

Exact results on the out-of-equilibrium Kondo model

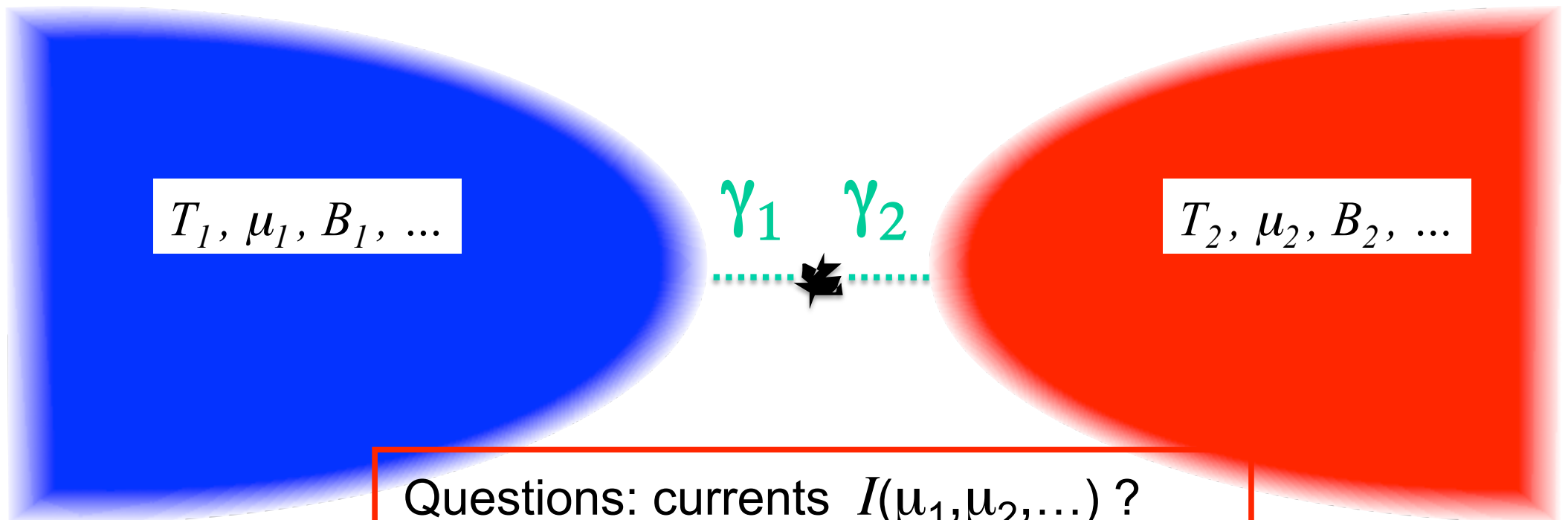
Edouard Boulat

Laboratoire MPQ, Université Paris 7 - Paris Diderot

with Loic Freton
see *PRL 112, 216802 (2014)*

Impurity out-of-equilibrium

- Several baths (macroscopic, at equilibrium)
- Out-of-equilibrium forcing
- Flow (of charge, spin, energy, ...) through impurity

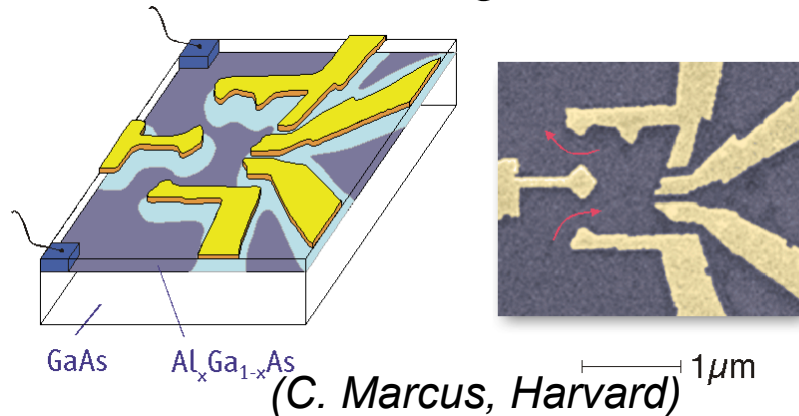


Questions: currents $I(\mu_1, \mu_2, \dots)$?
fluctuations $\Delta I(\mu_1, \mu_2, \dots)$?
state of the system ?

« Quantum impurities »

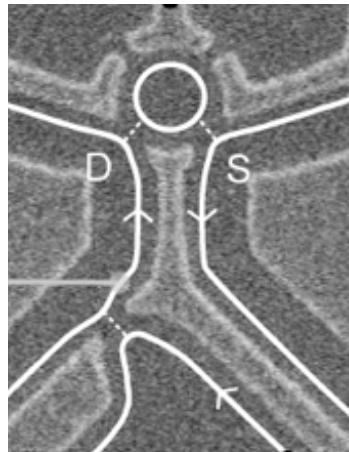
- Quantum dots:

- 2D electron gas



- Quantum Hall edge states

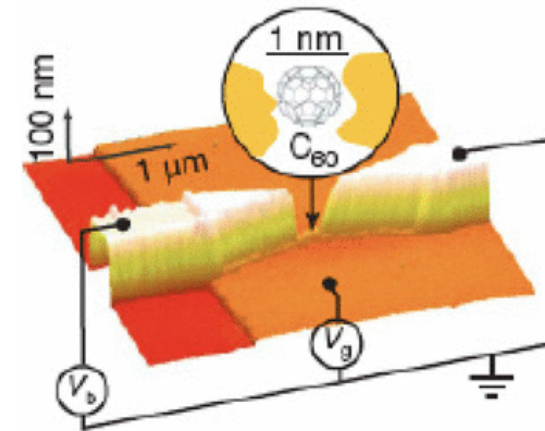
(F. Pierre, LPN)



- Molecules: metallic electrodes

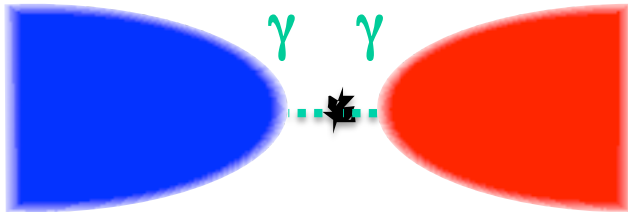
- break junctions

- electromigration



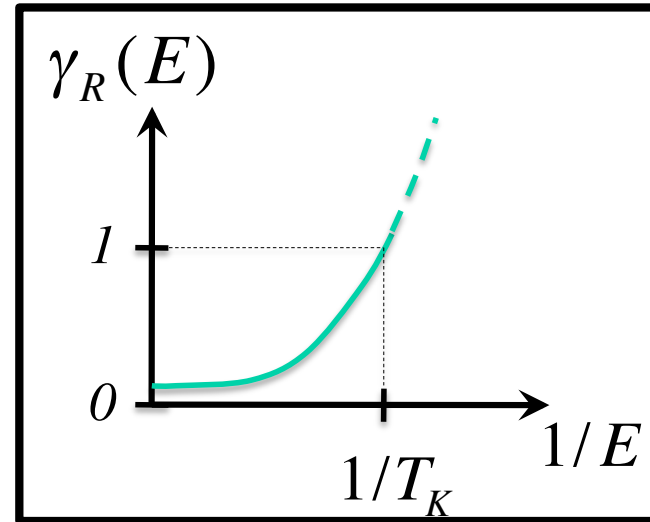
(W. Wernsdorfer, Institut Néel)

