

The Standard Model, the Higgs and beyond

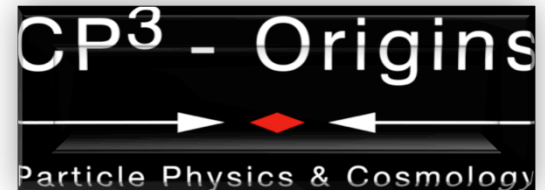
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My life as a boson: by Professor Peter Higgs

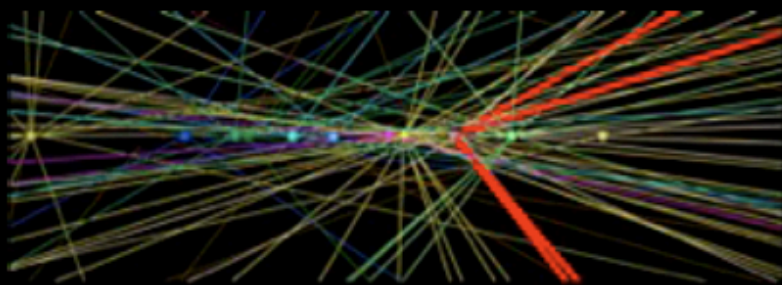
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livestreaming/higgs-boson/](http://www.swan.ac.uk/media-centre/livestreaming/higgs-boson/)



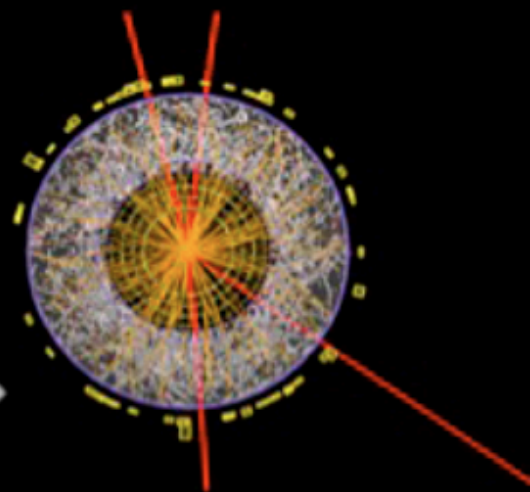
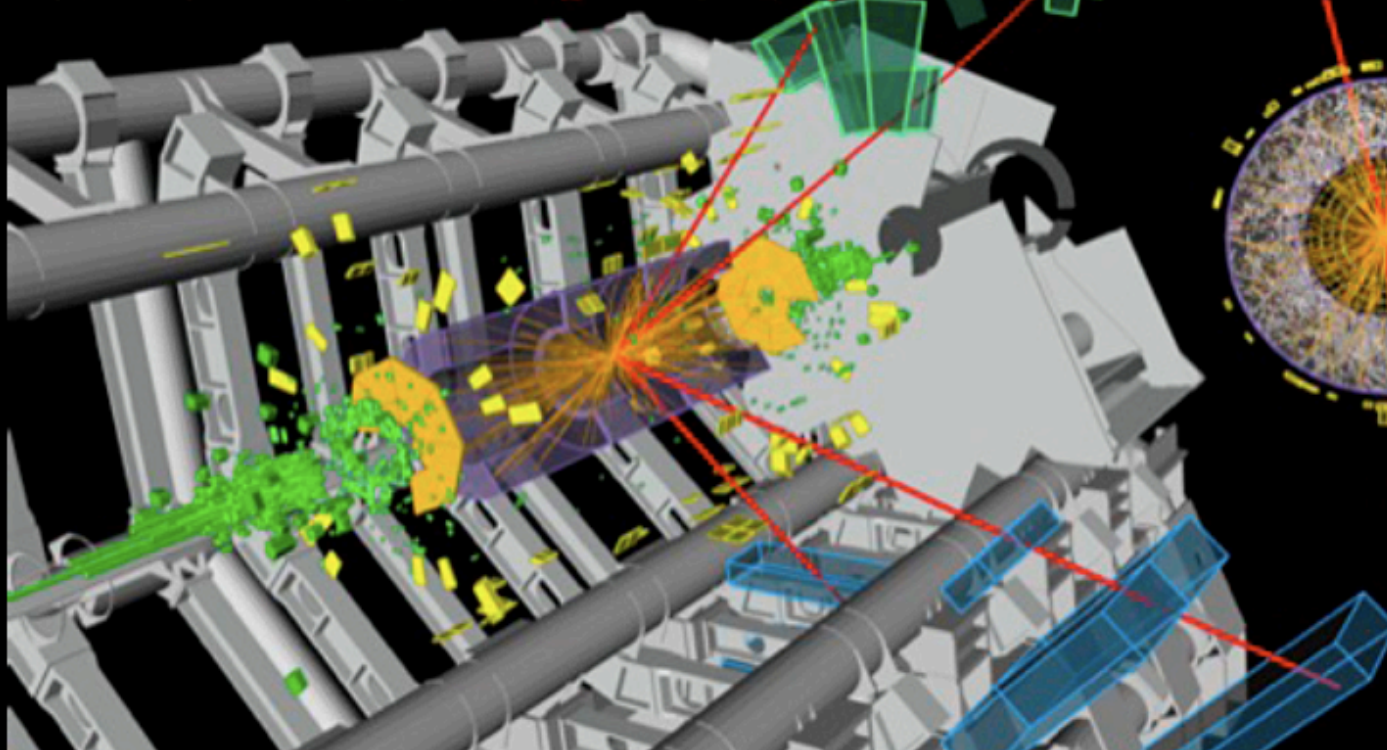
Based in part on lectures by Professor Graham Ross

