

A holographic Kondo model at strong coupling

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Kondo effect

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Screening of a magnetic impurity by
conduction electrons at low temperatures

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conduction electrons at low temperatures

Metals: Fermi liquid + impurities: $\rho \sim \rho_0 + T^2$

In the presence of magnetic impurities: $\rho \sim -\log(T)$

Holographic Kondo Model: Motivation

- Original Kondo problem well-understood in field theory
- Open question: Magnetic impurity coupled to Luttinger liquid
- Here: Realization in gauge/gravity duality
- Describes RG flow, condensation process
- Possible extensions:
Time dependence, Kondo lattices

Holographic Kondo Model

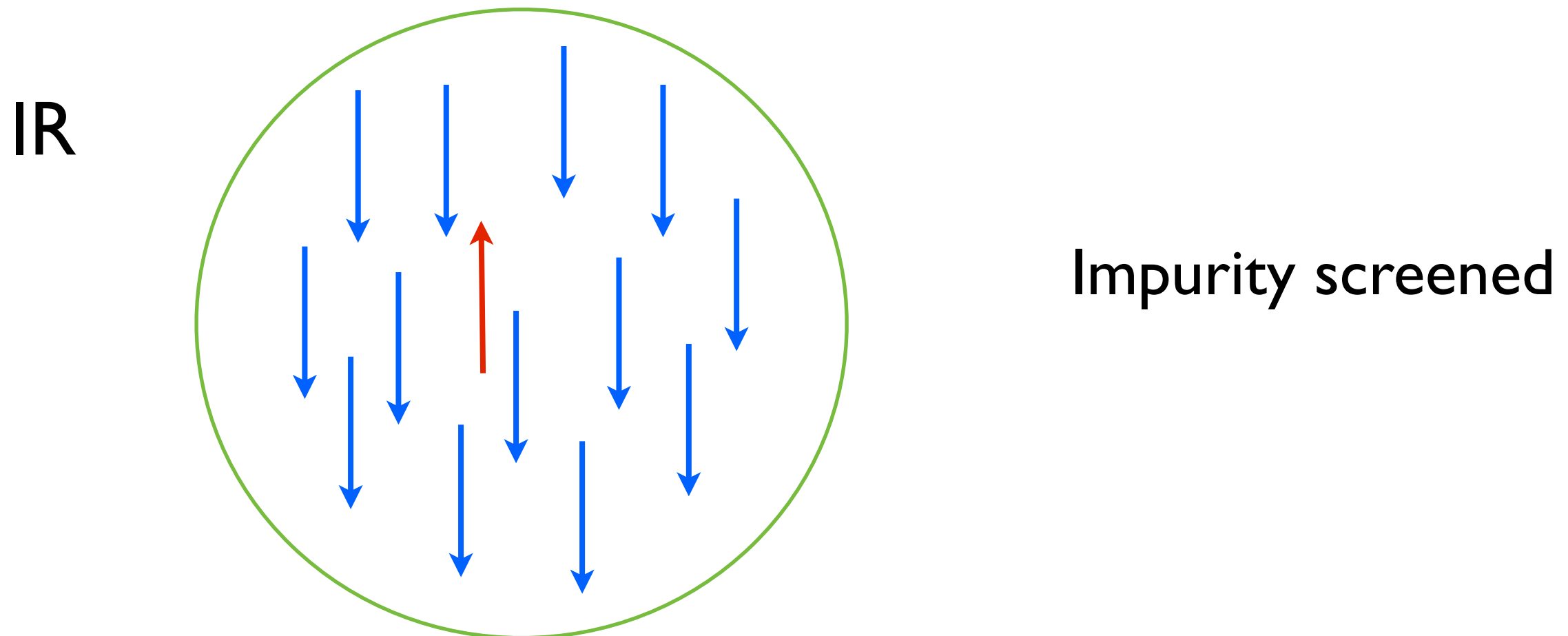
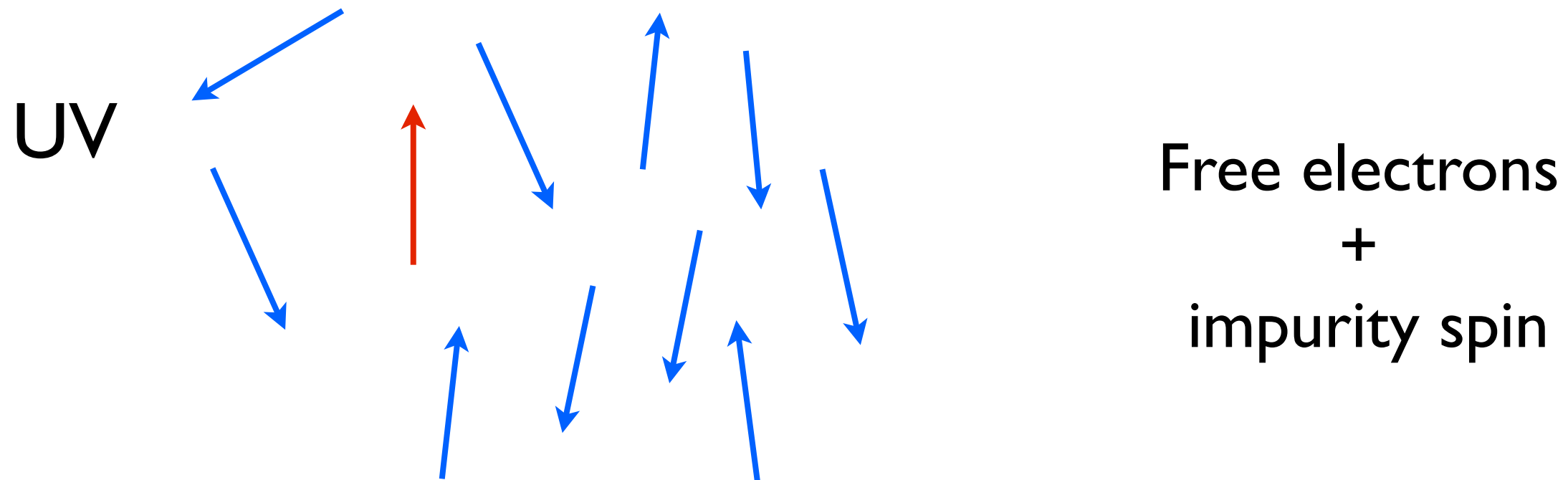
1. Kondo effect: Physics, CFT approach, large N
2. Holographic model: Similarities and differences
3. Time dependence, backreaction

Based on joint work with

C. Hoyos (Tel Aviv Univ.), A. O'Bannon (DAMTP Cambridge),
J. Wu (NCTS Taiwan)

arXiv 1310.3271, published in JHEP

Physics of the Kondo effect



Kondo model

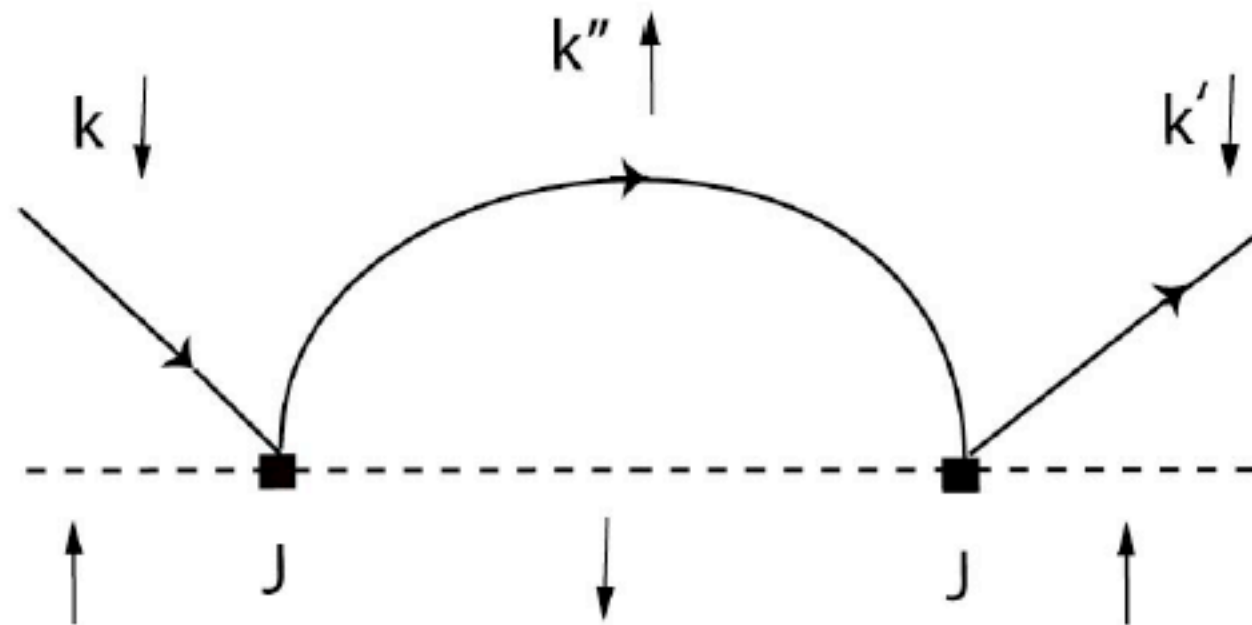
Kondo '64, Affleck+Ludwig '90s

$$H = \frac{v_F}{2\pi} \psi_L^\dagger i \partial_x \psi_L + v_F \lambda_K \delta(x) \vec{S} \cdot \psi_L^\dagger \frac{1}{2} \vec{\tau} \psi_L$$