Weak - Strong Coupling

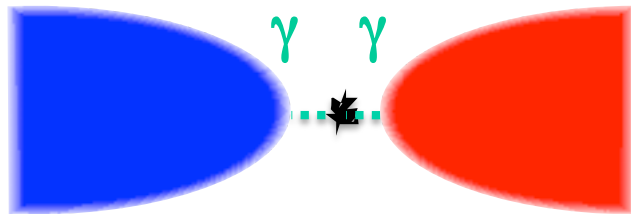


$$H = H_0 + \gamma H_K$$

baths impurity/bath coupling

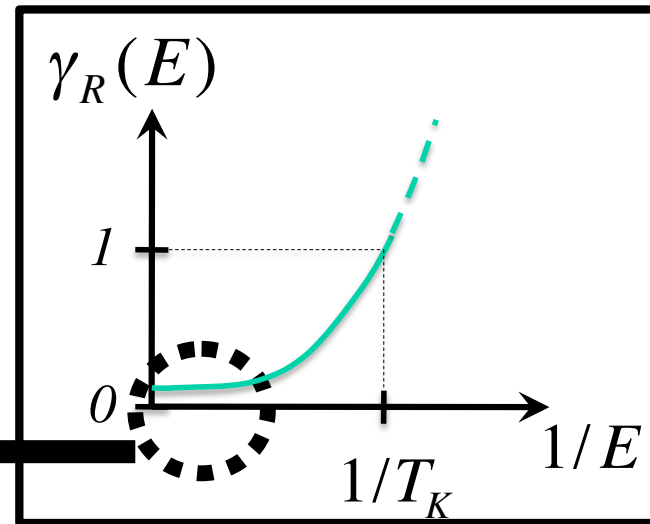


Weak - Strong Coupling



$$H = H_0 + \gamma H_K$$

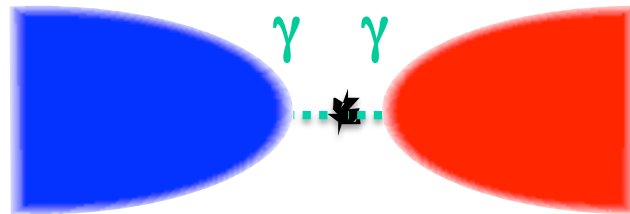
baths impurity/bath coupling



Weak coupling limit, $E \gg T_K$

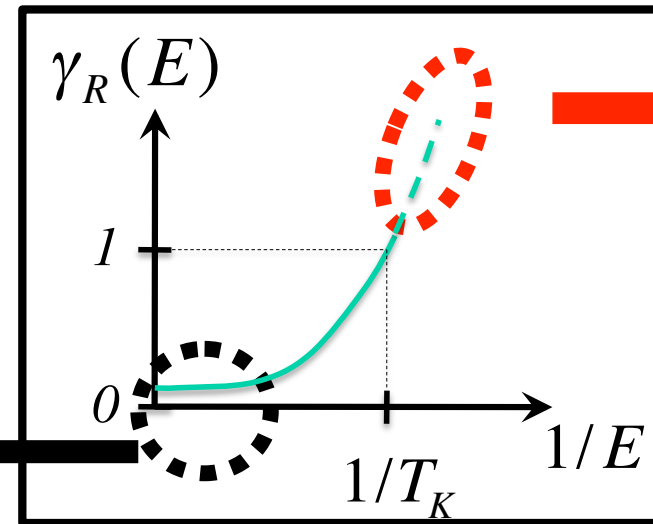


Weak - Strong Coupling



$$H = H_0 + \gamma H_K$$

baths impurity/bath coupling



Weak coupling limit, $E \gg T_K$

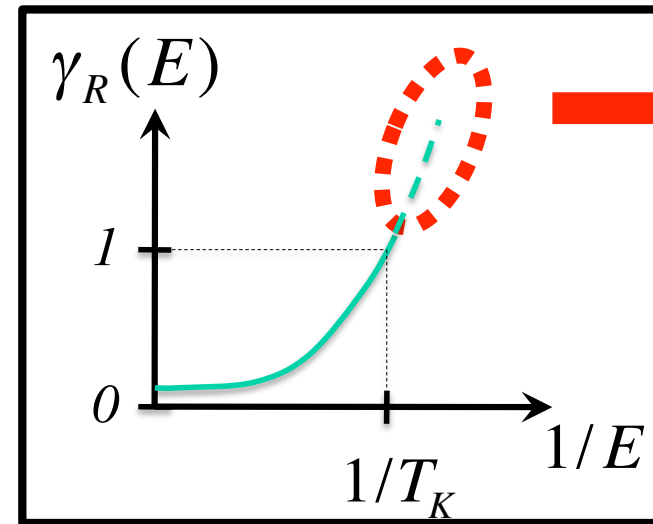
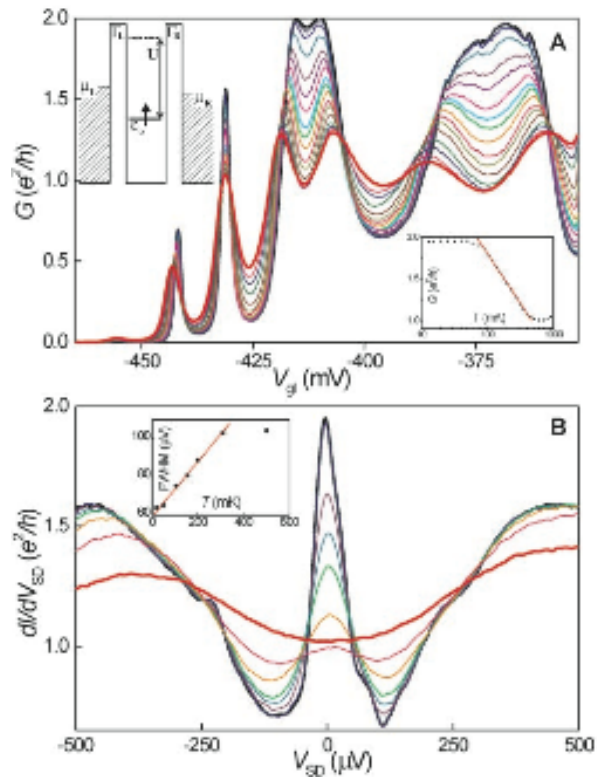


Strong coupling limit, $E \ll T_K$



At low energy :
strong coupling regime
“Physics is non perturbative”

Weak - Strong Coupling



Strong coupling limit, $E \ll T_K$

At low energy :
strong coupling regime
"Physics is non perturbative"

« Kondo resonance is a strong coupling phenomenon »

'Standard' perturbation theory

➤ Keldysh method:

- allows for a formal expression of the out-of-equilibrium density matrix

$$\hat{\rho}(t) = \mathcal{U}(0,t) \hat{\rho}(0) \mathcal{U}(0,t)^{-1} \quad \mathcal{U}(0,t) = \mathcal{P} e^{-i \gamma \int_0^t dt' H_B(t')}$$

- but how to evaluate/resum the perturbative expansion?

Fails in the strong coupling regime

➤ How to control approximate methods / other approaches?

- truncated EOM, diagrammatic methods,
- real-time RG, FRG,

➤ Integrability provides a non-perturbative approach

Non-perturbative approaches

a few available solutions !

- Dressed TBA
 - Quantum Hall edge states tunneling (P.Fendley, A.W.W.Ludwig, H.Saleur 1995)
 - Self-dual Interacting Resonant Level Model (E.B., P.Schmitteckert, H.Saleur 2008)
 - Map to equilibrium problem
 - Boundary sine Gordon model (V.Bazhanov, S.Lukyanov, A.B.Zamolodchikov 1999)
 - Effectively non-interacting systems (map to free fermions)
 - 1-ch Kondo (A. Schiller, U. Hershfield 1998)
 - Luttinger Liquid (A. Komnik, O. Gogolin 2003)
 - 2-ch Kondo (E. Sela, I. Affleck 2009) QCP & vicinity
- } Toulouse point
- Out-of-equilibrium forcing generically destroys integrable quasiparticles!
 - Dynamical forcing (AC...) ? Heat transport ?

The game is not over

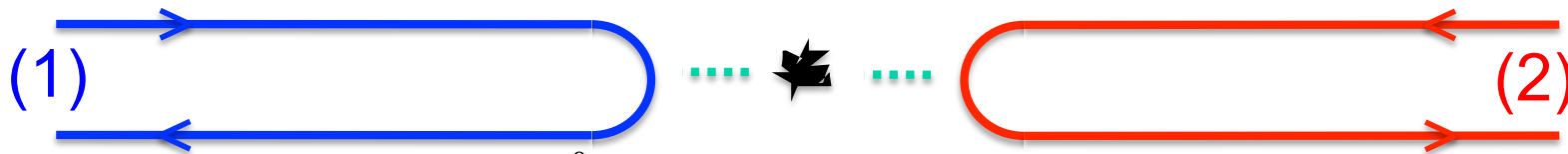
- Integrable theories have nevertheless a rich structure:
 - Infinite number of conserved quantities
 - Renormalization group flow is controlled non-perturbatively
- Can one use this rich structure to develop a controlled expansion out of equilibrium, in the strong coupling regime ?

 **Yes (at least in some cases)**

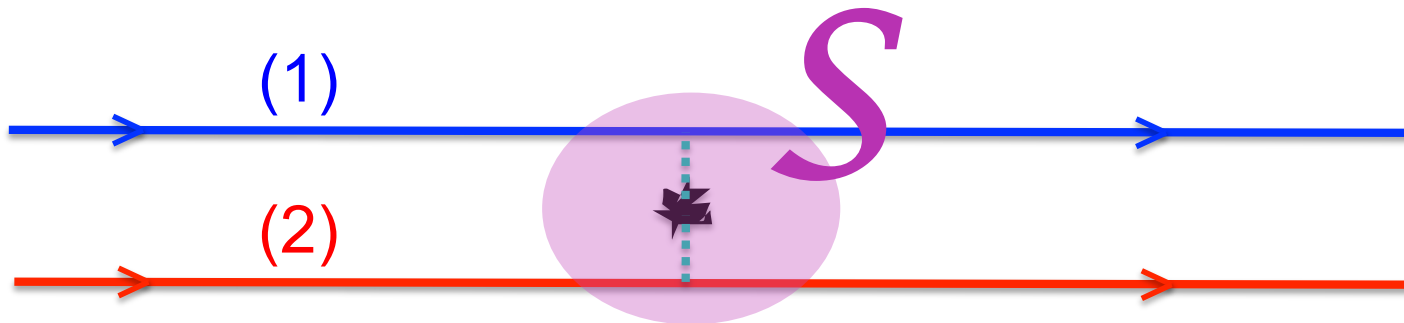
 Integrable Strong Coupling Expansion $V, T, \omega, \dots \leq T_K$

Modeling the baths

- Modes that couple to the impurity are 1D (conduction channel)
- Linearize the spectrum



$$H_0[\Psi] = \sum_{a=1,2} -iv_F \int_{-\infty}^0 dx \left[\Psi_{aR}^\dagger(x) \partial_x \Psi_{aR}(x) - \Psi_{aL}^\dagger(x) \partial_x \Psi_{aL}(x) \right]$$

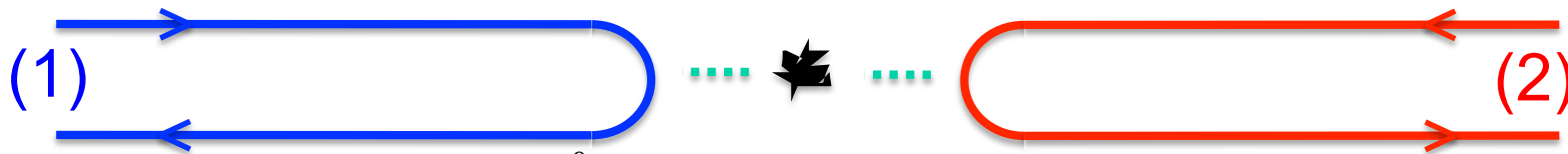


$$H_0[\Psi] = \sum_{a=1,2} -iv_F \int_{-\infty}^{\infty} dx \Psi_a^\dagger(x) \partial_x \Psi_a(x)$$

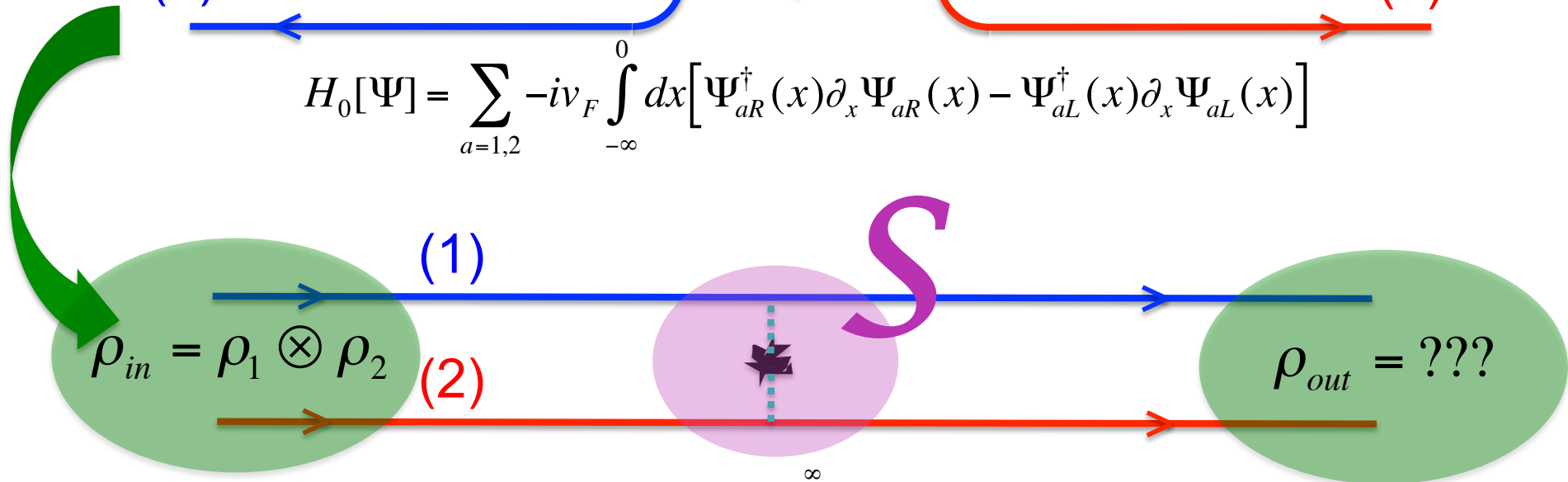
Chiral theory involving only right-moving fields: **scattering problem**

