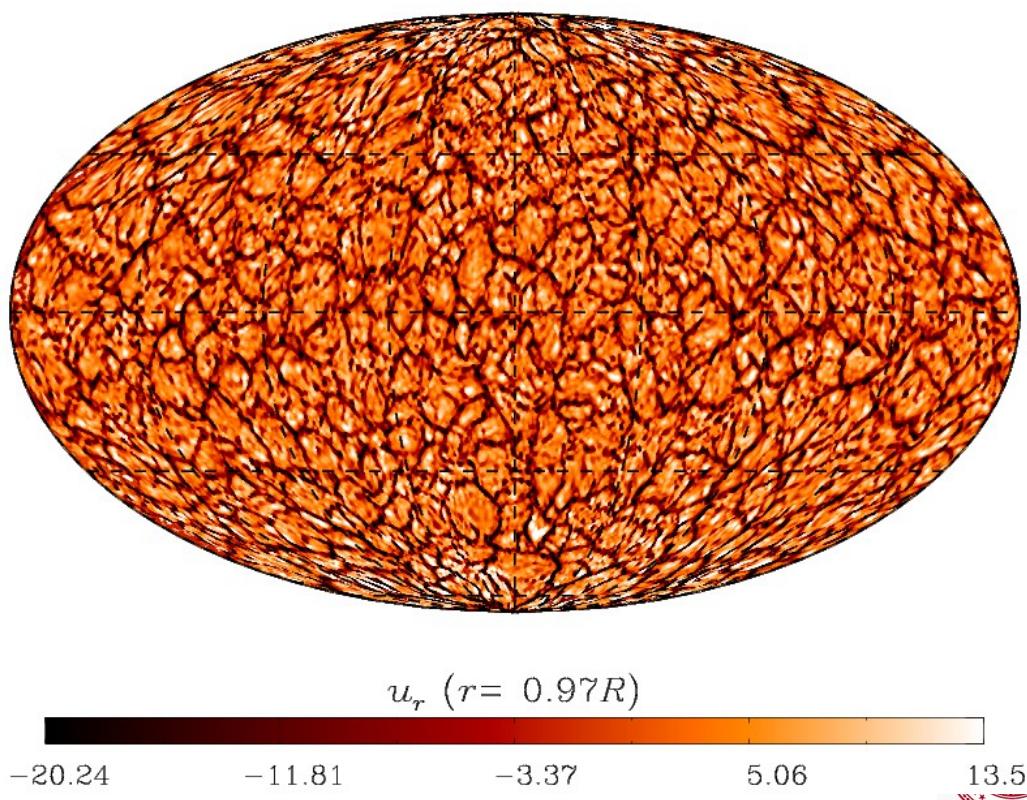
$$\nu_{SGS} = C_s \Delta (\pi_{ij} \pi_{ij})^2$$

C_s = ctn (Smagorinsky)

$C_s = Cs(\mathbf{U})$ (Dynamic Smagorinsky)

Implicit sub-grid scale modeling (ISGS, ILES): the non-linear computation of the truncation error (i.e., numerical viscosity) is identified with the SGS contribution (EULAG code),
(Margolin & Rider, 2002)

(Guerrero et al. 2013)



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-20.24

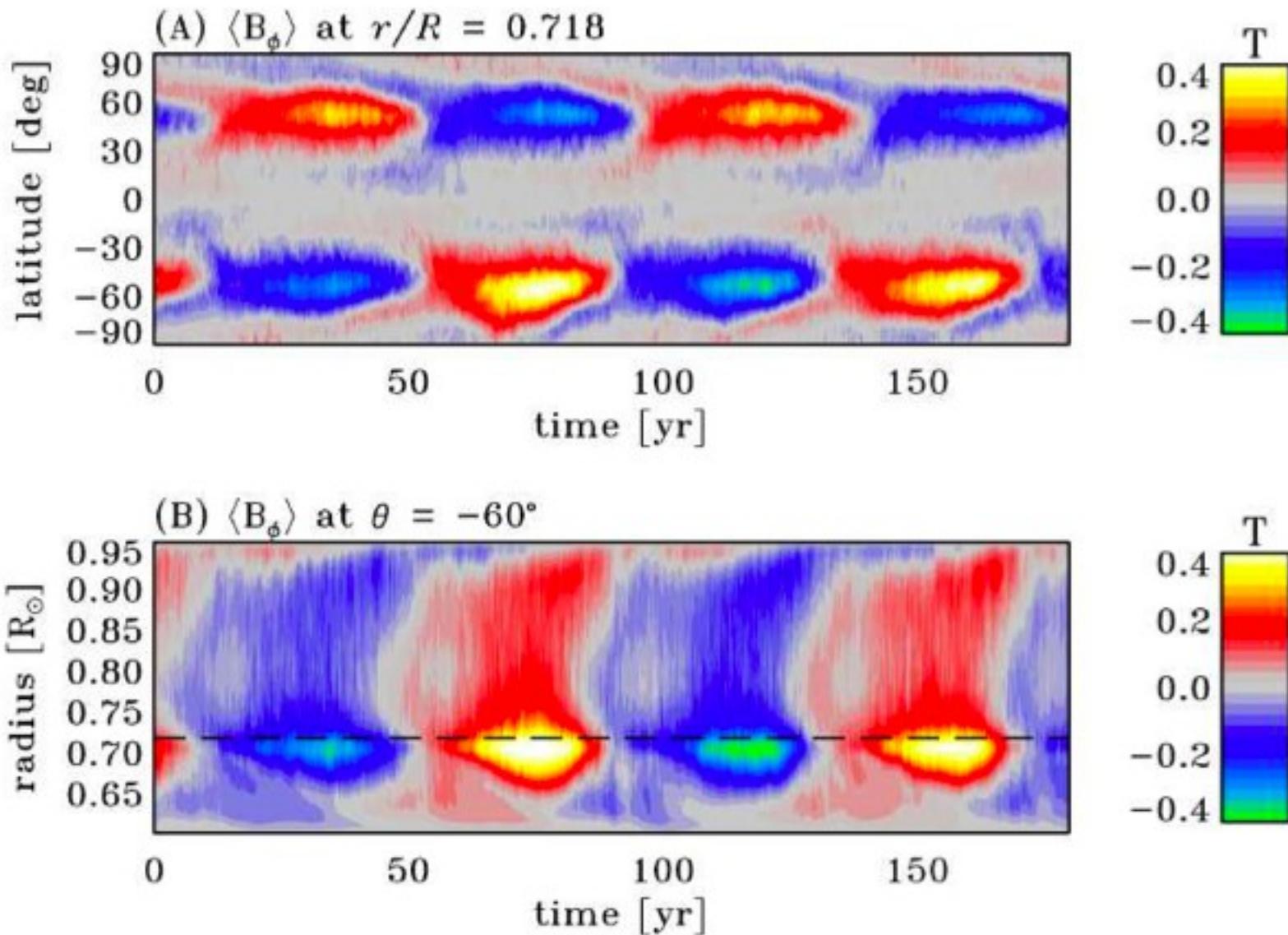
-11.81

-3.37

5.06

13.50





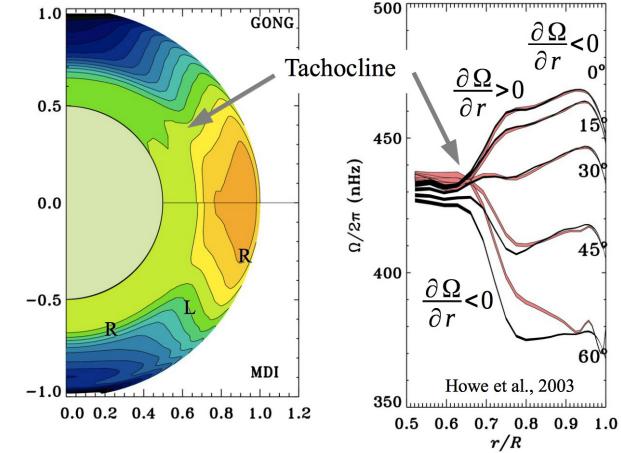
Ghizaru, Charbonneau & Smolarkiewicz (2010)

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3. Anelastic simulations with EULAG code

