A Tool for Dark Matter and New Physics

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Dark Matter Codes



Any model of New Physics/DM, relic density, direct and indirect

 $detection, \dots$



P. Gondolo, J. Edsjo, L. Bergstrom, P. Ullio, M. Schelke, T. Baltz, T. Bringmann and

G. Duda; Supersymmetry, relic density, direct and indirect detection,...





$\mathbf{\Omega}$

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" O atomes intelligents, dans qui l'Etre éternel s'est plu à manifester son adresse et sa puissance, vous devez sans doute goûter des joies bien pures sur votre globe car, ayant si peu de matière...," Voltaire, Micromegas, chapitre septième, conversation avec les hommes

Tools from particle physics

Need powerful, modular and versatile tools

micrOMEGAs: Tools for DM/ Collider Physics/ Flavour for a general NP

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Aim



Aim

Relic Density

direct detection Xenon, CDMS, Edelweiss... (CoGeNT, Dama/Libra...)

production at colliders

γ from annihil in galactic center or halo and from synchrotron emission Fermi, ICT, radio telescopes..

indirect e

from annihil in galactic halo or center PAMELA, Fermi, HESS, AMS, balloons... from annihil in galactic halo or center

d from annihil in galactic halo or center $_{
m GAPS}$

 $u, \overline{
u}$ from annihil in massive bodies

SK, Icecube, Km3Net

micrOMEGAS: Guiding Principle

Guiding principles

Humans make mistakes - computers do not Automation (with auto checks)

- Several groups are developing specialized codes Link them
- Users might want to improve one aspect Modularity
- We do not know what DM is made of Possibility to include different DM candidates
- Models are often complex with huge parameter space Speed of execution
- Ready made, stand-alone package for the non-expert User friendly

LANHEP

micrOMEGAs











- given any set of parameters it can identify LSP, NLSP, generate and calculate Ωh^2
- Model defined in Lanhep (more later)
- \blacktriangleright Fed into CalcHEPtree-level, some 3000 processes could be needed.
- Higgs sector (SUSY): improved Higgs masses/mixings (read from FeynHiggs, for example) but interpreted in terms of an effective scalar potential (GI), following FB and A. Semenov (PRD 02)
- Effective Lagrangian also includes important RC (Higgs couplings, Δm_b effects,..)
- Interfaced with Isajet, Suspect, SoftSUSY parameters at high scale run down to the ew scale
- $(g-2)_{\mu}, b \to s\gamma, B_s \to \mu^- \mu^+, \dots$
- NMSSM (with C. Hugonie), CP violation (with S. Kraml), UED, Dirac, host of others
 - "open source": procedure to define your own model
- powerful generalised direct detection module
- indirect detection cross sections, interface with propagation, polarisation completed
 - SLHA compliant, MCMC interface for parameter scans

given any set of parameters it can identify LSP, NLSP, generate and calculate Ωh^2 , direct and indirect

cross sections are generated on the fly. Only those needed are generated. These are then stored, if needed in the future procedure is speedy

Generalisation to other models easy, same principle: Needs a model file and a quantum number based on same Z_2 (in SUSY $Z_2 = R_p$) or even $\dots Z_N$. Classification is then straightforward.

set a switch (that can be changed by the user) so that even co-annihilation processes are generated on the fly. The code decides on *its OWN* when to include these co-annihilations.

DarkSusy

comes with a (large) set of codes for cross sections in SUSY.

you may well have a weirdo SUSY model for which some (co-annihilation) cross sections are not provided, you will have to code them. In principle you could code a _ cross section from any model

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micrOMEGAS: History

History: more than 10 years of micrOMEGAs

- 2001: First release, initial goal : compute the relic density in the MSSM including all processes
- 2004 : Exact numerical solution of relic density + improved MSSM (in particular loop corrected particle masses)
- 2006 : Direct detection
- 2008 : Generalization to arbitrary Particle Physics model
- 2010 : Indirect detection signatures
- 2013 : generalization relic density, Higgs observables, neutrino signature