Modeling the baths

- Modes that couple to the impurity are 1D (conduction channel)
- Linearize the spectrum



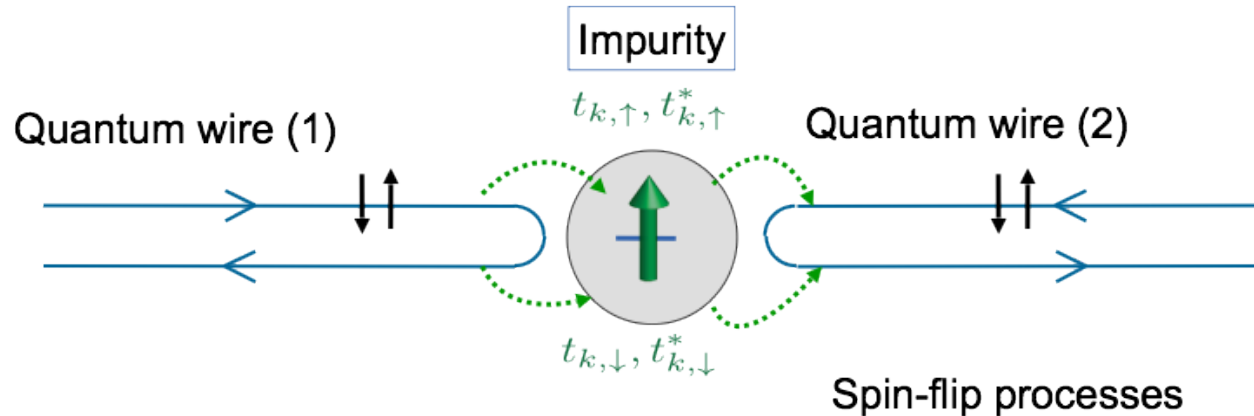
$$H_0[\Psi] = \sum_{a=1,2} -iv_F \int_{-\infty}^0 dx \left[\Psi_{aR}^\dagger(x) \partial_x \Psi_{aR}(x) - \Psi_{aL}^\dagger(x) \partial_x \Psi_{aL}(x) \right]$$



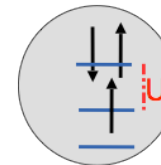
$$H_0[\Psi] = \sum_{a=1,2} -iv_F \int_{-\infty}^{\infty} dx \Psi_a^\dagger(x) \partial_x \Psi_a(x)$$

Chiral theory involving only right-moving fields: **scattering problem**

The $s=1/2$ Kondo model



Kondo stems from Anderson model



Parameters: - bare exchange coupling J
- anisotropy of couplings to the wires θ

Strategy

Want to describe the strong coupling regime $T, V, \omega \dots \leq T_K$

1. Incorporate out-of-equilibrium forcing AT the fixed point
→ yields a deformed CFT

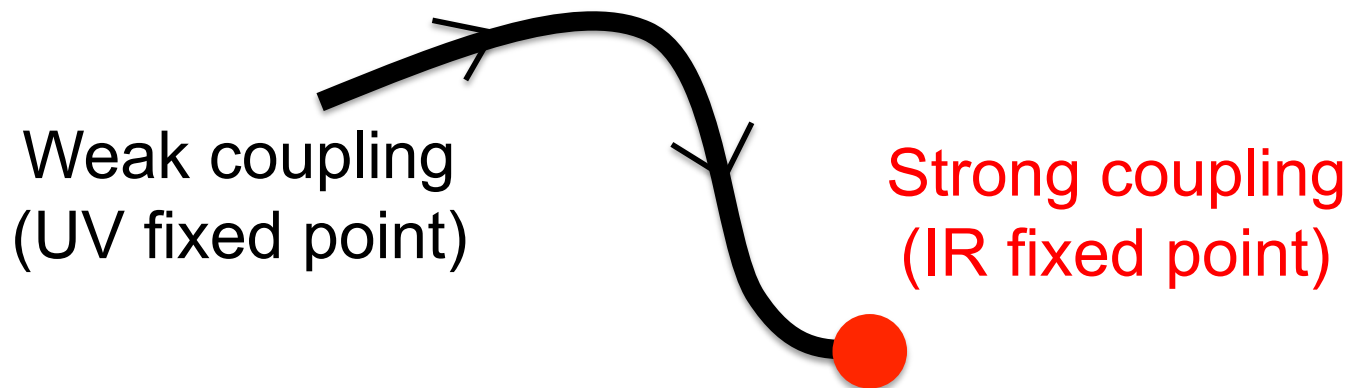
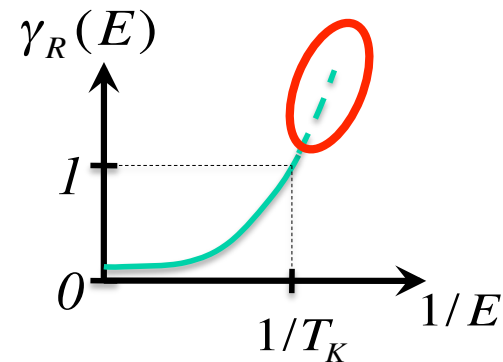
2. Use integrability to build the (many-body) S-matrix
→ incorporate (many-nody) back scattering

3. Expand in inverse powers of T_K

→ Net result: Taylor expansion of the universal scaling functions for local observables, at arbitrary order in principle

Strong coupling fixed point

- Perturbation is **relevant**
- Strong coupling fixed point described by BCFT



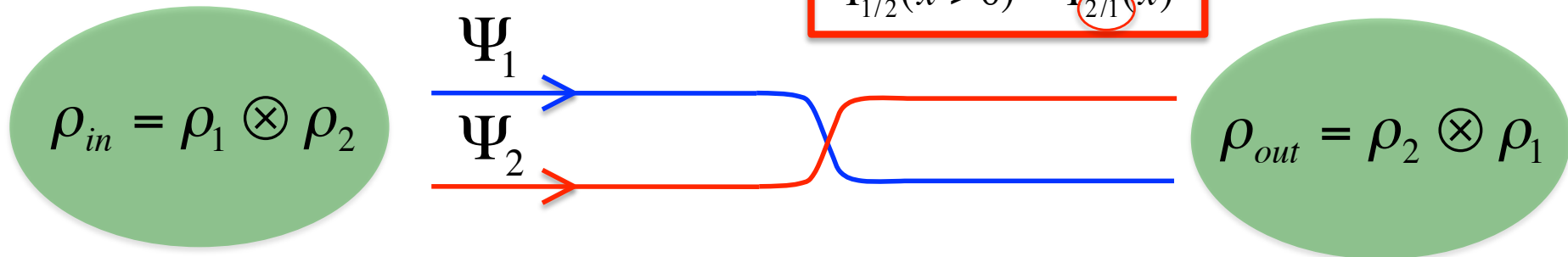
- Step 1: Out-of-equilibrium SC fixed point ($T_K = \infty$)

Strong coupling fixed point

- Boundary conditions: $\Phi(x = 0^-) = B \cdot \Phi(x = 0^+)$
- “Transparent fields”**: $\tilde{\Phi}(x < 0) = \Phi(x)$; $\tilde{\Phi}(x > 0) = B \cdot \Phi(x)$
They don't see the impurity!

BC for fermions: $\pi/2$ phase shift

$$\begin{aligned} \tilde{\Psi}_{1/2}(x < 0) &= \Psi_{1/2}(x) \\ \tilde{\Psi}_{1/2}(x > 0) &= \Psi_{2/1}(x) \end{aligned}$$



- Forcing out-of-equilibrium easily represented!
- Amounts to a gauge transformation $\mathcal{U}_{N.Equ}(z)$ for the transparent fields

$$\rho_{in} \propto e^{-\frac{H_0[\Psi_1] - \mu_1 Q_1}{T_1}} \otimes e^{-\frac{H_0[\Psi_2] - \mu_2 Q_2}{T_2}} \quad \longrightarrow \quad \langle I \rangle = \left(2e^2/h\right) (\mu_1(t) - \mu_2(t))$$

Recover the linear regime for the charge current

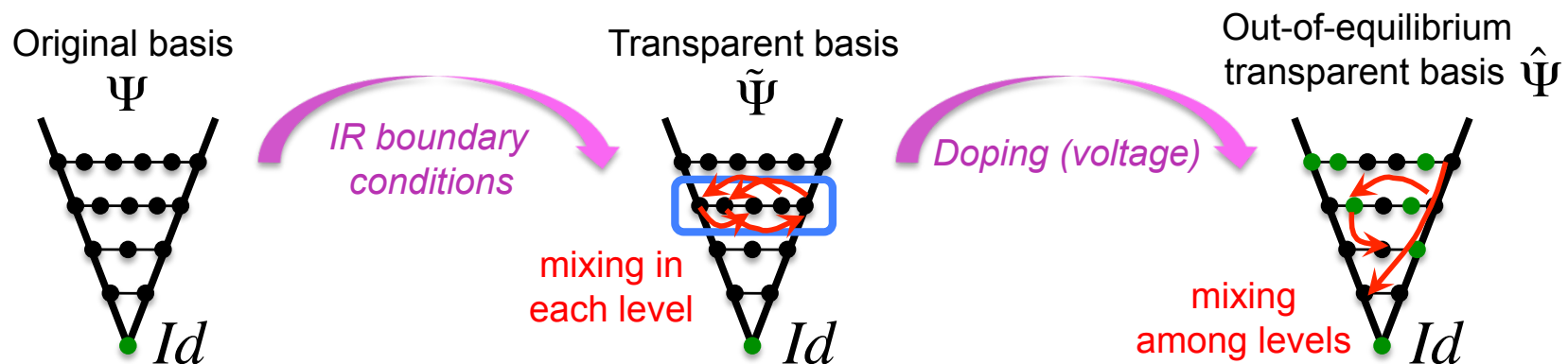
“Doping” a CFT

- The strong coupling fixed point has conformal symmetry ; transparent fields $\tilde{\Phi}$ are holomorphic (functions of $z = i(t-x)$)
- The forcing out of equilibrium can be absorbed by a gauge transformation (« doping ») $\hat{\Psi}(z) = \mathcal{U}_{N.Equ}(z) \cdot \tilde{\Psi}(z)$

$$\mathcal{U}_{N.Equ}(z) = e^{\int_z^{if dw} \Xi_a(w) \tilde{Q}_a(w)} \quad ; \quad \Xi_a(z = i(t-x)) = \int_0^{t-x} dt' \mu_a(t')$$

