



# Higgs properties in a Stealth Doublet Model

Johan Rathsman, Lund university  
based on work with Rikard Enberg and Glenn Wouda,  
arXiv:1304.1714 [hep-ph], accepted for publication in JHEP

Nordita, 2013-07-27

- 1 Higgs discovery
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# Higgs discovery



# Observations in agreement with standard model

CMS

ATLAS

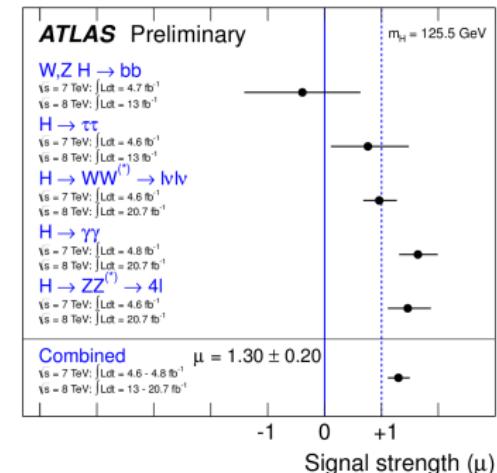
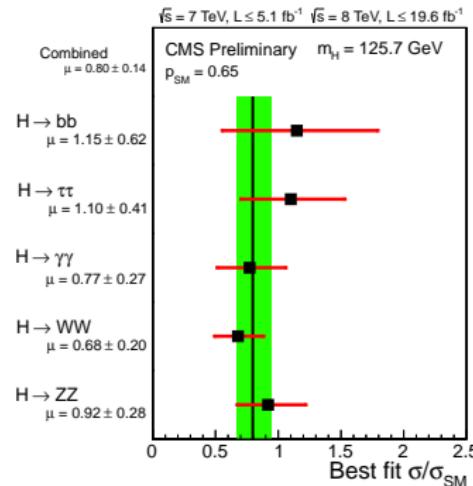
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What (if anything) is beyond?

Bottom up approach in two Higgs Doublet model



# Constraints from flavour sector - neutral currents

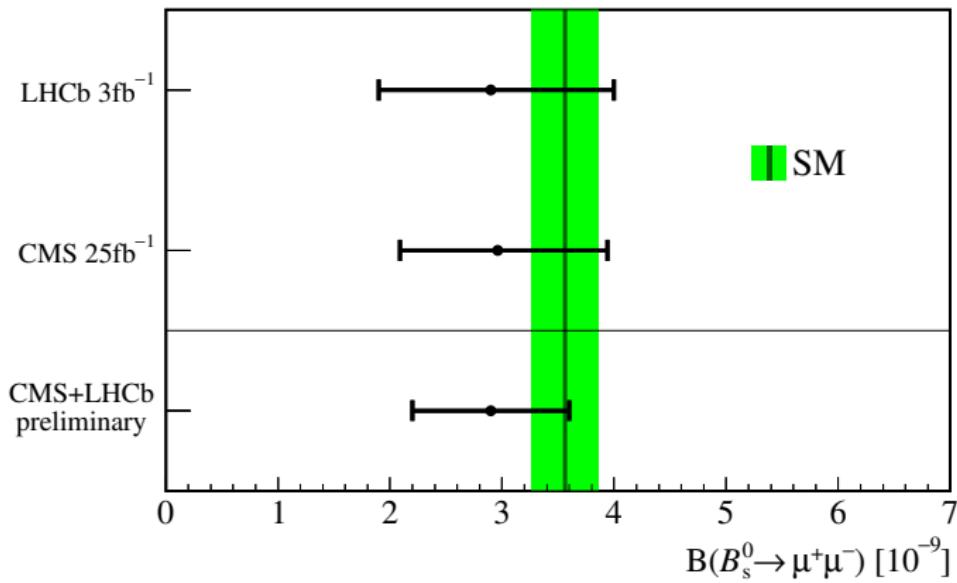
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also from  $b \rightarrow s\gamma$

