



Mean field models of Meridional flows and differential rotation in stellar convection zones

Manfred Küker

Leibniz-Institut für Astrophysik Potsdam (AIP)

Günther Rüdiger, Leonid Kitchatinov

Flux transport dynamo

- differential rotation
- α effect (helicity)
- diffusion
- meridional flow acts as conveyor belt
- cycle time determined by flow speed!

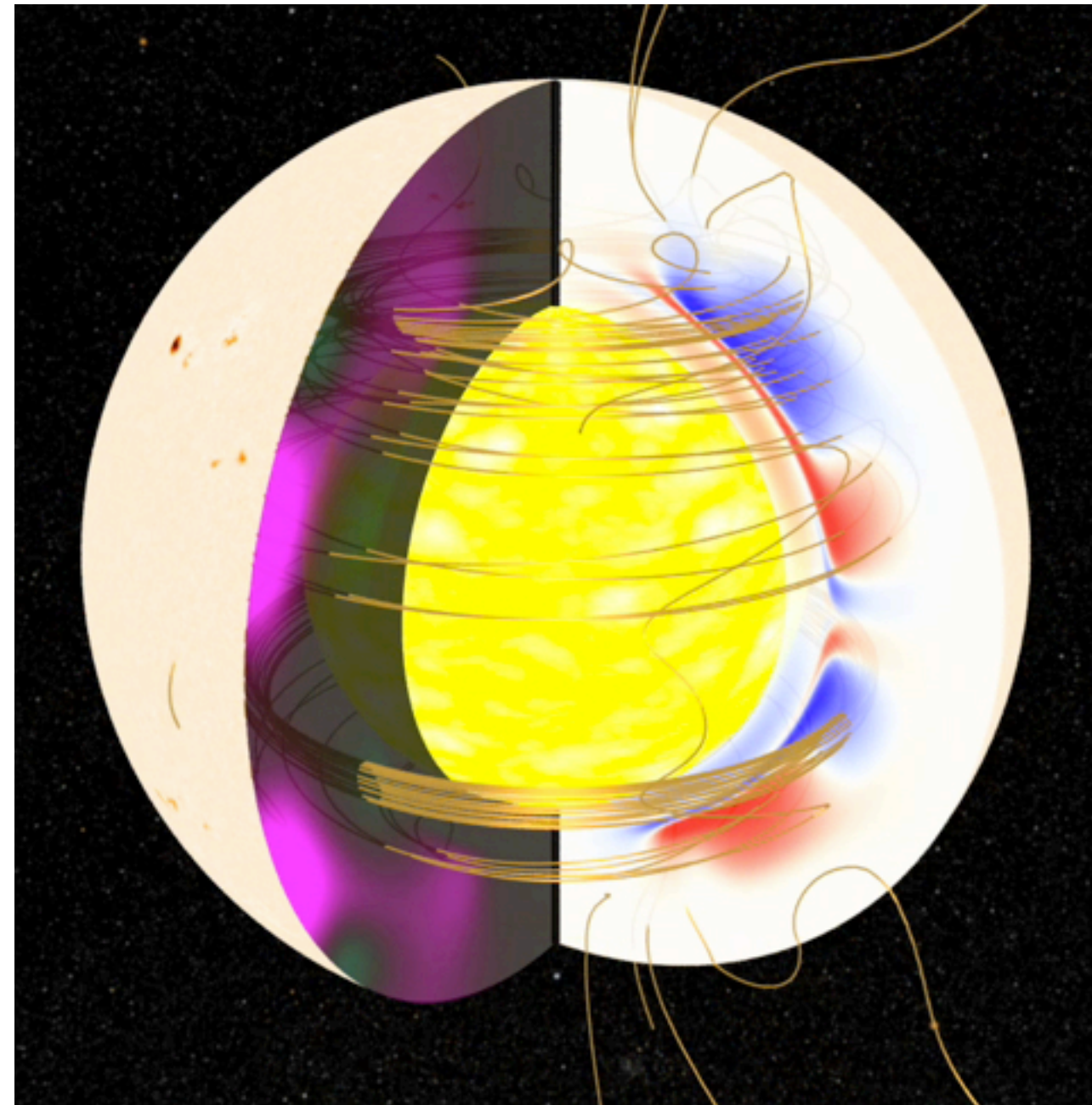
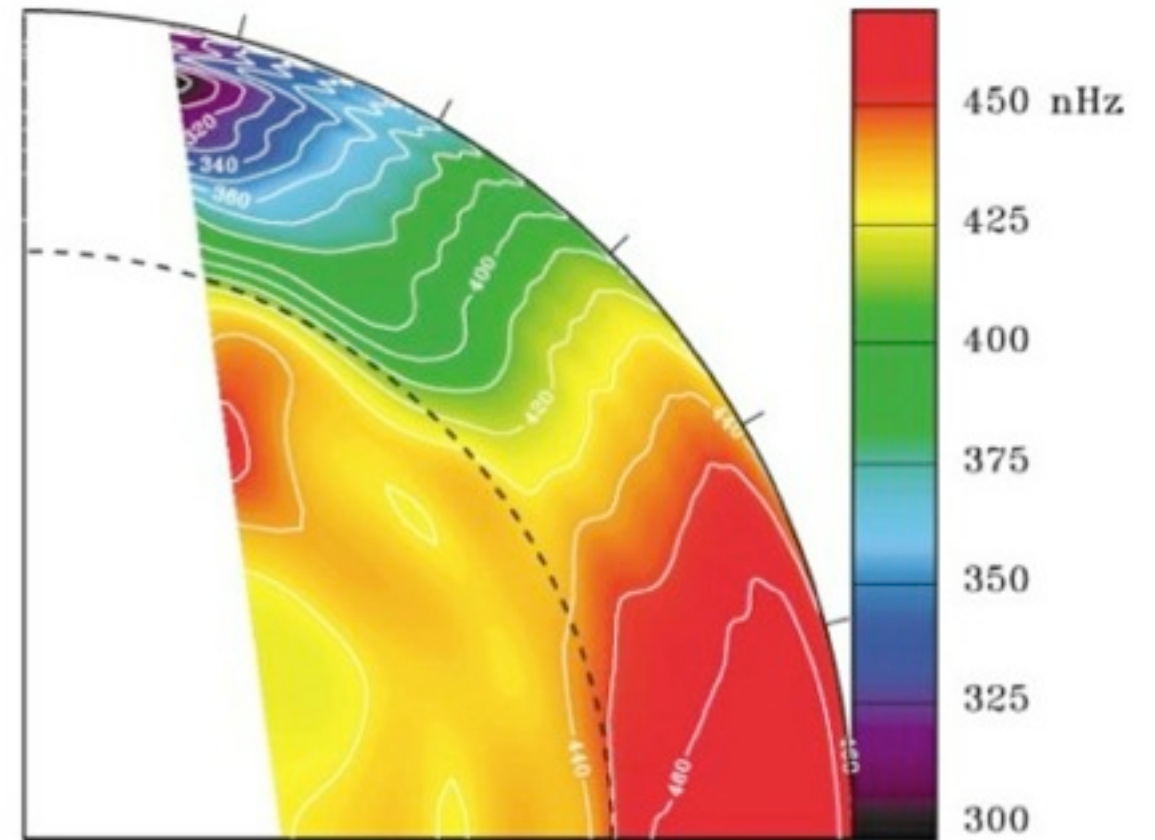
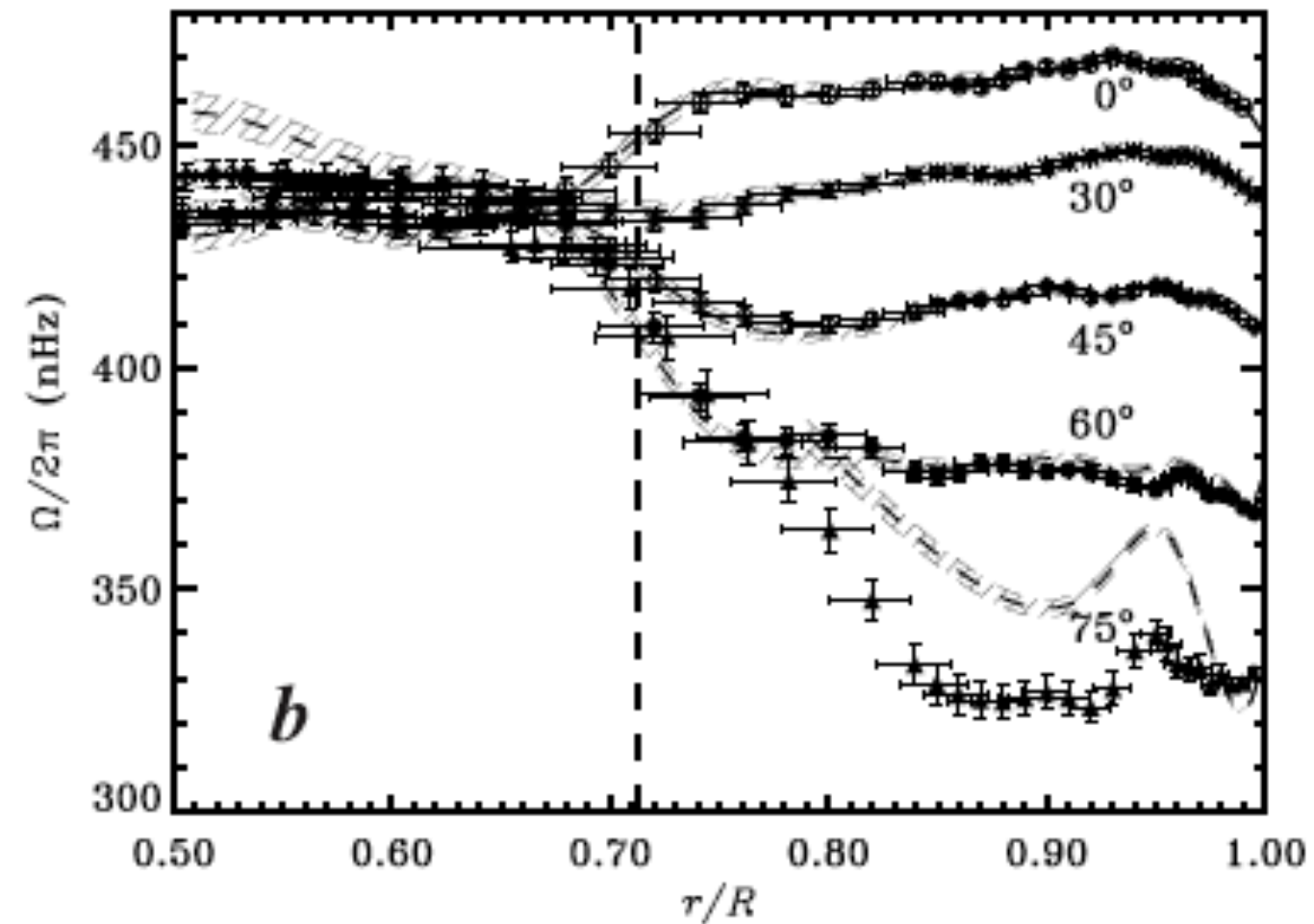