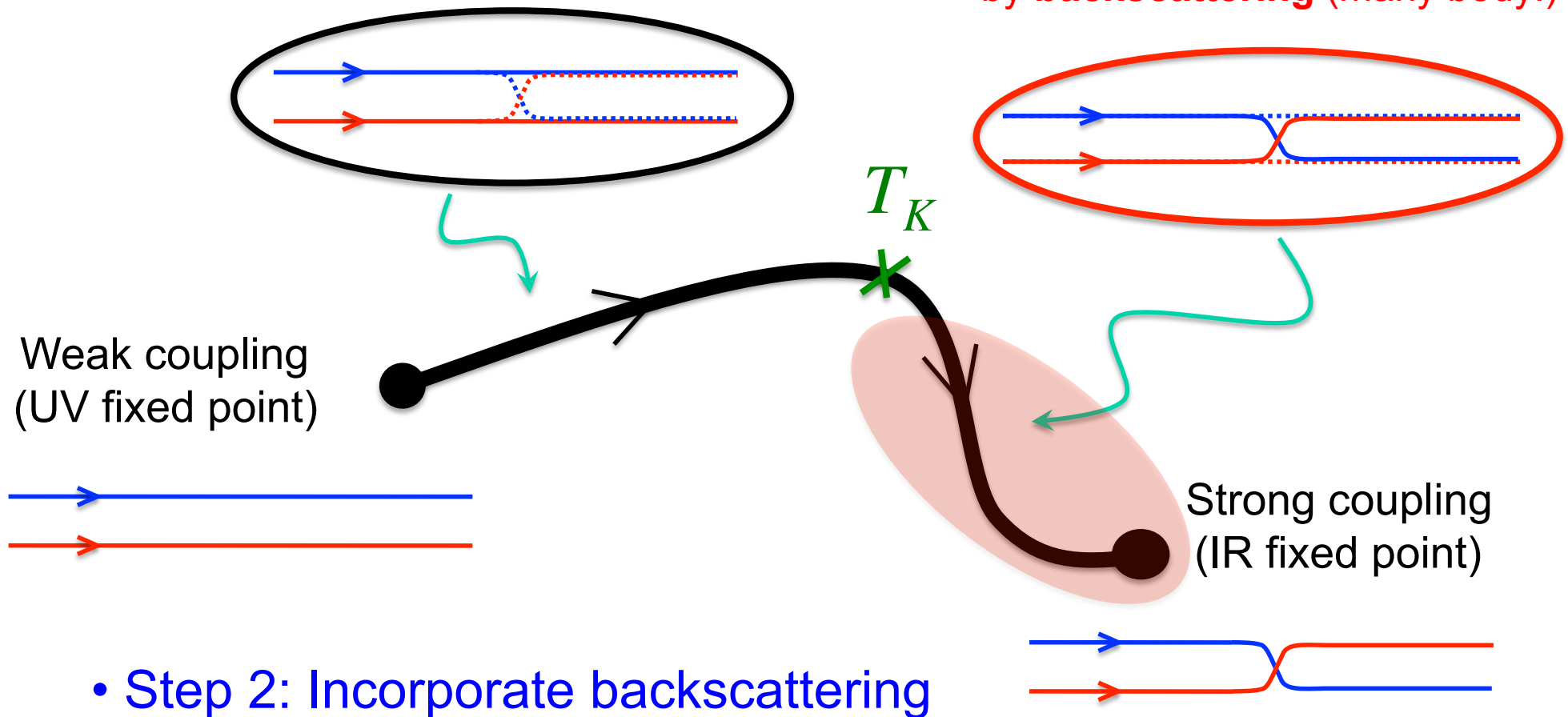
$$\langle A_1(x_1, t_1) A_2(x_2, t_2) \dots \rangle_{N.Equ} = \langle \hat{A}_1(x_1, t_1) \hat{A}_2(x_2, t_2) \dots \rangle_{Equ.} \quad ; \quad \hat{A} = \mathcal{U}_{N.Equ} \cdot A$$

- It's a deformation of the CFT (no geometrical interpretation unlike finite temperature CFT)

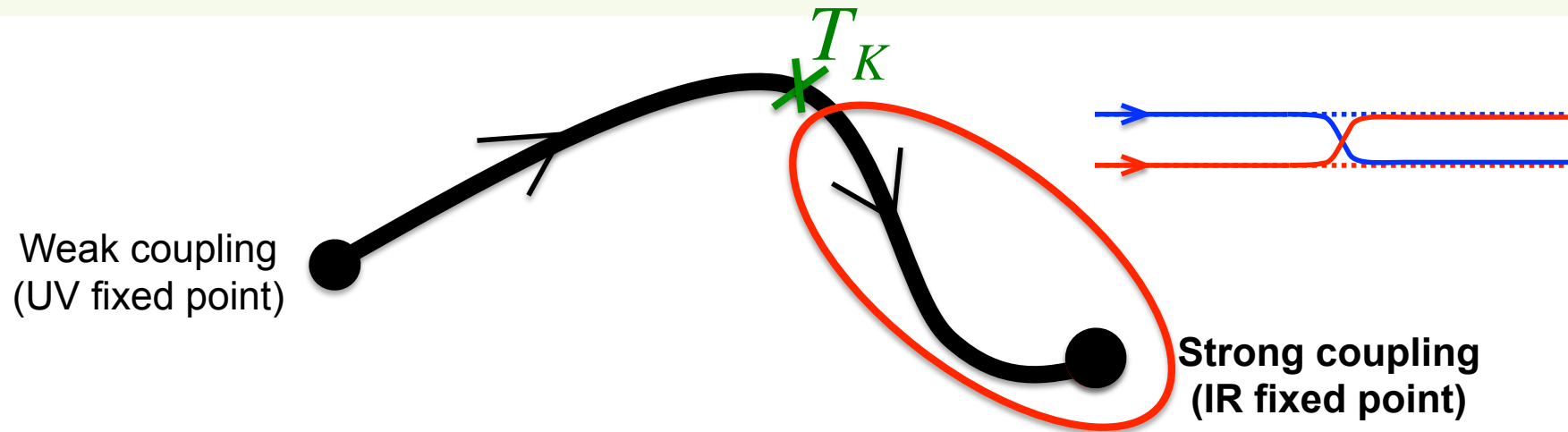


From weak to strong coupling

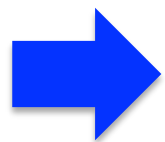
Physics is controlled
by **backscattering** (many body!)



Backscattering: dual theory



Integrability completely fixes the RG flow



The full approach to the IR fixed point can be described exactly by a **dual** theory. (F. Lesage, H.Saleur 1999)

$$H = H_0^{SC} + H_B^{SC} \quad H_B^{SC} = \sum_{n=1}^{\infty} \frac{g_{2n}}{(T_K)^{2n-1}} \hat{O}_{2n}(x=0)$$

Dual theory (2)

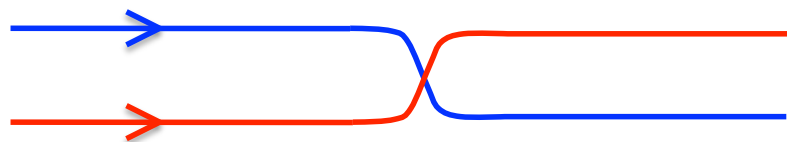
$$H = H_0^{SC} + H_B^{SC} \quad H_B^{SC} = \sum_{n=1}^{\infty} \frac{g_{2n}}{(T_K)^{2n-1}} \hat{O}_{2n}(x=0)$$

- The operators O_{2n} are the (infinitely many) conserved quantities stemming from integrability.
- The couplings g_n are pure numbers, **fixed** by integrability.
- Fermi liquid: the least irrelevant operator is $O_{2=T}$, an energy momentum tensor.
- Higher order processes have **integer** dimensions = 4,6,8,...

Backscattering transfers integer charges (electrons)
“SUPER FERMIL LIQUID”

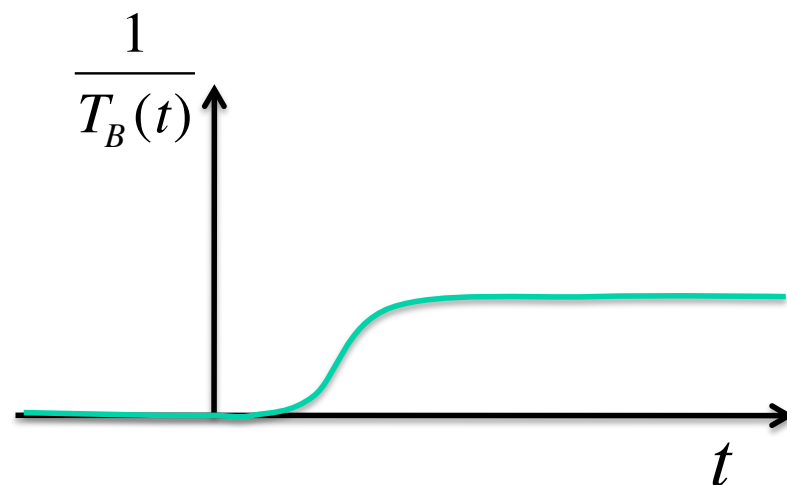
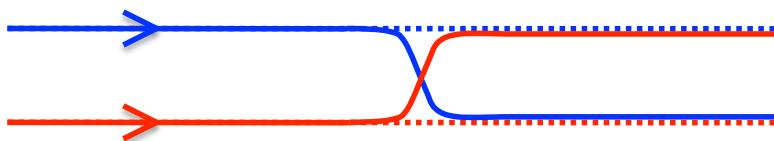
Keldysh expansion

