

# Fanning out of the $f$ -mode in the presence of nonuniform magnetic fields

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NORDITA

With

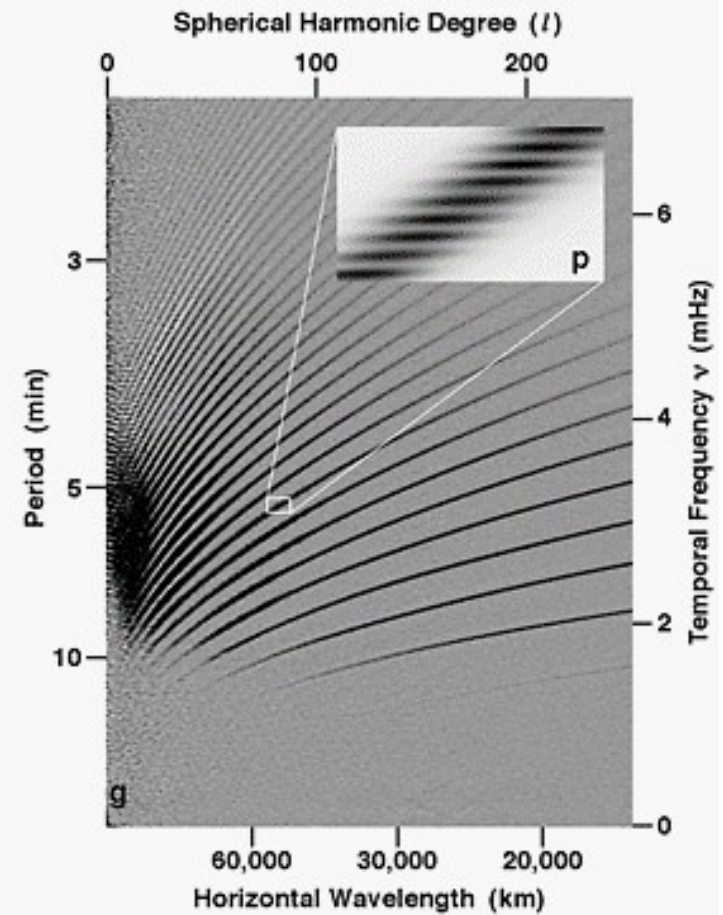
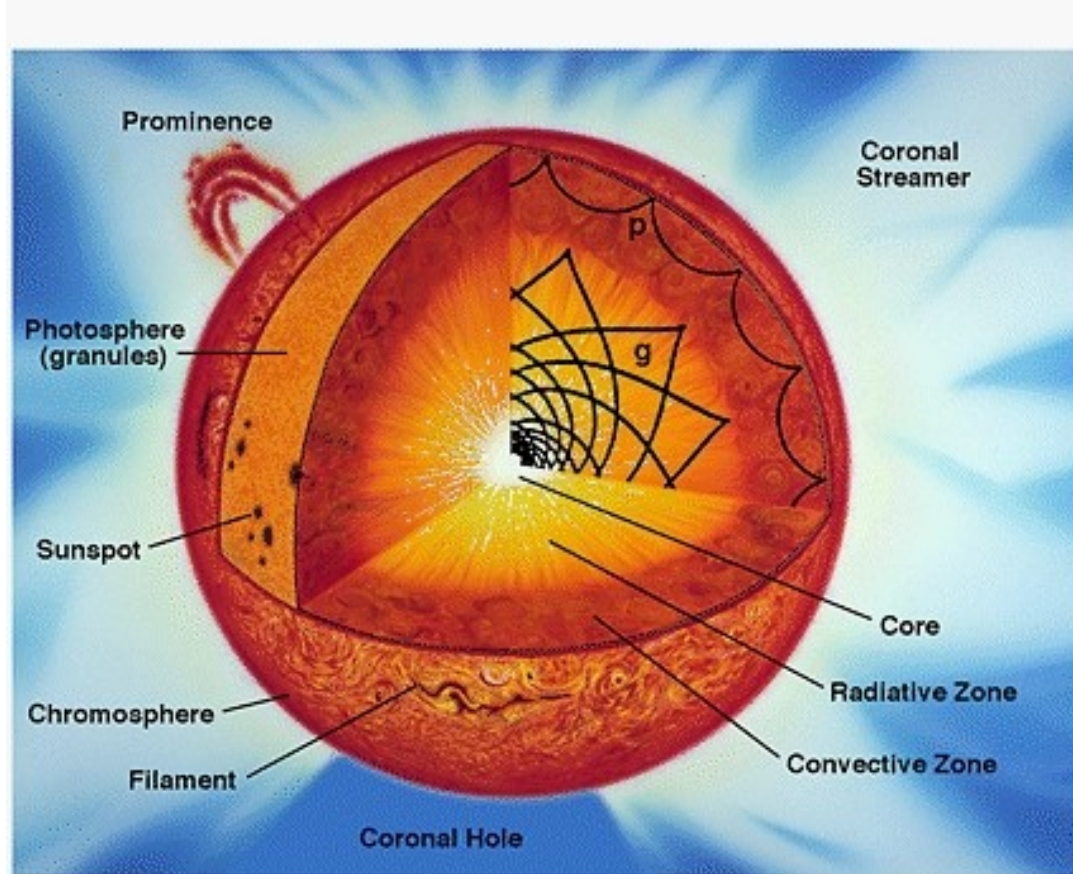
Axel Brandenburg, S. M. Chitre, Matthias Rheinhardt, Harsha Raichur

Simulations: <https://code.google.com/p/pencil-code/>

Observations: SDO/HMI

Pencil Code User Meeting, Trondheim, 14 MAY 2015

# The Sun and Helioseismology

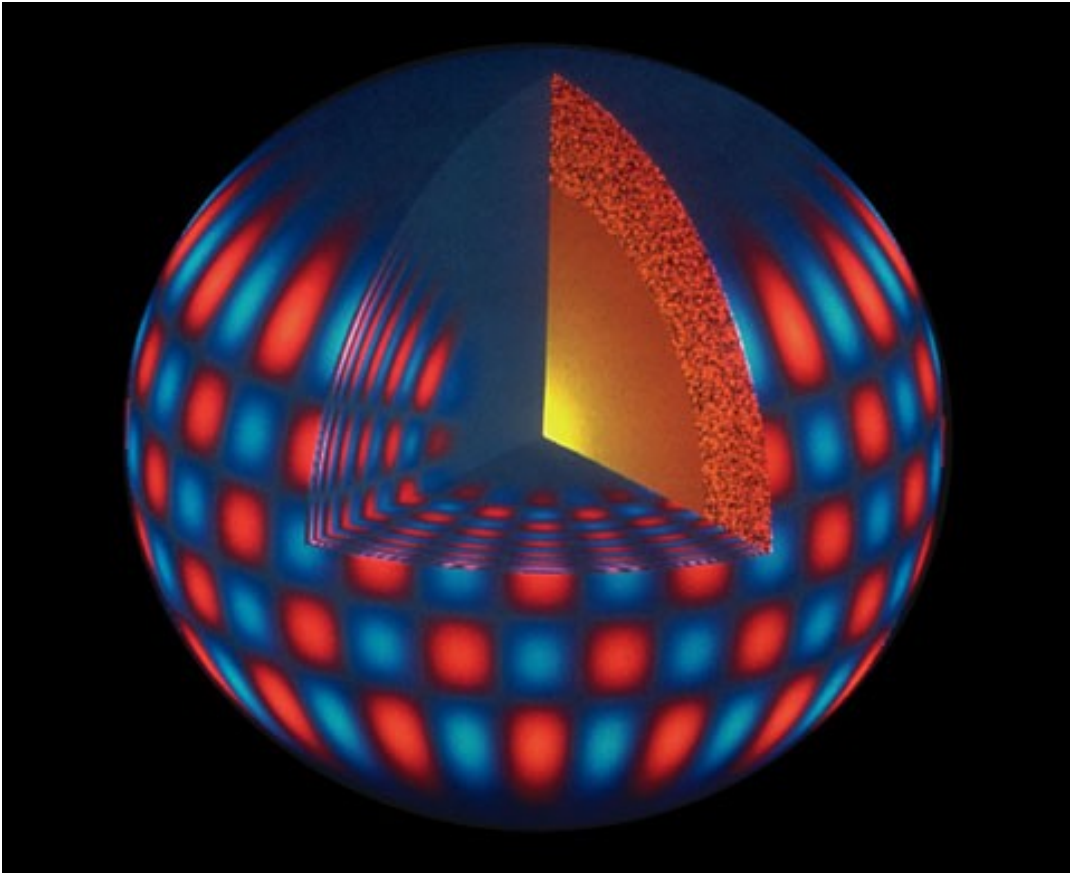


The Sun supports a wide variety of waves

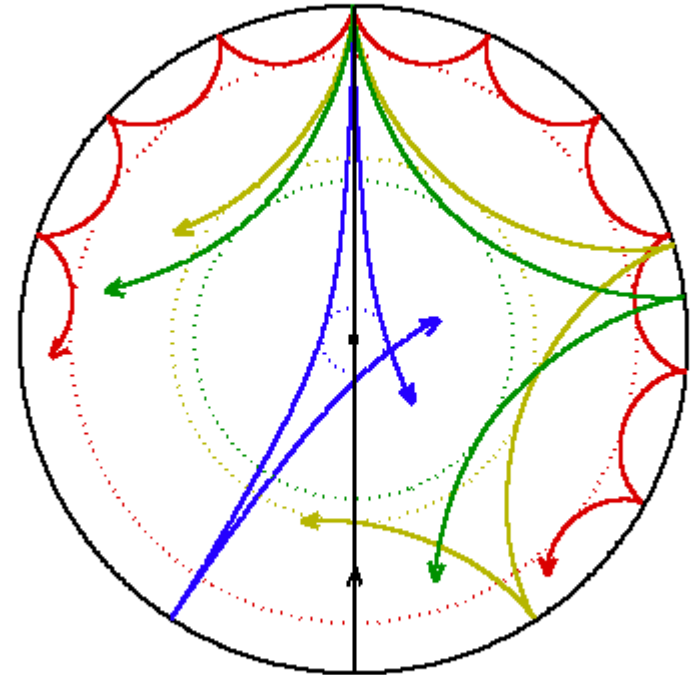
Helioseismology: science studying the wave oscillations in the Sun

Goal is to determine the physical conditions in the Solar interior

# Oscillations



*About 250,000 out of 10 million expected modes have been identified*



Internal ray paths of acoustic waves

- Refraction of waves
- Low  $l$ 's or large wavelengths penetrate to deeper layers

$$r_{\min} = \frac{c_s}{\omega} \sqrt{l(l+1)}$$

# The modes: dispersion relations

- ***g*-modes**  Gravity waves (Buoyancy driven)

$$\omega_{\text{BV}}^2 = \frac{g}{c_p} \frac{ds}{dz}$$

- In stably stratified medium
- Expected to exist in the core of the Sun
- Not yet observed in Sun !

