



# Phase Transitions in Constraint Satisfaction Problems

Lenka Zdeborová (LPTMS, Orsay, France)



In collaboration with:

- F. Krzakala (ESPCI Paris)
- G. Semerjian (ENS Paris)
- A. Montanari (Standford & ENS)
- F. Ricci-Tersenghi (La Sapienza Roma)

#### **Example of CSP: Coloring of random graphs**

#### **Coloring = Antiferromagnetic Potts model at zero temperature**

N=100 vertices, M=218 edges, average degree c=2M/N=4.36



# Why is random graph coloring interesting?

- Colorable threshold at average degree  $C_S$ 
  - a.s. colorable for  $c < c_s$  and a.s. uncolorable for  $c > c_s$
  - (A part of ) Proof of existence (Friedgut 1997, Achlioptas, Friedgut, 1999)
  - Exact (but not rigorous) location via cavity method (Mezard, Zecchina, 2002 + Mulet, Pagnani, Weight, Zecchina, 2002)
- Computationally hard region near to the colorable threshold

Why are the typical instances near to the threshold hard?

Is there a way how to make them easy?

#### **Insight From Statistical Physics of Spin Glasses**

**Prediction of a glassy (clustered) phase in the colorable region** Mezard, Zecchina, Parisi, 2002, Biroli, Monasson, Weigt, 1999)

Clusters (roughly said): groups of nearby solutions which are in some sense disconnected from each other (more precisely later in the presentation).

Suggestion: Clustering responsible for the onset of hardness.

Accomplishments (Mezard, Zecchina, Parisi, 2002):

 $\checkmark$  The exact colorable threshold computed.

✓ Survey Propagation algorithm designed.



#### Loose ends to be tied up

- Computed value of the clustering transition did not correspond to the empirically obtained values of the onset of hardness (linear performance of Walk-SAT etc.)
- Clustering transition continuous or discontinuous?
- Frozen variables present in clusters or not? (Frozen variable takes the same color in all the solutions belonging to the cluster.)
- SP works with frozen variables, but finds solutions without them.



## What do we mean by clusters?

- Roughly said: Lumps (groups) of nearby solutions which are in some sense disconnected from each other.
- For mathematical physicist: "Extremal Gibbs measures = pure states".
- For computer scientist: Fixed points of belief propagation.
- For spin glass physicist: Solutions of TAP equations.

#### Why do we need to speak about clusters above $C_d$ ?

- The Gibbs measure cease to be extremal. (Mézard, Montanari, 2005)
- The point-to-set correlations do not decay.
- Different from: A non-trivial solution of the survey propagation exists!

### **Refining the structure of clusters**

- Entropy (size) of a cluster s: logarithm of the number of solutions belonging to the cluster (divided by the number of variables).
- Complexity function  $\Sigma(s)$ : logarithm of the number of clusters of size s (divided by the number of variables).
- The zero temperature entropic 1RSB cavity method allows us to compute the Legendre transform  $\Phi(m)$  of  $\sum(s)$ . Main idea (Mézard, Palassini, Rivoire, 2005) weight each cluster by its size to the power m:

$$e^{N\Phi(m)} = \sum_{\alpha} (e^{Ns_{\alpha}})^m = \int e^{N[ms + \Sigma(s)]} ds$$
$$\Phi(m) = ms + \Sigma(s), \quad \frac{\partial \Sigma(s)}{\partial s} = -m$$

Note: the approach of Mézard, Zecchina, Parisi 2002; Mulet, Pagnani, Weigt, Zecchina 2002 was at m=0.

# Solve (mostly numerically) the 1RSB cavity equations

+ Work out the several special cases when the equations simplify (m=1, m=0, frozen variables, regular graphs ...)





### **Our results (+ their meaning)**

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# Learning from $\Sigma(s)$

Example of 6-coloring, regular graphs, degrees 17, 18, 19, 20









very low degree

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degree c=17

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degree c=18











degree c=19

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degree c=20

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## The important phase transitions: Overview



 $\therefore$  Clustering (dynamical) transition  $c_d(3) = 4, c_d(4) = 8.35, c_d(5) = 12.84$ 

- Glassy solution appears at m=1.
- The replica symmetric entropy still correct, no non-analyticity.
- Entropy dominated by exponentially many states.
- $\therefore$  Condensation (Kauzmann) transition  $c_c(3) = 4$ ,  $c_c(4) = 8.46$ ,  $c_c(5) = 13.23$ 
  - Discontinuity in the second derivative of entropy.
  - Entropy dominated by finite number of the largest states.
- $\therefore$  Rigidity transition  $c_r(3) = 4.66, c_r(4) = 8.83, c_r(5) = 13.55$ 
  - Frozen variables appears in the dominating states.
  - Frozen clusters disappear abruptly.
  - Minimal rearrangements diverge (Semerjian, 2007).

Moreover: The entropically dominating clusters are 1RSB stable in the colorable phase (at least for q>3)

Colorable threshold:  $c_s(3) = 4.69, c_s(4) = 8.90, c_s(5) = 13.67$ 

#### Large number of colors (analytical results)



#### **Algorithms beyond clustering transition**

- Hard to uniformly sample solutions with Monte Carlo (equilibration time diverges).
- > Not necessarily hard to find solutions.
- Belief Propagation gives correct marginals up to the condensation transition. Walk-COL works in linear time (empirical) beyond clustering (clusters without frozen variables cannot "end" at positive energy).



Is the rigidity more significant for hardness?

**Two arguments:** 

- All the solutions we are able to find on large graphs (N>10 000) belong to clusters without frozen variables.
- Minimal rearrangments (Semerjian, 2007) diverge if and only if frozen variables present. Local search should fail.

### Conclusions

- The dynamical transition redefined with respect to the dominating clusters and located.
- Final Provide the second secon
- Algorithmic consequences understood better, but more work have to be done in this direction.

### References

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