Free electrons

Local interaction
with magnetic impurity

Scattering with magnetic impurities



$$\rho \sim \rho_0 \left(1 + \kappa \log \frac{T}{|\epsilon - \epsilon_F|} \right)$$

Antiferromagnetic coupling $\kappa < 0$

Logarithmic behaviour at low temperatures

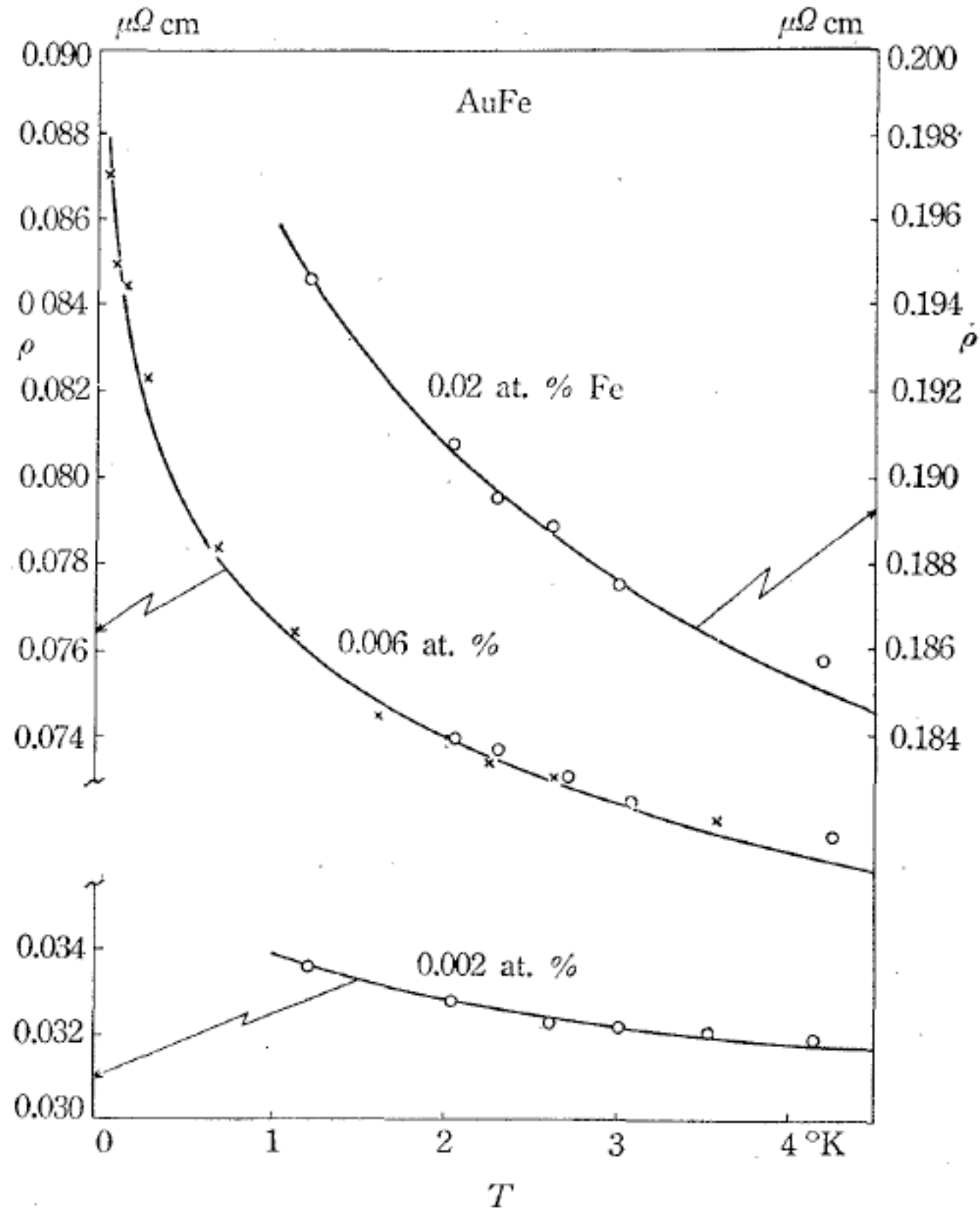


Fig. 1. Comparison of experimental and theoretical ρ - T curves for dilute AuFe alloys.

Jun Kondo:

- Progress of Theoretical Physics
- Volume 32, Issue 1
- Pp. 37-49

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T_K : Kondo temperature

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T_K : Kondo temperature

IR theory is strongly coupled!

$$T_K \sim \Lambda_{\text{QCD}}$$

RG flow

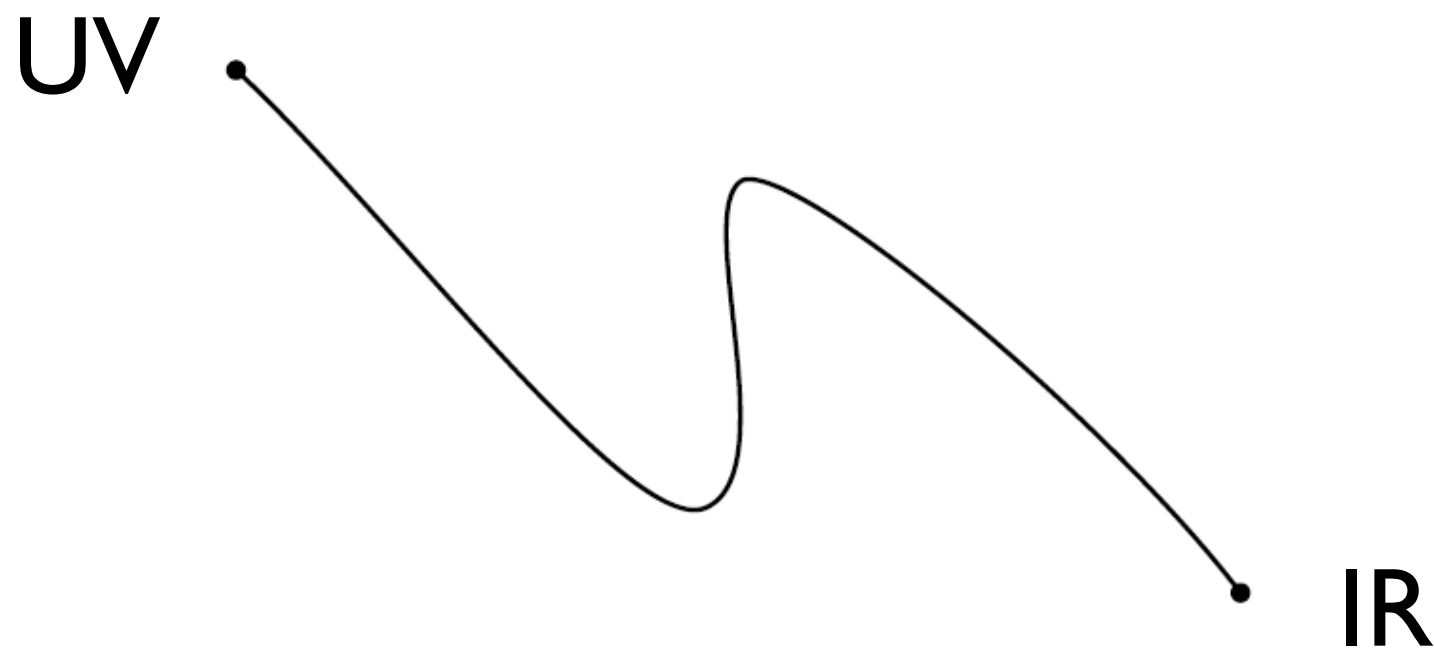
The Kondo model was decisive in the development of the RG formalism. (Wilson)

Negative beta function: $\beta_\lambda \propto -\lambda^2 + \mathcal{O}(\lambda^3)$

UV fixed point: Free theory

Asymptotic freedom

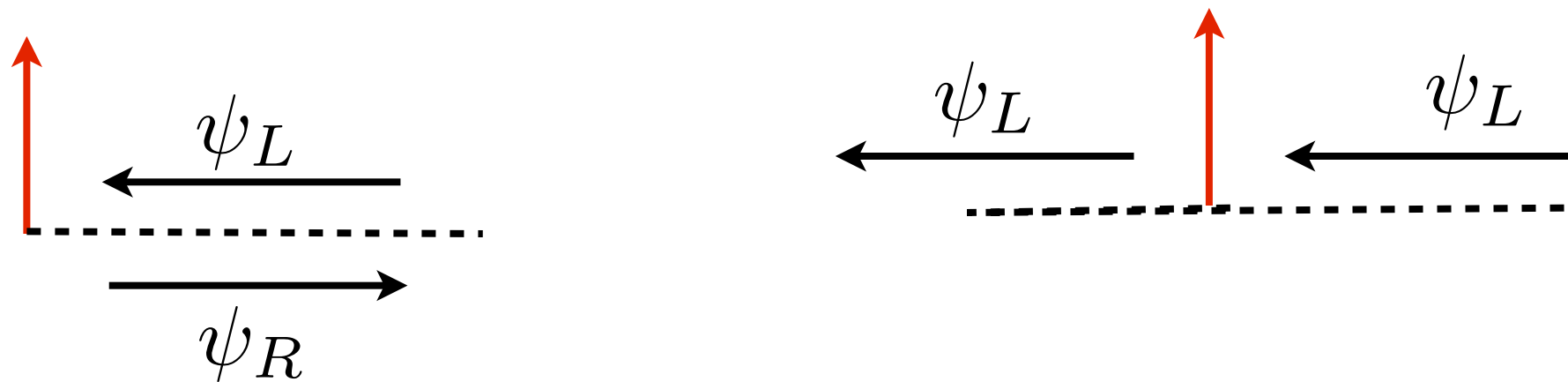
In some cases, the model flows to a strongly coupled IR fixed point.