<http://indico.cern.ch/scripts/SSLPdisplay.py?stdat=2008-06-30&nbweeks=6>

Visit www.cp3-origins.dk for various lectures and info

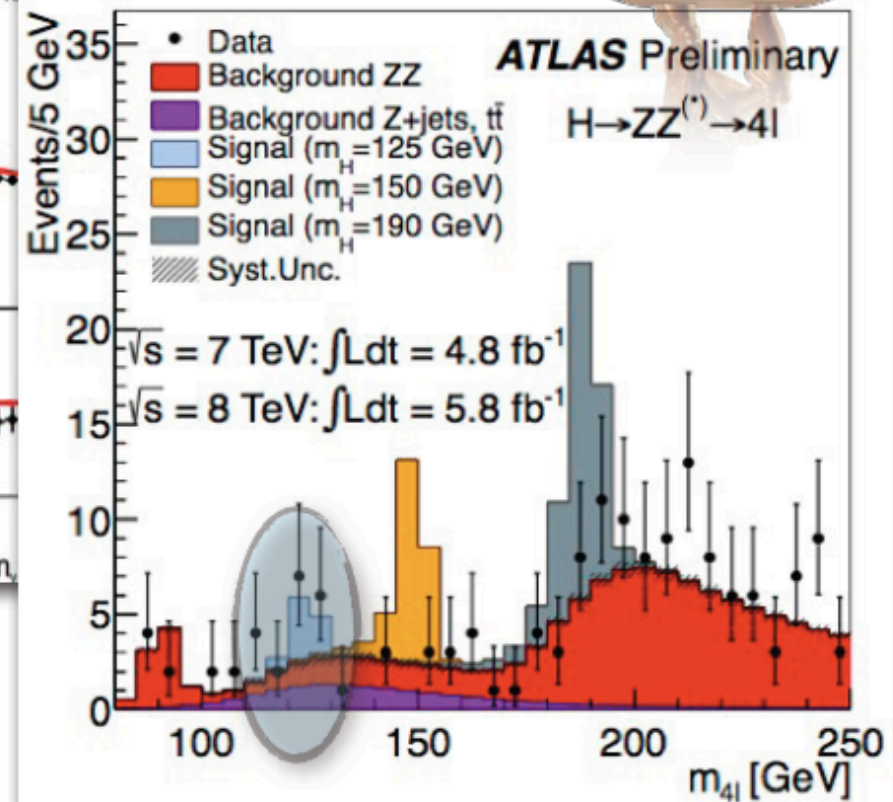
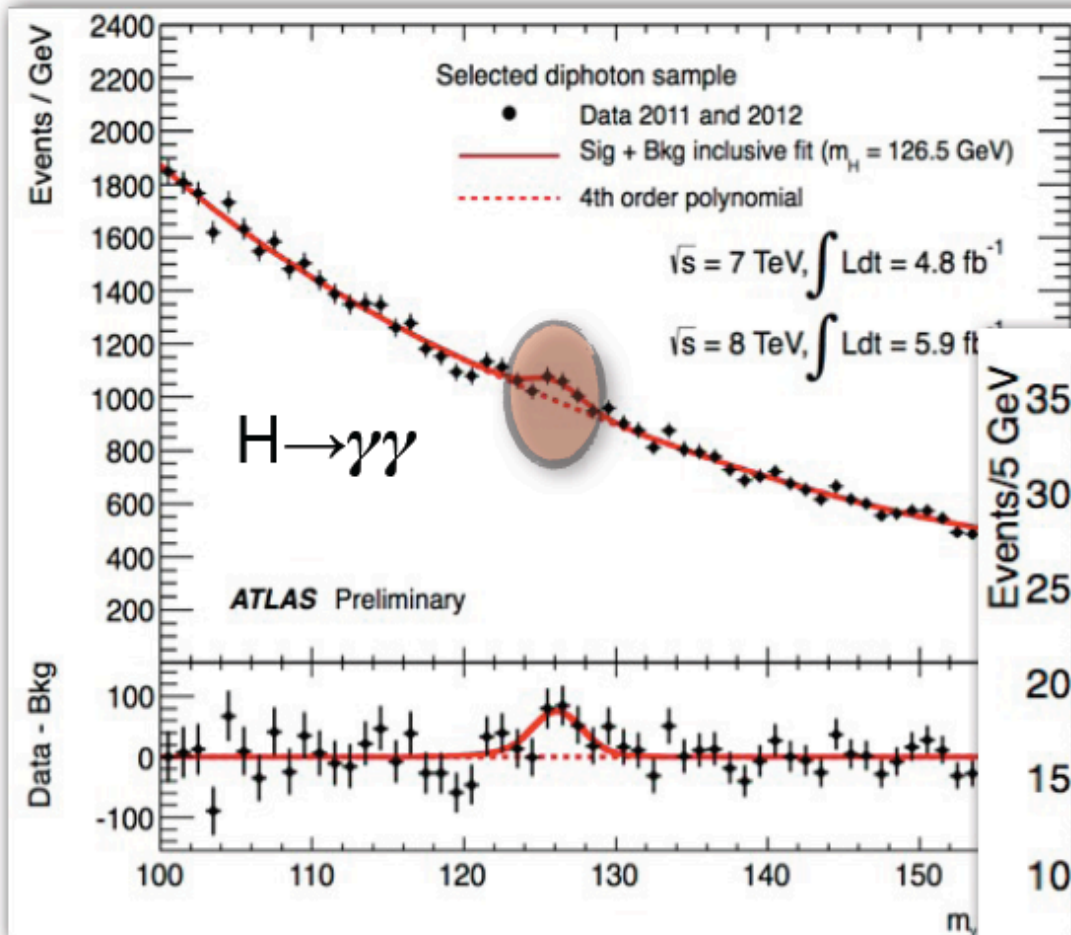


ATLAS
EXPERIMENT
<http://atlas.ch>

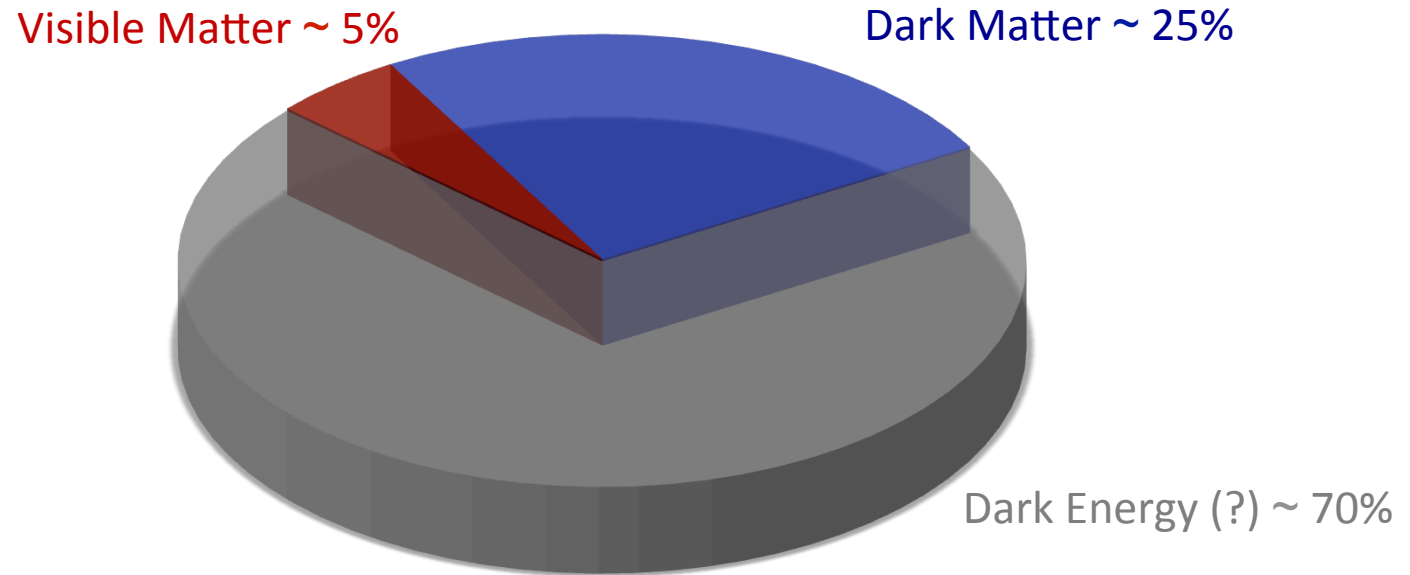


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2 bumps



What is the world made of?



