

Towards a Field Theory over Tensor Network States

Andrew G. Green



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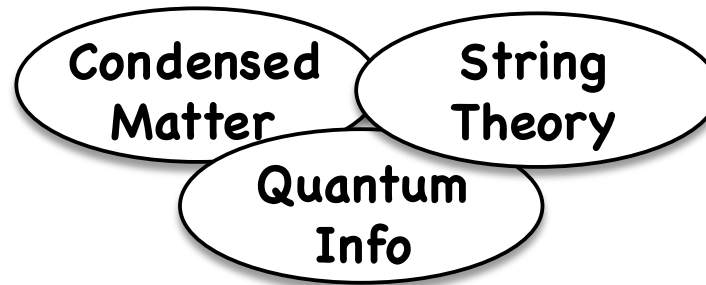
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Background:

- Convergence of ideas

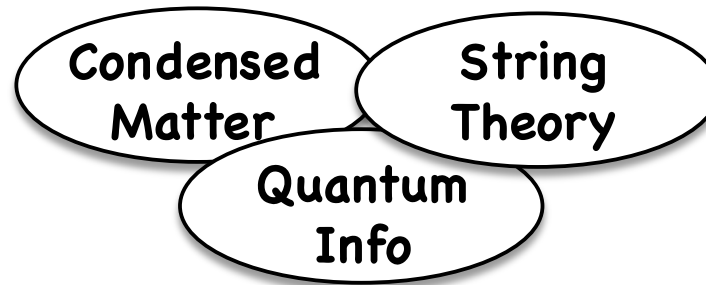


- Suggestive links
 - Holography vs hierarchical tensor networks
 - Is Holography an entanglement ansatz?
 - Can we capture new strongly-correlated phases?
- Ideas formulated in very different language



Background:

- Convergence of ideas



- Suggestive links

Holography vs hierarchical tensor networks

Is Holography an entanglement ansatz?

Need to develop a common Language
Import insights of tensor networks to field theory



Outline:

- Background

Variational States in Condensed Matter
Tensor Networks
RG and Hierarchical Tensor Networks

- Towards a Field Theory Over Tensor Networks

Goal and Key steps
Formulating the Field Theory
Potential Applications
Extensions to Higher Dimensions
Extension to Critical Systems

- Conclusions



I Background:

Variational States in Condensed Matter
Tensor Networks and Strong Correlation
Tensor Networks and AdS/CFT

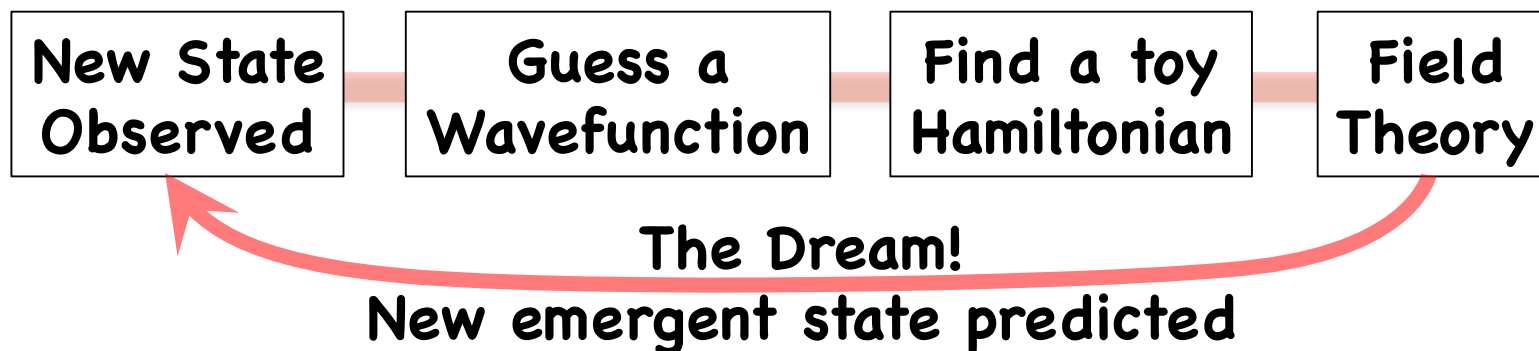


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Role of Variational States in Condensed Matter



Examples:

FQHE: Expt[DPvK]→Exact Diag→Wavefunction[Laughlin]→Hamiltonian [Haldane]
→Composite Fermion Picture[Jain/Read]→Field Theory[Lopez/Fradkin]
BCS: Experiment → Toy Hamiltonian → Guess wavefunction → Field Theory

Exceptions:

Haldane Conjecture → Field Theory → Expt
Topological Insulator → Toy Hamiltonian → wavefunction → experiment
Bethe Ansatz: variational wavefunction → tremendous power in 1d
Conformal Field Theory, Renormalisation Group... etc. etc.

Ab Initio Ideal/Myth

Ab Initio
Hamiltonian

Big
Computer

Experimental
Prediction

The Theory of Everything

$$i\hbar \frac{\partial \Psi}{\partial t} = \hat{H} \Psi$$

$$\hat{H} = - \sum_j^N \frac{\hbar^2}{2m} \nabla_j^2 - \sum_\alpha^M \frac{\hbar^2}{2M_\alpha} \nabla_\alpha^2 - \sum_j^N \sum_\alpha^M \frac{Z_\alpha e^2}{|r_j - R_\alpha|} \\ + \sum_{j < k}^N \frac{e^2}{|r_j - r_k|} + \sum_{\alpha < \beta}^M \frac{Z_\alpha Z_\beta e^2}{|R_\alpha - R_\beta|}$$

* Air	* Steel	* Paper	* Vitamins
* Water	* Plastic	* Dynamite	* Ham Sandwiches
* Fire	* Glass	* Antifreeze	* Ebola Virus
* Rocks	* Wood	* Glue	* Economists
* Cement	* Asphalt	* Dyes	* ...