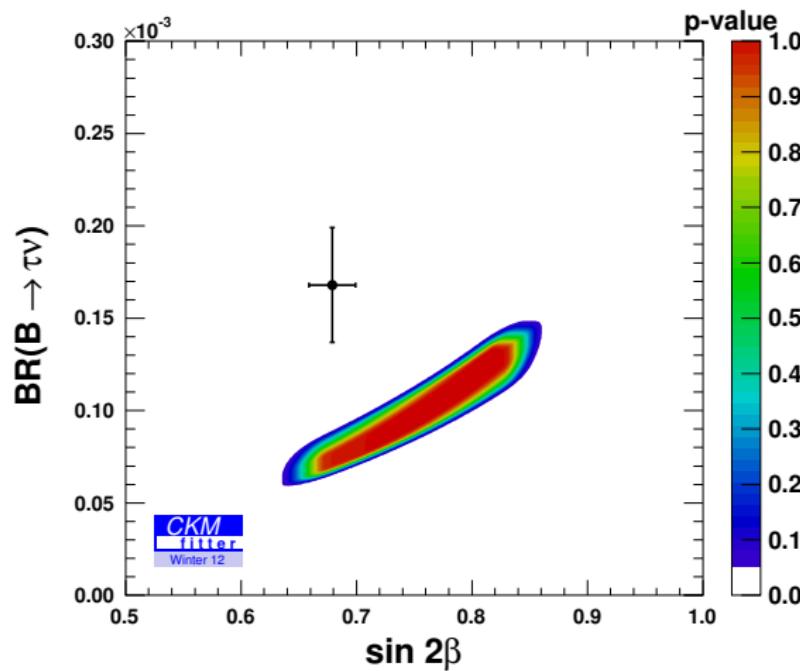


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# Constraints from flavour sector - charged currents

$B^+ \rightarrow \tau^+ \nu_\tau$ , winter 2012

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experimental errors at  $1\sigma$

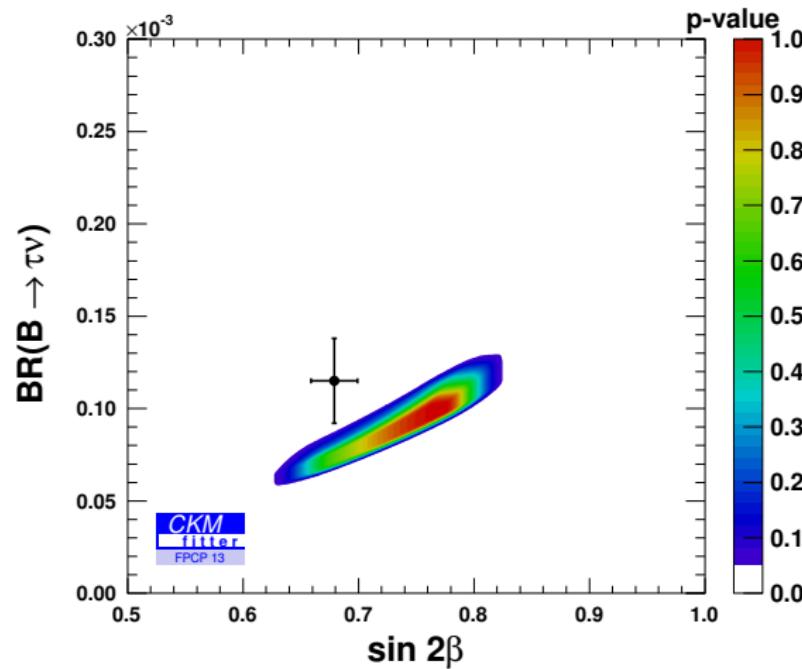


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# Constraints from flavour sector - charged currents

$B^+ \rightarrow \tau^+ \nu_\tau$ , summer 2013

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2HDM potential

Yukawa sector

EWSB

Higgs fermion  
couplings

Mass spectrum

Soft  $\mathbb{Z}_2$  breaking

TH constraints

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# The Stealth Doublet Model



# General two Higgs doublet model potential

- Two complex  $SU(2)_L$  doublets with hypercharge  $Y=1$ :  $\Phi_1, \Phi_2$
- Invariance under global  $SU(2)$ :  $\Phi_a \rightarrow U_{ab} \Phi_b$

## General potential

$$\begin{aligned}\mathcal{V} = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - \left[ m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h.c.} \right] + \frac{1}{2} \lambda_1 \left( \Phi_1^\dagger \Phi_1 \right)^2 \\ & + \frac{1}{2} \lambda_2 \left( \Phi_2^\dagger \Phi_2 \right)^2 + \lambda_3 \left( \Phi_1^\dagger \Phi_1 \right) \left( \Phi_2^\dagger \Phi_2 \right) + \lambda_4 \left( \Phi_1^\dagger \Phi_2 \right) \left( \Phi_2^\dagger \Phi_1 \right) \\ & + \left\{ \frac{1}{2} \lambda_5 \left( \Phi_1^\dagger \Phi_2 \right)^2 + \left[ \lambda_6 \left( \Phi_1^\dagger \Phi_1 \right) + \lambda_7 \left( \Phi_2^\dagger \Phi_2 \right) \right] \left( \Phi_1^\dagger \Phi_2 \right) + \text{h.c.} \right\}\end{aligned}$$