See lectures by Elgarøy for discussions of dark matter and dark energy

What is the world made of?

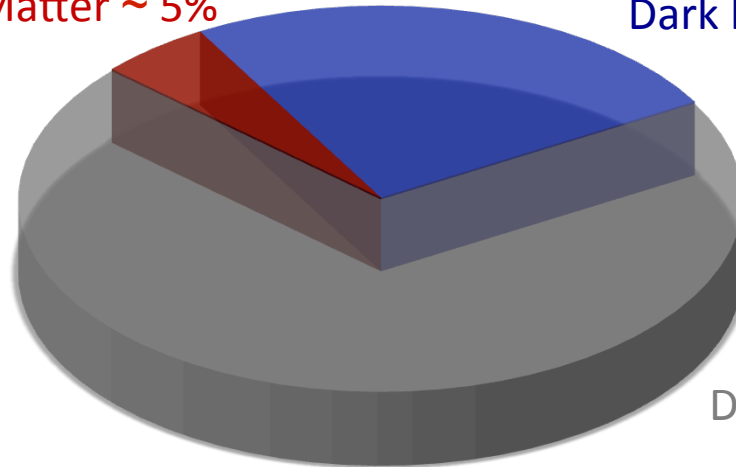
Two outstanding problems of the Standard Models of particle physics and cosmology:

What is the *origin* of mass?

What is the *nature* of Dark Matter?

Visible Matter ~ 5%

Dark Matter ~ 25%



Dark Energy ~ 70%

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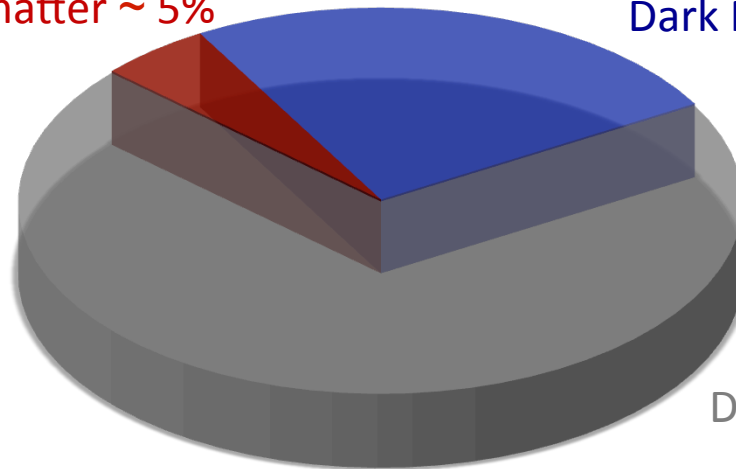
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No obvious particle physics scale associated with the problem of dark energy

Many proposed particle physics scales associated with the nature and origin of DM

A definite particle physics scale related to (part of!) the origin of mass, i.e. Higgs-mechanism
...though not necessarily the Higgs boson!

What is the world made of?

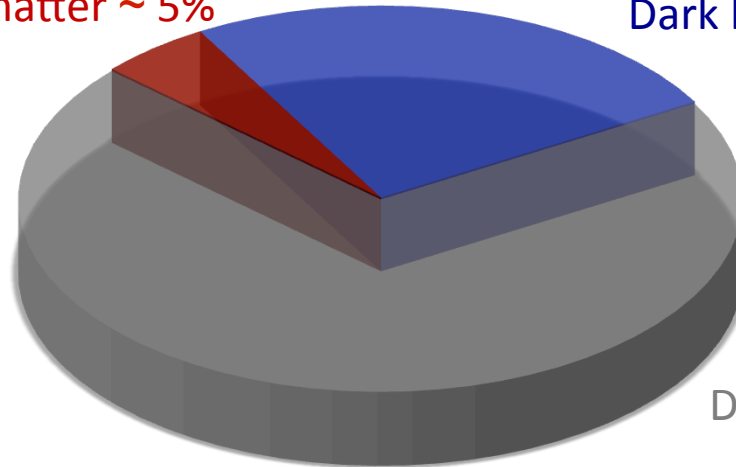
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In these lectures we will discuss the SM describing the visible matter in the universe, and the origin of mass for the SM particles – the Higgs mechanism

