Magnetic & Rayleigh Dark Matter

Itay Yavin

McMaster University

Perimeter Institute

MITP

Schloss Waldthausen

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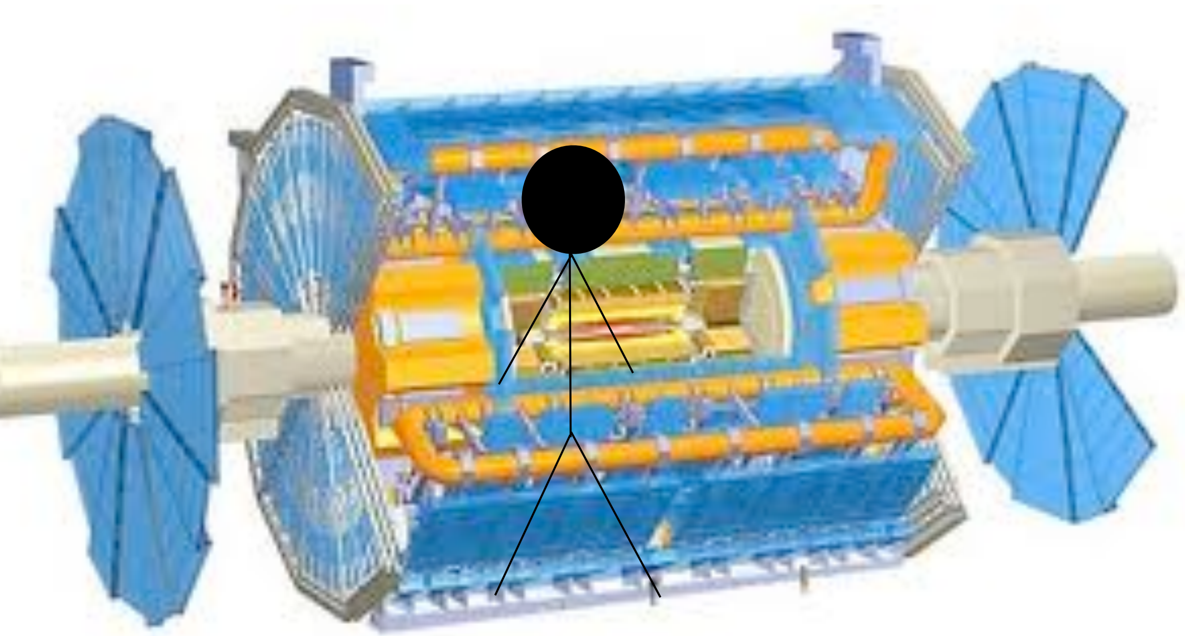
Thinking outside the SUSY box

(only one slide of philosophy, I promise . . .)



How dark is Dark Matter

(effective interactions with light)

Dark, but not too dark

From Effective Field Theory point of view, a natural reason why DM is dark is that its interactions with light are all coming through irrelevant operators.

$$\mathcal{L} \quad = \quad \text{massive fermion } \chi$$

$$+ \quad \frac{1}{2}\mu_\chi \; \bar{\chi}\sigma_{\mu\nu}\chi B^{\mu\nu}$$

Chang, Weiner, and IY arXiv:1007.4200

Dark, but not too dark

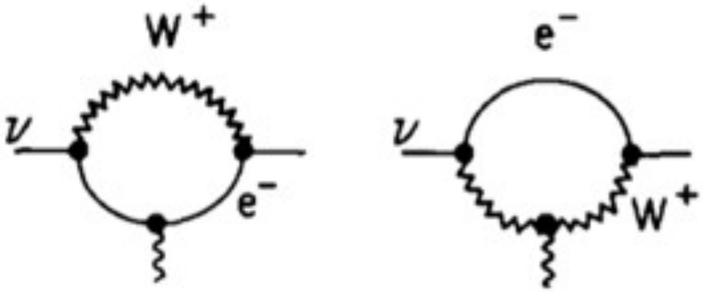
We have many examples of neutral objects with a magnetic dipole:

Neutrons

Charged constituents result in a magnetic dipole moment.

Neutrinos

Virtual cloud of charged matter results in magnetic dipole.



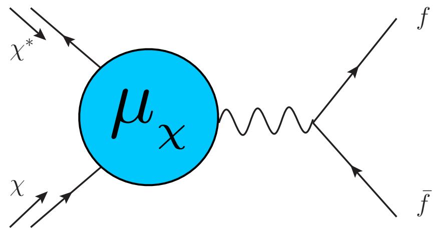
From Duncan, Grifols, Mendez, and Uma Sankar, Phys.Lett.B 191 (1987) 304

Notice that the mass difference of neutrinos may allow one neutrino to decay to another via a photon emission.

Phenomenology

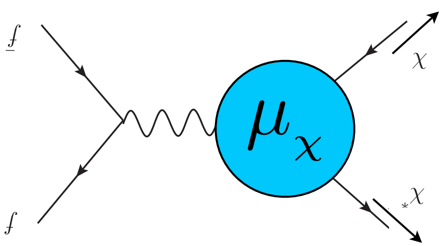
This single vertex contributes to a variety of observable processes

Annihilation



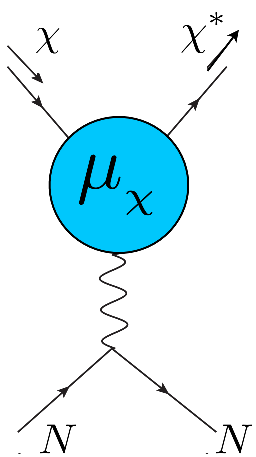
This process in the early universe can lead to the correct relic abundance of dark matter.

Production



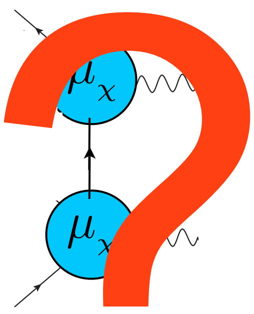
Dark matter in colliders

Direct detection



This process can be searched for in direct detection experiments looking for dark matter in the lab.