- ***p*-modes**  Sound waves (Pressure driven)

$$\omega^2 = \omega_c^2 + c_s^2 (k_x^2 + k_z^2)$$

(Isothermal)

- Excited by convective motions
- Observed and most extensively used
- The Sun is transparent to acoustic radiation
- Energy travel time through the Sun ~ few hrs

- ***f*-mode**



Surface waves (On interface)

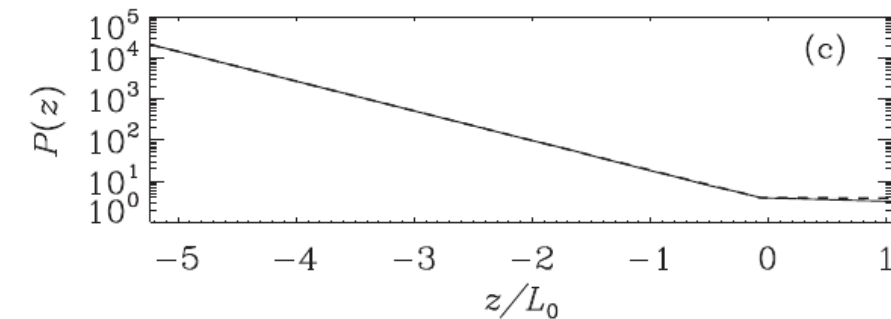
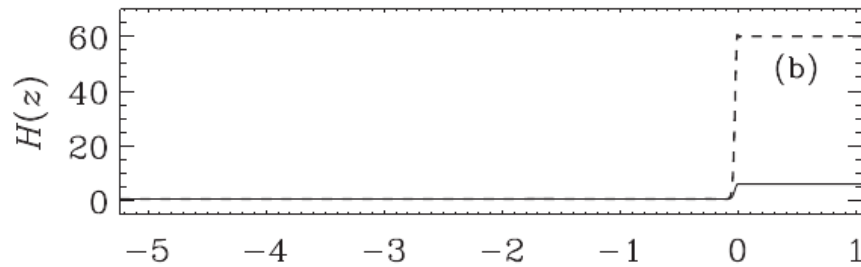
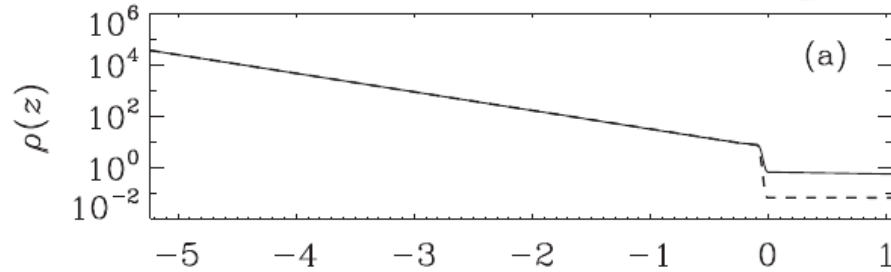
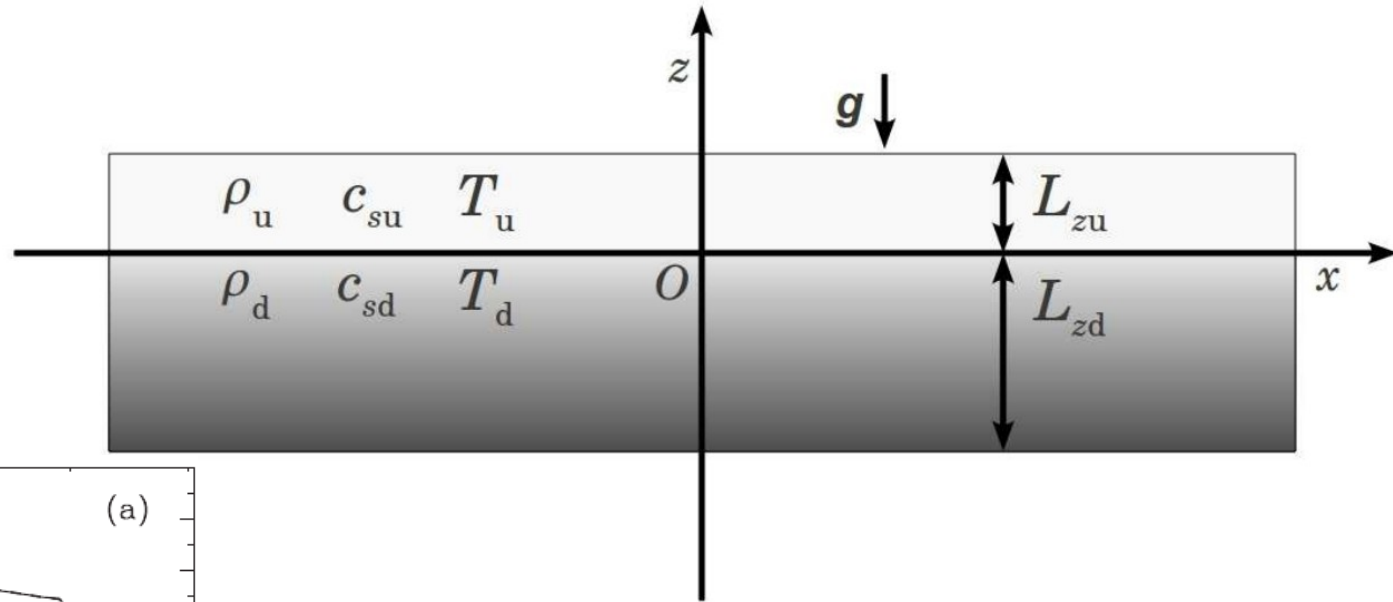
$$\omega_f^2 = gk_x \frac{1 - q}{1 + q}$$

- Water (surface) waves; dispersive
- Observed on solar surface
- Traditionally thought to be of less diagnostic value
- Significant frequency shifts are detected
- Shifts could tell us about the inhomogeneous structure

with  $q < 1$

# Our model setup and initial conditions

- 2-D Cartesian section
- Stably stratified under  $g$
- Piecewise isothermal
- Interface at  $z = 0$
- Bulk ( $z < 0$ )
- Atmosphere ( $z > 0$ )



- Sharp jump at the interface:

$$q = \rho_u(0)/\rho_d(0) = c_{sd}^2/c_{su}^2$$

- Background density:

$$\rho_{d,u}(z) = \rho_{d,u}(0) \exp(-z/H_{d,u})$$

- Scale height:

$$H_{d,u} = (c_p - c_v)T_{d,u}/g$$

$$L_0 = \gamma H_d$$



# Wave excitation and measurements

- Weak stochastic forcing  $\mathbf{f}$  in a homogeneous and isotropic fashion
- Stirring scale  $\ll$  the box dimensions (we chose,  $k_f/k_1 = 20$ )
- No effect seen when restricted the forcing to lower subdomain
- We measure the vertical velocity at the interface  $z=0$ ,  $u_z(x, 0, t)$
- Fourier transforming this in  $x$  and  $t$ , we get  $\hat{u}_z(k_x, \omega)$

## Useful dimensionless quantities

$$\text{Re} = \frac{u_{\text{rms,d}}}{\nu k_f} \quad \text{Ma} = \frac{u_{\text{rms,d}}}{c_{\text{sd}}} \quad \mathcal{F} = f_0 \text{Rc} = \frac{f_0 c_{\text{sd}}}{\nu k_f}$$

We define,

$$\tilde{k}_x = k_x L_0, \quad \tilde{\omega} = \frac{\omega}{\omega_0}, \quad \text{where} \quad \omega_0 = \frac{g}{c_{\text{sd}}}$$

$$\tilde{P}(k_x, \omega) = \frac{|\hat{u}_z|}{D^2} = \frac{|\hat{u}_z|}{L_0^2} \frac{c_{\text{sd}}^2}{u_{\text{rms,d}}^2}$$

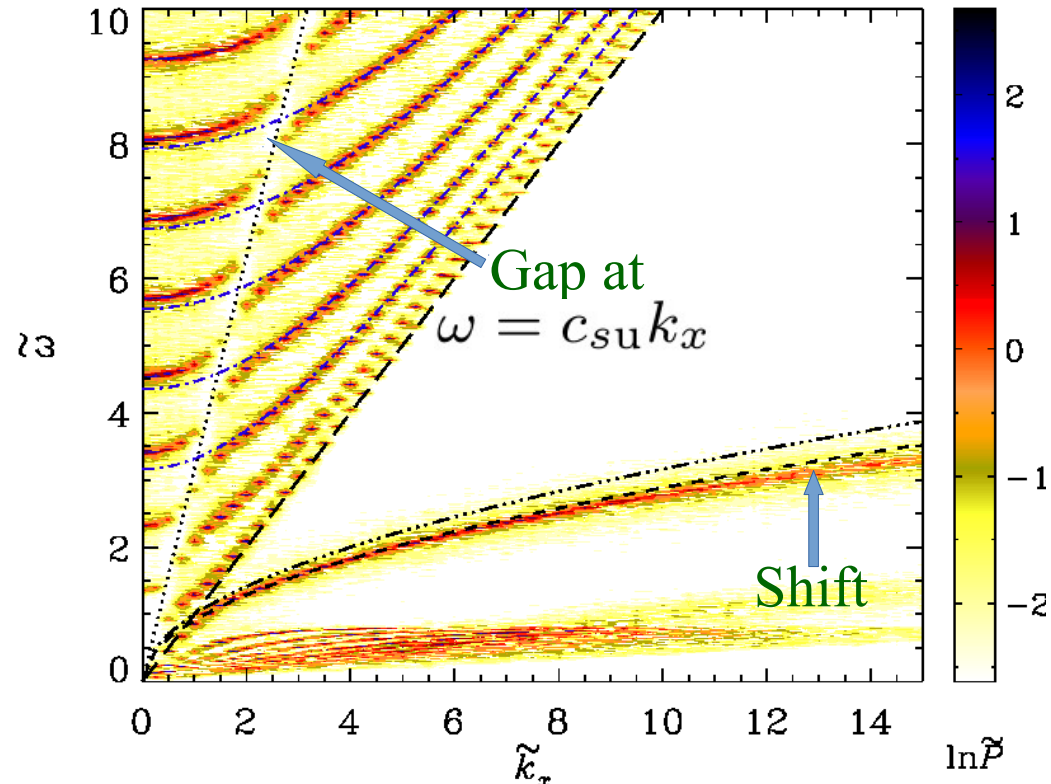
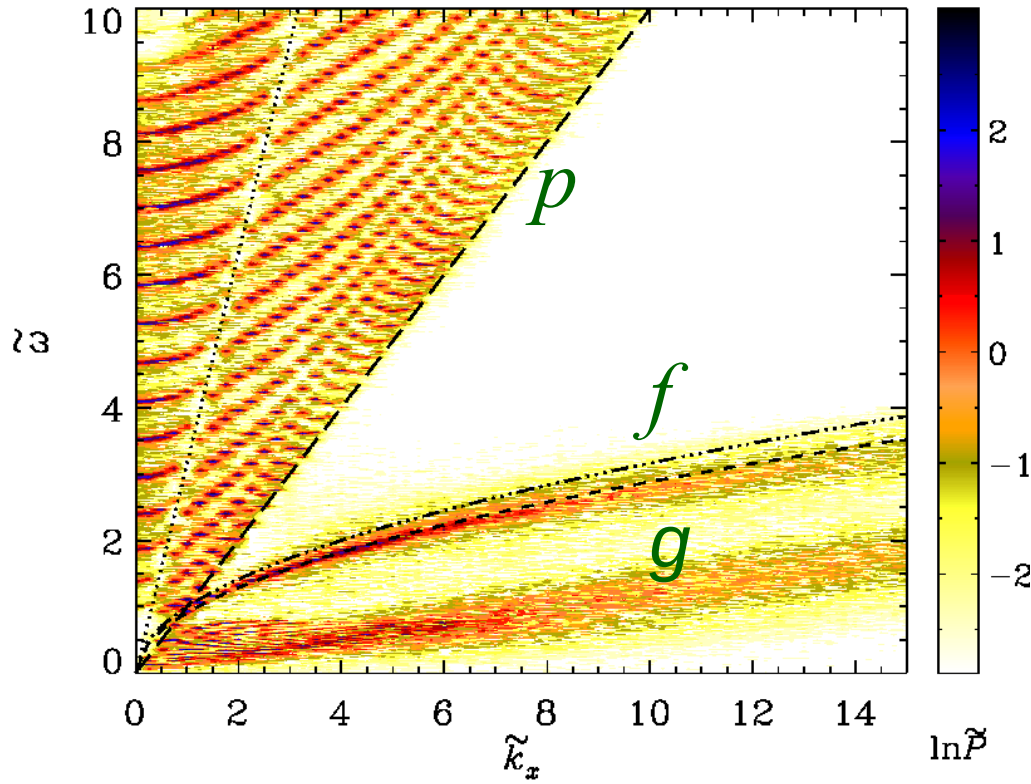
# Non-magnetic case

power\_xy  
(Matthias)