- Potential real  $\Rightarrow \{m_{11}^2, m_{22}^2, \lambda_{1-4}\}$  real,  $\{m_{12}^2, \lambda_{5-7}\}$  complex
- No explicit CP-violation  $\Rightarrow \{m_{12}^2, \lambda_{5-7}\}$  real

## Softly broken $\mathbb{Z}_2$ symmetry

- Potential symmetric under  $\Phi_1 \rightarrow \Phi_1, \Phi_2 \rightarrow -\Phi_2$   
 $\Rightarrow m_{12}^2 = 0, \lambda_{6-7} = 0$  in general basis
- $m_{12}^2 \neq 0 \Rightarrow$  soft breaking



# Yukawa sector

General Yukawa couplings for SM fermions (neutrinos massless)

$$\begin{aligned} -\mathcal{L}_{\text{Yuk}} = & \kappa_0^L \bar{L}_L \Phi_1 E_R + \kappa_0^U \bar{Q}_L \tilde{\Phi}_1 U_R + \kappa_0^D \bar{Q}_L \Phi_1 D_R \\ & + \rho_0^L \bar{L}_L \Phi_2 E_R + \rho_0^U \bar{Q}_L \tilde{\Phi}_2 U_R + \rho_0^D \bar{Q}_L \Phi_2 D_R \end{aligned}$$

where  $\tilde{\Phi}_i = -i\sigma_2 \Phi_i^*$

Impose even  $\mathbb{Z}_2$  parities on all fermions  $\Rightarrow \rho^F = 0$  at tree level



# Electroweak symmetry breaking

Fermions only couple to  $\Phi_1 \Rightarrow$  Higgs basis physically realized

- EW symmetry broken by non-zero vev of  $\Phi_1$  ( $v \approx 246$  GeV)

$$\Phi_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2}G^+ \\ v - h \sin \alpha + H \cos \alpha + iG^0 \end{pmatrix}$$

$$\Phi_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2}H^+ \\ h \cos \alpha + H \sin \alpha + iA \end{pmatrix}$$

- Three Goldstone bosons:  $G^\pm, G^0 \Rightarrow$  masses to  $W$  and  $Z$
- Two CP-even Higgs boson states:  $h, H$  with mixing angle  $\alpha$
- Three scalar states: one CP-odd  $A$  and two charged  $H^\pm$
- Higgs-gauge couplings from  $s_\alpha \equiv \sin \alpha, c_\alpha \equiv \cos \alpha$
- $\tan \beta$  not defined



# Higgs fermion couplings

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Yukawa couplings for SM fermions with mass eigenstates

$$D = \{d, s, b\}, U = \{u, c, t\}, L = \{e, \mu, \tau\}$$

$$-\mathcal{L}_{\text{Yuk}} = \frac{1}{v} \left( \sum_D \overline{D} m_D D + \sum_U \overline{U} m_U U + \sum_L \overline{L} m_L L \right) (s_\alpha h - c_\alpha H)$$



# Mass spectrum

Minimization:

$$m_{11}^2 = -\frac{1}{2}v^2\lambda_1$$

$$m_{12}^2 = \frac{1}{2}v^2\lambda_6$$

$$m_{22}^2 \quad \text{unconstrained}$$

## Masses of scalars

$$m_A^2 = m_{22}^2 + \frac{1}{2}v^2(\lambda_3 + \lambda_4 - \lambda_5)$$

$$m_{H^\pm}^2 = m_{22}^2 + \frac{1}{2}v^2\lambda_3$$

Mass-matrix of Higgses ( $\lambda_{345} = \lambda_3 + \lambda_4 + \lambda_5$ )

$$\mathcal{M}^2 = \begin{pmatrix} \lambda_1 v^2 & \lambda_6 v^2 \\ \lambda_6 v^2 & m_{22}^2 + \lambda_{345} v^2 \end{pmatrix}$$

Diagonalization:  $\sin 2\alpha = \frac{2v^2\lambda_6}{m_H^2 - m_h^2} = \frac{4m_{12}^2}{m_H^2 - m_h^2}$

- $\sin \alpha \propto m_{12}^2$  ( $m_{12}^2 = 0$  restores  $Z_2$  symmetry)



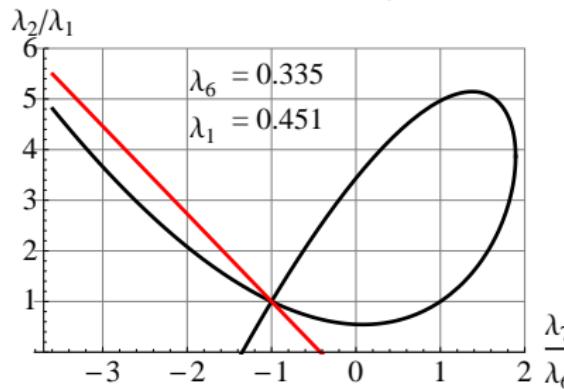
# Conditions for soft $\mathbb{Z}_2$ breaking

$m_{12}^2 = \frac{1}{2}\nu^2\lambda_6 \Rightarrow m_{12}^2 \neq 0$  gives hard  $\mathbb{Z}_2$  breaking?