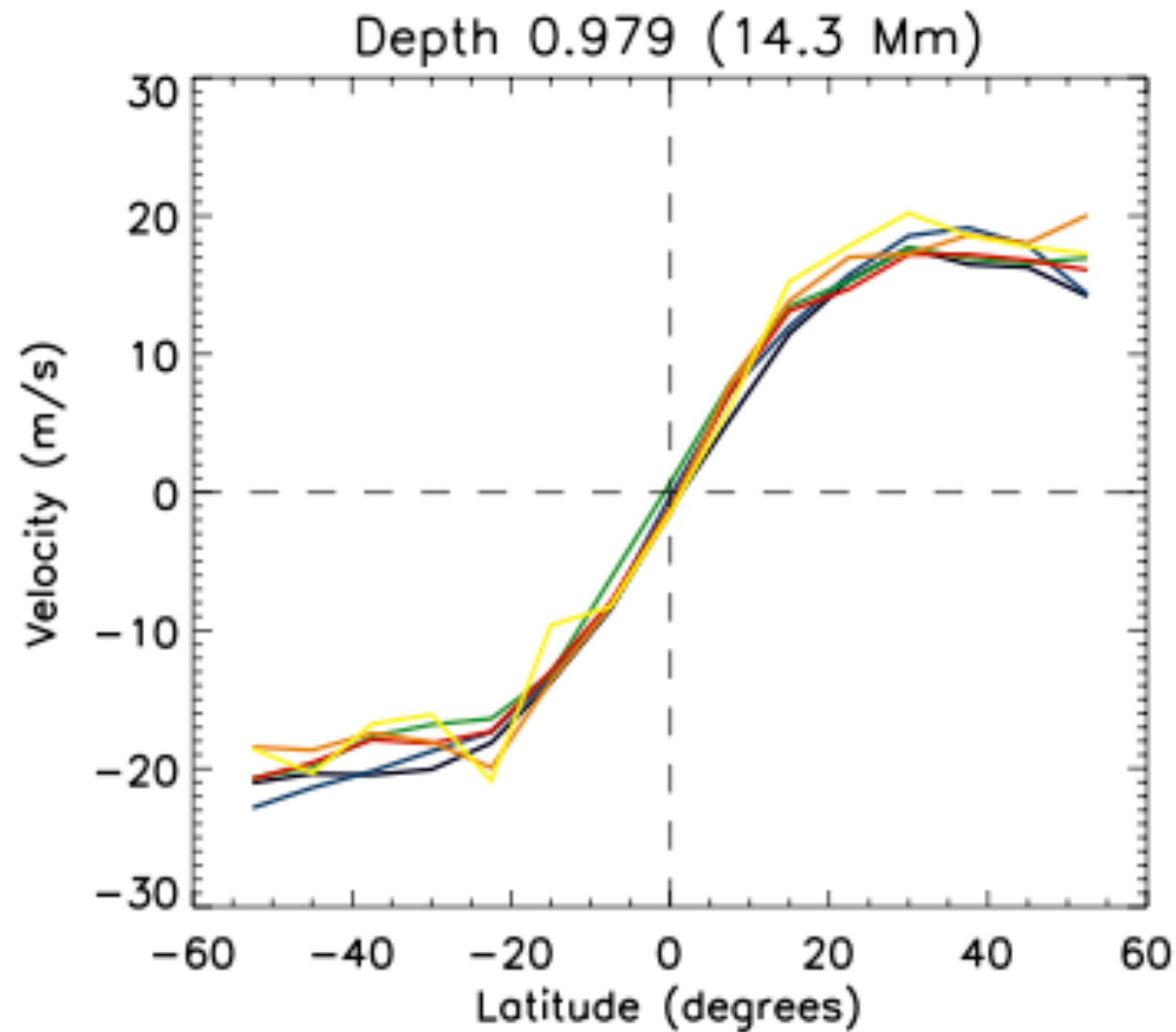
Image: NASA

Solar differential rotation: helioseismology



Thompson et al. 2003

Solar meridional flow



Gonzalez-Hernandez et al. 2008

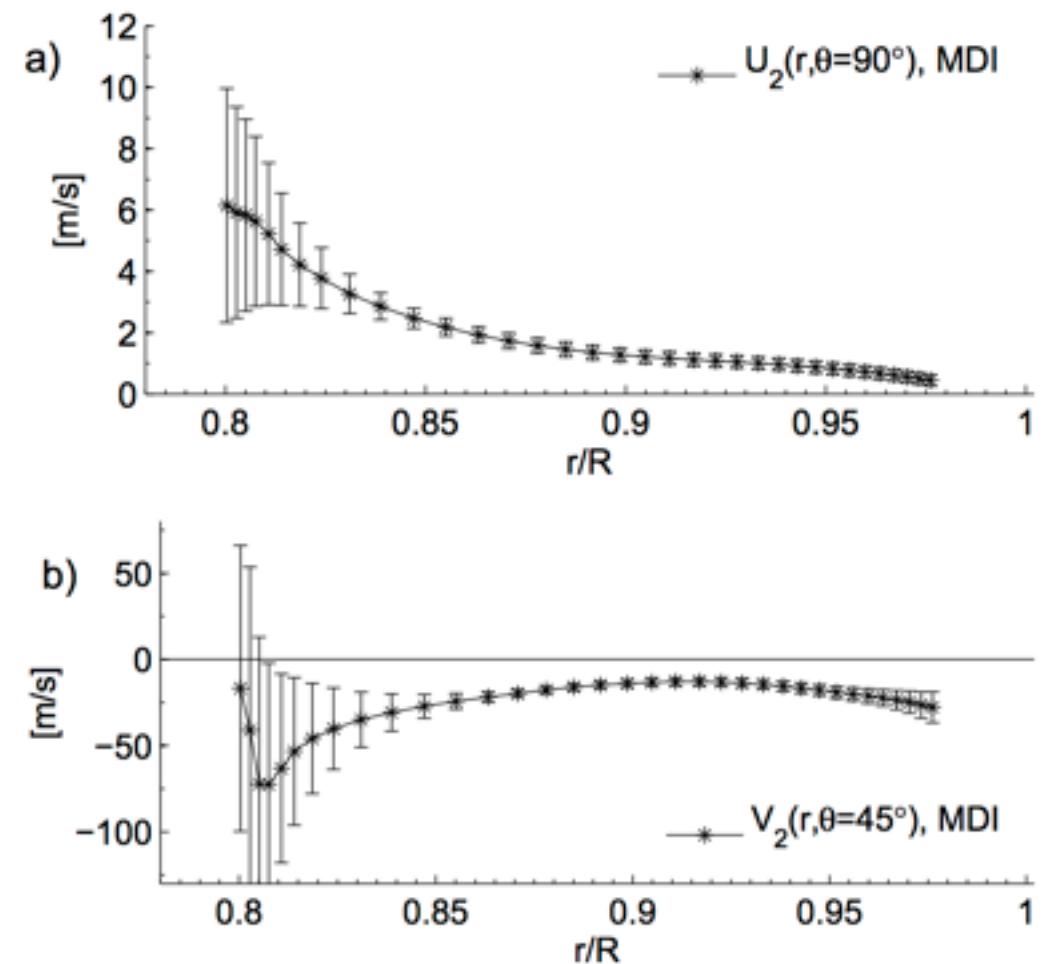


Fig. 2: a) Radial flow speed U_2 at the equator and b) horizontal flow speed V_2 at mid-latitude with 1σ -error bars as a function of r/R for the $s = 2$ component of the meridional flow estimated from data from MDI. Positive values of $U_2(V_2)$ indicate an outward (southward) directed flow.

Schad et al. 2012

Mean field hydrodynamics

Mean field ansatz: $\mathbf{u} = \bar{\mathbf{u}} + \mathbf{u}'$ $T = \bar{T} + T'$

Reynolds equation: $\rho \left[\frac{\partial \bar{\mathbf{u}}}{\partial t} + (\bar{\mathbf{u}} \cdot \nabla) \bar{\mathbf{u}} \right] = -\nabla \cdot \rho \mathbf{Q} - \nabla P + \rho \mathbf{g}$

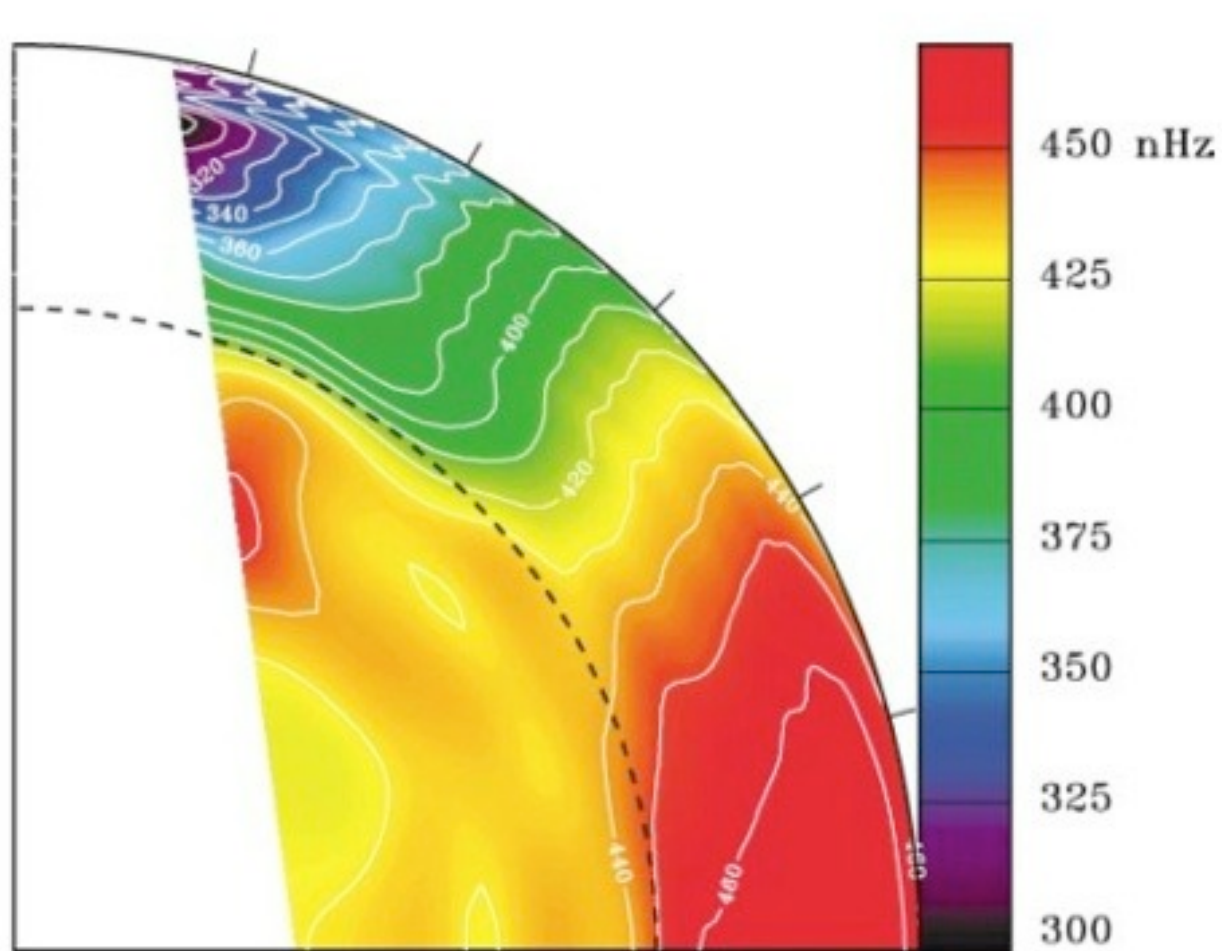
Heat transport equation: $\nabla \cdot (\mathbf{F}^{\text{conv}} + \mathbf{F}^{\text{rad}}) + \rho T \bar{\mathbf{u}} \cdot \nabla \bar{s} = 0$

Reynolds stress: $T_{ij} = -\rho Q_{ij}$

$$Q_{ij} = \langle u'_i u'_j \rangle = -\mathcal{N}_{ijkl} \frac{\partial \bar{u}_k}{\partial x_l} + \Lambda_{ijk} \Omega_k$$

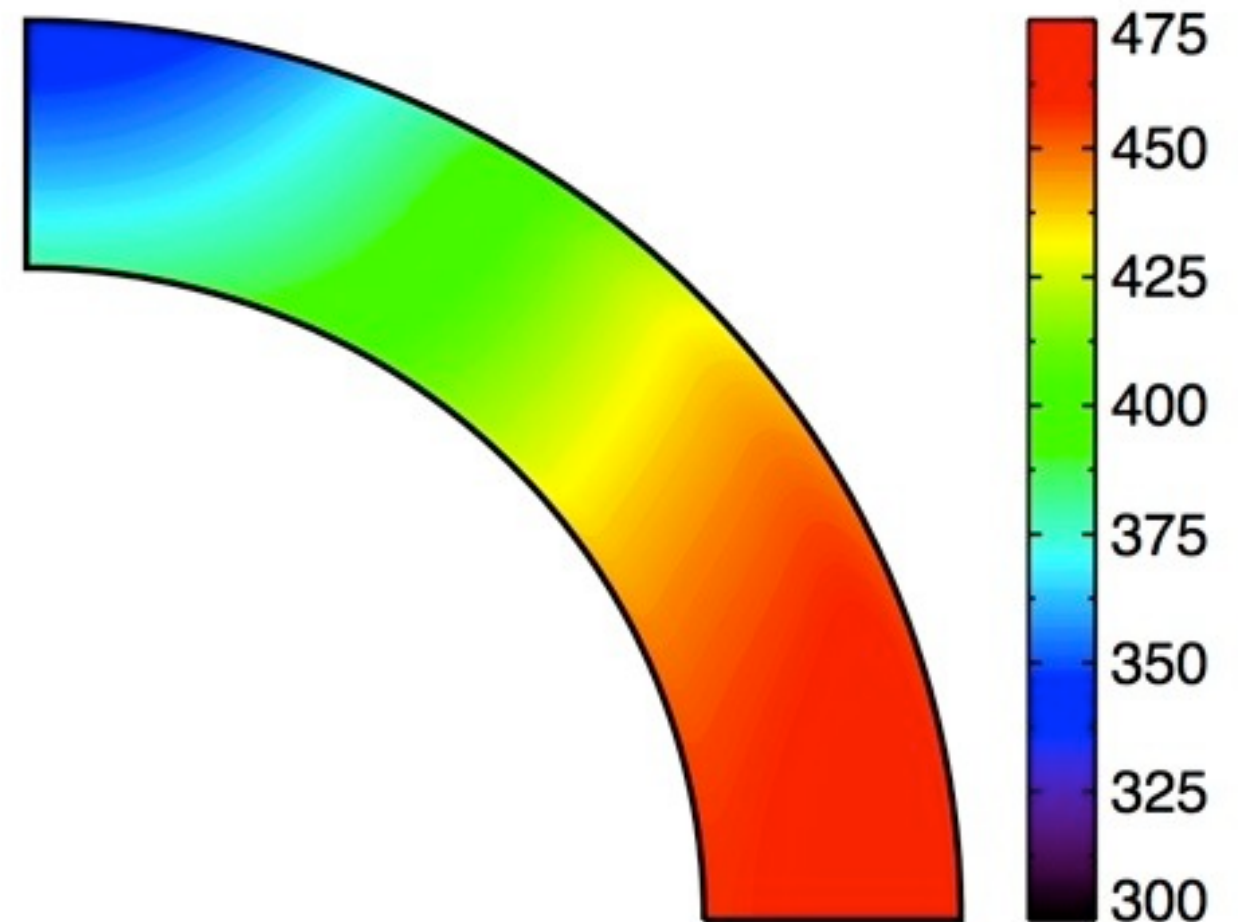
Convective heat flux: $F_i^{\text{conv}} = \rho c_p \langle u'_i T' \rangle = -\rho T \chi_{ij} \nabla_j \bar{s}$

Solar differential rotation: observed vs. model



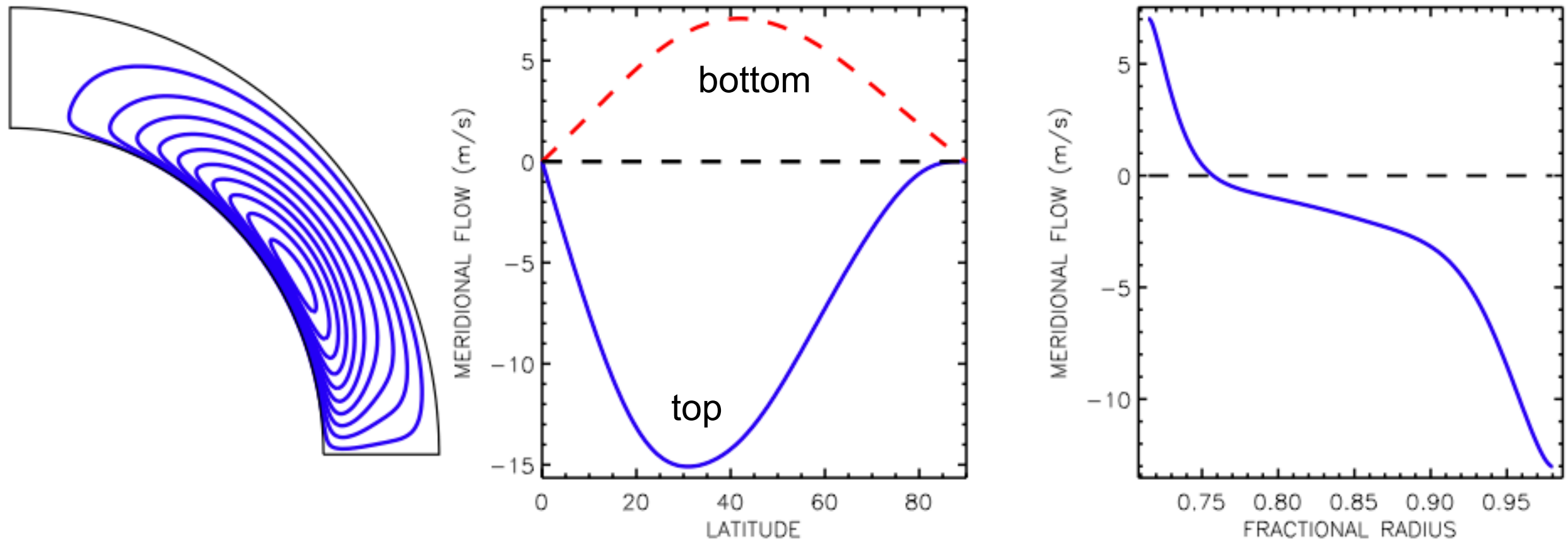
SOHO/MDI

Thompson et al. (2003)