CFT approach

Affleck+Ludwig, 90's

- One-impurity problem: Spherical symmetry
- Reduces to 1+1-dimensional system with boundary
- Consider only left-movers



CFT approach

Compare CFT's at UV and IR fixed points

UV: Free fermions, boundary condition $\psi_- = \psi_+$

IR: Free fermions, boundary condition $\psi_- = -\psi_+$

Change in boundary condition induces change in spectrum

Generalizations

Generalizations

Spin $SU(N)$, k channels $SU(k)$, Charge $U(1)$

Kac-Moody algebra: $SU(N)_k \times SU(k)_N \times U(1)$

$$[J_n^a, J_m^b] = if^{abc} J_{n+m}^c + k \frac{n}{2} \delta^{ab} \delta_{n,-m}$$

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Sugawara construction:

Separation of spin, channel and charge currents

$$H = \frac{1}{2\pi(N+k)} J^a J^a + \frac{1}{2\pi(k+N)} J^A J^A + \frac{1}{4\pi Nk} J^2 + \lambda_K \delta(x) \vec{S} \cdot \vec{J}$$

Large N: Slave fermions

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Spin of impurity: Young tableau with Q boxes

$S^a = \chi^\dagger T^a \chi$ Totally antisymmetric representation

$$\chi^\dagger \chi = q, \quad Q = q/N$$

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Kondo coupling as double-trace deformation:

$$\begin{aligned} \lambda_K \delta(x) J^a S^a &= \lambda_K \delta(x) \left(\psi_L^\dagger T^a \psi_L \right) \left(\chi^\dagger T^a \chi \right) \\ &= \frac{1}{2} \lambda_K \delta(x) \left(\psi_L^\dagger \chi \right) \left(\chi^\dagger \psi_L \right) \\ &= \frac{1}{2} \lambda_K \delta(x) \mathcal{O} \mathcal{O}^\dagger \end{aligned}$$

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\mathcal{O} $SU(N)$ singlet, charged under $U(N_f) \times SU(k) \times U(1)$

$$[\mathcal{O}] = 1/2$$

Large N: Slave fermions

Sachdev, Senthil, Vojta cond-mat/0209144

Kondo effect corresponds to condensation of $\mathcal{O} = \psi_L^\dagger \chi$

Mean field transition:

$$T > T_K, \quad \langle \mathcal{O} \rangle = 0, \quad SU(k) \times U(N_f) \times U(1)$$

$$T < T_K, \quad \langle \mathcal{O} \rangle \neq 0, \quad SU(k) \times U(N_f) \times U(1) \rightarrow U(1)_D$$

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Gravity dual for well-understood RG flow

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Note however: Holographic model and standard Kondo model have significant differences

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Extensions: Kondo lattices, time dependence, entanglement entropy

Impurities in string holographic models

Supersymmetric defects with localized fermions

- D5/D3 $AdS_2 \subset AdS_5$
[Kachru, Karch, Yaida] [Harrison, Kachru, Torroba]
- M2/D2 in ABJM $AdS_2 \subset AdS_4$
[Jensen, Kachru, Karch, Polchinski, Silverstein]
- D6 in ABJM $AdS_2 \subset AdS_4$, [Benincasa, Ramallo]
with backreaction [Itsios, Sfetsos, Zoakos]
- $D(8 - p)$ in Dp background $S^{7-p} \subset S^{8-p}$ [Benincasa, Ramallo]
other sphere wrappings [Karaiskos, Sfetsos, Tsatis]
- Spectrum of Wilson loops
[Mueck] [Faraggi, Pando Zayas] [Faraggi, Mueck, Pando Zayas]

New in our model:

- Model for entire RG flow
- Double-trace deformation
- Kondo temperature arises naturally
- Impurity screening
- Phase shift
- Power-law scalings at low T

Top-down probe brane model

- based on D7- and D5-probe branes in D3-brane background

	0	1	2	3	4	5	6	7	8	9
N_c D3	X	X	X	X						
N_7 D7	X	X			X	X	X	X	X	X
N_5 D5	X				X	X	X	X	X	

Defect theory

Harvey and Royston 0709.1482, 0804.2854
Buchbinder, Gomis, Passerini 0710.5170

	0	1	2	3	4	5	6	7	8	9
N_c D3	X	X	X	X						
N_7 D7	X	X			X	X	X	X	X	X

N_7 (1+1)-dimensional chiral fermions ψ_L

$$S_{3-7} = \int dx^+ dx^- \psi_L^\dagger (i\partial_- - A_-) \psi_L$$

Preserves 1/4 of SUSY

Kac-Moody algebra

$$SU(N_c)_{N_7} \times SU(N_7)_{N_c} \times U(1)_{N_c N_7}$$

D5-brane probes

D5-branes: Impurity

	0	1	2	3	4	5	6	7	8	9
N_c D3	X	X	X	X						
N_5 D5	X				X	X	X	X	X	

Skenderis, Taylor hep-th/0204054

Camino, Paredes, Ramallo hep-th/0104082

Gomis and Passerini hep-th/0604007

(0+1)-dimensional fermions χ

$$\chi^\dagger \chi = q$$

Kondo interaction: Complex scalar

	0	1	2	3	4	5	6	7	8	9
N_5 D5	X				X	X	X	X	X	
N_7 D7	X	X			X	X	X	X	X	X

Dual operator: $\mathcal{O} \equiv \psi_L^\dagger \chi$

TACHYON

$$m_{\text{tachyon}}^2 = -\frac{1}{4\alpha'}$$

D5 becomes magnetic flux in the D7

Holography - Top-down model for Kondo

	x^0	x^1	x^2	x^3	x^4	x^5	x^6	x^7	x^8	x^9
N_c D3	●	●	●	●	—	—	—	—	—	—
N_7 D7	●	●	—	—	●	●	●	●	●	●
N_5 D5	●	—	—	—	●	●	●	●	●	—

- 3-7 strings = chiral fermions (current algebra)
- 3-5 strings = slave fermions
- 5-7 strings = bifundamental scalar (tachyon)

Near-horizon limit

D3: $\text{AdS}_5 \times S^5$

D7: $\text{AdS}_3 \times S^5 \longrightarrow \text{CS } A_\mu \text{ dual to } J^\mu = \psi^\dagger \sigma^\mu \psi$

D5: $\text{AdS}_2 \times S^4 \longrightarrow \left\{ \begin{array}{l} \text{YM } a_t \text{ dual to } \chi^\dagger \chi = q \\ \text{Scalar } \Phi \text{ dual to } \psi^\dagger \chi \end{array} \right.$

Bottom-up model

Action

$$S = S_{CS} + S_{AdS_2},$$

$$S_{CS} = -\frac{N}{4\pi} \int \text{tr} \left(A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right),$$

$$S_{AdS_2} = -N \int d^3x \delta(x) \sqrt{-g} \left[\frac{1}{4} \text{tr} f^{mn} f_{mn} + g^{mn} (D_m \Phi)^\dagger D_n \Phi + V(\Phi^\dagger \Phi) \right]$$

