

The Return of Newton-Cartan space-time

The Spacetime Odyssey Continues

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based on work with:

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1409.1519, 1409.1522, 1502.00228, 1504.07461 (& to appear)

and

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1311.4794 (PRD) & 1311.6471 (JHEP)

Outline

- Why Newton-Cartan (NC) ? (-> non-relativistic space-time)
 - holography,
 - field theory
 - gravity
- What is NC (& its torsionful generalization TNC) geometry ?
 - NC from gauging the Bargmann algebra
- What theory of gravity does one get when making TNC dynamical ?
 - connection to Horava-Lifshitz gravity
 - perspectives (cosmology)
- Outlook

Motivation (Holography)

AdS/CFT has been very successful tool in studying strongly coupled (conformal) relativistic systems

-> **holography beyond original AdS-setup** ?

- How general is holographic paradigm ?
(nature of quantum gravity, black hole physics, cosmology)
- Examples of potentially holographic descriptions based on **non-AdS space-times**:
Lifshitz, Schrödinger, warped AdS3 (Kerr/CFT), flat space-time.

- simplest example appears to be

Lifshitz spacetimes

$$ds^2 = -\frac{dt^2}{r^{2z}} + \frac{1}{r^2} (dr^2 + d\vec{x}^2)$$

characterized by **anisotropic (non-relativistic)**
scaling between time and space

$$t \rightarrow \lambda^z t, \quad \vec{x} \rightarrow \lambda \vec{x}.$$

[Kachru,Liu,Mulligan]

* introduced originally to study strongly coupled systems with critical exponent z

Motivation (Holography) cont'd

- for standard AdS setup: boundary geometry is Riemannian just like the bulk geometry
- not generic: in beyond-AdS holography **bdry. geometry typically non-Riemannian**

Christensen,Hartong,Rollier,NO (1311)

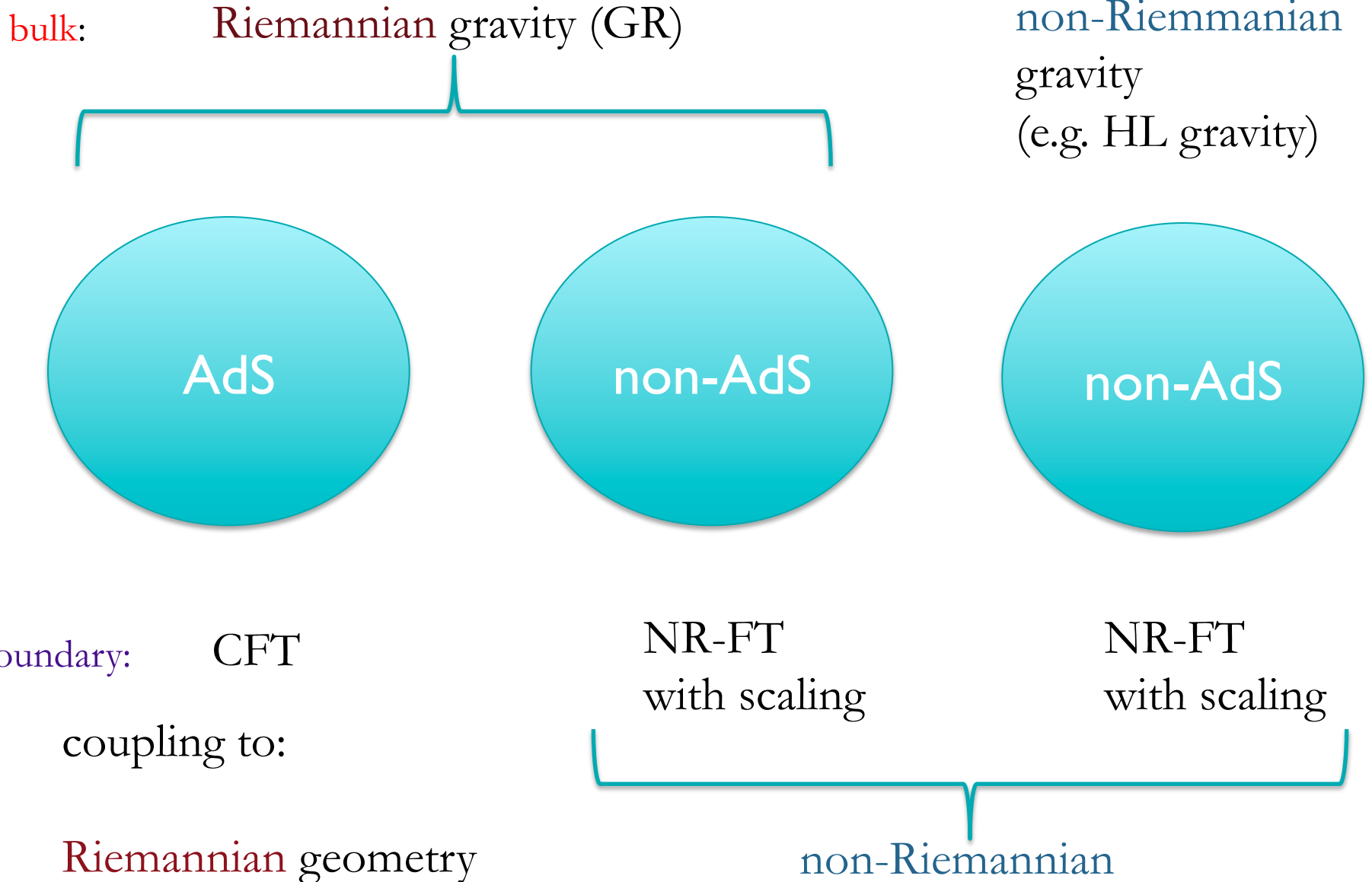
Hartong,Kiritsis,NO (1409)