Anelastic approximation
(Lipps & Helmer 1982, Lipps 1990)



$$\nabla \cdot (\rho_0 \mathbf{u}) = 0, \quad (1)$$

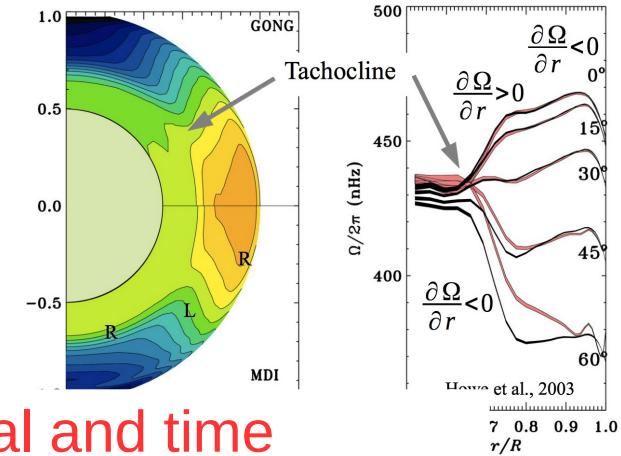
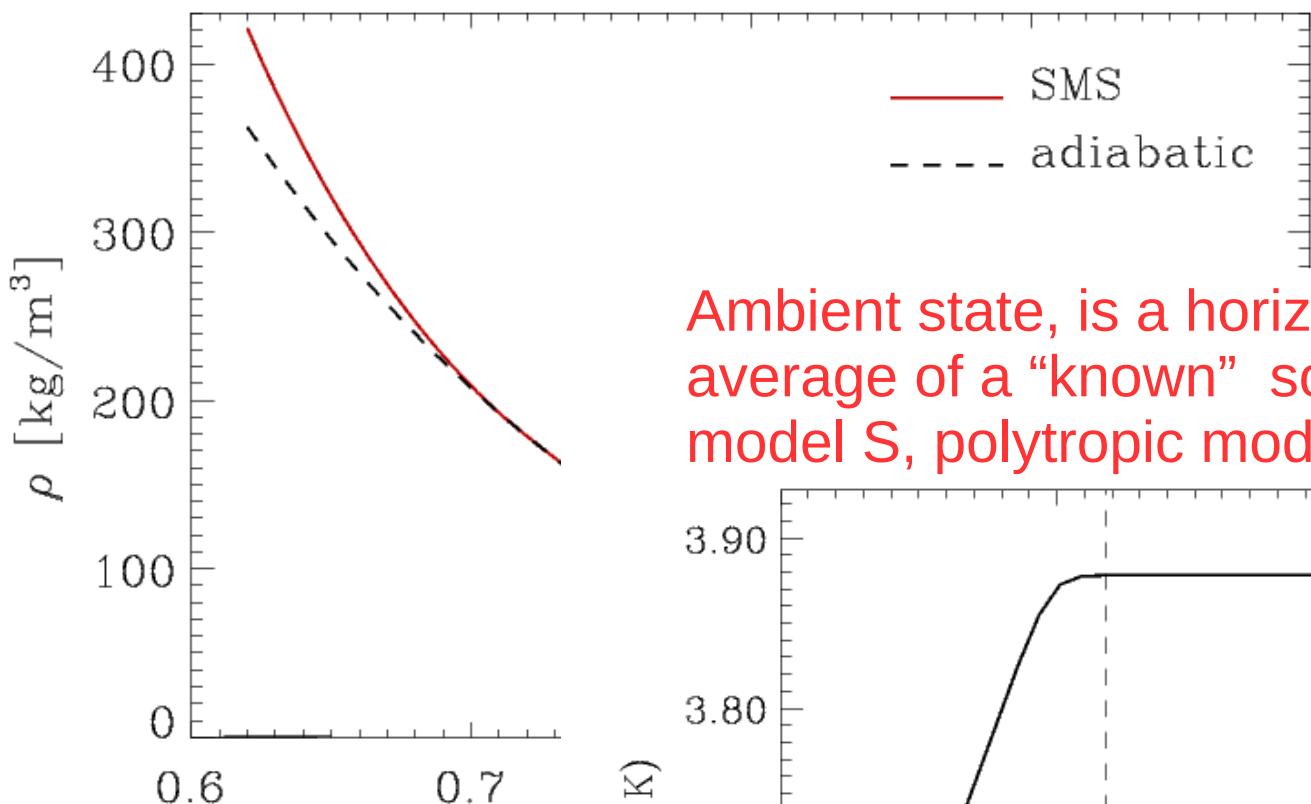
$$\frac{D\mathbf{u}}{Dt} - 2\Omega \times \mathbf{u} = -\nabla \left(\frac{p'}{\rho_0} \right) + \mathbf{g} \frac{\Theta'}{\Theta_0} + \frac{\mathbf{F}}{\rho_0}, \quad (2)$$

$$\frac{D\Theta'}{Dt} = -\mathbf{u} \cdot \nabla \Theta_e + \frac{1}{\rho_0} \mathcal{H}(\Theta') - \alpha \Theta' \quad (3)$$

