

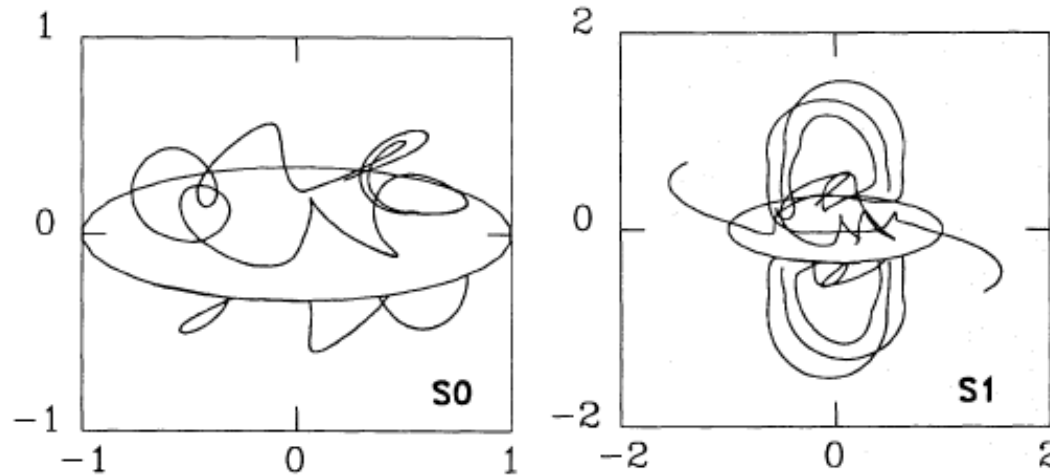
Faraday signature of helical and bi-helical magnetic fields

Axel Brandenburg (Nordita) & Rodion Stepanov (Perm)

- Compensating Faraday depolarization
 - By a helical field (single sense of twist)
 - By a bi-helical field & RM dependence
- Application to edge-on galaxies

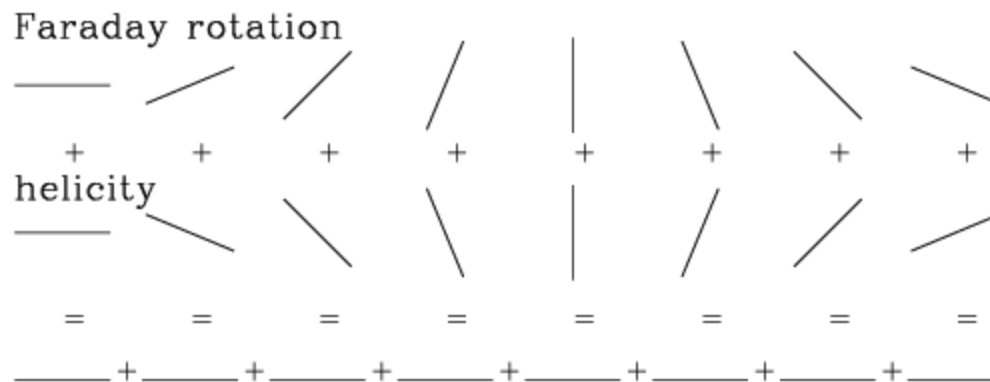


Helical (swirling) magnetic fields



Brandenburg &
Donner (1990)

Rotation from swirl compensates Faraday rotation



See also Sokoloff
et al. (1998)