Sufficient if there exists a basis where  $\lambda_6 = \lambda_7 = 0$  and  $m_{12}^2 \neq 0$   
(Davidson and Haber '05)

$$(\lambda_1 - \lambda_2) [\lambda_{345}(\lambda_6 + \lambda_7) - \lambda_2\lambda_6 - \lambda_1\lambda_7] - 2(\lambda_6 - \lambda_7)(\lambda_6 + \lambda_7)^2 = 0$$
$$(\lambda_1 - \lambda_2)m_{12}^2 + (\lambda_6 + \lambda_7)(m_{11}^2 - m_{22}^2) \neq 0$$

Solution  $\Rightarrow \lambda_2$  fixed,  $\lambda_7$  constrained (upper limit)



black line allowed  
red line not allowed

note:

$\lambda_2 = \lambda_1$ ,  $\lambda_7 = \lambda_6$   
always allowed

Parameterisation of potential:  $\{ \lambda_3, \lambda_7, m_h, m_H, m_A, m_{H^\pm}, s_\alpha \}$



# Theoretical constraints

## Positivity of potential

Demanding that the potential is bounded from below  $\Rightarrow$

$$\lambda_1 > 0, \quad \lambda_2 > 0, \quad \lambda_3 > -\sqrt{\lambda_1 \lambda_2}, \quad \lambda_3 + \lambda_4 - \lambda_5 > -\sqrt{\lambda_1 \lambda_2}$$

plus more complicated expressions

## Perturbativity

Cross-section for  $2 \rightarrow 2$  Higgs scattering processes  $\propto \frac{\lambda_{HHHH}^2}{16\pi^2}$

$\Rightarrow$  the quartic Higgs couplings  $\lambda_{HHHH}$  cannot be too large for the perturbative series to make sense

## Tree-level unitarity

requiring tree-level unitarity for  $HH$  and  $HV_L$  scattering  $\Rightarrow$  limits on eigenvalues of the corresponding scattering matrices



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Higgs signal

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# Mass spectrum

Agreement between LHC data and SM  $\Rightarrow$  two possibilities

- ①  $m_h=125 \text{ GeV}, s_\alpha \sim 0.9, m_H \gtrsim 300 \text{ GeV}$
- ②  $m_H=125 \text{ GeV}, s_\alpha \sim 0.1, m_h \sim 75 \text{ GeV}$

# Electroweak precision tests

- using 2HDMC <http://2hdmc.hepforge.org/> (JR, O. Stål)

$m_h/H = 125/300 \text{ GeV}$ ,  $s_\alpha = 0.9$

$m_h/H = 75/125 \text{ GeV}$ ,  $s_\alpha = 0.1$

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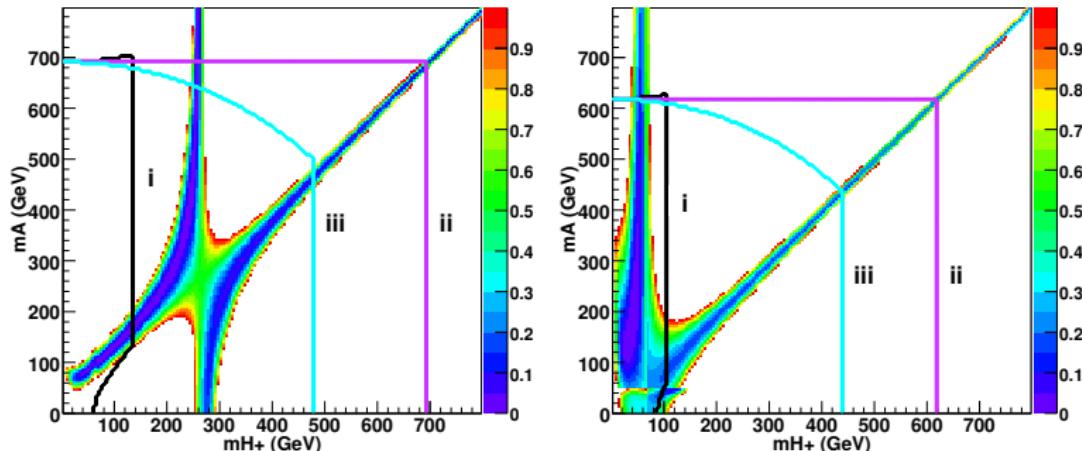
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- Allowed regions (90% C.L.) - custodial symmetry:  
 $m_A \approx m_{H^\pm}$  or  $m_{H^\pm}^2 \approx m_H^2 s_\alpha^2 + m_h^2 c_\alpha^2$
- Theoretical constraints (for  $\lambda_2 = \lambda_1$  and  $\lambda_7 = \lambda_6$ ):  
(i)  $\lambda_3 = 0$ , (ii)  $\lambda_3 = 2m_{H^\pm}^2/v^2$ , (iii)  $\lambda_3 = 4m_{H^\pm}^2/v^2$   
regions inside the lines are allowed