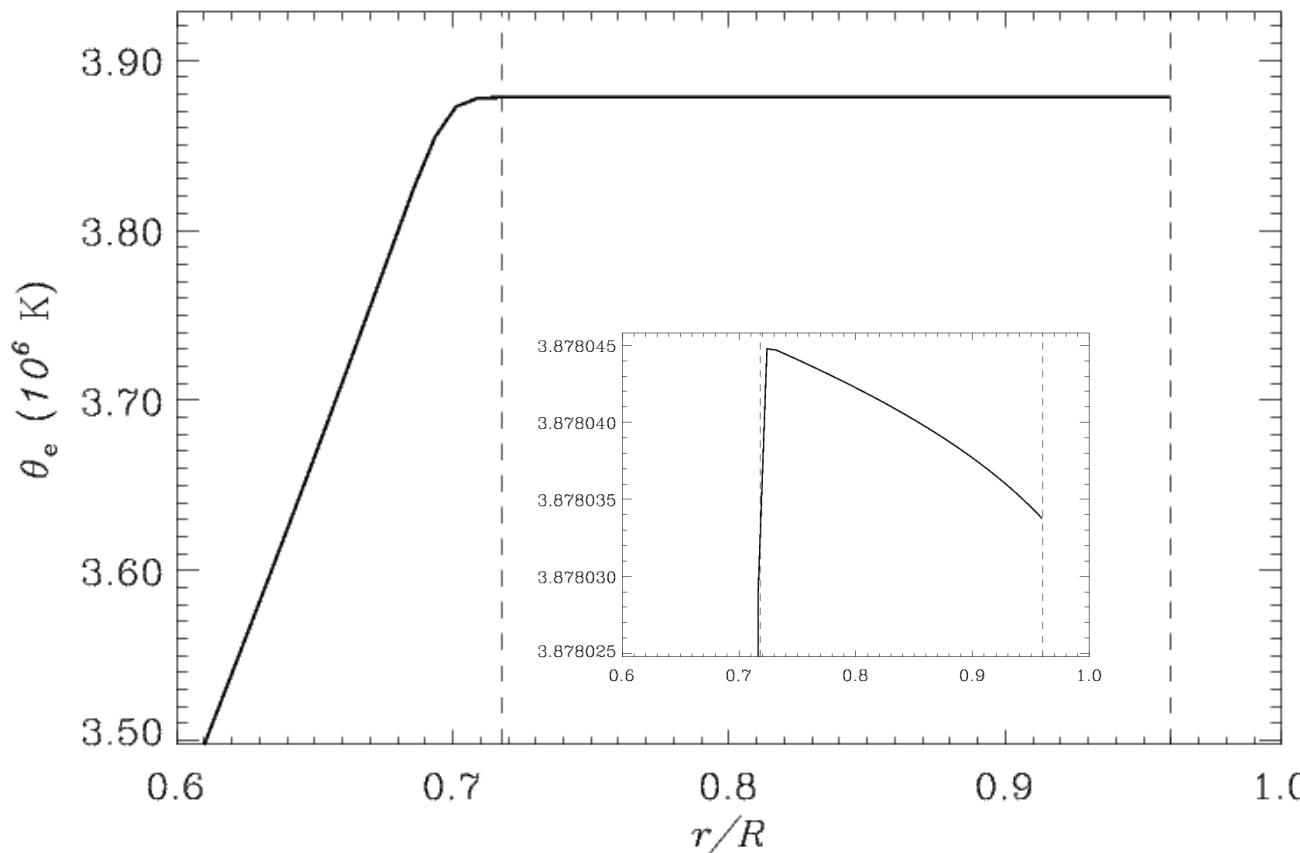
$$ds = c_p d \ln \Theta \quad (4)$$



The background state, ρ_0, p_0, Θ_0

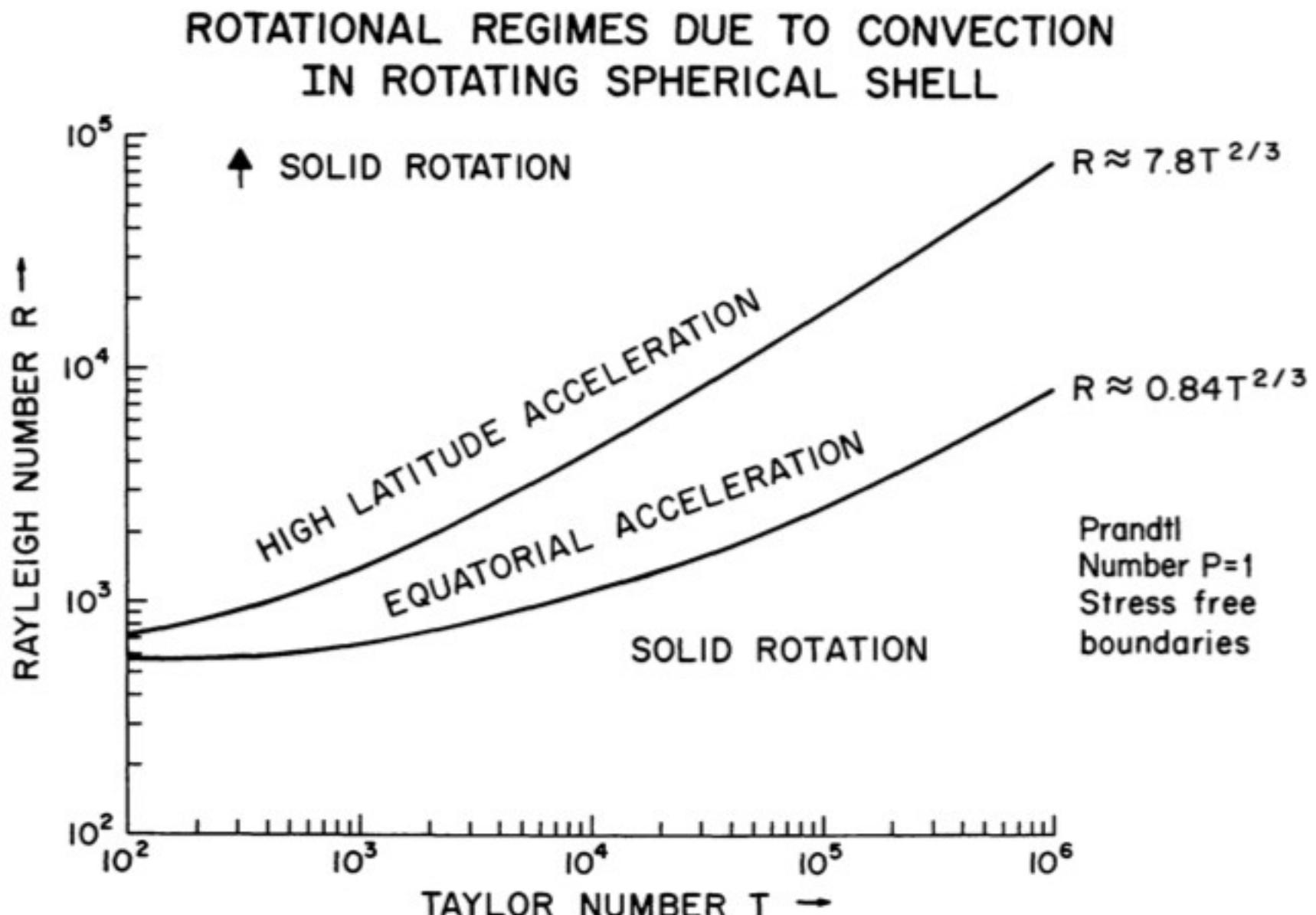
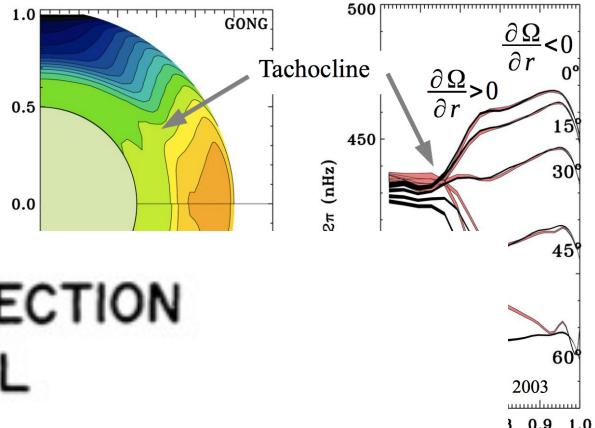


Ambient state, is a horizontal and time average of a “known” solution, e.g., solar model S, polytropic model.



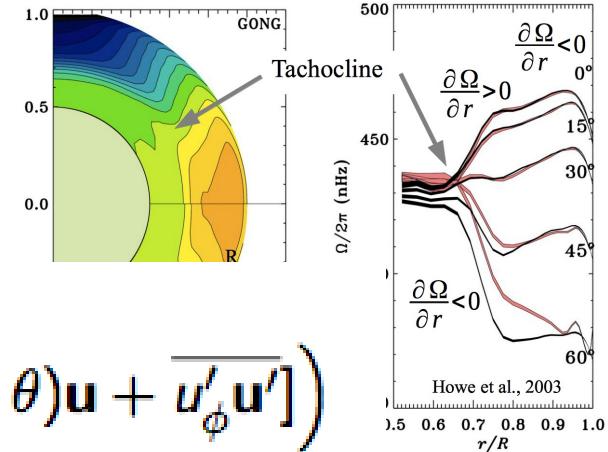
a) Angular momentum transport

Rotation vs. Convection



Angular momentum conservation

$$\frac{\partial(\overline{\rho u}_\phi)}{\partial t} = \frac{1}{r \sin \theta} \nabla \cdot (\overline{\rho r \sin \theta} [-\nu \nabla \overline{u}_\phi + (\overline{u}_\phi + \Omega_0 r \sin \theta) \mathbf{u} + \overline{u'_\phi u'}])$$



- Viscous flux: $F_{vis} = -\nu \nabla \overline{u}_\phi$
- Meridinal circulation: $F_{MC} = (\overline{u}_\phi + \Omega_0 r \sin \theta) \mathbf{u}$
- Reynold stresses flux: $F_{RS} = \overline{u'_\phi u'}$

$$(\lambda, \phi, z)$$

$$\lambda = r \sin \theta$$

$$z = r \cos \theta$$

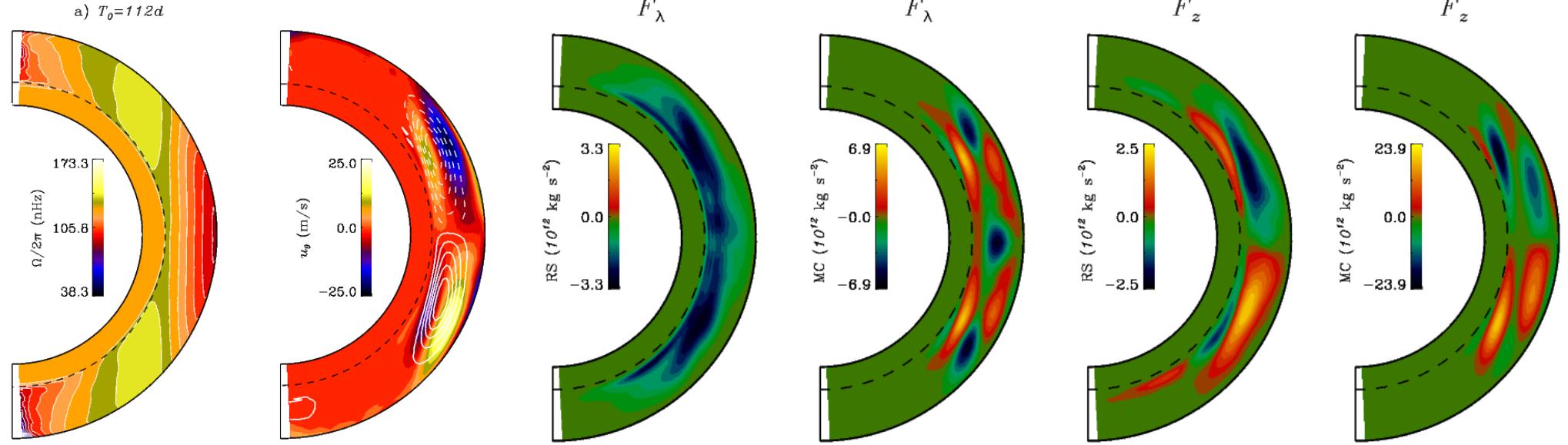
$$F_\lambda \hat{\lambda} = F_r \sin \theta \hat{r} + F_\theta \cos \theta \hat{\theta}$$

