Consequences of γ -softness of the core for some properties of chiral nuclei

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Outline

- Introduction
- Core-Particle-Hole-Coupling (CPHC) model
- S symmetry and its consequences
- Nonsymmetric cases
- Microscopic cores

Collaborators: Ch. Droste, S. G. Rohoziński and K. Starosta EPJ A 47 (2011) 90

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EPJ A 42 (2009) 72 + E. Grodner
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Characteristic features of chiral nuclei Partner bands:

- ,,almost" degenerate levels (200 – 300 keV) with the same parity, $\Delta I = 1$
- similar EM properties inside bands
- staggering of M1 transitions (also E2 with ΔI = 1)



Introduction

Experimental stimulus: ¹²⁸Cs, ¹²⁶Cs, ¹²⁴Cs



771 % (4*) / 543 % 779 % (4*) / 484 % (4*) / 543 % 779

278

(194

(15+)

(16+)

- E. Grodner et al. 128Cs as the Best Example Revealing Chiral Symmetry Breaking, PRL 97, 172501 (2006)
- E. Grodner et al. Partner bands of 126Cs first observation of chiral electromagnetic selection rules, PLB 703, 46 (2011)

^{126–128}Cs nuclei, EM transitions



Why γ -softness of the core



P. Moller et al., *Global Calculations of Ground-State Axial Shape Asymmetry of Nuclei*, PRL 97, 162502 (2006)

$$\begin{aligned} |\text{odd} - \text{odd}, i\rangle &= \sum_{r,k,m} U^{i}_{rkm} |\text{core}, r\rangle |\text{p}, k\rangle |\text{n}, m\rangle \\ H_{\text{o-o}} &= H_{\text{core}} - \chi Q q_{\text{p}} - \chi Q q_{\text{n}} - \chi' q_{\text{p}} q_{\text{n}} + h_{\text{p}} + h_{\text{n}} \\ Q, q_{\text{p,n}} &= \text{quadrupole operators} \end{aligned}$$

Cores:

- Bohr Hamiltonian (FQ, full quadrupole); (β, γ, Ω) or (α_{μ})
 - given by analytic formulas
 - calculated microscopically
- rigid rotor (DF, Davydov-Filipov); (Ω); [$\tilde{\beta}, \tilde{\gamma}$ fixed]

Proton-neutron part:

 $ph_{11/2} \otimes nh_{11/2}^{-1}, A \sim 130$

The Bohr Hamiltonian

$$H_{\text{GBH}}(\beta,\gamma,\Omega) = T_{\text{vib}} + T_{\text{rot}} + V$$

$$T_{\text{vib}}(\beta,\gamma) = -\frac{1}{2\sqrt{wr}} \left\{ \frac{1}{\beta^4} \left[\partial_\beta \left(\beta^4 \sqrt{\frac{r}{w}} B_{\gamma\gamma} \right) \partial_\beta - \partial_\beta \left(\beta^3 \sqrt{\frac{r}{w}} B_{\beta\gamma} \right) \partial_\gamma \right] + \frac{1}{\beta \sin 3\gamma} \left[-\partial_\gamma \left(\sqrt{\frac{r}{w}} \sin 3\gamma B_{\beta\gamma} \right) \partial_\beta + \frac{1}{\beta} \partial_\gamma \left(\sqrt{\frac{r}{w}} \sin 3\gamma B_{\beta\beta} \right) \partial_\gamma \right] \right\}$$

$$T_{\text{rot}}(\beta,\gamma,\Omega) = \frac{1}{2} \sum_{k=1}^3 I_k^2(\Omega) / J_k; \quad J_k(\beta,\gamma) = 4B_k(\beta,\gamma)\beta^2 \sin^2(\gamma - 2\pi k/3)$$

$$w = B_{\beta\beta} B_{\gamma\gamma} - B_{\beta\gamma}^2; \quad r = B_x B_y B_z$$

Simplest form of the kinetic energy

$$T_{\rm kin} = \frac{1}{2} B \sum_{\mu} |\dot{\alpha_{\mu}}|^2 \quad \leftrightarrow B_{\beta\beta} = B_{\gamma\gamma} = B_k = B, \ B_{\beta\gamma} = 0$$

