

Thoughts on Neutralino Dark Matter



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Old Results in Dark Matter

Observations

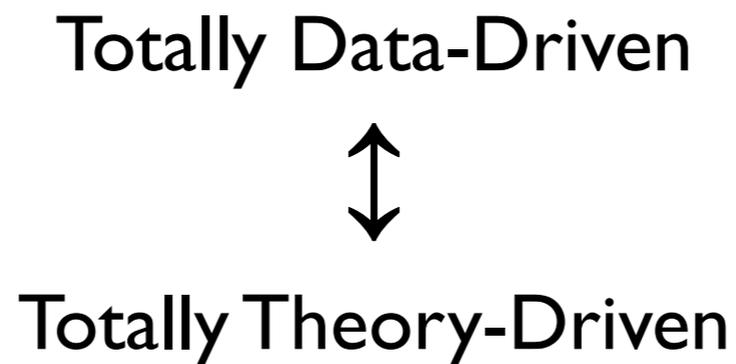
- Galactic Rotation Curves
- Cluster Dynamics (incl. collisions)
- Velocity dispersions of galaxies - dark matter extends beyond the visible matter
- Weak Gravitational Lensing (distribution of dark matter)
- CMB (+ Type IA SNe, plus BAO) all agree on LambdaCDM
- Structure Formation

Summary

- Some explanation is necessary for observed gravitational phenomena.
- It's largely non-relativistic (cold).
- Its abundance is $\Omega_{\text{DM}} \approx 0.26$.
- It's stable or very long-lived.
- It's non-baryonic (BBN+CMB, structure).
- It's neutral (heavy isotope abundances).

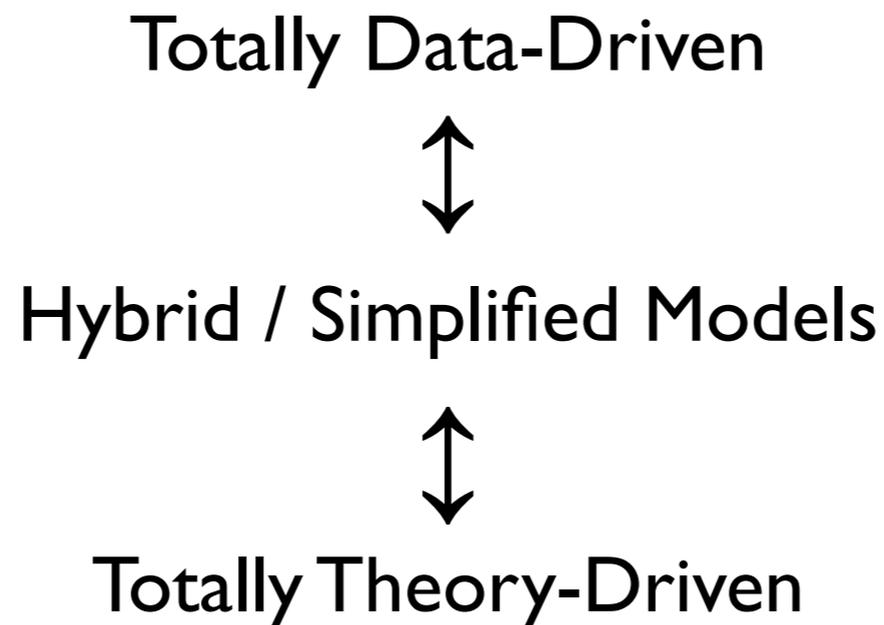
Current Situation

- **Abundance of experimental data!**
 - ▶ We're exploring dark matter with unprecedented and growing precision - both experimentally and theoretically.
- **Theoretical approaches:**



Current Situation

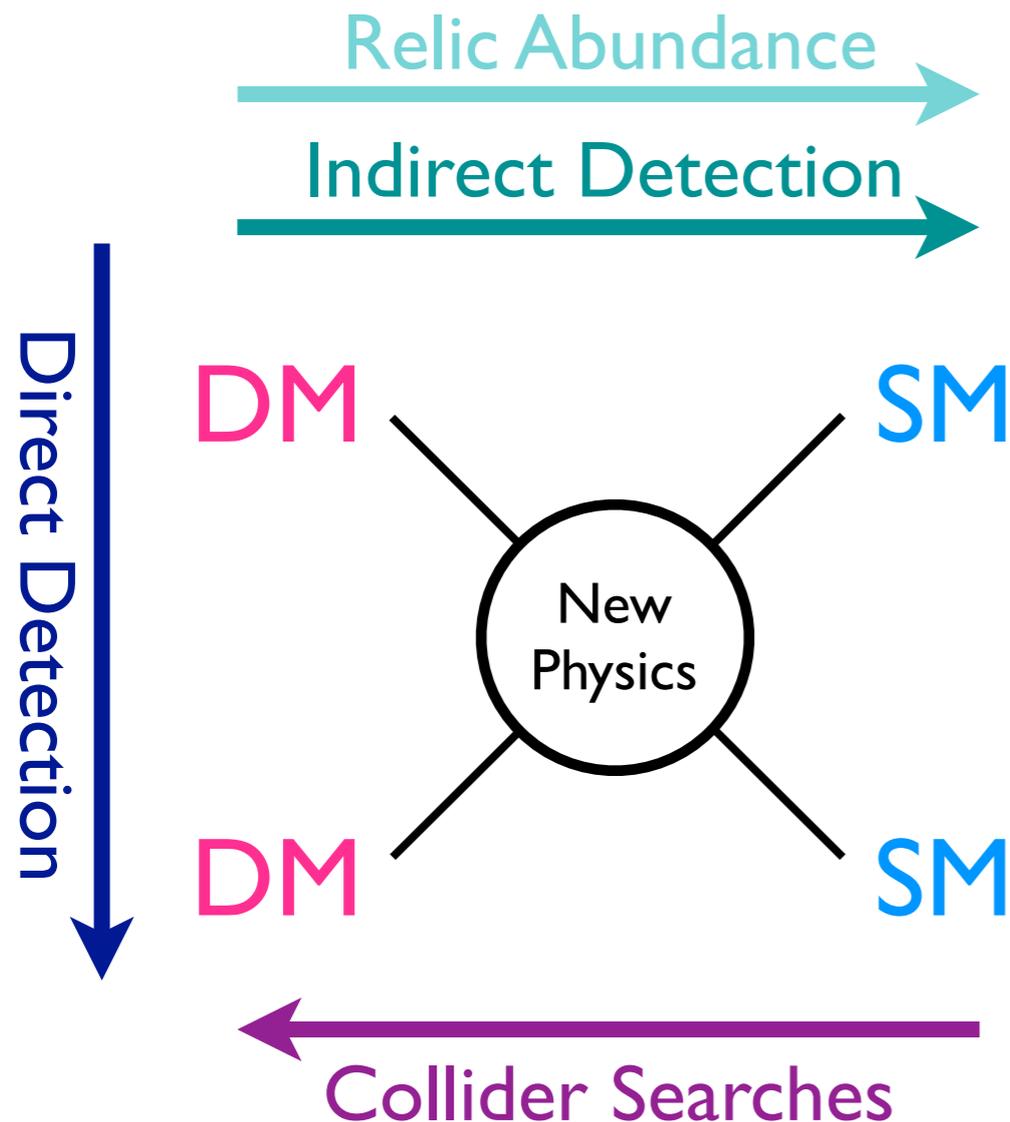
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 - ▶ We're exploring dark matter with unprecedented and growing precision - both experimentally and theoretically.
- Theoretical approaches:



Effective Theories

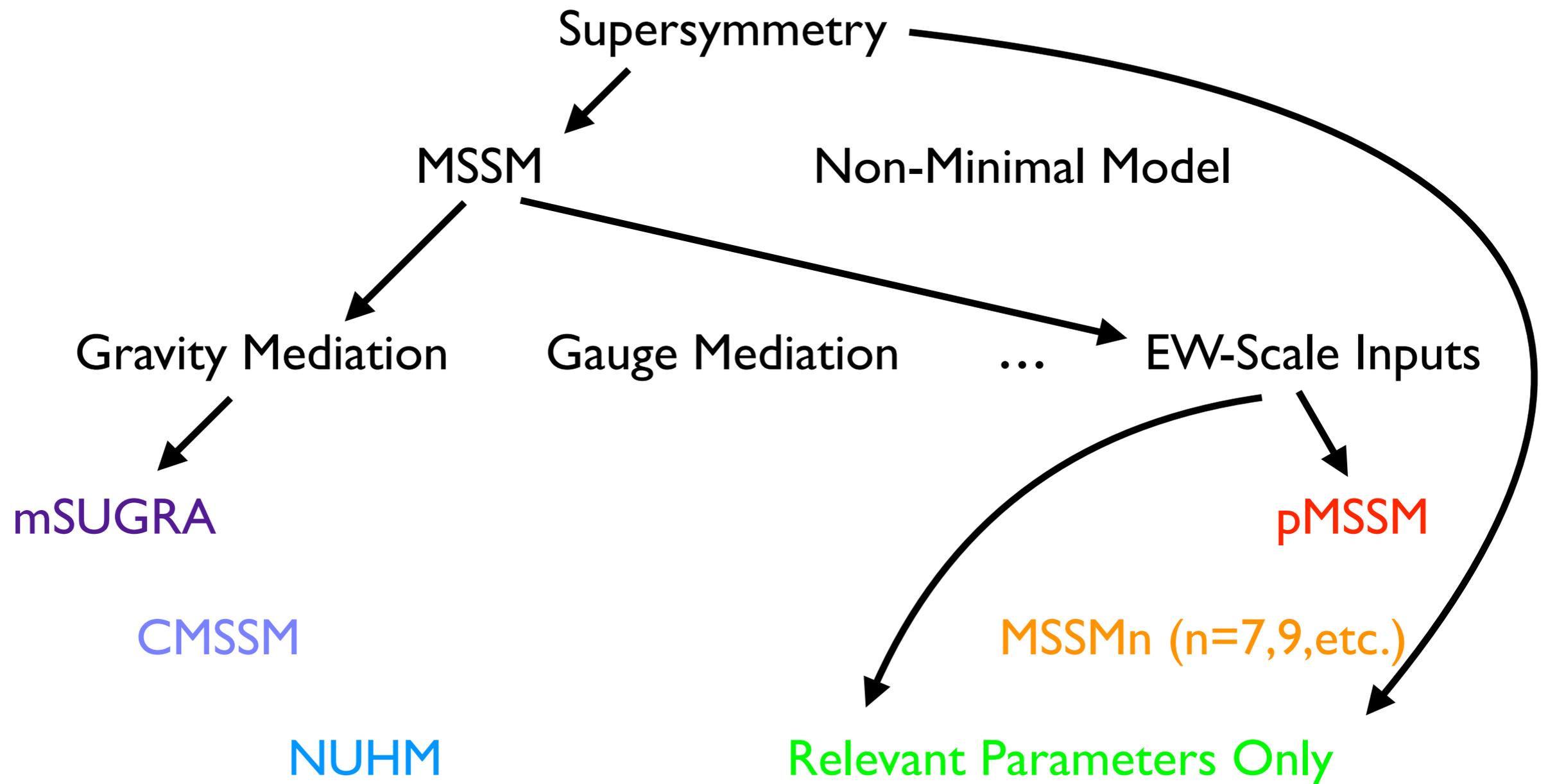
(McCabe talk on Monday)

$$\mathcal{L} = \mathcal{L}_{SM} + i\bar{X}\gamma^\mu\partial_\mu X - M_X\bar{X}X + \sum_q \sum_{i,j} \frac{G_{qij}}{\sqrt{2}} [\bar{X}\Gamma_i^X X] [\bar{q}\Gamma_q^j q]$$



- Idea: Reduce DM-SM interaction to a contact interaction.
- Universe of possible interactions is small (can enumerate)
- Utility in evaluating complementarity of detection techniques (good)
- Range of validity (careful)

Fundamental Theory

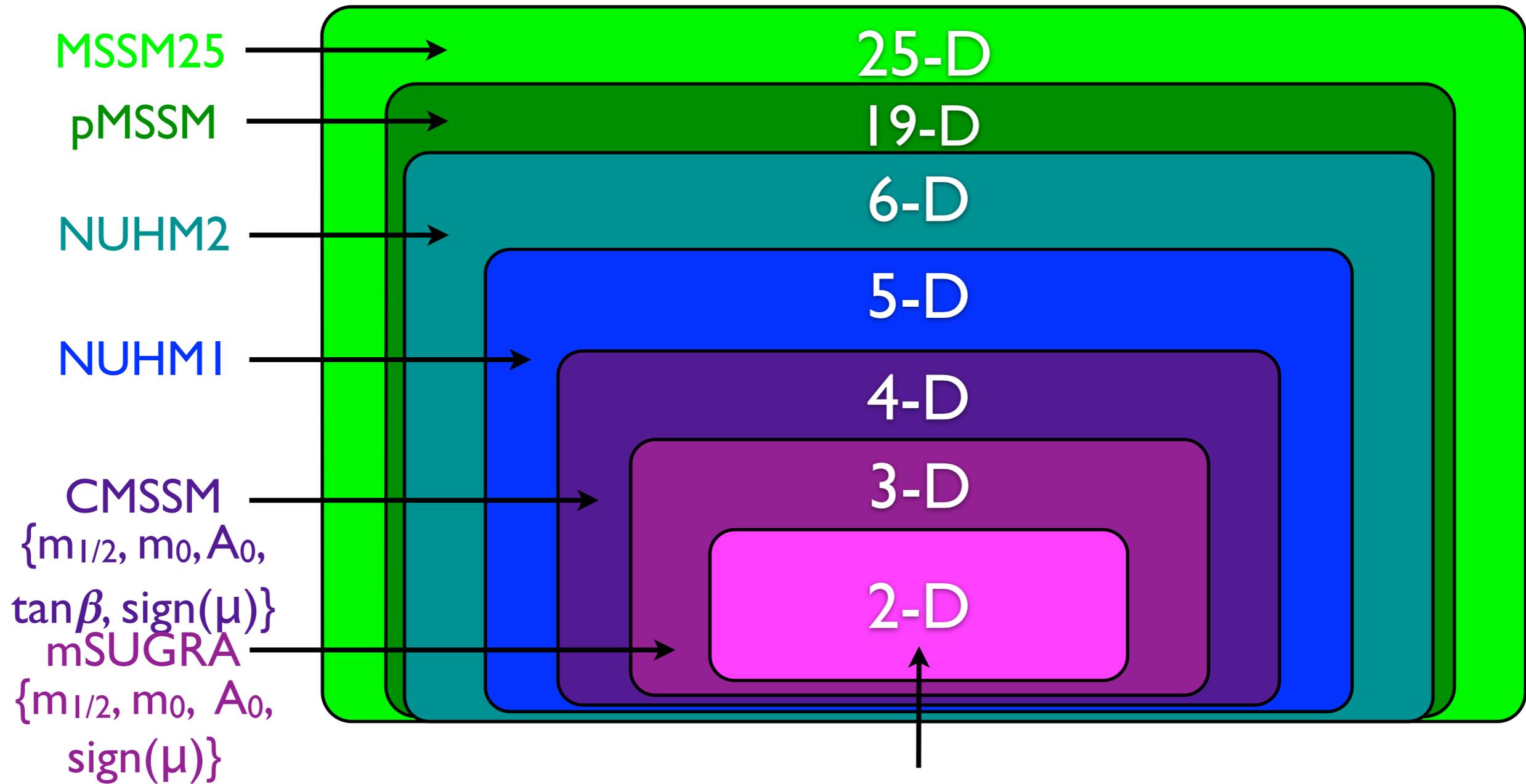


SUSY Dark Matter

1. **What is predicted** within the SUSY framework?
 - ▶ specific realization or more general possibilities
2. **What are the data really telling us?**
 - ▶ Priors on model → different interpretations
3. **When will we know for sure?**
 - ▶ Direct Dark Matter Searches

$$\tilde{\chi}_i = N_{i1}\tilde{B} + N_{i,2}\tilde{W} + N_{i3}\tilde{H}_d + N_{i,4}\tilde{H}_u$$