BTZ black hole

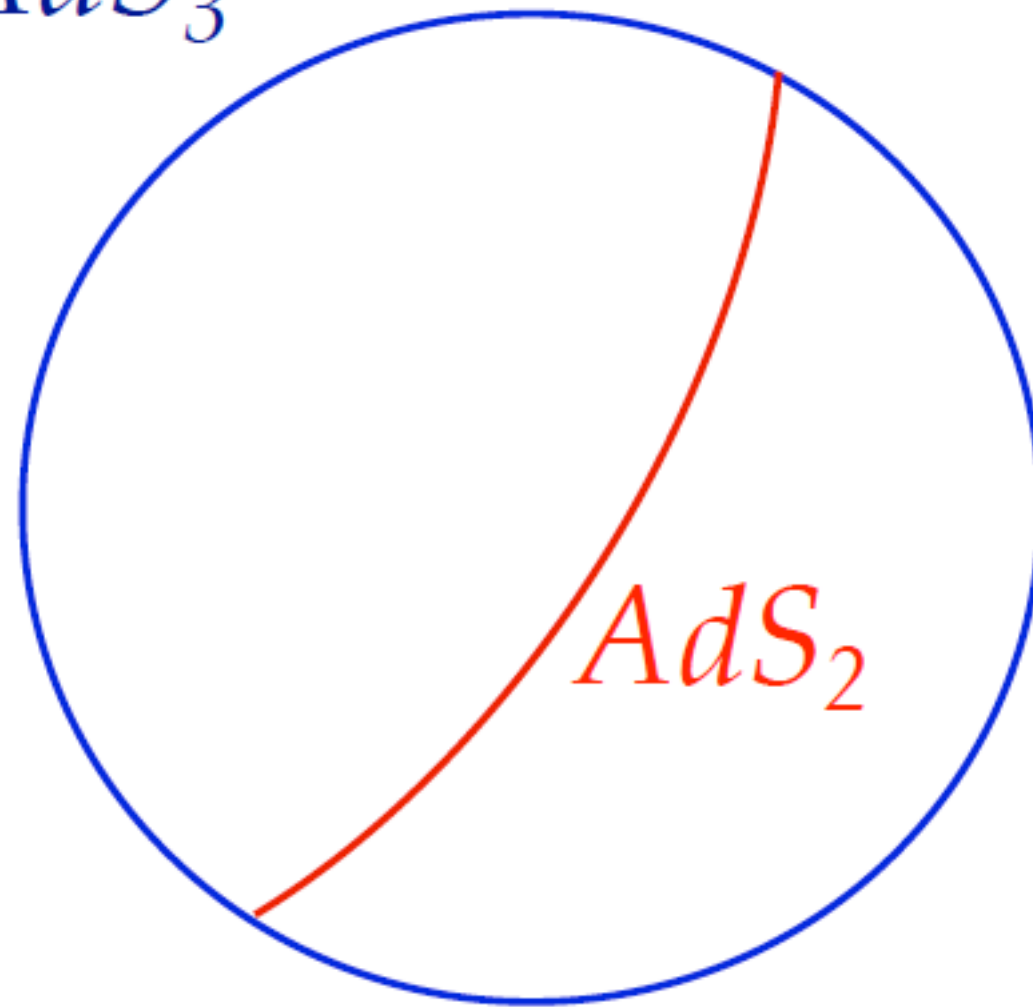
$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu = \frac{1}{z^2} \left(\frac{dz^2}{h(z)} - h(z) dt^2 + dx^2 \right)$$

$$h(z) = 1 - z^2/z_H^2$$

$$T = 1/(2\pi z_H)$$

Defect space

AdS_3



Bottom-up model

Spin group $SU(N)$ gauged

$U(k)_N$ Chern-Simons field dual to channel $SU(k)_N$, charge $U(1)$ current

$k=1$

Defect Yang-Mills field encodes impurity spin representation, $a_t(z) = \frac{Q}{z} + \mu$

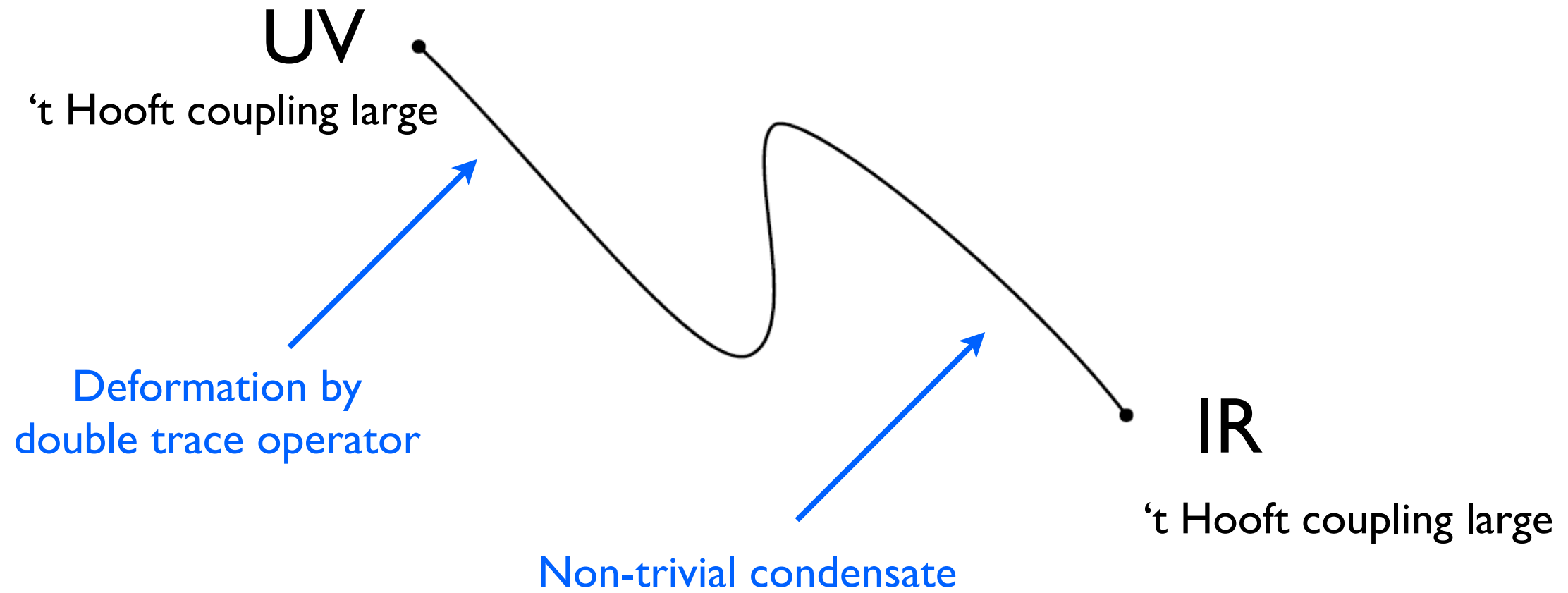
Φ complex scalar bifundamental under the two gauge fields, dual to $\mathcal{O} = \psi_L^\dagger \chi$

Potential $V(\Phi) = m^2 \Phi^\dagger \Phi$, mass at Breitenlohner-Freedman bound

Double-trace operator marginal

$$\Phi = z^{1/2}(\alpha \log(z) + \beta)$$

RG flow



Double trace deformation by $\mathcal{O}\mathcal{O}^\dagger$

$$\Phi = z^{1/2}(\alpha \log(z) + \beta)$$

Witten hep-th/0112258

$$\alpha = \kappa\beta$$

Renormalization:

$$\Phi = z^{1/2}\beta_0(\kappa_0 \log(\Lambda z) + 1) = z^{1/2}\beta(\kappa \log(\mu z) + 1)$$

Running coupling

$$\kappa = \frac{\kappa_0}{1 + \kappa_0 \ln\left(\frac{\Lambda}{\mu}\right)}$$

Dynamical scale: $\Lambda_K = \Lambda e^{1/\kappa_0}$

Kondo coupling

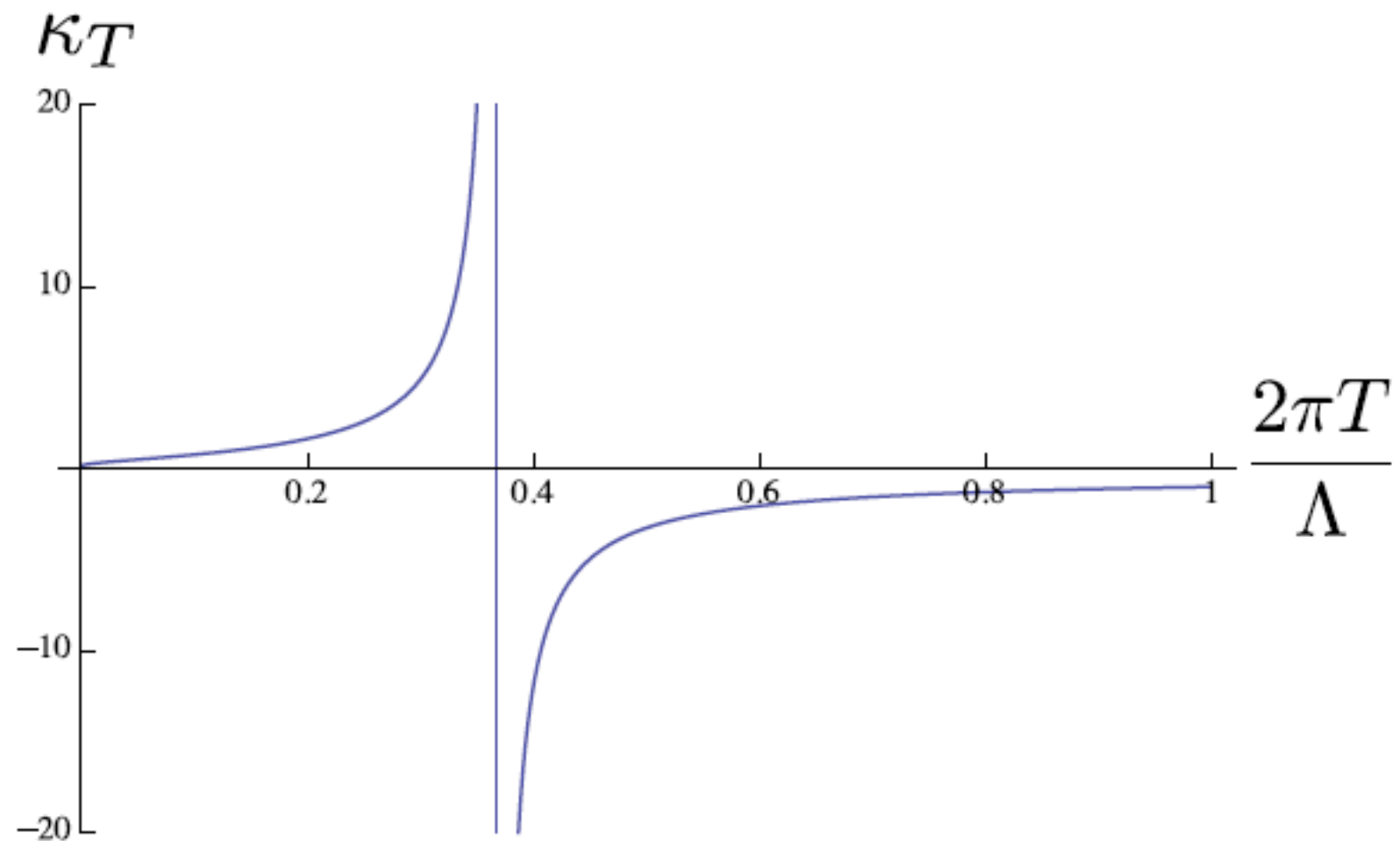
- Finite temperature solution:

$$\Phi = (z/z_H)^{1/2} \beta_T (\kappa_T \log(z/z_H) + 1) = \beta_0 (\kappa_0 \log(\Lambda z) + 1)$$