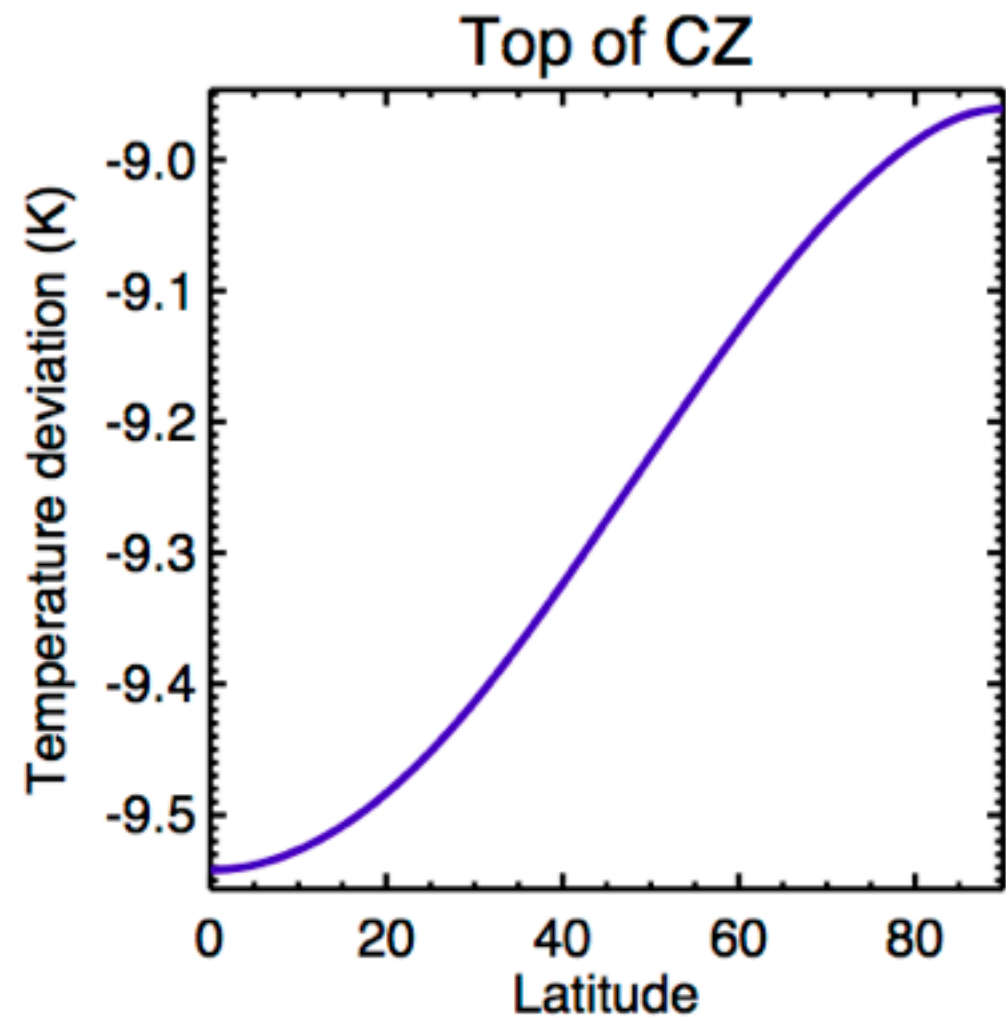
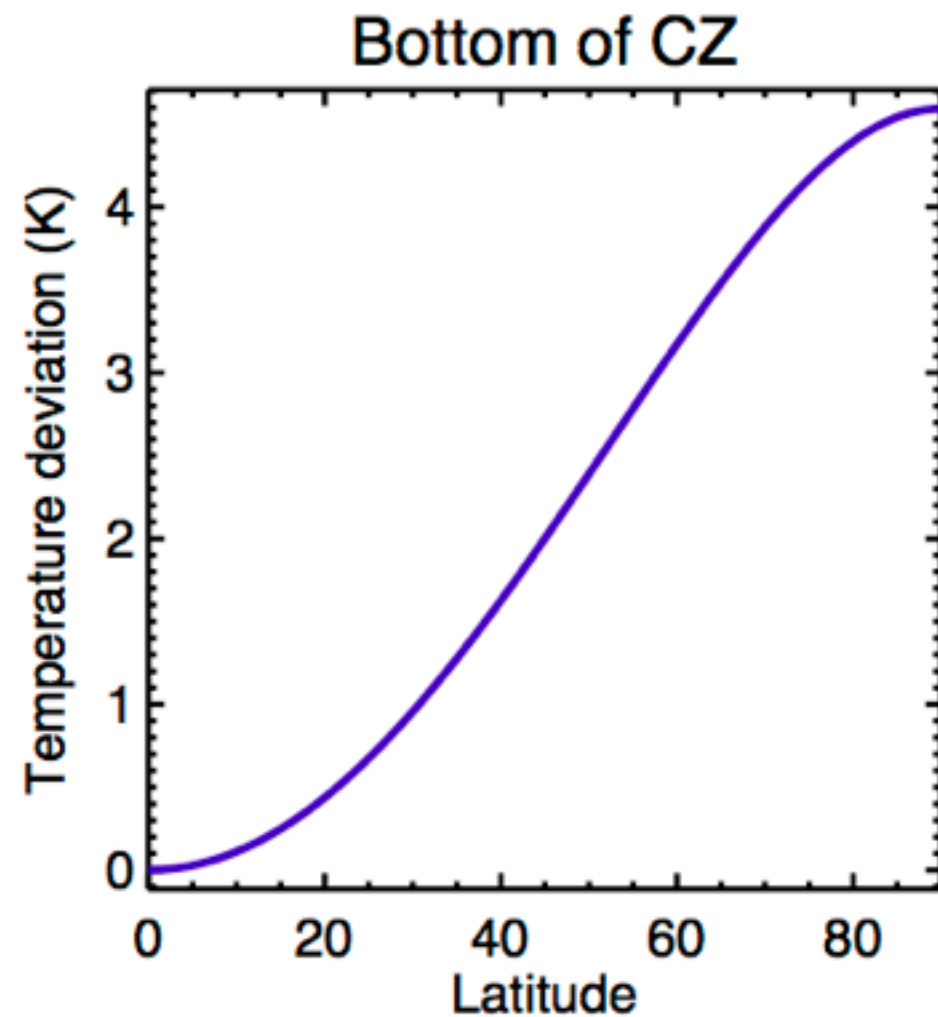


Model

Meridional flow: mean field model



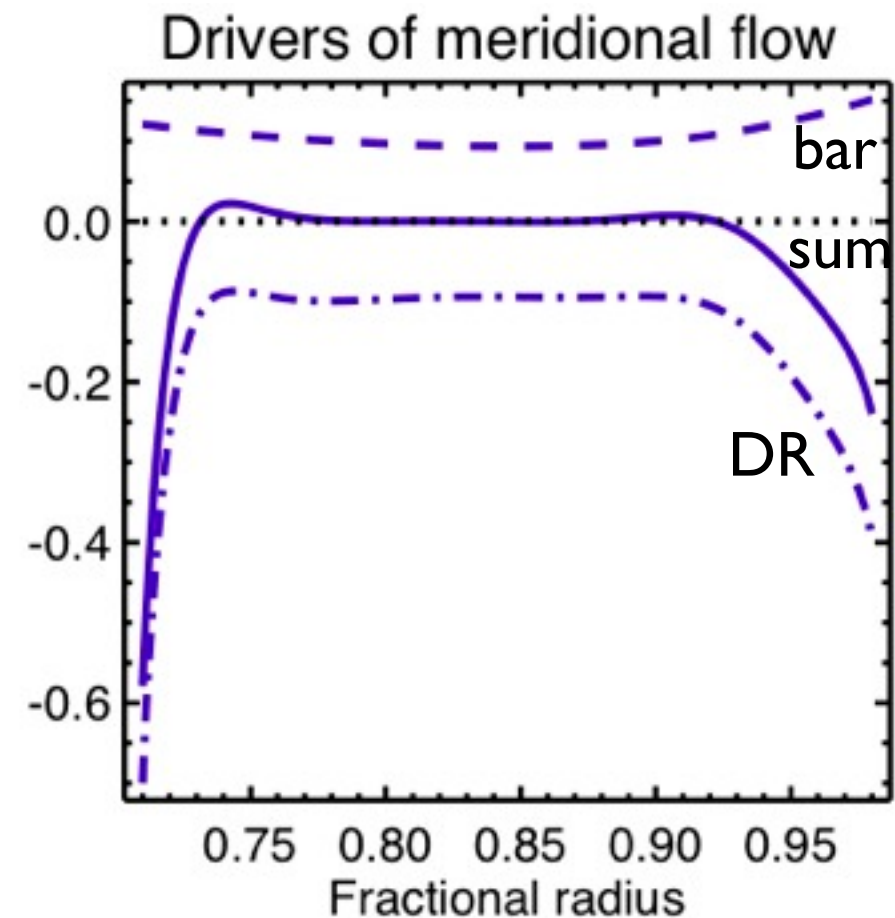
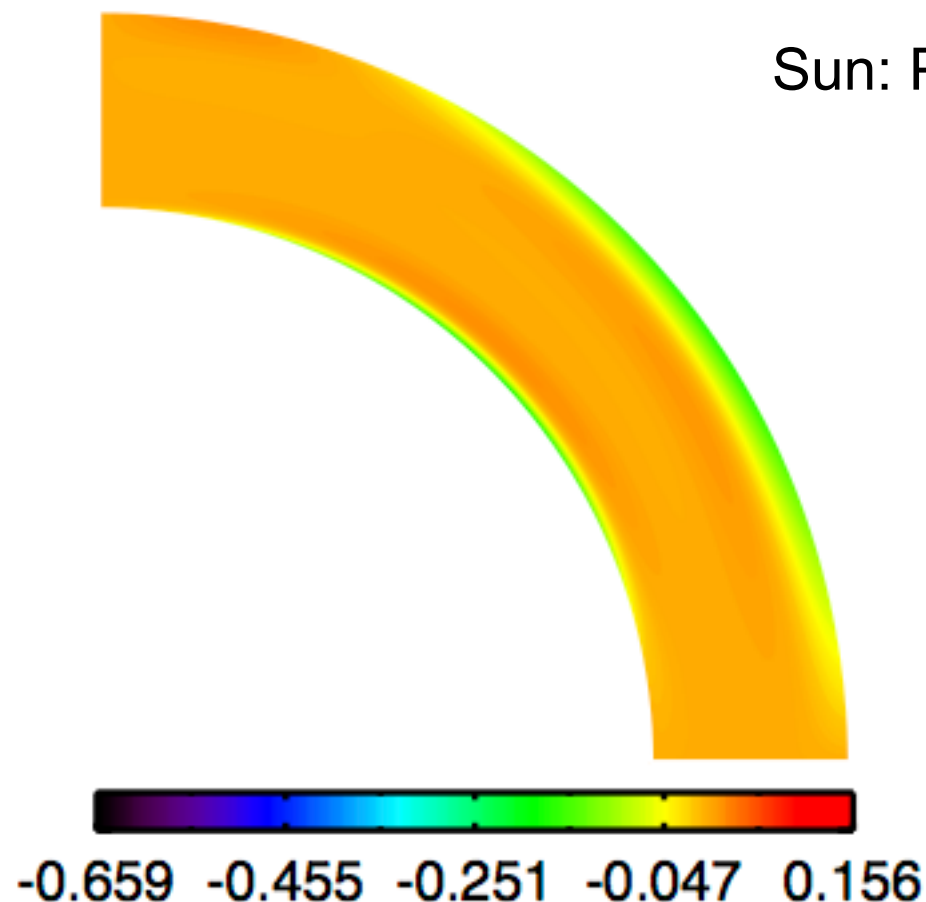
Deviation from adiabatic stratification



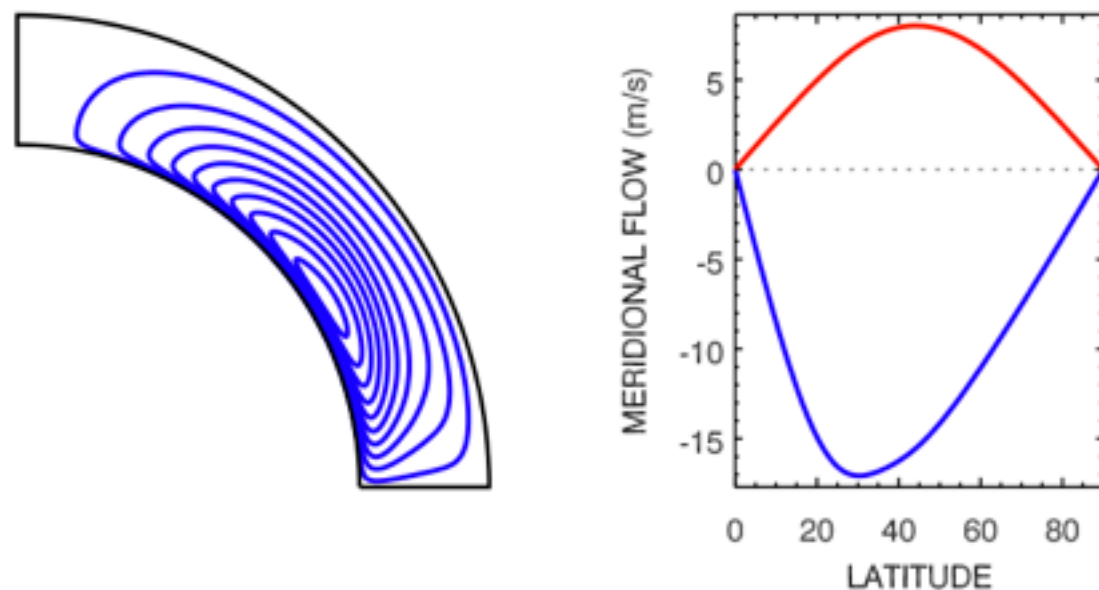
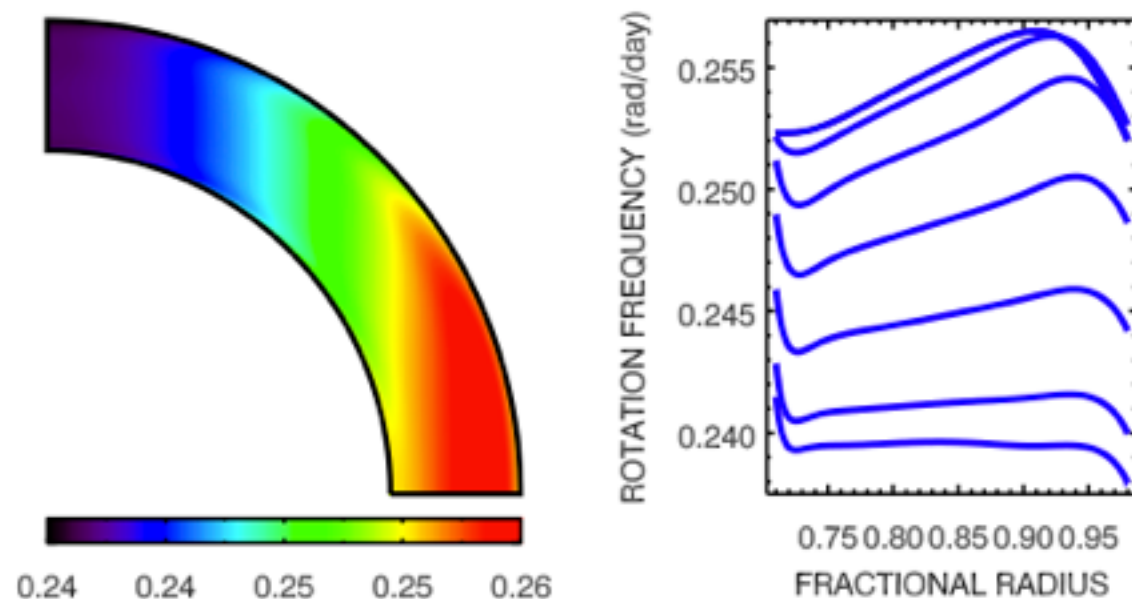
Thermal wind balance

Equation of motion:
$$-\left[\nabla \times \frac{1}{\rho} \nabla(\rho Q)\right]_{\phi} + r \sin \theta \frac{\partial \Omega^2}{\partial z} + \frac{1}{\rho^2} (\nabla \rho \times \nabla P)_{\phi} = 0$$