-> need new approach: prime (simplest) example to gain traction = Lifshitz
(lessons can subsequently be applied to other cases)

by making the resulting non-Riemannian geometry dynamical one gains access to **other bulk theories of gravity** (than those based on Riemannian gravity)

- apply holography (e.g. HL gravity)
- interesting in their own right

Different Holographic setups



Motivation (Field Theory)

- in **relativistic FT**: very useful to couple to **background (Riemannian) geometry**
 - > compute EM tensors, study anomalies, Ward identities, etc.
- background field methods for systems with **non-relativistic (NR) symmetries** require **NC geometry (with torsion)**
 - > there is full **space-time diffeomorphism** invariance when coupling to the right background fields
- Recent examples
 - * Son's approach to the **effective field theory for the FQHE** [Son, 2013], [Geracie, Son, Wu, Wu, 2014]
 - * non-relativistic (NR) **hydrodynamics** [Jensen, 2014]

Motivation (Gravity)

- interesting to make NC geometry dynamical
- > “new” theories of gravity

Hartong, Kiritsis, NO (1504)

will see:

dynamical Newton-Cartan (NC) = Horava-Lifshitz (HL) gravity



natural geometric framework with full diffeomorphism invariance
& possibly non-trivial consequences for HL gravity

such theories of gravity interesting as

- other bulk theories of gravity in holographic setups
- effective theories

- novel theories of gravity to examine cosmology
- effective theories for inflation ?

Newton-Cartan makes Galilean local

- NC geometry originally introduced by Cartan to geometrize Newtonian gravity

→ both Einstein's and Newton's theories of gravity admit geometrical formulations which are **diffeomorphism invariant**

- NC originally formulated in “metric” formulation
more recently: **vielbein formulation** (shows underlying sym. principle better)
Andringa, Bergshoeff, Panda, de Roo

Riemannian geometry: tangent space is **Poincare invariant**

Newton-Cartan geometry: tangent space is **Bargmann (central ext. Gal.) invariant**

- gives geometrical framework and extension to include torsion
i.e. as geometry to which non-relativistic field theories can couple
(boundary geometry in holographic setup is non-dynamical)

* will next consider **dynamical** (torsional) Newton-Cartan

From Poincare to GR

- make **Poincare local** (i.e. gauge the translations and rotations)

vielbein  spin connection 

$$A_\mu = P_a e_\mu^a + \frac{1}{2} J_{ab} \omega_\mu^{ab}$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu + [A_\mu, A_\nu] = P_a R_{\mu\nu}{}^a(P) + \frac{1}{2} J_{ab} R_{\mu\nu}{}^{ab}(J)$$

