Dark Energy Interactions – the Nordic Cosmology Workshop NORDITA, Stockholm, October, 1, 2014

DARK ENERGY ELECTRODYNAMICS

Balakin A.B.

Institute of Physics, Kazan Federal University, Kazan, Russia



Plan of the talk

- Motivation and analogies
- Mathematical formalism
- Exact analytical solutions
- Qualitative analysis and numerical simulation
- Cosmological applications: Unlighted epochs

Main idea: although Dark Energy seems to be an electrically neutral substratum, it can be considered as a cosmic medium, within which electromagnetic waves propagate; and we can search for fingerprints of a non-stationary Dark Energy in cosmic electrodynamic systems...

Dark Fluid

Dark Energy 72%

Gradient-type scheme of interaction

Archimedean – type force

Pyro-, piezo- and striction-type schemes of interaction

Non-minimal scheme of interaction Baryon Matter 5%

Optical activity-type scheme of interaction (via axionic Dark Matter)

Dark Matter 23%

PHOTONS



The talk is based on the results published in the papers

Balakin A.B., Bochkarev V.V. Archimedean-type force in a cosmic dark fluid.
 I. Phys. Rev. D 83, 024035 (2011) <u>arXiv:1012.2431</u>
 II. Phys. Rev. D 83, 024036 (2011) <u>arXiv:1012.2433</u>
 III. Phys. Rev. D 87, 024006 (2013) <u>arXiv:1212.4094</u>

Balakin A.B., Bochkarev V.V. and Lemos J.P.S. Non-minimal coupling...
 I. Phys. Rev. D 85, 064015 (2012) <u>arXiv:1201.2948</u>
 II. Phys. Rev. D 77, 084013 (2008) <u>arXiv:0712.4066</u>

• Balakin A.B., Dolbilova N.N. Electrodynamic phenomena induced by a dark fluid: Analogs of pyromagnetic, piezoelectric, and striction effects. Phys. Rev. D , 89, 104012 (2014) <u>arXiv:1404.5058</u>

Balakin A.B., Ni Wei-Tou. Non- minimal coupling of photons and axions
 I. Classical and Quantum Gravity, 27, 055003 (2010) arXiv:0911.2946
 II. Classical and Quantum Gravity, 31, 105002 (2014) arXiv:1403.6711



Electrodynamic equations

$$\nabla_k H^{ik} = -\frac{4\pi}{c} I^i \qquad \nabla_k F^{*ik} = 0$$

Excitation tensor

$$H^{ik} = \mathcal{H}^{ik} + \phi F^{*ik} + \mathcal{C}^{ikmn} F_{mn}$$

$${\cal H}^{ik}\equiv \pi^{ik}DW+{\cal D}^{ikpq}{\cal P}_{pq}$$

Total linear response tensor

$$\mathcal{C}^{ikmn} \equiv C^{ikmn}_{(0)} + \lambda^{ikmn} DW + Q^{ikmnpq} \mathcal{P}_{pq}$$

Convective derivative

 $D \equiv U^i \nabla_i$.

DE stress-energy
tensor
$$T_{ik}^{(DE)} \equiv WU_iU_k + \mathcal{P}_{ik}$$

Extended electric, magnetic permittivities and tensor of magneto-electric cross-effects

 $arepsilon^{im} = 2\mathcal{C}^{ikmn}U_kU_n$ = $arepsilon_{(0)}^{im} + \sigma^{im}DW + lpha^{im(pq)}\mathcal{P}_{pq}$

$$(\mu^{-1})^{ab} = -\frac{1}{2} \eta^{a}{}_{ik} \mathcal{C}^{ikmn} \eta_{mn}{}^{b}$$
$$= (\mu^{-1})^{ab}_{(0)} + \rho^{ab} DW + \beta^{ab(pq)} \mathcal{P}_{pq}$$

$$\nu^{am} = \eta^{a}{}_{ik} \mathcal{C}^{ikmn} U_{n}$$
$$= \nu^{am}_{(0)} + \omega^{am} DW + \gamma^{am(pq)} \mathcal{P}_{pq}$$

based on the standard decomposition of the linear response tensor

$$\begin{aligned} \mathcal{C}^{ikmn} &= \left(\varepsilon^{i[m} U^{n]} U^{k} - \varepsilon^{k[m} U^{n]} U^{i} \right) \\ &- \frac{1}{2} \eta^{ikl} (\mu^{-1})_{ls} \eta^{mns} + \eta^{ikl} U^{[m} \nu_{l}{}^{n]} + \eta^{lmn} U^{[i} \nu_{l}{}^{k]} \end{aligned}$$