Relic Density: Thermal average

must calculate all annihilation, co-annihilation processes. Each annihilation can consist of tens of cross sections...

$$\begin{split} \chi_{i}^{0}\chi_{j}^{0} &\to X_{SM}Y_{SM}, \chi_{1}^{0}\tilde{f}_{1} \to X_{SM}Y_{SM}, \dots \underbrace{v \times \sigma v}_{(x \times \sigma v)} \\ &< \sigma v > = \frac{\sum\limits_{i,j}^{g_{i}g_{j}} \int\limits_{(m_{i} + m_{j})^{2}} ds \sqrt{s} \underbrace{K_{1}(\sqrt{s}/T)}_{2T\left(\sum\limits_{i}g_{i}m_{i}^{2}K_{2}(m_{i}/T)\right)^{2}}, \\ p_{ij} \text{ is the momentum of the incoming particles in their center-of-mass frame.} \\ \\ \text{Origin of Boltzman factor} \underbrace{exp(-\delta M/T)}_{exp(-\delta M/T)} \end{split}$$

$$B_f = \frac{K_1((m_i + m_j)/T_f)}{K_1(2m_{\tilde{\chi}_1^0}/T_f)} \approx e^{-X_f \frac{(m_i + m_j - 2m_{\tilde{\chi}_1^0})}{m_{\tilde{\chi}_1^0}}} \quad if \ B_f < B_\epsilon \ \text{do not compute}$$

In micrOMEGAs $B_{\epsilon} = 10^{-6}$ by default. Most often $B_{\epsilon} = 10^{-2}$ enough for 1% accuracy, only a few processes computed.

In micrOMEGAs $< \sigma v > (T)$ is obtained from direct integration, through adaptive Simpson with two options (fast, default) and accurate (checks for accuracy).

Direct detection



ingredients/Modules: dark matter density and modulation, velocity distribution quark content in nucleon, Nuclear form factors,.....



Wimp-quark effective Lagrangian, $\psi_N \rightarrow \psi_q, \quad \chi \rightarrow \mathbf{DM}$

| | WIMP | Even Operators | Odd Operators | | |
|----|------|--|--|--|--|
| | Spin | $\hat{\mathcal{O}}_{q,e} \ \hat{\mathcal{O}}_{q,e}^{\prime}$ | $\hat{\mathcal{O}}_{q,o}\hat{\mathcal{O}}_{q,o}^{\prime}$ | | |
| | | $\hat{\mathcal{O}}_{q,e}$ | $\hat{\mathcal{O}}_{q,o}$ | | |
| | 0 | $2M_\chi \phi_\chi \phi_\chi^* \overline{\psi}_q \psi_q$ | $i(\partial_\mu \phi_\chi \phi_\chi^* - \phi_\chi \partial_\mu \phi_\chi^*) \overline{\psi}_q \gamma^\mu \psi_q$ | | |
| SI | 1/2 | $\overline{\psi_{\chi}}\psi_{\chi}\overline{\psi}_{q}\psi_{q}$ | $\overline{\psi}_{\chi}\gamma_{\mu}\psi_{\chi}\overline{\psi}_{q}\gamma^{\mu}\psi_{q}$ | | |
| | 1 | $2M_{\chi}A^*_{\chi\mu}A^{\mu}_{\chi}\overline{\psi}_{q}\psi_{q}$ | $+i(A_{\chi}^{*\alpha}\partial_{\mu}A_{\chi,\alpha}-A_{\chi}{}^{\alpha}\partial_{\mu}A_{\chi\alpha}^{*})\overline{\psi}_{q}\gamma_{\mu}\psi_{q}$ | | |
| | | | | | |
| | 1/2 | $\hat{\mathcal{O}}_{q,e}^{\prime} \ \overline{\psi}_{\chi} \gamma_{\mu} \gamma_{5} \psi_{\chi} \overline{\psi}_{q} \gamma_{\mu} \gamma_{5} \psi_{q}$ | $\hat{\mathcal{O}}_{q,o}' \ -rac{1}{2}\overline{\psi}_{\chi}\sigma_{\mu u}\psi_{\chi}\overline{\psi}_{q}\sigma^{\mu u}\psi_{q}$ | | |
| SD | 1 | $\sqrt{6}(\partial_{\alpha}A_{\chi\beta}^{*}A_{\chi\nu} - A_{\chi\beta}^{*}\partial_{\alpha}A_{\chi\nu})$ | $i\frac{\sqrt{3}}{2}(A_{\chi\mu}A^*_{\chi\nu} - A^*_{\chi\mu}A_{\chi\nu})\overline{\psi}_q\sigma^{\mu\nu}\psi_q$ | | |
| | | $\epsilon^{lphaeta u\mu}\overline{\psi}_q\gamma_5\gamma_\mu\psi_q$ | | | |
| | | | | | |

