## יוַעֵּילָם נְשָׂא אַשְׁפָה..." ישעיה כב

## On behalf of the ATLAS collaboration Hunting the Higgs Weizmann





### "And Eilam bare the quiver..."

Jesaia 22

# ועילם נשא אשפה... Higgs Hunters Independence Distribution of the second seco עיה **Eilam Gross**

### "And Eilam bare the quiver .... "

Jesaia 22

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### 17 september 2012

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Air Mail

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ISRAEL

FROM Peter Higgs 2 Darnaway Street Eduilourge EH3 6BG

Higgs (in a snail mail to me) thistory of SSB Order of contributions:-(i)1. Manbre (1960) Mambre & Jona-Lasinio (1961) 2. Goldstone (1961) 3. ljøldstore, Salam& Weinberg (1962) #: Anderson (1963) 5. Englest Brout ( aug. 1964) 6. Higge (Sep. & Oct 1964) 7. Geralnik, Hagen & Kibble (Mov. 1964) See the enclosed reprint for my account of phars 1 to be. Guralnik, Hagen & Tribble (7) showed how the Goldstone theorem is evaded in a Simple linear model. Mote that all six of usubere awarded the 2010 Sekerai Prize by the APS.

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Higgs (in a snail mail to me): Thistory of SSIS

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Order of contributions:-Landau (1960)960) Mambe & Jona Lasin co (1961) Goldstone(1961) Goldstone, Salam , Weiberg (1962) (1962) Anderson (1963) Englert & Brout (1964) ( aug (964) Higgs (1964) Star & Od 1966) Geralnik, Hagen & Kibble (1964) (her 1964)

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the ATD.

## A Prelude to the Nobel Prize

2010 Sakurai Prize awarded for 1964 Higgs Boson theory work to 0 Hagen, Guralnik, Kibble, Brout, Englert & Higgs



## Spontaneous Symmetry Breaking

- Spontaneously Symmetry Breaking was first introduced by Ginzburg & Landau (1950,1957) (in an attempt to explain superconductivity)
- The physics of the system (Lagrangian) posses some exact symmetry, but the vacuum (ground state) breaks this symmetry





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Nambu (1960) proposed for the first time that SSB is the source of fermion masses in elementary particle physics: "the existence of such a condensate (scalar field) would

break the symmetry of the model.... In particle physics, that would be a non-Abelian group containing the u(1) group associated with electric charge conservation as a subgroup"

## Spontaneous Symmetry Breaking



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Inspired by Nambu, Goldstone (1961) studies models featuring scalar fields and finds that all these models contains (under SSB) massless (Nambu-Goldstone) Bosons

Goldstone, Salam and Weinberg (1962) prove formally that Goldstone Bosons must occur whenever a symmetry ("like isospin or strangeness") is broken (Goldstone Theorem). But no such Bosons were observed experimentally.

Weinberg recalls in his Nobel lecture (1979) that he was so disappointed that he added a quote to the paper from king Lear: "Nothing will come out of nothing, speak again"

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Is Quantum Field Theory a one trick pony? Can it explain only long range interactions?



## Spontaneous Symmetry Breaking

Philip Anderson (1963) points out that in a superconductor the Goldstone mode becomes a massive plasmon-mode, due to its electromagnetic interaction.



Peter Higgs (Phys. Lett. July 1964) shows that one can evade Goldstone theorem. He shows that if the broken symmetry is local gauge symmetry (like electromagnetic U(1) gauge invariance), then, although the

Goldstone Bosons exist formally, and in some sense real, they can be eliminated by gauge transformation, so that they do not appear as physical particles. That explains why experiment fails to detect the massless Bosons.

The missing Gloldstone boson appears instead as helicity zero state of the massless boson which thereby acquire a mass.

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The massless boson eats the Goldstone Boson and acquires mass.



## The Higgs Mechanism

Based on field theory (using a lagrangian formalism) Higgs develops the formalism of the mechanism by which the Goldstone Boson is "eaten" by the photon and the pohoton becomes massive -> short range interaction

He sends the 3 pages paper to Physics Letter, the paper is rejected. Higgs: "I was rather shocked. I did not see why they would accept a paper that said this is a possible way to evade the Goldstone theorem, and then reject a paper that showed how you actually do it."

