Near-conformal BSM theories on the lattice Kieran Holland University of the Pacific Mass 2012, Nordita, June 13 2012

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summary

- Iattice role for BSM
- what can you do on the lattice?
- Imitations of simulations
- Iatest SU(3) results
- future focus

lattice role for BSM

- strong dynamics for EWSB: fix naturalness, triviality
- near-conformal ("walking"): fix FCNC/mass generation tension?
- new vista: reps beyond SU(3) fundamental, less tension with expt
- interesting models likely non-perturbative: lattice
- can lattice show if any given gauge theory conformal or not?
- would be sad if we (lattice) were not exploring possibility of new strong gauge theories in Nature

Uralness, triviality C/mass generation tension? amental, less tension with expt bative: lattice theory conformal or not? of exploring possibility of new

conformal window



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each TC fermion rep has window in flavor-color space where theory conformal

perturbative: find IR fixed point of beta-fn

amend: strong IRFP could break chiral symmetry

near-conformal just below window

lattice task: where exactly are the windows?

lattice



fundamental Lagrangian e.g. QCD: quarks, gluons fermions on sites, gauge fields on links U can fix gauge-invariance exact lose symmetry: translation, rotation, chirality recovered in continuum limit (zero spacing) non-perturbative no expansion

discretize space-time: regulator, lattice spacing aobservable $\langle B \rangle = \int DU \{ \text{Det}(D[U]) \}^{N_f} \exp(-S_g[U]) B[U]$ Monte Carlo simulation $\int DU \rightarrow \sum$ typical U

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limitations of simulations

1. generation of U configurations: numerical algorithm slow due to non-local fermion determinant

2. discrete lattice action - artifacts, difference from continuum improve lattice action: reduce artifacts (good) reduce speed (bad)

3. always simulate at finite volume, typically at finite mass control: continuum limit at fixed large physical volume, then take chiral limit - always saturate available computers

4. larger Nf or unusual representation - even harder than QCD!

cost

- USQCD: central US allocation of hep-lat computational resources
- BSM efforts: tens of millions of computer core-hours per year
- I core-year approx. 8,000 core hours
- similar resources as large-scale QCD efforts (MILC/HPQCD)
- since 2007, many lattice groups worldwide work on BSM
- In attice community not trying to do BSM on the cheap!



ep-lat computational resources

D efforts (MILC/HPQCD) orldwide work on BSM 3SM on the cheap!

lattice observables

- particle spectrum
- chiral condensate, Dirac operator eigenvalues
- running of renormalized gauge coupling (various schemes)
- detection of finite-temperature phase transitions
- Sparameter
- RG flow in bare coupling space

which theories

my talk: (a) 2-flavor SU(3) 2-index symmetric rep (sextet/NMWT) (b) 12-flavor SU(3) fundamental rep why? (a) sextet attractive for BSM: 3 Goldstones match W&Z, few d.o.f. (b) much work on fundamental rep for various number flavors

next talk: Francis Bursa on SU(2) reps

particle spectrum

expectations:

if chirally broken

 $M_{\pi}^{2} = a_{1}m + a_{2}m^{2}$ $F_{\pi} = F + b_{1}m$ $M_{\text{nuc}} = M_{\text{nuc},0} + c_{1}m$ $\langle \bar{\Psi}\Psi \rangle = d_{0} + d_{1}m + d_{2}m^{2}$

if conformal

 $M_{\text{hadron}} = c_{\text{hadron}} m^{1/(1+\gamma)}$ $\langle \bar{\Psi}\Psi \rangle = c_{\text{cond}} m^{(3-\gamma)/(1+\gamma)} + c_1 m$

notation $y_m = 1 + \gamma$ anomalous dimension

polynomials chiPT motivated but not sensitive to logs + large Nf issues in chiPT

universal critical exponent no mass gap in chiral limit dimension

2-flavor SU(3) sextet

simulation details: staggered fermion discretization - fast, remnant chiral symmetry but flavor symmetry broken partial fix - improved fermion and gauge lattice action

bare parameter choice: set fermion mass and lattice spacing largest lattice size $48^3 \times 96$ lightest fermion mass $1/(m_{\pi}a) \approx 7$ several lattice volumes, fermion masses, 2 lattice spacings each simulation: 1-2 thousand gauge configurations - expensive

Fodor et al. arXiv:1205.1878

results: spectrum



$$M_{1}=M_{1}+2, m_{1}\beta=3.2$$

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again, chiPT-like fits to data are good non-zero mass gaps in chiral limit