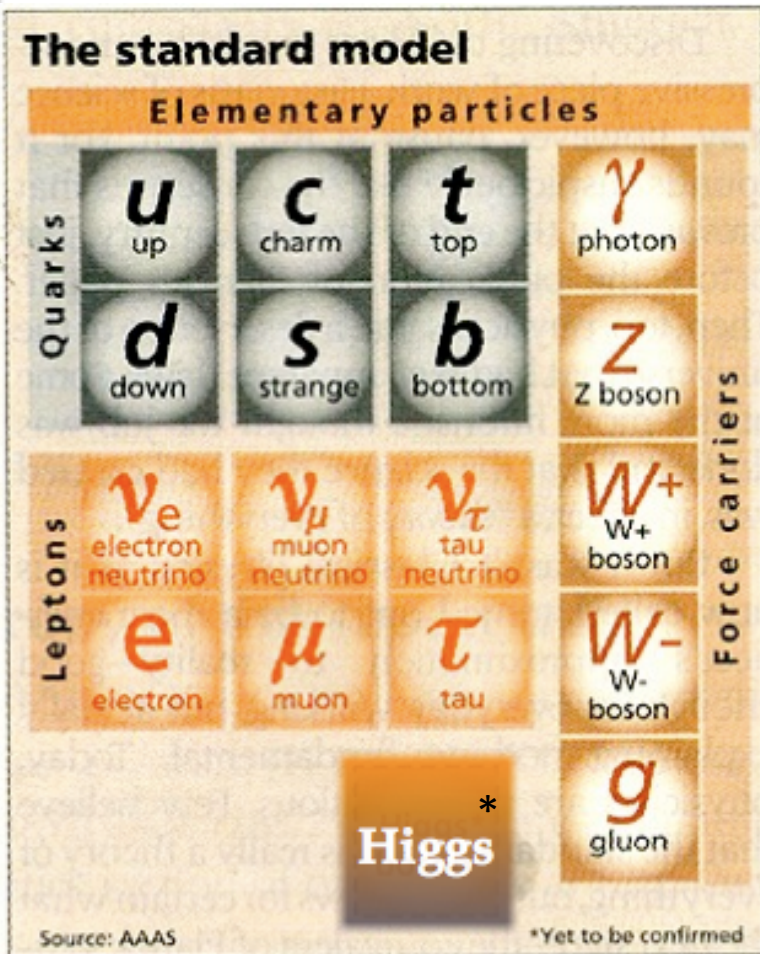
The SM, the Higgs and beyond

Tentative Outline

- Lecture 1 – Fundamental particles and interactions, Higgs teaser
- Lecture 2 - Symmetries, Quantum Field Theory, Gauge Theory
- Lecture 3 - The SM, the Higgs Mechanism of the Standard Model
- Lecture 4 – Beyond the SM, curing the ills of the SM Higgs boson?

The accompanying notes and exercises, which you should work through will provide the details omitted in the lecture slides

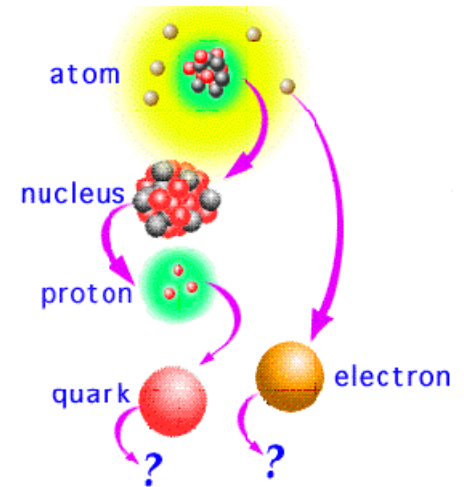
The Standard Model on a stamp



U(1)

SU(2)

SU(3)



4 known Fundamental Interactions

	Strength
Strong	$\alpha_s = \frac{g_s^2}{4\pi\hbar c} \sim 1^\dagger$
Electromagnetic	$\alpha_{em} = \frac{e^2}{4\pi\hbar c} \sim \frac{1}{137}$
Weak	$G_F m_p^2 \sim 10^{-5}^\dagger$
Gravitational	$G_N m_p^2 \sim 10^{-36}$

† *Short range* (Do you know why?)

The Standard Model

The standard model

Elementary particles			
Quarks	u up	c charm	t top
	d down	s strange	b bottom
Leptons	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino
	e electron	μ muon	τ tau
* Higgs			g gluon

Force carriers: **γ** photon, **Z** Z boson, **W^+** W^+ boson, **W^-** W^- boson

Source: AAAS *Yet to be confirmed

Each particle defined by a few charges or *quantum numbers*:

U(1) Electric-, weak-, strong charge, mass and in addition spin.

SU(2) For each charge there is a corresponding force acting on the charge – described by a gauge theory

Matter particles are spin-1/2

SU(3) *Force* particles are spin -1
Higgs Boson in the SM is spin-0

Fundamental Interactions

Strength

Strong – SU(3) $\alpha_s = \frac{g_s^2}{4\pi\hbar c} \sim 1^{1\dagger}$

EM – U(1) $\alpha_{em} = \frac{e^2}{4\pi\hbar c} \sim \frac{1}{137}$

Weak – SU(2) $G_F m_p^2 \sim 10^{*5}$

Gravitational $G_N m_p^2 \sim 10^{-36}$

THE PERIODIC TABLE

Leptons		Quarks (each in 3 “colors”)		
Particles like the electron (fermions, spin 1/2)	e 0.511 MeV -1	ν_e < 0.000003 0	d 7 -1/3	u 3 2/3
	μ 106 -1	ν_μ < 0.2 0	s 120 -1/3	c 1200 2/3
	τ 1777 -1	ν_τ < 20 0	b 4300 -1/3	t 175,000 2/3
				← charge

Particles like the photon (bosons, spin 1)	γ photon 0	“electromagnetism”
	g gluon 0 (8 “colors”)	“strong interaction”
	W^\pm Z^0 80,420 91,188	“weak interaction”

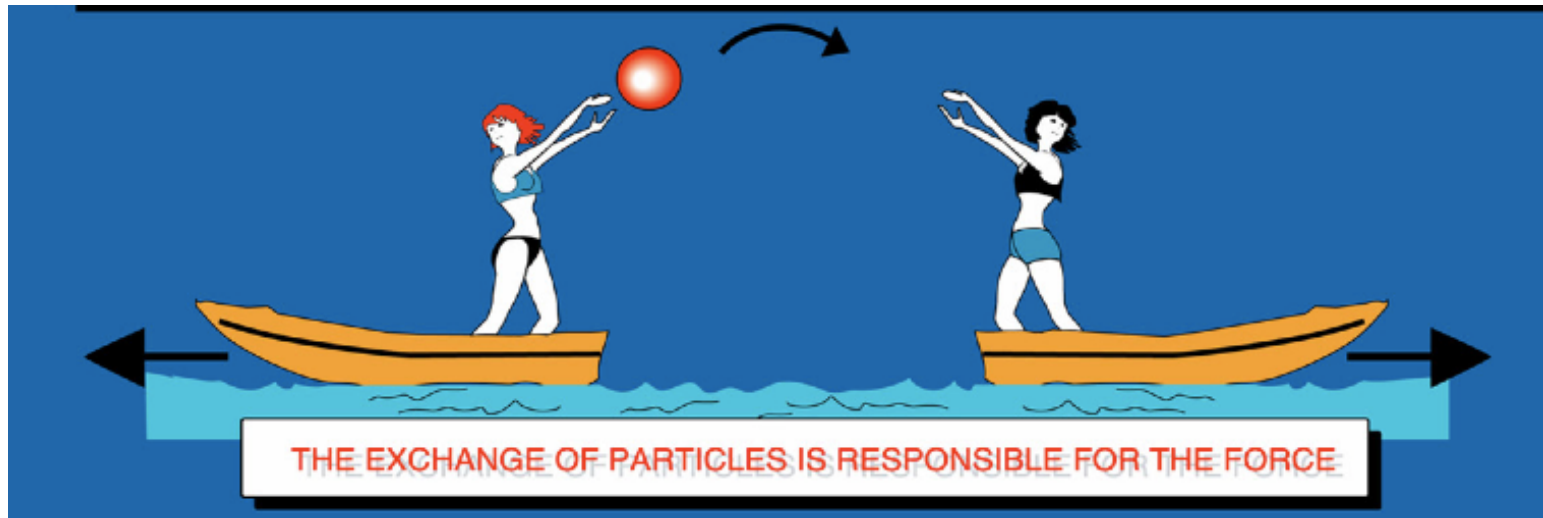
Forces

Can we understand forces as particle exchange clasically?