Model of spatially isotropic non-stationary Dark Energy

$$\mathcal{P}_{ik} = -P\Delta_{ik}$$
 $\varepsilon^{im} = \Delta^{im}\varepsilon$ $(\mu^{-1})_{ab} = \frac{1}{\mu}\Delta_{ab}$ $\nu^{am} = 0$

$$C_{(0)}^{ikmn} = \frac{1}{2\mu_{(0)}} \left[g^{ikmn} + (\varepsilon_{(0)}\mu_{(0)} - 1)(g^{ikmn} - \Delta^{ikmn}) \right]$$
Reduced standard part of linear response tensor

Auxiliary tensors

 $\Delta^{ikmn} \equiv \Delta^{im} \Delta^{kn} - \Delta^{in} \Delta^{km} \quad \Delta^l_i \equiv \delta^l_i - U^l U_i \quad \eta^{ikl} \equiv \epsilon^{ikls} U_s$

$$\begin{aligned} Q^{ikmnpq} &= \frac{1}{2} \left[\alpha_{(1)} \Delta^{pq} (g^{ikmn} - \Delta^{ikmn}) \right. \\ &\quad + \alpha_{(2)} U_l U_s (g^{iklp} g^{mnsq} + g^{iklq} g^{mnsp}) \\ &\quad + \beta_{(1)} \Delta^{pq} \Delta^{ikmn} - \beta_{(2)} (\eta^{ikp} \eta^{mnq} + \eta^{ikq} \eta^{mnp}) \right] \end{aligned}$$

1

 $g^{ikmn} \equiv g^{im}g^{kn} - g^{in}g^{km}$

Reduced tensor of striction coefficients

Dark Energy fingerprints in a cosmic electrodynamics: Unlighted epochs in the Universe history

[epochs for which the square of an effective refraction index of the cosmic medium is negative, and thus propagation of electromagnetic waves is stopped]

Phase velocity

 $V_{\rm ph} \equiv$

$$\frac{\omega}{k} = \frac{1}{n(t)} \qquad V_{\rm gr} = \frac{2n}{n^2 + t^2}$$

Group velocity

Refraction index $n^2(t) \equiv \varepsilon(t)\mu(t)$ $n^2(t) < 0$

How does the non-stationary Dark Energy influence the refraction index ? We elaborated in detail two schemes of indirect interactions.

I. Archimedean (gradient – type) scheme with non-minimal coupling

FLRW-type space-time

$$ds^{2} = dt^{2} - a^{2}(t)[(dx^{1})^{2} + (dx^{2})^{2} + (dx^{3})^{2}]$$

Gravity field equations

$$H^2 = \frac{8\pi G}{3}(\rho + E),$$

Dark Energy Dark Matter energy densities

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} [(\rho + E) + 3(\Pi + P)],$$

Compatibility condition

Dark Energy Dark Matter pressures

 $\dot{\rho} + \dot{E} + 3H(\rho + E + \Pi + P) = 0$

Cosmological constant is included into the DE energy density and pressure **Kinetic description of the Dark Matter**

Kinetic equation

$$\frac{p^{i}}{m_{(a)}} \left(\frac{\partial}{\partial x^{i}} - \Gamma_{il}^{k} p^{l} \frac{\partial}{\partial p^{k}} \right) f_{(a)} + \frac{\partial}{\partial p^{i}} \left[\mathcal{F}_{(a)}^{i} f_{(a)} \right] = 0$$

Relativistic four-gradient-type force (Archimedean-type force), which acts on the DM particles in a DE reservoir

$$\mathcal{F}_{(a)}^{i} = m_{(a)} \mathcal{V}_{(a)} \left[g^{ik} - \frac{p^{i} p^{k}}{p^{l} p_{l}} \right] \nabla_{k} \Pi$$