 $S = P_5(\text{core})C_{\text{pn}}(\text{proton-neutron})$

FQ core

LAB system

$$\begin{split} P_{5}\alpha_{\mu} &= -\alpha_{\mu} \quad \alpha\text{-parity of the core} \\ \text{INT (principal axes) system} \\ \{\alpha_{\mu}\} &\xrightarrow{R(\Omega)} \{\tilde{\alpha}_{0}, \tilde{\alpha}_{1} = \tilde{\alpha}_{-1} = 0, \tilde{\alpha}_{2} = \tilde{\alpha}_{-2} \} \\ &\tilde{\alpha}_{0} = \beta \cos \gamma, \\ &\tilde{\alpha}_{2} = \tilde{\alpha}_{-2} = \beta \sin \gamma / \sqrt{2} \\ & P_{5}(\beta, \gamma, \Omega) = (\beta, \gamma \pm \pi, \Omega) \\ & P_{5}(\beta, \gamma, \Omega) = (\beta, \pi/3 - \gamma, R_{x}(\pi/2)\Omega) \end{split}$$



DF core

$$P_5 \rightarrow R_x(\pi/2)$$

S symmetry, pn part

Proton-neutron part: C_{pn}

"Almost" the same (naturally isomorphic) spaces for proton and neutron

$$C_{pn}|(\mathbf{p},i)(\mathbf{n},k)\rangle = |(\mathbf{p},k)(\mathbf{n},i)\rangle$$

Proton and neutron at the same j level

$$C_{pn}|((p,j)(n,j))^{L}\rangle = (-1)^{L+1}|((p,j)(n,j))^{L}\rangle$$

General property

$$P_5^2 = I_{core}, \quad C_{pn}^2 = I_{pn}$$

 $S^2 = (P_5 C_{pn})^2 = I$

For the DF core cf. T. Koike, K. Starosta, I. Hamamoto, PRL 93 (2004) 172502

P₅ symmetry. The Bohr Hamiltonian

Conditions imposed by the P_5 symmetry

$$\begin{split} f(\beta, \pi/3 - \gamma) &= f(\beta, \gamma) \quad \text{for } f = B_{\beta\beta}, B_{\gamma\gamma}, B_x \\ B_{\beta\gamma}(\beta, \pi/3 - \gamma) &= -B_{\beta\gamma}(\beta, \gamma) \\ B_y(\beta, \pi/3 - \gamma) &= B_z(\beta, \gamma) \\ V(\beta, \pi/3 - \gamma) &= V(\beta, \gamma) \end{split}$$



Quadrupole operators and angular momenta

$$P_{5}Q_{\nu}P_{5}^{+} = -Q_{\nu}$$

$$P_{5}I_{k}P_{5}^{+} = I_{k}$$

$$C_{pn} q_{\nu,p,n} C_{pn}^{+} = -q_{\nu,n,n}$$

$$C_{pn} j_{k,p,n} C_{pn}^{+} = j_{k,n,p}$$

Odd-odd nucleus Hamiltonian $H_{o-o} = H_{core} - \chi Q q_p - \chi Q q_n - \chi' q_p q_n$

$$H_{\text{core}} \text{ is } P_5 \text{-symmetric} \longrightarrow SH_{\text{o-o}}S^+ = H_{\text{o-o}}$$

New additional label $s = \pm 1$ for eigenstates of H_{o-o}

S symmetry. E2 & M1 transition operators

E2

$$\begin{split} T(E2) &= \kappa Q + e_{\mathrm{p}}q_{\mathrm{p}} + e_{\mathrm{n}}q_{\mathrm{n}} \\ SQS^{+} &= -Q \\ e_{\mathrm{p}}q_{\mathrm{p}}, e_{\mathrm{n}}q_{\mathrm{n}} &\longrightarrow \forall J, s | T(E2) | J', s \rangle \approx 0 \end{split}$$