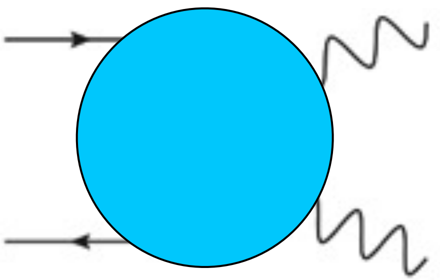
Gamma-rays



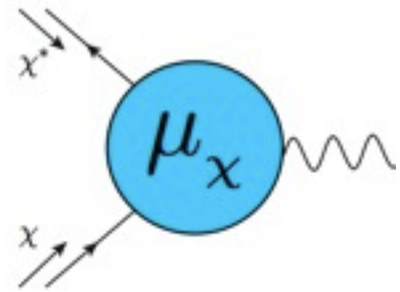
This process [illegible]n be searched for in gamma-ray lines in astrophysical observations (e.g. galactic center).

Rayleigh Dark Matter

Neutral particles also have two-photon interactions leading to Rayleigh scattering (the blue sky. . .) Weiner and IY arXiv:1206.2910



difficult to avoid ←



$$\mathcal{L}_{\mathrm{MiDM}} = \left(\frac{\mu_\chi}{2}\right) \bar{\chi}^* \sigma_{\mu\nu} B^{\mu\nu} \chi + c.c.$$

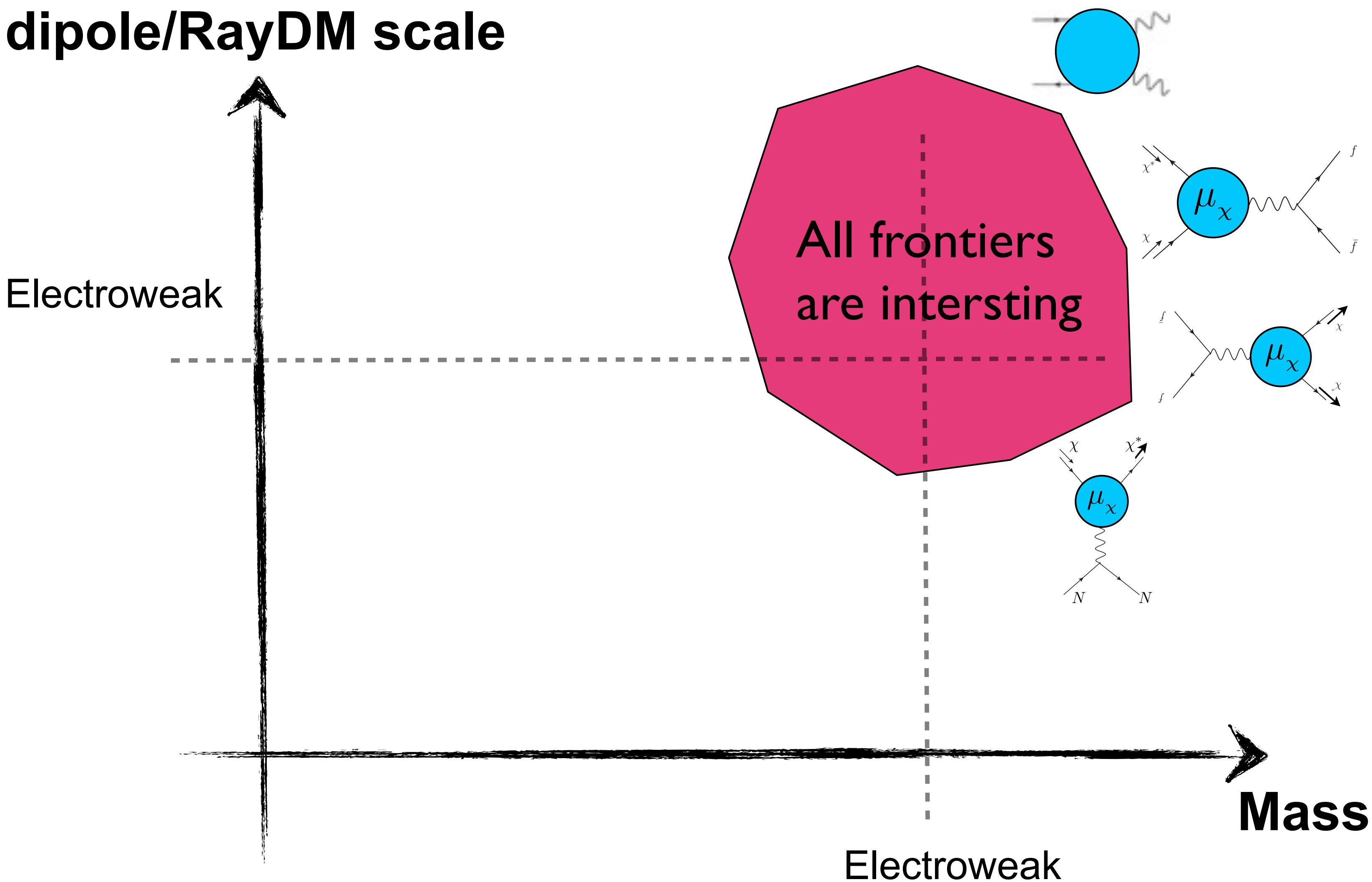
Magnetic (inelastic) Dark Matter

$$\mathcal{L}_{\mathrm{RayDM}} = \frac{1}{4\Lambda_R^3} \bar{\chi}\chi \left(\cos\theta_\chi B_{\mu\nu}B^{\mu\nu} + \sin\theta_\chi \mathrm{Tr} W_{\mu\nu}W^{\mu\nu}\right)$$

Rayleigh Dark Matter

$$+ \frac{i}{4\tilde{\Lambda}_R^3} \bar{\chi}\gamma_5\chi \left(\cos\theta_\chi B_{\mu\nu}\tilde{B}^{\mu\nu} + \sin\theta_\chi \mathrm{Tr} W_{\mu\nu}\tilde{W}^{\mu\nu}\right)$$