Higgs adds an epilogue to the paper: "it is worth noting that an essential feature of this type of theory is the prediction of incomplete multiplets of scalar and vector bosons" and sends the revised version to PRL.

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## The Higgs Mechanism

Higgs: "The referee who, I discovered later, was Nambu, drew my attention to a paper by Englert and Brout that they had just published in Physical Review Letters". Higgs is asked to cite Englert & Brout and the paper is accepted (August 1964)

Guralnik, Hagen and Kibble (1964). Guralnik (2009): "As we were literally placing the manuscript in the envelope to be sent to PRL, Kibble came into the office bearing two papers by Higgs and the one by Englert and Brout. These had just arrived in the then very slow and unreliable... Imperial College mail. We were very surprised and even amazed."

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## The Higgs Mechanism Higgs (in a snail mail to me):

To evade my first paper on theired how the goldstone theorem. Englest & Brout showed how a gauge field interaction terns yoldstone massless spin. O bosons (elementary on composite) into helicity 0 states of makine spin-1 particles. They at Started from Slepunan diagrams and didn't discuss the remaining massive Spine O particles. Snung second paper I used Lagrangian field theory explicitly with elementary scales field's (à la goldstone) compled to a gange massive spin- 0 boson gield, so the wes an obvious peakere, to which I drew attention all three of us tried without success

## The Higgs Mechanism

Higgs (in a snail mail to me): In my first paper I outlined how to evade the Goldstone theorem. Calastane theorem Englert & Brout showed how a gauge field interaction turns Goldstone massless bosons (elementary OR composite) into helicity-0 states of massive spin-1 particles. They started from Feynmann diagrams and didn't discuss the remaining massive spin-0 particles. In my second paper I used Lagarangian field theory explicitly with elementary scalar fields (a' la Goldstone) coupled to a gauge field, so the massive spin-0 boson was an obvious feature, to which I drew attention.

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I drew attention

wer an obvious seature, to which

all three of us tried without success

The Birth of the Standard Model Glashow (1961) suggests that the symmetry of the Electro-Weak interaction is SU(2)XU(1) and is broken to U(1) em. But Glashow puts the masses of the force carriers by hand and his theory is therefore nonrenormalizable



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Weinberg (1967) implements Higgs mechanism to Glashow's  $SU(2) \times U(1)$  and writes the most

quoted paper in the history of particle phsyics

(one of the most quoted . .... >8000 citations).

Weinberg predicts that the mass of the weak interaction force carriers is mW=80 GeV and mZ=90 GeV, but it took another 14 years to confirm it experimentally.

## The Birth of the Standard Model





wrong.

Is this model renormalizable? We usually do not expect non-Abelian gauge theories to be renormalizable if the vector-meson mass is not zero, but our  $Z_{\mu}$  and  $W_{\mu}$  mesons get their mass from the spontaneous breaking of the symmetry, not from a mass term put in at the beginning. Indeed, the model Lagrangian we start from is probably renormalizable, so the question is whether this renormalizability is lost in the reordering of the perturbation theory implied by our redefinition of the fields.

The (theoretical) story was completed when 'tHooft (& Veltman) proved the renormalizability of Yang-Mills theories with masses generated by spontaneous symmetry breaking in a scalar field system in 1971.

All that is left is to find the mass generator, the Higgs Boson

How Elementary Particles Acquire Mass • A mass term is given by  $m\overline{\psi}_L\psi_R$ 

Only left handed fields carry weak charge.

Via SSB the Higgs field "charges" the vacuum with a weak charge and the symmetry is preserved ("hidden")

$$g_{H\psi}H_{L}\overline{\psi}_{L}\psi_{R} - > g_{H\psi}\langle H_{L}\rangle\overline{\psi}_{L}\psi_{R} = g_{H\psi}\nabla\overline{\psi}_{L}\psi_{R}$$

$$m_{\psi} = g_{H\psi} v, \qquad g_{H\psi} = -\frac{\psi}{v}$$



The coupling of the Higgs to particles is proportional to the particles' mass

m

The Higgs Boson will therefore decay with a higher probability to the heaviest particle kinematically available

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## Seeing The Higgs Boson – How To? Proton-Proton Collisions



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## Seeing Particles



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### The BIG challenge in 2012: PILE-UP