- Temperature-dependent coupling

$$\kappa_T = \frac{\kappa_0}{1 + \kappa_0 \ln \left(\frac{\Lambda}{2\pi T} \right)}$$

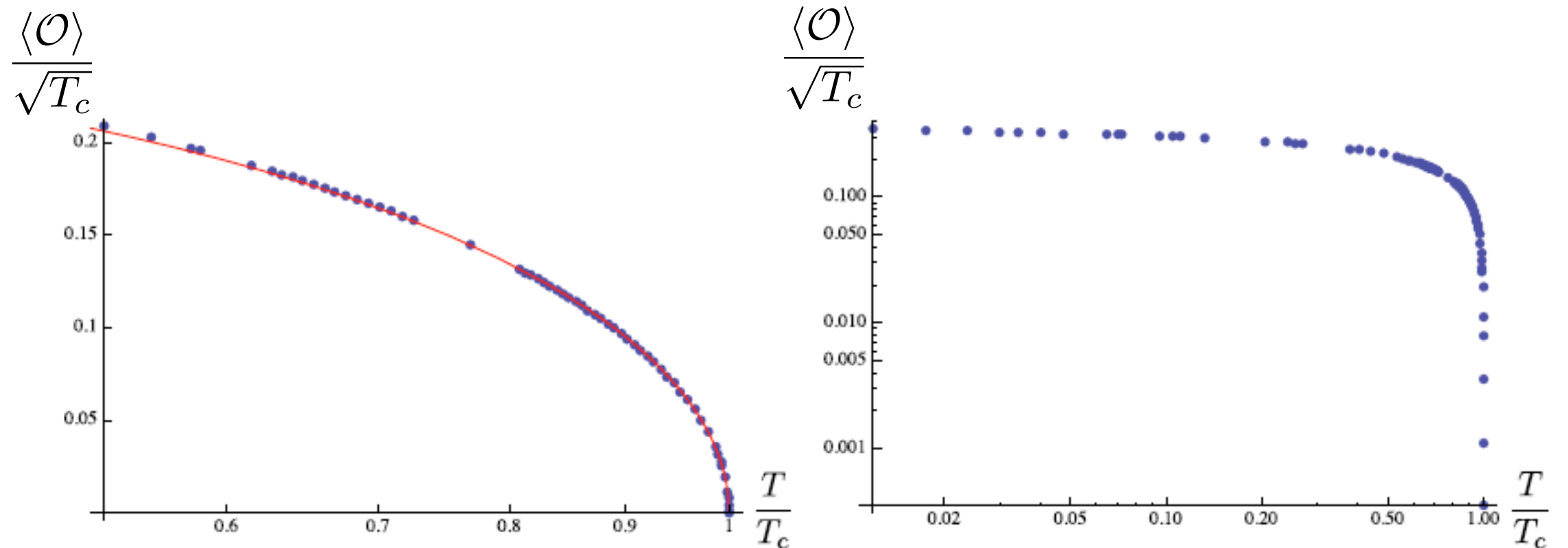
Scale generation



Divergence of Kondo coupling determines Kondo temperature

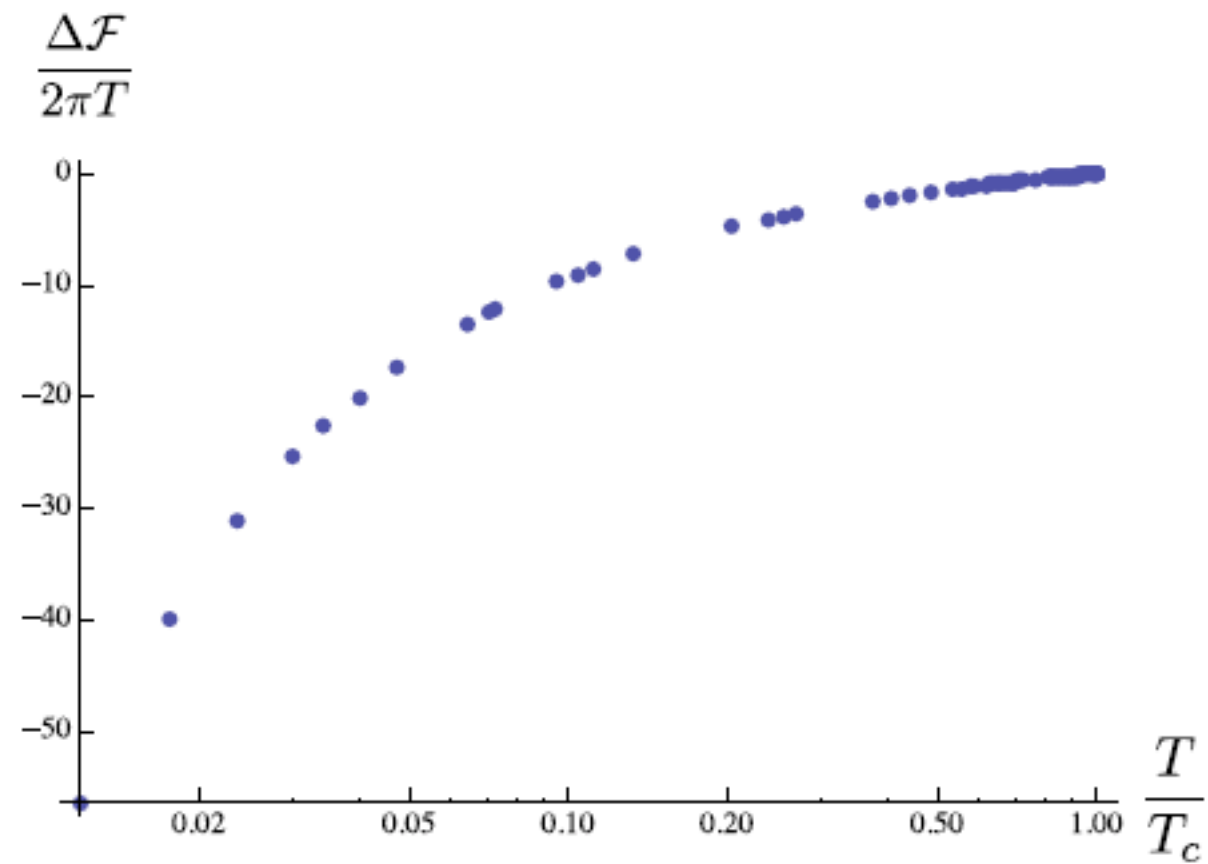
Below this temperature, scalar condenses

Condensate



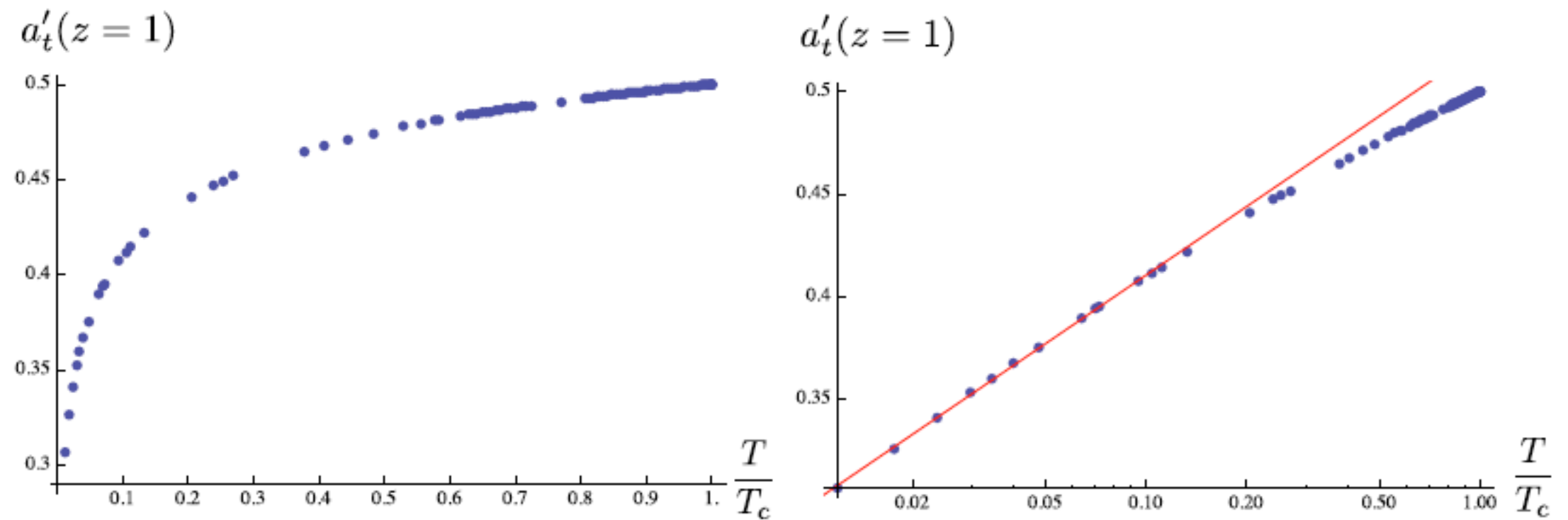
Mean field transition
 $\langle \mathcal{O} \rangle$ approaches constant for $T \rightarrow 0$