Singly helical field

Stokes Q and U parameters $P = Q + iU$

$$Q = p_0 \int_{-\infty}^{\infty} \varepsilon \cos 2(\psi + \phi \lambda^2) dz$$

$$U = p_0 \int_{-\infty}^{\infty} \varepsilon \sin 2(\psi + \phi \lambda^2) dz$$

Intrinsic polarized emission from B

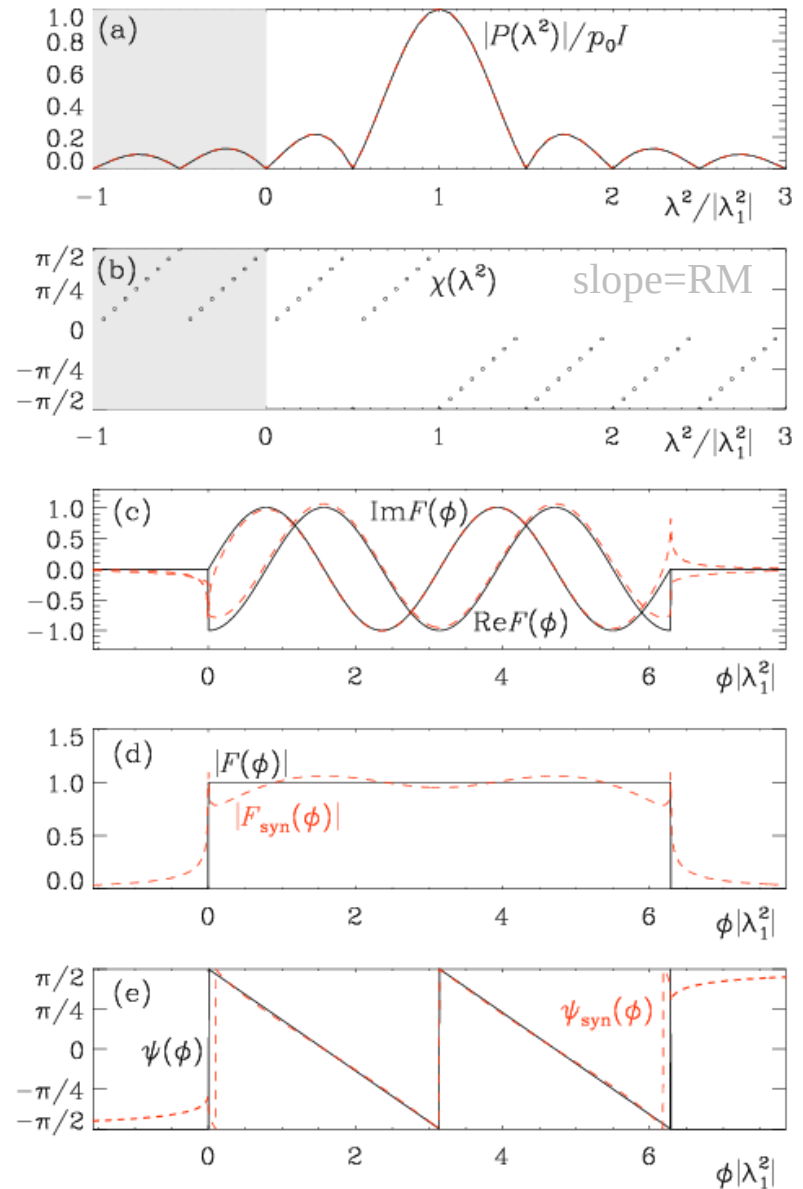
$$B_{\perp} = |B_{\perp}| e^{i\psi_B}, \quad \psi = \psi_B + \frac{1}{2} \pi$$

Cancellation condition

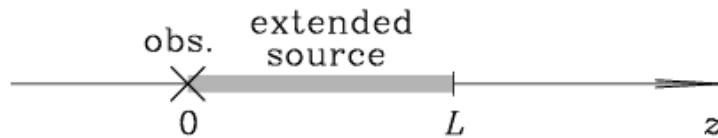
$$\psi = -kz, \quad \phi = -Kn_{\text{th}} B_0 z$$

Helical field w/
positive helicity

$$\mathbf{B} = \begin{pmatrix} B_1 \cos kz \\ -B_1 \sin kz \\ B_0 \end{pmatrix}$$



Only works if $RM > 0$ and $k > 0$



Peak determined by single parameter

$$\lambda_1^2 = -k / Kn_{\text{th}} B_0 \propto k / RM$$

But difficult/impossible to recover $F(\phi)$

(Burn 1966)