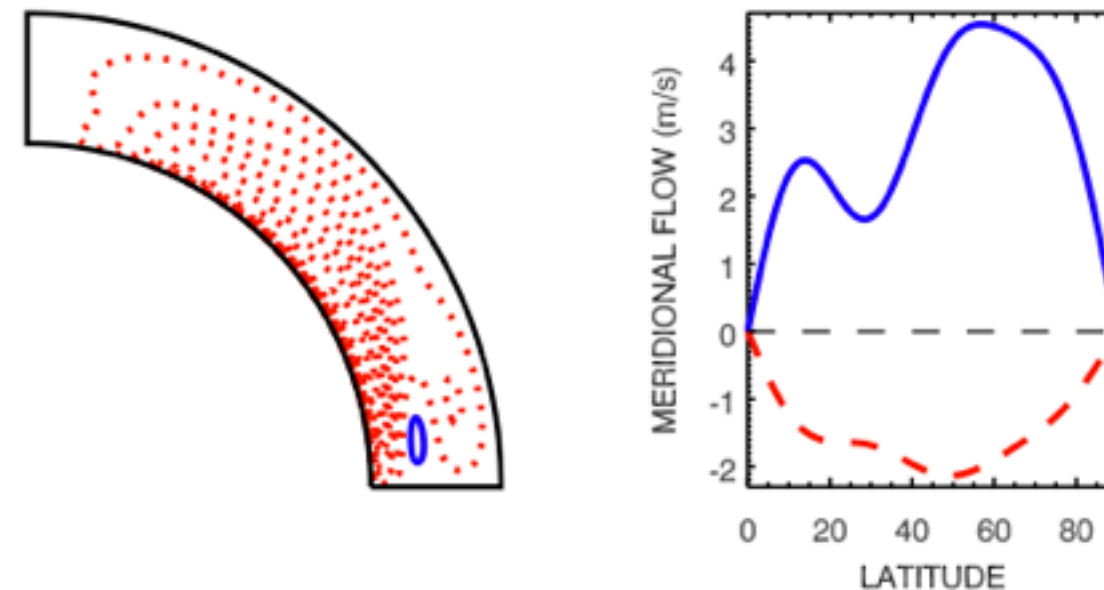
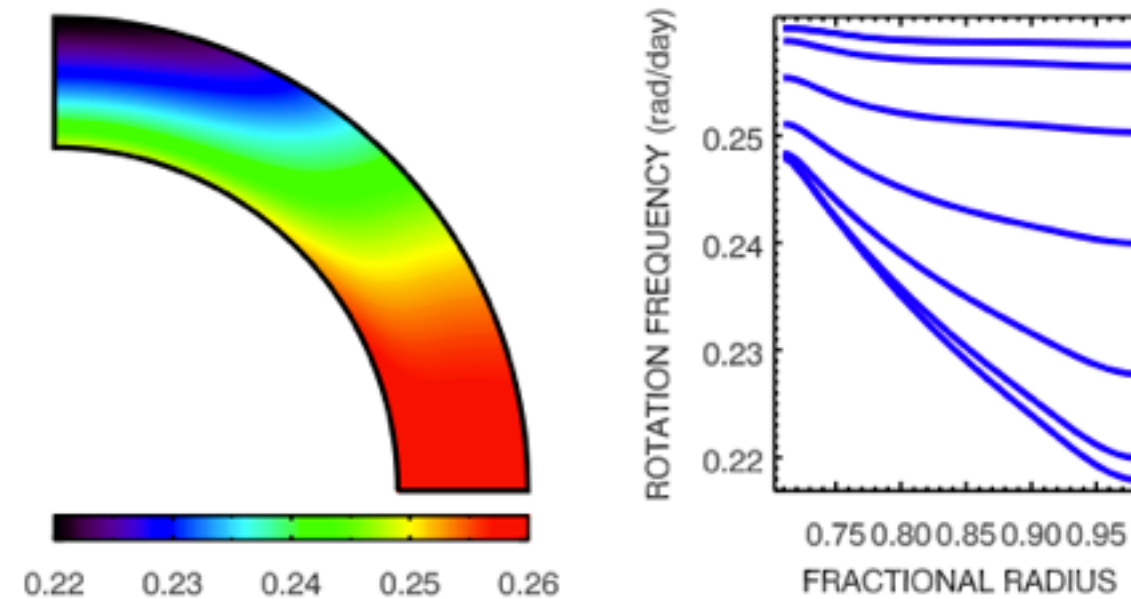
Fast rotation:
$$r \sin \theta \frac{\partial \Omega^2}{\partial z} - \frac{g}{rc_p} \frac{\partial \bar{s}}{\partial \theta} \approx 0 \quad \text{not valid in boundary layers!}$$



Reynolds stress vs. baroclinic flow

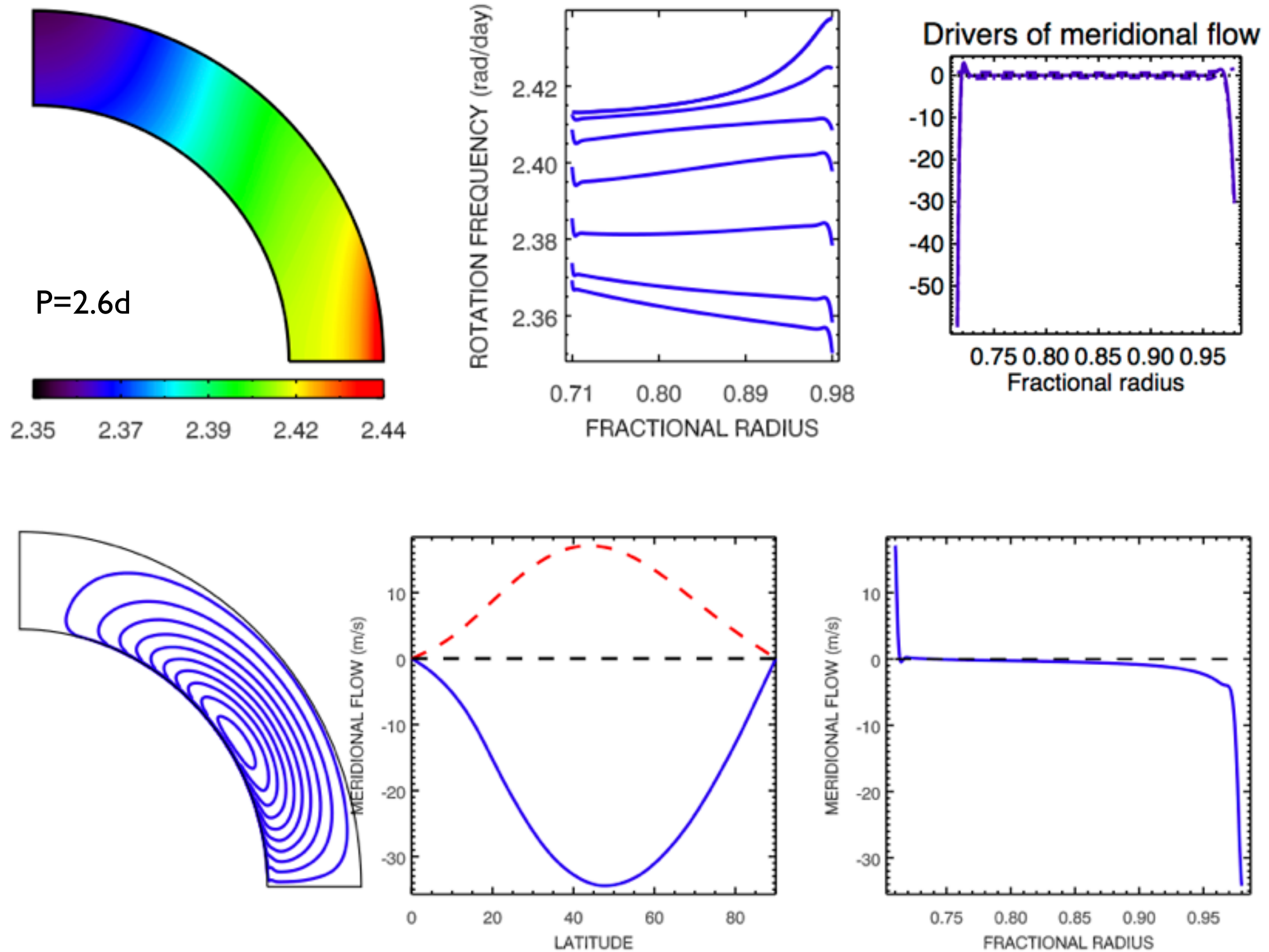


no baroclinic flow



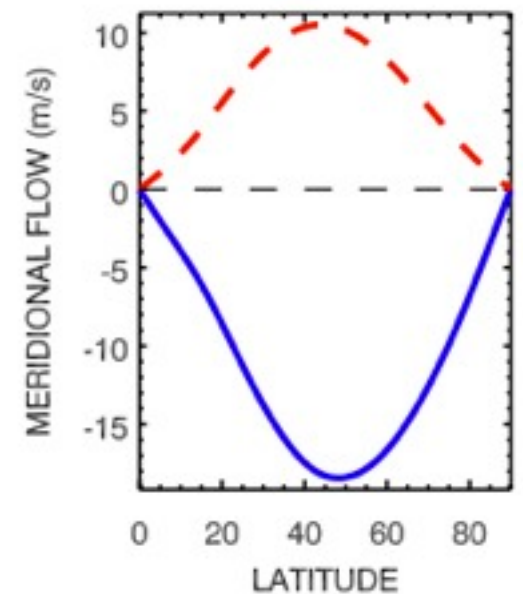
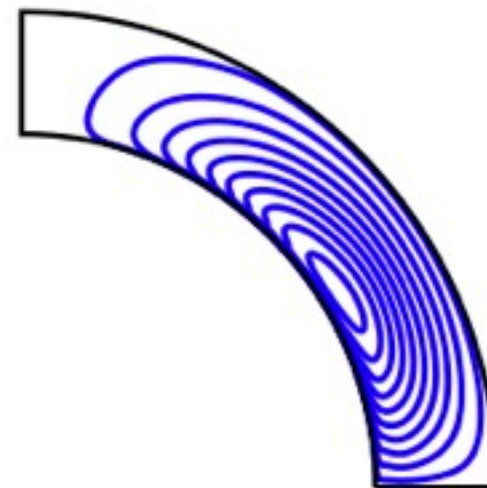
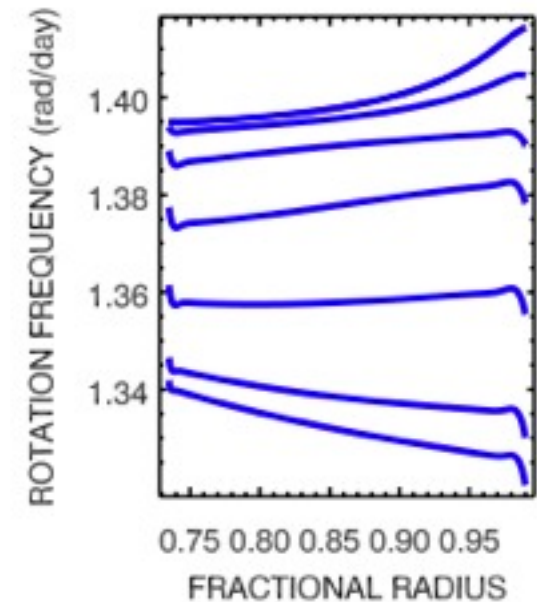
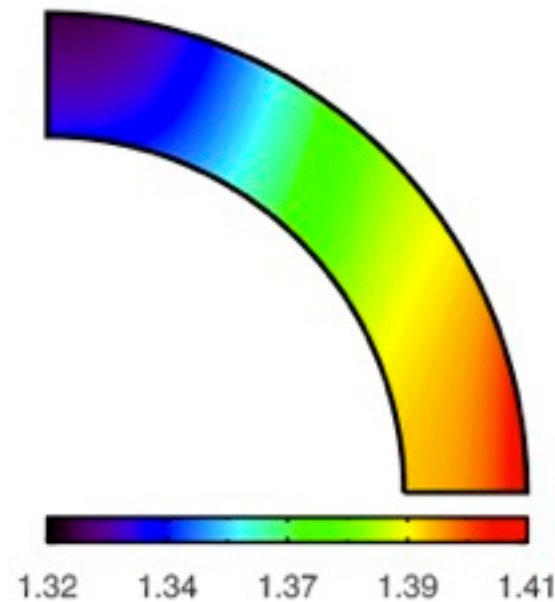
no Λ effect

Rapidly rotating Sun

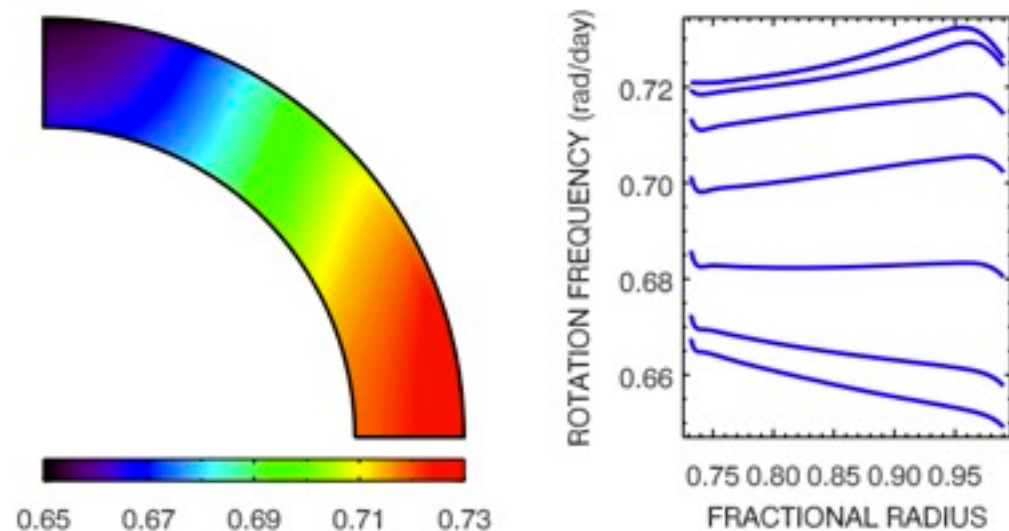


CoRoT-2a

- young solar twin (0.5 Gyr)
- rotation period: 4.5d
- spot modeling (Fröhlich et al. 2009):
 - solar-type DR
 - $\delta\Omega=0.11$ rad/day
- theory:
 - solar-type DR
 - $\delta\Omega=0.09$ rad/day
 - solar-type flow

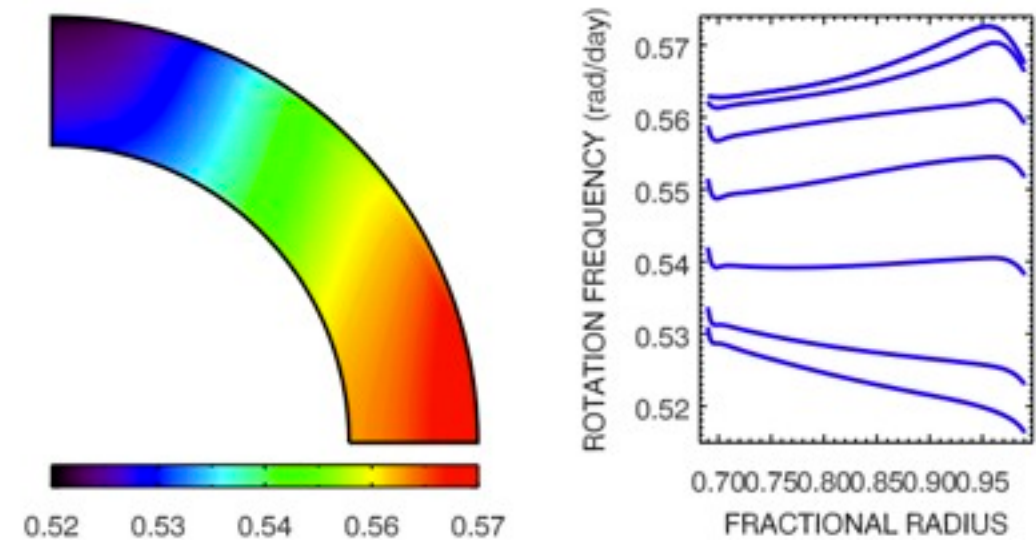


MOST stars



κ I Cet (G5 V)

- $P_{\text{rot}}=8.8\text{d}$
- surface DR:
- photometry: $\delta\Omega=0.064$ rad/day
- model: $\delta\Omega=0.077$
- surface meridional flow speed: 19.3 m/s