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# Compatibility with Higgs signal

Signal strength compared to SM ( $\mathcal{H} = h/H$ )

$$\mu_{\mathcal{H}XX} = \frac{\sum_i \sigma_i(pp \rightarrow \mathcal{H}) \text{BR}(\mathcal{H} \rightarrow XX)}{\sum_i \sigma_i(pp \rightarrow H_{\text{SM}}) \text{BR}(H_{\text{SM}} \rightarrow XX)}$$

universal rescaling of couplings to fermions and  $W/Z$  bosons

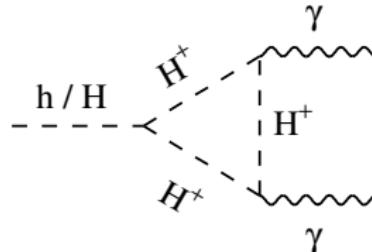
$$\mu_{h\gamma\gamma} = s_\alpha^2 \frac{\text{BR}(h \rightarrow \gamma\gamma)}{\text{BR}(H_{\text{SM}} \rightarrow \gamma\gamma)}, \quad \mu_{hff} = \mu_{hVV} = s_\alpha^2$$

$$\mu_{H\gamma\gamma} = c_\alpha^2 \frac{\text{BR}(H \rightarrow \gamma\gamma)}{\text{BR}(H_{\text{SM}} \rightarrow \gamma\gamma)}, \quad \mu_{Hff} = \mu_{HVV} = c_\alpha^2$$

assuming  $m_A, m_{H^\pm} > m_h/2$  for  $\mu_{Hff}, \mu_{HVV}$



# Enhancement of $h/H \rightarrow \gamma\gamma$ from charged scalar



$$g_{hH^+H^-} = -iv(-s_\alpha\lambda_3 + c_\alpha\lambda_7)$$

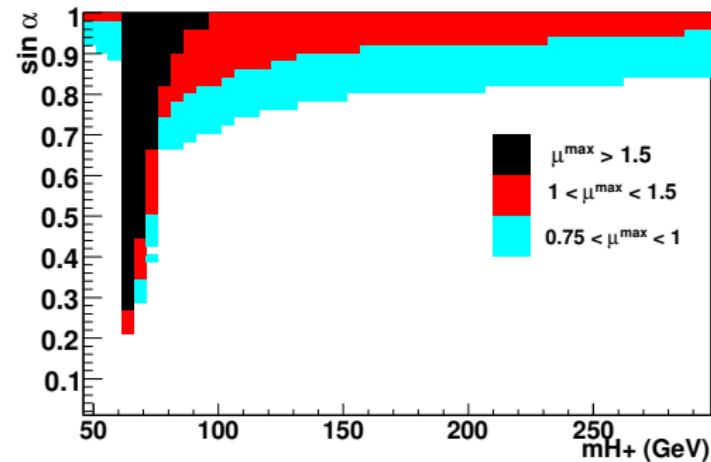
$$g_{HH^+H^-} = -iv(c_\alpha\lambda_3 + s_\alpha\lambda_7)$$

scan:  $-5 \leq \lambda_3 \leq 5$ ,  $-5 \leq \lambda_7 \leq \min\{5, \lambda_7^{\max} = \lambda_6 + \frac{(\lambda_1 - \lambda_{345})^2}{8\lambda_6}\}$

$\mu_{h\gamma\gamma}$

$m_h = 125$  GeV  
 $m_H = 300$  GeV  
 $m_A = m_{H^\pm} + 50$  GeV  
 $\lambda_2 = \lambda_1$