$$\hat{\mathcal{L}}_{eff}(x) = \sum_{q,s} \lambda_{q,s} \hat{\mathcal{O}}_{q,s}(x) + \xi_{q,s} \hat{\mathcal{O}}'_{q,s}(x)$$

In model files of micrOMEGAs (CalcHEP) these operators are added

- In the usual approach these low energy operators and their coefficients are extracted by computing WIMP-quark *amplitudes* from Feynman diagrams and using Fierz transformations,..
- In micrOMEGAs all operators are defined and only need to extract coefficients automatically
- \checkmark we compute $\chi q \rightarrow \chi q$ at $q^2 = 0$ as a normal cross section but...
- Interference between one projection operator and an effective vertex singles out SI or SD
- If with the S-matrix, $\hat{S} = 1 i\mathcal{L}$ obtained from the complete Lagrangian at the quark level

$$\lambda_{q,e} + \lambda_{q,o} = \frac{-i\langle q(p_1), \chi(p_2) | \hat{S}\hat{\mathcal{O}}_{q,e} | q(p_1), \chi(p_2) \rangle}{\langle q(p_1), \chi(p_2) | \hat{\mathcal{O}}_{q,e}\hat{\mathcal{O}}_{q,e} | q(p_1), \chi(p_2) \rangle}$$

$$\lambda_{q,e} - \lambda_{q,o} = \frac{-i\langle \overline{q}(p_1), \chi(p_2) | \hat{S}\mathcal{O}_{q,e} | \overline{q}(p_1), \chi(p_2) \rangle}{\langle \overline{q}(p_1), \chi(p_2) | \hat{\mathcal{O}}_{q,e}\hat{\mathcal{O}}_{q,e} | \overline{q}(p_1), \chi(p_2) \rangle}$$

warning: couplings proportional to light quark masses must be kept

Indirect Detection



- ν from the Sun and the Earth (**New in 3_0**)
- \checkmark γ , ν , charged cosmic rays ($e^+, \overline{p}, \overline{D}$) from annihilation in the galactic halo (\overline{D} in future)
- other signals: synchrotron emission,....? (in progress)

Features in the Indirect Detection Module of micrOMEGas

- Annihilation cross sections for all 2-body tree-level processes for all models.
- Annihilation cross sections including radiative emission of a photon for all models.
- Annihilation cross sections into polarized gauge bosons.
- Annihilation cross sections for the loop induced processes $\gamma\gamma$ and γZ^0 in the MSSM and NMSS (more general models possible)
- Modelling of the DM halo with a general parameterization and with the possibility of including DM clumps.
- Integrals along lines of sight for γ -ray signals.
- Propagation: Work with Pierre Salati and Sylvie Rosier (2-zone diffusion model with Green's function and tabulation. Diffusion parameters compatible with B/C.)
 Computation of the propagation of charged particles through the Galaxy, including the possibility to modify the propagation parameters.
- Effect of solar modulation on the charged particle spectrum.
- Model independent predictions of the indirect detection signal