### Proton Runs 2010-12

Highest luminosity =  $7.73 \cdot 10^{33}$  cm<sup>-2</sup>s<sup>-1</sup>

Recorded luminosity =  $25.62 \text{ fb}^{-1}$ 



ATLAS: Status of SM Higgs searches, 4/7/2012

## Higgs Production @ the LHC Higgs hardly couples to u & d quarks (which make protons)

To produce a Higgs Boson in P-P collisions 4 processes are used: ggF, VBF, Associate Production and ttH



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## Higgs hardly couples to u & d quarks (which make protons)

To produce a Higgs Boson in P-P collisions 4 processes are used: ggF, VBF, Associate Production and ttH





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## Higgs hardly couples to u & d quarks (which make protons)

To produce a Higgs Boson in P-P collisions 4 processes are used: ggF, VBF, Associate Production and ttH



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## Higgs Decay Modes

- The Higgs Boson decays to the heaviest kinematically available particles pair
- A light Higgs (mH~125 GeV) decays to ττ and mainly to a pair of bottom
   Quarks (bb)
- But H->bb is hard to detect or trigger on (unless produced in association with a W or a Z)
- Leptons (electrons or muons) and photons are easy to trigger on and detect.
- Though BR(H->gamma gamma)~10<sup>-3</sup>,
   H->gamma gamma is the favorite channel for a Higgs with mH~110–130





## Higgs Decay Modes

Once the Z and W channels are open (mH>120) it decays to ZZ\* and WW\*

The Higgs decay modes are classified according to the decays of the daughter bosons, thus the main decay modes are

the golden channel 41=4 leptons

and other WW or ZZ channels







## Higgs Decay Rates





## m<sub>H</sub>=125 GeV





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### m<sub>H</sub>=125 GeV

## For a channel to be usable, we must be able to trigger it

Most efficient and clean triggers are photon or lepton based







## m<sub>H</sub>=125 GeV

## For a channel to be usable, we must be able to trigger it

Most efficient and clean triggers are photon or lepton based



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## m<sub>H</sub>=125 GeV Normalized Usable BRs



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### Electroweak measurements are Higgs backgrounds



- Good agreement with theory , W, Z, tt become a challenge for theory
- Systematics dominate
- Higgs cross section same order of magnitude as Di-Boson production (WW,WZ,ZZ)


### m<sub>H</sub>=125 GeV







### m<sub>H</sub>=125 GeV Channels Weight





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# H->yy Probing LEP 114 GeV





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# H->yy Probing LEP 114 GeV

Clean signature: 2 energetic isolated photons->narrow mass peak  $E^{T}(\gamma 1, \gamma 2) > 40, 30 \text{ GeV}$ 

A narrow peak is searched for over a large, smooth background.

Data are split into 9+1(VBF) categories based on direction of photons (detector region), conversion mode (which affect  $\gamma \gamma$  mass resolution, which is excellent) and  $p^{T}_{\gamma \gamma}$  perpendicular to  $\gamma \gamma$  thrust axis

A fit is performed to the background side band under the BG only hypothesis (an exponential or polynomial carefully chosen to avoid bias in EACH category) (only data is considered)

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H-> $\gamma\gamma$  Experimental Aspects .Needs a powerful  $\gamma$  /jet separation to suppress  $\gamma$  j and jj background with jet ->  $\pi^{0}$  faking single  $\gamma$  $m_{\gamma_{1}\gamma_{2}}^{2} = 2E_{\gamma_{1}}E_{\gamma_{2}}\left(1 - cos \sphericalangle(\gamma_{1}, \gamma_{2})\right)$ 

The fine longitudinal and lateral segmentation and pointing geometry of the ATLAS EM calorimeter enable good yy angular separation and better Z-vertex determination. This is crucial in high pile up environment and in identifying fake photons from pions



Present understanding of calorimeter E response (from tag&probe Z->ee,  $J/\psi$  ->ee, W->e $\nu$  data

and MC)-> Excellent mass resolution Eilam Gross, Higgs Discovery,







# H->yy Results

 $m_{\gamma \gamma}$  was fit (per category)

with exponential or polynomials functions for background plus a sum of Crystal Ball and Gaussian (for tails) for signal.