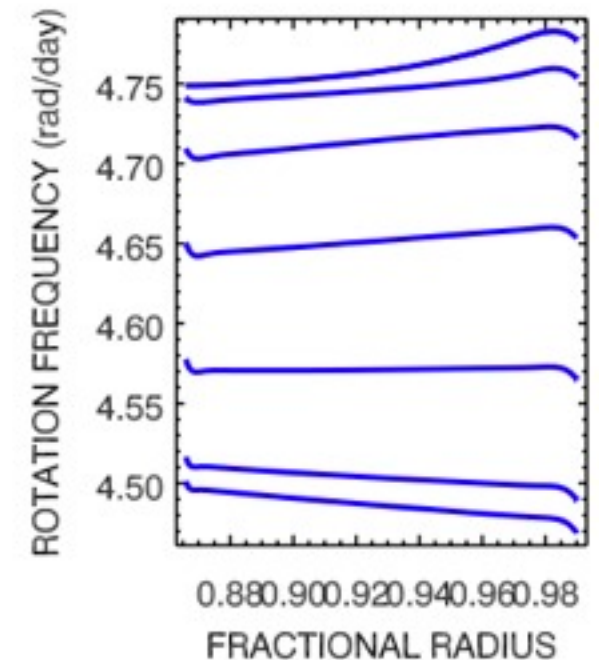
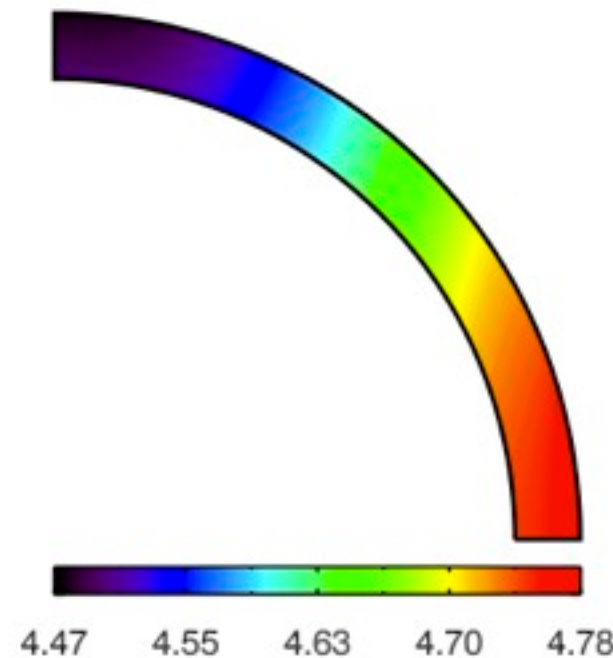
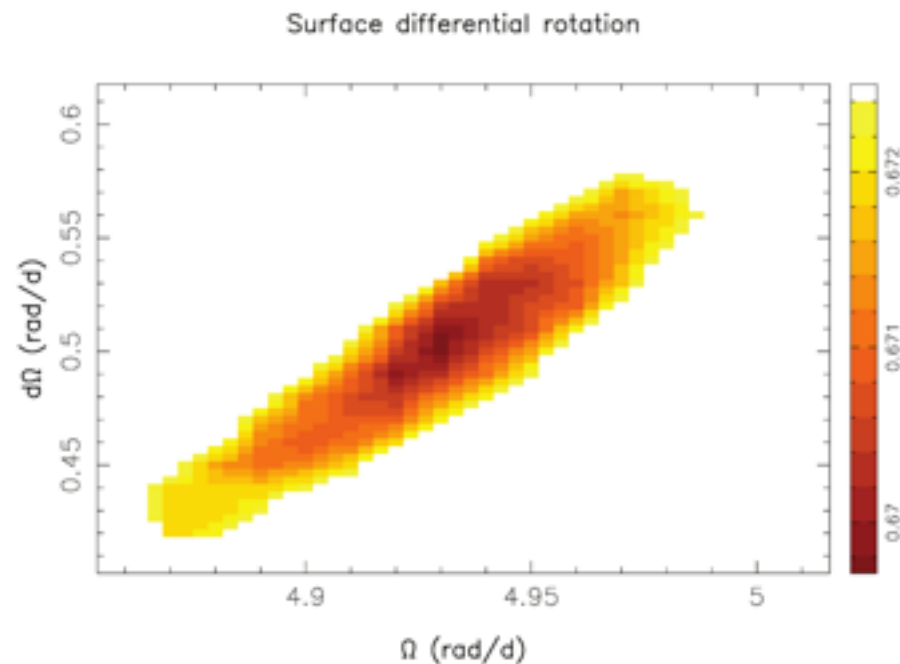


ϵ Eri (K2 V):

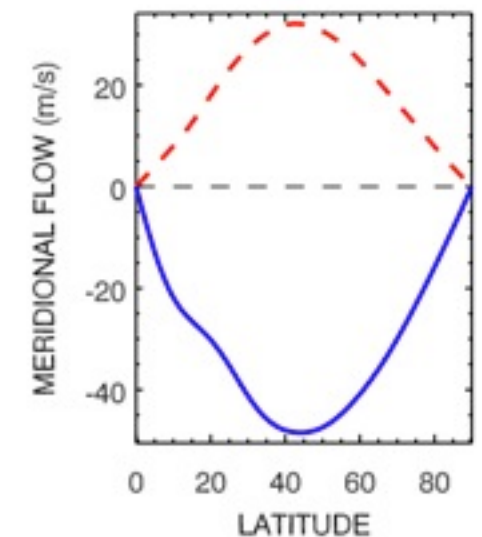
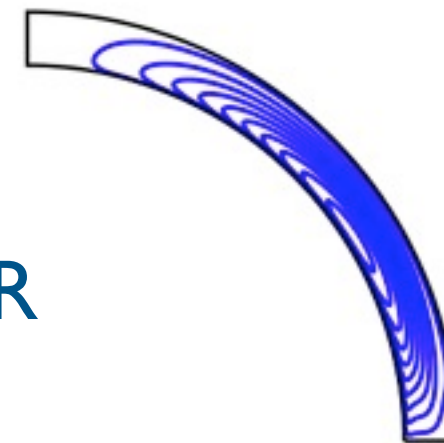
- $P_{\text{rot}}=11.2\text{d}$
- surface DR:
- photometry: $\delta\Omega=0.062$ rad/day
- model: $\delta\Omega=0.051$ rad/day
- surface meridional flow speed: 12.5 m/s

HDI 71488 (V899 Her)

Jeffers & Donati 2008



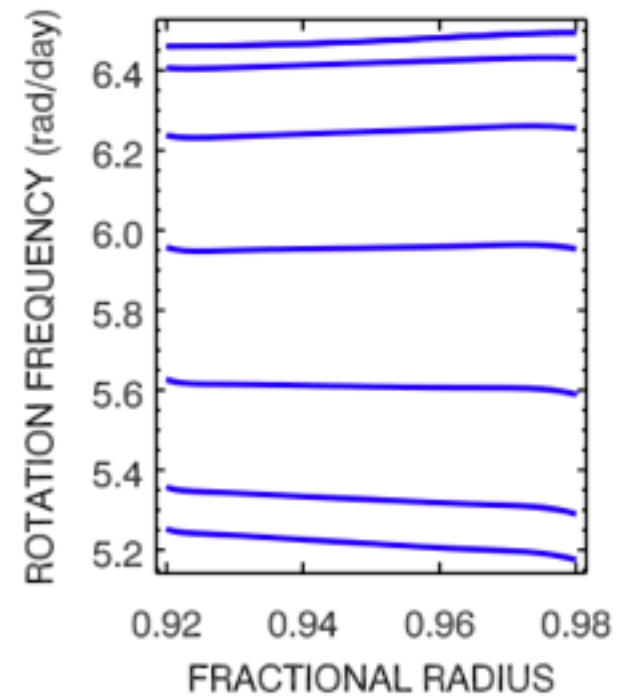
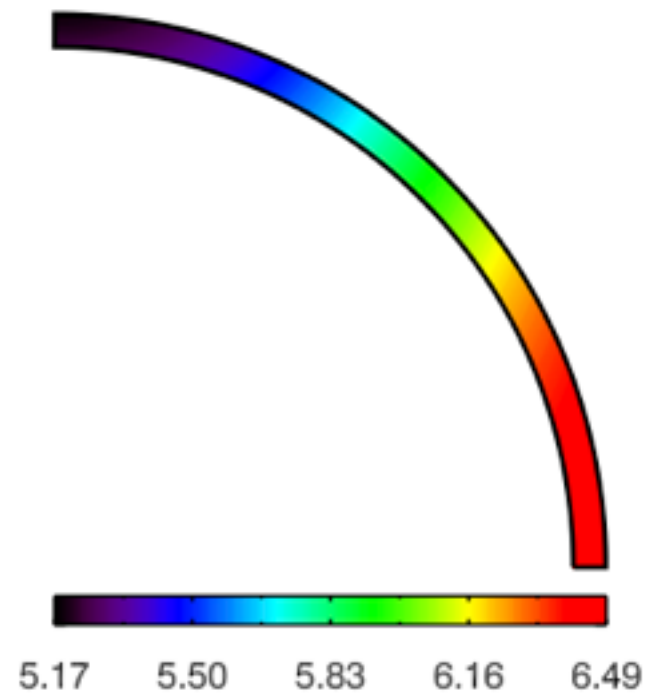
- fast-rotating young G dwarf
- 1.1 solar masses
- Doppler-Zeeman imaging: strong DR
- model: $\delta\Omega=0.3$ rad/day



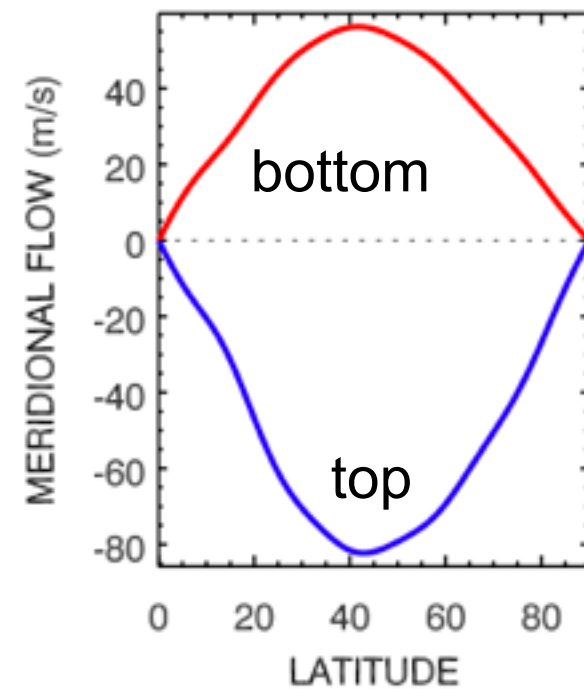
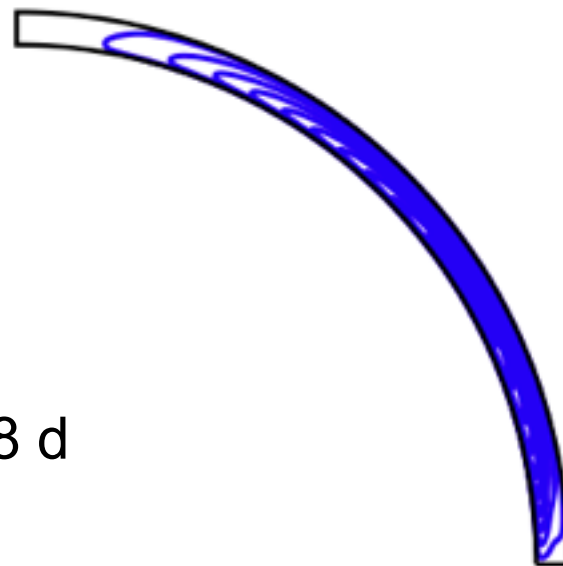
Shallow convection zone

$P=1\text{d}$

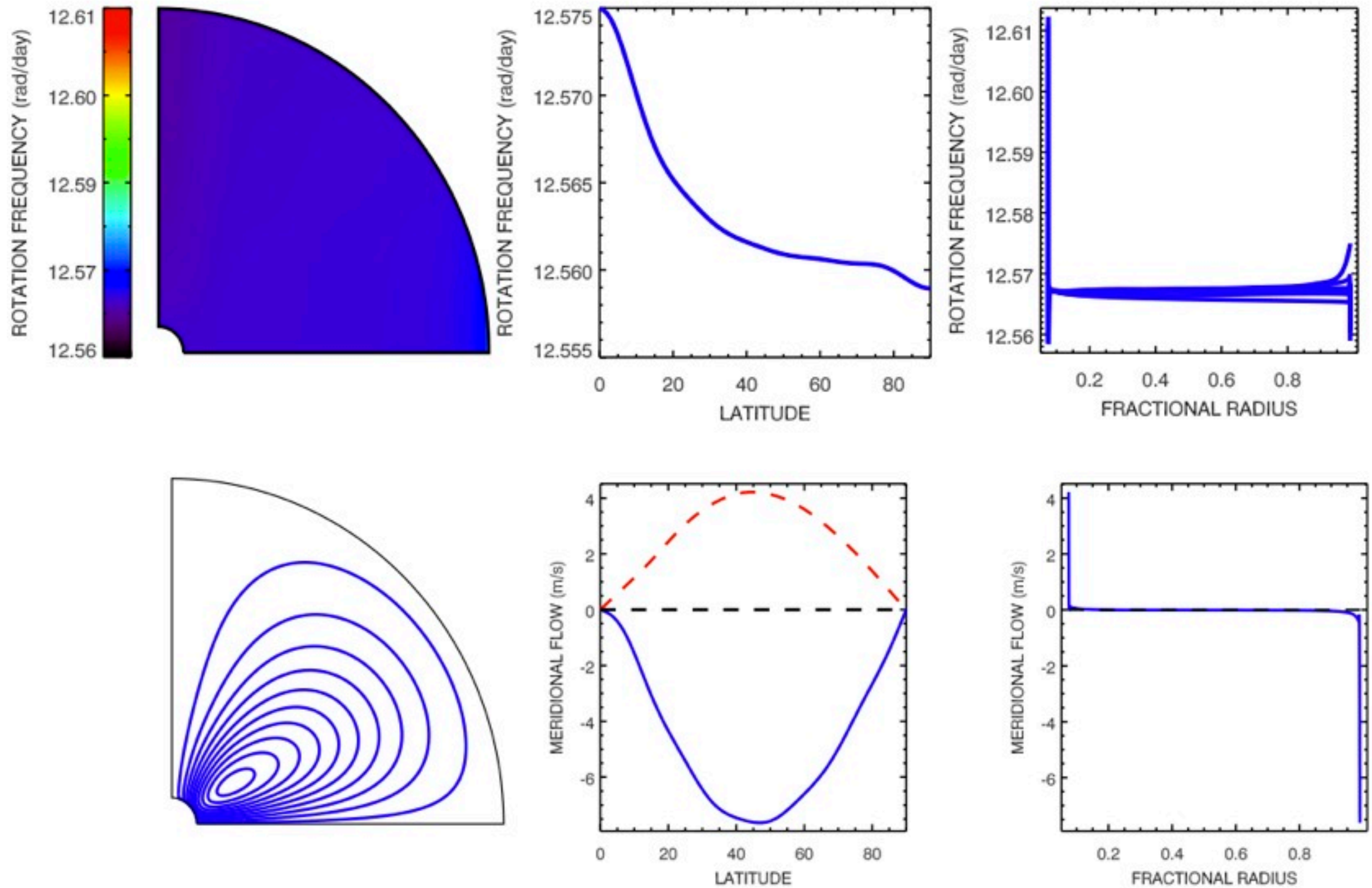
$1.4M_{\text{sol}}$



$\delta\Omega=1.3 \text{ rad/d}$
lapping time=4.8 d



M dwarf (0.3 sol. masses)