$8\pi \times 2\pi$  (1024 x 600)

$q = 0.1$

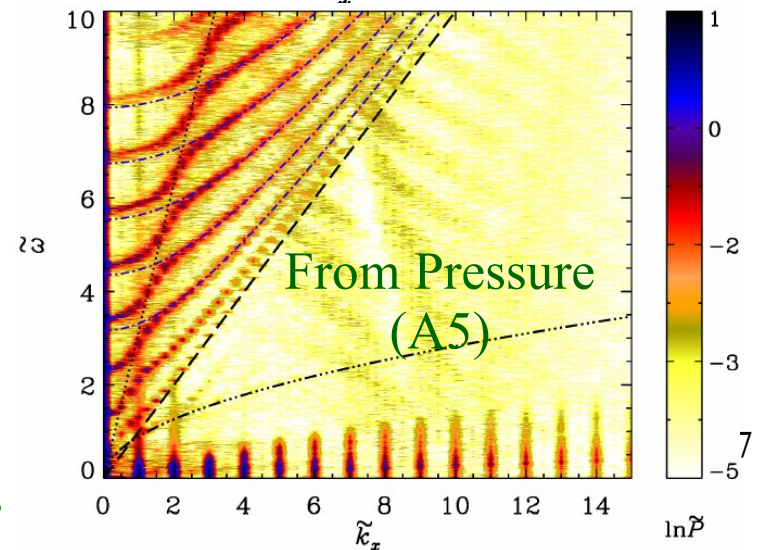
$8\pi \times \pi$  (1024 x 300)



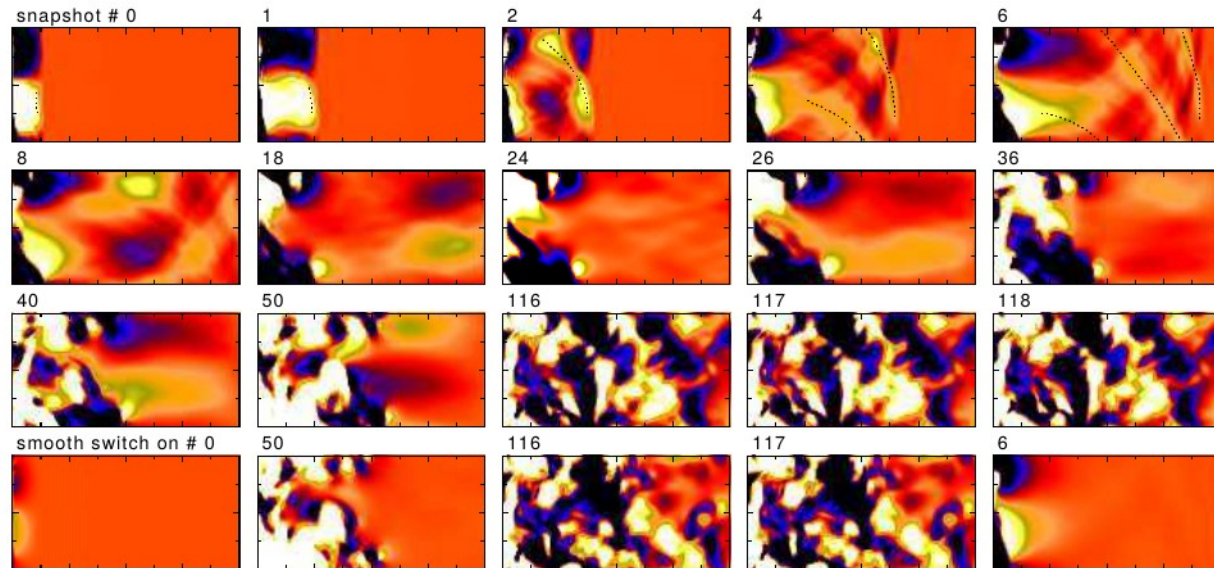
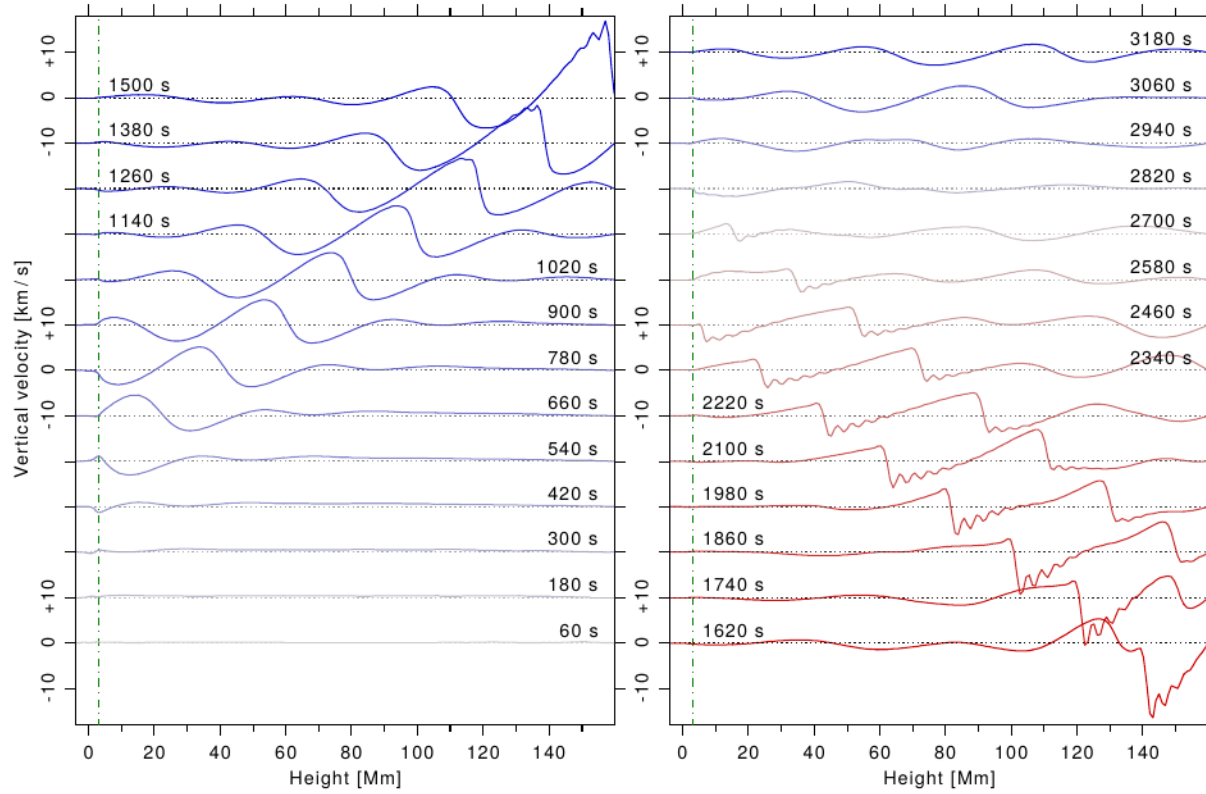
- Second family of  $g$ -modes due to hot layer above Hindman & Zweibel (94); Wagner & Schmitz (07)

- Avoided crossings  $\Rightarrow$  “ $a$ -mode”
- Acoustic mode localized in the upper hot layer
- Discussed/seen in more realistic setup of HZ94
- Non-detection in solar context  $\Rightarrow$