Free energy



Phase with scalar condensate more stable below critical temperature

Electric flux at horizon



$$\sqrt{-g} f^{tr} \Big|_{\partial AdS_2} = q = \chi^\dagger \chi$$

Impurity is screened

Phase shift

Below T_c , scalar transfers electric flux from 2d YM to 3d CS field

Wilson loop for 3d gauge field: $W(z) \equiv \oint dx A_x(z)$

Leads to phase shift e^{iW} for chiral fermions

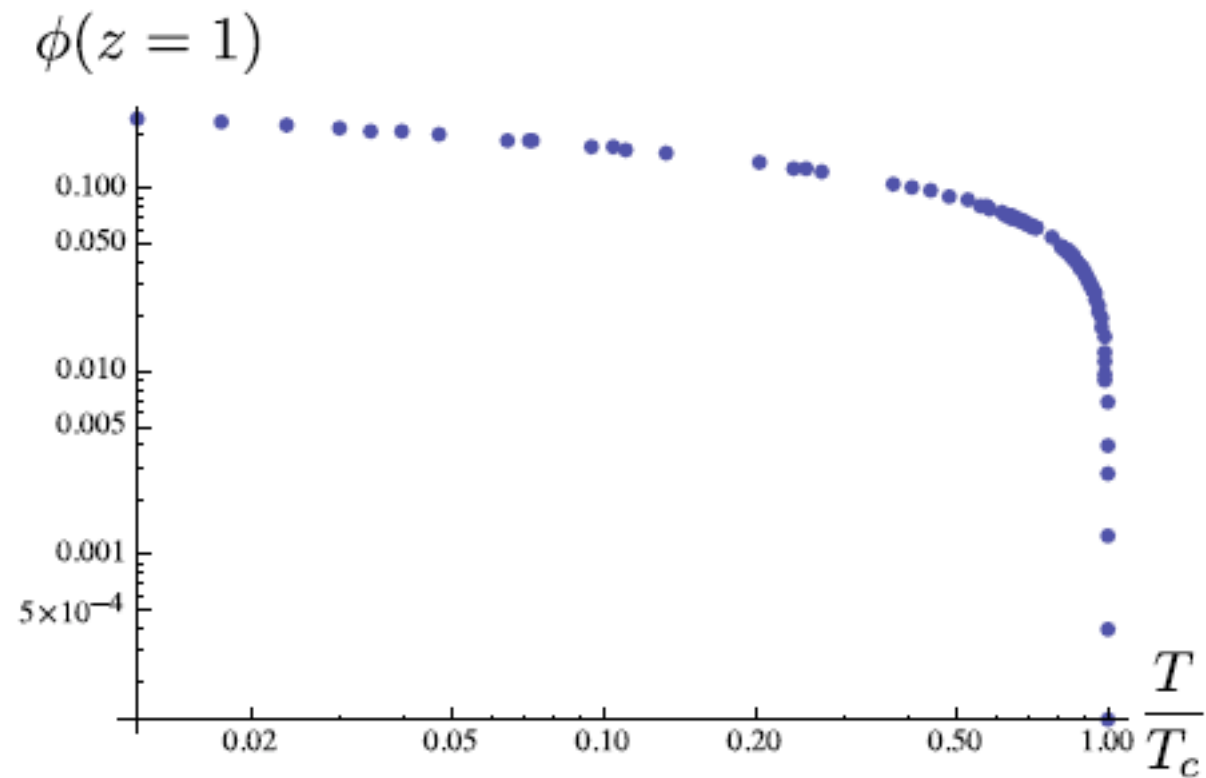
Resistivity from leading irrelevant operator

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No log behaviour due to strong coupling

Resistivity from leading irrelevant operator

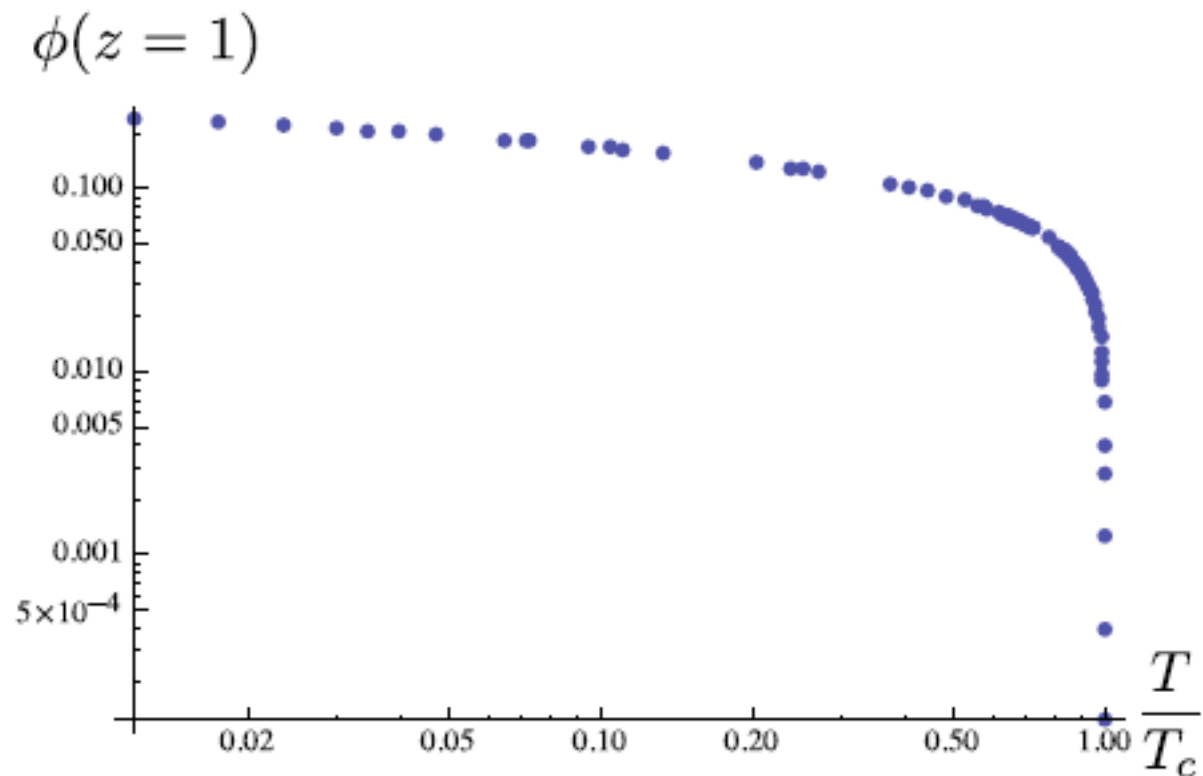
No log behaviour due to strong coupling



$$\phi_0(z=1) = \phi_\infty$$

Resistivity from leading irrelevant operator

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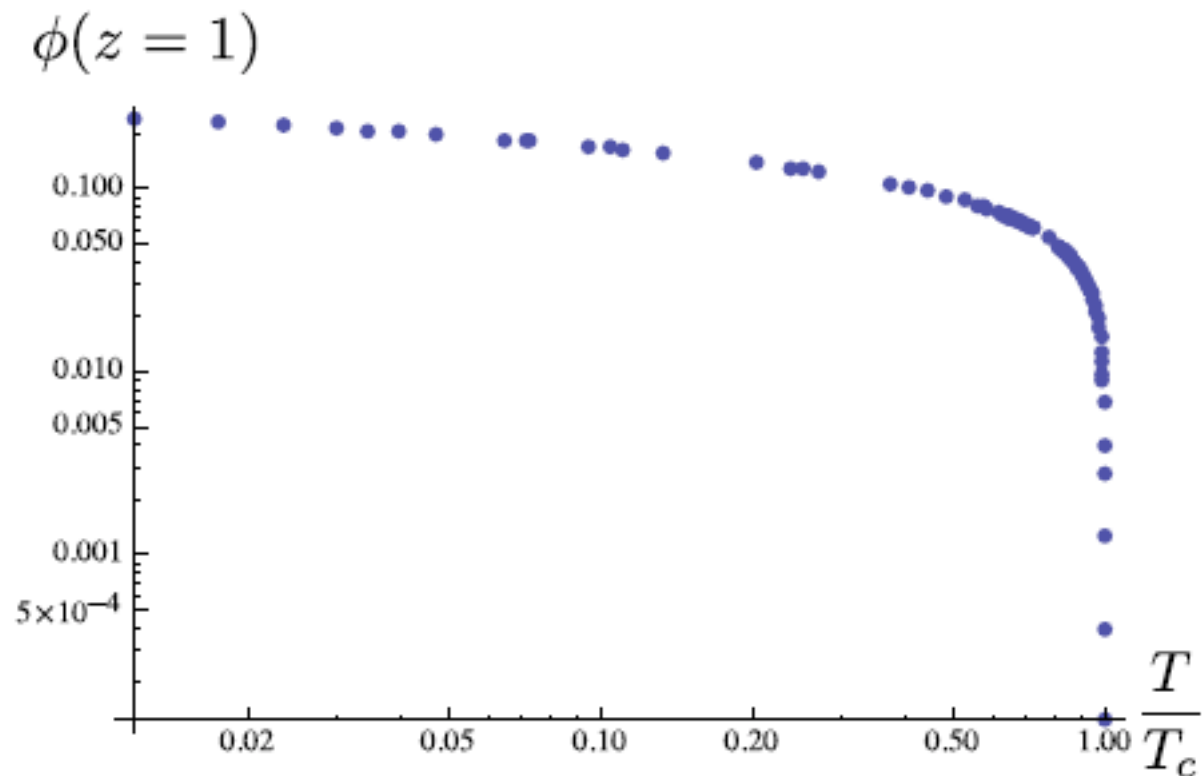