$P=0.5d$

Stellar differential rotation

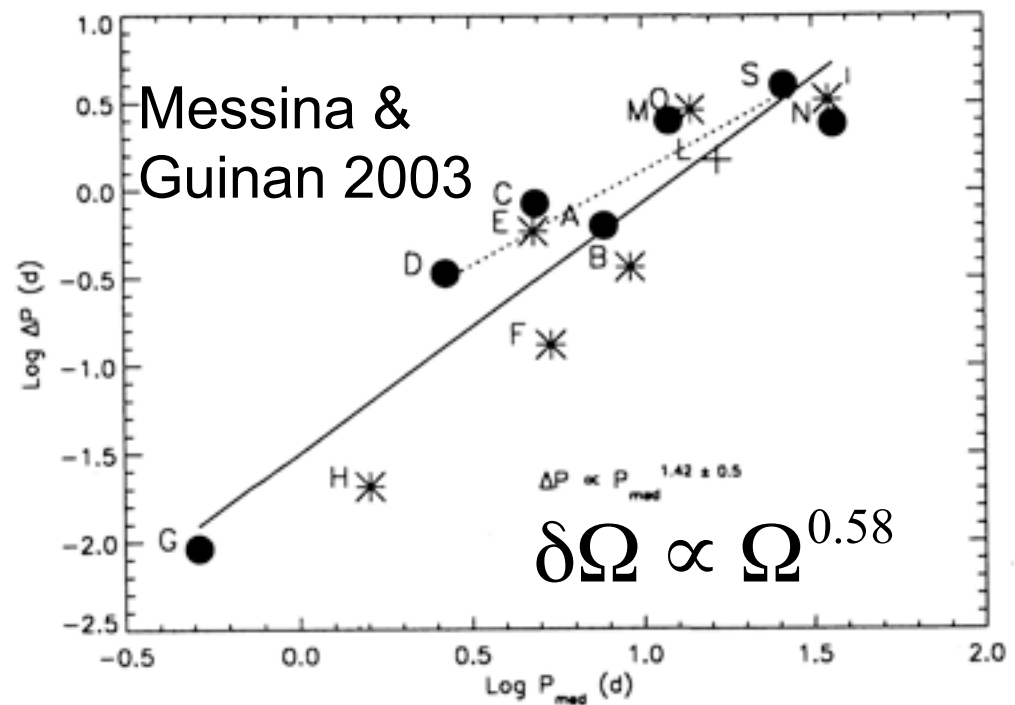
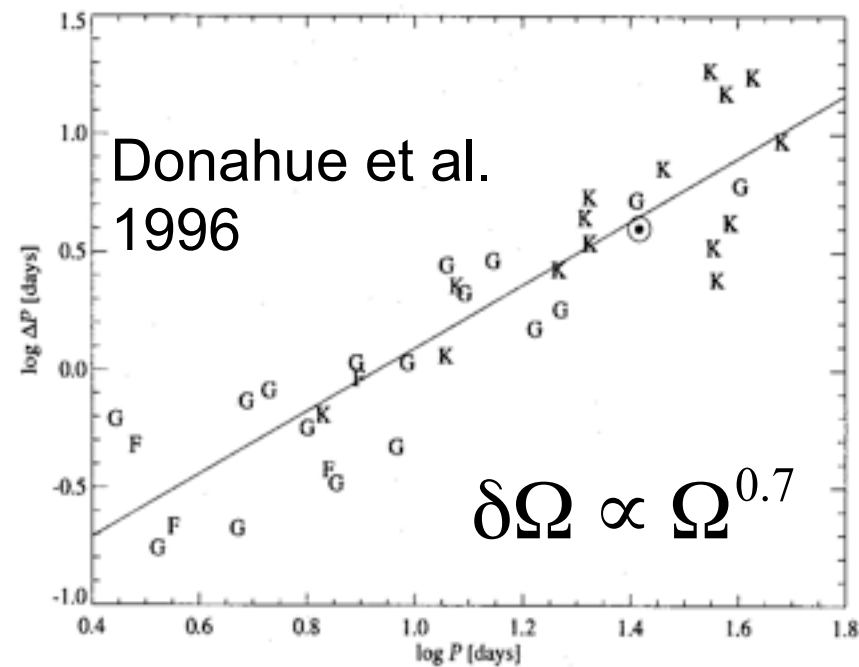
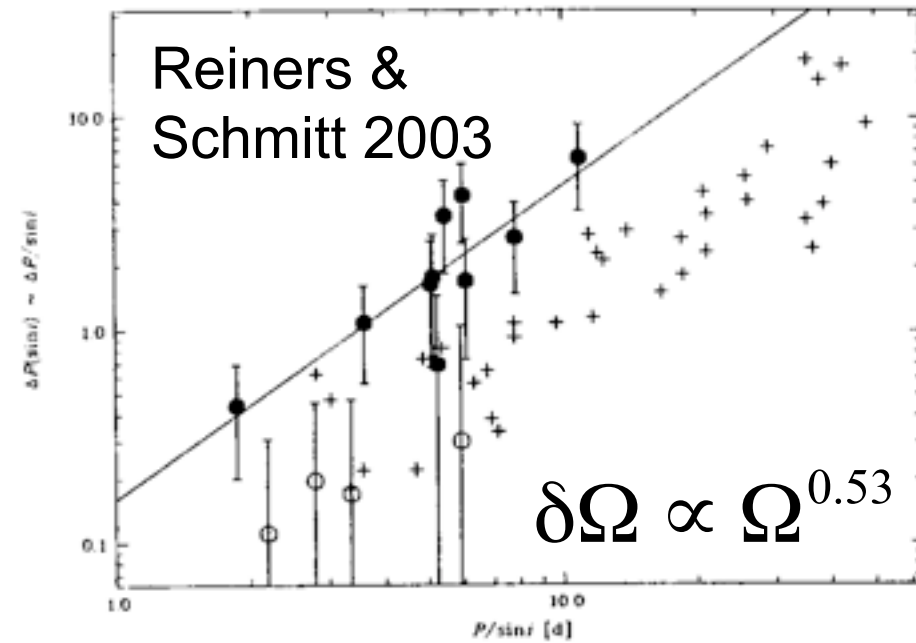
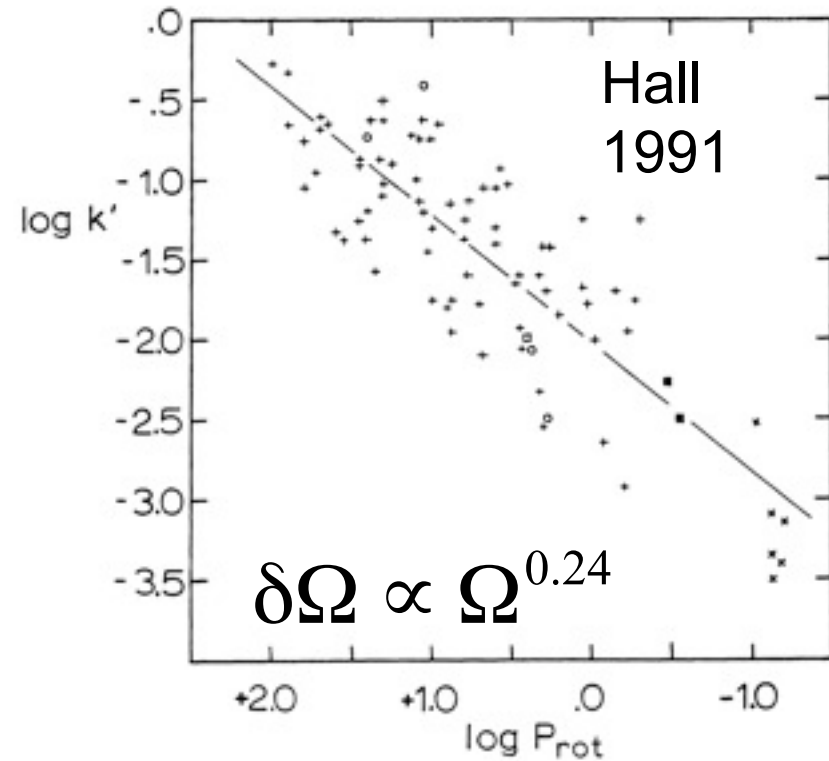
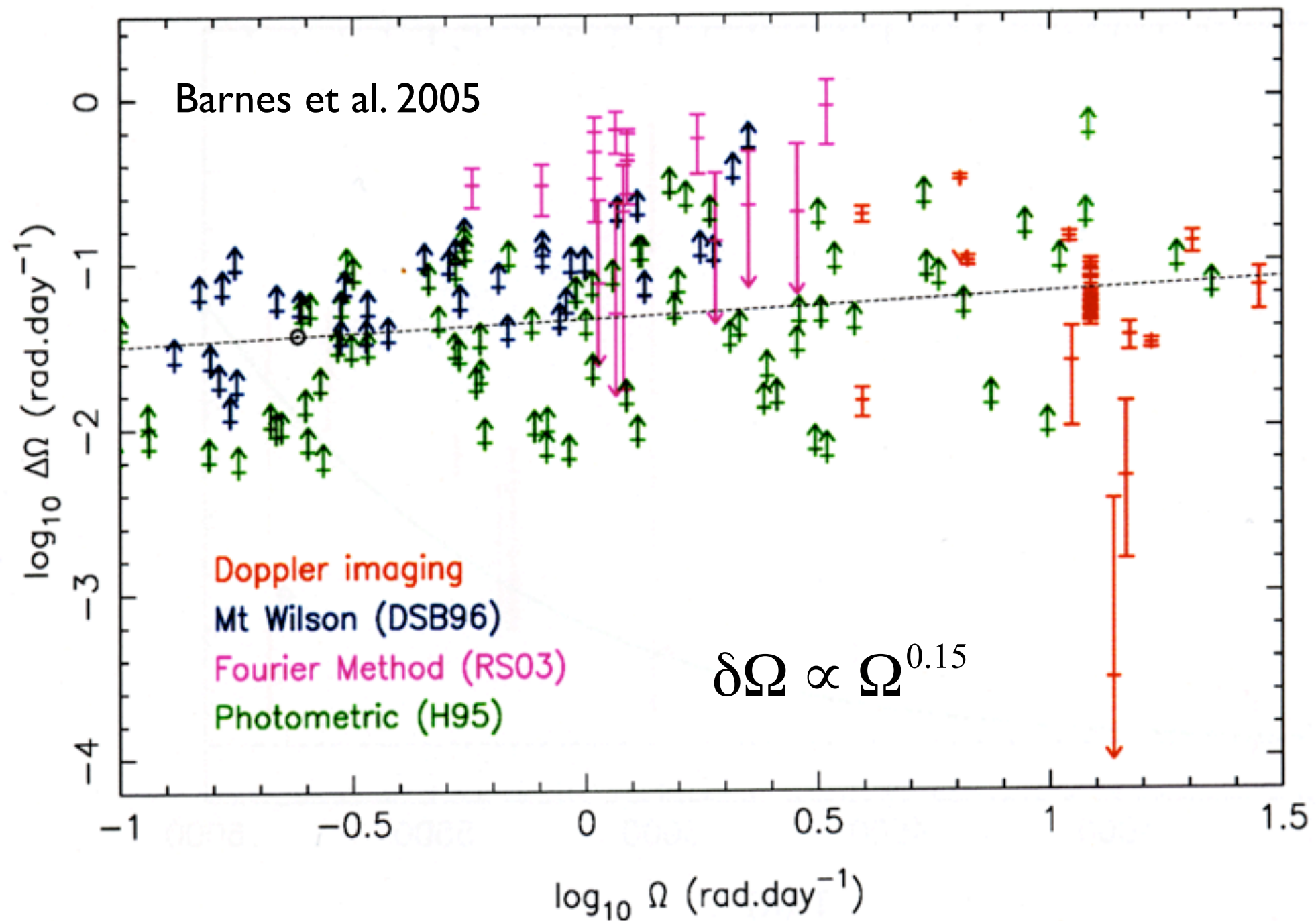
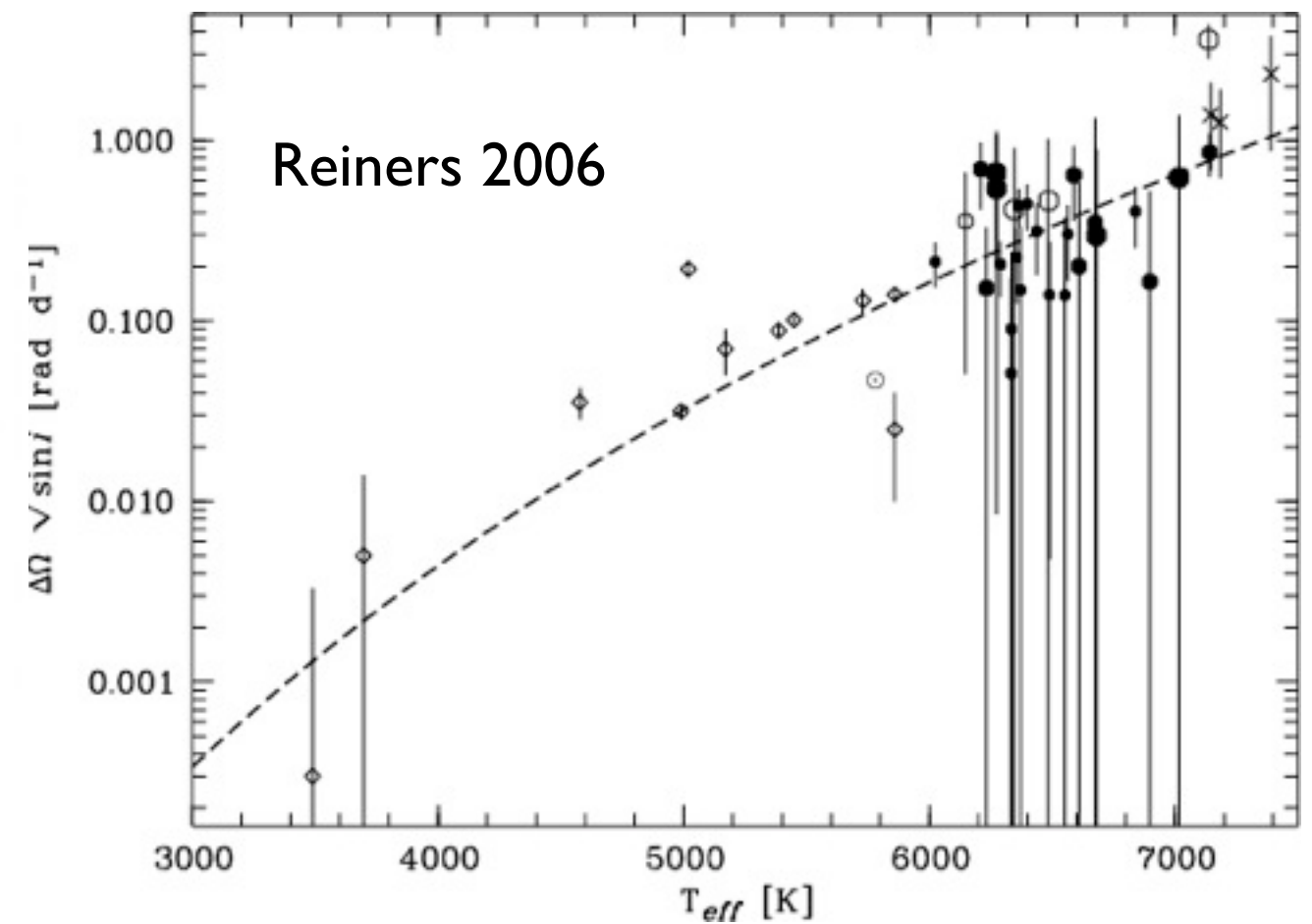
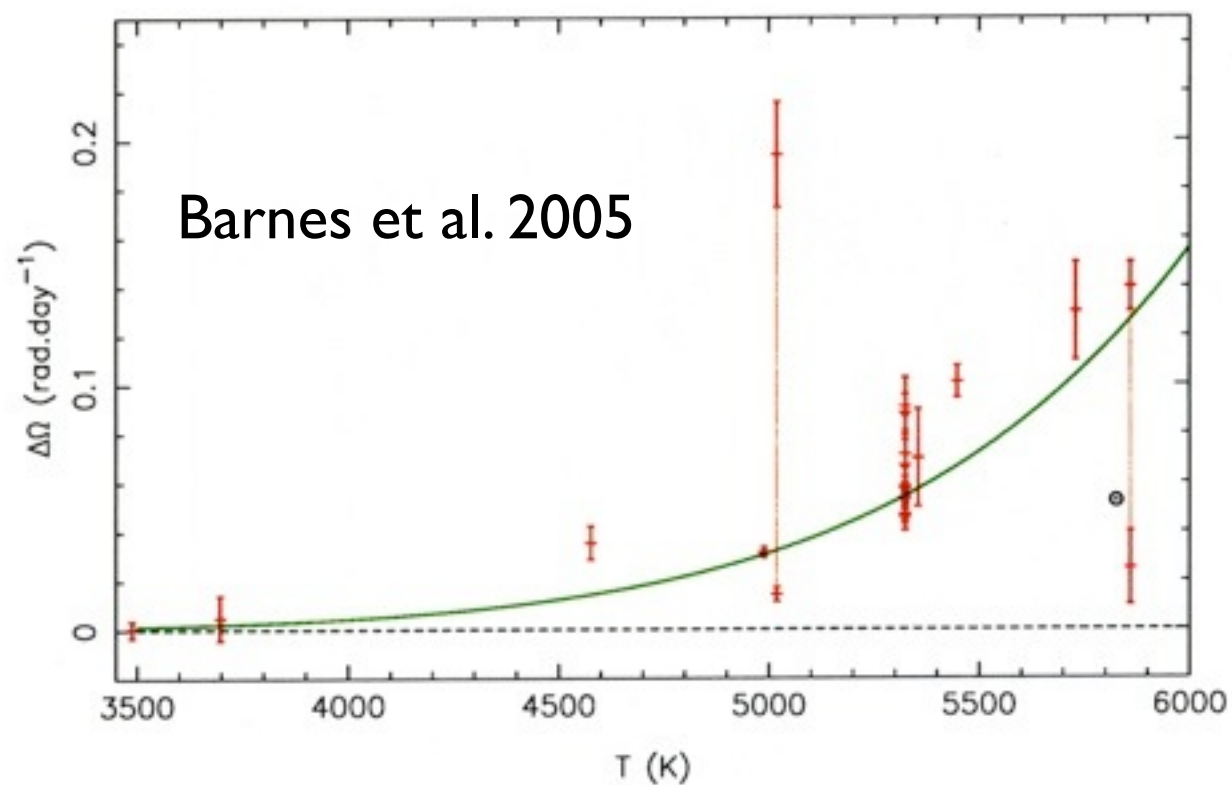


FIG. 3.—Range of observed rotation periods vs. $\log \langle P \rangle$. Least-squares fit of these data yields $\Delta P \propto \langle P \rangle^{1.3 \pm 0.1}$ (correlation coefficient $r = 0.90$).