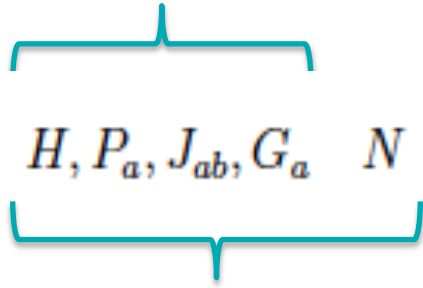
$$\delta A_\mu = \partial_\mu \Lambda + [A_\mu, \Lambda], \quad \Lambda = \xi^\mu A_\mu + \frac{1}{2} J_{ab} \lambda^{ab}$$

$$R_{\mu\nu}{}^a(P) = 0 \left\{ \begin{array}{l} \text{spin connection expressed in terms of vielbein} \\ \delta A_\mu \rightarrow \delta e_\mu^a = \mathcal{L}_\xi e_\mu^a + e_\mu^b \lambda_b^a \\ \text{covariant derivative defined via vielbein postulate} \\ R_{\mu\nu}{}^{ab}(J) = \text{Riemann curvature 2-form} \end{array} \right.$$

- GR is a **diff invariant theory** whose tangent space invariance group is the **Poincaré group**
- * **Einstein equivalence principle** -> **local Lorentz invariance**

Gauging the Bargmann algebra

Galilean



H, P_a, J_{ab}, G_a, N

Bargmann

(Galilean algebra is c to infinity limit of Poincare)

$$[H, G_a] = P_a$$

$$[P_a, G_b] = 0$$



$$[P_a, G_b] = N\delta_{ab}$$

gauge Bargmann and impose curvature constraints

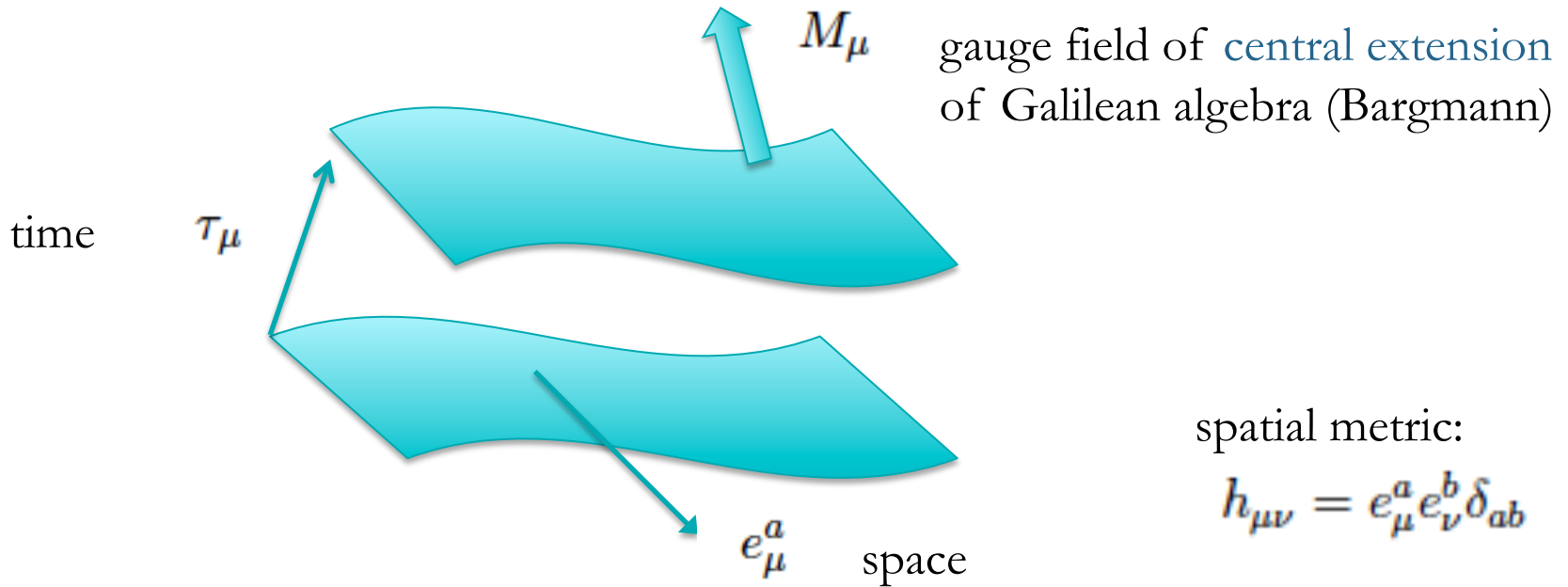
$$R_{\mu\nu}(H) = R_{\mu\nu}{}^a(P) = R_{\mu\nu}(N) = 0.$$