$$F_z \hat{z} = F_r \cos \theta \hat{r} - F_\theta \sin \theta \hat{\theta}$$

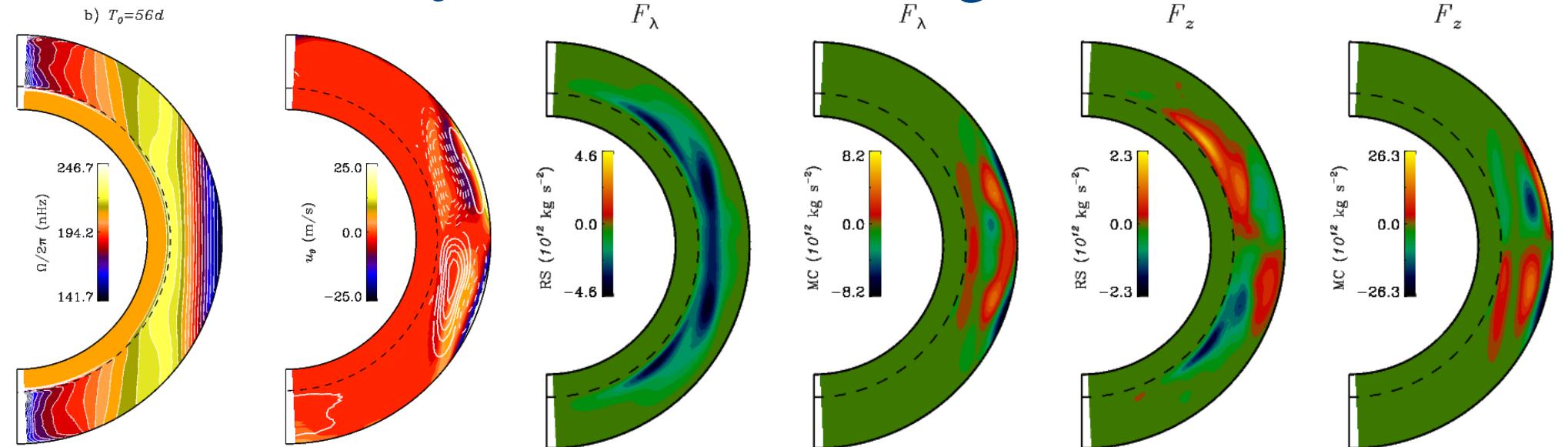
$$Q_{\phi r}, Q_{\phi \theta}, Q_{r \theta}$$



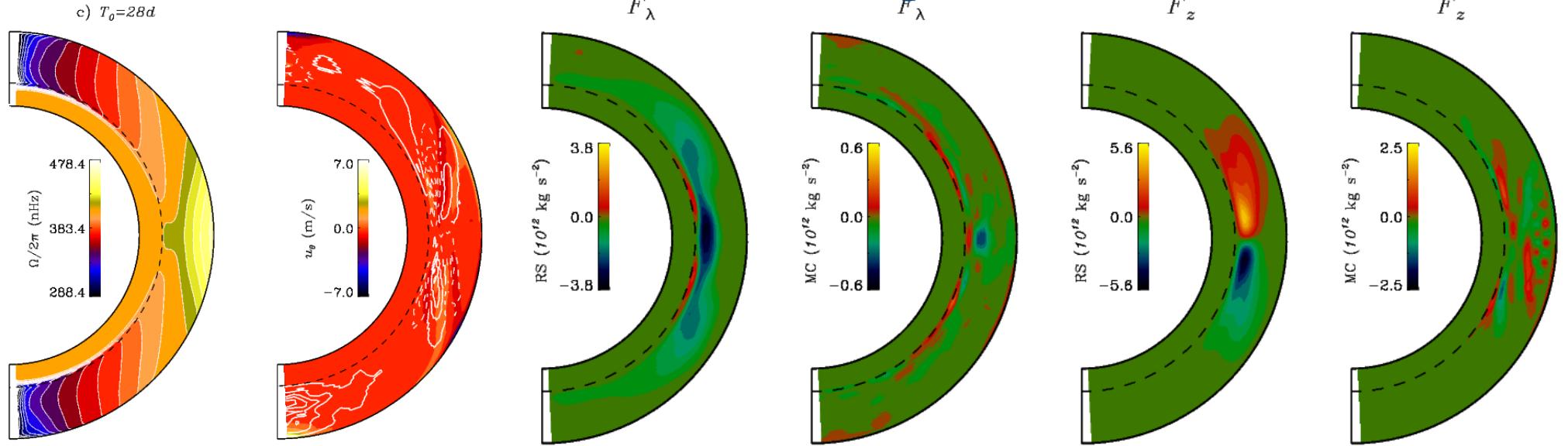
a) $\Omega_o=112\text{days}$ ($1/4 \Omega_\odot$, $Re \sim 60$)



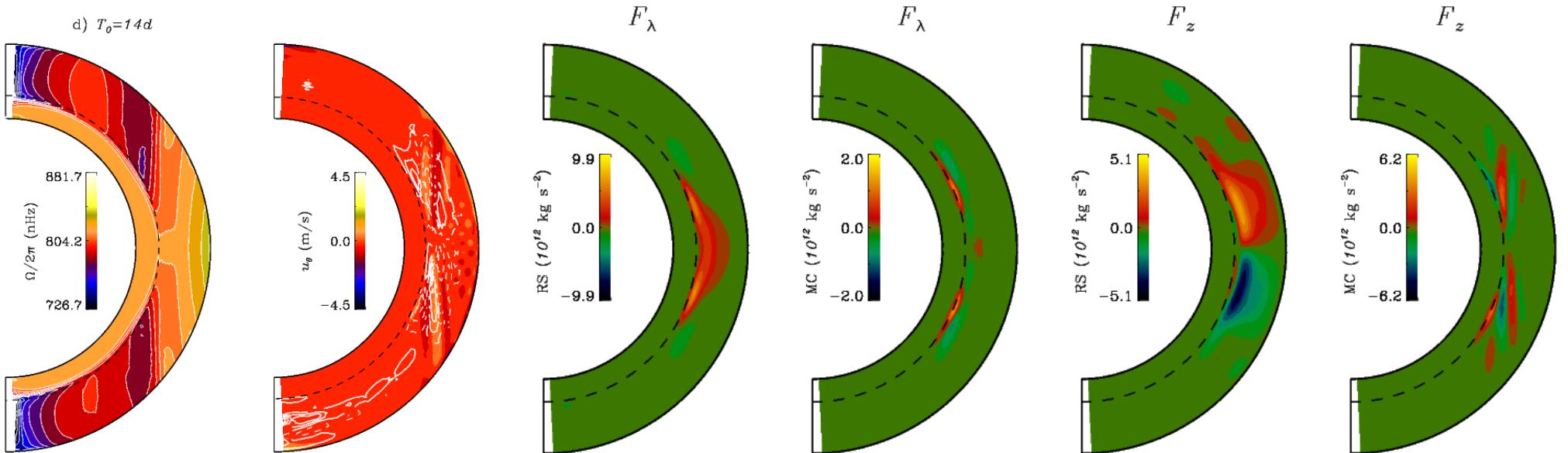
b) $\Omega_o=56\text{days}$ ($1/2 \Omega_\odot$, $Re \sim 60$)