- Background was fitted from data
- ~170 signal events are expected in 10 fb<sup>-1</sup> for m<sub>H</sub>=126 GeV
- ~6340 expected in m<sub>H</sub>=125
   GeV mass window ->
   S/B~3% in signal mass window
   (20% in VBF)



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### Exclusion: CLs



->CLs measures the compatibility of the data with the signal hypothesis.
->If CLs<5% the signal hypothesis is excluded at the 95% CL</li>

-> $\mu_{up}$  is the signal strength for which CLs=5%

-> If  $\mu_{up} < 1 \Rightarrow \sigma(m_H) < \sigma_{SM}$ 

 $=>m_H$  is excluded at the 95% Confidence Level

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Discovery: po

$$q_0 = -2\log \frac{max_{\{b\}}L(b)}{max_{\{\mu,b\}}L(\mu s(m_H) + b)}$$

# ->po measures the compatibility of the data with the NO-HIGGS hypothesis.

->If p\_0=0.025 the NO-HIGGS hypothesis is rejected at the 2 
$$\sigma$$
 level

$$p_0 = Prob(q_0 > q_0^{obs} \mid H_0)$$

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#### Consistency of data with background-only expectation



Data sample m<sub>H</sub> of max deviation local p-value local significance expected from SM Higgs

2011	126 GeV	3×10-4	3.5 σ	1.6 σ	
2012	127 GeV	3×10 <sup>-4</sup>	3.4 σ	1.9 σ	
2011+2012	126.5 GeV	2x10 <sup>-6</sup>	4.5 σ	2.4 σ	
	Global 2011+2012 (ind	cluding LEE ove	er 110-150 GeV ro	ange): 3.6 σ	





2jet/VBF category brings ~ 3% gain in expected sensitivity; observed gains in data are 10-15% (both years):

- $\rightarrow$  is this a fluctuation or something more fundamental?
- $\rightarrow$  further studies needed (next step !)

Caveat: 2jet category affected by largest systematics (~ 40% on signal yield)





# m<sub>yy</sub>=126.9 GeV



E EXPERIMEN

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### myy=126.9 GeV



Figure 31: Event display of a diphoton event candidate where both photon candidates are unconverted. The event number is 56662314 and it was recorded during run 203779 at  $\sqrt{s} = 8$  TeV. The leading photon has  $E_T$ =62.2 GeV and  $\eta$  =0.39. The subleading photon has  $E_T$ =55.5 GeV and  $\eta$  =1.18. The measured diphoton mass is 126.9 GeV. The  $p_T$  and  $p_{Tt}$  of the diphoton are 9.3 GeV and 6.5 GeV, respectively. Only reconstructed tracks with  $p_T$ >1 GeV, hits in the pixel and SCT layers and TRT hits with a high threshold are shown.



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### m<sub>yy</sub>=127.0 GeV



Figure 32: Event display of a diphoton with a jet event candidate where both photon candidates are unconverted. The event number is 36443051 and it was recorded during run 203195 at  $\sqrt{s} = 8$  TeV. The leading photon has  $E_T$ =63.0 GeV and  $\eta$  =0.50. The subleading photon has  $E_T$ =56.1 GeV and  $\eta$  = -0.96. The measured diphoton mass is 127.0 GeV. The  $p_T$  and  $p_{Tt}$  of the diphoton are 83.9 GeV and 83.3 GeV, respectively. The jet has  $E_T$ =113 GeV and  $\eta$  = -0.9. Only reconstructed tracks with  $p_T$ >1 GeV, hits in the pixel and SCT layers and TRT hits with a high threshold are shown.



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# $\hat{\mu} = \left\{ \mu \left| L(\mu s(m_H) + b) \right| = \max L(\mu, b) \right\}$ $\sigma / \sigma SM$ з Signal strength (μ)

#### Fitted signal strength

Normalized to SM Higgs expectation at given  $m_H(\mu)$ 

Best-fit value at 126.5 GeV:  $\mu$ =1.9 ± 0.5





Consistent results from various categories within uncertainties (most sensitive ones indicated)

### The Golden Channel H->ZZ->41



 $\left(s_i / \sqrt{s_i + b_i}\right)$ 

 $W_i \simeq -$ 

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# The Golden Channel: H->ZZ->41