New coupling constants

Classical 3-dimensional analog: the standard Archimedean force

$$\vec{F}_{(Arch)} = -V_0 \vec{\nabla} P_{(Pascal)}$$

Macroscopic moments of distribution functions of the DM particles

Menergy density
$$E(x) = \sum_{(a)} \frac{4\pi m_{(a)}^4}{x^3} \int_0^\infty q^2 dq f_{(a)}^0(q^2) \sqrt{1 + q^2 F_{(a)}(x)}$$

D

DM pressure

$$P(x) = \sum_{(a)} \frac{4\pi m_{(a)}^4 F_{(a)}(x)}{3x^3} \int_0^\infty \frac{q^4 dq f_{(a)}^0(q^2)}{\sqrt{1 + q^2 F_{(a)}(x)}}$$

Coupling functions

$$F_{(a)}(x) = \frac{1}{x^2} \exp\{2\mathcal{V}_{(a)}[\Pi(1) - \Pi(x)]\} \quad x = \frac{a(t)}{a(t_0)}$$

DM balance equations

$$\dot{E} + 3H(E+P) = -Q$$
 $Q \equiv 3\dot{\Pi} \sum_{(a)} \mathcal{V}_{(a)} P_{(a)}$

Hydrodynamic description of Dark Energy

Extended constitutive equation for DE

$$\rho(t) = \rho_0 + \sigma \Pi + \frac{\xi}{H(t)} \dot{\Pi}$$

Rheologic-type contribution

Balance equations for DE

$$\dot{\rho} + 3H(\rho + \Pi) = Q$$
 $Q = 3\Pi \sum_{(a)} \mathcal{V}_{(a)} P_{(a)}$

Key equation for the DE pressure

 $\xi x^2 \Pi''(x) + x \Pi'(x) (4\xi + \sigma) + 3(1 + \sigma) \Pi + 3\rho_0 = \mathcal{J}(x)$

Nonlinear Archimedean source in the key equation of DE evolution describing dark matter backreaction on the DE

$$\mathcal{J}(x) = -\sum_{(a)} E_{(a)} \frac{[x^2 F_{(a)}(x)]'}{2x^4} \int_0^\infty \frac{q^4 dq e^{-\lambda_{(a)}} \sqrt{1+q^2}}{\sqrt{1+q^2} F_{(a)}(x)}$$

Analytic example I: the case of massless (or ultrarelativistic) DM

$$\mathcal{J}_{(0)}(x) = E_{(0)} \frac{\mathcal{V}_{(0)} \Pi'(x)}{x^3} \exp\{\mathcal{V}_{(0)}[\Pi(1) - \Pi(x)]\}$$

Analytic example II: the case of cold DM

$$\mathcal{J}_{(C)}(x) = \frac{3N_{(C)}T_{(C)}\mathcal{V}_{(C)}\Pi'(x)}{x^4}e^{2\mathcal{V}_{(C)}[\Pi(1)-\Pi(x)]}$$

Exact solution of the anti-Gaussian type

Exact logarithmic solutions for DE pressure and energy density

$$\Pi(x) = \Pi(1) - \frac{4}{V_{(0)}} \log x \qquad \rho(x) = \rho(1) + \frac{4}{V_{(0)}} \log x \qquad E_{(0)} = \left[\frac{3\xi - 1}{V_{(0)}} - \frac{3}{4}\rho_{0}\right]$$
Hubble function
$$H^{2}(x) = \frac{8\pi G}{3} \left[\rho(1) + E_{(0)} + \frac{4}{V_{(0)}} \log x\right] \qquad E_{(a)} = \frac{N_{(a)}m_{(a)}\lambda_{(a)}}{K_{2}(\lambda_{(a)})}$$

$$\lambda_{(a)} = \frac{m_{(a)}}{k_{(B)}T_{(a)}}$$
Scale factor is regular
$$a(t) = a^{*} \exp\left\{\frac{8\pi G}{3V_{(0)}}(t - t^{*})^{2}\right\} \qquad H(t) = \frac{16\pi G}{3V_{(0)}}(t - t^{*})$$
Acceleration $-q(t) = \frac{\ddot{a}}{aH^{2}} = 1 + \frac{3V_{(0)}}{16\pi G(t - t^{*})^{2}} \ge 1$
Perpetually accelerated Universe