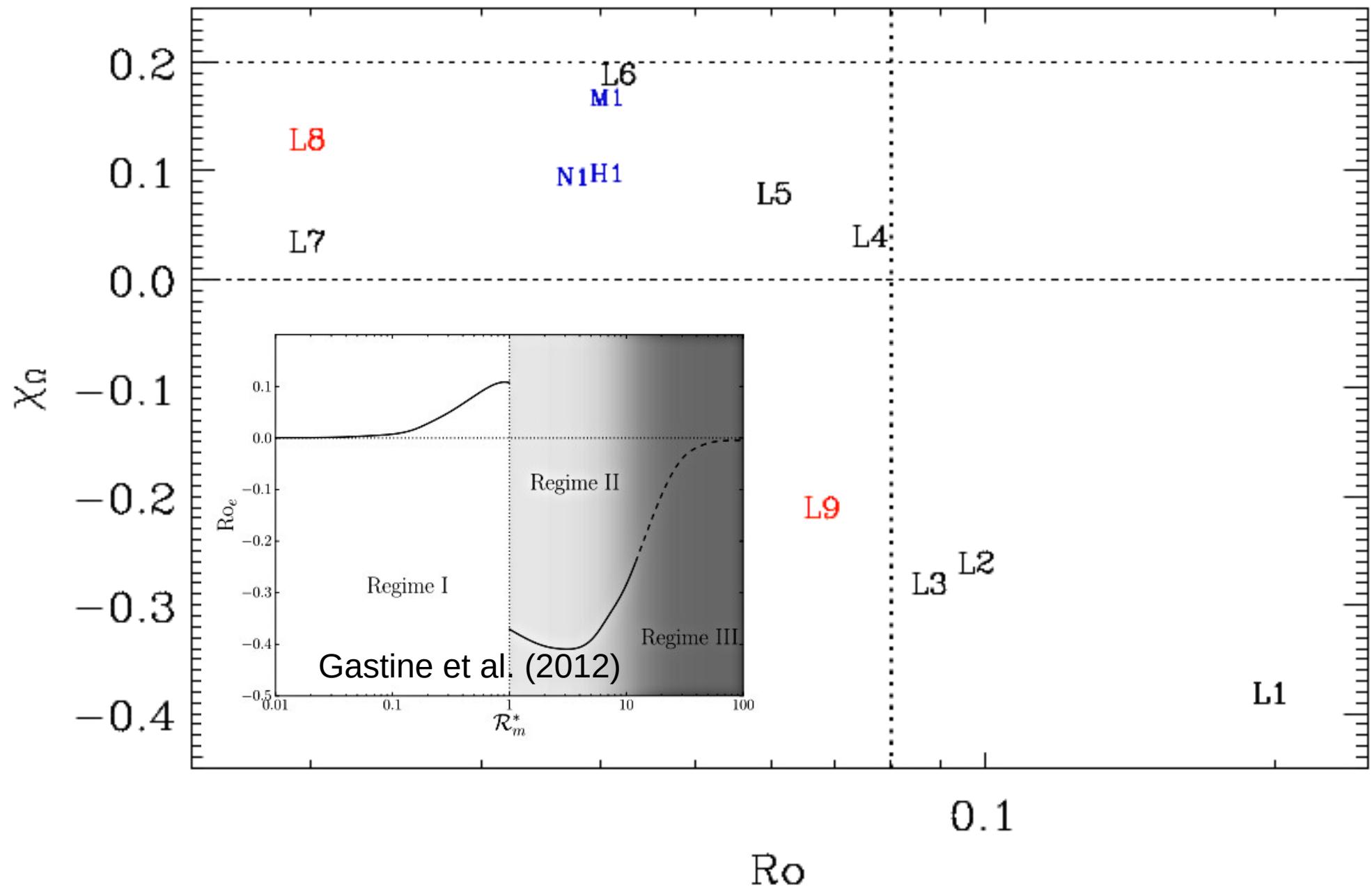


c) $\Omega_\theta=28\text{days}$ (Ω_\odot , $Re \sim 60$)



d) $\Omega_\theta=14\text{days}$ ($2\Omega_\odot$, $Re \sim 60$)



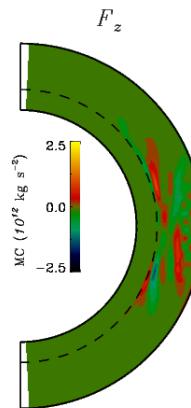
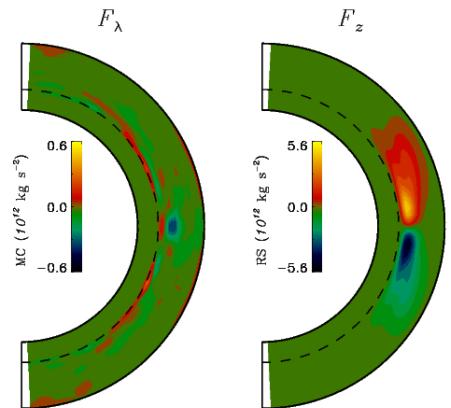
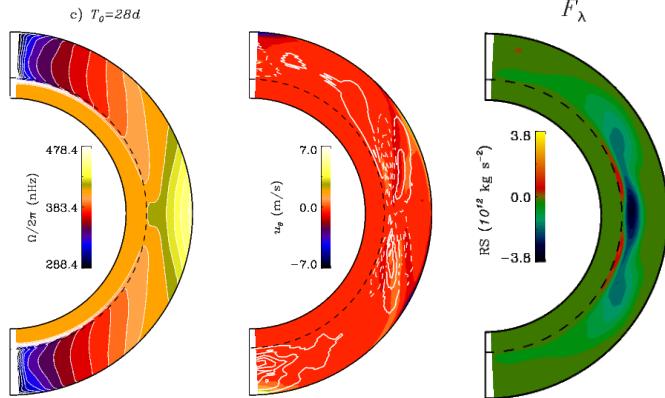


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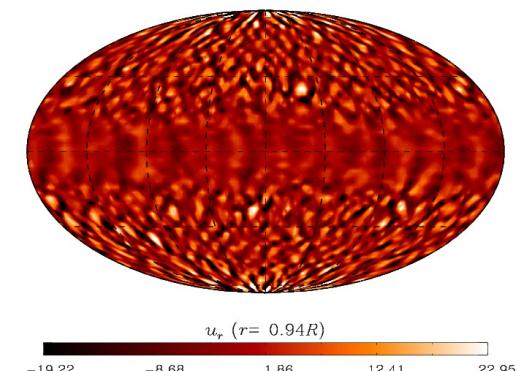


b) Convergence: $x=128 \times 64 \times 47$

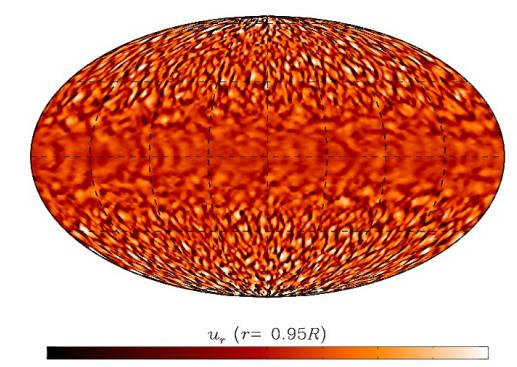
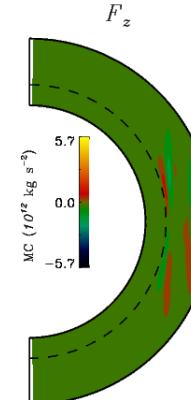
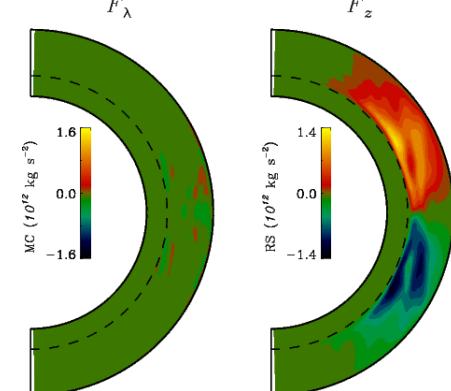
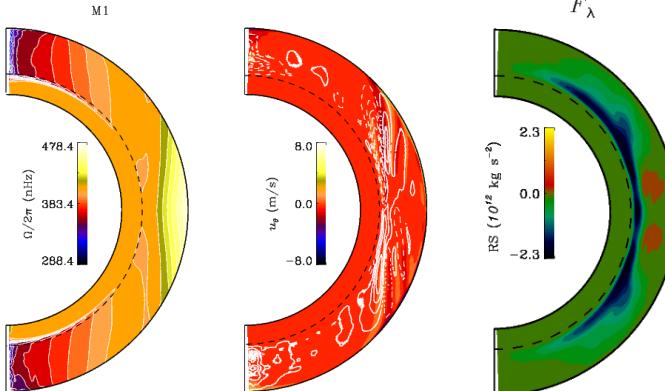
1x