- Start at time $t = -\infty$ at the SC fixed point ($T_K = \infty$)



$$\rho(-\infty) = \rho_{SC} = e^{-H_0^{SC} / k_B T}$$

- Switch on backscattering at time $t=0$

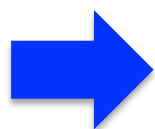
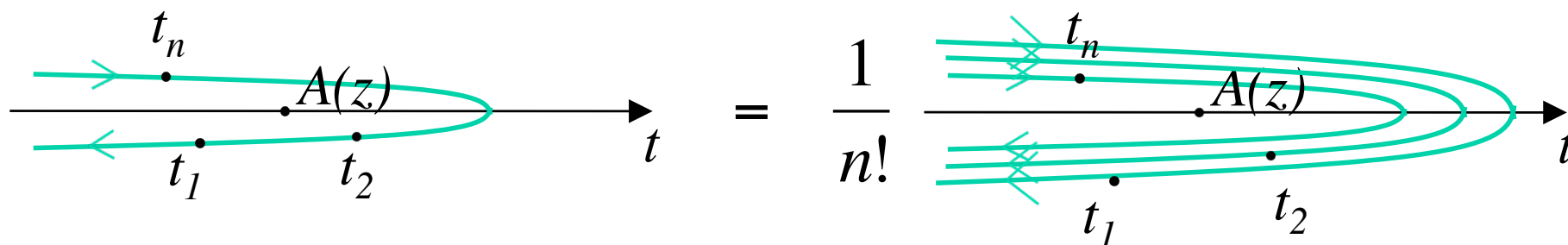


$$\rho(t) = U(t) \rho_{SC} U(t)^\dagger$$

$$U(t) = \mathcal{P}_K e^{-i \gamma \int_{-\infty}^t dt' H_B^{SC}(t')}$$

Effective operators

In a **super Fermi liquid**, the Keldysh expansion bears a simple form:



Each (local) operator can be replaced by an *effective* operator:

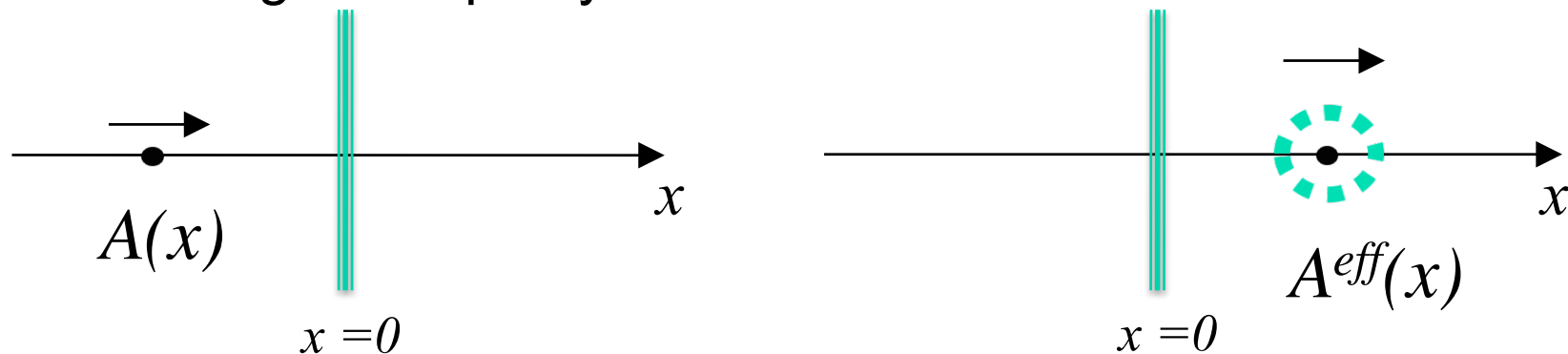
Complete many-body scattering

$$A^{eff}(z) = \mathcal{U}_{BS}(z) \cdot A(z)$$

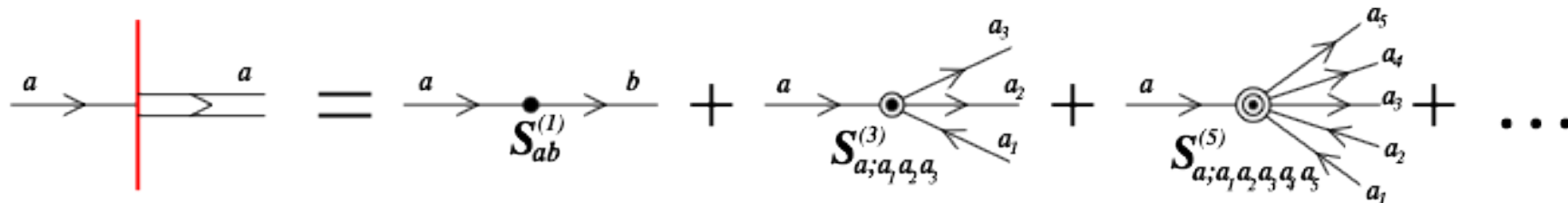
$$= e^{-i \int_z dt H_B(t)} \cdot A(z) = \sum_{n=0}^{\infty} \frac{(-i)^n}{n!} \oint_z dt_1 \dots \oint_z dt_n$$

Effective operators

- Operators undergo a **dressing by scattering** when crossing the impurity:



$$\langle A(x,t) \rangle_{N.Equ} = \langle \mathcal{U}_{N.Equ}(z) \cdot \mathcal{U}_{BS}(z) \cdot A(z) \rangle_0$$



Summary

A simple formula

$$\langle A(t) \rangle_{\text{n.eq.}} = \langle A_{\text{eff}} \rangle_0$$

can be expanded in powers of $1/T_K$: non-linear effects to arbitrary order !

Contours + CFT relations :

→ Algebraic (computer-friendly) reformulation



PT is *finite*: No UV divergence !

Exact formula ! **Directly gives universal results**

Versatile formulation $V, \mu, \omega, T_i, t, B, \epsilon_d \dots$

Integrable models with integer conserved quantities only



Not always easy to find conserved quantities and couplings

The price to pay...

Electrical current operator : **initial fermions**

$$I_{BS} = 1/2((\psi_{1\uparrow}^\dagger \psi_{1\uparrow}) - (\psi_{2\uparrow}^\dagger \psi_{2\uparrow}) + (\psi_{1\downarrow}^\dagger \psi_{1\downarrow}) - (\psi_{2\downarrow}^\dagger \psi_{2\downarrow}))$$

The price to pay...

Electrical current operator : **Transparent fermions**

Zeroth order

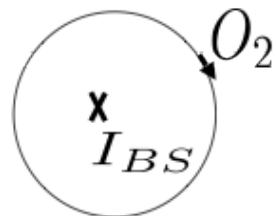
\times
 I_{BS}

$$\begin{aligned}
 & -\frac{1}{2} \sin(\theta) \cos(\theta) (\cos(\xi) - 1) + \frac{1}{2} i \sin(\theta) \sin(\xi) ((\Psi_{1\uparrow})^\dagger (\Psi_{2\uparrow})) \\
 & + \frac{1}{2} \sin(\theta) \cos(\theta) (\cos(\xi) - 1) + \frac{1}{2} i \sin(\theta) \sin(\xi) ((\Psi_{1\uparrow}) (\Psi_{2\uparrow})^\dagger) \\
 & \quad + \frac{1}{2} (\sin^2(\theta) \cos(\xi) + \cos^2(\theta)) ((\Psi_{1\uparrow})^\dagger (\Psi_{1\uparrow})) \\
 & \quad + \frac{1}{2} (-\sin^2(\theta) \cos(\xi) - \cos^2(\theta)) ((\Psi_{2\uparrow})^\dagger (\Psi_{2\uparrow})) \\
 & + -\frac{1}{2} \sin(\theta) \cos(\theta) (\cos(\xi) - 1) - \frac{1}{2} i \sin(\theta) \sin(\xi) ((\Psi_{1\downarrow})^\dagger (\Psi_{2\downarrow})) \\
 & + \frac{1}{2} \sin(\theta) \cos(\theta) (\cos(\xi) - 1) - \frac{1}{2} i \sin(\theta) \sin(\xi) ((\Psi_{1\downarrow}) (\Psi_{2\downarrow})^\dagger) \\
 & \quad + \frac{1}{2} (\sin^2(\theta) \cos(\xi) + \cos^2(\theta)) ((\Psi_{1\downarrow})^\dagger (\Psi_{1\downarrow})) \\
 & \quad + \frac{1}{2} (-\sin^2(\theta) \cos(\xi) - \cos^2(\theta)) ((\Psi_{2\downarrow})^\dagger (\Psi_{2\downarrow}))
 \end{aligned}$$

The price to pay...