*lower limit on chromospheric temperature ?*



# Switch-on effects: Bourdin (2014)



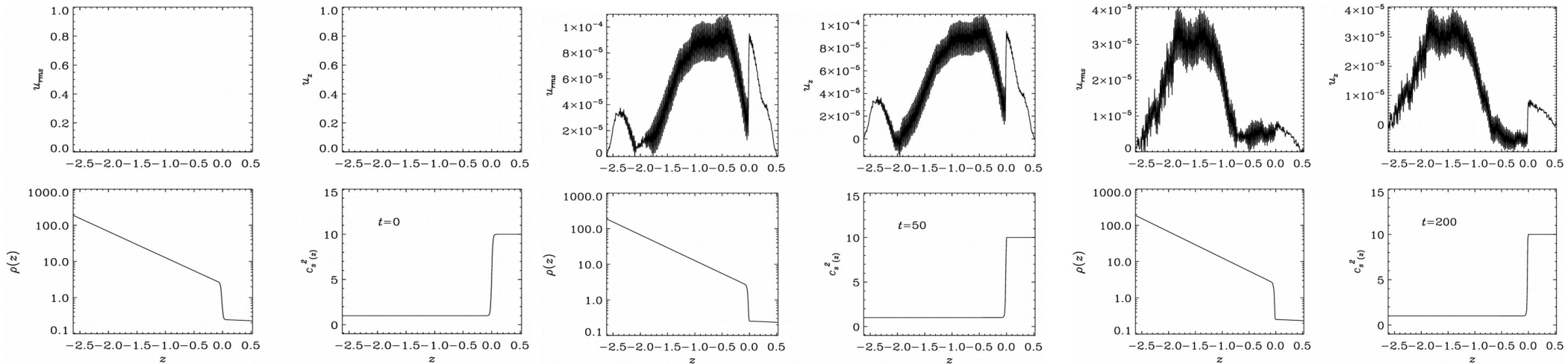


# Experiment (Philippe and Matthias)

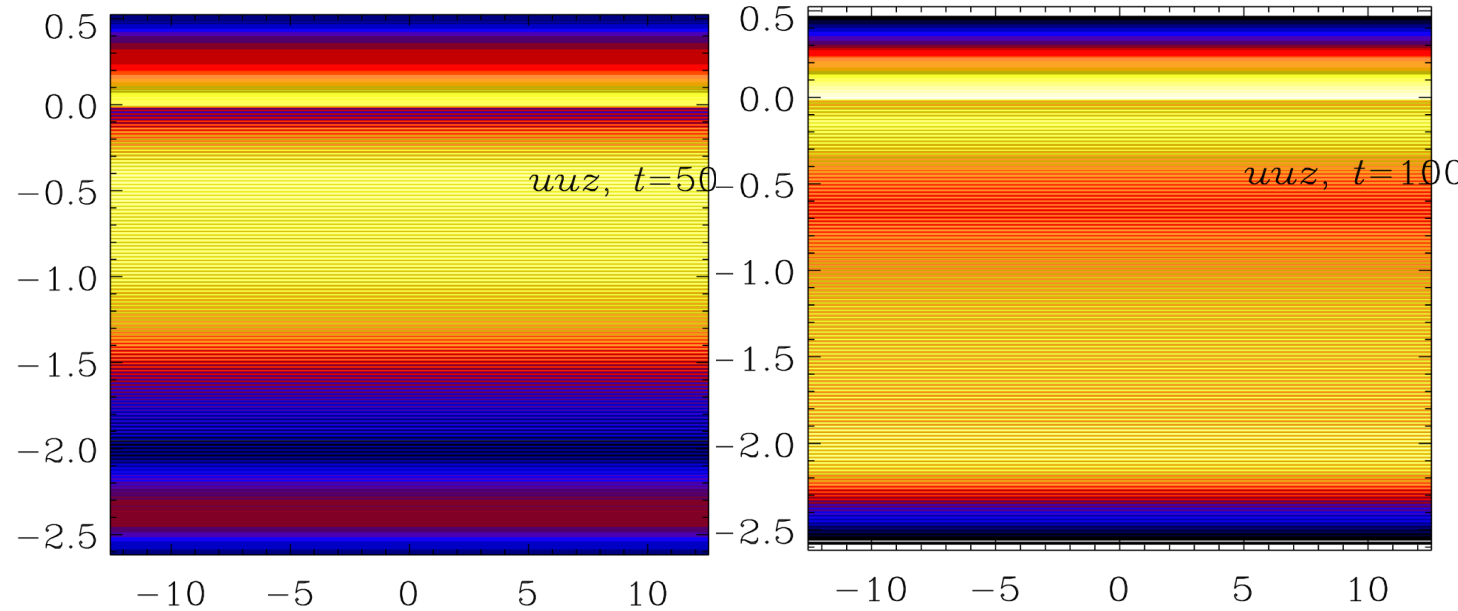
$t=0$

$t=50$

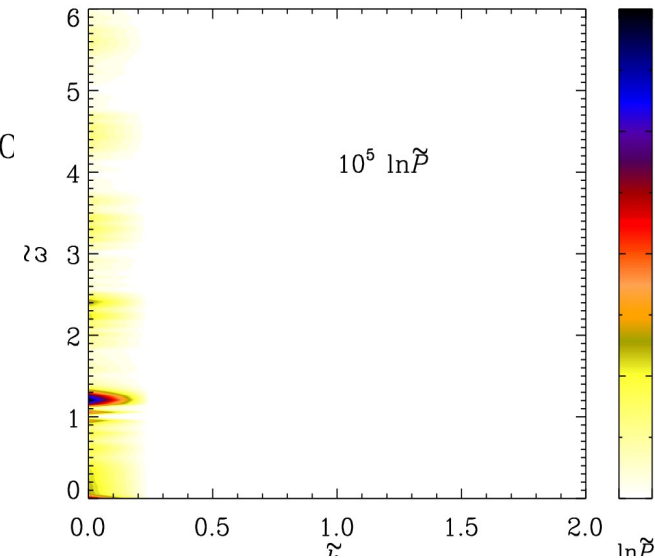
$t=200$



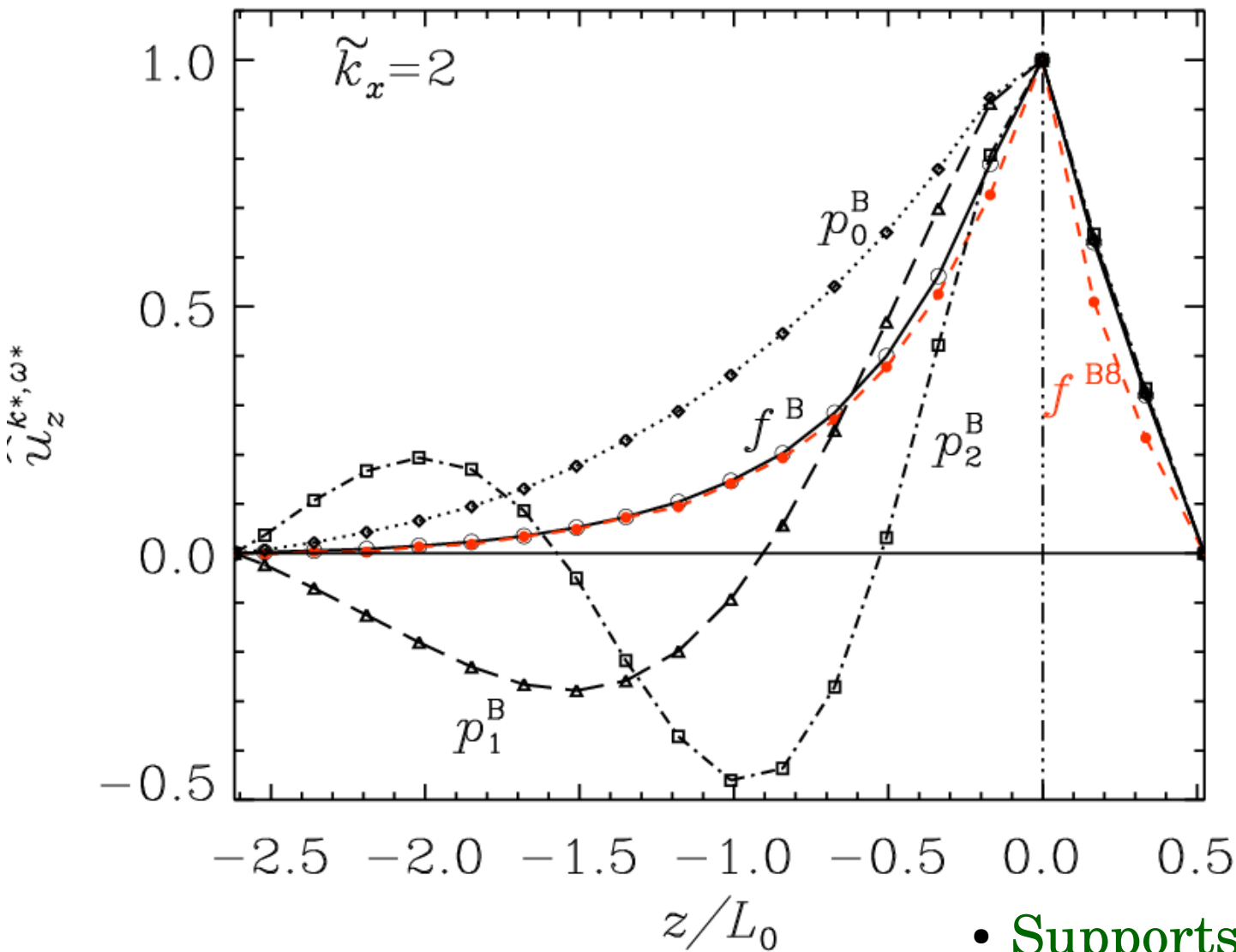
➔ Time



k-omega



# Eigenfunctions ( $f$ AND $p_0$ ) and Duvall's Law



- **Duvall's Law:**

$$k_z = \frac{\pi(n_D + \alpha)}{L_{zd}}$$

$$n_D = 1, 2, 3, \dots$$

**Best Fit:**  $\alpha = 3/2$

- **These simulations:**

$$k_z = \frac{\pi(n + 1/2)}{L_{zd}}$$

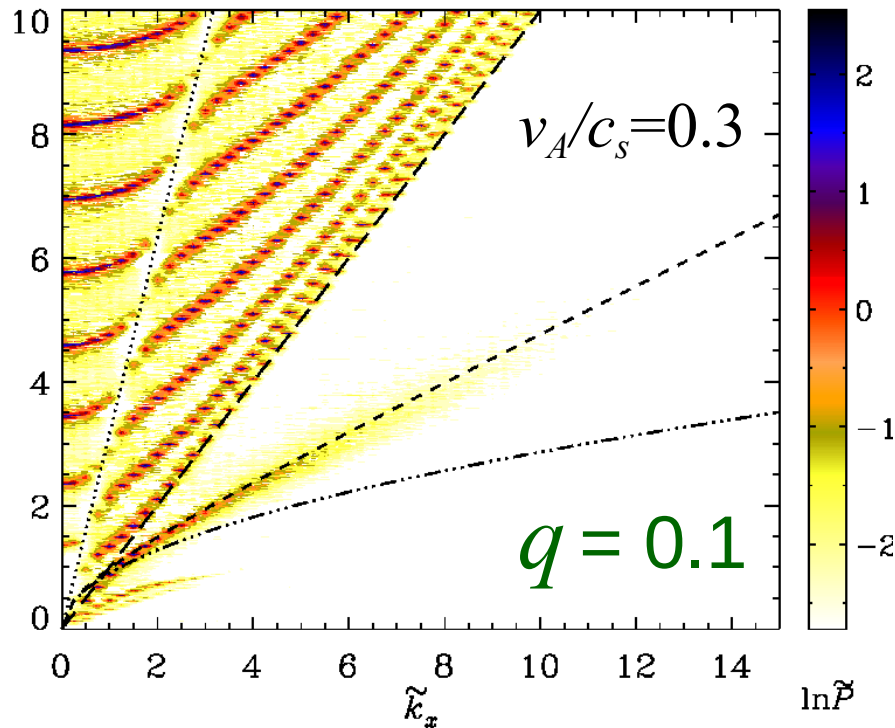
$$n = 0, 1, 2, \dots$$

- **Supports quarter waves as discussed in earlier works**

# Horizontal magnetic field

- We impose a uniform horizontal magnetic field in the entire domain

- Alfven speeds in the two subdomains:  $v_{Ax\ d,u} = \frac{B_x}{\sqrt{\mu_0 \rho_{d,u}}}$



- Phase speed at a magnetic interface (no gravity):

$$c_{fm}^2 = \frac{\rho_u v_{Axu}^2 + \rho_d v_{Axd}^2}{\rho_u + \rho_d}$$

$$= \frac{2\rho_u v_{Axu}^2}{\rho_d + \rho_u} = \frac{2\rho_d v_{Axd}^2}{\rho_d + \rho_u}$$