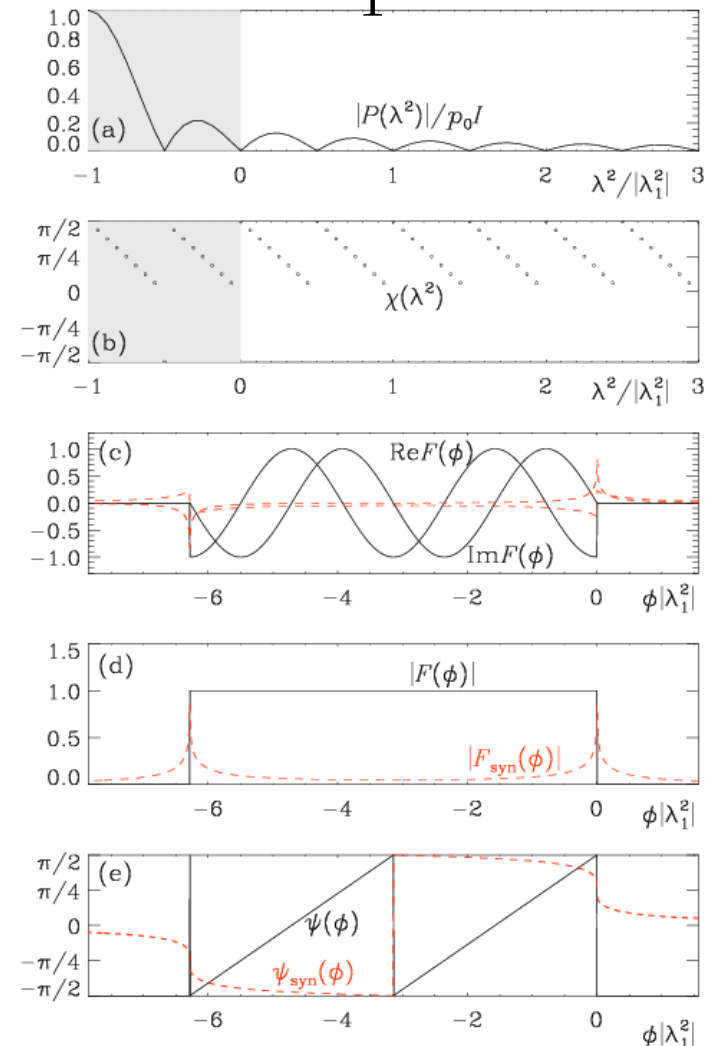
$$P(\lambda^2) = \int_{-\infty}^{\infty} F(\phi) e^{2i\phi\lambda^2} d\phi$$

$$F(\phi) = \frac{1}{2\pi} \int_{-\infty}^{\infty} P(\lambda^2) e^{-2i\phi\lambda^2} d(2\lambda^2)$$

Positivity:

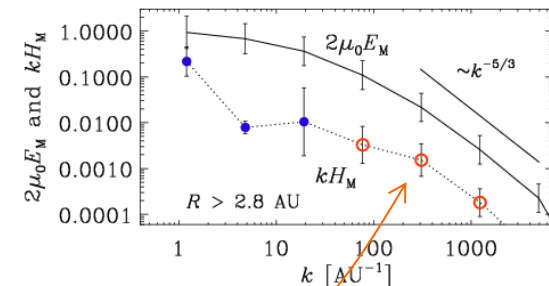
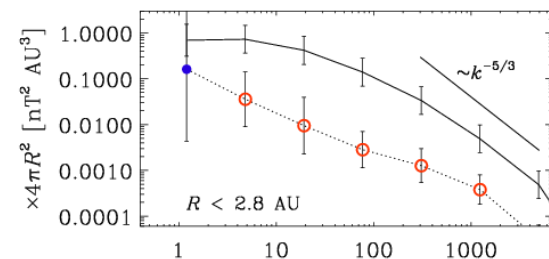
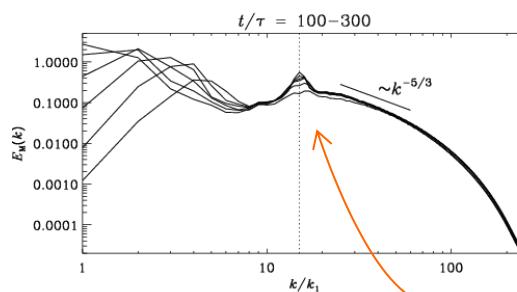
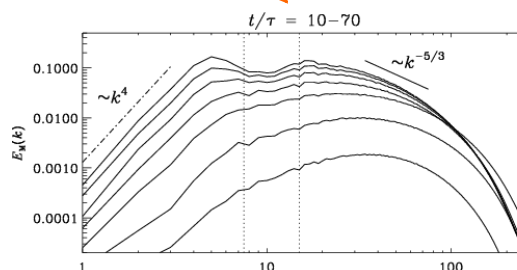
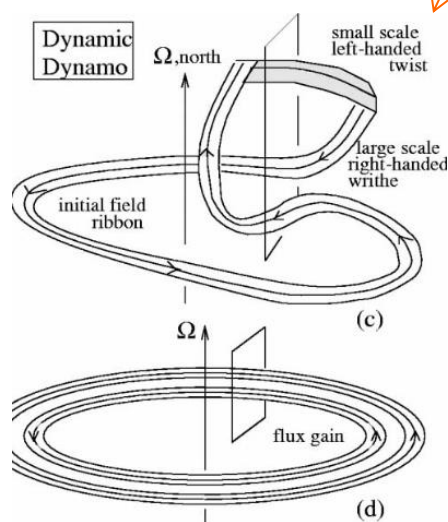
$$F_{\text{syn}}(\phi) = \frac{1}{2\pi} \int_0^{\infty} P(\lambda^2) e^{-2i\phi\lambda^2} d(2\lambda^2)$$