Interesting Parameter Space



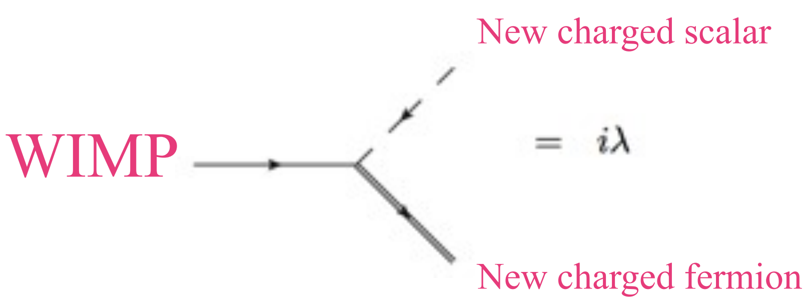
Weiner and IY arXiv:1206.2910

New charged states

(the microscopic origin of these interactions)

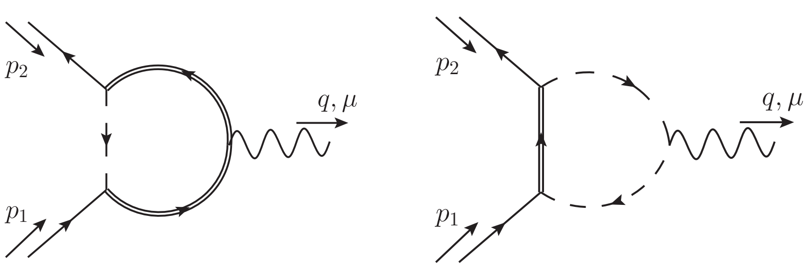
New Charged States

The MiDM and RayDM operators arise from integrating out charged matter that couples to the WIMP. Charged matter at the electroweak scale is necessary,

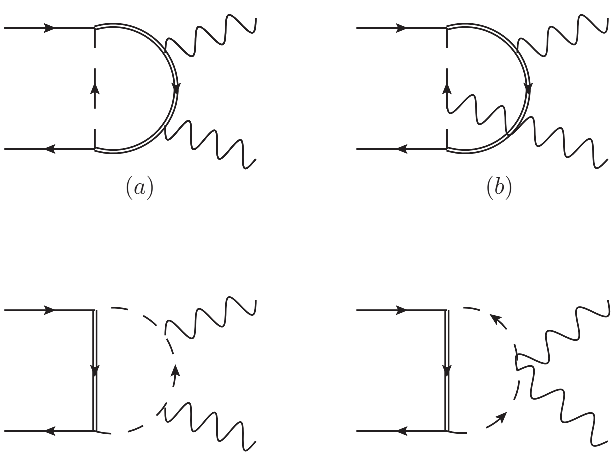


$$\begin{aligned}\mathcal{L} = {} & \bar{\chi}\left(i\partial\!\!\!/ - m_\chi\right)\chi + \bar{\psi}\left(iD\!\!\!\!/ - M_f\right)\psi \\ & + (D^\mu\varphi)^\dagger D_\mu\varphi - M_s^2\varphi^\dagger\varphi \\ & + \lambda\bar{\psi}\chi\varphi - \frac{\kappa}{4}\left|\varphi^\dagger\varphi\right|^2 + \text{h.c.}\end{aligned}$$

Magnetic (inelastic) Dark Matter

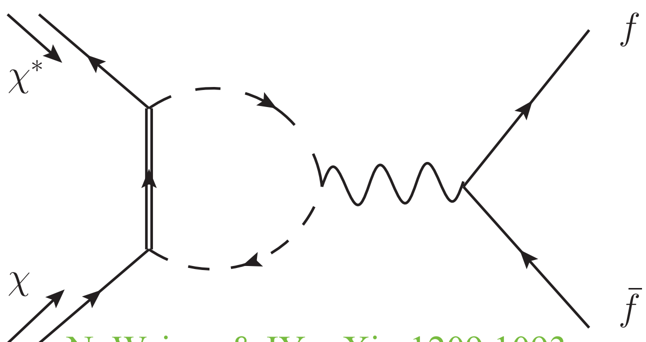


Rayleigh Dark Matter



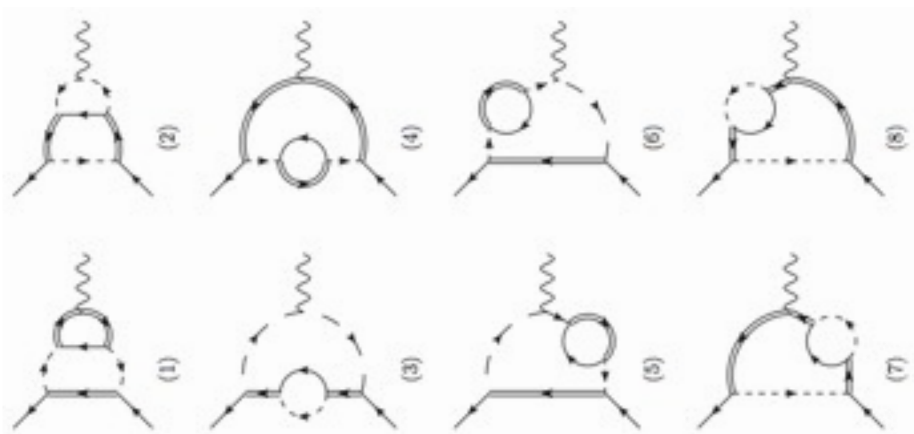
Relic Abundance

The dark matter density observed today is determined by the annihilation of DM into charged pairs through the dipole operator in the early Universe,