The neutrino spectrum originating from dark matter annihilation is also computed. With in version

3.0 Capture in the Sun and the Earth

N. Baro, FB, G. Chalons, G. Drieu La Rochelle, S. Hao, Ninh Le Duc, A. Semenov, (D. Temes)



SloopS, examples for relic

| | | | Tree | $A_{	au	au}$ | $\overline{\mathrm{DR}}$ | | | | |
|--|---------------------------------------|---|---------|--------------|--------------------------|--|--|--|--|
| $\overline{	ilde{\chi}_1^0	ilde{\chi}_1^0}$ – | $\rightarrow W^+W^-$ [26%] | a | +11.84 | +4.3% | +5.1% | | | | |
| | | b | +4.17 | +12.7% | +13.4% | | | | |
| $	ilde{\chi}^0_1 	ilde{\chi}^+_1$ - | $\rightarrow u \overline{d} \ [12\%]$ | a | +15.28 | +6.8% | +7.0% | | | | |
| | | b | -5.31 | +30.4% | +30.7% | | | | |
| $-	ilde{\chi}^0_1	ilde{\chi}^0_1$ – | $\rightarrow Z^0 Z^0 [9\%]$ | a | +4.28 | +10.4% | +9.6% | | | | |
| | | b | +1.83 | +12.7% | +12.0% | | | | |
| $	ilde{\chi}^0_1 	ilde{\chi}^+_1$ - | $\rightarrow Z^0 W^+ [6\%]$ | a | +6.99 | +1.7% | +2.1% | | | | |
| | | b | -0.51 | +85.6% | +86.5% | | | | |
| $\Omega_\chi h^2$ | | | 0.00931 | 0.00909 | 0.00908 | | | | |
| $rac{\delta\Omega_\chi h^2}{\Omega_\chi h^2}$ | | | | -2.4% | -2.5% | | | | |
| Tree-level values of the s-wave (a) and p-wave (b) coefficients in units $10^{-26} \text{cm}^3 \text{s}^{-1}$ in a | | | | | | | | | |
| higgsino scenario | | | | | | | | | |
| Baro, Chalons, Sun Hao FB 2009 | | | | | | | | | |

Form Factor Approach

Recent Work: FB, G. Drieu la Rochelle, A. Mariano (see arXiv 2011-2014) **Effective form factors** Can we express most of the corrections within a form factor approach. A Born-Improved approximation. Build up a library of form factors to replace tree-level vertices. One can then use tree-level automated calculators and hence easy interface with MicrOMEGAS.



Example of the one-loop vertex diagrams that can be cast into an effective $\tilde{\chi}_1^0 \tilde{\chi}_j^0 Z$ vertex.

Performance (difficult one $\chi_1^0 \chi_1^0 \to ZZ$, performance better for $\chi_1^0 \chi_1^0 \to \mu^+ \mu^-$)



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micrOMEGAS_3.0: new features

Generalization of relic density computation

3-body final state Discrete symmetries other than Z_2 Asymmetric dark matter Different universe expansion

Cross sections and decays: Higgs sector Higgs decays into virtual particles Loop-induced Higgs decay

DM detection

Direct detection

Neutrino signal from DM capture (reduce uncertainty from strange content, use new lattice results)

Models distributed

MSSM NMSSM (with C. Hugonie) CPV-MSSM (with S. Kraml) Right-handed neutrino (with G. Servant) Little Higgs (A. Belyaev) Inert Doublet Model Inert doublet+singlet Z3 (with K. Kannike, M. Raidal)

Models not yet public

SUSY N=2 (with K. Benakli, M. Goodsell)
UED (with M. Kakizaki)
MSSM+RHneutrino (with M. Kakizaki, S. Kraml, E.K. Park)
UMSSM (with J. DaSilva,)
BMSSM (with G. Drieu La Rochelle)