$$\lambda_1^2 < 0$$

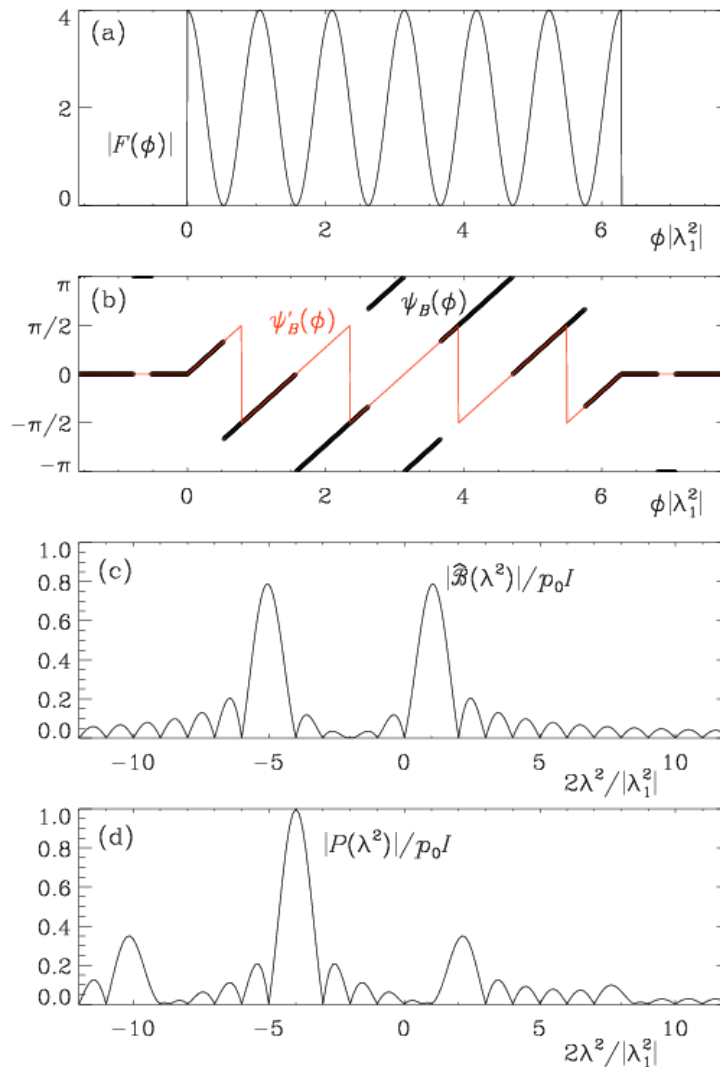


Expect bi-helical fields

- Magnetic helicity conserved
- Inverse cascade produces small-scale waste!
- Opposite sign of helicity (or k)



← ambiguity lead to “line splitting”