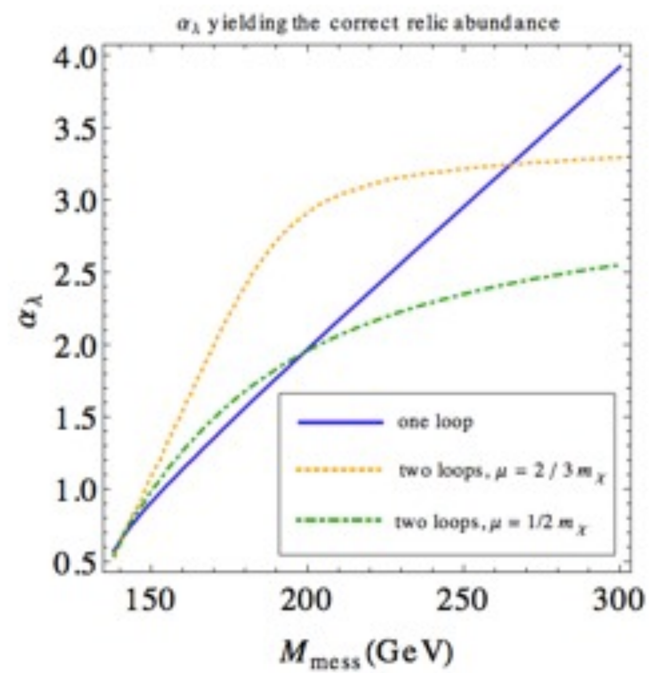


N. Weiner & IY arXiv:1209.1093

$$\mathcal{L} = \bar{\chi}\left(i\partial\!\!\!/ - m_\chi\right)\chi + \bar{\psi}\left(iD\!\!\!\!/ - M_f\right)\psi + \left(D^\mu\varphi\right)^\dagger D_\mu\varphi - M_s^2\varphi^\dagger\varphi + \lambda\bar{\psi}\chi\varphi - \frac{\kappa}{4}\left|\varphi^\dagger\varphi\right|^2 + \text{h.c.}$$

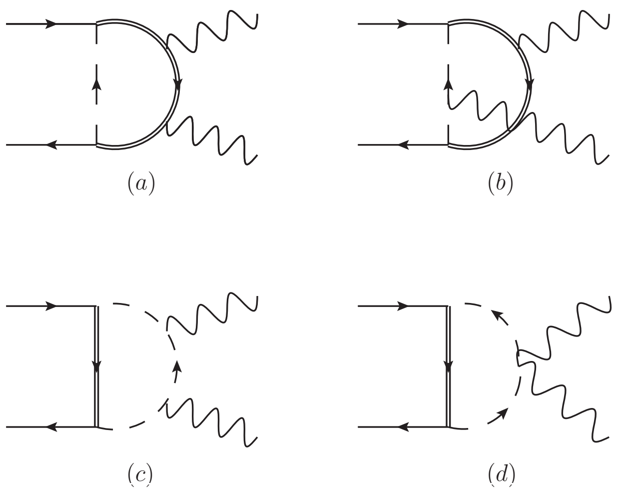


C. Tamarit & IY arXiv:1305.2951

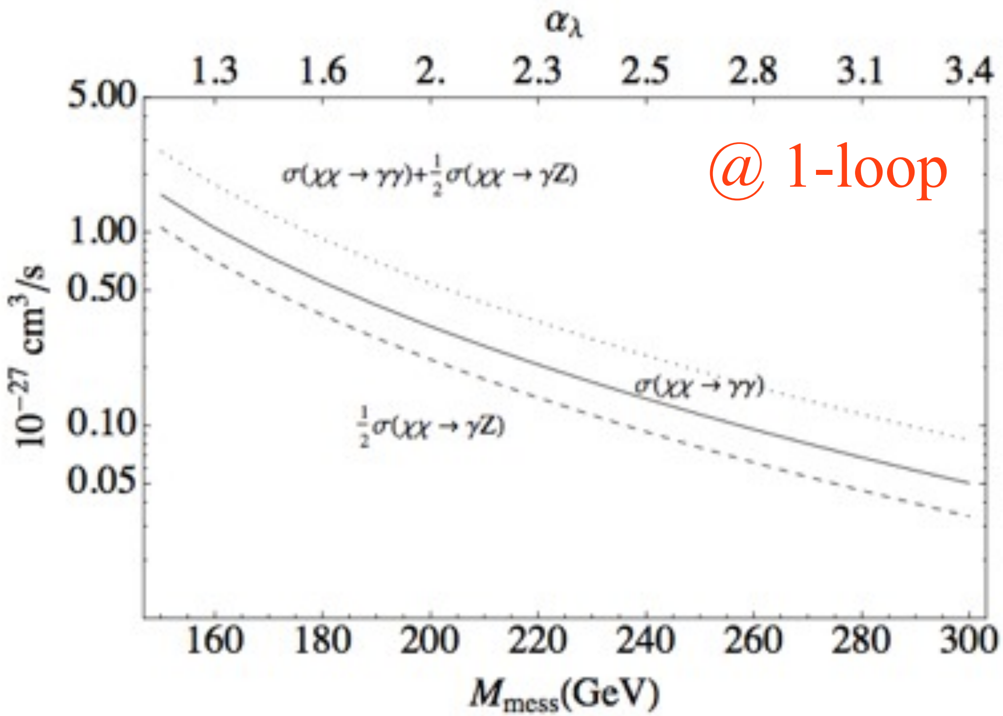


Gamma Rays

Annihilation today is dominated by the Rayleigh operator (no excessive continuum),



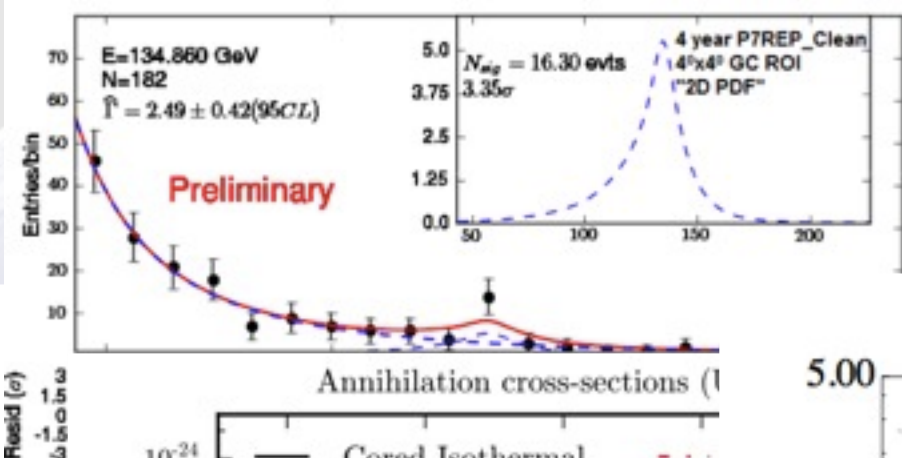
By electroweak gauge invariance we expect annihilations into photons as well as Z-bosons.