Forces

Can we understand forces as particle exchange classically?

Repulsive force as particle exchange:



What about an attractive force?

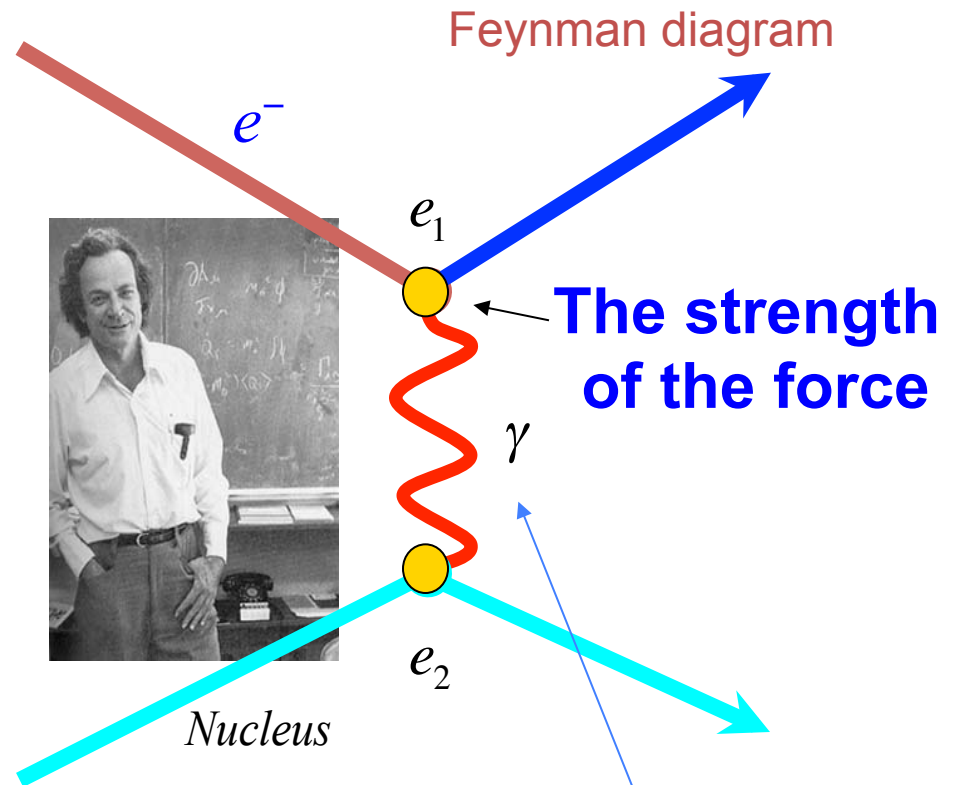
Exchange forces

Electromagnetism

$$V_{em}(r) = \frac{e_1 e_2}{4\pi} \frac{1}{r}$$

In momentum space :

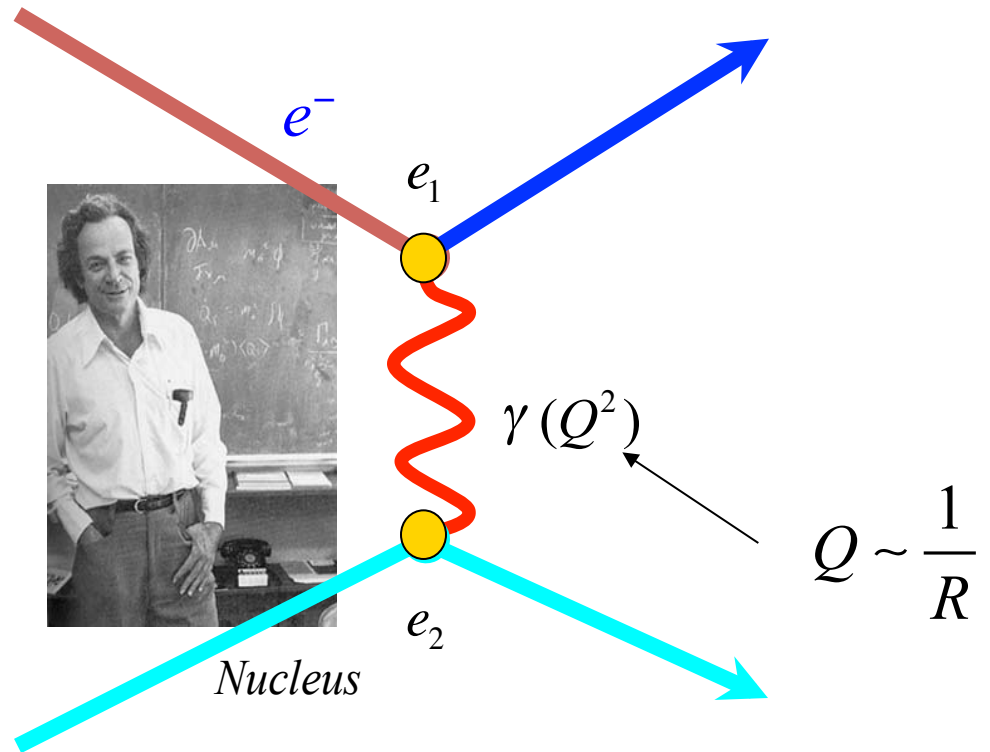
$$V_{em}(|\mathbf{q}|) \sim \int V_{em}(|\mathbf{r}|) e^{-i\mathbf{k} \cdot \mathbf{r}} d^3\mathbf{r} \sim \frac{\alpha}{|\mathbf{k}|^2}$$



Exchange forces

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Experiments conducted in momentum space :

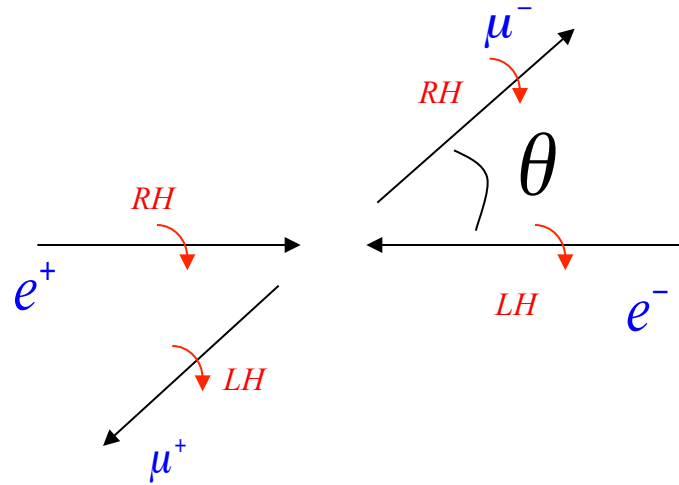
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$$Q^2 \equiv -\mathbf{k}^2$$