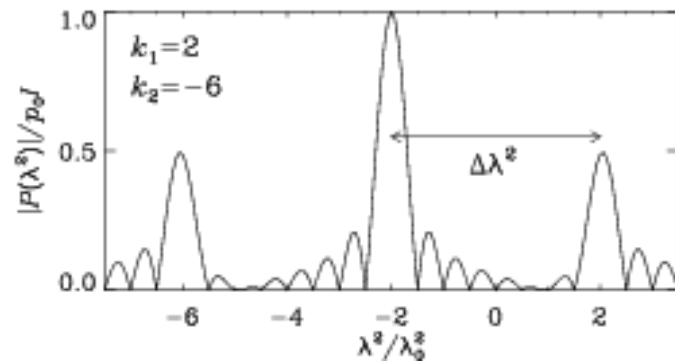
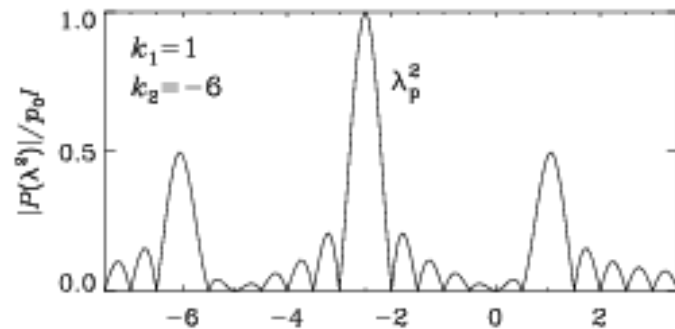
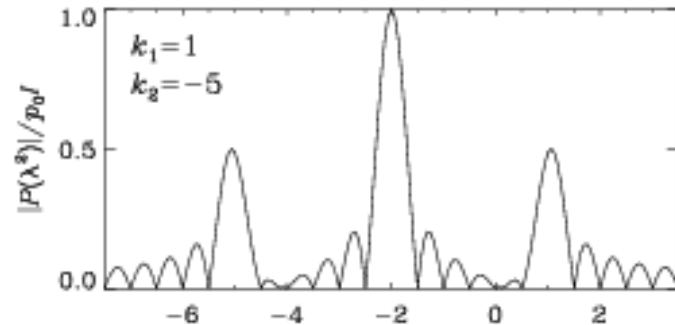
Peaks at $k_1=1$ and $k_2=-5$

translate to $k_1+k_2 = -4$
and to $k_1-k_2 = 6$

(i) peak in P at -4
peak separation 6

(ii) in Faraday dispersion:
frequency 6
-2x phase gradient -4

← ambiguity: other examples



Peaks at $k_1=1$ and $k_2=-5$
1 and -6, or, 2 and -6

translate to $k_1+k_2=-4$
and $k_1-k_2=6$

Table 1. Summary of the three examples shown in Fig. 4 for $\text{RM} > 0$ with bi-helical magnetic fields of wavenumbers k_1 and k_2 , the corresponding values of $k_{\pm} = k_1 \pm k_2$, the peak wavenumber k_p , the peak separation Δk , the phase gradient (ϕ derivative, indicated by ∇ for brevity), and corresponding values for λ_p^2 and $\Delta\lambda^2$. All values of k are normalized by k_0 and all values of λ^2 are normalized by λ_0^2 .

k_1	k_2	k_+	k_-	k_p	Δk	$\nabla\psi'_B$	λ_p^2	$\Delta\lambda^2$	$\nabla\psi$
1	-5	-4	6	-4	6	2	-2	3	2
1	-6	-5	7	-5	7	2.5	-2.5	3.5	2.5
2	-6	-4	8	-4	8	2	-2	4	2