Dimensionality

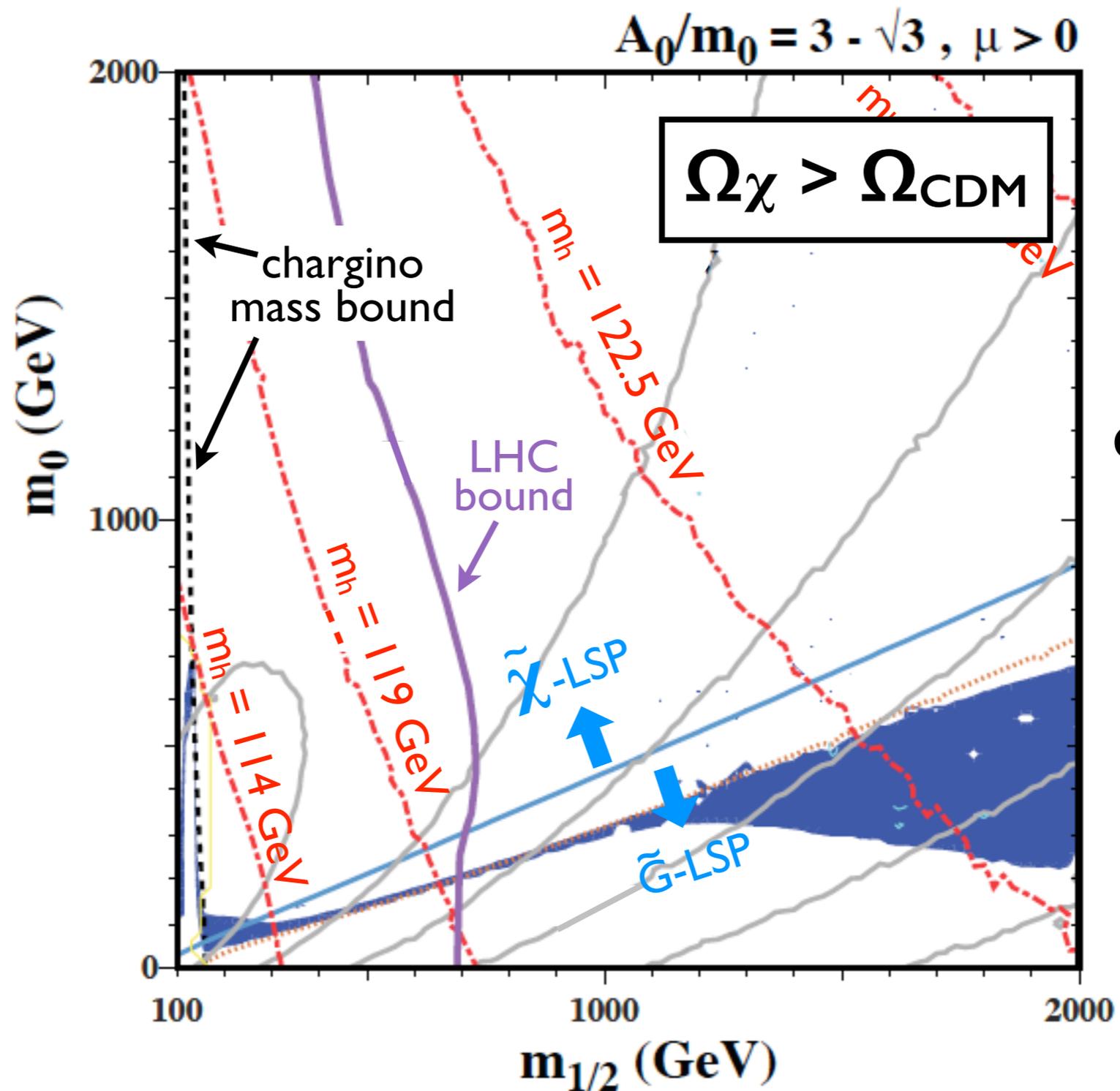


CMSSM
 $\{m_{1/2}, m_0, A_0, \tan\beta, \text{sign}(\mu)\}$
 mSUGRA
 $\{m_{1/2}, m_0, A_0, \text{sign}(\mu)\}$

Polonyi mSUGRA

$$A_0 = (3 - \sqrt{3})m_{3/2} \rightarrow \{m_{1/2}, m_0, \text{sign}(\mu)\}$$

Polonyi mSUGRA

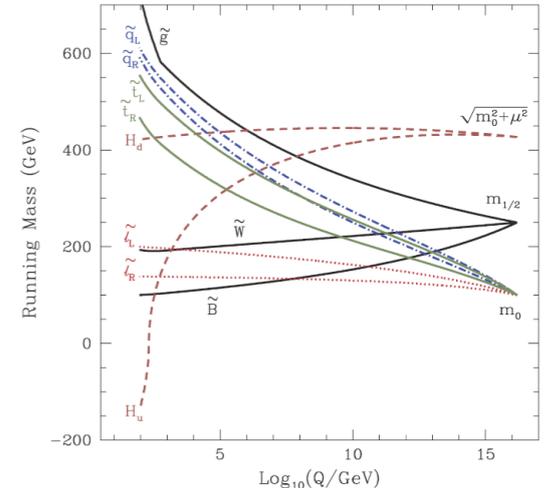


Increase
dimensionality?

Ellis, Luo, Olive,
Sandick (2013)

Universality Scale

- Input universality scale, M_{in} , assumed to be M_{GUT}
- Could be larger: “superGUT”
- SUSY breaking and mediation characterized by Planck or string scale



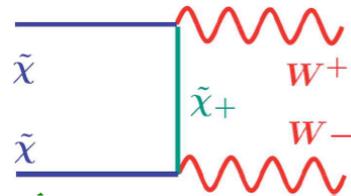
Polonsky & Pomarol (1994)

For recent analyses, see Ellis, Mustafayev, & Olive (2010,2011)

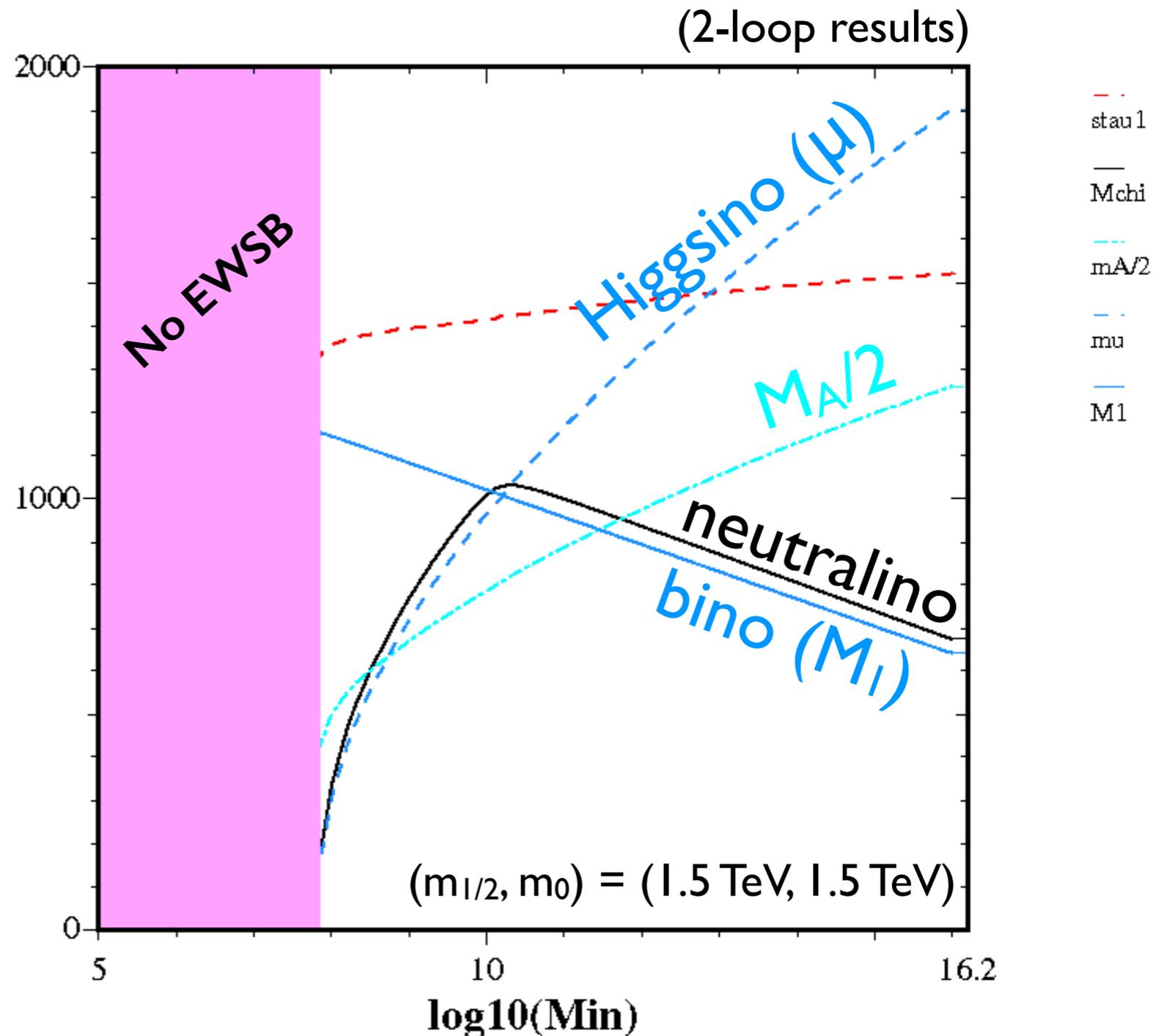
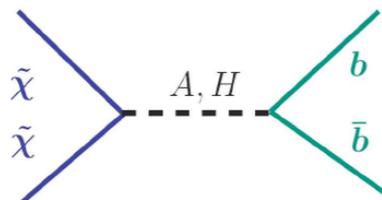
- Could be smaller: “subGUT/GUTless”, “Mirage”, or “TGM”
 - Ellis, Olive, & Sandick (2006, 2007, 2008); Ellis, Luo, Olive, & Sandick (2013)
 - Choi et al. (2004, 2005), Kachru et al. (2003), and others
 - Monaco et al. (2011)
- Lowest dynamical scale in the Polonyi/hidden sector where SUSY is broken, or scale of interactions that transmit breaking to observable sector

Dark Matter Abundance

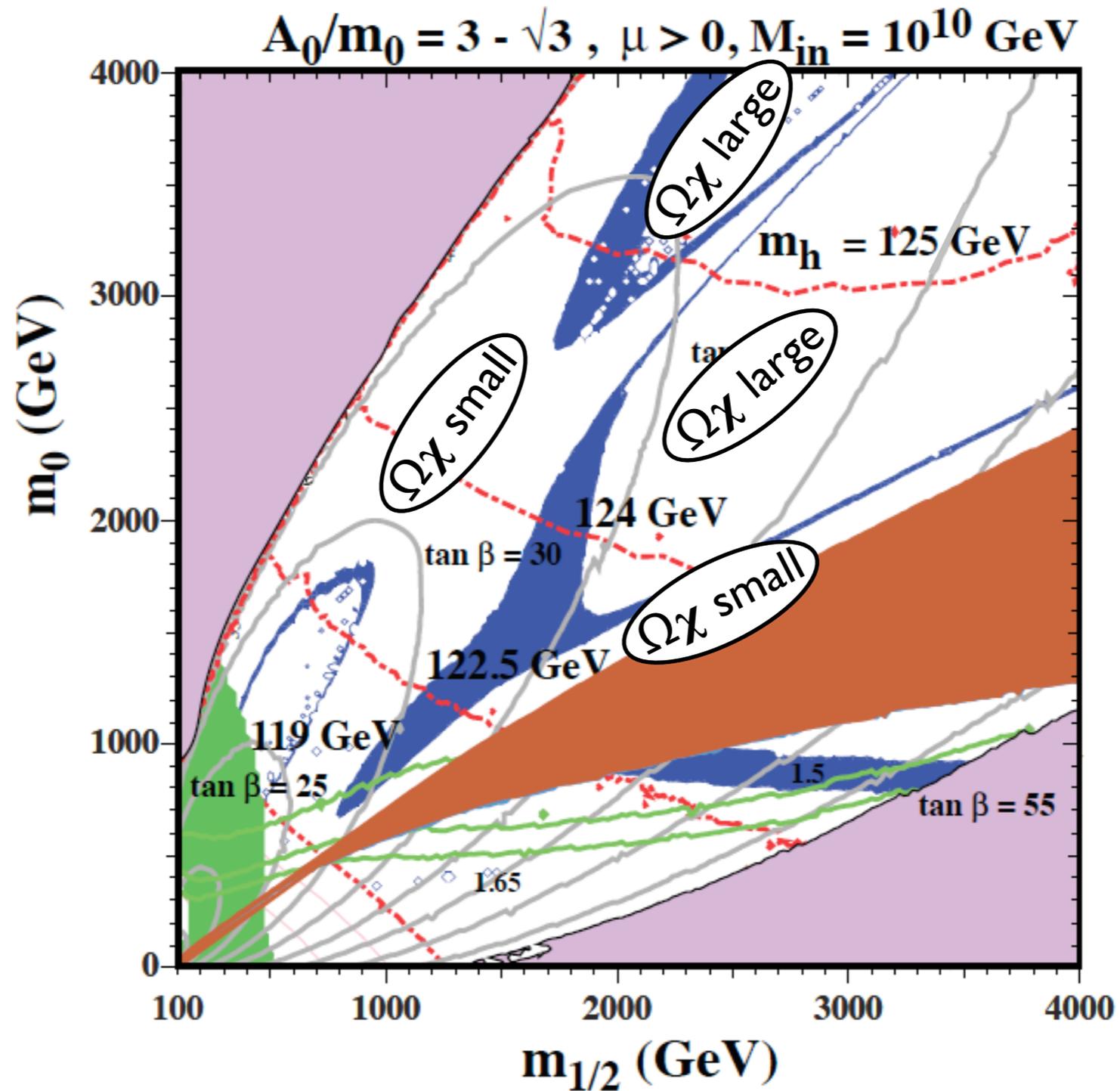
- neutralino LSP becomes Higgsino-like at low M_{in}