I Background:

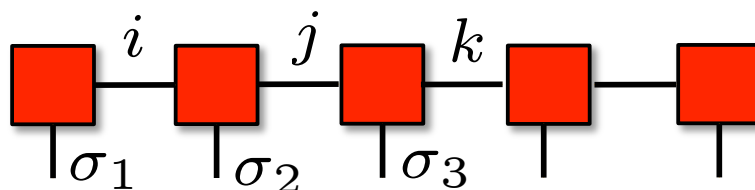
Variational States in Condensed Matter
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Tensor Networks

- Class of variational wavefunctions
- Embody insights about entanglement structure
- Describe groundstates of local Hamiltonians efficiently
- Exact for some model H (AKLT, Majumdar-Ghosh, etc)
- **Matrix Product States (1d tensor network)**

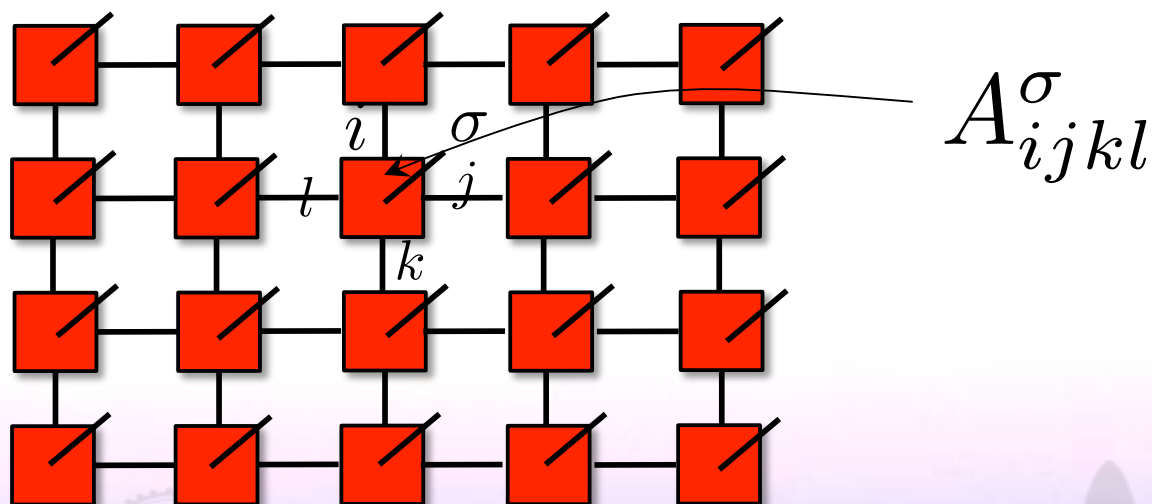
$$|\phi\rangle = \sum_{\{\sigma\}} A_i^{\sigma_1} A_{ij}^{\sigma_2} A_{jk}^{\sigma_3} A_{kl}^{\sigma_4} \dots |\sigma_1, \sigma_2, \sigma_3, \sigma_4, \dots\rangle$$



- MPS is a restricted sum of product states
- MPS dense on Hilbert space

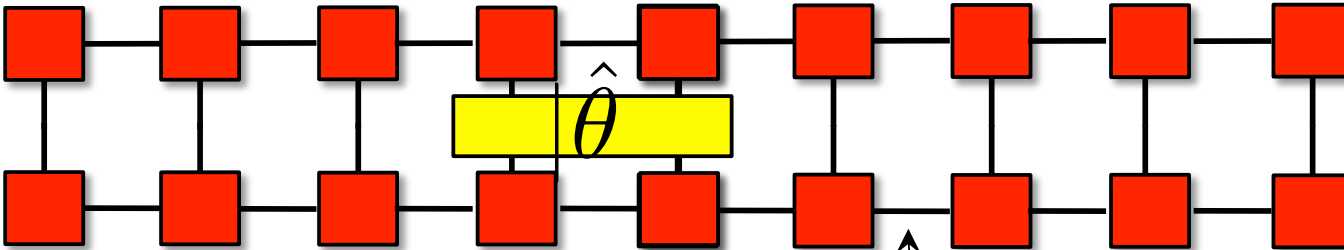
Tensor Networks

- Class of variational wavefunctions
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- Exact for some model H (AKLT, Majumdar-Ghosh, etc)
- PEPS (projected entangled pair state - 2d tensor network)



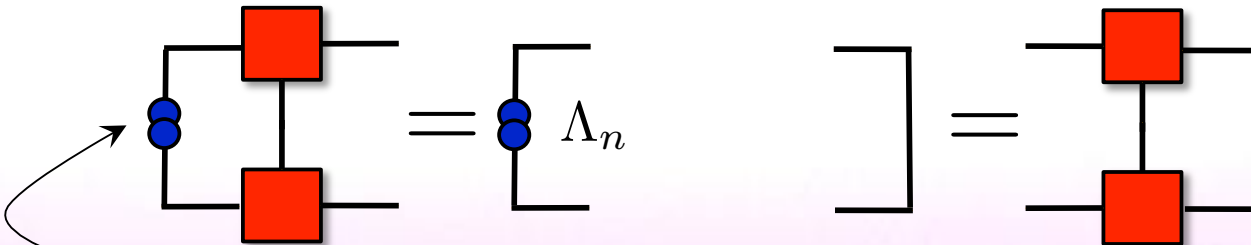
Tensor Networks – Locality and Gauge Fixing

- **Q.** Is the expectation of local operators local in terms of A s?
- **A.** Not in general. In 1d can always gauge fix to make it so.

$$\langle \psi | \hat{\theta} | \psi \rangle =$$


$UU^\dagger = 1$

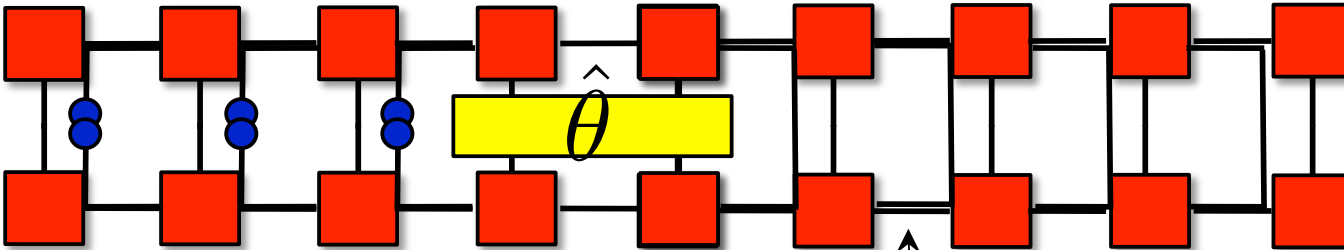
Canonical Form:



$\Lambda_{n-1} = \text{diag}\{\lambda_1, \lambda_2, \dots\}$

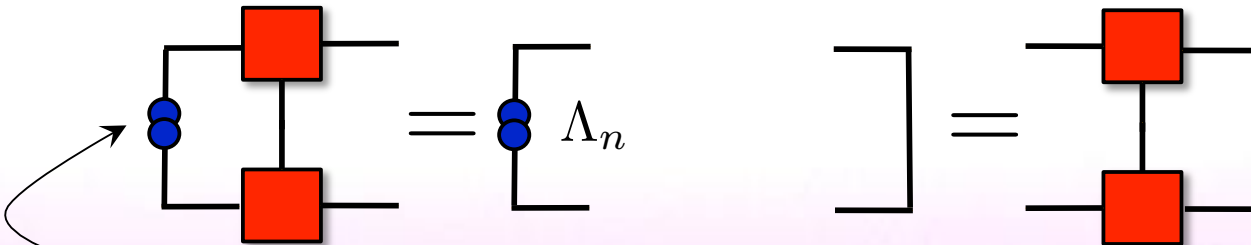
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Tensor Networks – Locality and Gauge Fixing

- **Q.** Is the expectation of local operators local in terms of A s?
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Canonical Form:

- In canonical form, MPS is a Schmidt decomposition

$$|\phi\rangle = \sum_{\alpha} \lambda_{\alpha} |\phi_L^{\alpha}\rangle |\phi_R^{\alpha}\rangle \text{ of each bond [Vidal,PRL91,147902,(2003)]}$$

$$|\phi\rangle = \sum_{\{\sigma\}} A_i^{\sigma_1} A_{ij}^{\sigma_2} A_{jk}^{\sigma_3} A_{kl}^{\sigma_4} \dots |\sigma_1, \sigma_2, \sigma_3, \sigma_4, \dots\rangle$$

vs

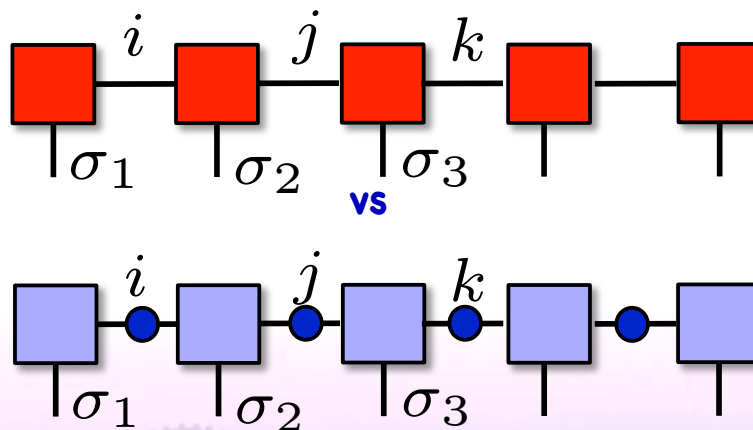
$$|\phi\rangle = \sum_{\{\sigma\}} \Gamma_i^{\sigma_1} \lambda_i^1 \Gamma_{ij}^{\sigma_2} \lambda_i^2 \Gamma_{jk}^{\sigma_3} \lambda_i^3 \Gamma_{kl}^{\sigma_4} \lambda_i^4 \dots |\sigma_1, \sigma_2, \sigma_3, \sigma_4, \dots\rangle$$

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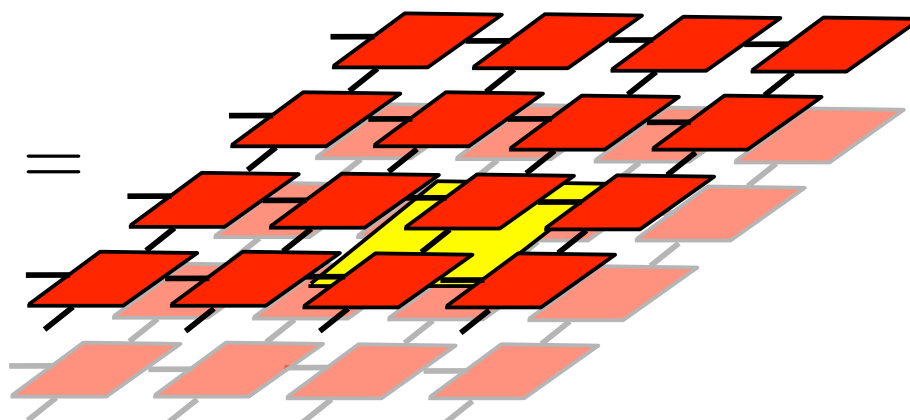
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Tensor Networks – Locality and Gauge Fixing

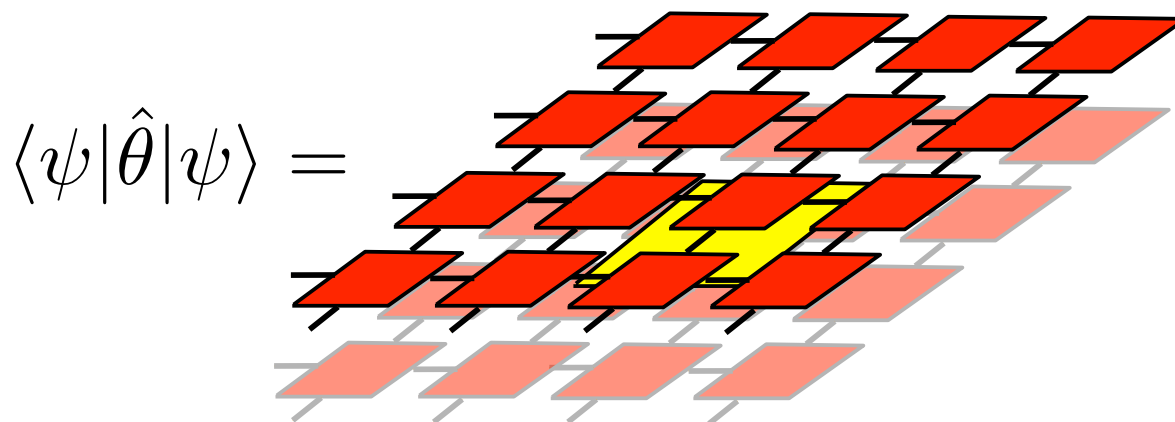
- **Q.** Is the expectation of local operators local in terms of As?
- **A.** Not in general. In 2d (unlike 1d) there is no exact local form

$$\langle \psi | \hat{\theta} | \psi \rangle =$$


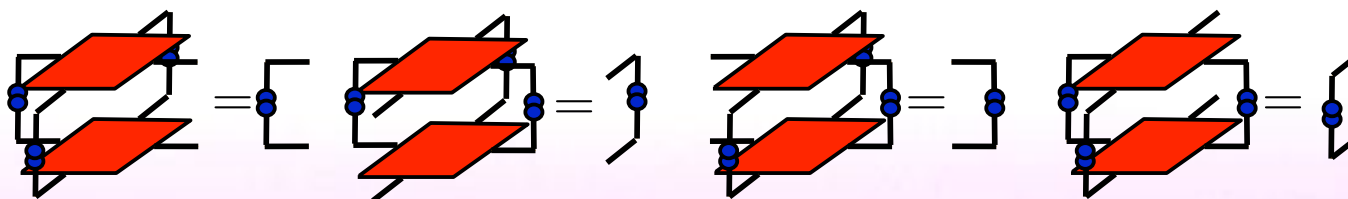
- Physically, expect entanglement to decay, so should be local
- Obey area laws by construction
- Numerically, various approximation schemes seem to work
- Similar to maximally localised Wannier orbitals? [Mazari et al
Rev Mod Phys (2012)]