Electrical current operator : Transparent fermions

First order: $1/T_K$



$$\begin{aligned}
 & -(\sin \theta * (2 * \cos[\theta/2]^2 * (\cos \xi - i * \sin \xi) + 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi - i * \sin \xi) + \\
 & +(\sin \theta * (-2 * \cos[\theta/2]^2 * (\cos \xi + i * \sin \xi) + 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi + i * \sin \xi) + \sin \theta^2 * (\cos \xi - i * \sin \xi) \\
 & + - (\sin \theta * (-2 * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi) + \\
 & +(\sin \theta * (2 * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) + \sin \theta^2 * (\cos \xi - i * \sin \xi) \\
 & + - ((\cos[\theta/2]^2 - \cos \theta - \sin[\theta/2]^2) * \sin \theta^2 * (\cos \xi - i * \sin \xi) - \sin[\theta/2]^2 * \sin \theta^2 * (\cos \xi + i * \sin \xi) \\
 & + - (\sin \theta^2 * (-\cos \xi + 2 * \cos[\theta/2]^2 * (\cos \xi - i * \sin \xi) - \cos \theta * (\cos \xi - i * \sin \xi) \\
 & +2 * \cos[\theta/2]^2 * (\cos \xi + i * \sin \xi) - \cos \theta * (\cos \xi + i * \sin \xi) \\
 & +(i/16) * \sin \theta^2 * (i * (\cos \xi - i * \sin \xi) + i * \cos \theta * (\cos \xi - i * \sin \xi) + (2 * i) * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi)) - (i/16) * \sin \theta^2 * ((-i) * (\cos \xi + i * \sin \xi) + i * \cos \theta * (\cos \xi + i * \sin \xi) + (2 * i) * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi)) \\
 & +(\sin \theta * (2 * \cos[\theta/2]^2 * (\cos \xi - i * \sin \xi) + 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi - i * \sin \xi)))/16 - (\sin \theta * (2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi + i * \sin \xi)))/16 + (\sin \theta^3 * (\cos \xi - i * \sin \xi))/16 - (\sin \theta^3 * (\cos \xi + i * \sin \xi))/16 \\
 & + - (\sin \theta * (-2 * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi))/16 + (\sin \theta * (-2 * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi))/16 - (\sin \theta^3 * (\cos \xi - i * \sin \xi))/16 + (\sin \theta^3 * (\cos \xi + i * \sin \xi))/16 \\
 & +(\sin \theta^2 * (\cos \xi - \cos \theta * (\cos \xi - i * \sin \xi) - i * \sin \xi))/16 + (\sin \theta^2 * (\cos \xi + \cos \theta * (\cos \xi + i * \sin \xi) - i * \sin \xi))/16 - (\sin \theta^2 * (-\cos \xi - \cos \theta * (\cos \xi + i * \sin \xi) - i * \sin \xi))/16 - (\sin \theta^2 * (-\cos \xi + \cos \theta * (\cos \xi + i * \sin \xi) - i * \sin \xi))/16 \\
 & +(\sin \theta^2 * (\cos \xi + \cos \theta * (\cos \xi - i * \sin \xi) - i * \sin \xi))/16 - (\sin \theta^2 * (-\cos \xi + \cos \theta * (\cos \xi - i * \sin \xi) - i * \sin \xi))/16 - (\sin \theta^2 * (-\cos \xi - \cos \theta * (\cos \xi - i * \sin \xi) + i * \sin \xi))/16 - (\sin \theta^2 * (-\cos \xi + \cos \theta * (\cos \xi - i * \sin \xi) + i * \sin \xi))/16 \\
 & + - (\sin \theta * (-2 * \cos[\theta/2]^2 * (\cos \xi - i * \sin \xi) + 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi - i * \sin \xi) + \sin \theta^2 * (\cos \xi - i * \sin \xi)))/16 + (\sin \theta * (2 * \cos[\theta/2]^2 * (\cos \xi + i * \sin \xi) + 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi + i * \sin \xi) + \sin \theta^2 * (\cos \xi - i * \sin \xi)))/16 \\
 & + - (\sin \theta * (2 * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi) + \sin \theta^2 * (\cos \xi - i * \sin \xi)))/16 + (\sin \theta * (-2 * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) + \sin \theta^2 * (\cos \xi - i * \sin \xi)))/16 \\
 & +(\sin \theta^2 * (\cos \xi + 2 * \cos[\theta/2]^2 * (\cos \xi - i * \sin \xi) - \cos \theta * (\cos \xi - i * \sin \xi) - i * \sin \xi))/16 - (\sin \theta^2 * (-\cos \xi + 2 * \cos[\theta/2]^2 * (\cos \xi + i * \sin \xi) - \cos \theta * (\cos \xi + i * \sin \xi) - i * \sin \xi))/16 \\
 & +(\sin \theta^2 * (-\cos \xi + \cos \theta * (\cos \xi - i * \sin \xi) + 2 * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi) + i * \sin \xi))/16 - (\sin \theta^2 * (\cos \xi + \cos \theta * (\cos \xi + i * \sin \xi) + 2 * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) + i * \sin \xi))/16 \\
 & + - (\sin \theta * (-2 * \cos[\theta/2]^2 * (\cos \xi - i * \sin \xi) + 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi - i * \sin \xi)))/16 + (\sin \theta * (2 * \cos[\theta/2]^2 * (\cos \xi + i * \sin \xi) + 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi + i * \sin \xi)))/16 - (\sin \theta^3 * (\cos \xi - i * \sin \xi))/16 + (\sin \theta^3 * (\cos \xi + i * \sin \xi))/16 \\
 & +(\sin \theta * (2 * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi)))/16 - (\sin \theta * (-2 * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi)))/16 + (\sin \theta^3 * (\cos \xi - i * \sin \xi))/16 - (\sin \theta^3 * (\cos \xi + i * \sin \xi))/16
 \end{aligned}$$

The price to pay...

Electrical current operator : Transparent fermions

$$\begin{aligned} & * \sin \xi) - i * \sin \xi) / 16 + (\sin \theta^2 * (-\cos \xi + \cos \theta * (\cos \xi - i * \sin \xi) + i * \sin \xi)) / 16 - (\sin \theta^2 * (\cos \xi + \cos \theta * (\cos \xi + i * \sin \xi) + i * \sin \xi)) / 16 * (\psi_{1\uparrow}(\psi_{2\uparrow}^\dagger(\psi_{1\downarrow}^\dagger\psi_{2\downarrow}))) \\ & - i * \sin \xi) + i * \sin \xi) / 16 - (\sin \theta^2 * (\cos \xi - \cos \theta * (\cos \xi + i * \sin \xi) + i * \sin \xi)) / 16 - (\sin \theta^2 * (\cos \xi + \cos \theta * (\cos \xi + i * \sin \xi) + i * \sin \xi)) / 16 * (\psi_{1\uparrow}(\psi_{2\uparrow}^\dagger(\psi_{1\downarrow}^\dagger\psi_{2\downarrow}))) \\ & - i * \sin \xi) / 8 + (\cos[\theta/2]^2 * \sin \theta^2 * (\cos \xi - i * \sin \xi)) / 8 + (\cos[\theta/2]^2 * \sin \theta^2 * (-\cos \xi + i * \sin \xi)) / 8 - (\cos[\theta/2]^2 * \sin \theta^2 * (\cos \xi + i * \sin \xi)) / 8 * (\psi_{1\uparrow}^\dagger(\psi_{1\uparrow}(\psi_{1\downarrow}^\dagger\psi_{1\downarrow}))) \\ & \xi - i * \sin \xi) / 8 - (\sin[\theta/2]^2 * \sin \theta^2 * (\cos \xi - i * \sin \xi)) / 8 + (\cos[\theta/2]^2 * \sin \theta^2 * (\cos \xi + i * \sin \xi)) / 8 + (\sin[\theta/2]^2 * \sin \theta^2 * (\cos \xi + i * \sin \xi)) / 8 * (\psi_{1\uparrow}^\dagger(\psi_{1\uparrow}(\psi_{2\downarrow}^\dagger\psi_{2\downarrow}))) \\ & 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi + i * \sin \xi)) / 16 - (\sin \theta^3 * (\cos \xi - i * \sin \xi)) / 16 + (\sin \theta^3 * (\cos \xi + i * \sin \xi)) / 16 * (\psi_{1\uparrow}^\dagger(\psi_{1\uparrow}(\psi_{1\downarrow}^\dagger\psi_{2\downarrow}))) \\ & 2 * \cos[\theta/2]^2 * (\cos \xi + i * \sin \xi) + 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi + i * \sin \xi)) / 16 + (\sin \theta^3 * (\cos \xi - i * \sin \xi)) / 16 - (\sin \theta^3 * (\cos \xi + i * \sin \xi)) / 16 * (\psi_{1\uparrow}^\dagger(\psi_{1\uparrow}(\psi_{1\downarrow}^\dagger\psi_{2\downarrow}))) \\ & \xi - i * \sin \xi) / 8 + (\sin[\theta/2]^2 * \sin \theta^2 * (\cos \xi - i * \sin \xi)) / 8 - (\cos[\theta/2]^2 * \sin \theta^2 * (\cos \xi + i * \sin \xi)) / 8 - (\sin[\theta/2]^2 * \sin \theta^2 * (\cos \xi + i * \sin \xi)) / 8 * (\psi_{2\uparrow}^\dagger(\psi_{2\uparrow}(\psi_{1\downarrow}^\dagger\psi_{1\downarrow}))) \\ & 2]^2 * \sin \theta^2 * (-\cos \xi - i * \sin \xi)) / 8 - (\sin[\theta/2]^2 * \sin \theta^2 * (\cos \xi - i * \sin \xi)) / 8 - (\sin[\theta/2]^2 * \sin \theta^2 * (-\cos \xi + i * \sin \xi)) / 8 + (\sin[\theta/2]^2 * \sin \theta^2 * (\cos \xi + i * \sin \xi)) / 8 * (\psi_{2\uparrow}^\dagger(\psi_{2\uparrow}(\psi_{2\downarrow}^\dagger\psi_{2\downarrow}))) \\ & n \xi)) / 16 - (\sin \theta * (-2 * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi))) / 16 + (\sin \theta^3 * (\cos \xi - i * \sin \xi)) / 16 - (\sin \theta^3 * (\cos \xi + i * \sin \xi)) / 16 * (\psi_{2\uparrow}^\dagger(\psi_{2\uparrow}(\psi_{1\downarrow}^\dagger\psi_{2\downarrow}))) \\ & \sin \xi)) / 16 + (\sin \theta * (2 * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi))) / 16 - (\sin \theta^3 * (\cos \xi - i * \sin \xi)) / 16 + (\sin \theta^3 * (\cos \xi + i * \sin \xi)) / 16 * (\psi_{2\uparrow}^\dagger(\psi_{2\uparrow}(\psi_{1\downarrow}^\dagger\psi_{2\downarrow}))) \\ & * \cos \theta * (\cos \xi - i * \sin \xi) + \sin \theta^2 * (\cos \xi - i * \sin \xi)) / 16 - (\sin \theta * (2 * \cos[\theta/2]^2 * (\cos \xi + i * \sin \xi) + 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi + i * \sin \xi) + \sin \theta^2 * (\cos \xi + i * \sin \xi))) / 16 * (d\psi_{1\downarrow}^\dagger\psi_{2\downarrow}) \\ & \theta/2]^2 * (\cos \xi - i * \sin \xi) + \sin \theta^2 * (\cos \xi - i * \sin \xi)) / 16 - (\sin \theta * (-2 * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) + \sin \theta^2 * (\cos \xi + i * \sin \xi))) / 16 * (\psi_{1\downarrow}^\dagger d\psi_{2\downarrow}) \\ & + ((\cos[\theta/2]^2 - \cos \theta - \sin[\theta/2]^2) * \sin \theta^2 * (\cos \xi - i * \sin \xi)) / 8 - ((\cos[\theta/2]^2 - \cos \theta - \sin[\theta/2]^2) * \sin \theta^2 * (\cos \xi + i * \sin \xi)) / 8 * (\psi_{1\downarrow}^\dagger(\psi_{1\downarrow}(\psi_{2\downarrow}^\dagger\psi_{2\downarrow}))) \\ & \xi + 2 * \cos[\theta/2]^2 * (\cos \xi - i * \sin \xi) - \cos \theta * (\cos \xi - i * \sin \xi) - i * \sin \xi)) / 16 - (\sin \theta^2 * (-\cos \xi + 2 * \cos[\theta/2]^2 * (\cos \xi + i * \sin \xi) - \cos \theta * (\cos \xi + i * \sin \xi) - i * \sin \xi)) / 16 * (\psi_{1\downarrow}^\dagger d\psi_{1\downarrow}) \\ & \cos \xi + \cos \theta * (\cos \xi - i * \sin \xi) + 2 * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi) + i * \sin \xi)) / 16 - (\sin \theta^2 * (\cos \xi + \cos \theta * (\cos \xi + i * \sin \xi) + 2 * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) + i * \sin \xi)) / 16 * (d\psi_{2\downarrow}^\dagger\psi_{2\downarrow}) \\ & \cos \theta * (\cos \xi - i * \sin \xi) + \sin \theta^2 * (\cos \xi - i * \sin \xi)) / 16 - (\sin \theta * (-2 * \cos[\theta/2]^2 * (\cos \xi + i * \sin \xi) + 2 * \cos[\theta/2]^2 * \cos \theta * (\cos \xi + i * \sin \xi) + \sin \theta^2 * (\cos \xi + i * \sin \xi))) / 16 * (d\psi_{1\downarrow}\psi_{2\downarrow}^\dagger) \\ & [\theta/2]^2 * (\cos \xi - i * \sin \xi) + \sin \theta^2 * (\cos \xi - i * \sin \xi)) / 16 - (\sin \theta * (2 * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) - 2 * \cos \theta * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi) + \sin \theta^2 * (\cos \xi + i * \sin \xi))) / 16 * (\psi_{1\downarrow} d\psi_{2\downarrow}^\dagger) \\ & \cos \xi + 2 * \cos[\theta/2]^2 * (\cos \xi - i * \sin \xi) - \cos \theta * (\cos \xi - i * \sin \xi) + i * \sin \xi)) / 16 + (\sin \theta^2 * (\cos \xi + 2 * \cos[\theta/2]^2 * (\cos \xi + i * \sin \xi) - \cos \theta * (\cos \xi + i * \sin \xi) + i * \sin \xi)) / 16 * (d\psi_{1\downarrow}^\dagger\psi_{1\downarrow}) \\ & * \cos \theta * (\cos \xi - i * \sin \xi) - (2 * i) * \sin[\theta/2]^2 * (\cos \xi - i * \sin \xi)) + (i/16) * \sin \theta^2 * (i * (\cos \xi + i * \sin \xi) - i * \cos \theta * (\cos \xi + i * \sin \xi) - (2 * i) * \sin[\theta/2]^2 * (\cos \xi + i * \sin \xi)) * (\psi_{2\downarrow}^\dagger d\psi_{2\downarrow}) \end{aligned}$$