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results: chiral condensate



chiral condensate also fit reasonably, with non-zero chiral limit 2nd independent measurement of subtracted condensate two methods have consistent values in chiral limit - very good

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volume-dependence



volume-dependence of spectrum significant $M_{\pi}L > 5$ to be within 1% of infinite-volume value

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 $\gamma = 2.13 \pm 0.18$

$\gamma = 1.091 \pm 0.0034$

fit data to power-like conformal behavior $M_{\rm had} = c_{\rm had} m^{1/1 + \gamma}$ anomalous dimension γ individual fits reasonable; combined fit unacceptable

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fit states together: bad $\gamma = 1.47 \pm 0.26$

large

other 2-flavor sextet work

Kogut & Sinclair, arXiv:1111.2319 finite-tem



critical temperature

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- finite-temperature chiral phase transition vary physical size of (periodic) time direction via lattice spacing (bare coupling) chiral susceptibility develops peak at critical temperature
- peak more pronounced at lighter mass consistent with chiSB at zero temperature want to repeat closer to continuum limit

other 2-flavor sextet work



0.2

0.3

u = 1/g

0.4

0.5



vary lattice action IR fixed point ex also find small v

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-0.05

0

01

vary lattice action and operator

IR fixed point existence varies with method

also find small value $\gamma < 0.45$

12-flavor SU(3) fund

simulation details: staggered fermion discretization - fast, remnant chiral symmetry but flavor symmetry broken partial fix - improved fermion and gauge lattice action

bare parameter choice: set fermion mass and lattice spacing largest lattice size $48^3 \times 96$ lightest fermion mass $1/(m_{\pi}a) \approx 6$ several lattice volumes, fermion masses, 1 lattice spacing each simulation: 1-2 thousand gauge configurations - expensive

Fodor et al. arXiv:1205.1878,1104.3124

spectrum



chiPT-like fits of data reasonable

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i5Pion fit plotted with pion, ijPion, and scPion
0.2

$$M_{i5Pion}^2 = c_1 m + c_2 m^2$$

0.15
 $C_2 = 69.9 \pm 6.7$
 $\chi^2/dof = 3.57$



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alternative: conformal?

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$M_{\pi}(L) = f(x), \quad x = Lm^{1/y_m}$

fferent for each state

ge x i.e. $M_{\pi} = c_1 m^{1/y_m}$

as $x \to 0$

nass

easonable do not agree

ho: 0.300(17)

alternative: conformal?

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add more states: situation worse

data appears not to allow universal value of exponent

conformal description of data appears worse than chiSB form

need accurate data to test this

lattice activity SU(3) 12-flavor

- Fodor et al.
- Appelquist et al., Lattice Strong Dynamics
- Pallante, Lombardo, Deuzeman
- Lin et al.
- Yamawaki et al.
- Jin & Mawhinney
- Hasenfratz et al.

running coupling

renormalized coupling as function of lattice size L $g^2(L)$ run coupling by varying lattice size e.g. L to 2L repeat at various lattice spacings, find continuum limit $a/L \rightarrow 0$ coupling flows to IR fixed point - conformal

Appelquist, Fleming, Neil, arXiv:0901.3766

static quark potential

measure quark potential on same simulations for particle spectrum potential has confining linear behavior at intermediate separation $M_{
m nucleon}/\sqrt{\sigma}$ appears non-zero in chiral limit, as does nucleon mass consistent with chiSB in spectrum, inconsistent with IRFP

Fodor et al. arXiv:1104.3124

the story so far

- In attice methods and machines can be applied to new models
- very important to have many signals, discretizations, people
- 2-flavor sextet theory looks non-conformal, but remaining issues
- 12-flavor fundamental is difficult many puzzles
- many more studies done on fundamental rep learning phase
 given high cost, may wish to focus resources on sextet model

extra slides

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running coupling

 $q^2(L)$ Polyakov loop correlator scheme again, run coupling by varying lattice size e.g. L to 2L lattice artifacts require continuum limit $a/L \rightarrow 0$ indication of IR fixed point, systematics issues **Kieran Holland**

Lin et al., arXiv:1205.6076

static quark potential

quark potential runs faster than perturbation theory inconsistent with existence of IR fixed point

 $V(R) - V(R_0) = C_F \int_{R_0}^{R} \frac{\alpha_{qq}(R')}{R'^2} dR'$

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Fodor et al. arXiv:1104.3124

predict RHS with n-loop pert thy

near-conformal TC

near-zero of beta function gauge coupling walks, not runs separate scales - useful phenomenology