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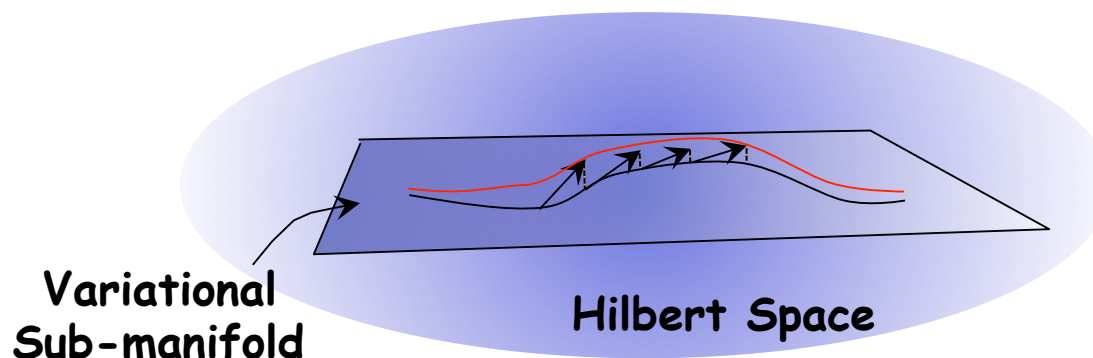


Quasi-Canonical Form:



Time dependence of Tensor Network States

- Fixed bond order – restricted sub-manifold of Hilbert space
- Higher bond order – higher dimension sub-manifold



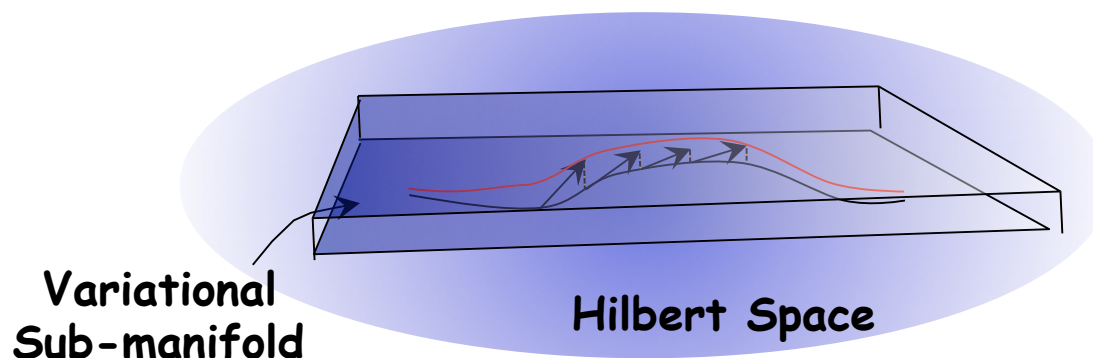
Time Dependent Variational Principle:

- Bond-order grows under Hamiltonian evolution
- Continually Project back to fixed bond order

$$\langle \partial_{\bar{A}_i} \psi | \partial_{A_j} \psi \rangle \dot{A}_j = \langle \partial_{\bar{A}_i} \psi | \hat{\mathcal{H}} | \psi \rangle$$

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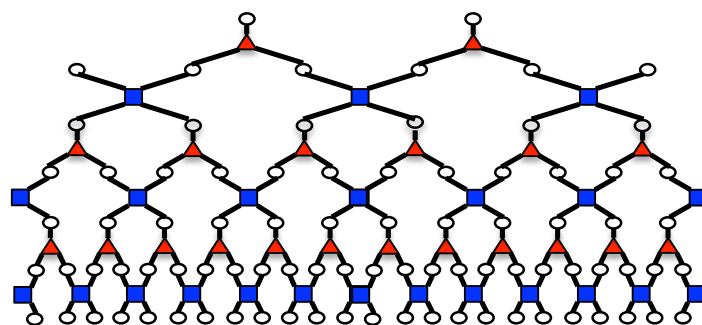


RG and Hierarchical Tensor Networks

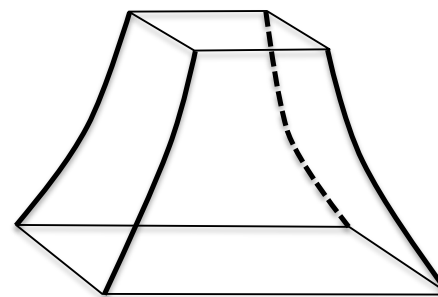
- Critical systems don't obey area laws
- Exponentially large bond order required
- Scaling suggests a more efficient way to encode

MERA (multi-scale entanglement renormalisation ansatz)

[Vidal Phys Rev Lett 101, 110501 (2008)]



VS



Similarity to AdS [Swingle Phys Rev D86, 065007 (2012); ArXiv1209.3304]

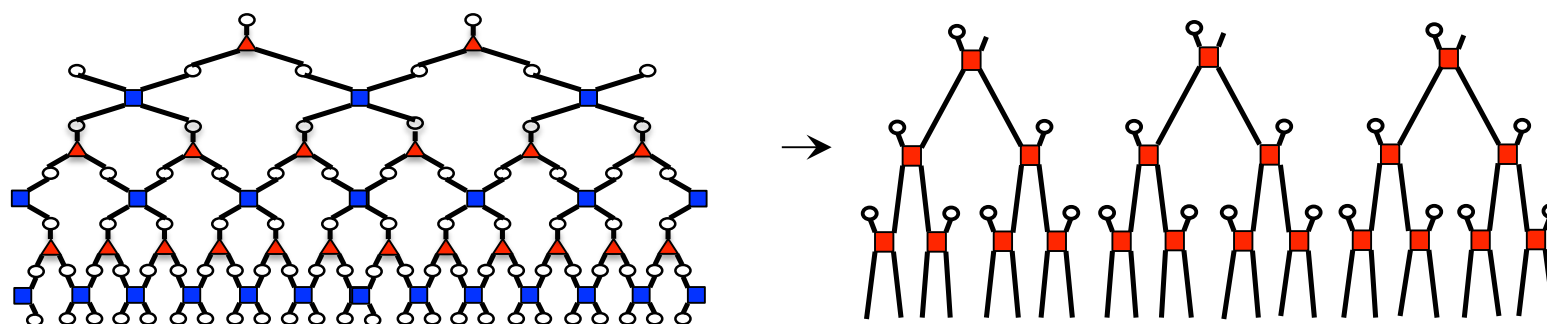
- Extends 1d-2d
- Extra dimension - entanglement RG scale
- Entanglement minimal surface [Ryu, Takayanagi PRL 96, 181602 (2006)]
- Finite T \rightarrow finite extent