- m_A decreases with M_{in} → appearance of rapid annihilation funnel



sub-GUT mSUGRA



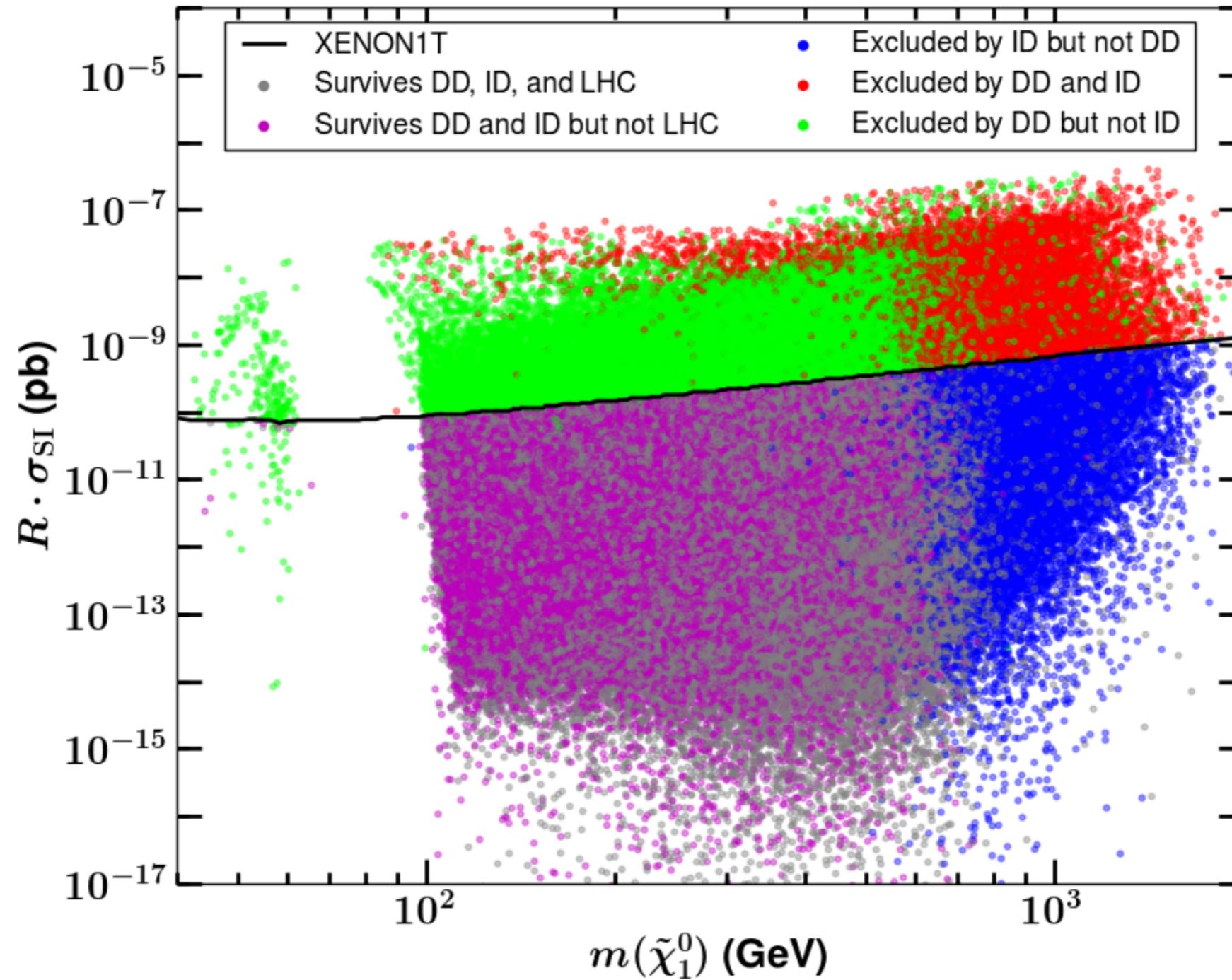
Ellis, Luo, Olive,
Sandick (2013)

To Higher Dimensions...

- There are still viable, few-parameter models motivated by high-scale physics.
- Strength: one of these models might actually describe our Universe!
- Strength: understand how observables change in the parameter space!
- Weakness: may be missing important model classes
- Higher dimensional models more fully explore the possible combinations of observables (if sampling of the model space is adequate!).

pMSSM

Cahill-Rowley et al. (2013)

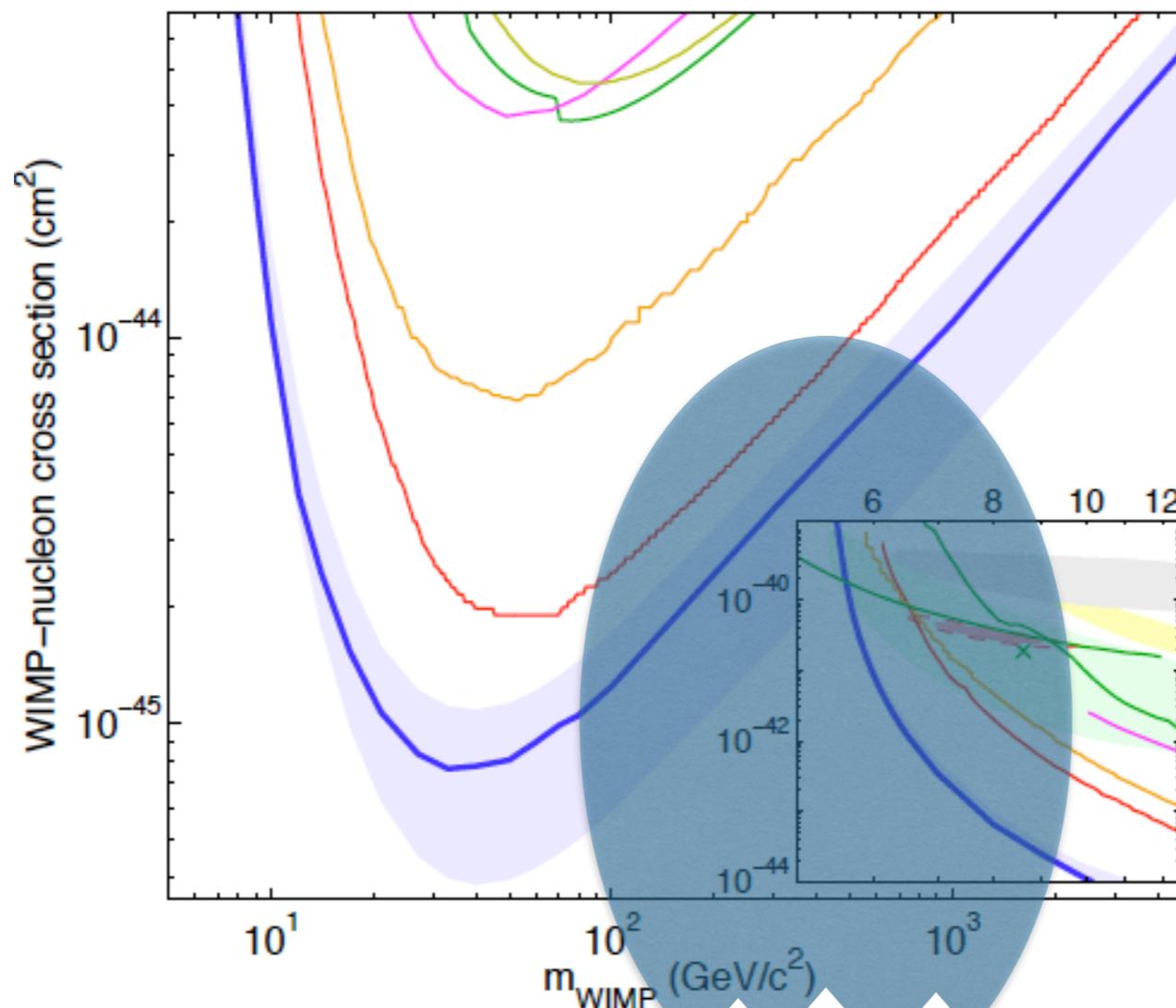


**3. When will we
know for sure?**

Future Prospects

- Timeline for discovery/exclusion?

Akerib et al. (LUX Collaboration), 2013

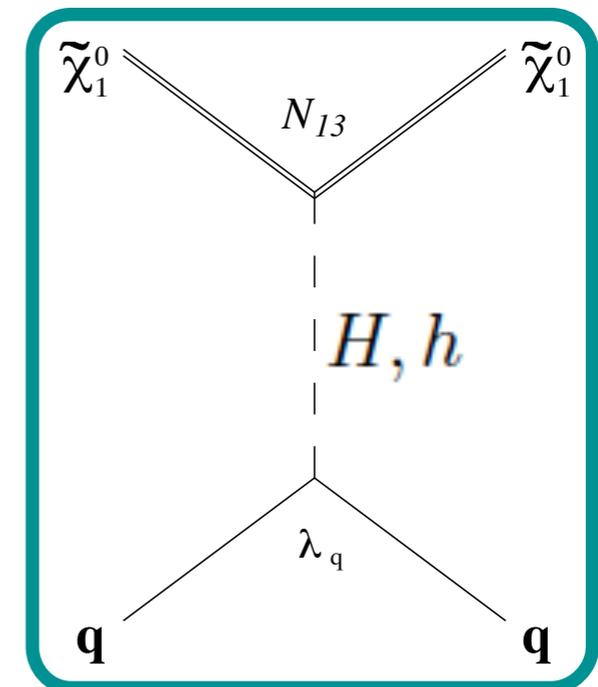
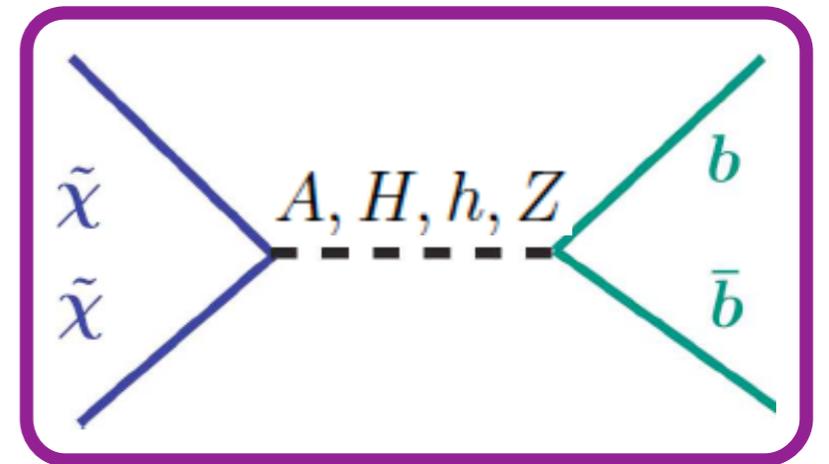


Could answer within a low-dimensional model (not general), or within the MSSM (not conclusive).

Simplified models can help you construct a definite, model-independent answer.