"virtual photon"

Application to a scattering processes

$$e^+ e^- \rightarrow \mu^+ \mu^-$$

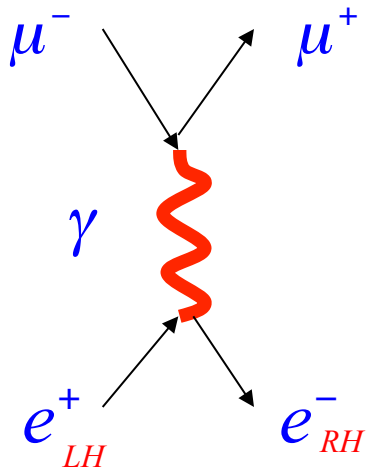


$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 E_{CM}^2} |M|^2$$

Feynman diagram

QM : Transition amplitude

$$\langle \text{final state} | H_I | \text{initial state} \rangle$$



$$M \propto \langle \mu^+ \mu^- | H_I | \gamma \rangle^\alpha \langle \gamma | H_I | e^+ e^- \rangle_\alpha$$

$$\langle \gamma | H_I | e^+ e^- \rangle^\alpha \propto e(0, 1, i, 0)$$

$$\langle \mu^+ \mu^- | H_I | \gamma \rangle^\alpha \propto e(0, \cos \theta, i, \sin \theta)$$

$$M(RL \rightarrow RL) = e^2 (1 + \cos \theta)$$

The Standard Model

The standard model

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Higgs Boson in the SM is spin - 0

Fundamental Interactions

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<i>Strong</i> – SU(3)	$\alpha_s = \frac{g_s^2}{4\pi\hbar c} \sim 1^{1\dagger}$
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Higgs Boson looks special but more generally
Mass is a special charge!...
can you give a reason why?

Mass and the elementary particles

i) There is no anti-charge for mass!

But here we completely neglect gravity which is unimportant for elementary particle processes and we don't have a quantum theory of gravity...there's plenty to do!

Mass and the elementary particles

i) There is no anti-charge for mass!

But here we completely neglect gravity which is unimportant for elementary particle processes and we don't have a quantum theory of gravity...there's plenty to do!

ii) There is another peculiar feature of mass in a quantum theory :

Scattering of massive spin one-states:

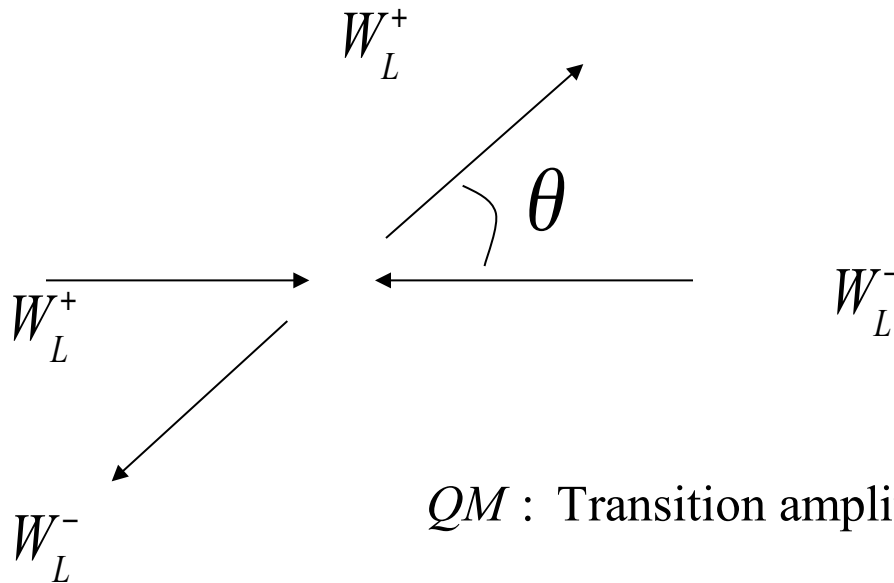
Do you remember any spin-1 states (Lorentz vector) from classical or quantum physics?

How many *degrees of freedom*?

How many *physical* degrees of freedom does a massless spin-1 field have?

How about a massive spin-1 vector. Can you explain why?

Scattering of massive W-bosons

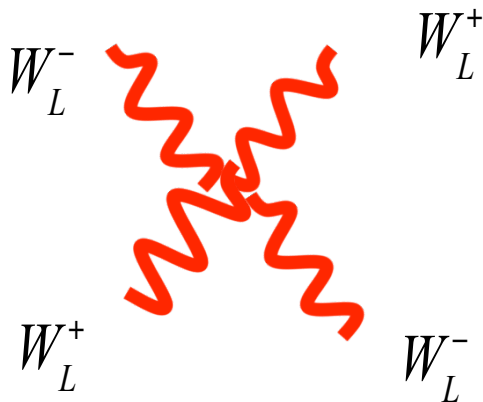


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Feynman diagram

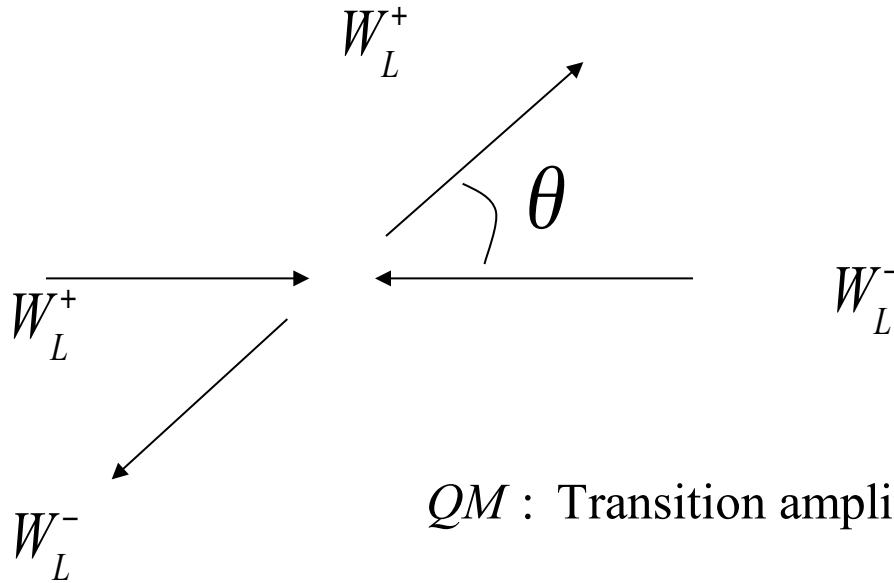


$$M \propto \langle W^+ W^- | H_I | W^+ W^- \rangle$$

$$M(W_L^+ W_L^- \rightarrow W_L^+ W_L^-) \sim E^2 / m_W^2 !$$

Where E is the energy...what does this imply for High energy scatterings such as at LHC?