RG and Hierarchical Tensor Networks

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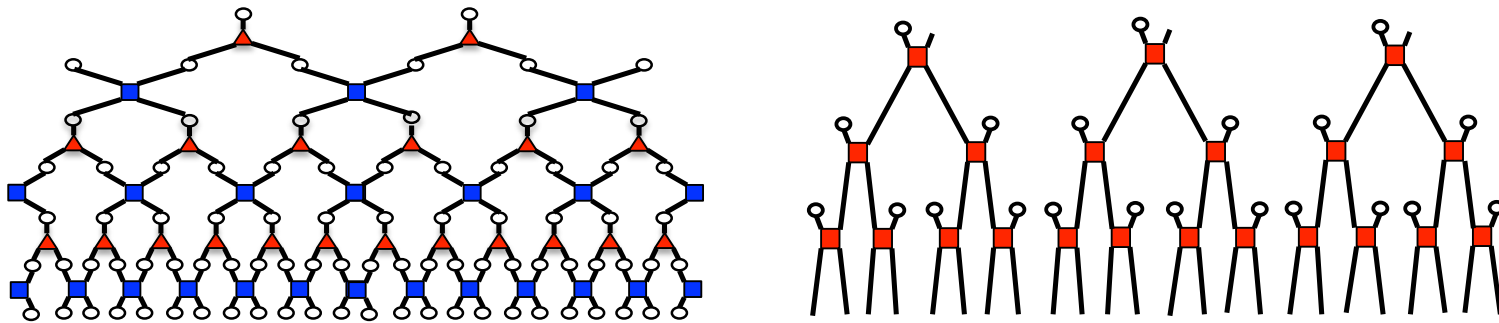
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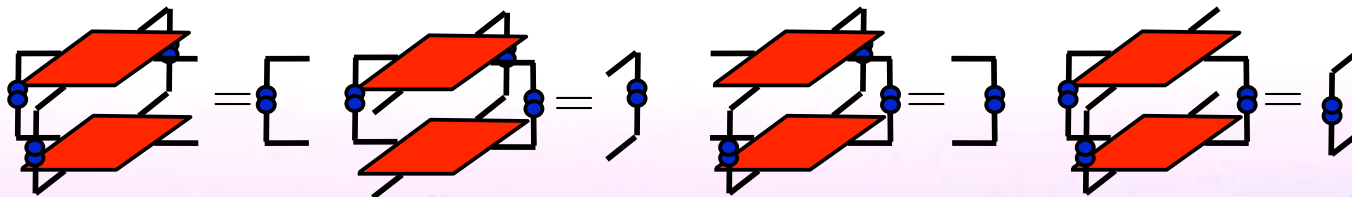
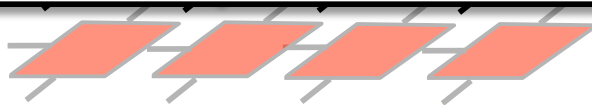
Exact Holographic mapping [Xiao-Liang Qi arXiv:1309.6282]

- Unitary transformation to disk
- Wavelet transform on Cayley tree
- Residual entanglement \rightarrow metric





**Tensor Networks harbour deep insights
Import to field theory**



II Towards a Field Theory over Tensor Network States:

**Goal and Key Steps
Formulating the Field Theory
Special Cases
Potential Applications
Extensions to Higher Dimensions
and Critical Systems**



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Goal

Import insights from tensor networks into a
Field theory over tensor network states

Info

$$\mathcal{Z} = \text{Tr} e^{-\beta \mathcal{H}} = \int [DA] e^{-S[A]}$$

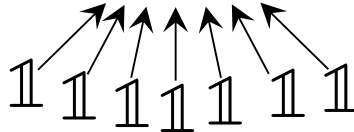
MPS data

Advantages

- Saddle points=> TDVP (time-dependent variational principle)
- Fluctuations - expansion about saddle or increase bond order
- Field theory treatment of gauge freedoms in MPS?
- Extension to higher dimensions?
- Various potential applications...

Key Steps

Import insights from tensor networks into a
Field theory over tensor network states

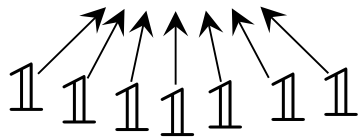
$$\begin{aligned} \mathcal{Z} &= \text{Tr} e^{-\beta \mathcal{H}} \\ &= \int D\psi |\psi\rangle \langle \psi| \int d\tau [\langle \psi | \partial_\tau \psi \rangle - \langle \psi | \hat{\mathcal{H}} | \psi \rangle] \end{aligned}$$


- Insert resolutions of identity over over-complete set
- Usually $|\psi\rangle$ product states
- Can we do the same with matrix product states?

Key Steps

Import insights from tensor networks into a
Field theory over tensor network states

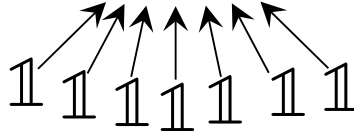
$$\begin{aligned} \mathcal{Z} &= \text{Tr} e^{-\beta \mathcal{H}} \\ &= \int DA e^{\int d\tau [\langle A | \partial_\tau A \rangle - \langle A | \hat{\mathcal{H}} | A \rangle]} \end{aligned}$$



- Insert resolutions of identity over over-complete set
- Usually $|\psi\rangle$ product states
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Key Steps

Import insights from tensor networks into a Field theory over tensor network states

$$\mathcal{Z} = \text{Tr} e^{-\beta \mathcal{H}}$$


$$\mathbb{1} = \int DA |A\rangle \langle A|$$

$$= \int DA e^{\int d\tau [\langle A | \partial_\tau A \rangle - \langle A | \hat{\mathcal{H}} | A \rangle]}$$