Kruskal & Schwarzschild 1954  
Miles & Roberts 1989

- Dispersion relation in the presence of gravity:

$$\omega_{fm}^2 = c_{fm}^2 k_x^2 + g k_x \frac{1 - q}{1 + q}$$

# Vertical/oblique magnetic field

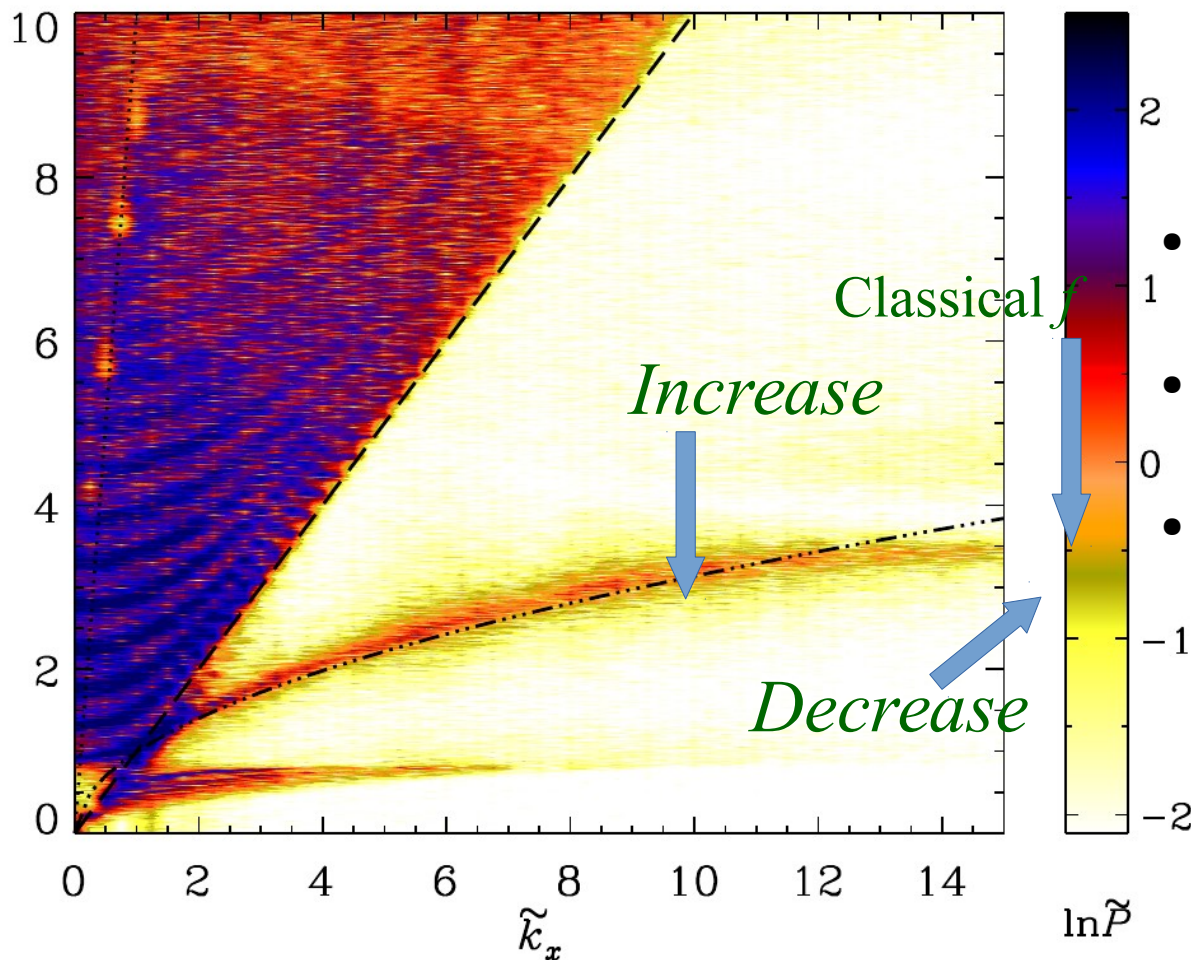
- We impose a uniform magnetic field in the entire domain
- Alfvén speeds (vertical component) in the two subdomains

$$v_{Azd,u} = \frac{B_{z0}}{\sqrt{\mu_0 \rho_{d,u}}}$$

Vertical

$$q = 0.01^{23}$$

$$v_A/c_s = 0.12$$



Some discussions:

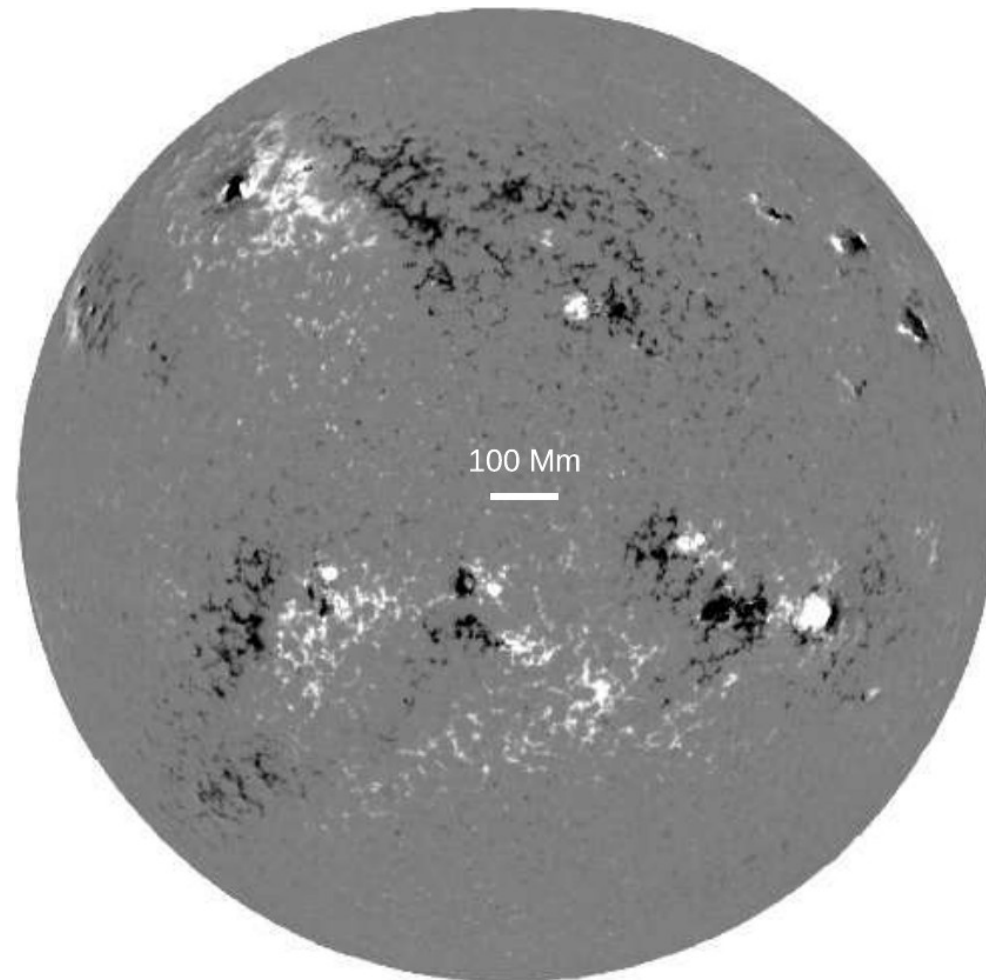
- Libbrecht et al (90)
- Fernandes et al (92)
- Hindman & Zweibel (94)



# *Non-uniform magnetic field*

## *Fanning of the f-mode*

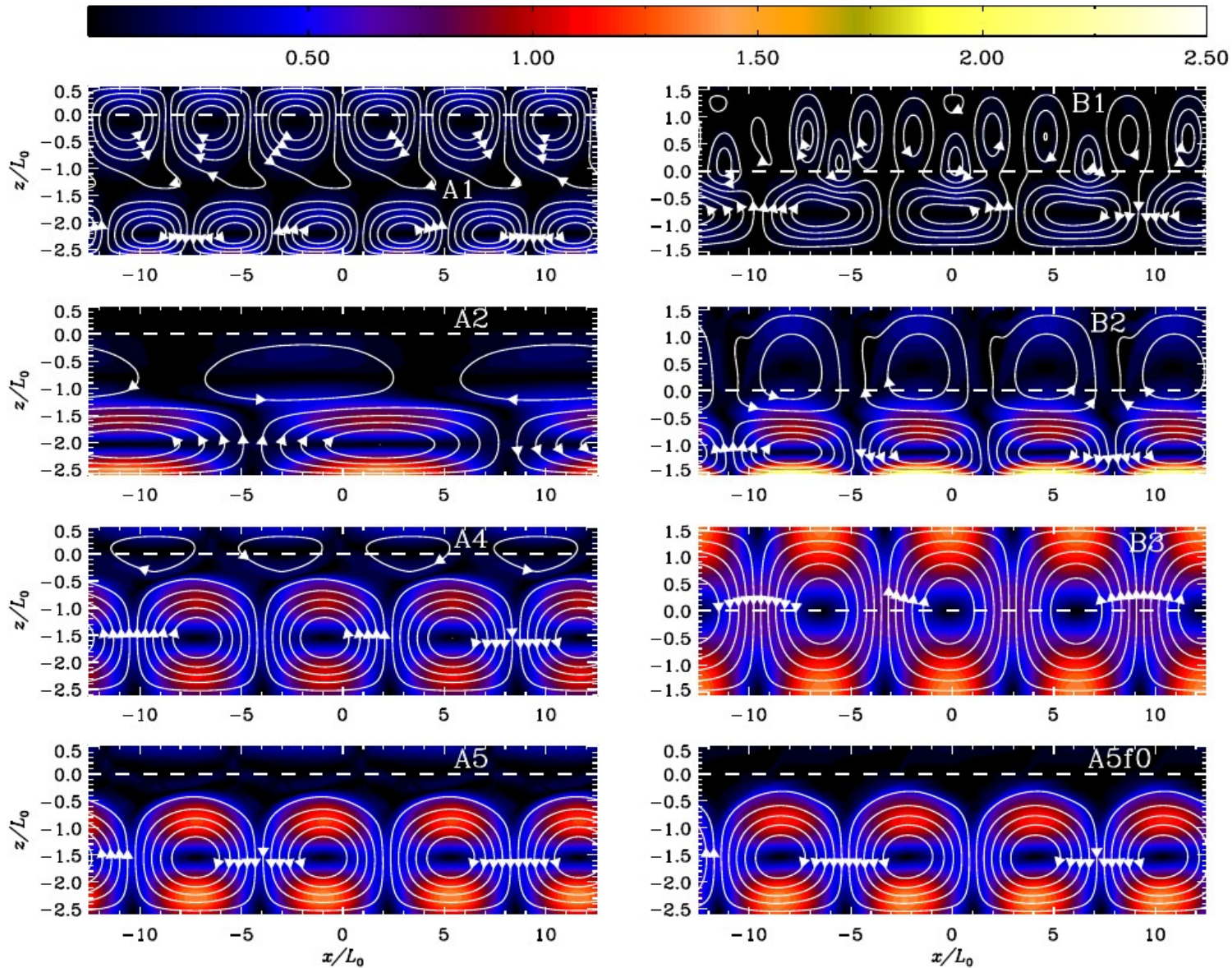
*Seeking motivation from an active phase of the Sun*



Magnetogram (MDI)