- $\odot$  CLEAN but very low rate ( $\sigma$ ~2-5fb), yet probably most trustable
- All information is available, one can fully reconstruct the kinematics and the masses (m<sub>21</sub>, m<sub>41</sub>)
- Signature: Two pairs of same flavor opposite charged isolated leptons, one or both compatible with Z ->narrow peak
- Main backgrounds:
  - ZZ\* (irreducible)
  - for m<sub>H</sub><2m<sub>Z</sub>, Zbb, Z+jets, tt
- Suppress backgrounds with isolation and impact parameters cuts on two softest leptons





### The Golden Channel: H->ZZ->41







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#### $H \rightarrow 41$ mass spectrum after all selections: 2011+2012 data



m(41) > 160 GeV (dominated by ZZ background): 147 ± 11 events expected 191 observed

~ 1.3 more ZZ events in data than SM prediction  $\rightarrow$  in agreement with measured ZZ cross-section in 41 final states at  $\int s = 8$  TeV

#### $H \rightarrow 41$ mass spectrum after all selections: 2011+2012 data







	In the region 125 ± 5 GeV								
Dataset		2011	20	12	20	)11+20	12		
xpected B only		2±0.3		3±0.4		5.1±0.8			
xpected S m <sub>µ</sub> =125 GeV		<u>+0.3</u>	3±	3±0.5		5.3±0.8			
Observed	in the data	4	9		1	3			
2011+ 201	12	4 <sub>1</sub>	1 2	le2µ		4e			
Data		6		5		2			
Expected	IS/B	1.6		1	(	D.5			
Reducible/	/total B	5%	45	5%	55	5%			

 $4\mu$  candidate with  $m_{4\mu}$ = 125.1 GeV

 $p_T$  (muons)= 36.1, 47.5, 26.4, 71.7GeV  $m_{12}$ = 86.3 GeV,  $m_{34}$ = 31.6 GeV 15 reconstructed vertices



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4e candidate with  $m_{4e}$ = 124.6 GeV

#### $p_T$ (electrons)= 24.9, 53.9, 61.9, 17.8 GeV. $m_{12}$ = 70.6 GeV, $m_{34}$ = 44.7 GeV 12 reconstructed vertices



 $2e2\mu$  candidate with  $m_{2e2\mu}$ = 123.9 GeV

 $p_{T}$  (e,e,µ,µ)= 18.7, 76, 19.6, 7.9 GeV, m (e<sup>+</sup>e<sup>-</sup>)= 87.9 GeV, m(µ<sup>+</sup>µ<sup>-</sup>) = 19.6 GeV 12 reconstructed vertices





# Consistency of the data with the background-only expectation



#### Fitted signal strength

Best-fit value at 125 GeV:  $\mu$ =1.3 ± 0.6



# "TEVATRON" Channel H->WW





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# $H \rightarrow WW \rightarrow l \nu l \nu$ : $t \rightarrow e \mu$ background

Event display of a top pair e-mu dilepton candidate with two btagged jets. The electron is shown by the green track pointing to a calorimeter cluster, the muon by the long red track intersecting the muon chambers, and the missing ET direction by the dotted line on the XY view. The secondary vertices of the two b-tagged jets are indicated by the orange ellipses on the zoomed

vertex region view.



tag



- The cannel is challenging
   2 neutrinos- no mass reconstruction ->mT
- Signature: 2 high p<sub>T</sub> opposite sign isolated leptons with large E<sub>T</sub><sup>miss</sup>->Understanding of E<sub>T</sub><sup>miss</sup> is crucial
- Main background from WW, top, Z+jets, W+jets ->Use of control regions to estimate fakes
- A control region is defined rich in the measured BG (e.g. WW or top), contaminations are being subtracted and then the BG is extrapolated to the signal region (mostly using MC) Example: b-tag is inverted to estimate Top BG
- -> large E<sub>T</sub><sup>miss</sup>, m<sub>ll</sub> incompatible with m<sub>z</sub> (DY),
   -> b jet veto (tt),
   ->Topological cuts against irreducible WW
   (ΔΦ<sub>ll</sub>)
- Jet bins: +0j, +1, +2jet (VBF)
- Discriminating variable  $m_T = \sqrt{(E_T^{ll} + E_T^{miss})^2 + (p_T^{ll} + p_T^{miss})^2}$

#### $m_T = \sqrt{(E_T^{ll} + E_T^{miss})^2 + (p_T^{ll} + p_T^{miss})^2}$



 $N_{iets}$  with  $p_T > 25 \text{ GeV}$ 

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### $H \rightarrow WW \rightarrow l \nu l \nu$



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## The Fermionic Channels





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## H->bb: W/ZH->W/Zbb

 H->bb is the dominant decay of a low mass Higgs.
 It also extremely important

It also extremely important to measure Higgs couplings.