The price to pay...

order $1/(T_K)^6$, 10 GB, 1 km²

order $1/(T_K)^4$, 1 MB, 100 m²



order $1/(T_K)^2$, 4KB, 0.2m²

.

order $1/T_K$, 400 Bytes, 0.01m²

Results

DC bias applied to the two baths

Net analytical result : current

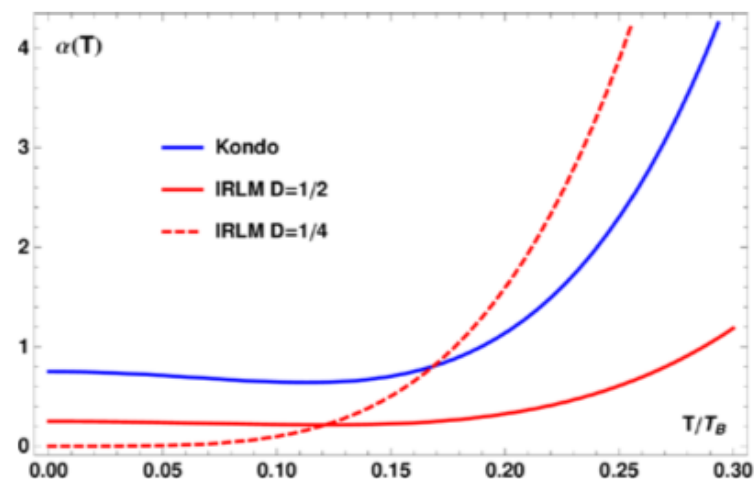
$$I(V) = 2 \sin^2 \theta V - \frac{\sin^2 \theta V^3}{4T_B^2} + \frac{(252\sqrt{3} + 65\pi + 5\pi \cos 2\theta) \sin[\theta]^2 V^5}{1920\pi T_B^4} + O[V^7/T_B^7]$$

linear response

Fermi Liquid result

Temperature dependence of non linearity captured !

$$I(V)/I_0 V = (1 + \alpha(T) V^2/T_B^2 + \dots)$$



Results

DC bias: Universal differential conductance $G=dI/dV$

Rescaled quantities

$$\bar{X} = X / T_K$$

$$\frac{G(\bar{V}, \bar{T}, \bar{h})}{G_0} = \underbrace{1 - c_T \bar{T}^2 - \alpha_V c_T \bar{V}^2 - \alpha_h c_T \bar{h}^2}_{\text{Fixed point}} \underbrace{\left[\begin{aligned} &+ \bar{V}^2 c_T^2 (\gamma_V \bar{V}^2 + \gamma_T \bar{T}^2 + \gamma_h \bar{h}^2) + \rho_{Th} c_T^2 \bar{T}^2 \bar{h}^2 + \rho_h c_T^2 \bar{h}^4 + \rho_T c_T^2 \bar{T}^4 \\ &- \bar{V}^4 c_T^3 (\kappa_V \bar{V}^2 + \kappa_T \bar{T}^2 + \kappa_h \bar{h}^2) - \bar{T}^4 c_T^3 (\beta_V \bar{V}^2 + \beta_T \bar{T}^2 + \beta_h \bar{h}^2) \\ &- \bar{h}^4 c_T^3 (\chi_V \bar{V}^2 + \chi_T \bar{T}^2 + \chi_h \bar{h}^2) - \gamma_{Th} c_T^3 \bar{V}^2 \bar{T}^2 \bar{h}^2 + \mathcal{O}(T_K^{-7}) \end{aligned} \right]}_{\text{Fermi liquid corrections}}$$

NEW !

$$\alpha_V = 3/2\pi^2 \approx 0.15 \quad c_T = \pi^2/4 \approx 2.46 \quad \alpha_h = 1/\pi^2 \approx 0.10$$

$$\gamma_V = \frac{252\sqrt{3} + 65\pi + 5\pi \cos 2\theta}{48\pi^5} \approx 0.04$$

$$\rho_{Th} = 2(5\sqrt{3} + \pi)/\pi^3 \approx 1.52$$

$$\gamma_T = \frac{72\sqrt{3} + 17\pi + \pi \cos 2\theta}{4\pi^3} \approx 1.44^{\pm 0.03}$$

$$\rho_T = \frac{5}{3} + \frac{36\sqrt{3}}{5\pi} \approx 5.64$$

$$\gamma_h = 3(5\sqrt{3} + \pi)/\pi^5 \approx 0.12$$

$$\rho_h = (6\sqrt{3} + \pi)/3\pi^5 \approx 0.02$$

Results

DC bias: Universal differential conductance $G=dI/dV$

$$\frac{G(\bar{V}, \bar{T}, \bar{h})}{G_0} = \underbrace{1 - c_T \bar{T}^2 - \alpha_V c_T \bar{V}^2 - \alpha_h c_T \bar{h}^2}_{\text{Fixed point}} + \underbrace{\bar{V}^2 c_T^2 (\gamma_V \bar{V}^2 + \gamma_T \bar{T}^2 + \gamma_h \bar{h}^2) + \rho_{Th} c_T^2 \bar{T}^2 \bar{h}^2 + \rho_h c_T^2 \bar{h}^4 + \rho_T c_T^2 \bar{T}^4}_{\text{Fermi liquid corrections}}$$

$$- \bar{V}^4 c_T^3 (\kappa_V \bar{V}^2 + \kappa_T \bar{T}^2 + \kappa_h \bar{h}^2) - \bar{T}^4 c_T^3 (\beta_V \bar{V}^2 + \beta_T \bar{T}^2 + \beta_h \bar{h}^2)$$

$$- \bar{h}^4 c_T^3 (\chi_V \bar{V}^2 + \chi_T \bar{T}^2 + \chi_h \bar{h}^2) - \gamma_{Th} c_T^3 \bar{V}^2 \bar{T}^2 \bar{h}^2 + \mathcal{O}(T_K^{-7})$$

NEW !

$$\kappa_V = \frac{12960 + 19025\sqrt{5} + 7308\sqrt{3}\pi + 420\pi^2}{720\pi^8} \approx 0.01$$

$$\kappa_T = \frac{5265 + 7375\sqrt{5} + 2826\sqrt{3}\pi + 150\pi^2}{36\pi^6} \approx 1.11$$

$$\kappa_h = \frac{1215 + 1850\sqrt{5} + 756\sqrt{3}\pi + 35\pi^2}{120\pi^8} \approx 0.08$$

$$\beta_V = \frac{4455 + 6300\sqrt{5} + 2376\sqrt{3}\pi + 121\pi^2}{15\pi^4} \approx 22.36$$

$$\beta_T = \frac{2(14580 + 21250\sqrt{5} + 8316\sqrt{3}\pi + 427\pi^2)}{315\pi^2} \approx 71.77$$

$$\beta_h = \frac{2(3240 + 4600\sqrt{5} + 1674\sqrt{3}\pi + 75\pi^2)}{45\pi^4} \approx 10.66$$

$$\chi_V = \frac{8(117 + 175\sqrt{5} + 56\sqrt{3}\pi + 2\pi^2)}{\pi^8} \approx 0.70$$

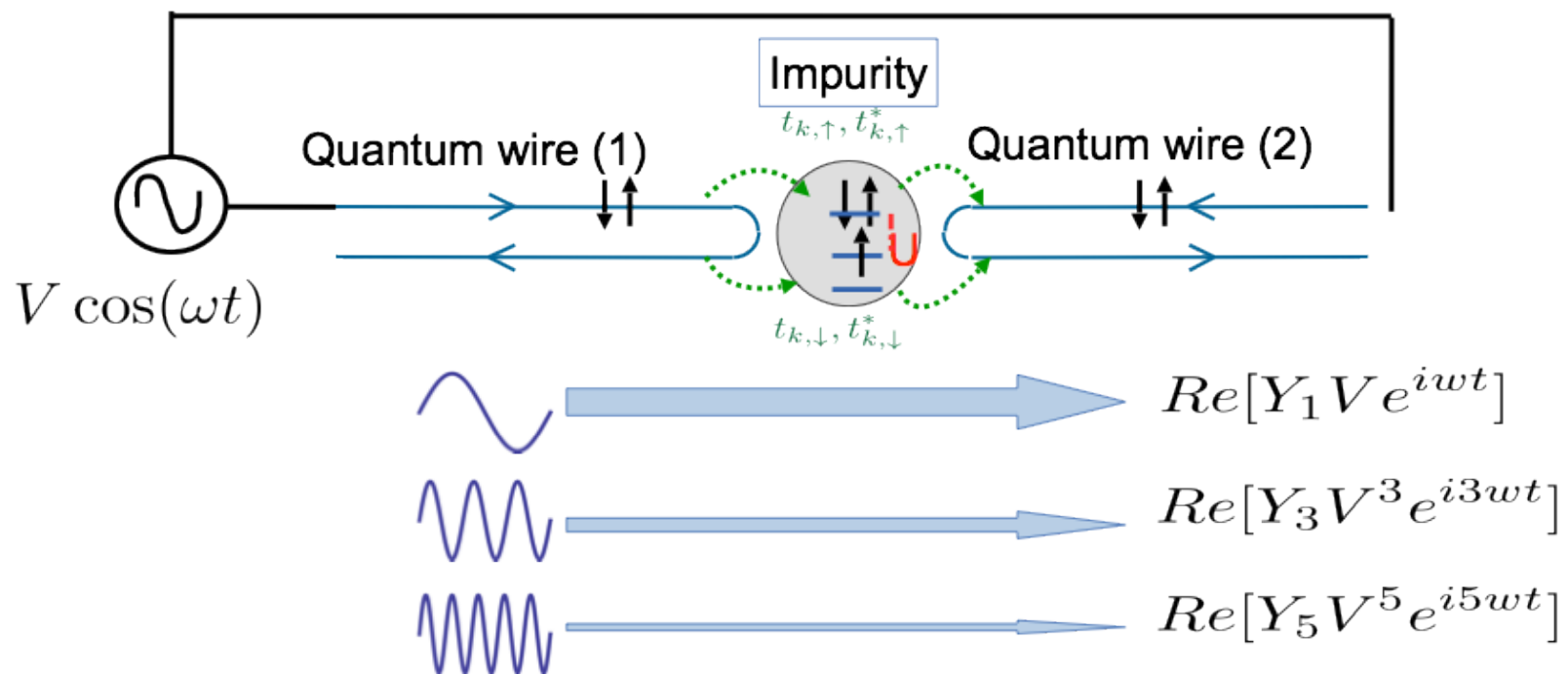
$$\chi_T = \frac{16(117 + 175\sqrt{5} + 56\sqrt{3}\pi + 2\pi^2)}{3\pi^6} \approx 4.62$$