independent fields: τ_μ, e_μ^a, m_μ

$$h_{\mu\nu} = e_\mu^a e_\nu^b \delta_{ab}$$

= gauge fields of Hamiltonian, spatial translations and central charge

Newton-Cartan geometry



NC geometry = no torsion

$$\longrightarrow \tau_\mu = \partial_\mu t$$

notion of absolute time

TTNC geometry = twistless torsion $\longrightarrow \tau_\mu = \text{HSO}$

preferred foliation in equal time slices

TNC geometry no condition on τ_μ

- in TTNC: torsion measured by $a_\mu = \mathcal{L}_{\hat{v}} \tau_\mu$
 geometry on spatial slices is Riemannian

Adding torsion to NC

Christensen,Hartong,Rollier,NO
Hartong,Kiritsis,NO/Hartong,NO
Bergshoeff,Hartong,Rosseel

- inverse vielbeins

$$(v^\mu, e_a^\mu)$$

$$v^\mu \tau_\mu = -1, \quad v^\mu e_\mu^a = 0, \quad e_a^\mu \tau_\mu = 0, \quad e_a^\mu e_\mu^b = \delta_a^b$$

can build Galilean boost-invariants

$$\hat{v}^\mu = v^\mu - h^{\mu\nu} M_\nu,$$

$$\bar{h}_{\mu\nu} = h_{\mu\nu} - \tau_\mu M_\nu - \tau_\nu M_\mu,$$

$$\tilde{\Phi} = -v^\mu M_\mu + \frac{1}{2} h^{\mu\nu} M_\mu M_\nu,$$

-introduce Stueckelberg scalar chi
(to ensure N-invariance):

$$M_\mu = m_\mu - \partial_\mu \chi.$$



affine connection of TNC (inert under G,I,N)

$$\Gamma_{\mu\nu}^\rho = -\hat{v}^\rho \partial_\mu \tau_\nu + \frac{1}{2} h^{\rho\sigma} (\partial_\mu \bar{h}_{\nu\sigma} + \partial_\nu \bar{h}_{\mu\sigma} - \partial_\sigma \bar{h}_{\mu\nu})$$

with torsion $\Gamma_{[\mu\nu]}^\rho = -\frac{1}{2} \hat{v}^\rho (\partial_\mu \tau_\nu - \partial_\nu \tau_\mu)$

$$\nabla_\mu \tau_\nu = 0, \quad \nabla_\mu h^{\nu\rho} = 0,$$

analogue of metric compatibility

intermezzo: geodesics on NC space-time

- worldline action of non-rel particle of mass m on NC background

$$S = \int d\lambda L = \frac{m}{2} \int d\lambda \frac{\bar{h}_{\mu\nu} \dot{x}^\mu \dot{x}^\nu}{\tau_\rho \dot{x}^\rho}$$

[Kuchar],
[Bergshoeff et al]

- gives the geodesic equation with NC connection $\frac{d^2 x^\mu}{d\lambda^2} + \Gamma_{\nu\rho}^\mu \frac{dx^\nu}{d\lambda} \frac{dx^\rho}{d\lambda} = 0$,

* reduces to **Newton's law** $\frac{d^2 x^i}{dt^2} + \delta^{ij} \partial_j \Phi = 0$,

provided we take

$$\begin{aligned} M_t &= \partial_t M + \Phi, \\ M_i &= \partial_i M, \end{aligned}$$

for flat NC space-time: **zero Newtonian potential**

symmetries of flat NC = conformal Killing vectors (spanning **Lifshitz**) + **extra**

Dynamical Newton-Cartan geometry

so far: (T)NC geometry was non-dynamical:

- what happens when we allow it to fluctuate ?
- what is **the theory of gravity that incorporates local Galilean symmetry ?**
(Einstein equivalence principle, but applied to Galilean instead of Lorentz)

recently shown that:

[Hartong,NO]

- dynamical NC geometry = projectable HL gravity
- dynamical TTNC geometry = non-projectable HL gravity

* Horava-Lifshitz gravity was originally introduced as non-Lorentz invariant and renormalizable UV completion of gravity

- phenomenologically viable ?
- interesting theoretically as **alternate bulk gravity theories**
relevant to i) holography for strongly coupled non-relativistic systems
ii) alternate theories in cosmology