Stellar Differential Rotation: Dependence on rotation?

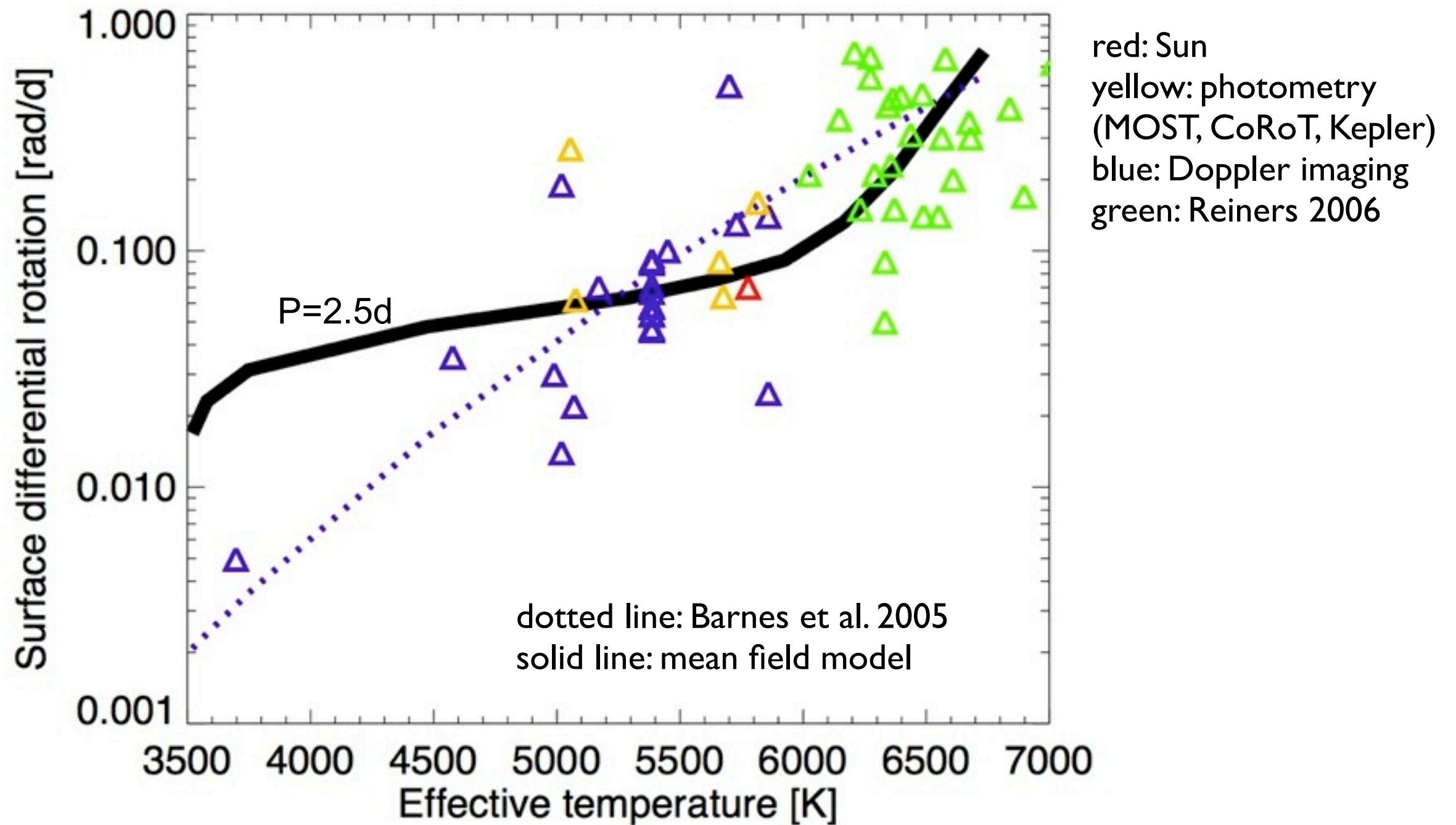


Stellar Differential Rotation: dependence on temperature

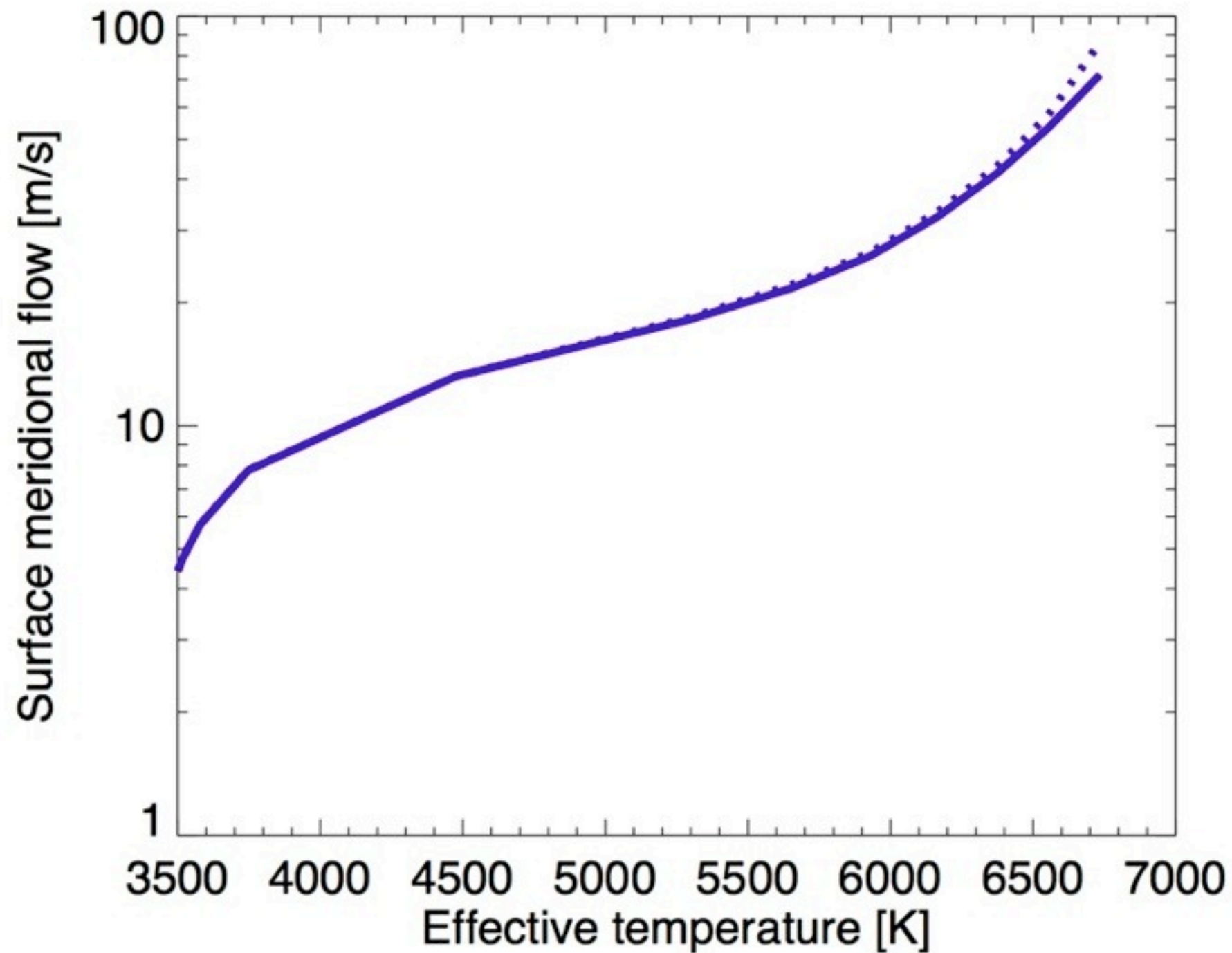


$$\delta\Omega \propto T^{8.9}$$

DR: dependence on temperature

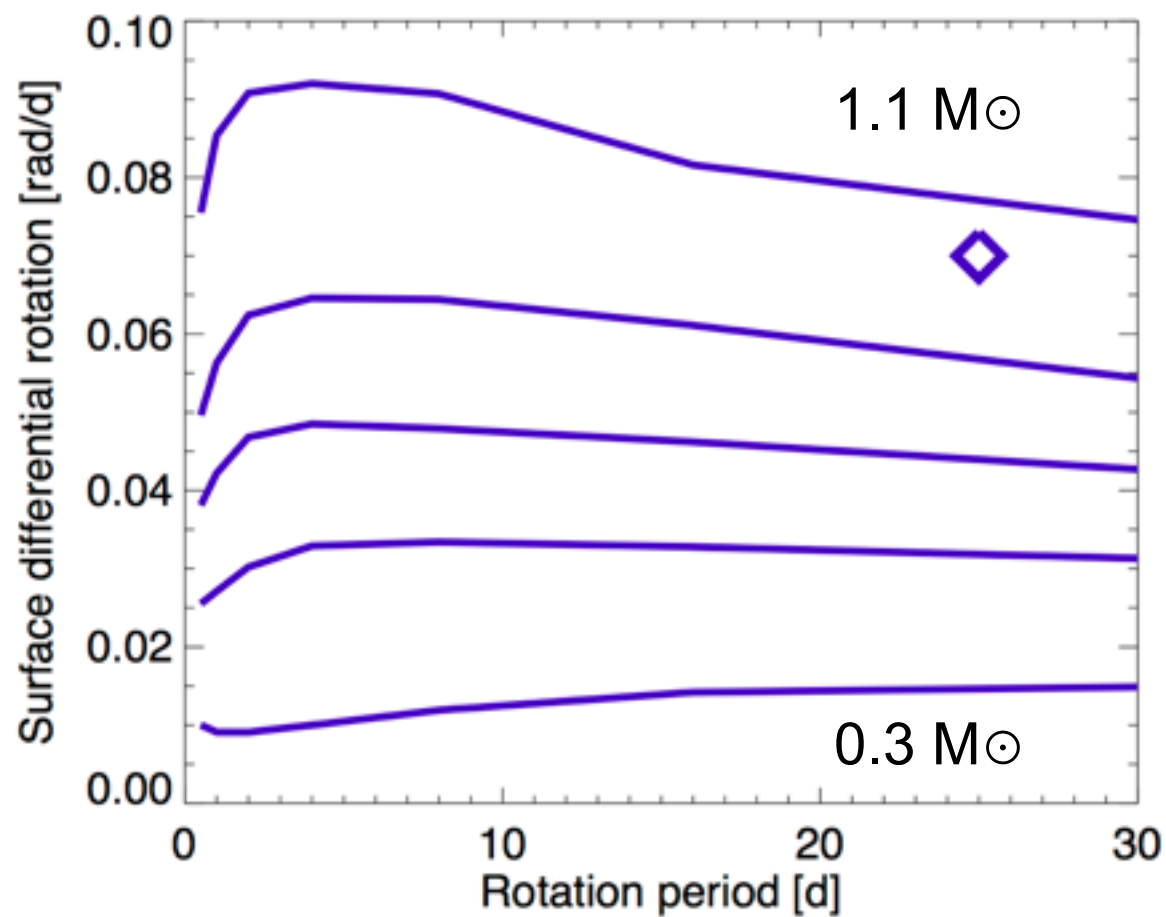


Meridional flow: temperature dependence

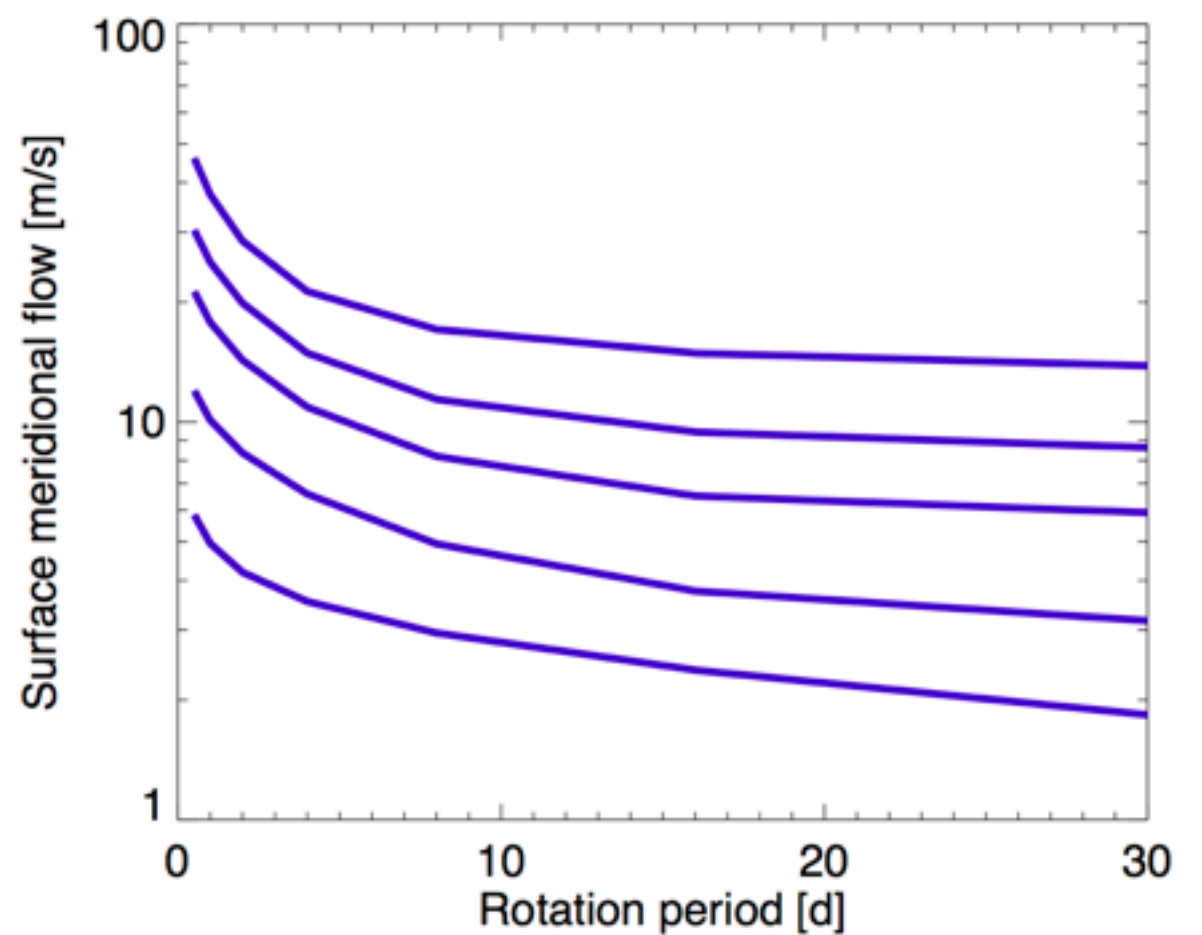


Dependence on rotation period

Differential rotation



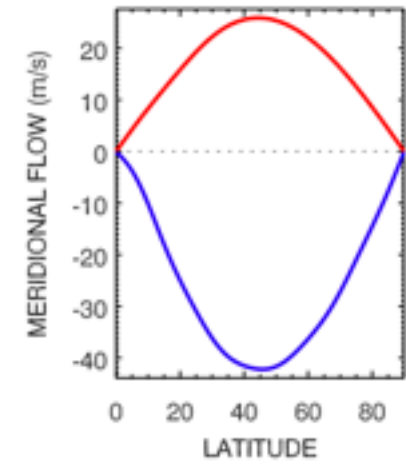
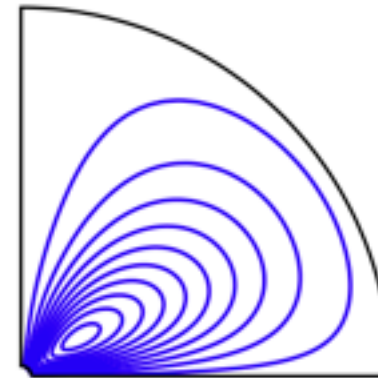
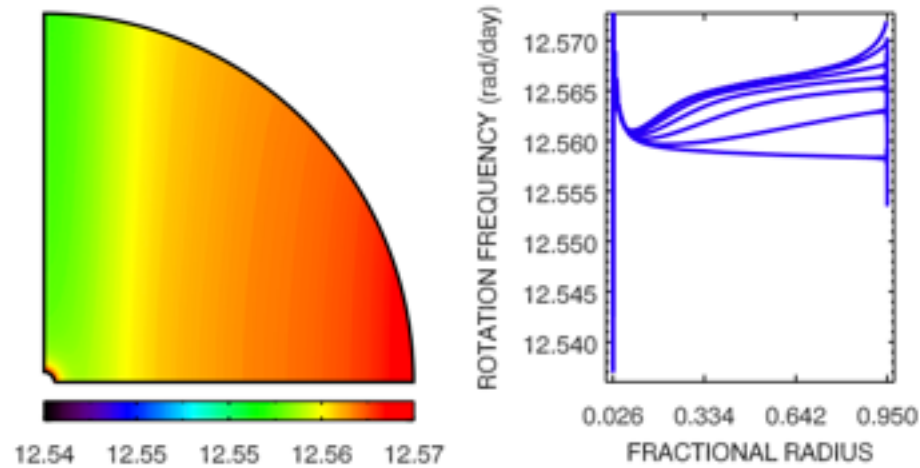
Meridional flow



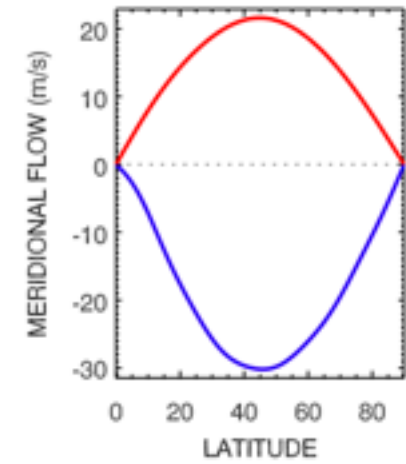
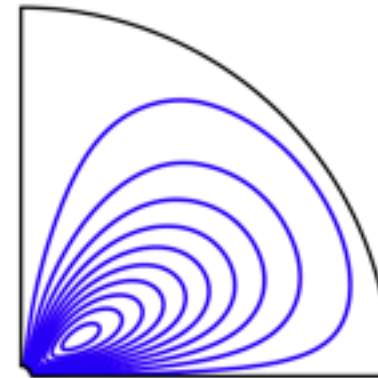
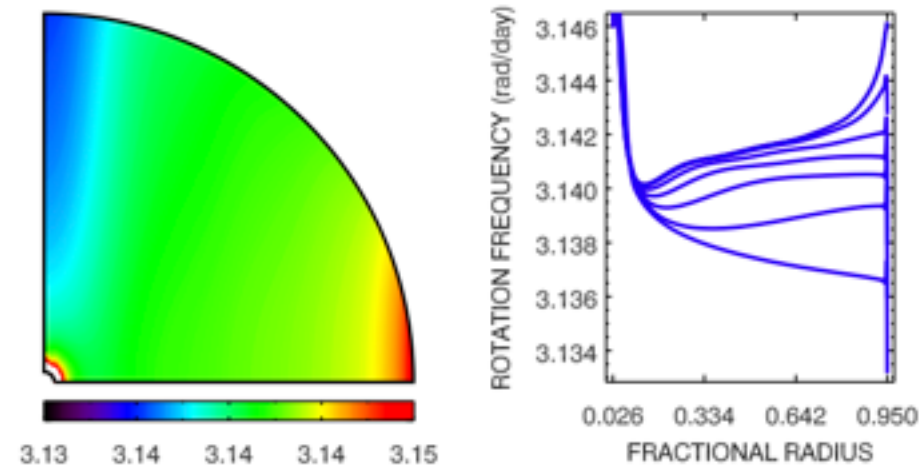
T Tauri star

1 Myr
2.33 R_{\odot}
1.94 L_{\odot}

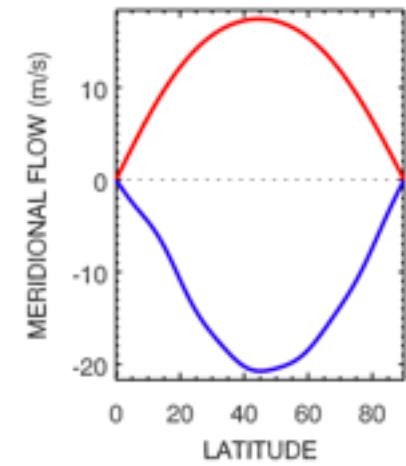
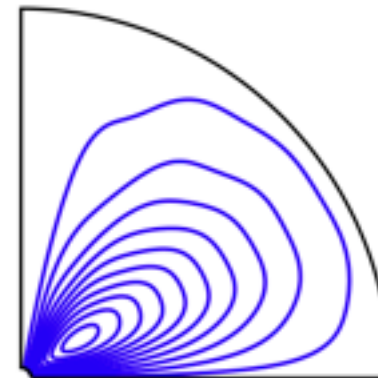
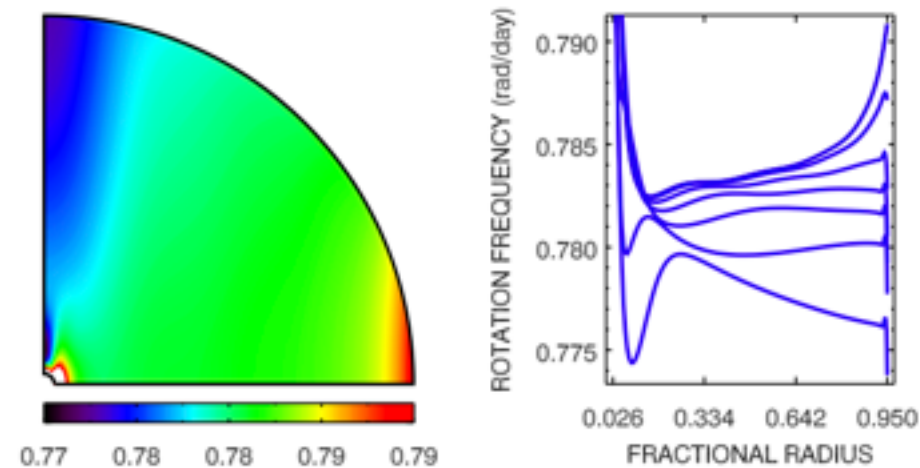
P=0.5d



P=2d



P=4d



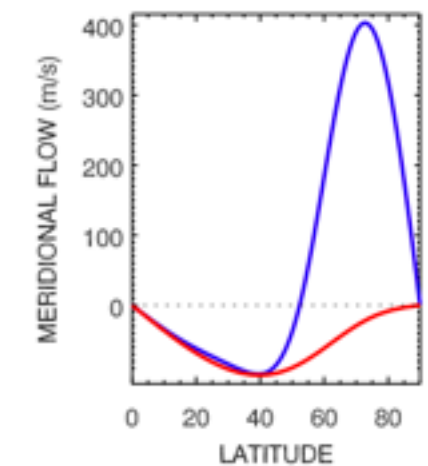
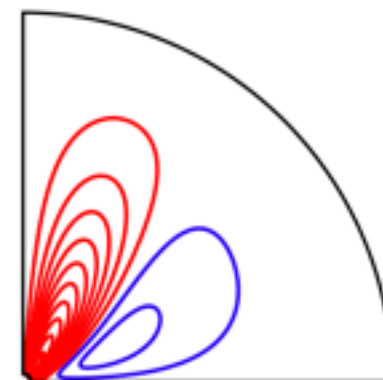
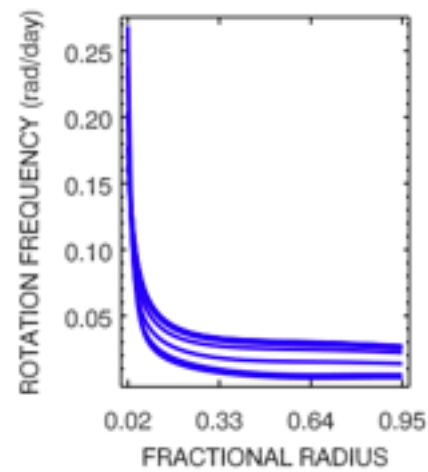