# Background magnetic field: nonuniform

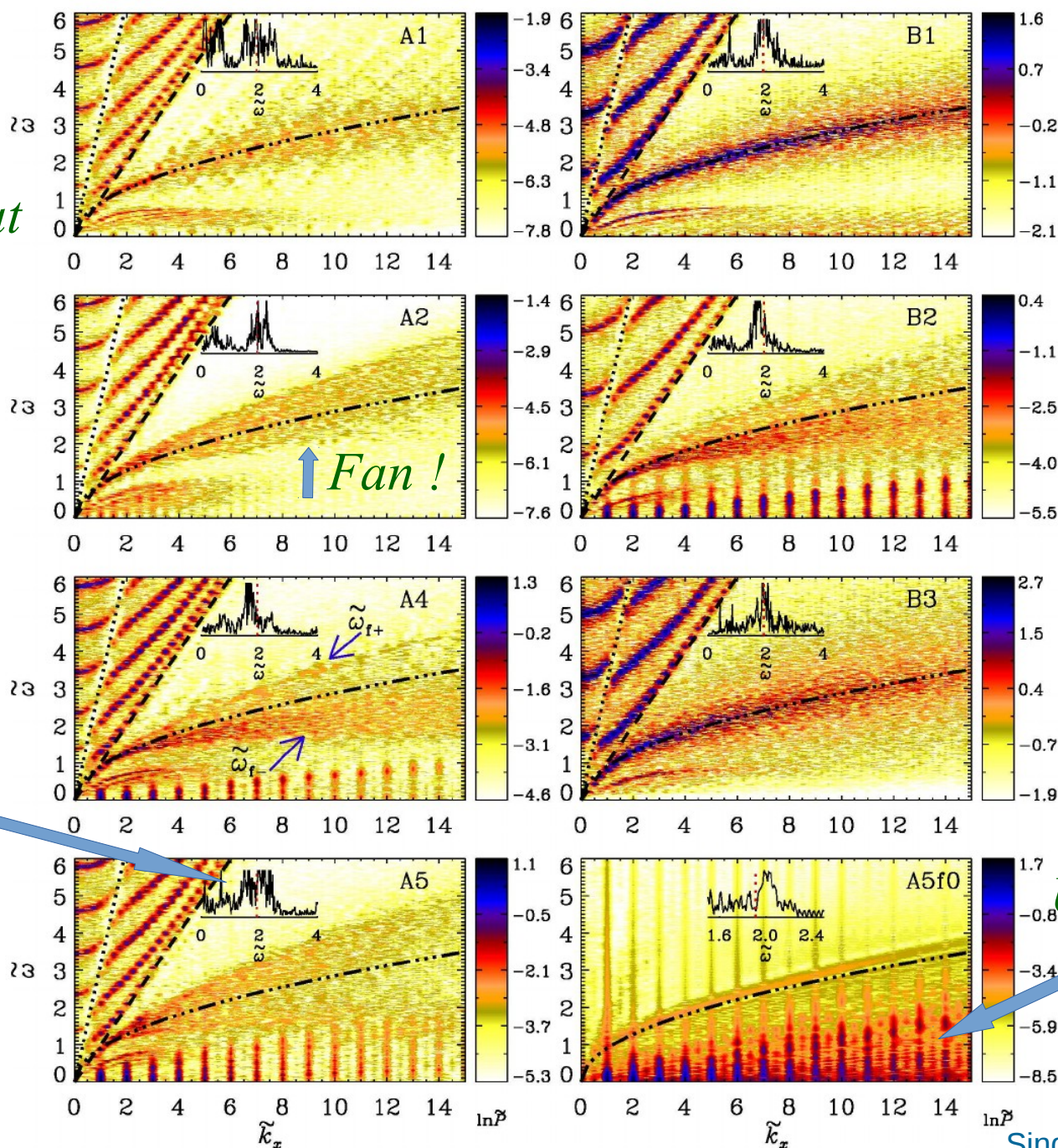
- Consider harmonically varying magnetic field in space
- Force the induction equation with  $\mathcal{E}_{0y} = \hat{\mathcal{E}}_0 \cos(k_x^B x) \cos(k_z^B z)$



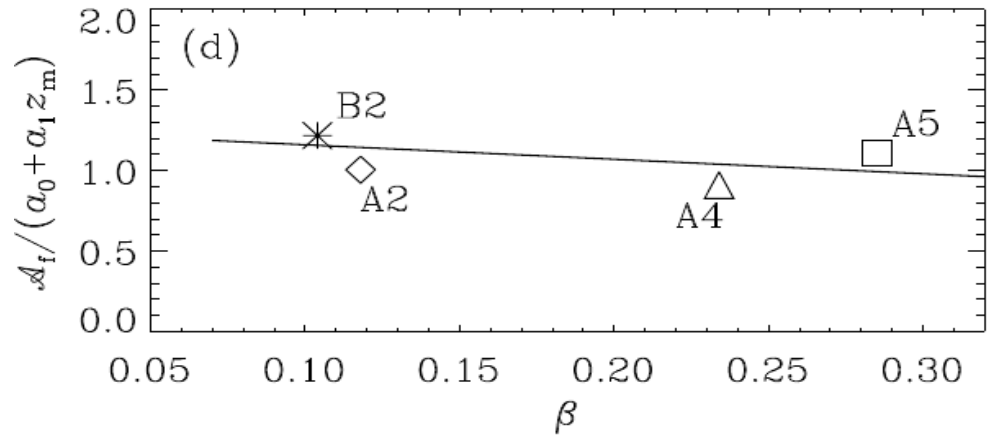
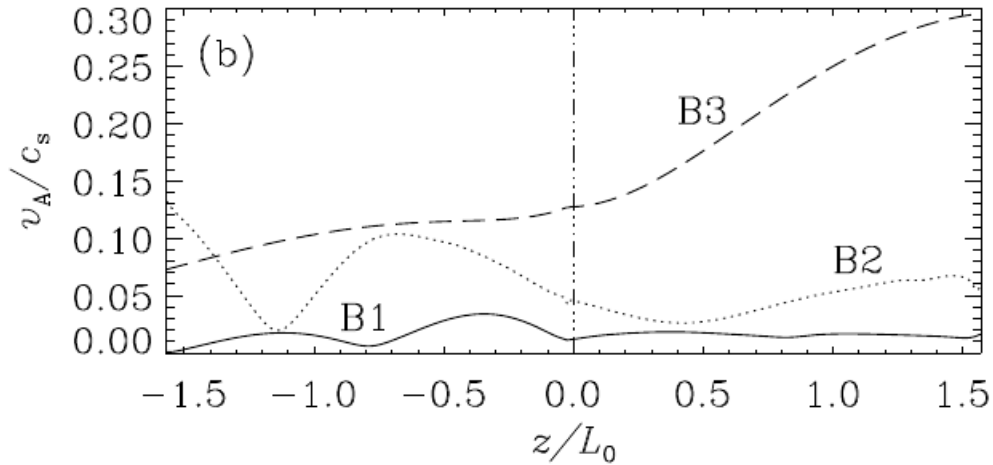
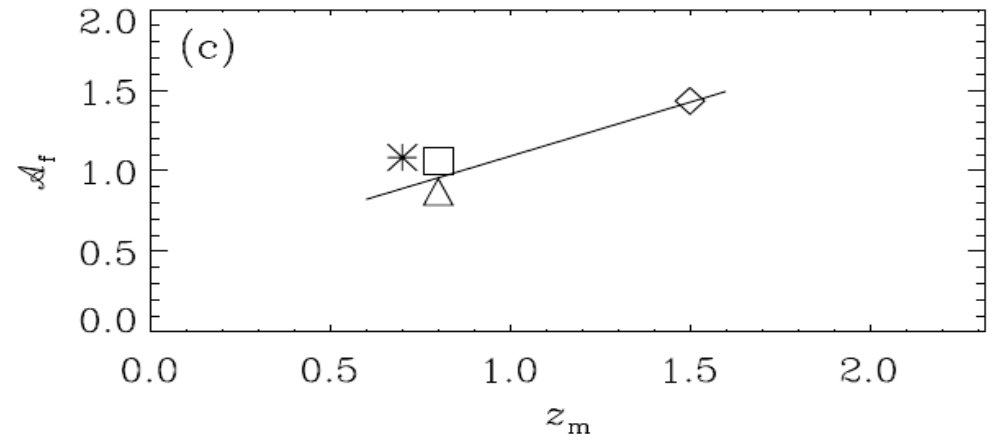
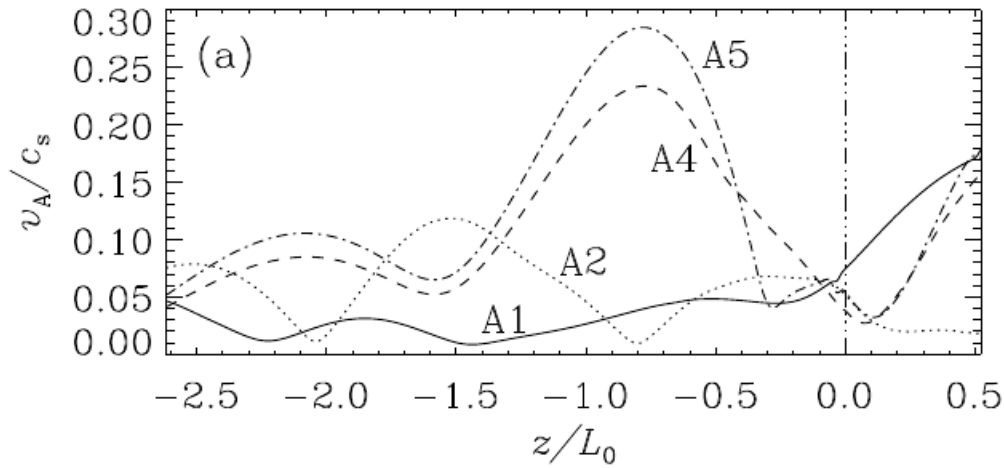


# Fanning out of the $f$ -mode: a precursor transient?

Inset:  
 $f$ -mode profile at  
 $\tilde{k}_x = 4.5$



# Subsurface concentration and asymmetry

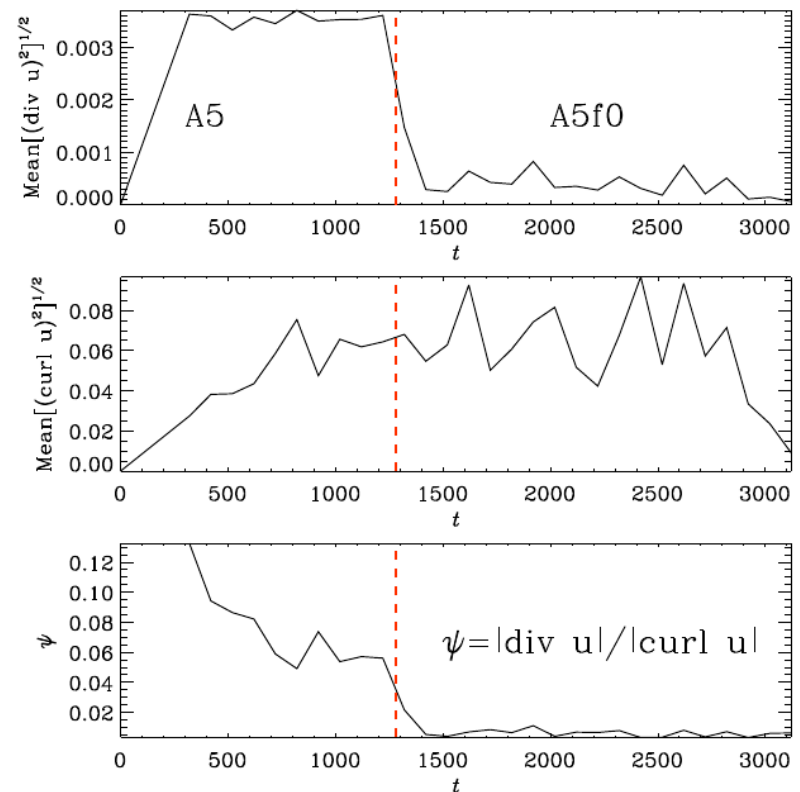


$$v_A(z) = \sqrt{\frac{\langle B^2 \rangle_x(z)}{\mu_0 \rho(z)}}, \quad \beta = \left( \max_{z \leq 0} v_A(z) \right) / c_{sd}$$