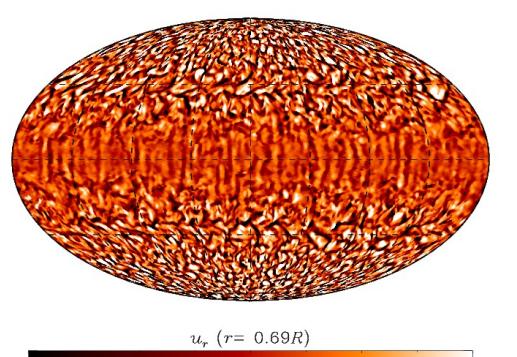
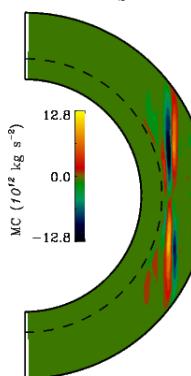
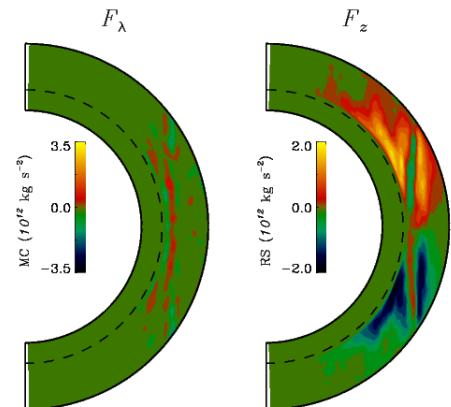
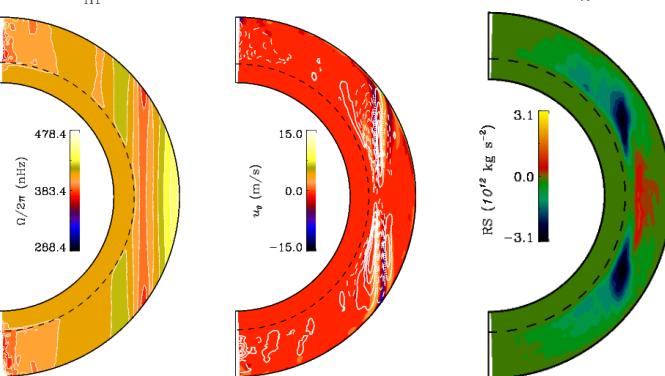
Vertical velocity



2x

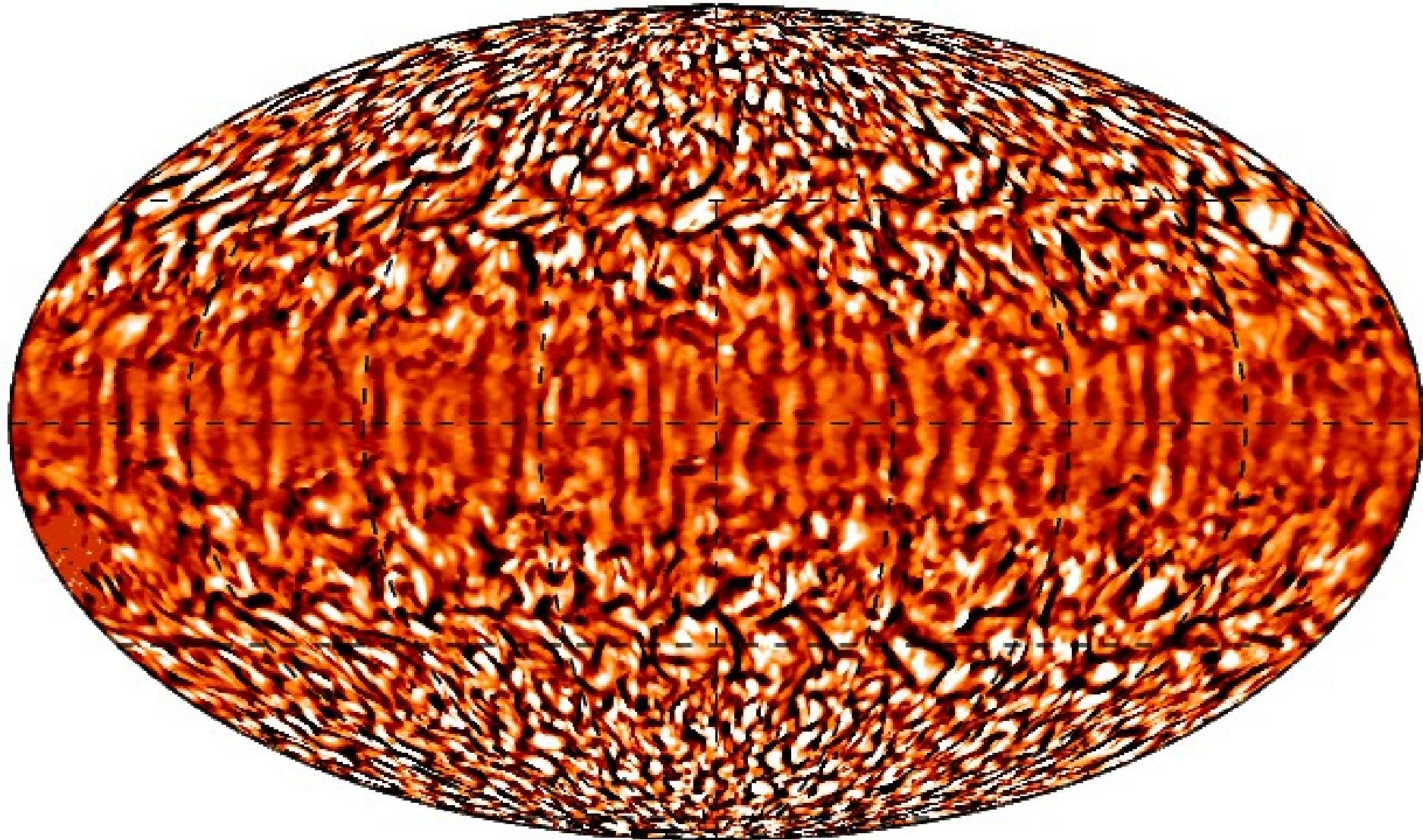


4x



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$u_r \ (r= 0.95R)$

-24.00

-13.25

-2.50

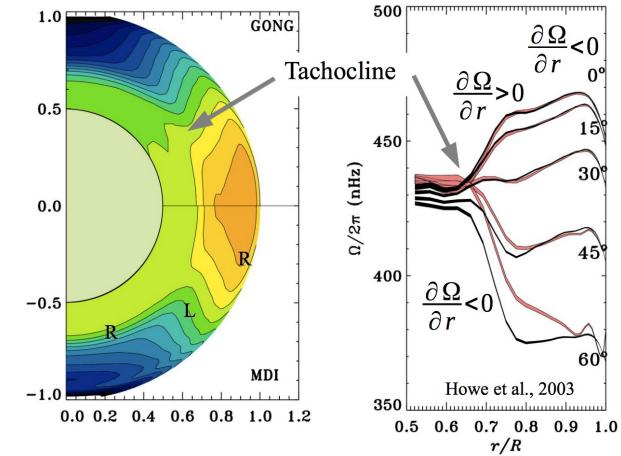
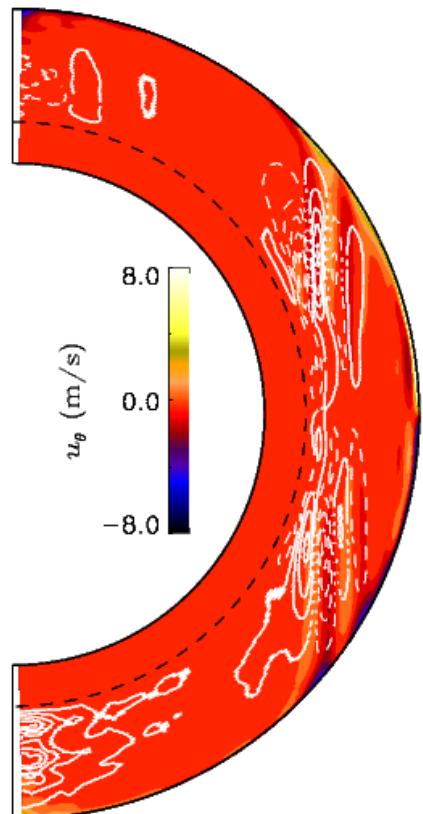
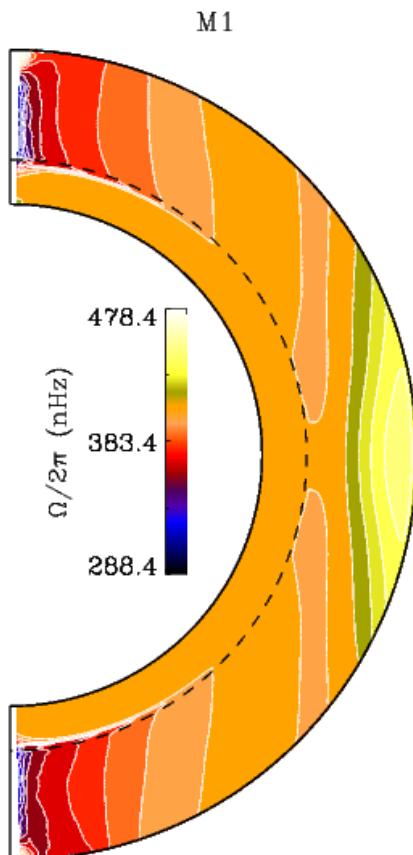
8.25