 $\sigma = -1$

Explicit example of the Little Rip – type solution

Exact super-exponential solution

$$\sigma = -1 \ \xi = \frac{1}{3} \ \mathcal{V} = 0$$
 $\Pi(x) = \Pi(1) + 3[\rho(1) + \Pi(1) - \rho_0] \log x - \frac{9}{2} \rho_0 \log^2 x$
 $\rho(x) = \rho(1) - 3[\rho(1) + \Pi(1)] \log x + \frac{9}{2} \rho_0 \log^2 x$

 Scale factor is super-exponential
 $\frac{a(t)}{a(t_0)} = \exp\left\{\sqrt{\frac{2\rho(1)}{9\rho_0}} \sinh[\sqrt{12\pi G\rho_0}(t - t_0)]\right\}$

 Hubble function
 $H(t) = \sqrt{\frac{8\pi G\rho(1)}{3}} \cosh[\sqrt{12\pi G\rho_0}(t - t_0)]$

 Acceleration parameter
 $-q(t) = 1 + \sqrt{\frac{9\rho_0}{2\rho(1)}} \frac{\sinh[\sqrt{12\pi G\rho_0}(t - t_0)]}{\cosh^2[\sqrt{12\pi G\rho_0}(t - t_0)]}$

Explicit example of the Little Rip – type solution

Unlighted epochs in the Universe history

Refraction index calculated using nonminimal interaction scheme

$$n^{2}(t) = \frac{1 - QH^{2}(t)[1 + q(t)]}{1 + QH^{2}(t)[1 + q(t)]} = \frac{1 + Q\dot{H}(t)}{1 - Q\dot{H}(t)}$$

The model operates with one nonminimal coupling parameter and time derivative of the Hubble function

For more details see, e.g.,

Balakin A.B., Bochkarev V.V., Lemos J.P.S. Phys. Rev. D 85, 064015 (2012)

Numerical study of the Archimedean-type model: Solution without transition points (perpetually accelerated Universe)



Numerical study of the Archimedean-type model: Periodic solution with infinite number of transition points



Numerical study of the Archimedean-type model: Solution with one transition point





Numerical study of the Archimedean-type model: Solution with two transition points



Numerical study of the Archimedean-type model: Quasi-periodic solution with finite number of transition points



Refraction index calculated using the interaction scheme involving analogy with electro- and magneto-striction Balakin A.B., Dolbilova N.N. Phys. Rev. D , 89, 104012 (2014)

$$\varepsilon = \varepsilon_{(0)} + \lambda_1 DW - \alpha P$$

$$\frac{1}{\mu} = \frac{1}{\mu_{(0)}} + \lambda_2 DW - \beta P$$

DE pressure

1

Coupling parameters describing striction-type interaction of DE with spatially isotropic electrodynamic system

1

$$n^2 \equiv \varepsilon \mu = \frac{n_{(0)}^2 + \mu_{(0)}(\lambda_1 DW - \alpha P)}{1 + \mu_{(0)}(\lambda_2 DW - \beta P)}$$

Typical behavior of the squared effective refraction index for two exact analytic solutions: anti-Gaussian and super-exponential



Unlighted epochs (periods with negative squared refraction index) of the I type (continuous), II type (discontinuous simply connected) and III type (discontinuous double-connected)

Conclusions

- The main idea of the work is to try to find Dark Energy fingerprints in the cosmic electrodynamic systems; we discussed the DE fingerprints associated with the so-called unlighted (dark) epochs in the Universe history, during which electromagnetic waves cannot propagate and thus cannot scan Universe interior.
- The unlighted epochs could be formed due to non-minimal coupling of photons and gravitons, and due to a striction-type coupling of an electrodynamic system with the Dark Energy, when the Universe expansion is non-stationary and the Universe history contains epochs of accelerated and decelerated expansion.
- Specific non-stationary Universe evolution displaying the unlighted epochs can be provided by an Archimedean-type interaction between Dark Energy and Dark Matter, which plays a role of effective energy-momentum re-distributor between DM and DE components of the Dark Fluid.

Thank you for the attention!