Many more models implemented by users

micrOMEGAs_3.0

- Annihilation cross sections for some selected 3-body processes in addition to the 2-body tree-level processes. The 3-body option can be included in the computation of the relic density and/or for annihilation of dark matter in the galaxy.
- Possibility of using different tables for the effective degrees of freedom $(g_{eff}(T), h_{eff}(T))$ in the early Universe
- Annihilation cross sections for the loop induced processes $\gamma\gamma$ and γZ^0 in the NMSSM and the CPVMSSM
- New function for incorporating DM clumps
- New function to define the strange quark content of the nucleon
- The LanHEP source code for new models is included
- ${}$ New models with scalar DM (Inert doublet model and model with Z_3 symmetry)
- New implementation of the NMSSM which uses the Higgs self-couplings and the particle spectrum calculated in NMSSMTools⁴.0
- New versions of spectrum generators used in the MSSM (Suspect_2.4.1) and in the CPVMSSM (CPsuperH2.3)
- Extended routines for flavour physics in the MSSM
- Update in interface tools to read files produced by other codes, this allows easy interface to other codes

micrOMEGAs_3.0: Asymmetric dark matter

We consider a dark matter particle that is not self-conjugate. This occur for example when the DM particle is a Dirac fermion or a complex scalar. We define Y^+ and Y^- as the abundances of the DM particle/anti-particle $(\chi \overline{\chi})$. At very high temperature we assume that an asymmetry has taken place such that $Y^+ \neq Y^-$ ($\Delta Y = Y^+ - \neq Y^-$) but that at later time the system is at equilibrium. Assume that some symmetry only allows annihilation $\chi^+\chi^- \to SM$ (no annihilation $\chi^+\chi^+$ for example). The relic density is always larger than in the case of no asymmetry. We can have a situation where only χ^+ survives and hence no indirect signal survives

$$\Omega h^{2} = \frac{8\pi}{3H_{100}^{2}} \frac{m_{\chi}}{M_{\text{Planck}}} \frac{\sqrt{Y_{0}^{2} + \Delta Y^{2}}}{s_{0}} > \Omega h^{2} (\Delta Y = 0).$$

micrOMEGAs_3.0: Asymmetric dark matter



 Ωh^2 as a function of ΔY for different DM masses in the 4th generation Dirac neutrino dark matter model for DM masses of 3, 6, 10, 50, 100 GeV as indicated.

$\texttt{micrOMEGAS}__{3.0:}$ Annihilation into Three-body and Four-body

Simple observation: $\chi \chi' \to XV \ (V, W, Z)$ closed but $\chi \chi' \to XV^{\star} \to Xll'$ open



 Ωh^2 as a function of M_{DM} in the MSSM (full) and relative difference between the 3-body and 2-body value (dashed).

microMEGAS_3.0: DM with non Z_2 discrete symmetry: $Z_3, \cdots Z_N$

DM stabilized with a larger symmetry than Z_2 Ex: Z_3 custodial $SU(2) W^{+\prime}W^{-\prime} \rightarrow Z'H$ semi-annihilation.

Example of $\chi_i \chi_j \to \chi_k A$, χ 's stable A unstable (decays eventually to SM). Lead to coupled Boltzman equations.

$$\frac{dn}{dt} = -\langle v\sigma^{\chi\overline{\chi}\to XX} \rangle \left(n^2 - n_{\rm eq}^2 \right) - \frac{1}{2} \langle v\sigma^{\chi\chi\to\overline{\chi}X} \rangle \left(n^2 - n n_{\rm eq} \right) - 3Hn.$$

micrOMEGAS_3.0: Gamma-Ray line:

 $\gamma\gamma$ and $Z\gamma$ in MSSM and NMSSM this is done through <code>SloopS</code>

$micrOMEGAs__{3.0:}$ Dark matter profile and clumps

The DM density is given as the product of the local density at the Sun with the halo profile function

$$\rho(r) = \rho_{\odot} F_{halo}(r)$$

Possibility to set different halo profiles setProfileZhao ($\alpha, \beta, \gamma, rc$)

setProfileEinasto(α)