$$\phi_0(z=1) = \phi_\infty$$

IR fixed point stable: Flow near fixed point governed by operator dual to a_t

Resistivity from leading irrelevant operator

No log behaviour due to strong coupling



$$\phi_0(z=1) = \phi_\infty$$

IR fixed point stable: Flow near fixed point governed by operator dual to a_t

Dimension

$$\Delta_+ = \frac{1}{2} + \sqrt{\frac{1}{4} + 2\phi_\infty^2}$$

Resistivity from leading irrelevant operator

Entropy density $s = s_0 + c_s \lambda_{\mathcal{O}}^2 T^{-2+2\Delta_+}$

Resistivity $\rho = \rho_0 + c_+ \lambda_{\mathcal{O}}^2 T^{-1+2\Delta_+}$

Time dependence

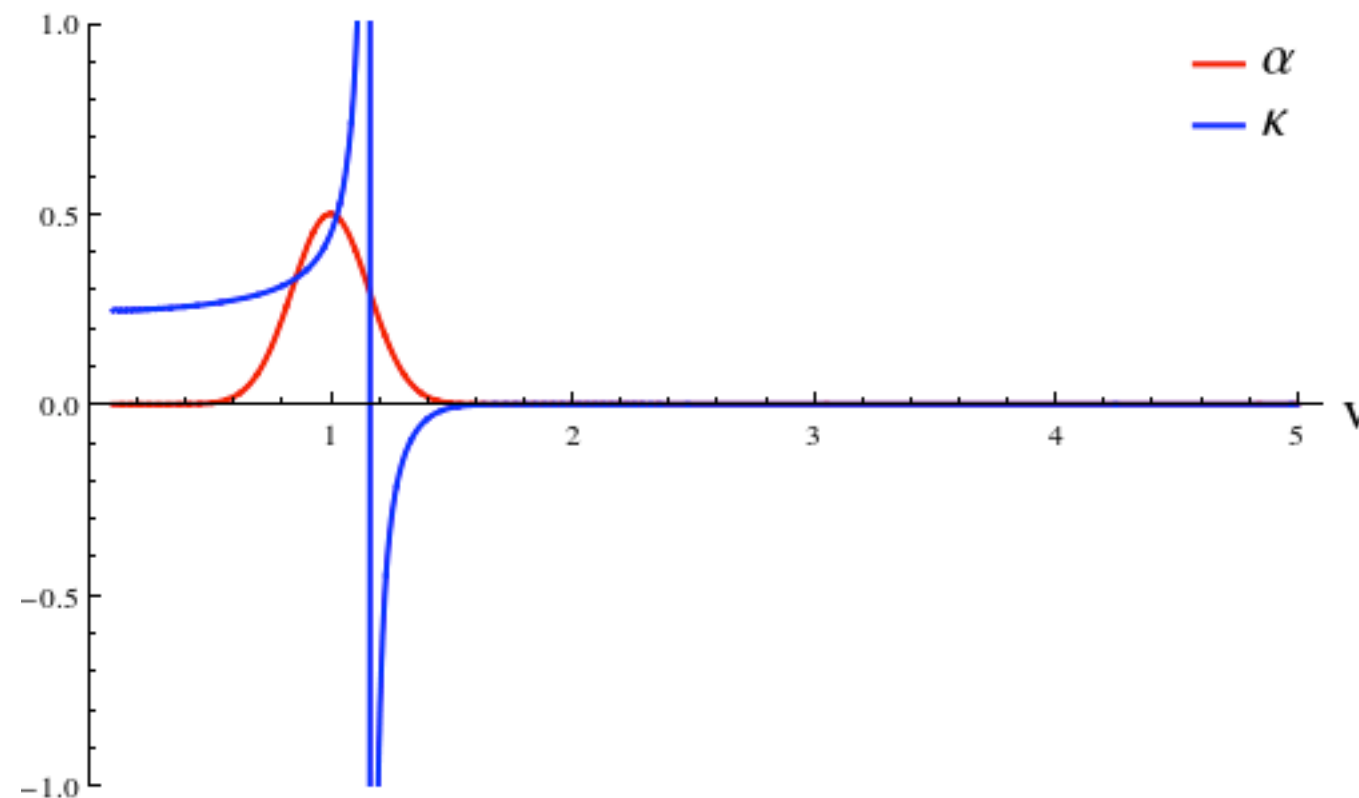
J.E., Flory, Newrzella, Strydom, Wu in progress

Look for time-dependent solutions of the equations of motion

$$A_x(z, x, t), a_t(z, t), a_z(z, t), \phi(z, t), \psi(z, t) \neq 0$$

modelling the evolution of the system
after turning on the Kondo interaction at $t = t_0$

T_c : Kondo coupling generating a Gaussian condensate pulse



J.E., Flory, Newrzella, Strydom, Wu (preliminary)

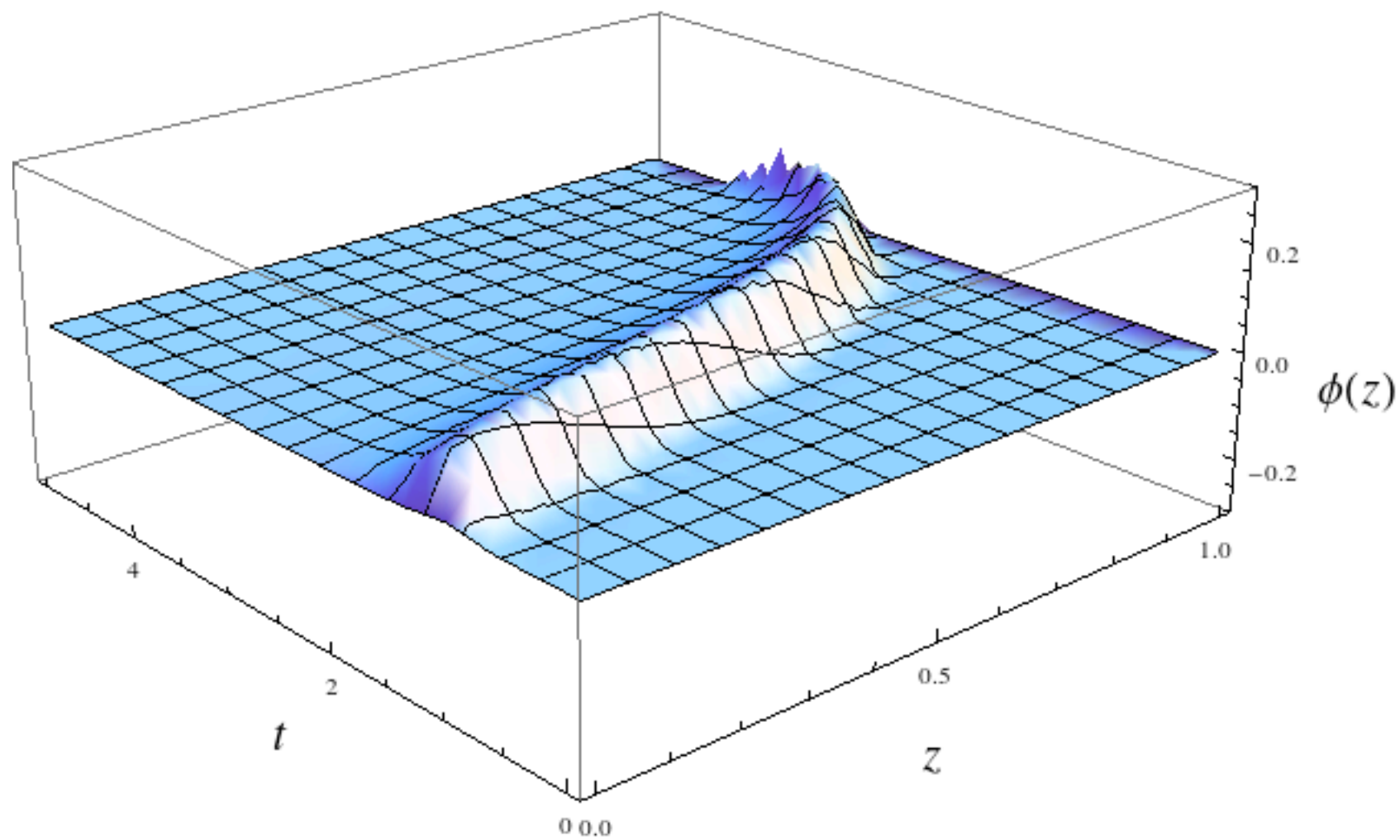
$$\phi(z) = \alpha z^{1/2} \ln(\Lambda z) + \beta z^{1/2} + \mathcal{O}\left(z^{3/2} \log(\Lambda z)\right)$$