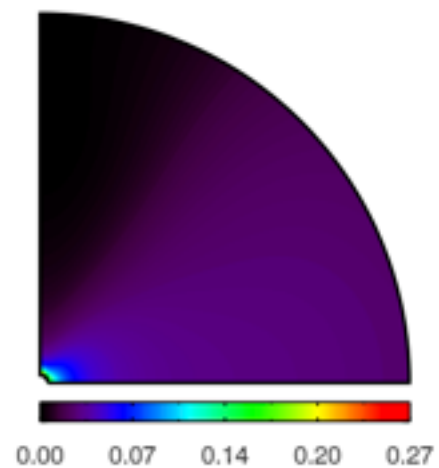
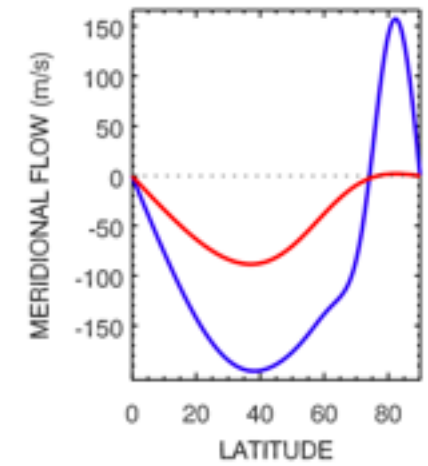
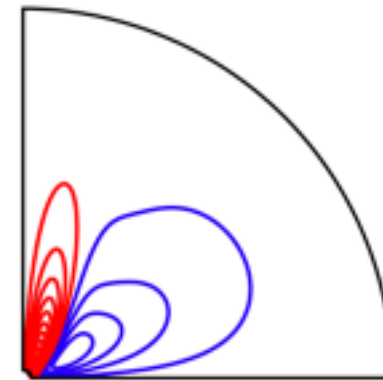
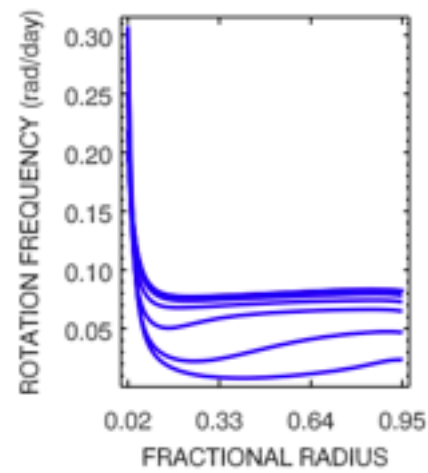
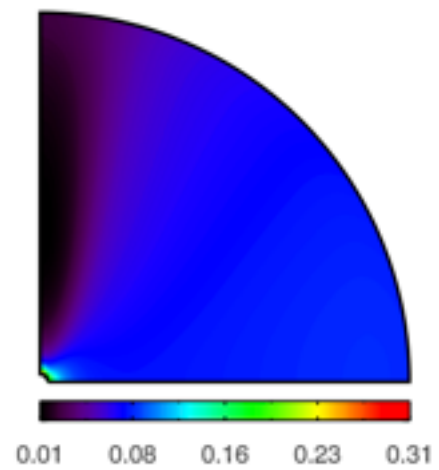
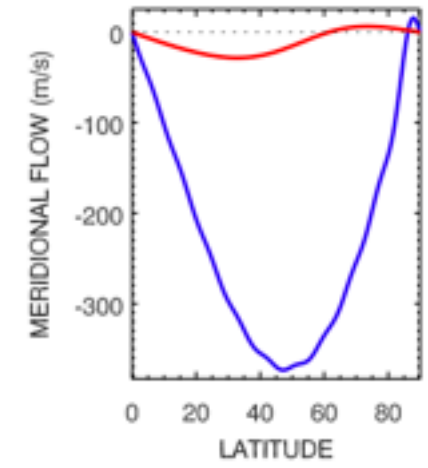
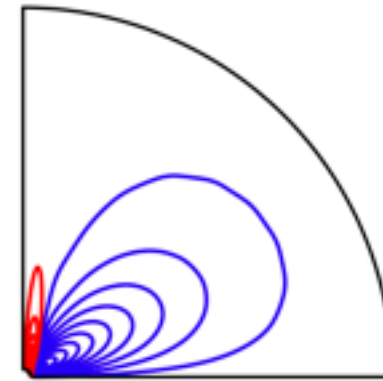
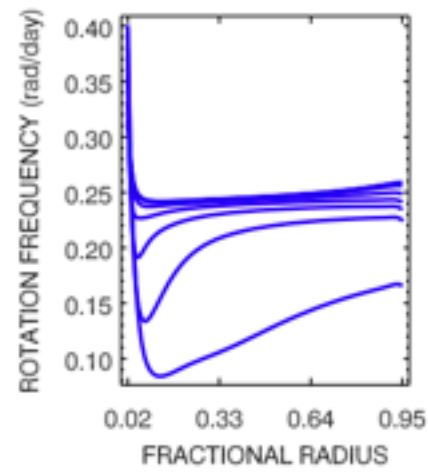
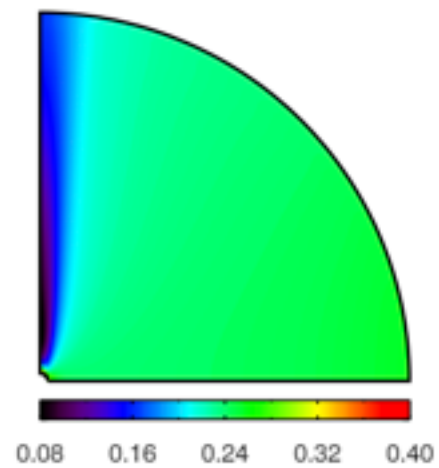
Giant star

11.78 Gyr
37.8 R_{\odot}
386 L_{\odot}

P=25d

P=80d

P=250d



Conclusions

- mean field model reproduces solar differential rotation and surface meridional flow
- always solar-type DR and meridional flow for MS stars
- surface DR and meridional flow of lower MS stars depend more on temperature than rotation period
- stagnation point of MF close to the bottom of CZ
- return flow not slow
- main sources of DR: Reynolds stress and baroclinic flow
- thermal wind balance does not hold near boundaries
- meridional flow in rapidly rotating convection zones is
 - driven by shear in boundary layers
 - concentrated in boundary layers
 - fast