Resonance Models

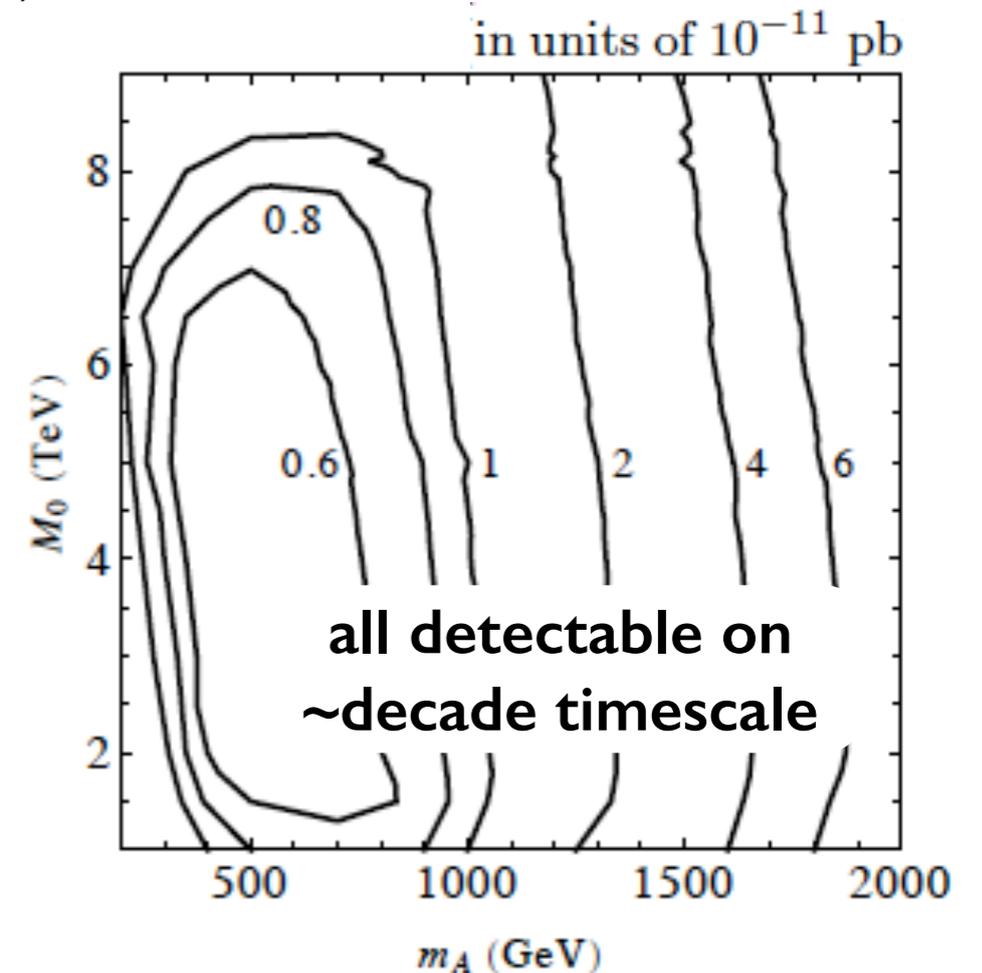
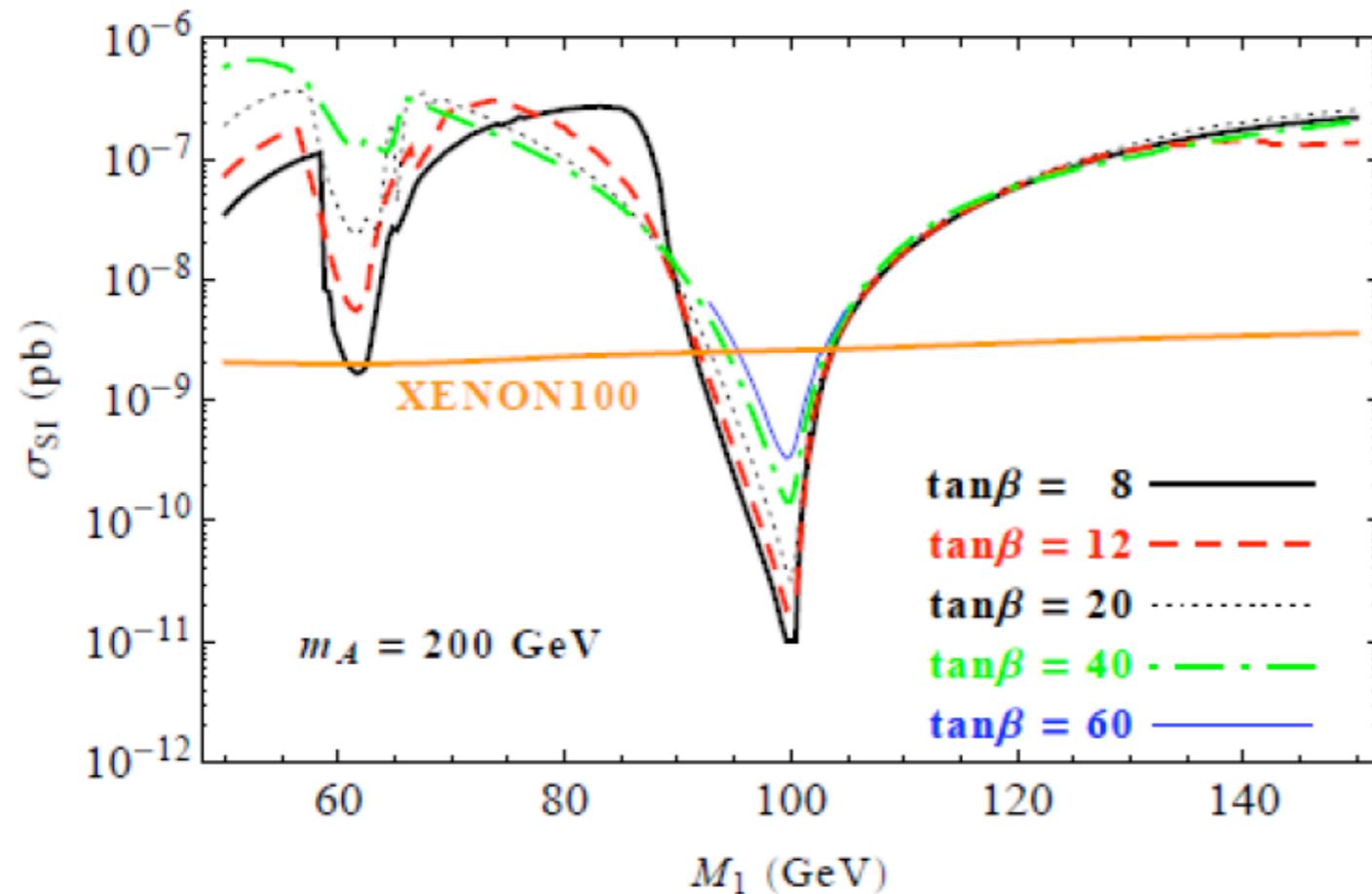
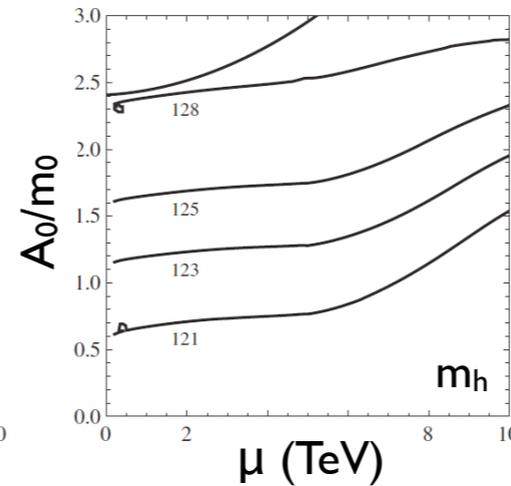
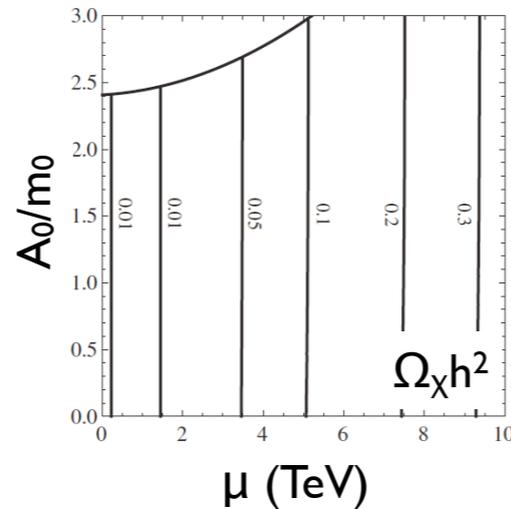
- **Neutralino:** $\tilde{\chi}_1^0 = \underbrace{\alpha \tilde{B}}_{M_1} + \beta \tilde{W}^0 + \underbrace{\gamma \tilde{H}_d^0 + \delta \tilde{H}_u^0}_{\mu}$
- **s-channel resonance annihilations** occur when $2m_{\tilde{\chi}_1^0} \approx m_{A,H,h,Z}$
- As $(\sigma_{ann.} v)$ increases, $\Omega_{\tilde{\chi}_1^0}$ decreases
 - If $\Omega_{\tilde{\chi}_1^0}$ too large, increase Higgsino content: μ
- **Scattering with quarks** is governed by $M_1, \mu, m_A, \tan \beta$
- **Relevant parameters:** $\{M_1, \mu, m_A, \tan \beta, M_0, A_0\}$



If DM abundance is achieved through a resonance, how small could σ_{SI} possibly be?

Hooper, Kelso, Sandick, & Xue, PRD 2013

- Relic Abundance: μ
- Higgs mass: A_0
- Free parameters: $(m_0, M_1, m_A, \tan\beta)$

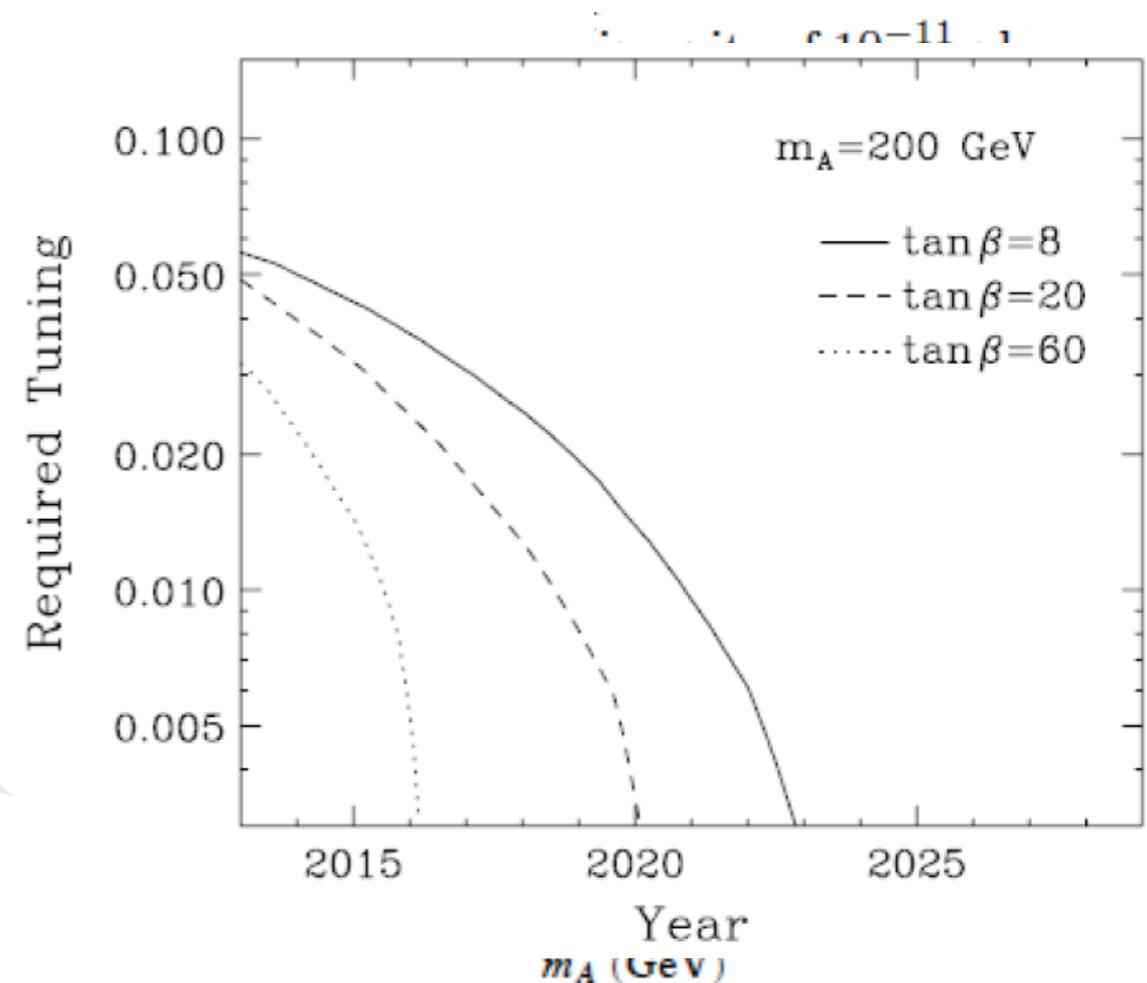
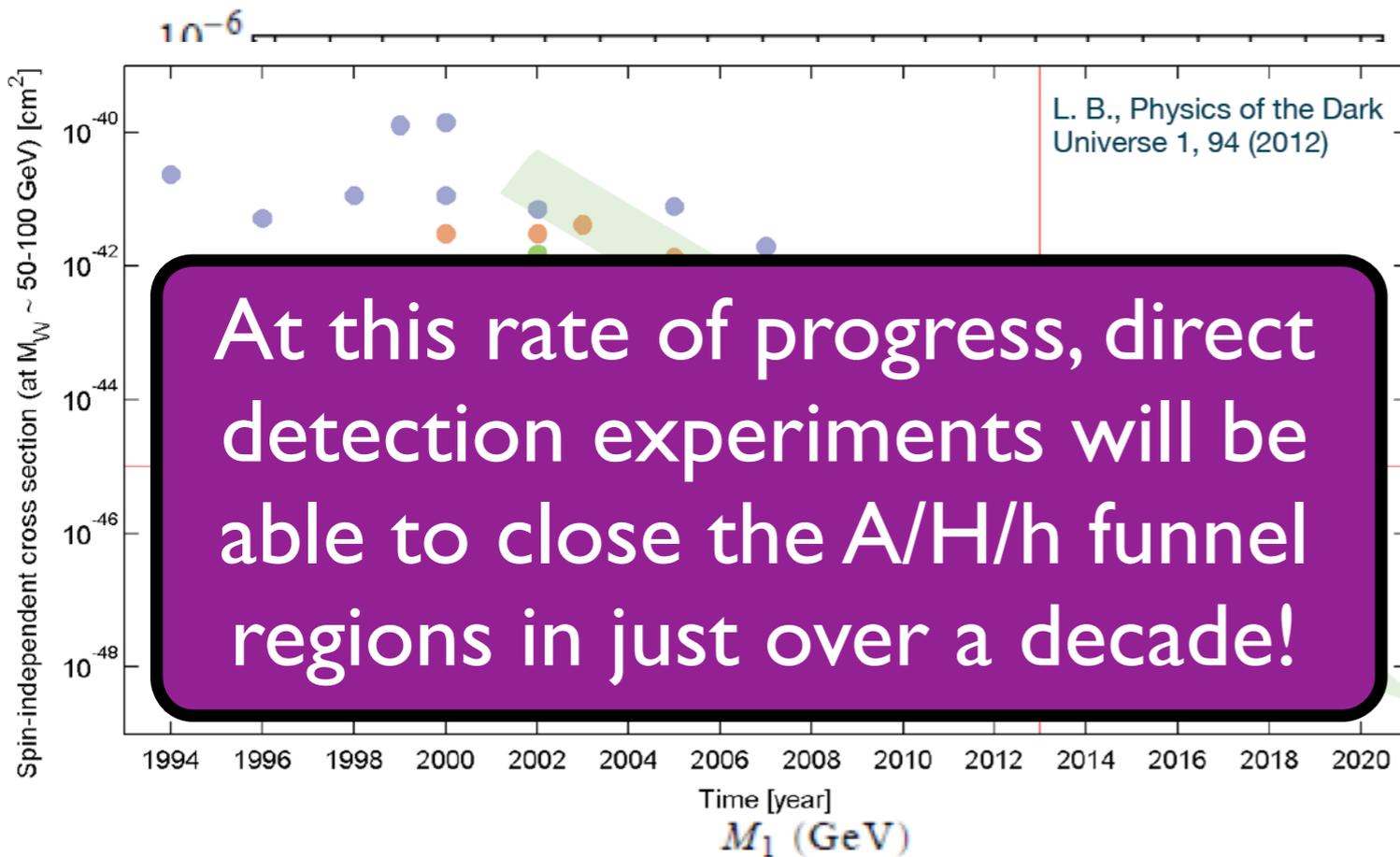


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- Free parameters: $(m_0, M_1, m_A, \tan\beta)$

If Nature is MSSM-like, and neutralino dark matter at a resonance makes up all the dark matter in the Universe, then direct detection experiments are pushing the resonance to be more and more exact.