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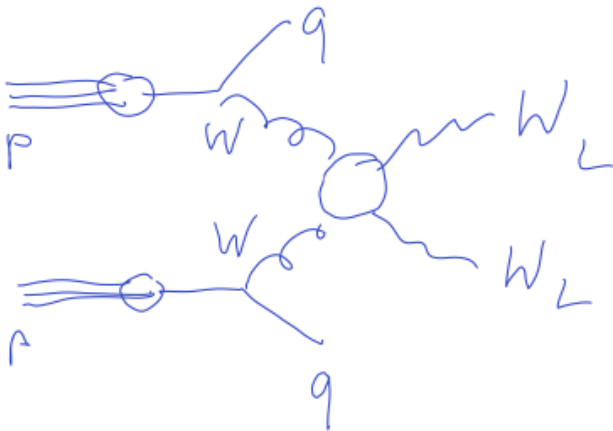


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Either QM breaks down or we discover something new at LHC...seems that we already are starting to see it

The Higgs mechanism

Gauge theories do not suffer from the *unitarity* problem, spin-1 gauge fields are massless

In fact we take *gauge symmetry* as a fundamental symmetry of the SM.
However, the *W/Z Bosons are still observed to be massive*

The way out is to invoke the Higgs-mechanism and mass as an emergent phenomenon via *Spontaneous symmetry breaking*

We will (hopefully) see towards the end how the Higgs boson can unitarize WW scattering

The Higgs mechanism - teaser

Consider a theory with a real scalar field $\phi(x,t)$ (you can think of it just as a classical particle
Next lecture will introduce Quantum Field Theory)

$$V(\phi) = \mu^2 \phi^2 + \frac{1}{2} \lambda^2 \phi^4$$

What *symmetry*, i.e. a non-trivial transformation of $\phi(x,t)$ leaves the potential invariant?:

$$P : \phi \rightarrow \phi' ; \quad V(\phi) = V(\phi')$$

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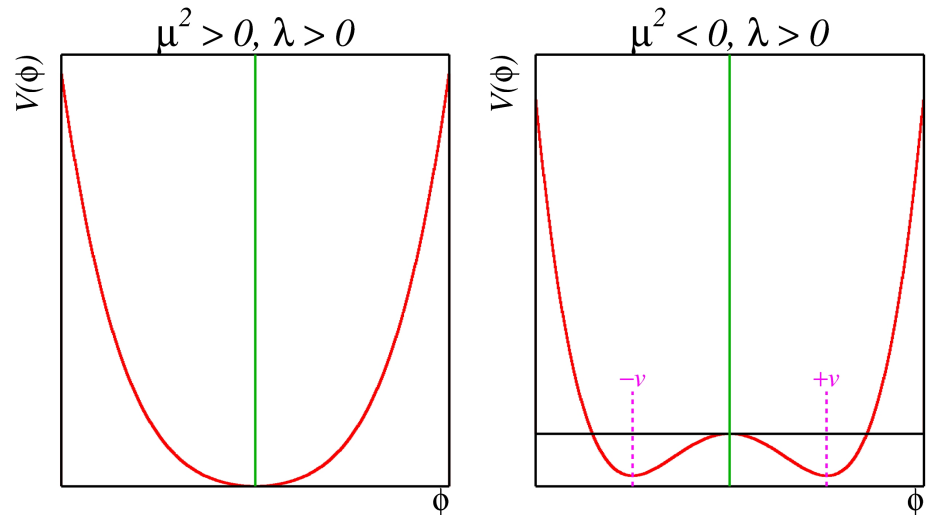
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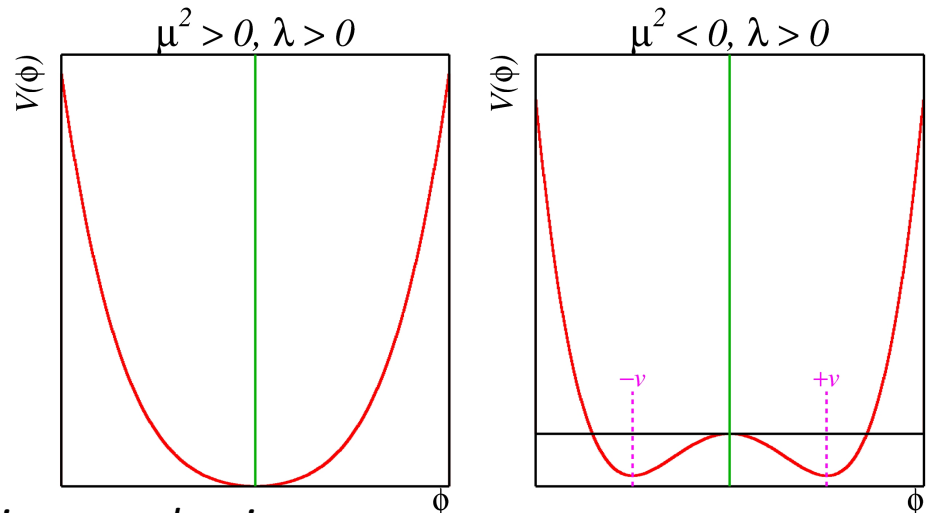
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This is our first example of