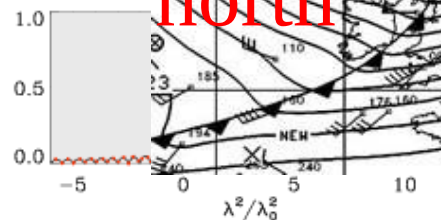
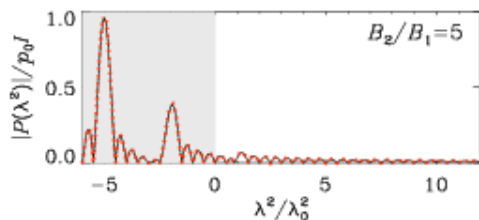
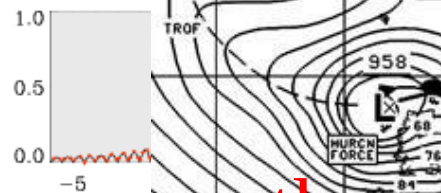
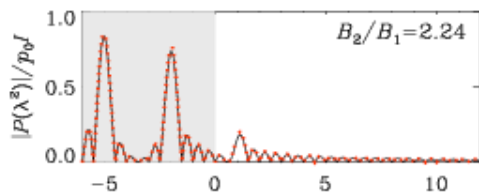
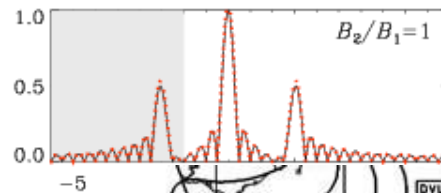
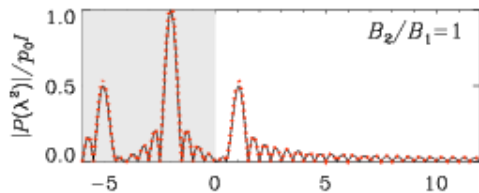
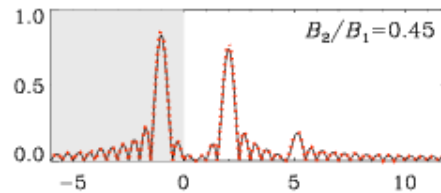
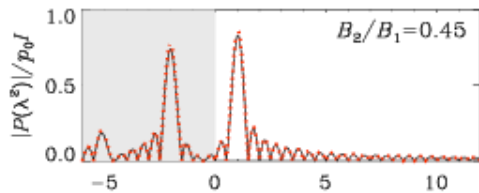
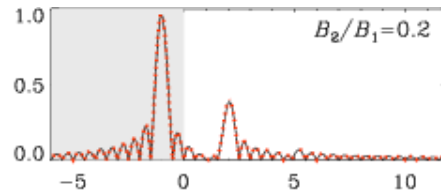
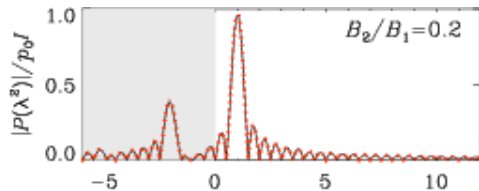
For $RM > 0$, only +ve helicity obs

$RM > 0$

$RM < 0$

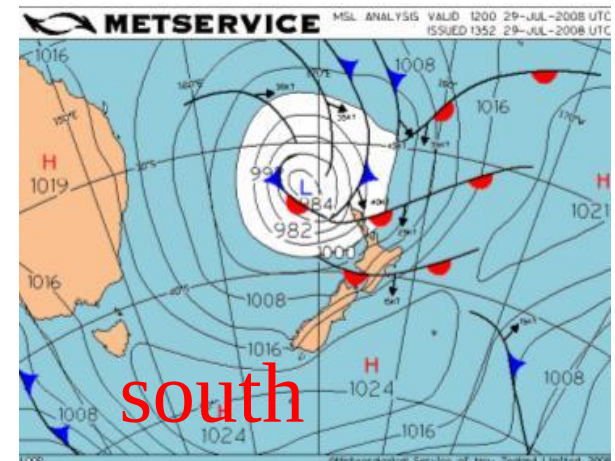
For $RM > 0$,
Left: north
Right: south

For $RM < 0$,
Left: south
Right: north



north

south



Scales and applications

$$z \leftrightarrow \phi$$

$$k \leftrightarrow \lambda^2$$

- $L = 1 \text{ kpc} \rightarrow k = 6 \text{ kpc}^{-1} \rightarrow \bullet = 30 \text{ cm}$
- $L < 0.1 \text{ kpc} \rightarrow k > 60 \text{ kpc}^{-1} \rightarrow \bullet = 1 \text{ m}$
- Assuming $B = 3 \text{ } \odot G$, $n_e = 0.03 \text{ cm}^{-3}$

λ coverage only possible with SKA: 2 cm – 6m

Conclusions

- RM synthesis: measure magnetic helicity
- Need line of sight component: edge-on galaxy
- Expect polarized intensity only in 2 quadrants
- 2 characteristic peaks