Data-Driven SUSY

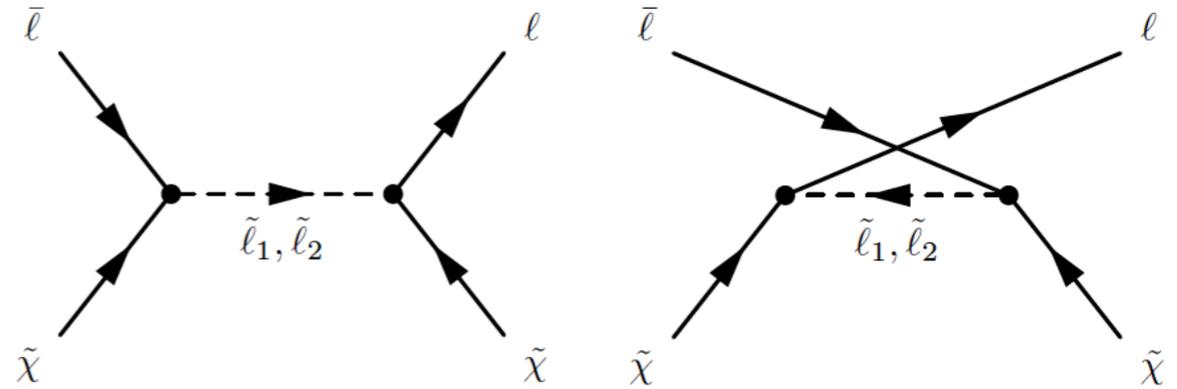
- 2. What are the data telling us?
 - ▶ Investigate parameter space near current constraints.
 - ▶ Dramatically different answers, depending on assumptions!
- What we really know about sparticles: sleptons, charginos, and 3rd gen. squarks heavier than ~ 100 GeV, 1st/2nd gen. squarks heavier than ~ 1.1 TeV, gluino heavier than ~ 1 TeV
- Other constraints: Higgs ~ 126 GeV, dark matter, rare B decays, electric dipole moments, anomalous magnetic moments
- Simple model: bino-like LSP and light sleptons (everything else heavy)
 - ▶ $m_{\tilde{\chi}_1}, m_{\tilde{l}_1}, m_{\tilde{l}_2}, \alpha, \varphi$

Light Sleptons

- Relic Abundance:

$$\langle \sigma v \rangle \sim c_0 + c_1 \left(\frac{T}{m_\chi} \right)$$

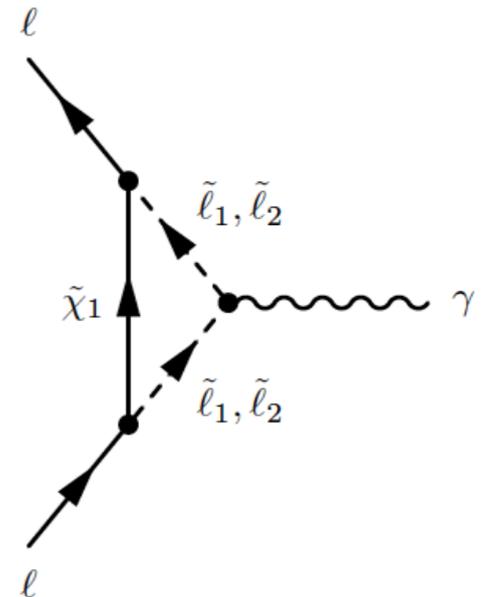
$$c_0 = \frac{m_{\tilde{\chi}}^2}{2\pi} g^4 Y_L^2 Y_R^2 \cos^2 \alpha \sin^2 \alpha \left(\frac{1}{m_{\tilde{\ell}_1}^2 + m_{\tilde{\chi}}^2} - \frac{1}{m_{\tilde{\ell}_2}^2 + m_{\tilde{\chi}}^2} \right)^2$$



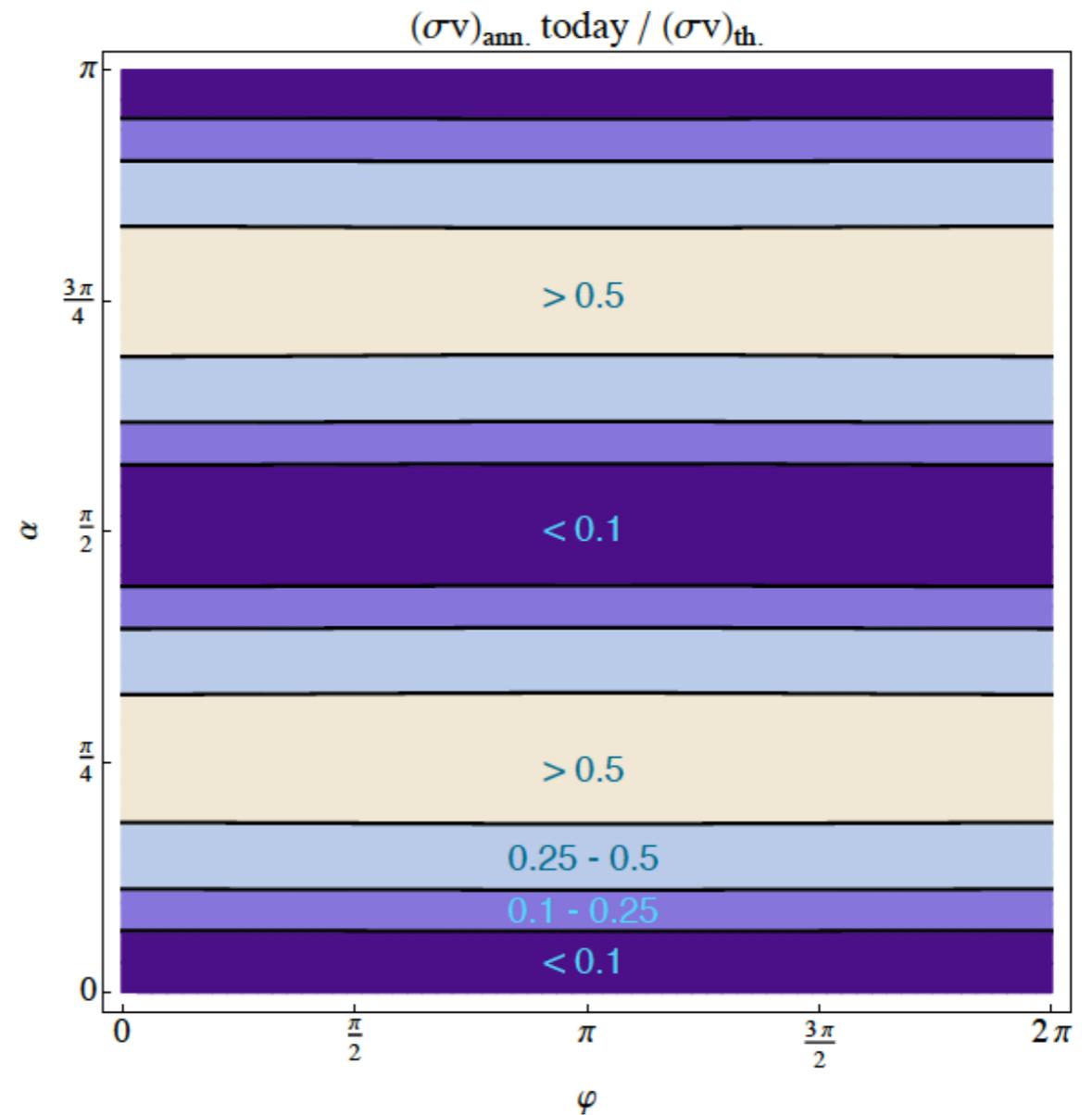
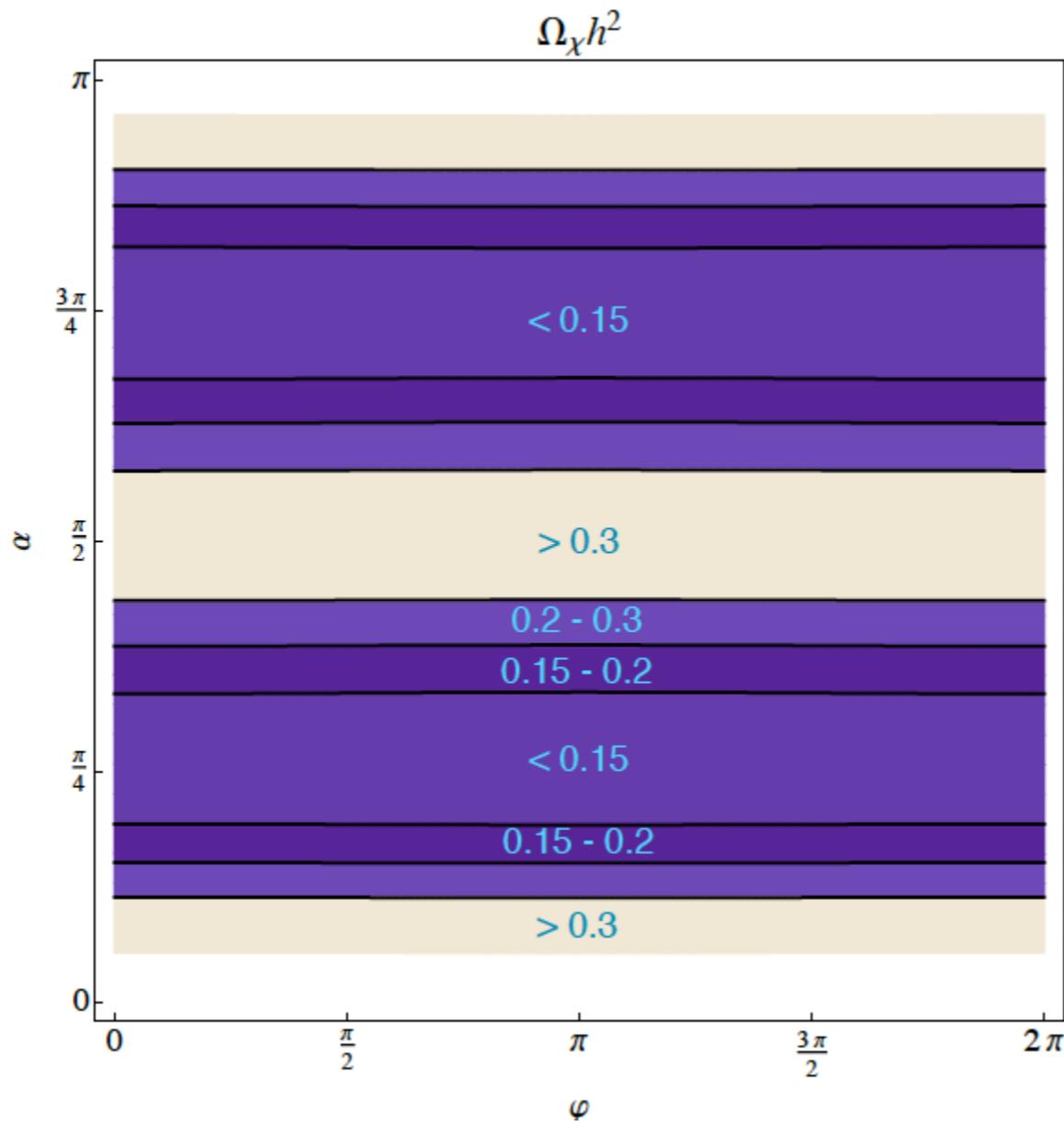
- Dipole Moments:

$$\Delta a = \frac{m_\ell m_{\tilde{\chi}}}{4\pi^2 m_{\tilde{\ell}_1}^2} g^2 Y_L Y_R \cos \varphi \cos \alpha \sin \alpha \left[\frac{1}{2(1 - r_{\tilde{\ell}_i})^2} \left(1 + r_{\tilde{\ell}_i} + \frac{2r_{\tilde{\ell}_i} \ln r_{\tilde{\ell}_i}}{1 - r_{\tilde{\ell}_i}} \right) \right] - (\tilde{\ell}_1 \rightarrow \tilde{\ell}_2)$$

$$\frac{d}{|e|} = \frac{m_{\tilde{\chi}}}{8\pi^2 m_{\tilde{\ell}_1}^2} g^2 Y_L Y_R \sin \varphi \cos \alpha \sin \alpha \left[\frac{1}{2(1 - r_{\tilde{\ell}_i})^2} \left(1 + r_{\tilde{\ell}_i} + \frac{2r_{\tilde{\ell}_i} \ln r_{\tilde{\ell}_i}}{1 - r_{\tilde{\ell}_i}} \right) \right] - (\tilde{\ell}_1 \rightarrow \tilde{\ell}_2)$$