Spontaneous symmetry breaking or the Higgs mechanism



The Higgs mechanism - teaser

Now consider the same theory coupled to another field ψ via

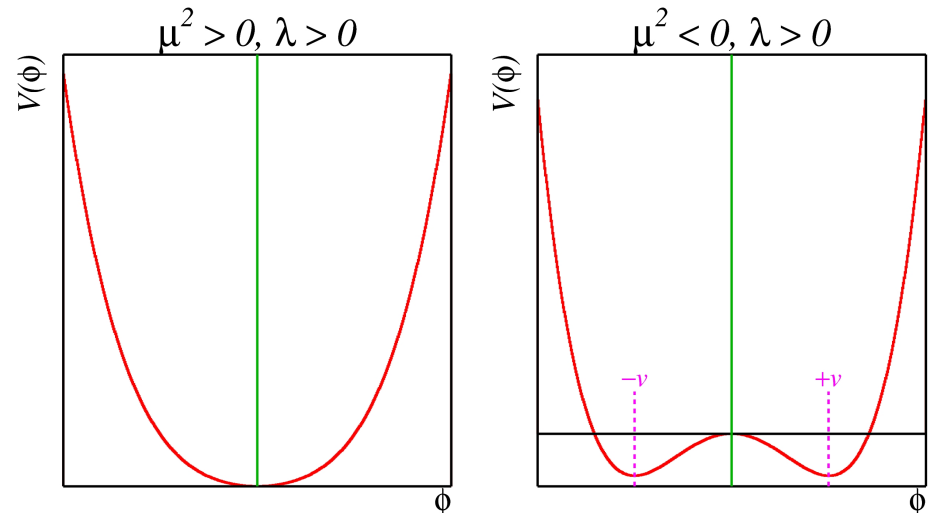
$$V(\phi) = \mu^2 \phi^2 + \frac{1}{2} \lambda^2 \phi^4 + \frac{1}{2} \psi^2 \phi$$

To do perturbation theory in the broken phase we expand $\phi = v + \delta\phi$

$$V(\phi) \rightarrow \frac{1}{2} \psi^2 v + \dots$$

The field ψ has acquired a mass!

This is our first example of
Spontaneous symmetry breaking
and the Higgs mechanism



v is a 'preferred direction' in field space – compare alignment of magnets, Hansson's lecture

Goldstone Theorem

It turns out that there is one essential feature need for the SM Higgs mechanism that our toy model does not capture – *Goldstone bosons*

Our example was of spontaneous breaking of a *discrete symmetry* we need to consider Spontaneous breaking of a *continuous symmetry*

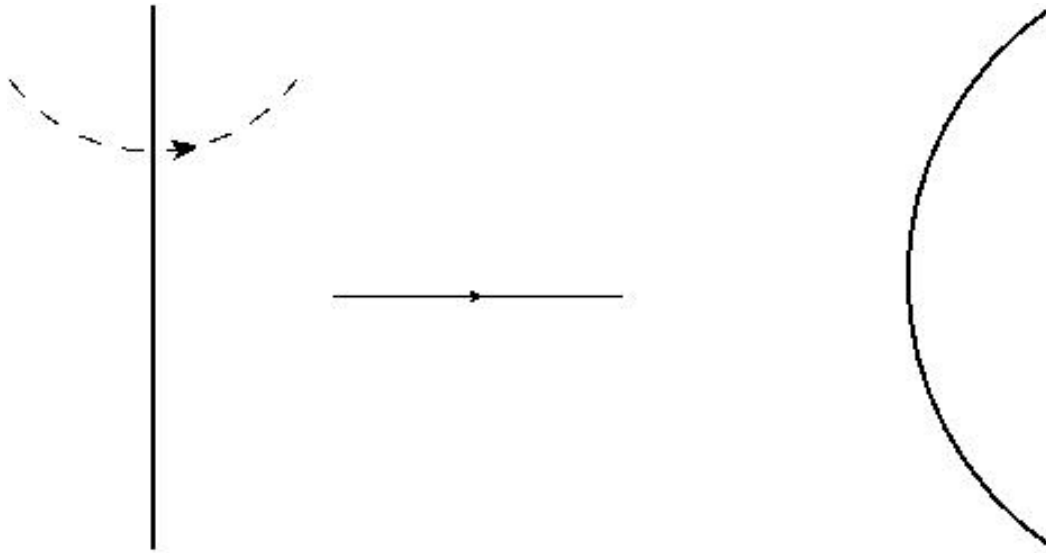
Goldstones Theorem:

For every *spontaneously broken global continuous* symmetry
there will be an associated *massless* particle

Goldstone Theorem - examples

This is simpler than it might sound!

What are the symmetries before and after we bend a (nearly) rigid rod attached at both ends?

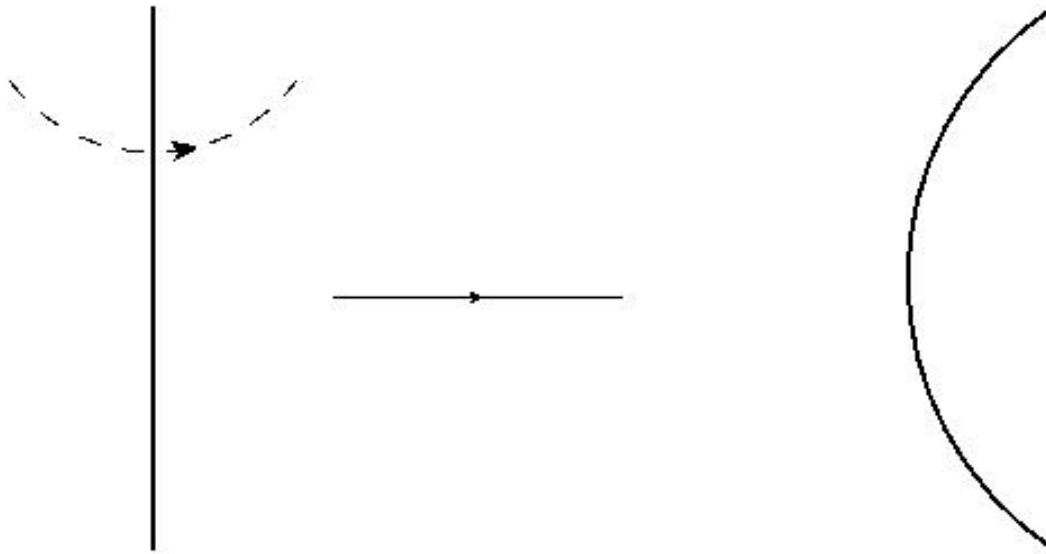


Can you identify a 'massless excitation' existing after but not before symmetry breaking?

Goldstone Theorem

This is simpler than it might sound!

What are the symmetries before and after we bend a (nearly) rigid rod attached at both ends?



Can you identify a ‘massless excitation’ existing after but not before symmetry breaking?

What about water freezing to ice, can you identify a ‘massless excitation’ there?

Generally related to phase transitions as in Hansson’s lectures –

What drives the Higgs mechanism? -> Beyond the Standard Model!

Goldstone Theorem

Goldstones Theorem:

For every *spontaneously broken global continuous* symmetry
there will be an associated *massless* particle

The Higgs Mechanism:

If the spontaneously broken global symmetries are *gauged* the corresponding gauge fields
will acquire mass.

In fact the longitudinal mode of those gauge fields 'are' the Goldstone bosons which are
eaten

All of these statements and toy examples are what we now aim to
concretize in the SM Quantum Field Theory and beyond!