- Insert
- Usually
- Can we

Q. What is the Measure?

set

Q. What is the Berry Phase?

Q. Is the theory local?

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and Critical Systems



Gauge Fixing, Locality and General Parameterization

Local Field Theory \Rightarrow Gauge Fix to Canonical Form

- **Canonical Equations**

$$\sum_{\sigma} A^{\sigma\dagger} \Lambda_{n-1} A^{\sigma} = \Lambda_n \quad \sum_{\sigma} A^{\sigma} A^{\sigma\dagger} = 1$$


- **General Parameterization**

$$A^{\sigma} = N^{\sigma} U^{\sigma}, \quad \sum_{\sigma} U^{\sigma\dagger} N^{\sigma\dagger} \Lambda_{n-1} N^{\sigma} U^{\sigma} = \Lambda_n$$

Diagonal matrix
of spin coherent
state spinors


$SU(N)/DU(N)$

Residual canonical equations

Gauge Fixing, Locality and General Parameterization

Local Field Theory \Rightarrow Gauge Fix to Canonical Form

- Canonical Equations

$$\sum_{\sigma} A^{\sigma\dagger} \Lambda_{n-1} A^{\sigma} = \Lambda_n \quad \sum_{\sigma} A^{\sigma} A^{\sigma\dagger} = 1$$


- Bond Order 2 Parameterization

$$N^{\sigma} = \begin{pmatrix} n_1^{\sigma} & 0 \\ 0 & n_2^{\sigma} \end{pmatrix}$$

$$U^{\sigma} = \cos \theta^{\sigma} / 2 + i \boldsymbol{\tau} \cdot \mathbf{u} \sin \theta^{\sigma} / 2$$

$$\mathbf{u} = (\cos \phi, \sin \phi)$$

$$d\Lambda_n = \sum_{\sigma} \Gamma^{\sigma} \cos \theta_{\sigma}$$

$$0 = \sum_{\sigma} \Gamma^{\sigma} \sin \theta_{\sigma}$$

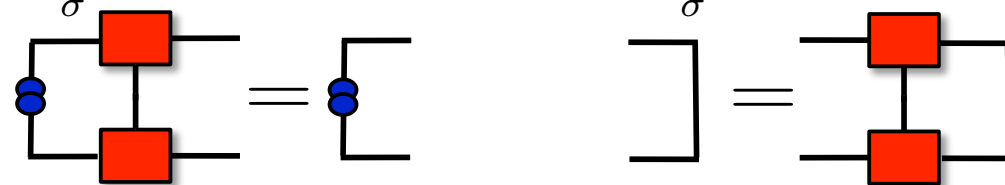
$$\Lambda_n = (\mathbb{1} + \tau_z d\Lambda_n) / 2$$

$$2\Gamma_n = [d\Lambda_{n-1}(1 + \sigma \bar{n}^z) + \sigma \Delta n^z / 2]$$

Gauge Fixing, Locality and General Parameterization

Local Field Theory \Rightarrow Gauge Fix to Canonical Form

- Canonical Equations**

$$\sum_{\sigma} A^{\sigma\dagger} \Lambda_{n-1} A^{\sigma} = \Lambda_n \quad \sum_{\sigma} A^{\sigma} A^{\sigma\dagger} = 1$$


- General Bond Order - split into SU(2) subgroups**

$$N^{\sigma} = \text{diag}(n_1^{\sigma}, n_2^{\sigma}, \dots)$$

$$U^{\sigma} = \prod_l e^{i\theta^{\sigma} \mathbf{v}_l^{\sigma} \cdot \boldsymbol{\tau} / 2}$$

$$d\Lambda_n^i = \sum_{\sigma} \Gamma^{\sigma,j} \left(\prod_l^{\rightarrow} e^{\theta^{\sigma} \mathbf{v}_l^{\sigma} \cdot \mathbf{f}} \right)_{j,i}$$

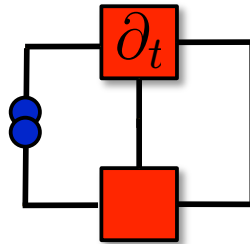
$$\Gamma^{\sigma,j} = \text{Tr} \left[\tau_j \Lambda_{n-1} N^{\sigma\dagger} N^{\sigma} \right] \quad [\tau_i, \tau_j] = 2if_{ijk} \tau_k \quad \Lambda_n = \frac{1 + \sum'_{\beta_d} d\Lambda_n^{\beta_d} \tau_{\beta_d}}{2}$$

Berry Phase

- Contribution from n^{th} site in chain

$$\langle \psi | \partial_t \psi \rangle_n = \sum_{\sigma} \text{Tr} \left[A_n^{\sigma\dagger} \Lambda_{n-1} \partial_t A_n^{\sigma} \right]$$

=



$$= \sum_{\sigma} \text{Tr} \left[U_n^{\sigma\dagger} N_n^{\sigma\dagger} \Lambda_{n-1} \partial_t (N_n^{\sigma} U_n^{\sigma}) \right]$$

$$= \sum_{\sigma} \text{Tr} \left[\Lambda_{n-1} N_n^{\sigma\dagger} \partial_t N_n^{\sigma} \right] + \sum_{\sigma} \text{Tr} \left[U_n^{\sigma\dagger} \left(\Lambda_{n-1} N_n^{\sigma\dagger} N_n^{\sigma} \right) \partial_t U_n^{\sigma} \right]$$

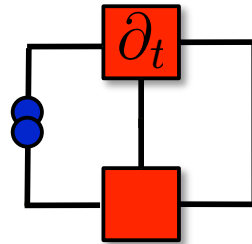


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=



Bond Order 2

$$= \sum_n \left[\sum_{\alpha=1}^{\chi} \Lambda_{n-1}^{\alpha} \langle \mathbf{n}^{\alpha} | \partial_t \mathbf{n}^{\alpha} \rangle + \partial_t \phi (d\Lambda_n - d\Lambda_{n-1}) \right]$$

General

$$= \sum_n \left[\sum_{\alpha=1}^{\chi} \Lambda_{n-1}^{\alpha} \langle \mathbf{n}^{\alpha} | \partial_t \mathbf{n}^{\alpha} \rangle + \sum_{i=1}^{\chi(\chi-1)} \partial_t \phi_i u_{\gamma}^i (d\Lambda_n^{\gamma} - d\Lambda_{n-1}^{\gamma}) \right]$$

γ Labels diag generators
 i Labels $SU(2)$ subgroups

$$\Lambda_n = \frac{1}{2} \left(\mathbb{1} + \sum_{i=1}^{\chi-1} \tau_i d\Lambda_n^i \right) \quad u_i^{\gamma} = \frac{1}{2} \text{Tr}[\tau_{\gamma}^z, \tau_i]$$