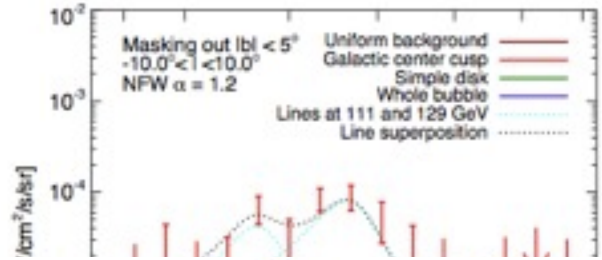
The Fermi Line

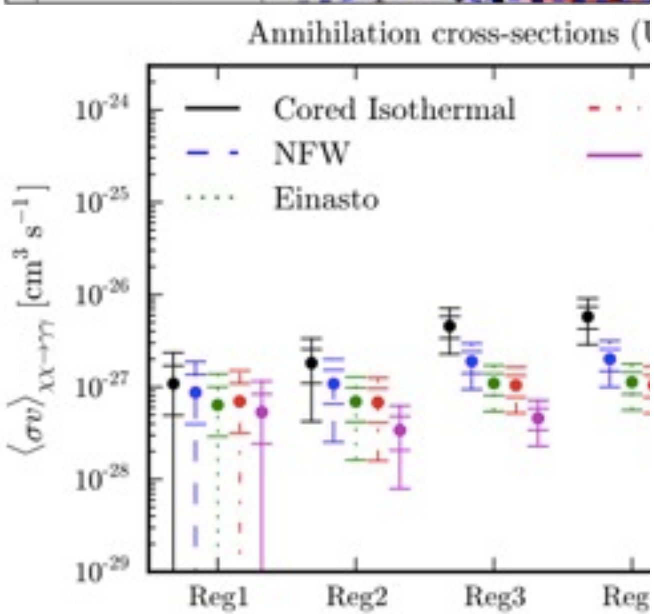
FERMI data reveals an excess at 135 GeV that can be interpreted as Dark Matter annihilation.