NC/TTNC gravity

TNC geometry is a natural geometrical framework underlying HL gravity

- NC quantities combine into: $g_{\mu\nu} = -\tau_\mu\tau_\nu + \hat{h}_{\mu\nu}$
- ADM parametrization of metric used in HL gravity:

$$ds^2 = g_{\mu\nu}dx^\mu dx^\nu = -N^2 dt^2 + \gamma_{ij} (dx^i + N^i dt) (dx^j + N^j dt)$$

relation:

$$\tau_\mu \sim \text{lapse} \quad , \quad \hat{h}_{\mu\nu} \sim \text{spatial metric} \quad , \quad m_\mu \sim \text{shift} + \text{Newtonian potential} \quad ,$$

some features:

- khronon field of BPS appears naturally $\tau_\mu = \psi \partial_\mu \tau$ Blas,Pujolas,,Sibiryakov (2010)

NC (no torsion): $N = N(t)$ projectable HL gravity

TTNC: $N = N(t, x)$ non-projectable HL gravity

- U(1) extension of HMT emerges naturally as Bargmann U(1)
- new perspective (via chi field) on nature of U(1) symmetry

Horava,Melby-Thompson,(2010)

Effective actions reproduce HL

- covariant building blocks:

- extrinsic curvature: $\hat{h}_{\nu\rho} \nabla_\mu \hat{v}^\rho = -K_{\mu\nu}$ spatial curvature $\overset{\vee}{R}_{\mu\nu\sigma}{}^\rho$.

- covariant derivative, torsion vector a_μ , inverse spatial metric $h^{\mu\nu}$

- tangent space invariant integration measure $e = \det(\tau_\mu, e_\nu^a)$

-> construct all terms that are **relevant or marginal** (up to dilatation weight $d+z$)

- in 2+1 dimensions for $1 < z \leq 2$

$$S = \int d^3x e [C (h^{\mu\rho} h^{\nu\sigma} K_{\mu\nu} K_{\rho\sigma} - \lambda (h^{\mu\nu} K_{\mu\nu})^2) - \mathcal{V}]$$

kinetic terms (2nd order)

potential:

$$-\mathcal{V} = 2\Lambda + c_1 h^{\mu\nu} a_\mu a_\nu + c_2 \mathcal{R} + \delta_{z,2} [c_{10} (h^{\mu\nu} a_\mu a_\nu)^2 + c_{11} h^{\mu\rho} a_\mu a_\rho \nabla_\nu (h^{\nu\sigma} a_\sigma) + c_{12} \nabla_\nu (h^{\mu\rho} a_\rho) \nabla_\mu (h^{\nu\sigma} a_\sigma) + c_{13} \mathcal{R}^2 + c_{14} \mathcal{R} \nabla_\mu (h^{\mu\nu} a_\nu) + c_{15} \mathcal{R} h^{\mu\nu} a_\mu a_\nu]$$

Perspectives

new perspectives on HL:

- different vacuum (flat NC space-time): reexamine issues with HL gravity
- IR effective theory for non-relativistic field theories
- insights into non-relativistic quantum gravity corner of $(\hbar, G_N, 1/c)$ cube ?
 - relevance for cosmology ?

alternate theories of gravity in cosmological scenarios, effective theories for inflation

e.g. aspects of HL cosmology:

- dark matter shows up as integration constant
- bouncing and cyclic universes
- scale-invariant perturbations ($z=3$, in $D=4$) without inflation

Outlook

- revisit **HL gravity** using TNC language-> implications for HL cosmology ?
 - gauging other groups: e.g. Carroll group > ultra-relativistic gravity [Hartong]
 - applications to non-rel. hydrodynamics (also relevant for cosmo):
fluid/gravity: black branes with zero/non-zero particle no. ? Galilean perfect fluids [Hoyos, Kim, Oz] [Jensen]
 - **NC supergravity, NC in string theory** [Bergshoeff et al]
-
- employ **similar techniques** to Schrödinger, warped AdS, flat space holography [Andrade, Keeler, Peach, Ross,]. [Hofman, Rollier] [Armas, Blau, Hartong (in progress)]
 - **adding charge (Maxwell in bulk),**
adding other exponents (hyperscaling, matter scaling) [Kiritsis, Goutereaux] [Gath, Hartong, Monteiro, NO] [Khveshchenko] [Karch] [Hartnoll, Karch]
 - effective TNC theories and their coupling to matter (e.g. QH-effect) [Son] et al