- Multi-jet background kills its inclusive production
- W/ZH is feasible for low Higgs mass channels: | v bb, | bb and v v bb
- Signature lepton,MET and b-tag
- Z/W+jets and tt BG can be reduced by requiring boosted Higgses with tight b-tag
- Sensitivity is still very low
   p0 expxted=0.15 obs= 0.64 (a deficit)

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- Observed  $H \rightarrow tt$ 

120

### $H \rightarrow \tau \tau$

#### VBF clean and sensitive

2 tagged back to back forward jets and two tagged taus



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Local p<sub>0</sub>

10

1

10-1

10-2

100

110

### All for one - Combine forces




## 4th July 2012: A long night



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## <u>A long night</u>



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# Peter Higgs



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# Englert



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## Where were you when....



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#### the Higgs boson

#### Global Effort → Global Success

Results today only possible due to extraordinary performance of accelerators – experiments – Grid computing

Observation of a new particle consistent with a Higgs Boson (but which one...?)

Historic Milestone but only the beginning

Global Implications for the future



#### Higgs: Really, I did not expect to see it in my lifetime

Rolf: I think we have it, you agree? We have a discovery. We should state it. We observed a new particle, consistent with the Higgs Boson. It is a milestone.

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## ATLAS Higgs Group



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## ATLAS Higgs Group



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## Disclaimer

- Correlated uncertainties (Jet energy scales, Luminosity etc... taken into account)
- When data driven methods are used, systematics are not correlated
- Theory uncertaintes are carefully taken into account across channels using the recommendation of the LHC Higgs cross section group









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Excellent consistency (better than 20!) of the data with the background-only hypothesis over full mass spectrum except in one region

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#### Evolution of the excess with time





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### A Prelude to the Nobel Prize

The 2004 Wolf prize, awarded by the Wolf Foundation, is often thought to be the most prestigious prize in physics after the Nobel prize (in the pic. Brout & Englert).



eilam gross, Bonn, December 2011

### A Prelude to the Nobel Prize

2010 Sakurai Prize awarded for 1964 Higgs Boson theory work to Hagen, Guralnik, Kibble, Brout, Englert & Higgs



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#### 4th July 2012

### Higgs Hunters Independence Day

### ATLAS T-shirt

This is the most incredible thing that have happned to me in my lifetime

P. Higgs 4th July 2012



eilam gross, Bonn, December 2011

### A Phenomenological Profile of the Higgs Boson



Nuclear Physics B106 (1976) 292-340 © North-Holland Publishing Company

#### A PHENOMENOLOGICAL PROFILE OF THE HIGGS BOSON

John ELLIS, Mary K. GAILLARD \* and D.V. NANOPOULOS \*\* CERN, Geneva

Received 7 November 1975

The situation with regard to Higgs bosons is unsatisfactory. First it should be stressed that they may well not exist. Higgs bosons are introduced to give intermediate vector bosons masses through spontaneous symmetry breaking. However, this symmetry breaking could be achieved dynamically [10] without elementary Higgs bosons. Thus the confirmation or exclusion of their existence would be an important constraint on gauge theory model building. Unfortunately, no way is known to calculate the mass of a Higgs boson, at least in the context of the popular Weinberg-Salam [11]

Eilam Gross, WIS, Freiburg Jan 2012

### A Phenomenological Profile of the Higgs Boson



Nuclear Physics B106 (1976) 292-340 © North-Holland Publishing Company

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Received 7 November 1975

We should perhaps finish with an apology and a caution. We apologize to experimentalists for having no idea what is the mass of the Higgs boson, unlike the case with charm [3,4] and for not being sure of its couplings to other particles, except that they are probably all very small. For these reasons we do not want to encourage big experimental searches for the Higgs boson, but we do feel that people performing experiments vulnerable to the Higgs boson should know how it may turn up.

Eilam Gross, WIS, Freiburg Jan 2012