Gauge Fixing

- Canonical constraints

$$d\Lambda_n = \sum_{\sigma} \Gamma^{\sigma} \cos \theta_{\sigma} , \quad 0 = \sum_{\sigma} \Gamma^{\sigma} \sin \theta_{\sigma}$$

Iteratively defines $d\Lambda_n$,
(together with def'n of Γ^{σ})

Fixes θ^{\downarrow} , given θ^{\uparrow}

$$d\Lambda_n \equiv d\Lambda_n(\{\mathbf{n}_1, \mathbf{n}_2, \theta^{\uparrow}\}) \quad \theta^{\downarrow} \equiv \theta^{\downarrow}(\{\mathbf{n}_1, \mathbf{n}_2, \theta^{\uparrow}\})$$

- Gauge fixing

$$\mathbb{1} = \int D\Lambda D\theta^{\downarrow} \delta[\theta^{\downarrow} - \theta^{\downarrow}(\{\mathbf{n}_1, \mathbf{n}_2, \theta^{\uparrow}\})] \delta[d\Lambda_n - d\Lambda_n(\{\mathbf{n}_1, \mathbf{n}_2, \theta^{\uparrow}\})]$$

$$D\mathbf{n} D\chi D\theta^{\sigma} D\psi \iff D\mathbf{n} D\chi D\theta^{\uparrow} D\psi \iff D\mathbf{n} D\chi D\Lambda D\psi$$

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Interesting Special Cases

- **Maximally Entangled States** $\Lambda \propto \mathbb{1} \Rightarrow A^\sigma = N^\sigma U$

$$\langle \psi | \partial_t \psi \rangle = \frac{1}{\chi} \sum_n \sum_{\alpha=1}^{\chi} \langle \mathbf{n}^\alpha | \partial_t \mathbf{n}^\alpha \rangle$$

$$\langle \psi | \hat{\mathcal{H}} | \psi \rangle = \left\langle \frac{1}{2} \sum_n J \sigma_n \cdot \sigma_{n+1} \right\rangle = \frac{1}{2} \sum_n J |U_{\alpha,\beta}|^2 \mathbf{n}_n^\alpha \mathbf{n}_{n+1}^\beta$$

- **Spatially Uniform**

$$\langle \psi | \partial_t \psi \rangle = \sum_n \sum_{\alpha=1}^{\chi} \Lambda^\alpha \langle \mathbf{n}^\alpha | \partial_t \mathbf{n}^\alpha \rangle$$

- In both Cases, Effectively χ replicas of system
- glued together with $SU(\chi)$ field
- No intrinsic dynamics for
 - Behaves as fancy Lagrange multiplier (maximally entangled)
 - Inherited through those of N^σ (spatially uniform case)



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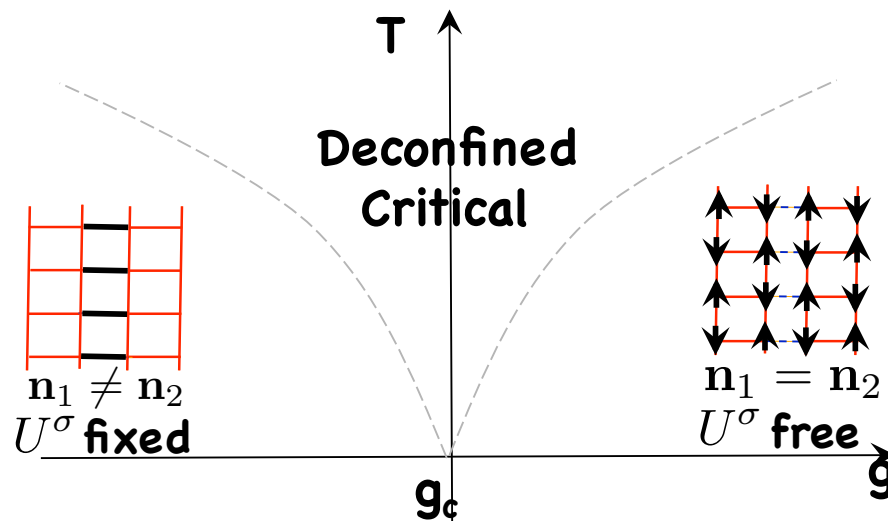
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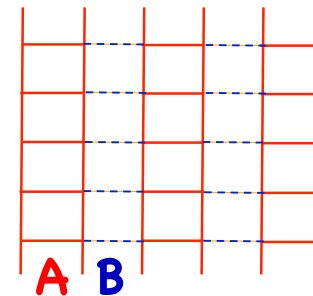
Extensions to Higher Dimensions
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Deconfined Criticality



$$\mathcal{H} = g \sum_{\langle ij \rangle}^A J \hat{\sigma}_i \cdot \hat{\sigma}_j + \sum_{\langle ij \rangle}^B J \hat{\sigma}_i \cdot \hat{\sigma}_j$$



- [Senthil et al, Science 303, 1490 (2004)]
- Critical theory not described by order parameter fluctuations
- Gauge fields/Lagrange multipliers determine critical behaviour
- MPS states may characterize both sides at low bond order
- Certain MPS degrees of freedom soften at transition

Q. Emergent fluctuations in Connection?

Potential Applications

Extended Truncated Wigner Approximation

- Propagates density matrix using saddle-point approx to Keldysh

$$\hat{\rho}(t) = \int D\psi_+ \psi_- \hat{\rho}(0) e^{-iS_+ + iS_-}$$

- All entanglement contained in $\hat{\rho}(0) = \sum_{\alpha} \lambda_{\alpha} |\psi_{\alpha}\rangle \langle \psi_{\alpha}|$
- Propagate $|\psi_{\alpha}\rangle$ independently
- Perturbative corrections increase Schmidt rank [Polkovnikov PRA68, 053604 (2003)]

Alternative

- Propagate with saddle points of MPS Keldysh
 - i. Decompose $\hat{\rho}(0)$ over product states – propagate with MPS saddle point
=>Allows entanglement to grow to some limit
 - ii. Decompose $\hat{\rho}(0)$ over MPS – propagate with MPS saddle point
=>Allows restructuring of entanglement
- Quench action formalism – decompose $\hat{\rho}(0)$ over Bethe states and propagate [Caux and Essler, PRL110, 257203, (2013)]

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Extended Truncated Wigner Approximation