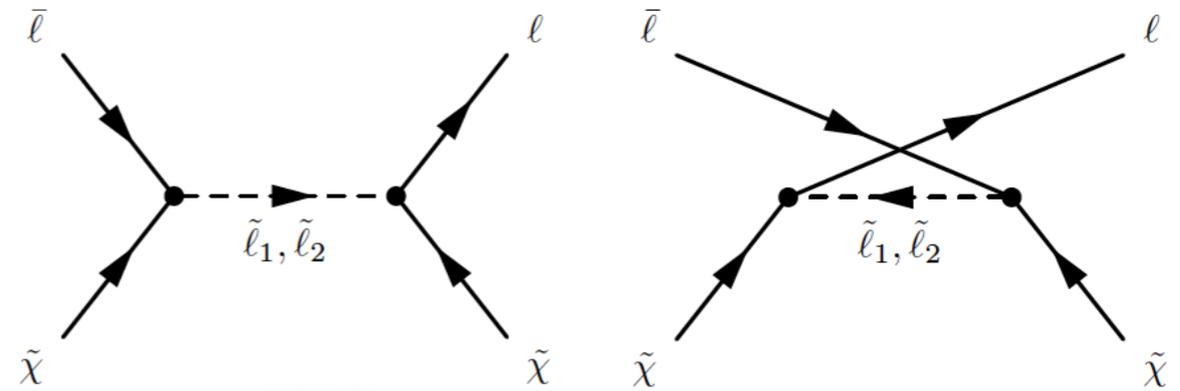
Light Smuons Scenario ($M_1 \neq M_2$)



Light Sleptons

- Annihilation Cross Section:

$$\langle \sigma v \rangle \sim c_0 + c_1 \left(\frac{T}{m_{\tilde{\chi}}} \right)$$



$$c_0 = \frac{m_{\tilde{\chi}}^2}{2\pi} g^4 Y_L^2 Y_R^2 \cos^2 \alpha \sin^2 \alpha \left(\frac{1}{m_{\tilde{\ell}_1}^2 + m_{\tilde{\chi}}^2} - \frac{1}{m_{\tilde{\ell}_2}^2 + m_{\tilde{\chi}}^2} \right)^2$$

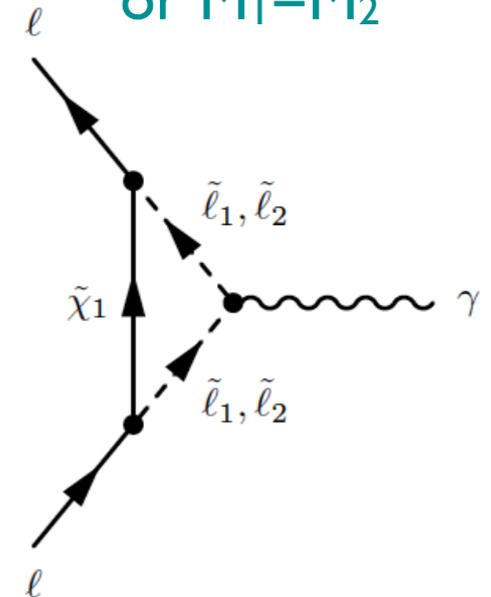
maximal when $\alpha = n\pi/4$, n odd

- Dipole Moments:

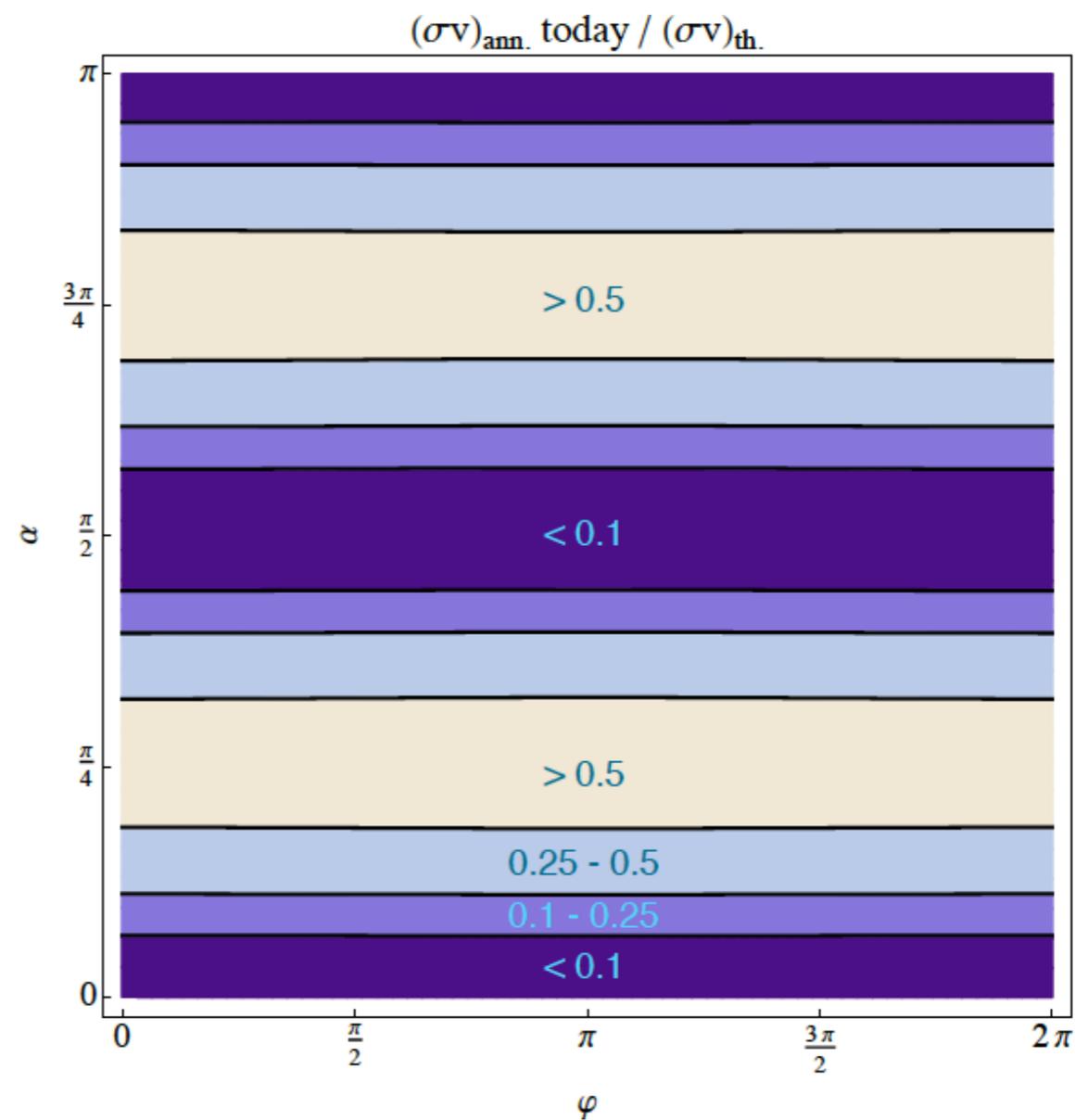
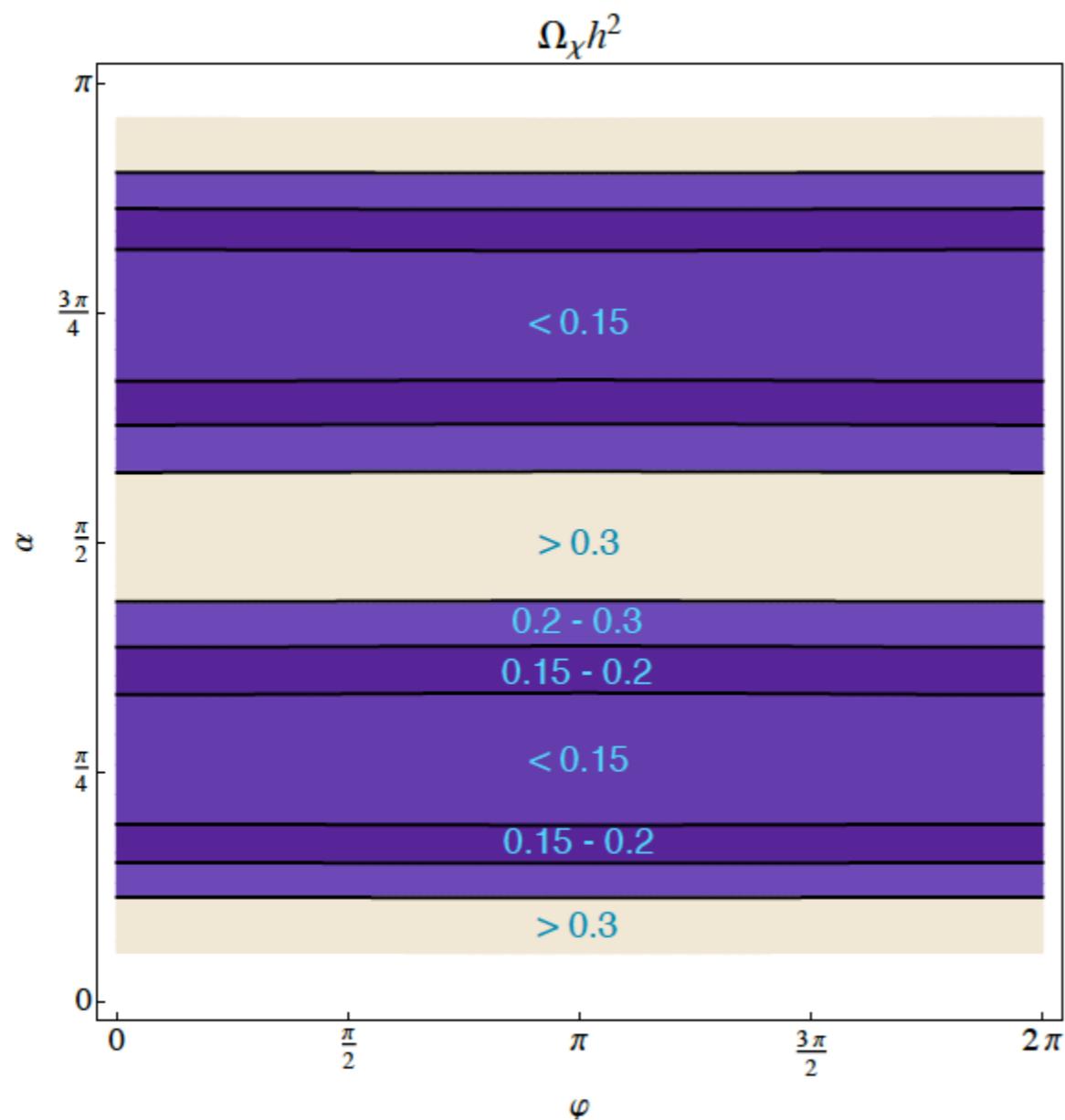
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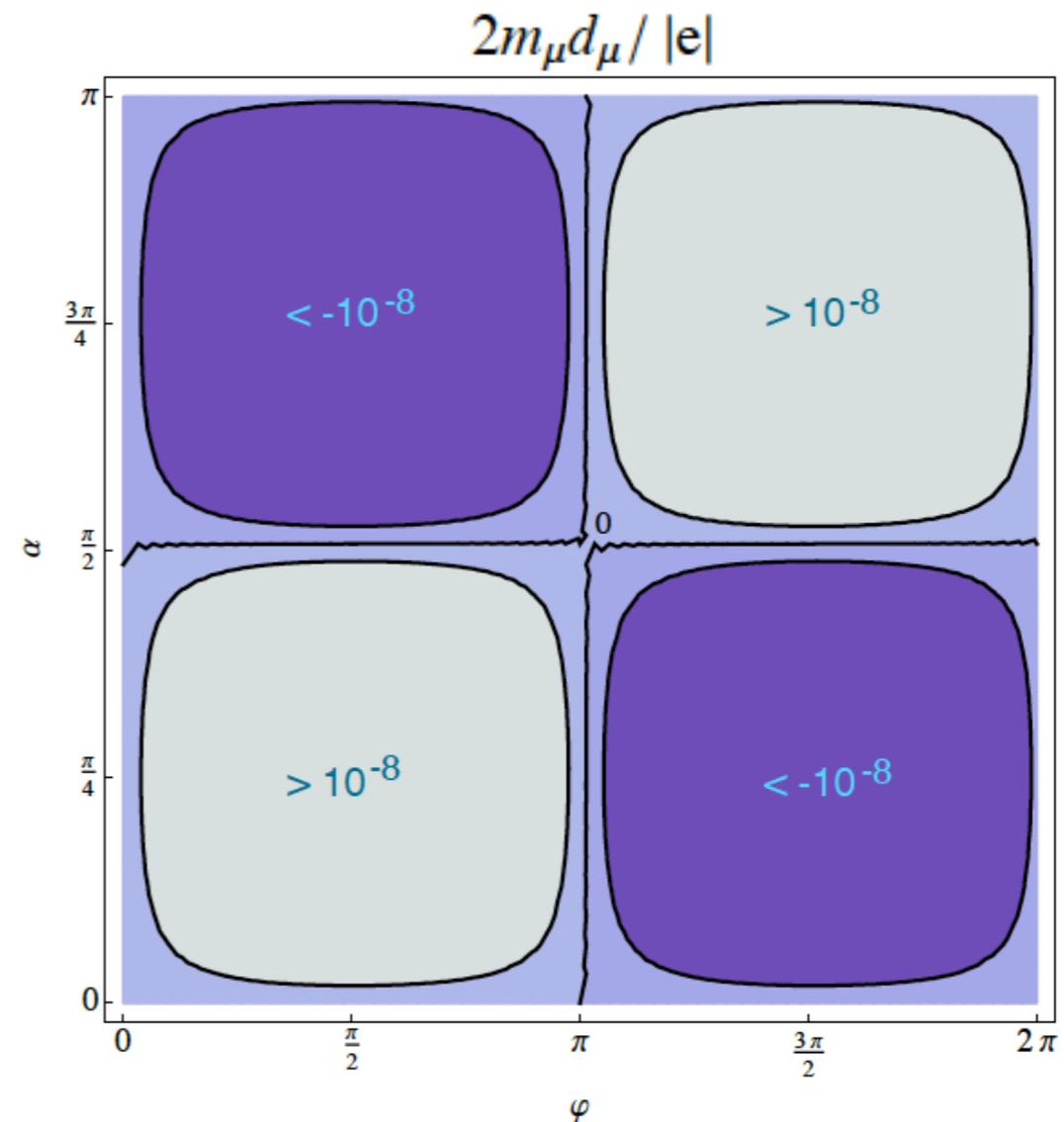
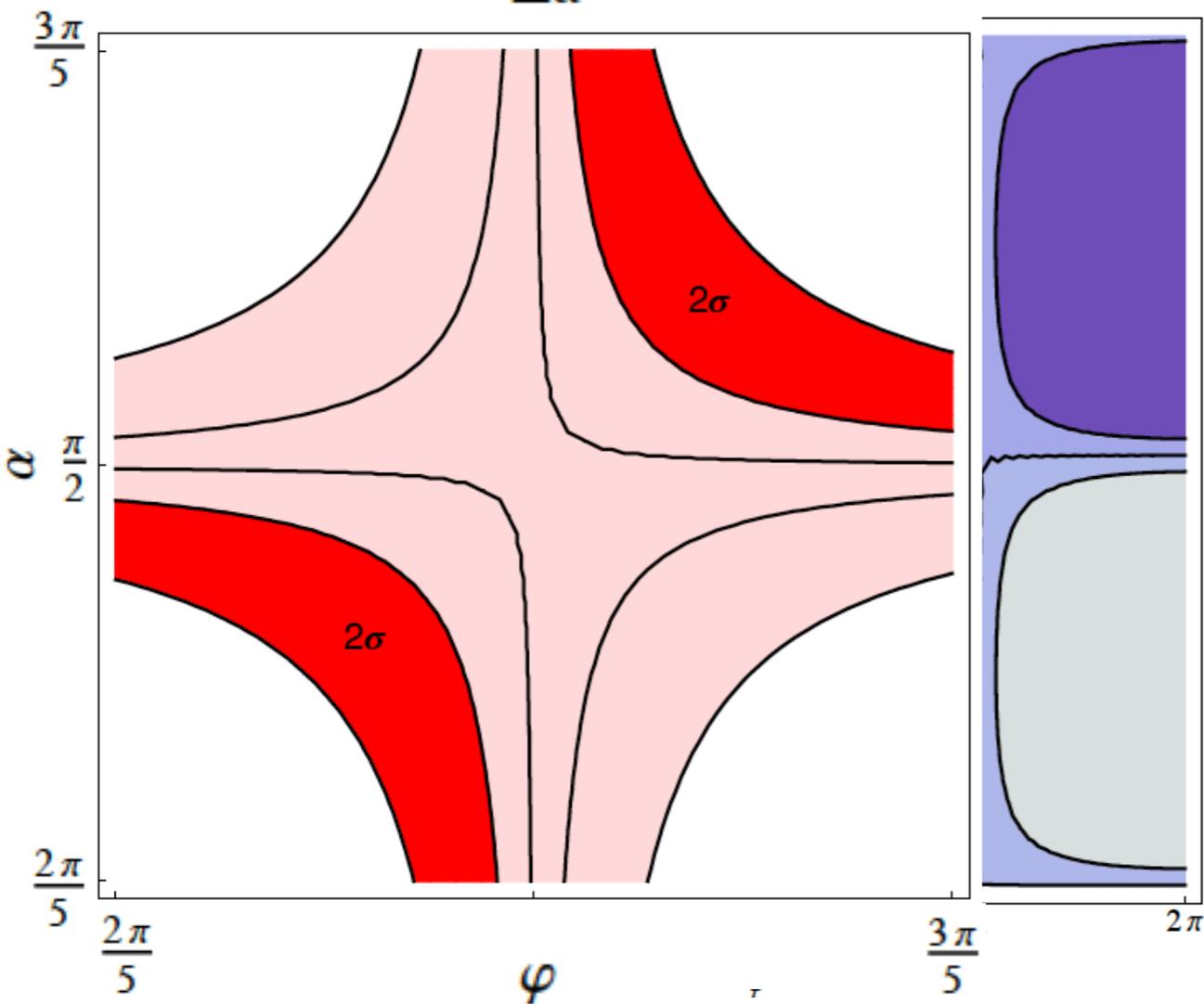
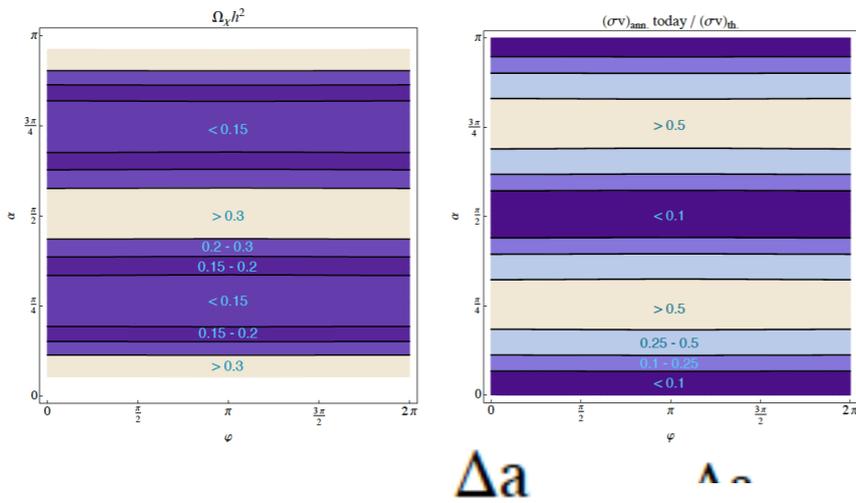
zero when $\alpha = n\pi/2$, n integer or $M_1 = M_2$