# Vertical stripes: unstable eigenmodes

- Persistent after switching off hydrodynamic forcing (A5f0)
- Can be identified to indicate at least one *unstable* eigenmode
- The ratio  $|\nabla \cdot \mathbf{u}| / |\nabla \times \mathbf{u}| \rightarrow$  everywhere small for A5f0
- Velocity field of the stripes thus close to *solenoidal*
- Their occurrence and amplitude depend strongly on  $B_0$
- Notice the difference of *f*-mode in A5 and A5f0: *Fan disappears!*



# Implications for the Sun

- For the Sun :  $\gamma H_d \approx 0.5 \text{ Mm}$

$H_d$  is the pressure scale height

- The fan is best observed at

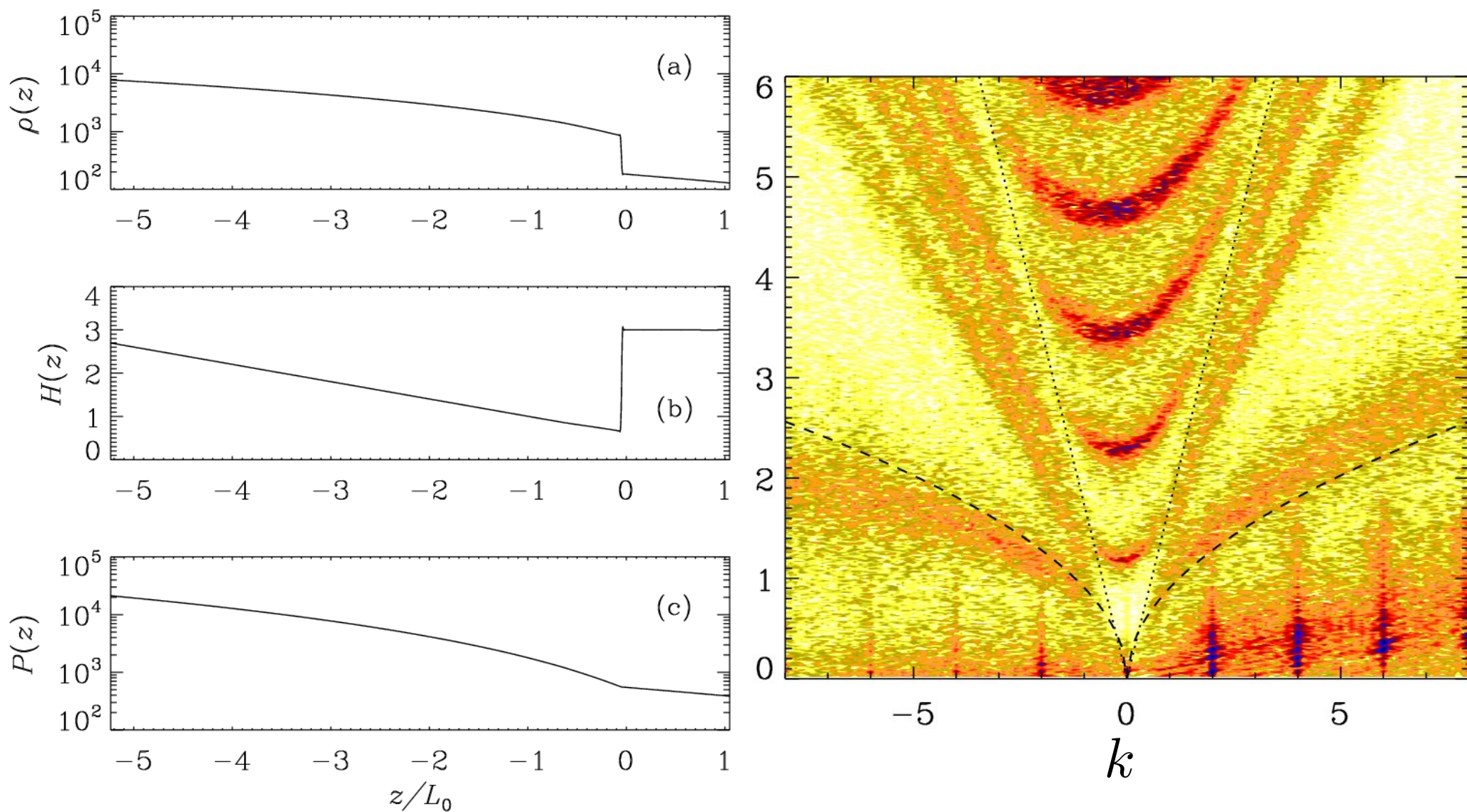
$$\tilde{k}_x \gtrsim 3 \text{ and } 1 \lesssim \tilde{\omega} \lesssim 3$$

- For Sun with radius  $R = 700 \text{ Mm}$ , this corresponds to

$$k_x \gtrsim 6 \text{ Mm}^{-1}, \text{ or } \ell > 4200$$

- Cadence requirement : less than a minute
- Next: Do more realistic models help?

# Isentropic stratification with sub-surface shear



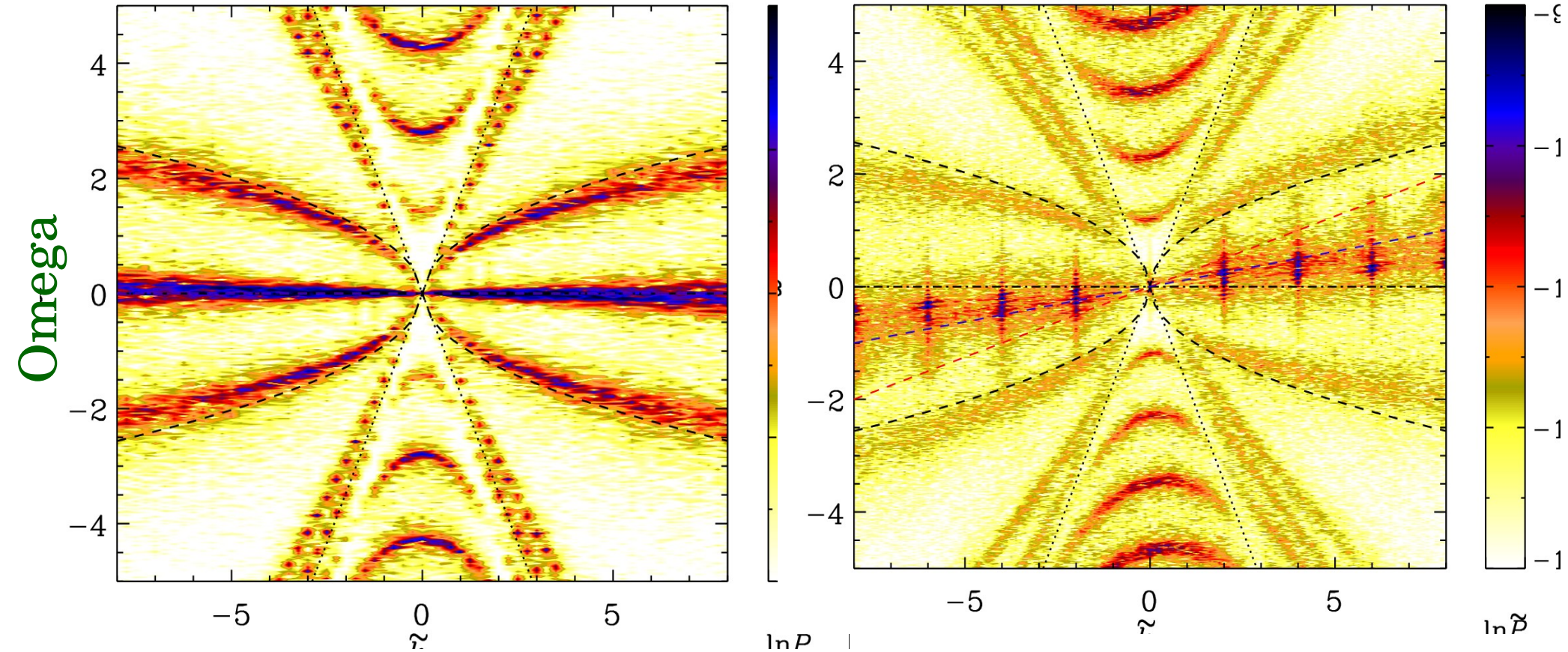
**Notice:** asymmetry of the  $f$ -mode, fan and  $\alpha$ -mode