Kondo coupling: κ

Condensate: $\alpha = \kappa\beta$

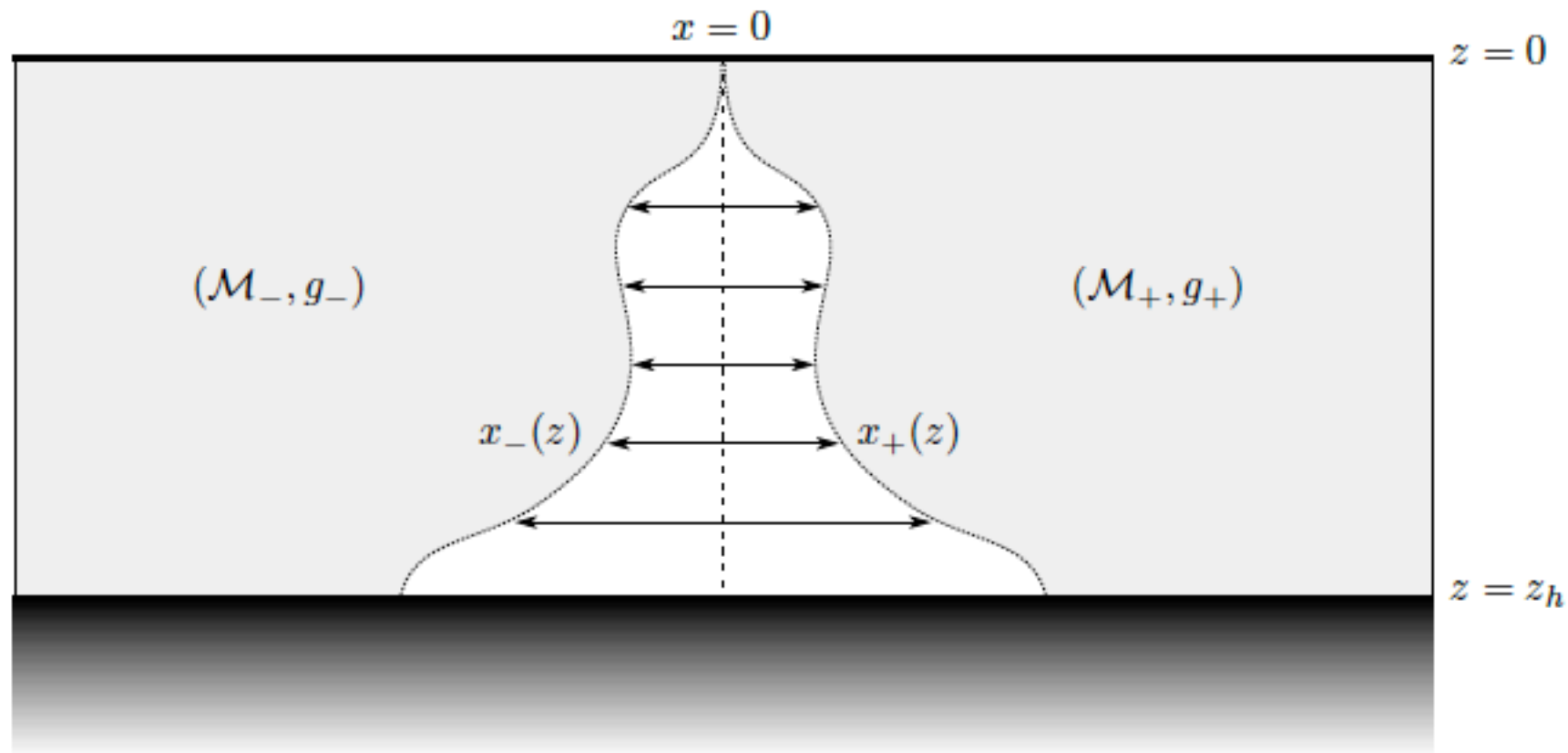
Scalar $\phi(z, t)$

Gaussian condensate pulse propagates to horizon
and falls into black hole



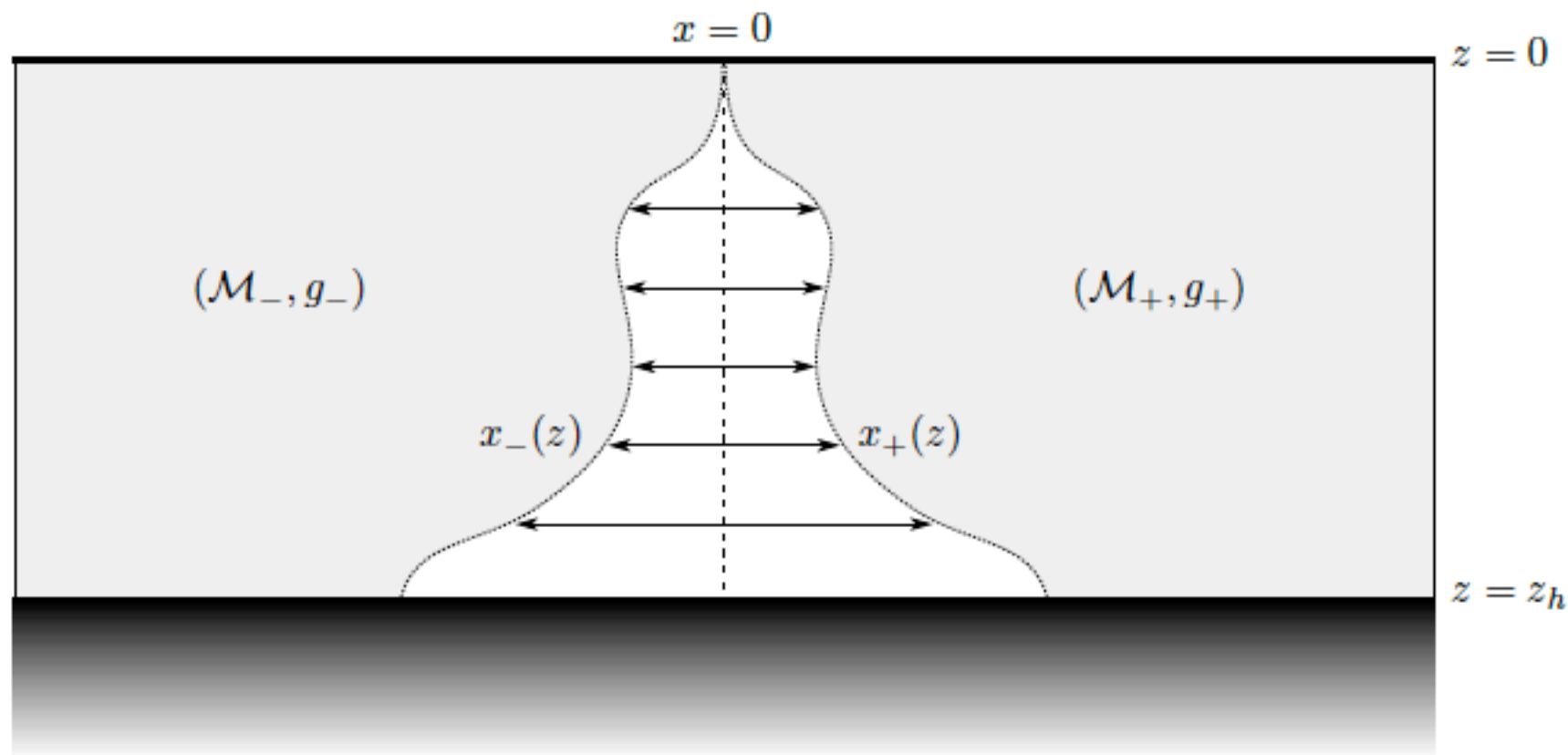
Including the backreaction

Flory, Newrzella, J.E., Hoyos, O'Bannon in progress



Including the backreaction

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Entanglement Entropy

Conclusion

- Kondo effect at large N : $(0+1)$ -dimensional superfluid
- **Holographic model:** $S = S_{\text{CS}} + S_{\text{AdS}_2}$
- **Two couplings:** 't Hooft coupling (large), coupling of double-trace operator (runs)
- RG flow, screening, phase shift, power-law scaling
- Time dependence and further extensions under investigation

Conference summary and outlook

Conference summary and outlook

Applied

QCD: Mesons, Veneziano limit
Time dependence
Anisotropies, shock waves

Condensed matter:
Lifshitz spaces
Non-relativistic hydrodynamics
Conductivities
Broken translation invariance
Quantum dots
Graphene
Quantum quenches
Experiments ?

Fundamental

(In)stability of AdS
Entanglement entropy
Higher order curvature terms
New explicit duality examples
Higher spin theories
Matrix models
Bootstrap
Quantum spectral curves

- Extremely active subject
- Diversity of examples and methods
- Wanted: Smoking gun example (experiment)
- Understand gauge/gravity duality?

Valentina, Lárus and Siggi

Thank you very much

for a wonderful workshop

in a magnificent location