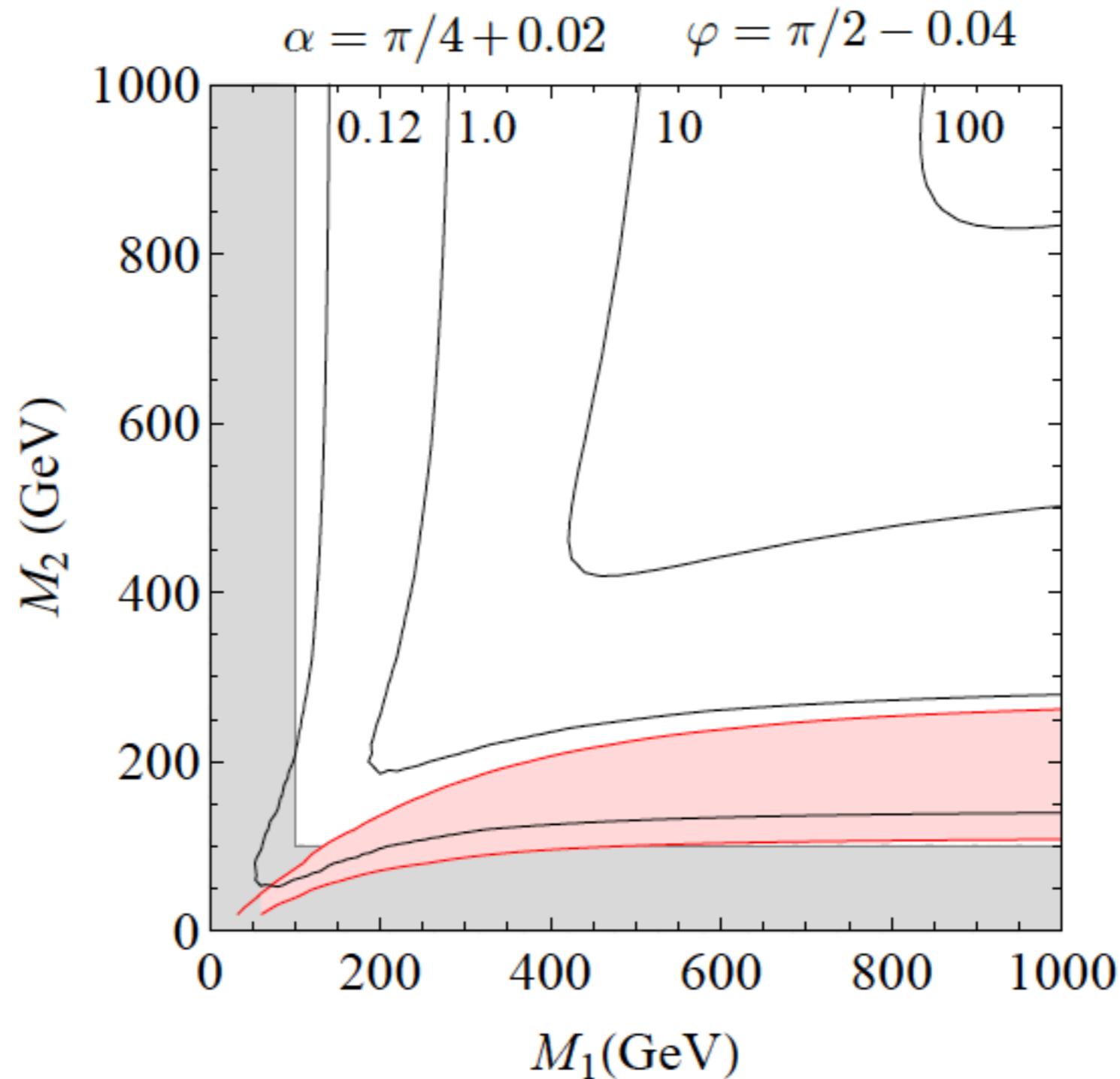
Light Smuons Scenario ($M_1 \neq M_2$)



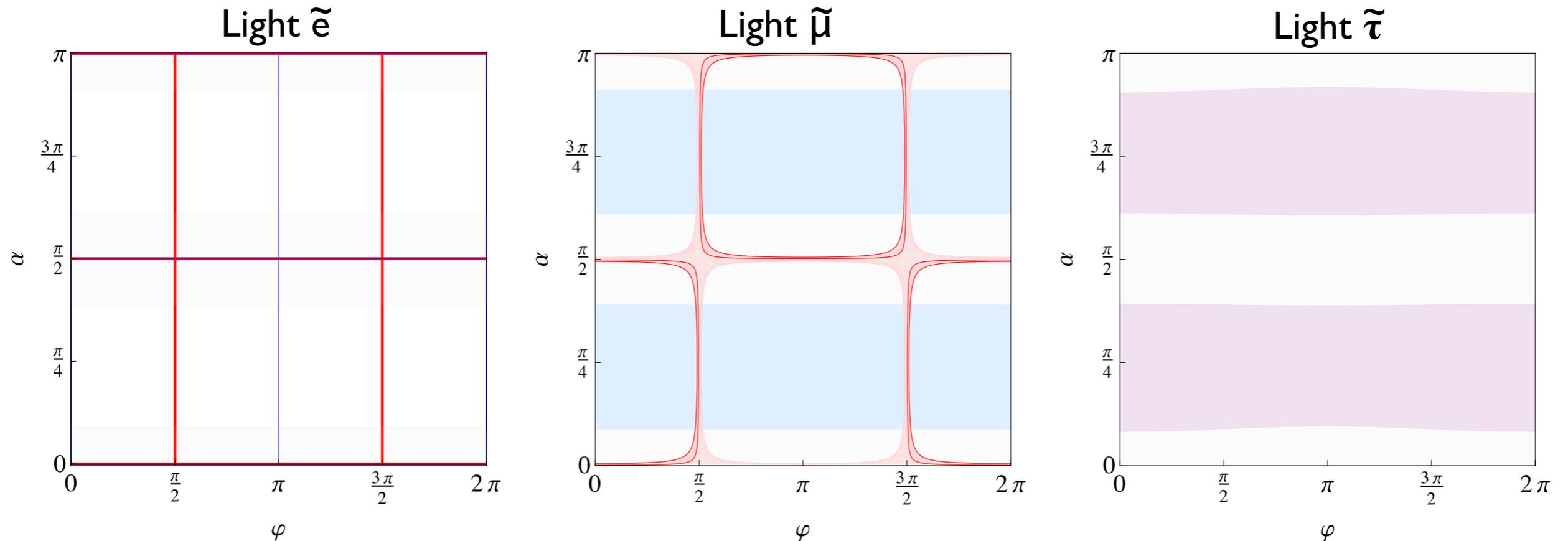
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Light Smuons Scenario



Light Sleptons ($M_1 \neq M_2$)

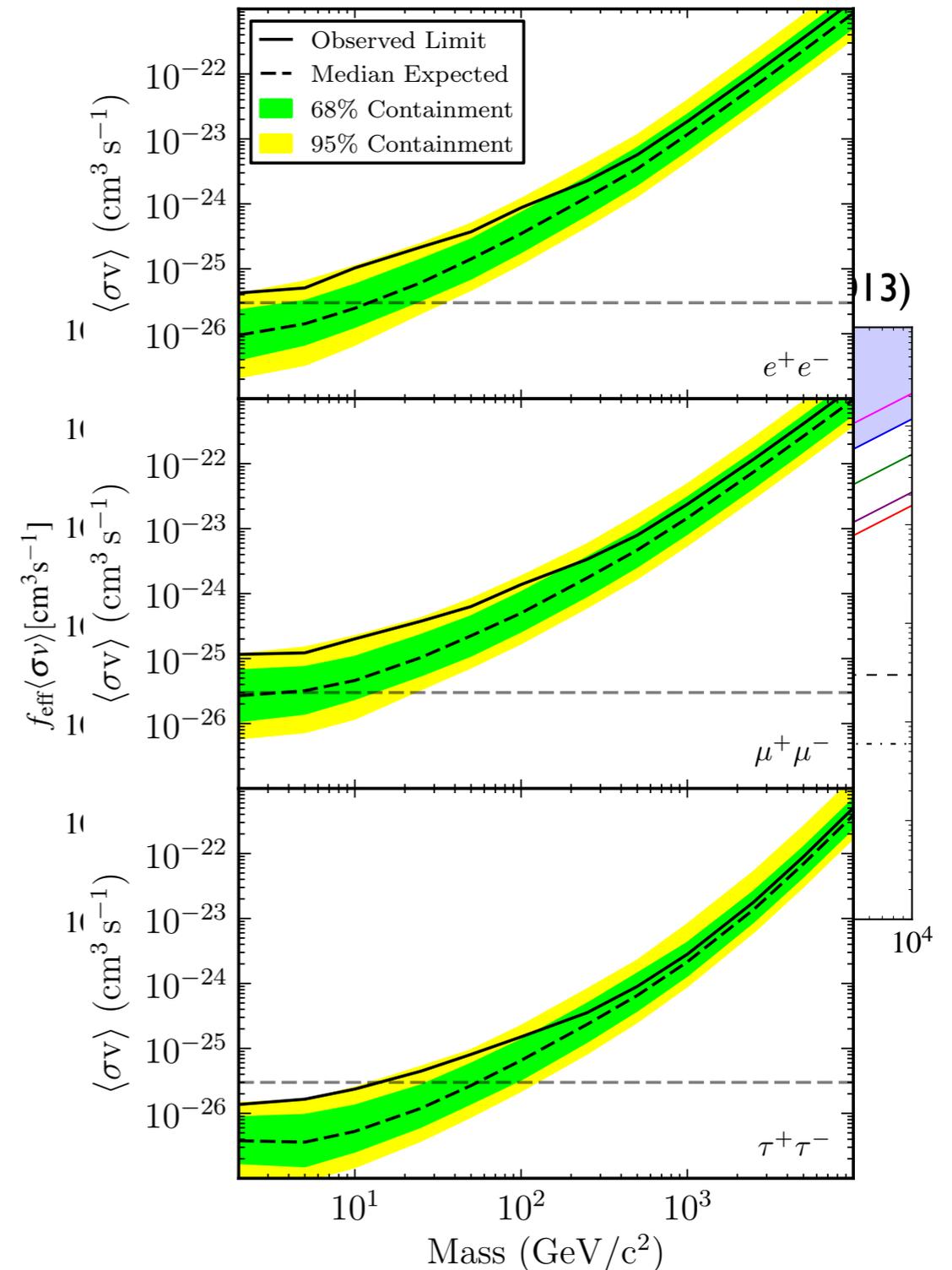


- If $M_1 = M_2$, dipole moments vanish, but too much dark matter.
- Light \tilde{e} : Angles must be tuned to $\alpha \approx 10^{-3}$ and $\varphi \approx 10^{-6}$, but relic abundance is too large
- Light $\tilde{\mu}$: Possible to obey dipole moment constraints (or explain Δa), and have thermal dark matter (for small range of φ)
- Light $\tilde{\tau}$: Relic abundance is the only constraint (see also Pierce, et al.; Hagiwara et al., 2013)