# Preliminary experiments (with/without shear)

Shear

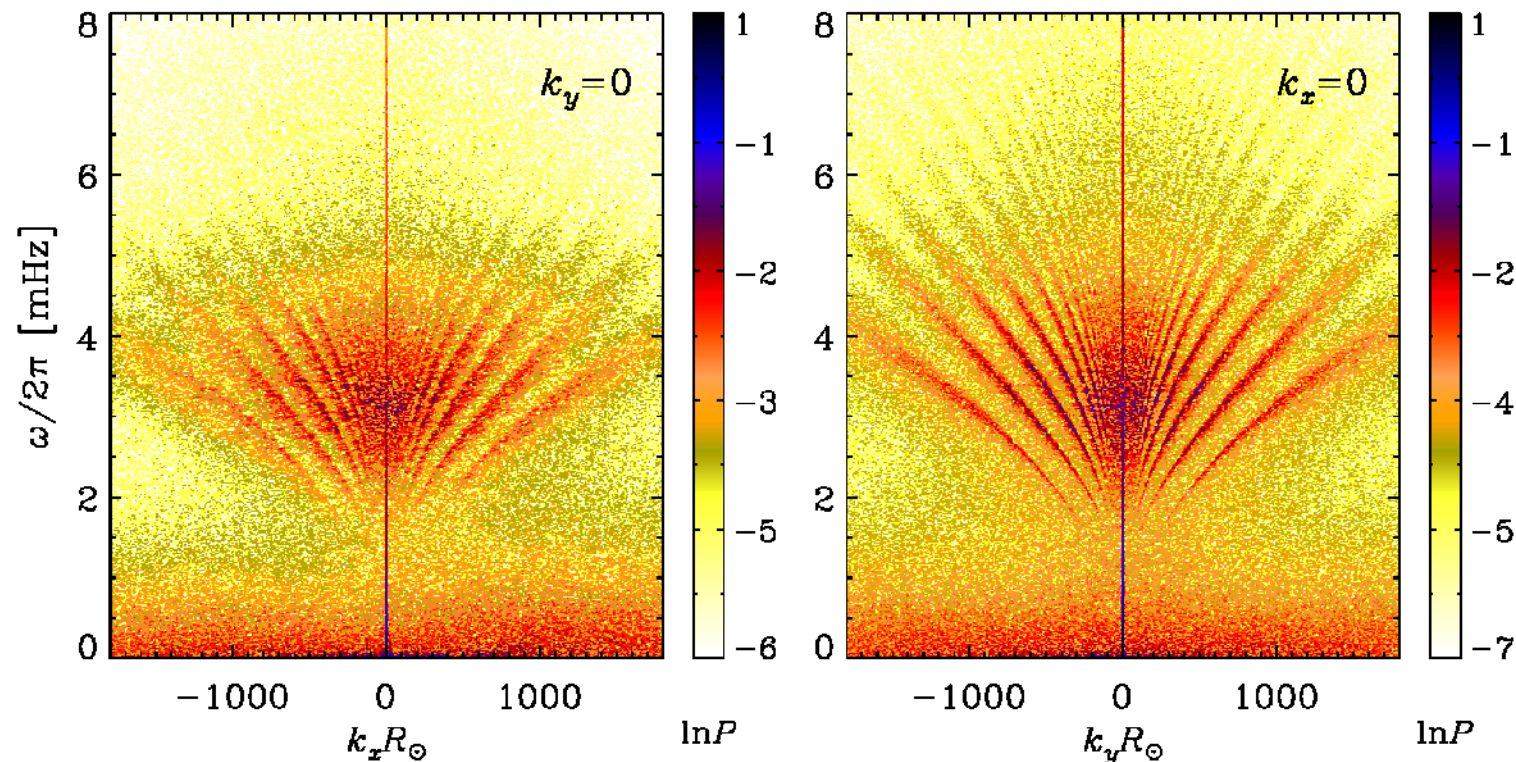
No Shear



Horizontal wavenumber



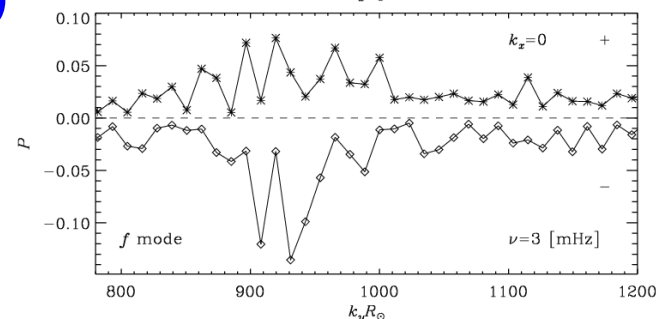
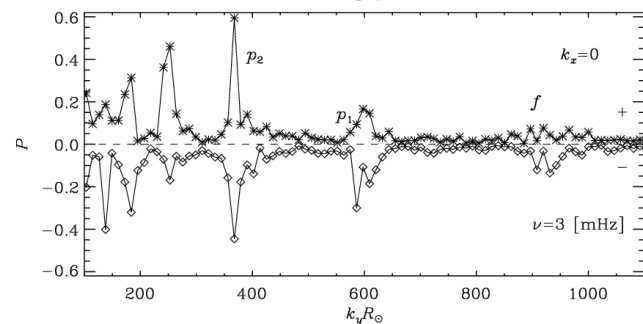
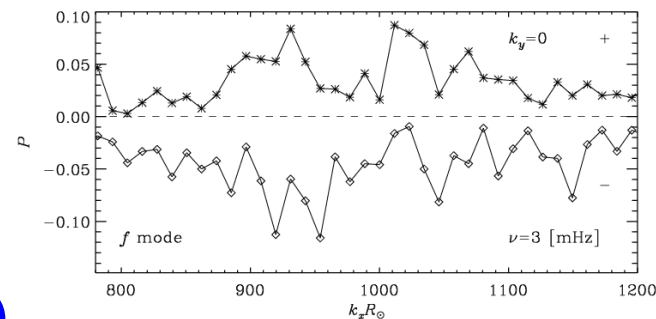
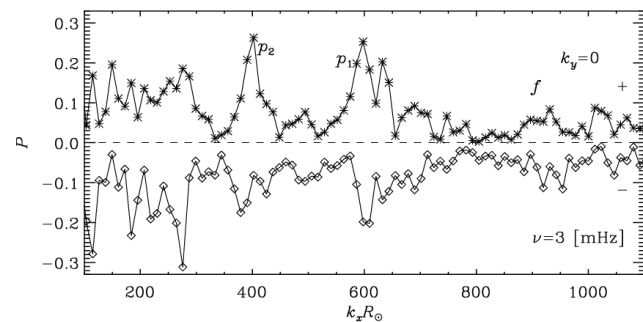
# Preliminary analysis using HMI data



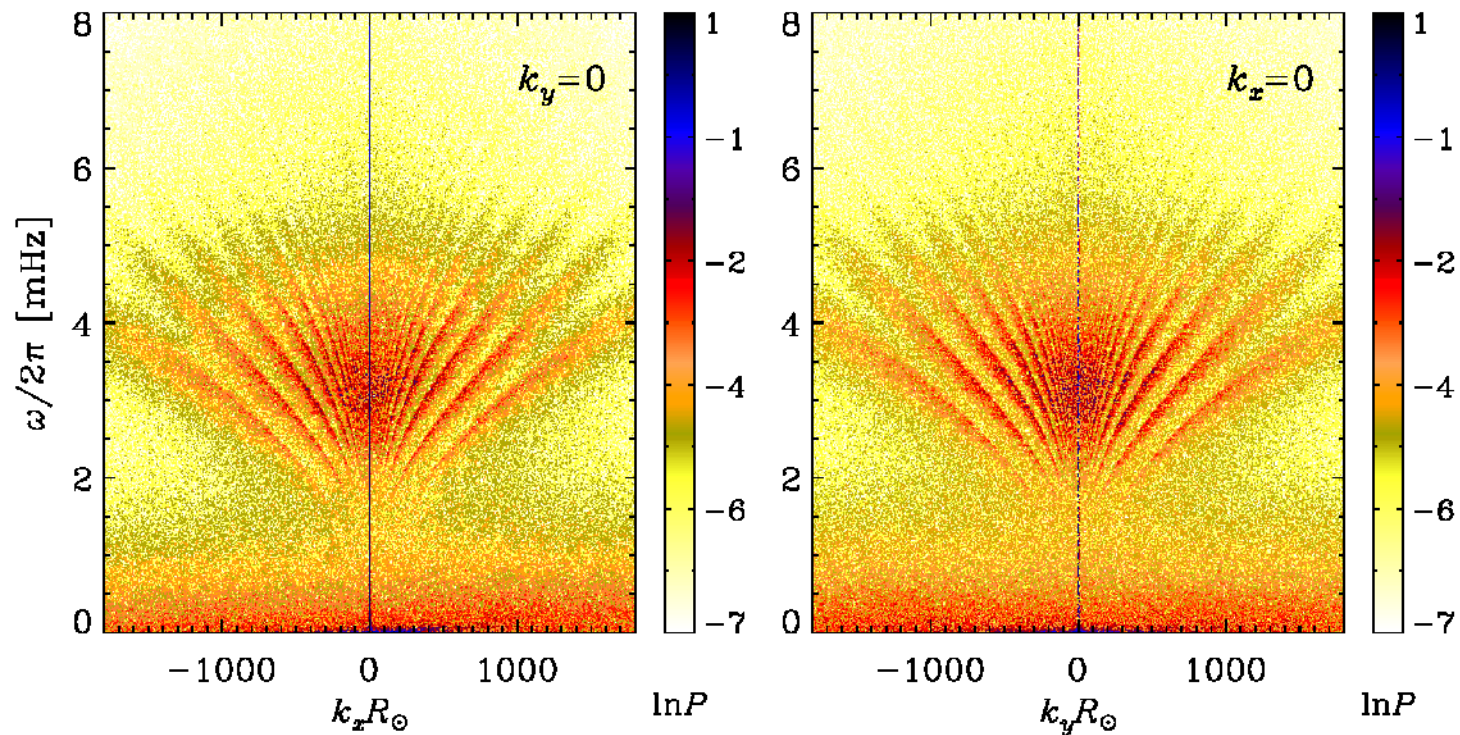
1024 x 1024 patch

24 hours of data (2010)

Line profiles at 3 mHz



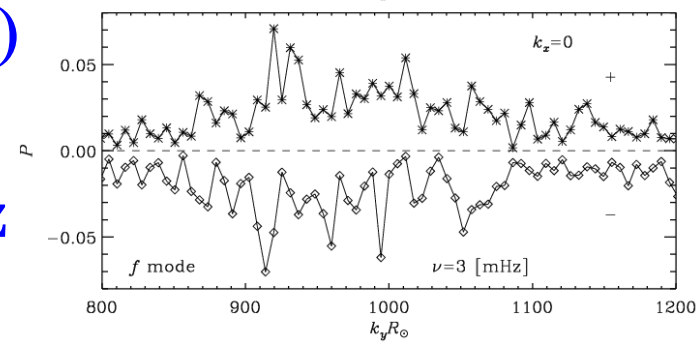
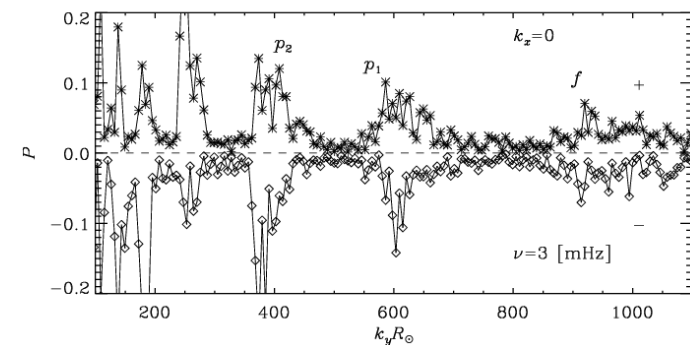
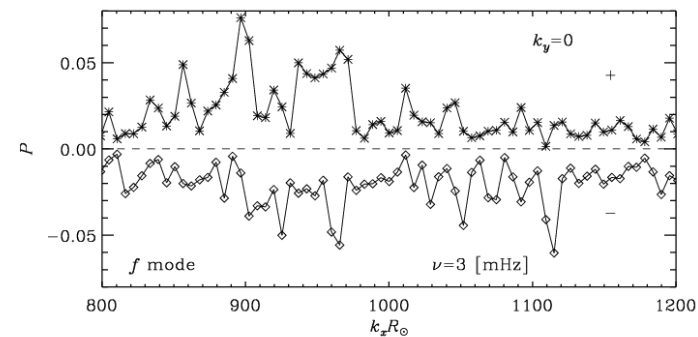
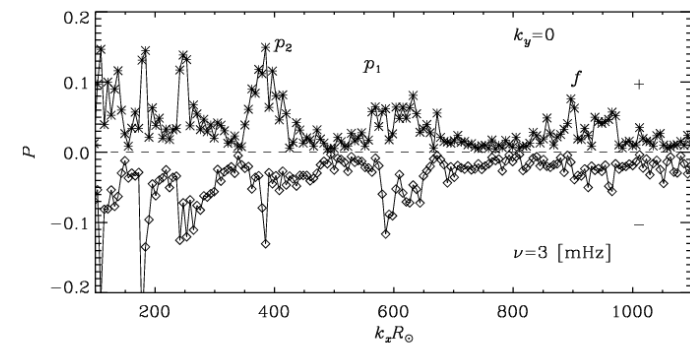
# Preliminary analysis using HMI data



2048 x 2048 patch

24 hours of data (2010)

Line profiles at 3 mHz



# Conclusions

- (Uniform) Horizontal and vertical imposed magnetic fields have distinguishing effects
- Goal: To identify diagnostic signatures of nonuniform  $\mathbf{B}$ -field
- Fanning out of the  $f$ -mode and pattern of vertical stripes seen
- Fanning and its asymmetry depend on the strength and the location of field concentration beneath the interface

*Measurements of such fanning and its asymmetry with respect to the classical  $f$ -mode have the potential to determine the subsurface flux concentrations in the Sun*



*Precursors to the flux emergence at the surface ?*