19.00

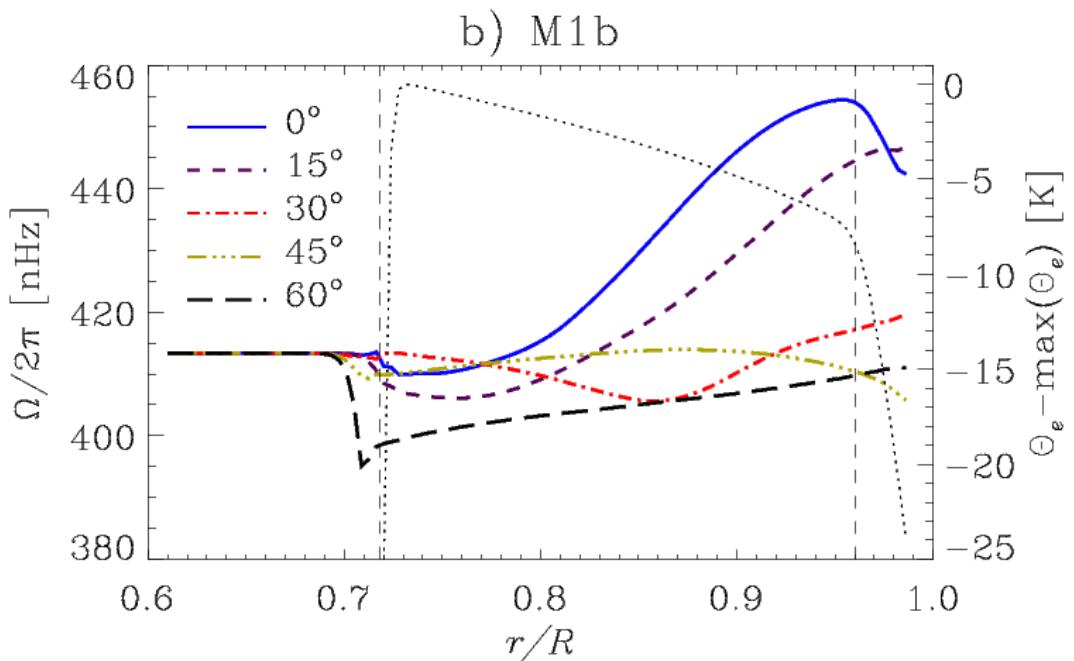


c) Near-surface shear layer

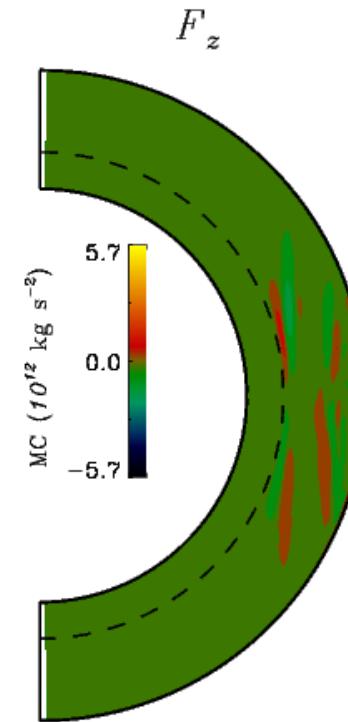
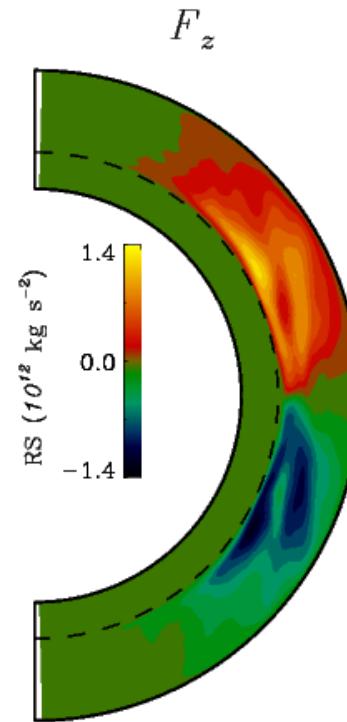
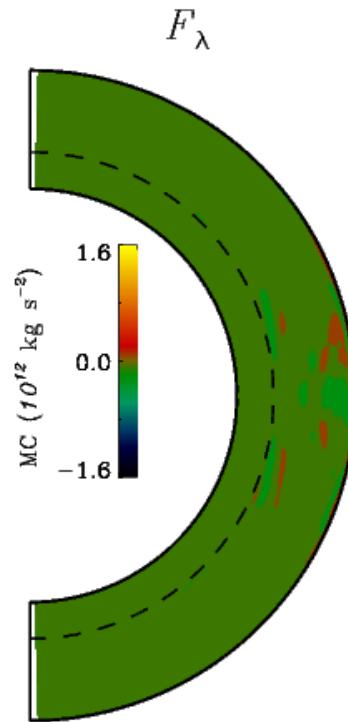
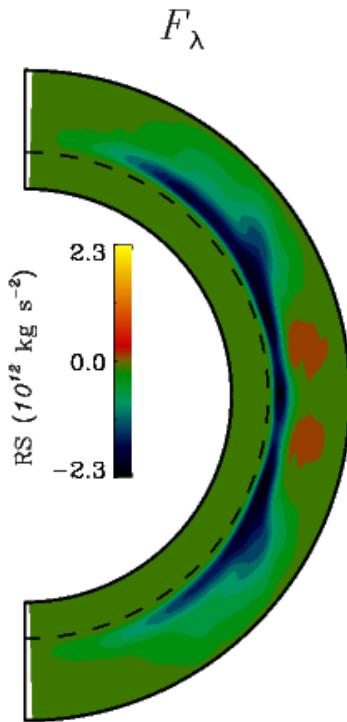
Based in models slow rotating
models where convection
dominates over rotation ...