Other channels  
 $s_\alpha^2 \gtrsim 0.7$



- Other Higgs constraints included using HiggsBounds

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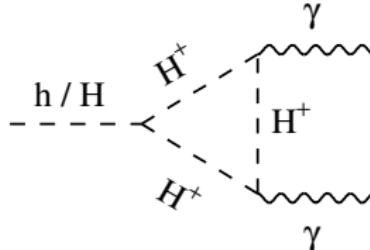
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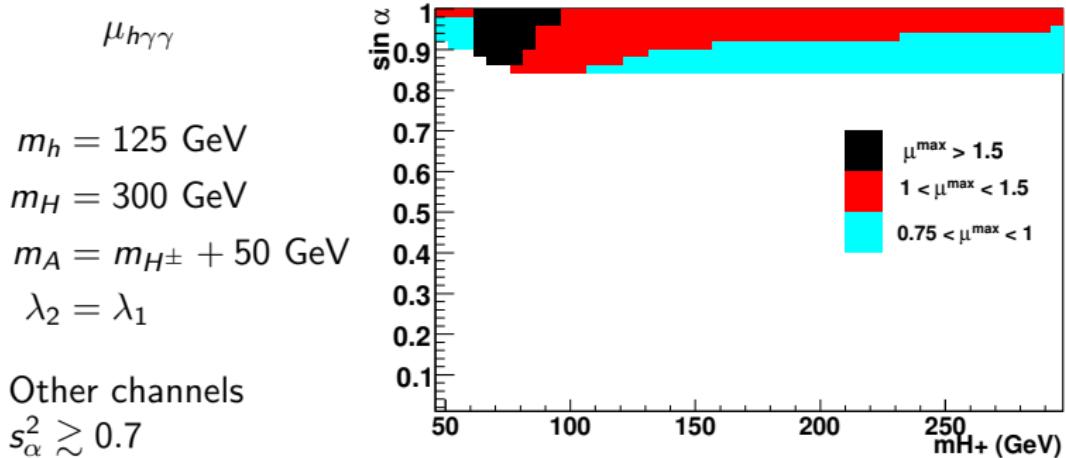
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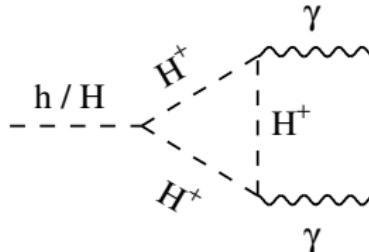
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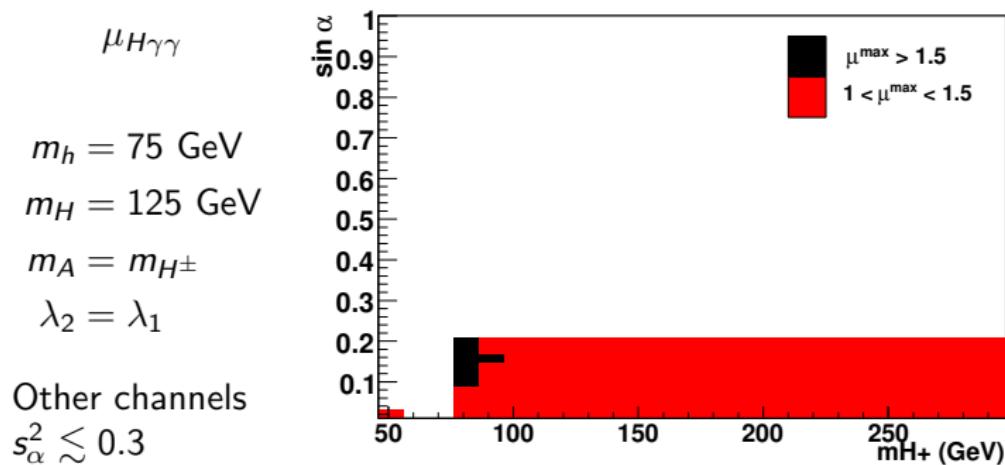
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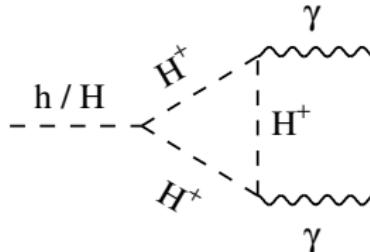
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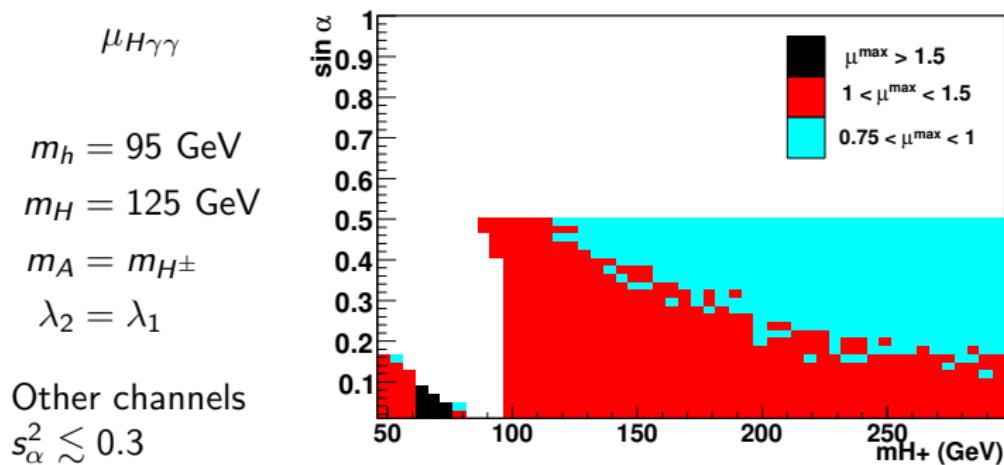
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scan:  $-5 \leq \lambda_3 \leq 5$ ,  $-5 \leq \lambda_7 \leq \min\{5, \lambda_7^{\max}\}$ ,  $\lambda_7^{\max} = \lambda_6 + \frac{(\lambda_1 - \lambda_{345})^2}{8\lambda_6}$