$$\chi_h = \frac{(45(3 + 5\sqrt{5}) + 60\sqrt{3}\pi + 2\pi^2)}{45\pi^8} \approx 0.002$$

$$\gamma_{Th} = \frac{4428 + 6050\sqrt{5} + 2160\sqrt{3}\pi + 96\pi^2}{12\pi^6} \approx 2.66$$

Results

AC bias applied to the two baths

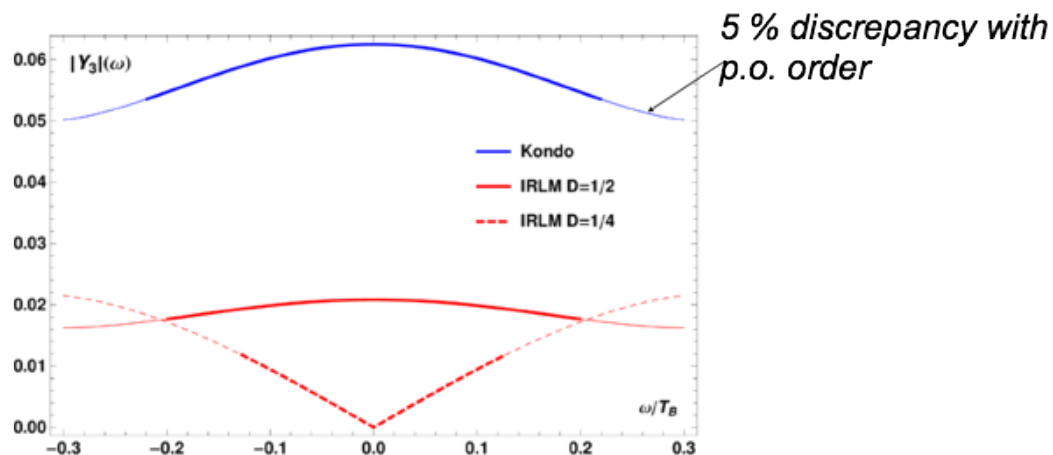


Higher harmonics can be captured ! $Y_a(V, \omega, T)$

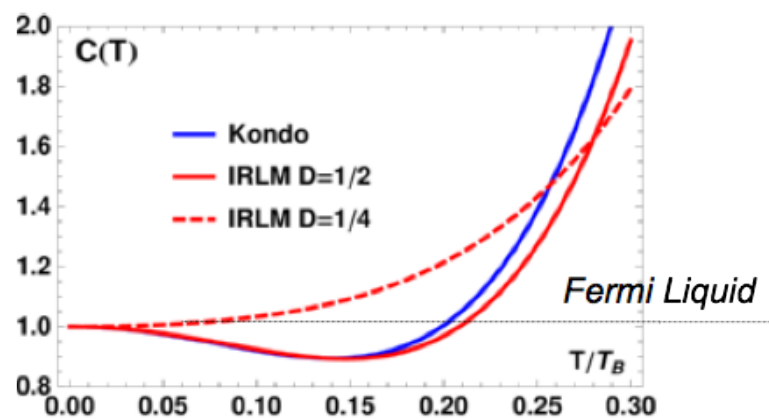
Results

AC bias applied to the two baths

3rd harmonic :



Capacitance $C(T)$, imaginary part of $Y_1(T)$



Results

Exact results for the noise !

Symmetric noise :

$$S(\tau) = (I(t + \tau)I(t) - I^2) + (I(t)I(t + \tau) - I^2)$$

Noise spectrum :

$$S(\omega) = \int e^{i\omega\tau} S(\tau) d\tau$$

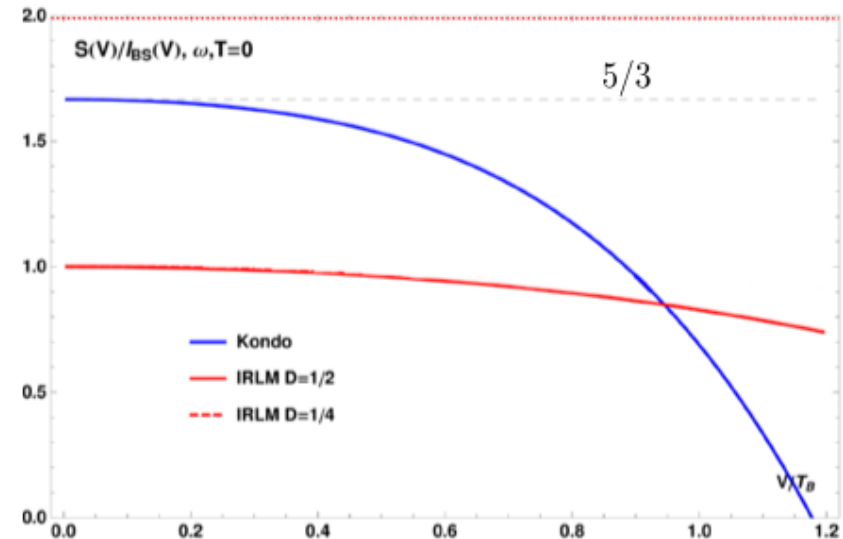
Results

Fano factor

$$F = \frac{1}{2e} \frac{S}{I_{BS}}$$

At $V=0$: $F = \frac{5}{3}$

Sela, Oreg, von Oppen, Koch, + Gogolin, Komnik: PRL 2006



Corrections : $C = \cos \theta$

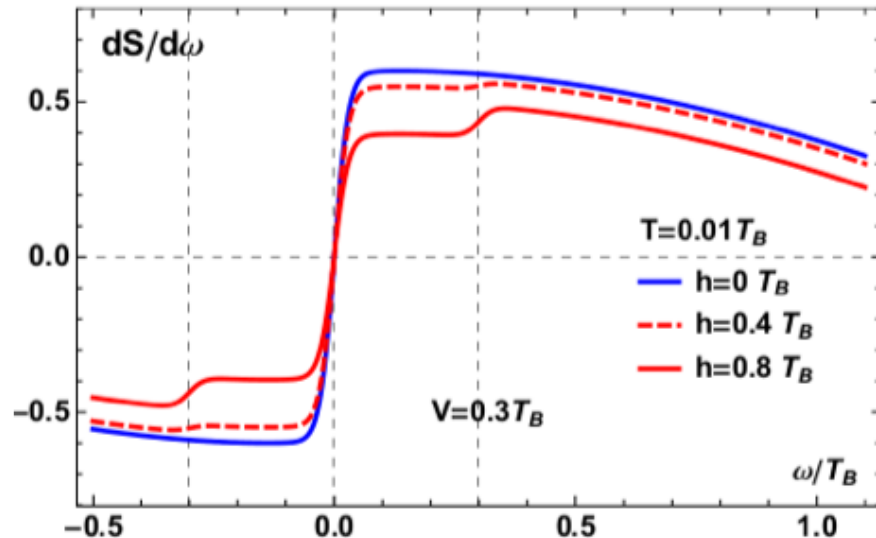
$$F = \frac{1}{3} (5 - 8C^2) + \frac{(C^2 - 1) (\pi (25C^2 + 153) - 18\sqrt{3}) V^2}{360\pi T_B^2} + \frac{((513297 + 782875\sqrt{5} + 376866\sqrt{3}\pi + 37275\pi^2) C^2 - 42630\pi^2 - 313866\sqrt{3}\pi - 475625\sqrt{5} - 276372) V^4}{604800\pi^2 T_B^4} + \frac{(875\pi^2 C^6 - 70(972 + 500\sqrt{5} - 882\sqrt{3}\pi - 163\pi^2) C^4) V^4}{604800\pi^2 T_B^4}$$

Results

At T=0 : non-analyticities

$$\begin{aligned}
 S^{(0)}(\omega, V, 0) = & \frac{4 \sin^2(\theta)}{\pi (W^2 + 4)^2} ((|V - \omega| + |V + \omega|) (4 \cos^2(\theta) + W^2) + 8|\omega| \sin^2(\theta)) \\
 & - \frac{W \sin^2(\theta) \cos(\theta)}{\pi (W^2 + 4)^2 T_B} (|V - \omega| (8V \cos(2\theta) + W^2(V - \omega) - 4(V + \omega)) \\
 & + |V + \omega| (8V \cos(2\theta) + W^2(V + \omega) - 4(V - \omega)) + 32V|\omega| \sin^2(\theta))
 \end{aligned}$$

*Derivative of the noise at low temperature :
Non analyticities at T=0 are rounded by temperature*



Conclusions

- ✓ A generic – **but perturbative** – method for some integrable systems: super Fermi liquids
- ✓ Gives exact (formal) expression for the out-of-equilibrium density matrix
- ✓ Yields exact results in variety of conditions:
 - ✓ Voltage (AC/DC)
 - ✓ Finite temperature(s)
 - ✓ Particle-hole asymmetry
 - ✓ Magnetic field
- ✓ Perspectives:
 - ✓ (Slow) quenches
 - ✓ Non Fermi liquid fixed points?