Albert for FERMI

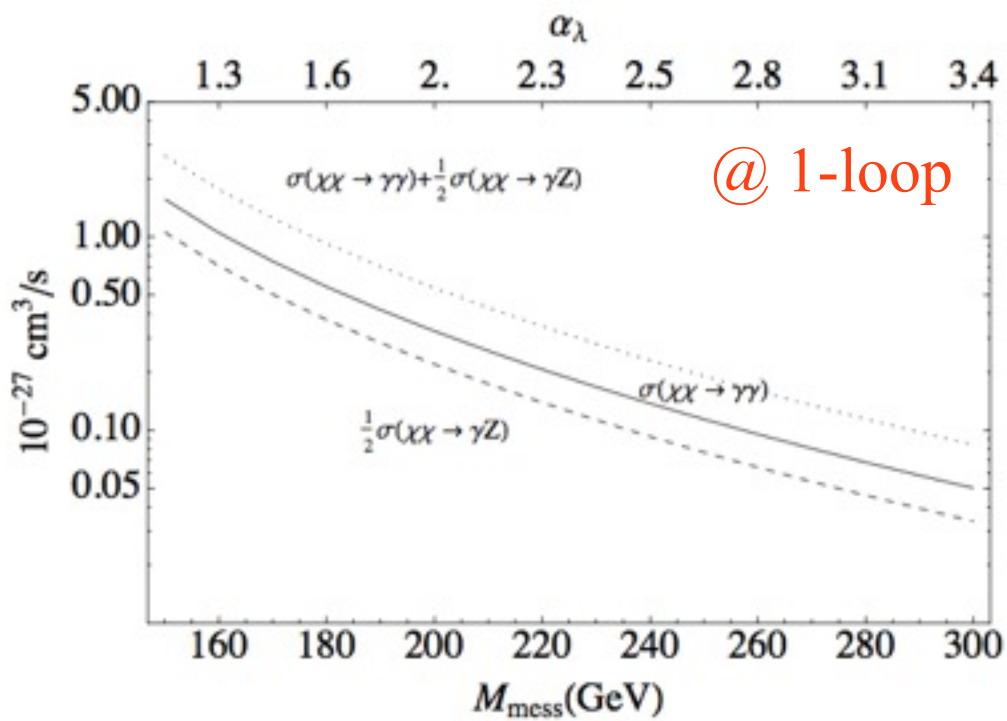


Finkbeiner & Su arXiv:1206.1616





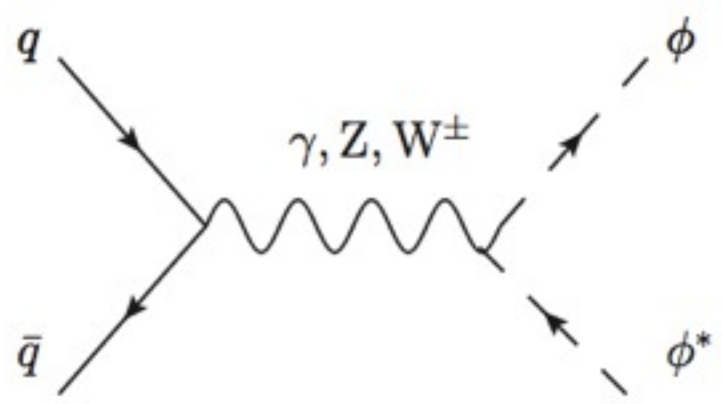
Weniger arXiv:1204.2797

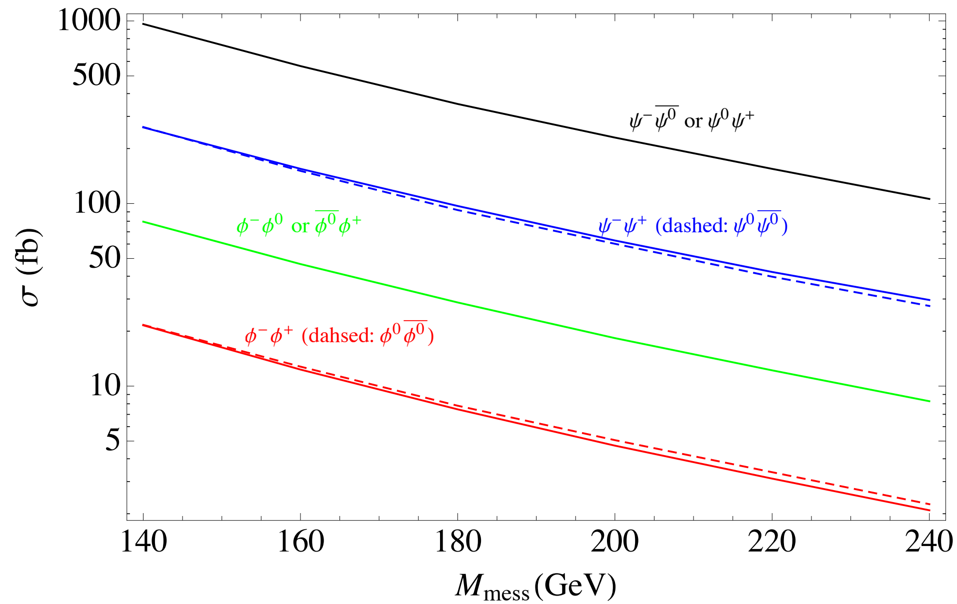


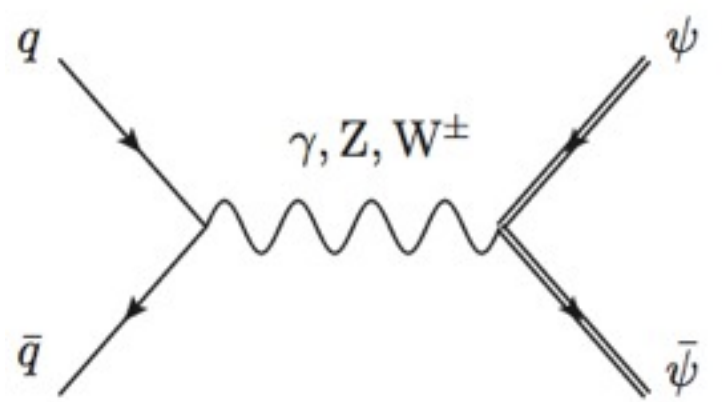
Already Produced at the LHC

These charged states were already produced at the LHC through their gauge interactions:

(Liu, Shuve, Weiner, IY arXiv:1303.4404)



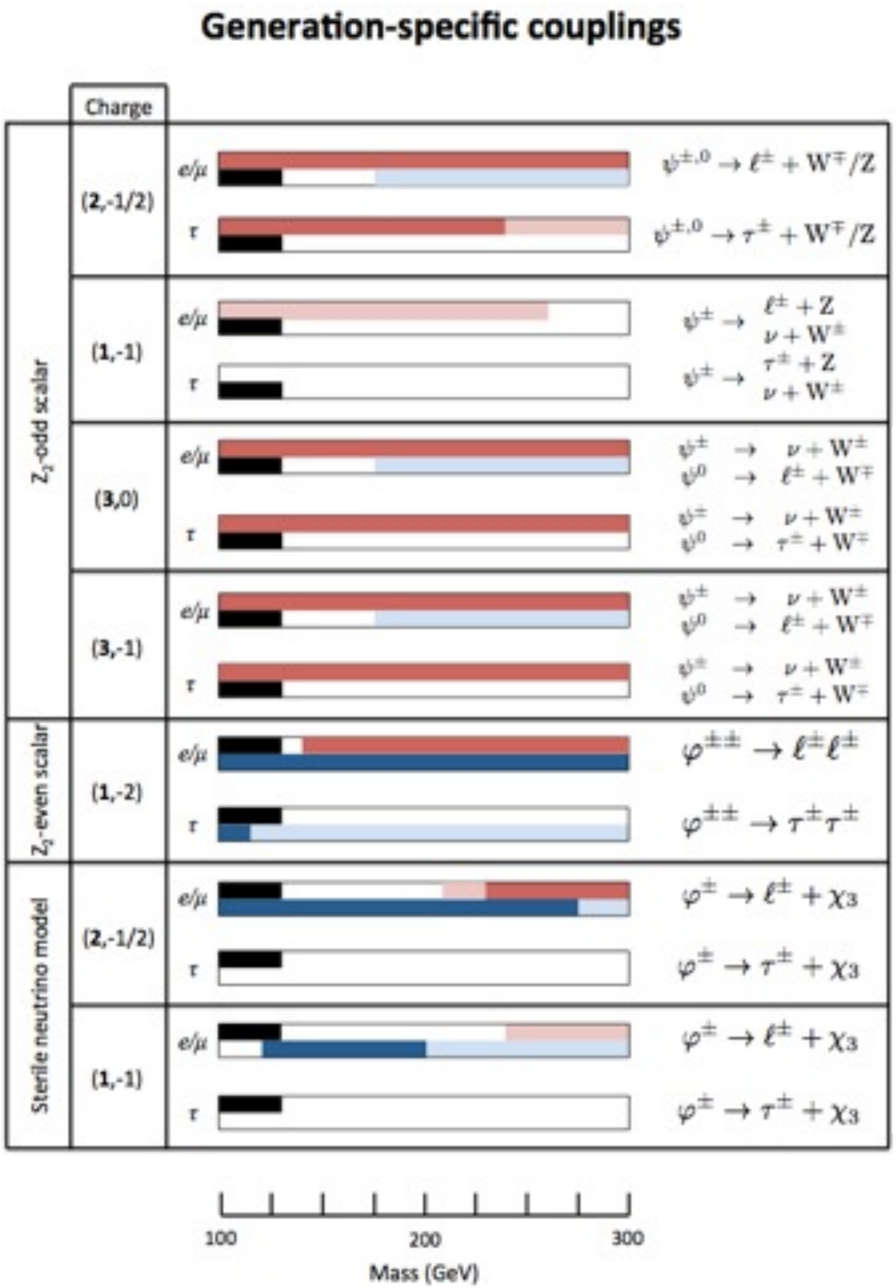




The production cross-section for the fermionic states is a lot larger.