Clumps $< \rho^2 >$ can be significantly larger than $< \rho >^2$ A simple implementation with all clumps occupy the same small volume V_{cl} with a constant density.

$$<
ho^2>(r)=
ho(r)(
ho(r)+
ho^{eff}_{clump}(r))$$

or

To be more general, one could assume that ρ_{cl} and f_{cl} depend on the distance from galactic center. The effect of clumping is then described by the equation

$$<\rho^2>(r)=\rho(r)(\rho(r)+\rho^{eff}_{clump}(r))$$

and the function

ullet setRhoClumps (ho_{clump}^{eff})

allows to implement a more sophisticated clump structure.

$\texttt{micrOMEGAS}__{3.0:}$ The Higgs sector at colliders

Loop-induced Higgs production and decay: $H \to \gamma \gamma$ and $gg \to H$ (N)NLO corrections to the generic coupling $HF_{\mu\nu}F^{\mu\nu}$ can be be generated for any model. Need to implement the Higgs to scalars/fermions/vector bosons

$$\mathcal{L} = g_{h\psi\psi}\overline{\psi}\psi h + ig'_{h\psi\psi}\overline{\psi}\gamma_5\psi h + g_{h\phi\phi}M_{\phi}h\phi\phi + g_{hVV}M_VhV_{\mu}V^{\mu}$$

Implementation of signal strengths and interface to HiggsBounds

$$\mu_{\gamma\gamma}^{ggF} = \frac{\sigma(gg \to h)_{NP} Br(h \to \gamma\gamma)_{NP}}{\sigma(gg \to h)_{SM} Br(h \to \gamma\gamma)_{SM}}$$
$$\mu_{\gamma\gamma}^{VBF} = \frac{\sigma(WW \to h)_{NP} Br(h \to \gamma\gamma)_{NP}}{\sigma(WW \to h)_{SM} Br(h \to \gamma\gamma)_{SM}}$$

micrOMEGAS_3.0 Dark matter models

The $\ensuremath{\texttt{micrOMEGAS}}$ distribution now comes with the following models

- MSSM and some of its extensions (NMSSM and CPVMSSM)
- the little Higgs model (LHM),
- a model with a right-handed neutrino (RHNM),
- the inert doublet model (IDM),
- a model with an extended scalar sector and a Z3 symmetry(Z3ID)
- model with a fourth generation of leptons (SM4).

microMEGAs_3.0: Neutrino Signals: Earth, Sun,

Capture Rate Earth and Sun, C_χ from Gould 1987
 DM-Nucleus cross section (ref. direct detection), form-factor
 DM velocity dist. and local density
 Number density of nucleus (Sun: Asplund, Grevesse, Sauval 2004. Earth:
 McDonough 2003)

- Annihilation of captured DM: $A_{\chi\chi}$
- Evaporation, E_{χ} . may be important for light (< 5 GeV DM)
- Neutrino Flux:

$$\dot{N}_{\chi} = C_{\chi} - A_{\chi\chi}N_{\chi}^2 - E_{\chi}N_{\chi} ,$$

Neutrino spectra: from decays of fermions, gauge bosons, hadrons,...cascade decays. Effect of neutrino oscillations, from Cirelli et al 2005 (we use their tables).

Future developments

Indirect:

Improved propagation (with P. Salati et al) Interface USINE, GALPROP, CLUMPY...

Constraints:

Improve Collider limits - most likely link to independent code for LHC limits on new particles (sModels, S. Kraml et al)

Automatisation of new physics contribution to B observables, g-2 etc..

- Include more DM models
- Beyond one-loop: Form Factors and limitations, Sommerfeld resummation, Sudakov and Weak boson radiation for very heavy DM.

Conclusions

- To understand the nature of dark matter clearly requires information and cross checks from cosmology, direct and indirect detection as well as from collider physics
- MicrOMEGAS is tool to perform these analyses in a generic model that aims at high precision and high level of efficiency

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Soon on your