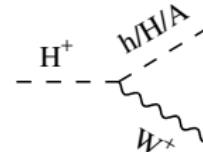


- Other Higgs constraints included using HiggsBounds



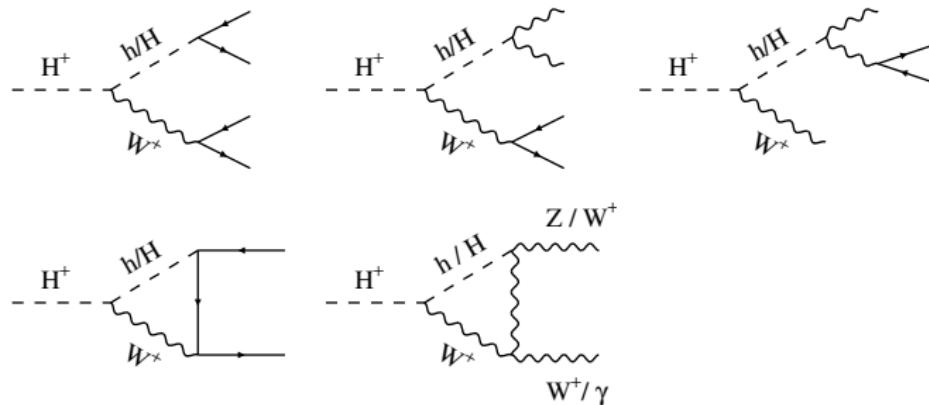
# Other signals at LHC - charged scalar

Basic decay vertex



$$\cos\alpha / \sin\alpha / 1$$

Decays into fermions and SM gauge bosons ( $m_{H^\pm} < m_A$ )



- Note: all diagrams proportional to  $\sin(2\alpha)$   $\Rightarrow$  vanish in no-mixing limit  $\sin\alpha \rightarrow 0$  or  $\cos\alpha \rightarrow 0$
- Proper calculation of loops require (on-shell) renormalisation

Tree-level dominates if open (on-shell or slightly off-shell)



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Examples:

$$m_h = 125 \text{ GeV}$$

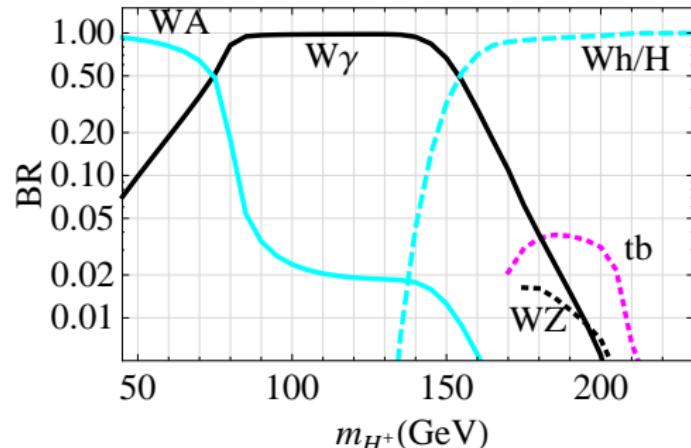
$$m_H = 300 \text{ GeV}$$

$$s_\alpha = 0.9$$

$$m_A = m_{H^\pm} - 10 \text{ GeV}$$

$$\lambda_3 = 2m_{H^\pm}^2/v^2$$

$$\lambda_2 = \lambda_1, \lambda_7 = \lambda_6$$



$$m_h = 125 \text{ GeV}$$

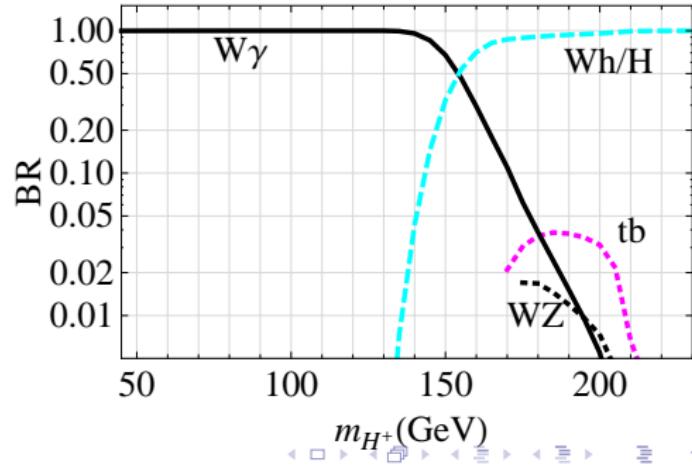
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$$m_A = m_{H^\pm}$$

$$\lambda_3 = 2m_{H^\pm}^2/v^2$$

$$\lambda_2 = \lambda_1, \lambda_7 = \lambda_6$$



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# Other signals at LHC - pseudoscalar

Similar to  $H^\pm$  except no decay to pair of SM gauge bosons

Example:

$$m_h = 125 \text{ GeV}$$

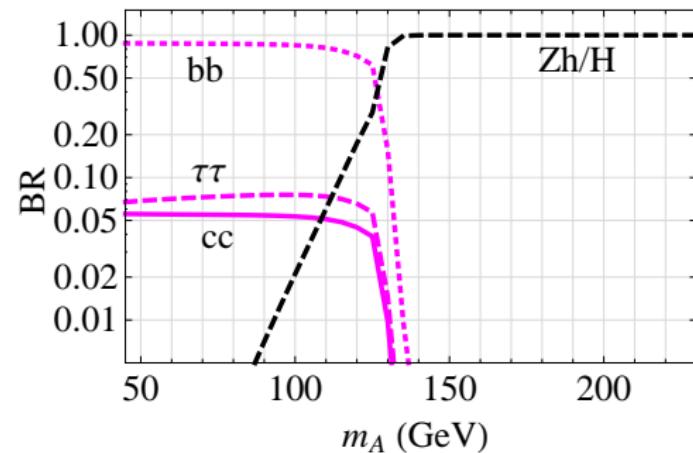
$$m_H = 300 \text{ GeV}$$

$$s_\alpha = 0.9$$

$$m_A = m_{H^\pm}$$

$$\lambda_3 = 0$$

$$\lambda_2 = \lambda_1, \lambda_7 = \lambda_6$$





# Production of $H^\pm$ and $A$ through Drell-Yann

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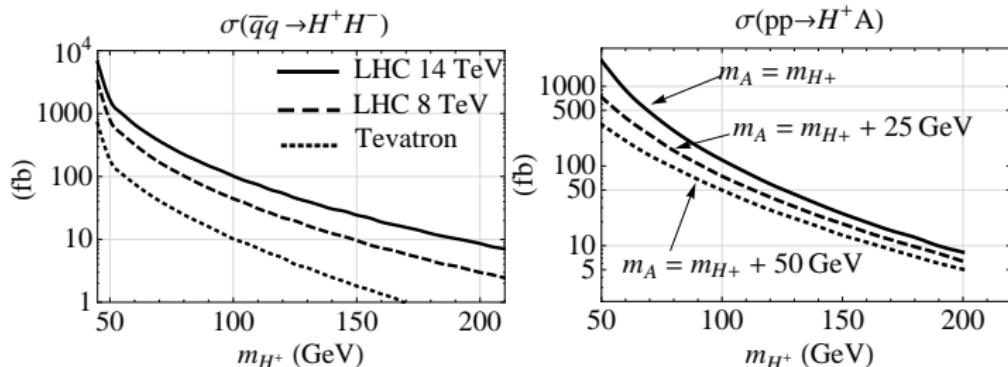
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# Conclusions

## Higgs discovery

- precision measurements probe for physics beyond the SM

## Stealth Doublet model

- naturally avoids flavour constraints
- makes definite predictions for Higgs properties
- predicts extra scalars with unusual decay modes

more work needed to investigate possibility of observing extra scalars