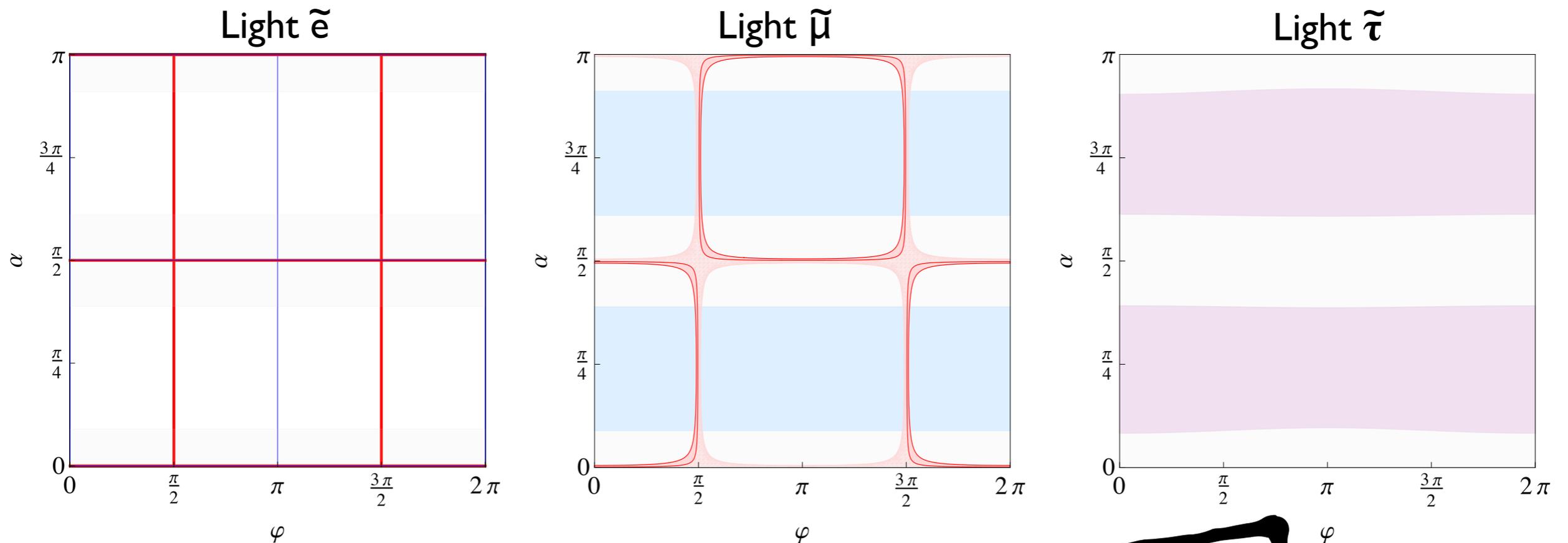
Other Signatures

Ackermann et al. (2014)

- Indirect Detection:
 - Best Case Scenario is annihilation to taus. Possibly within reach of Fermi (dSphs).
 - Annihilation to neutrinos: Unconstrained by dipole moments, but annihilation rate is small (relic abundance too large). Probably not detectable at IceCube or SK.
 - CMB: not currently constrained for annihilations to muons or taus, and will remain just out of reach, even for CVL experiment



Light Sleptons ($M_1 \neq M_2$)



- If $M_1=M_2$, dipole moments vanish, but too much dark matter
- Light \tilde{e} : Angles must be too large
- Light $\tilde{\mu}$: Angles must be too large
- Light $\tilde{\tau}$: Only constraint (see also Pierce, et al.; Hagiwara et al., 2013)

Viable scenarios with bino-like dark matter and light smuons/staus.

Summary

- Finite distinct ways to observe dark matter:
 - Abundance, Annihilation Today, Decay Today, Production at Colliders, Direct Detection
- A spectrum of theoretical approaches to particle dark matter

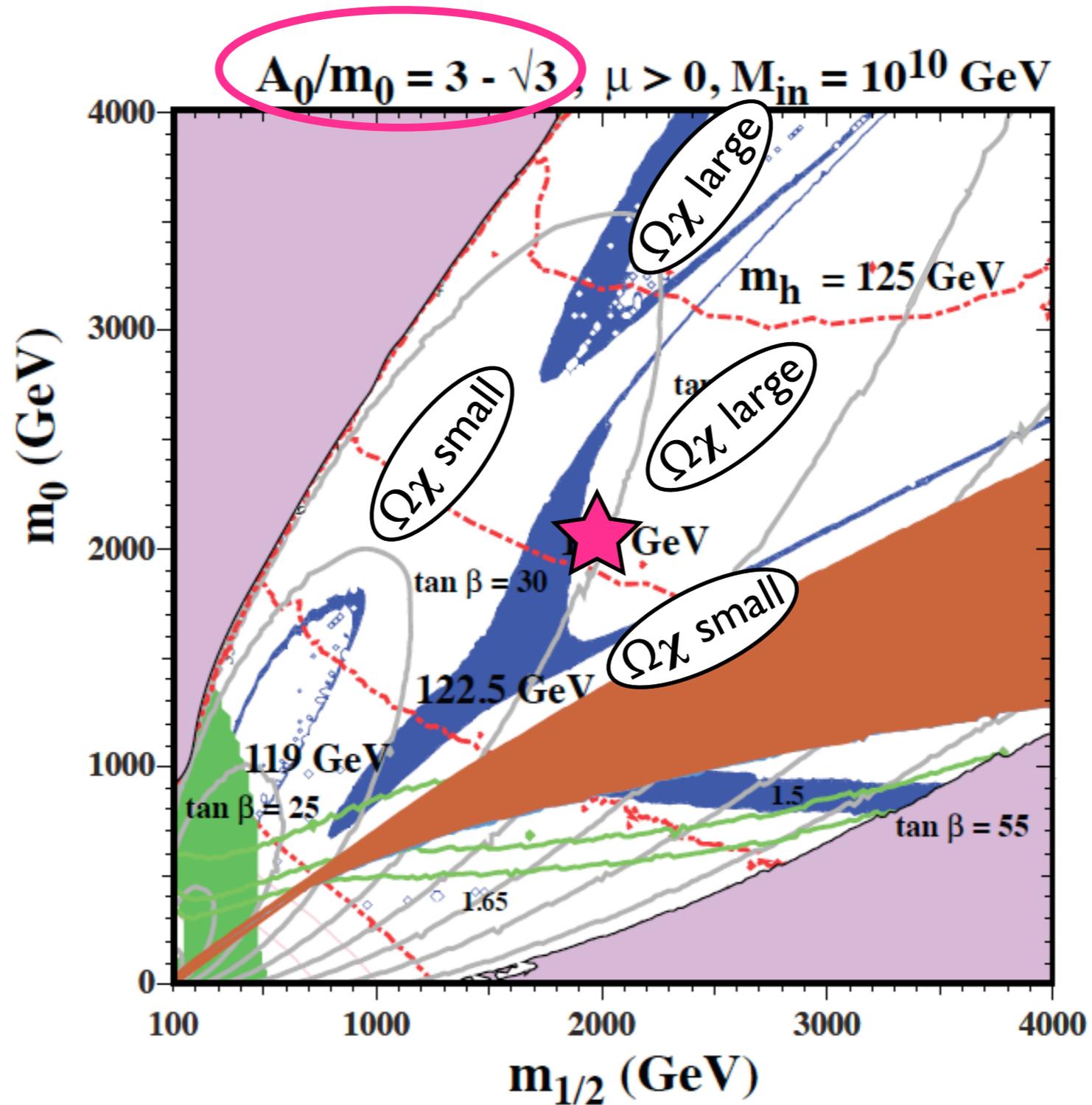


Extra Slides

Looking Forward

- Direct dark matter searches - towards the neutrino background! and directional searches!
- Indirect dark matter searches
 - Fermi, HAWK, VERITAS, AMS-02, GAPS, CTA GAMMA-400...
- LHC - SUSY/DM discovery potential at 14 TeV
- 100 TeV Hadron Collider
- Linear Collider - ILC at 500 GeV, CLIC at 3TeV

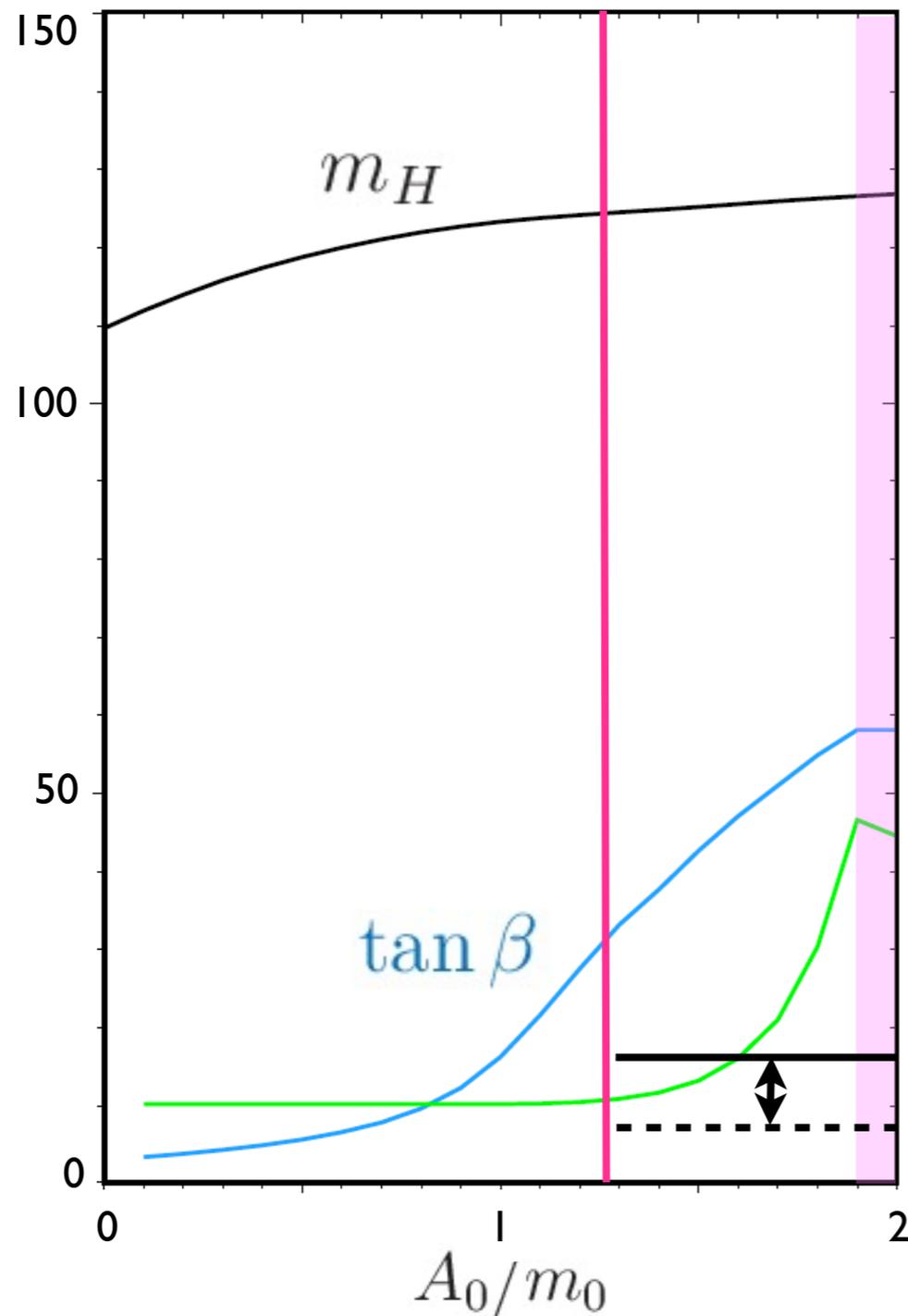
sub-GUT mSUGRA



Ellis, Luo, Olive,
Sandick (2013)

sub-GUT mSUGRA

$m_{1/2} = 2000 \text{ GeV}, m_0 = 2000 \text{ GeV}$



At large $\tan\beta$,

$$\text{BR}(B_s \rightarrow \mu^+ \mu^-) \sim \tan^6 \beta$$

A_0 { large enough: $m_H \approx 126 \text{ GeV}$
 small enough: $\text{BR}(B_s \rightarrow \mu^+ \mu^-) \leq 1.5 \text{ SM value}$

$$\frac{\text{BR}(B_s \rightarrow \mu^+ \mu^-)}{\text{BR}(B_s \rightarrow \mu^+ \mu^-)_{\text{SM}}} \times 10$$

95% CL

**Polonyi Model
 seems to be
 the sweet spot!**