- Propagates density matrix using saddle-point approx to Keldysh

$$\hat{\rho}(t) = \int D\psi_+ \psi_- \hat{\rho}(0) e^{-iS_+ + iS_-}$$

- All entanglement contained in $\hat{\rho}(0) = \sum_{\alpha} \lambda_{\alpha} |\psi_{\alpha}\rangle \langle \psi_{\alpha}|$
- Propagate $|\psi_{\alpha}\rangle$ independently
- Perturbative corrections increase Schmidt rank [Polkovnikov PRA68, 053604 (2003)]

Alternative

- Propagate with saddle points of MPS Keldysh
 - i. Decompose $\hat{\rho}(0)$ over product states – propagate with MPS saddle point
=> Allows entanglement to grow to some limit

ii. D

=>

- Quen
[Ca

Q. Late time hydrodynamics of eigenstate thermalization?

propagate

Potential Applications

Fluctuation corrections to MPS time evolution

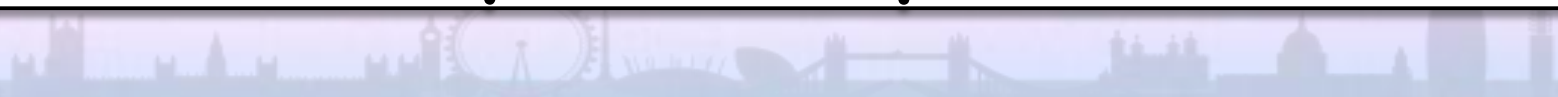
- Two ways to include fluctuations about saddle point
- Re-summing the effects of fluctuations (usual field theory)
- Increasing bond order
- Re-summing fluctuations may improve low bond-order saddle point and fidelity of time evolution over low bond order states

Q. Fluctuation Corrections to MPS?

Many body localization

- Product state field theory – saddle points describe low energy
- Many body localization is a dynamical phase transition through spectrum
- MPS field theory may permit description of dynamics higher in the spectrum.

Q. Field Theory of mid-spectrum states?



II Towards a Field Theory over Tensor Network States:

Goal and Key Steps
Formulating the Field Theory
Special Cases
Potential Applications
**Extensions to Higher Dimensions
and Critical Systems**

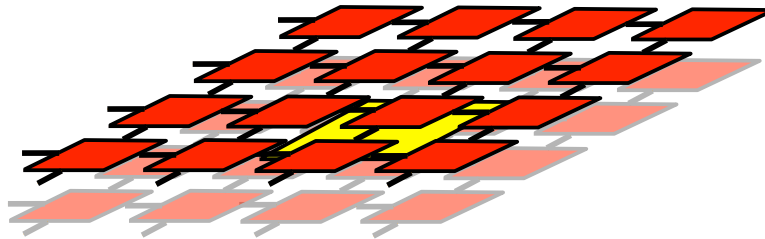


Extension to Higher Dimensions

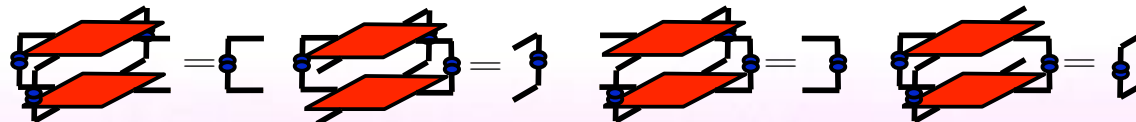
Fundamental difficulty...

- No canonical form in higher dimensions
- Action not local in terms of A
- Decay of entanglement suggests approximately true
- Quasi canonical form
- Connection to multi-band Wannier functions
- [Marzari et al RevModPhys84, 1419 (2012)]?

$$\langle \psi | \hat{\theta} | \psi \rangle =$$

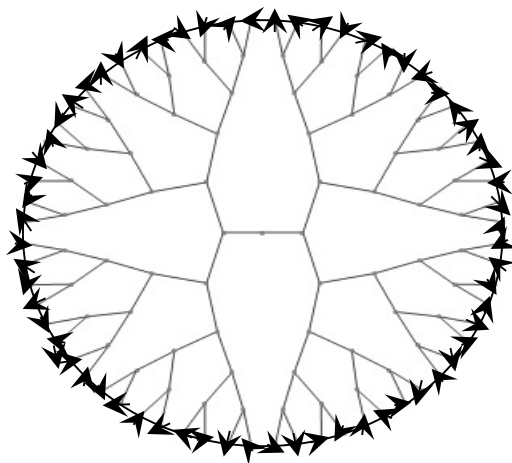


Quasi-Canonical Form:

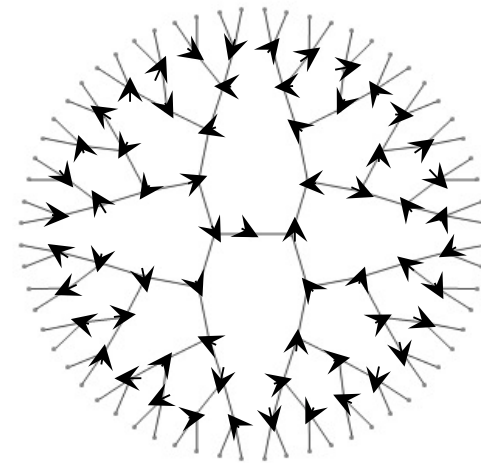


Extensions to Critical Systems

- Various RG Schemes
- MERA (multiscale entanglement renormalization ansatz) [Vidal, PRL99, 220405(2007)]
- TRG (tensor RG) [Verstrate et al, Adv. Phys 57, 143 (2008)]
- SRG (second RG) [Xie et al PRL103, 16069 (2009)]
- HOTRG (higher order TRG) [Xie et al PRB86, 045139 (2012)]
- Exact Holographic Mapping [Xiao-Liang Qi[ArXiv:1309.6282]



Wavelet/RG trans
of boundary $|\psi\rangle, \hat{\mathcal{H}}$
and $\hat{\theta}$ to bulk



- Relation to AdS/CFT [Swingle, Phys. Rev. D 86, 065007 (2012), ArXiv:1209.3304]
- Applying RG to field theory over MPS \rightarrow [S.-S. Lee, NPB 832, 56 (2010); 851, 143 (2011)]

Conclusions

So Far

- Field Theory over Tensor Network States

$$\mathcal{Z} = \text{Tr} e^{-\beta \mathcal{H}} = \int [DA] e^{-S[A]}$$

MPS data



- Imports insights about entanglement to field theory
- Managed for 1d – exploring potential applications

Future

- RG→ describing both renormalization of fluctuations and entanglement structure?
- Extension to higher dimensions and critical theories?
- Links to AdS/CFT?