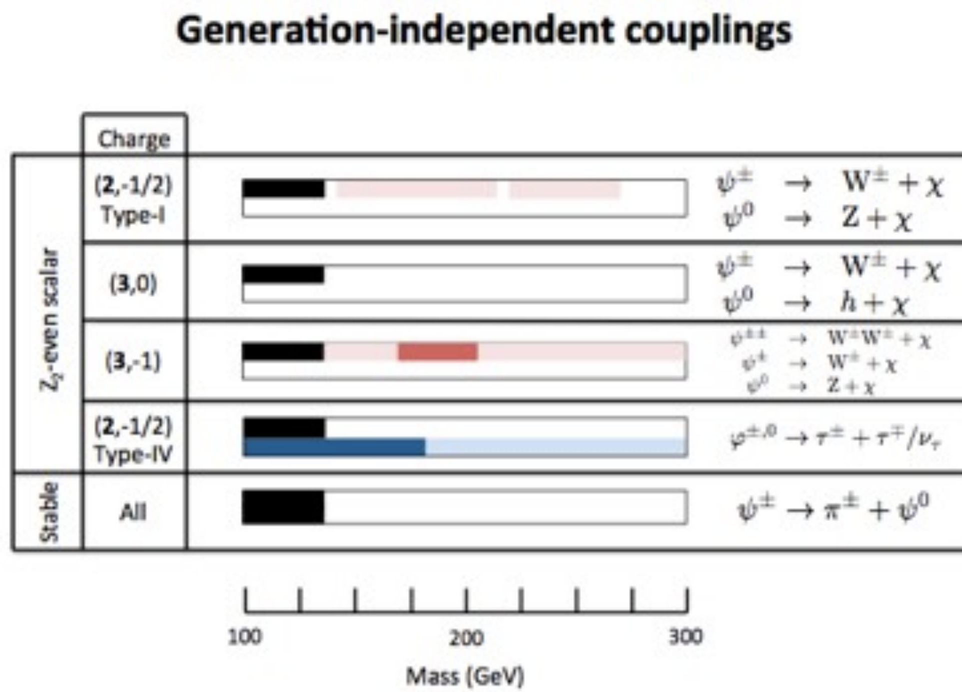
These cross-sections are now being probed at the LHC.

The discovery prospects really depend on how they decay.



The charged states should not remain stable. How to search for these states depends on their decays.

(Liu, Shuve, Weiner, IY arXiv:1303.4404)



So what is this all good for?

Concentrating on the leading operators, MiDM (dim-5) and RayDM (dim-7), motivates a variety of searches,

1) New electroweak states produced at the LHC with characteristic decays

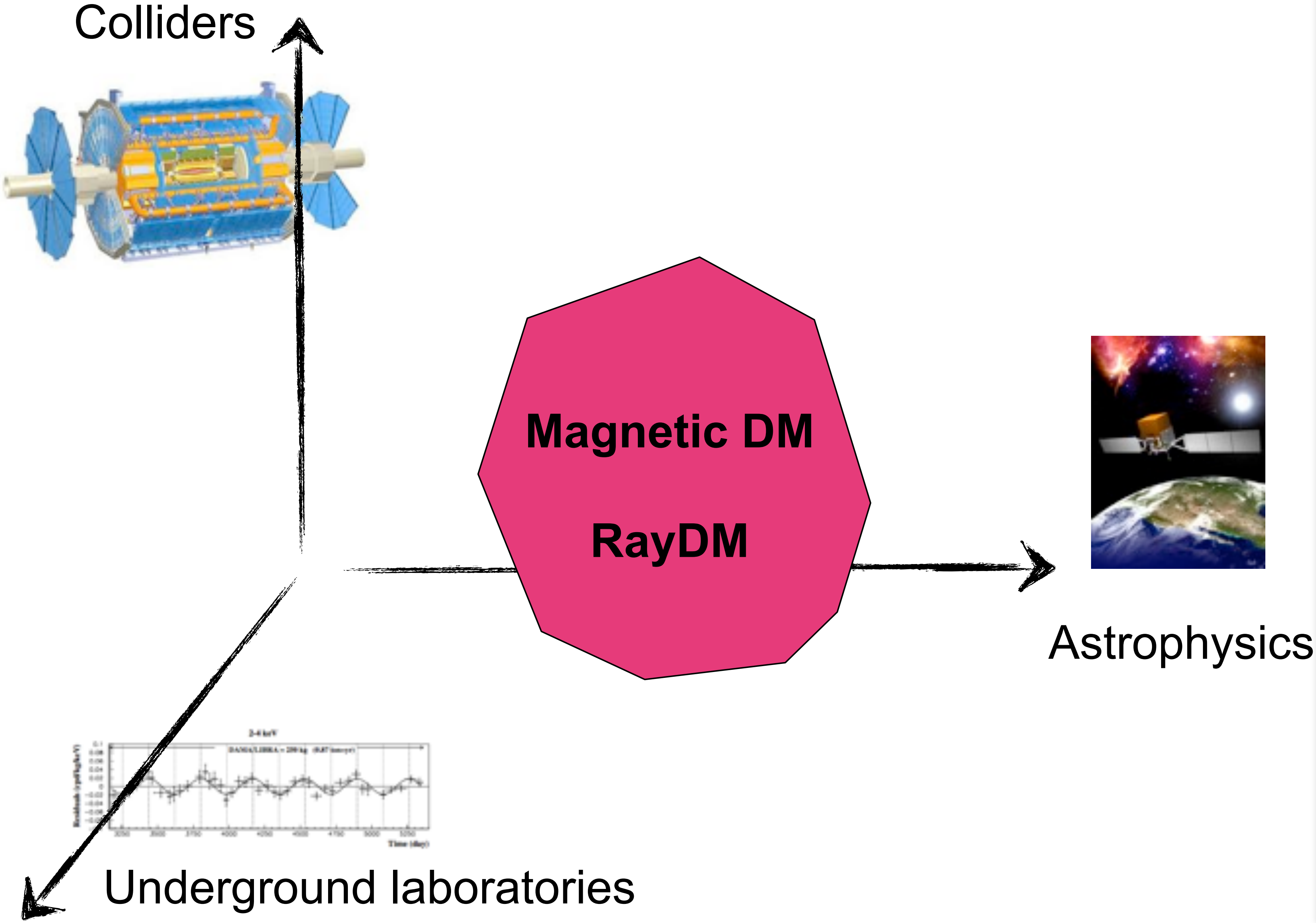
2) New types of searches for dark matter at colliders