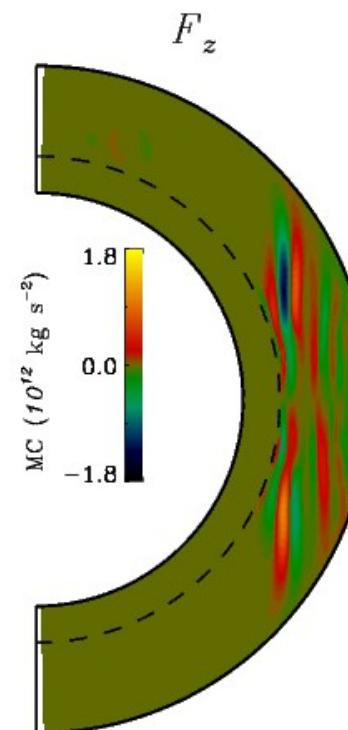
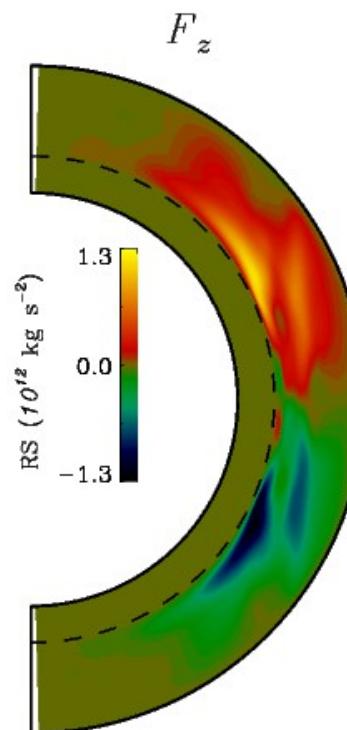
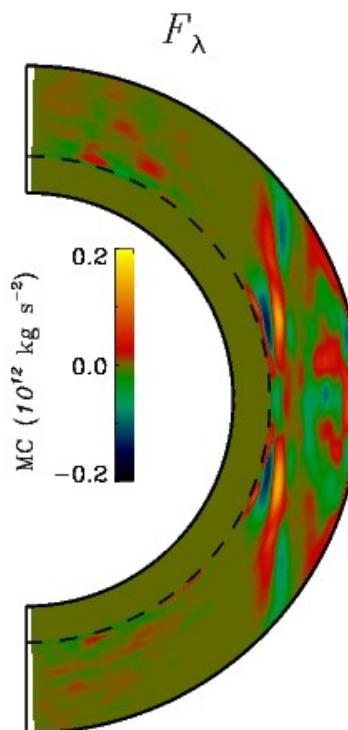
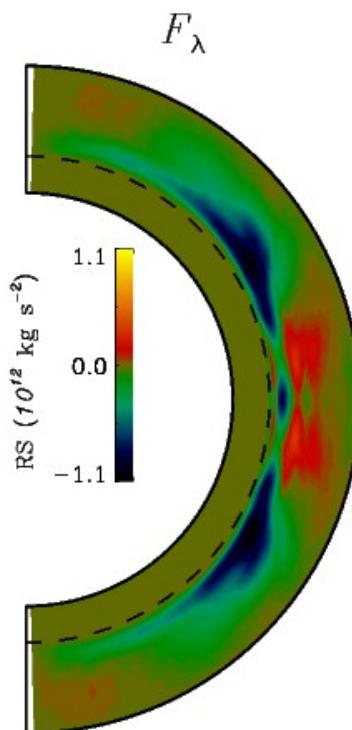
Negative radial shear appears at lower latitudes



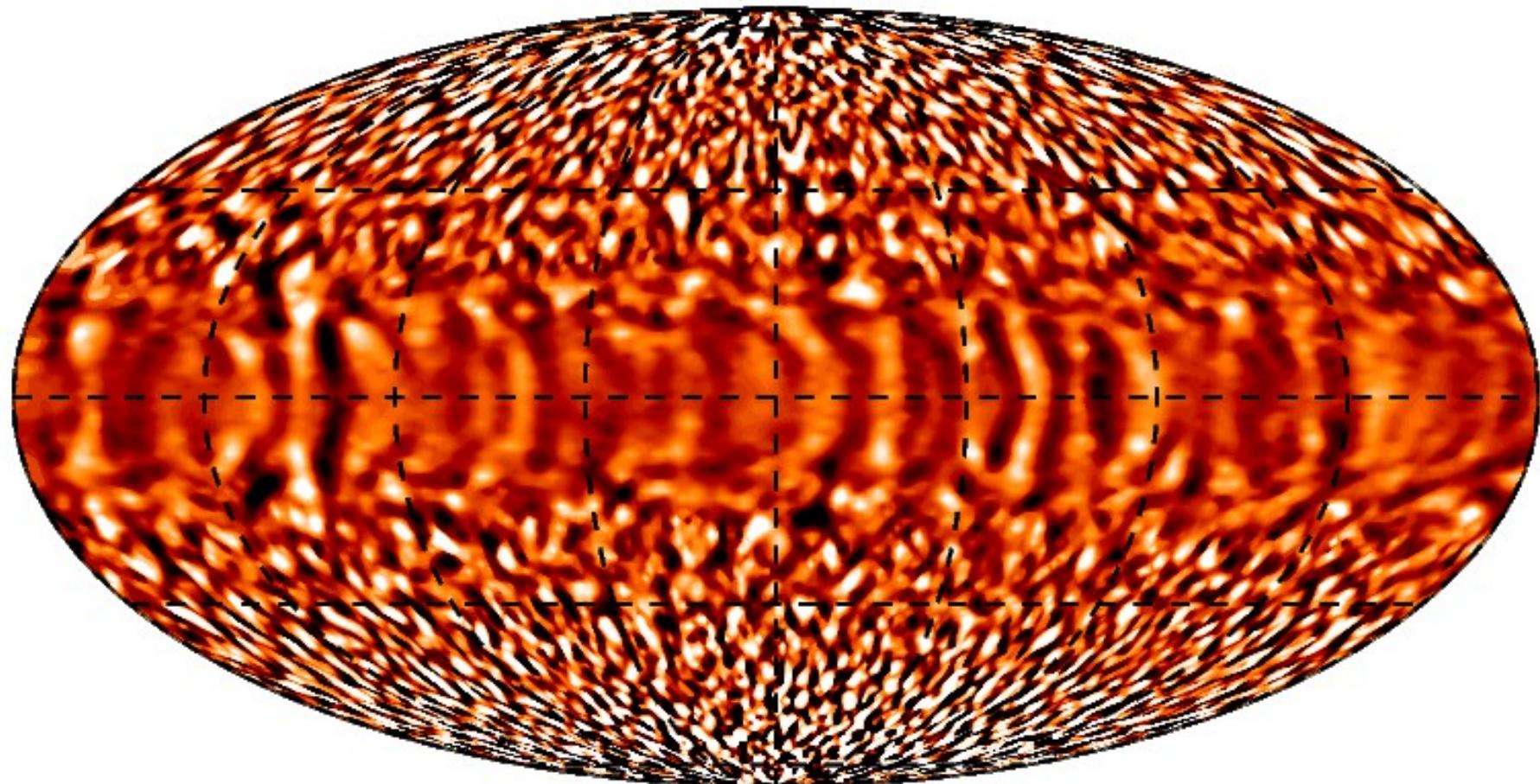
2x



2x-nssl



a) N1



u_r ($r = 0.94R_s$)

-15.64

-8.60

-1.56

5.47

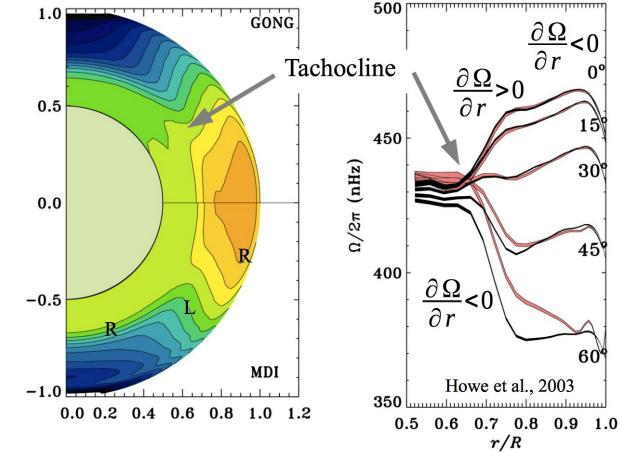
12.51

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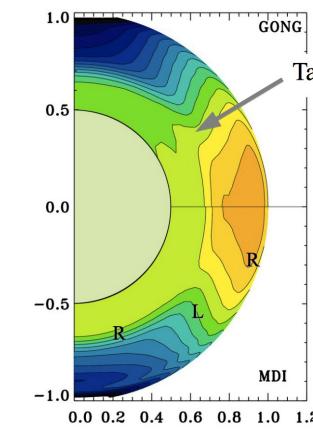
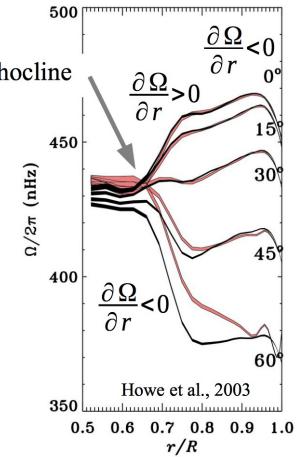
CONCLUSIONS

- EULAG code is a powerful tool to study solar/stellar rotation and dynamo, high Re are obtained for relatively low resolutions



- The resulting rotation profile depends on the balance between Coriolis and buoyancy forces. For the stratification considered here solar-like rotation (fast equator) is obtained at the solar rotation rate
- Due to the highly sub-adiabatic stable region, all the simulations exhibit the formation of a tachocline
- Taylor-Proudman columnar rotation profiles are persistent, indicating deficient latitudinal turbulent heat transfer





- For all the simulations at the solar rotation rate the meridional flow is multicellular
- Vigorous convection at the upper part of the convection zone leads to the formation of a “near-surface shear layer” at equatorial latitudes