3) Gamma-ray lines from galactic center

4) Gamma-rays from celestial objects

5) Can be seen in direct-detection experiments (possibly accompanied by a characteristic x-ray).

Conclusions



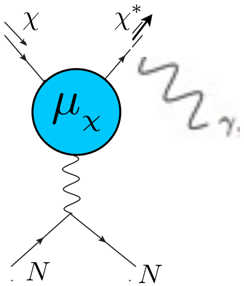
Danke schön

Thank You

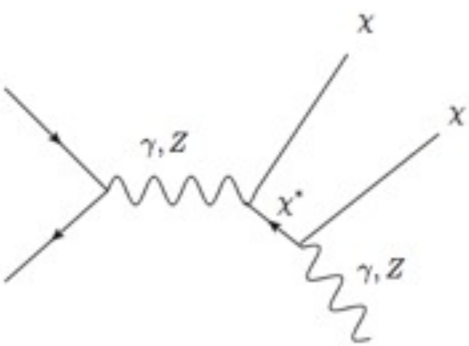
Other Searches

Is there any other testable phenomenology? More in the next section, but

MiDM (one-photon vertex)

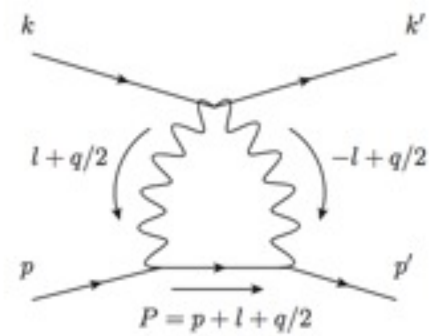


Excited state will decay. Look for x-ray lines.

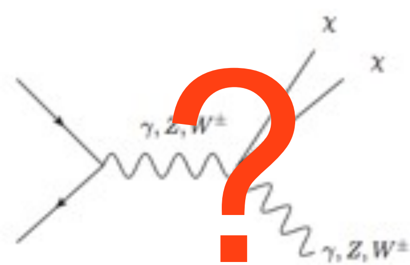


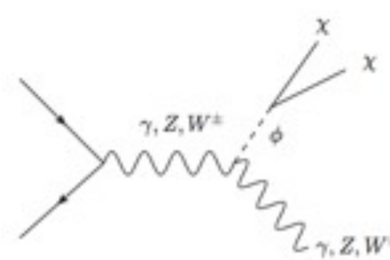
Production at LHC, monophoton from FSR.

RayDM (two-photon vertex)



Seems difficult, but forced us to find new ways (ala HQET) to calculate these effects.





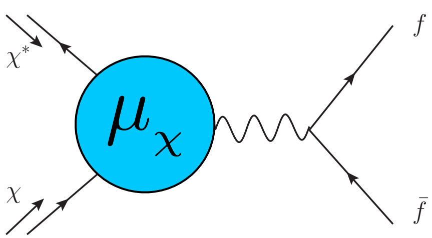
Motivates mono-W searches

Scaling Relations

For thermally produced MiDM, the direct and gamma-ray line indirect signatures are roughly independent of the size of the dipole

Duda, Gelmini, & Gondolo arXiv:hep-ph/0102200, Weiner and IY arXiv:1206.2910

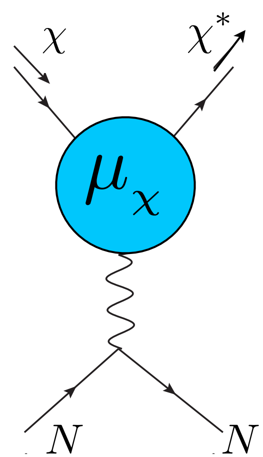
annihilation



$$\rho_{\text{MiDM}} = \rho_0 \times \frac{\mu^2_{\text{thermal}}}{\mu^2_\chi}$$

numerical value required for dark matter to be a thermal relic

direct detection

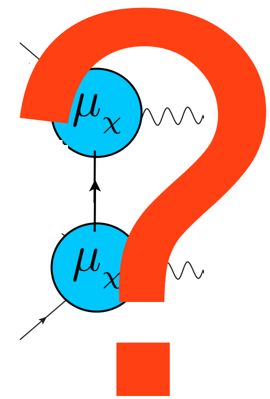


$$R_{DD} \propto n_\chi \mu^2_\chi =$$

$$\frac{\rho_0}{m_\chi}\frac{\mu^2_{\text{thermal}}}{\mu^2_\chi} \times \mu^2_\chi =$$

$$\frac{\rho_0}{m_\chi}\mu^2_{\text{thermal}}$$

gamma-rays



$$R_{\gamma\gamma} \propto n^2_\chi \mu^4_\chi =$$

$$\frac{\rho^2_0}{m^2_\chi}\frac{\mu^4_{\text{thermal}}}{\mu^4_\chi} \times \mu^4_\chi =$$

$$\frac{\rho^2_0}{m^2_\chi}\mu^4_{\text{thermal}}$$