M1

$$T(M1) = \sqrt{\frac{3}{4\pi}} \mu_N (g_I I_{\text{core}} + g_p j_p + g_n j_n)$$

$$P_5 I_{\text{core}} P_5^+ = I_{\text{core}}$$

$$g_I - (g_p + g_n)/2 = 0 \qquad \longrightarrow \langle J, s | T(M1) | J', s \rangle = 0, \quad J \neq J'$$

Example (reasonable values): $g_I = 0.44$, $g_p = 1.22$, $g_n = -0.21$

$$T(M1) \sim (g_I J + (g_{lp} - g_R)l_p + (g_{sn} - g_I)s_p) + (g_{ln} - g_R)l_n + (g_{sp} - g_I)s_n)$$

Selection rules



Selection rules



Schematic core Hamiltonians

- Simple kinetic energy, B const
- Potential energy

$$V = \frac{1}{2}V_C\beta^2 + [G + h_1\cos 3\gamma + h_2(\cos^2 3\gamma - 1)^{\kappa}] \times (\exp(-\beta^2/d^2) - 1)$$

Properties

$$\begin{split} & E(2^+_1) \approx 0.35 \text{ MeV} \\ & \beta_{\rm eq} \approx 0.25 \\ & B(E2,2_1 \rightarrow 0_1) \approx 0.28 \text{ eb}^2 \end{split}$$

S symmetric Hamiltonians

Wilets-Jean, potential well (PW), potential barrier (PB). Simple kinetic energy



Chiral bands in nuclei, NORDITA, Stockholm, April 2015

S symmetric cases, energies, E2



S symmetric cases, intraband transitions

M1 transitions





S symmetric cases, $s \rightarrow g$ transitions



Deviations from the S symmetry. Non-symmetric core

- Non-symmetric core (potential or/and kinetic energy)
- Different one-particle spaces for proton and neutron



Case 1. Non-symmetric potential energy

Non-symmetric core



Non-symmetric core, cont.

 $g \rightarrow g$ transitions



$s \rightarrow g$ transitions

Deviations from the S symmetry. Different orbitals for proton and neutron

Case 2. Proton $h_{11/2}$; neutron $g_{9/2}^{-1}$ or $f_{7/2}^{-1}$ core — WJ potential (P_5 symmetric)



Different orbitals for proton and neutron, cont.

Intraband transitions

$s \rightarrow g$ band transitions



$${}^{128}\text{Xe} + p \otimes n^{-1} \longrightarrow {}^{128}\text{Cs}$$
$${}^{128}\text{Ba} + p \otimes n^{-1} \longrightarrow {}^{128}\text{La}$$

- Mean field: Skyrme SIII and SLy4 RMF NL3
- Pairing: seniority force
- Calculations with constraints to get required deformation
- Kinetic energy ATDHFB type formulas

Deformation variables

$$\begin{split} \beta \cos \gamma &= cq_0 = c \langle \Psi | Q_0 | \Psi \rangle = \langle \Psi | \sum_{i=1}^A (3z_i^2 - r_i^2) | \Psi \rangle \\ \beta \sin \gamma &= cq_2 = c \langle \Psi | Q_2 | \Psi \rangle = \langle \Psi | \sum_{i=1}^A \sqrt{3} (x_i^2 - y_i^2) | \Psi \rangle; \quad c = \sqrt{\pi/5} / A \overline{r^2} \end{split}$$

HFB with constraints

$$\begin{split} &\delta \langle \Psi | H_{\text{micr}} - \lambda_0 Q_0 - \lambda_2 Q_2 | \Psi \rangle = 0 \\ &\langle \Psi | Q_0 | \Psi \rangle = q_0, \quad \langle \Psi | Q_2 | \Psi \rangle = q_2 \quad \longrightarrow \Psi(\beta, \gamma) \end{split}$$

Potential energy

$$V = \langle \Psi(\beta, \gamma) | H_{\text{micr}} | \Psi(\beta, \gamma) \rangle$$

ATDHFB, some formulas, kinetic energy

HF+BCS case

Then

Mass parameters (inertial functions)

$$B_{q_i q_j} = \hbar^2 (S_{(1)}^{-1} S_{(3)} S_{(1)}^{-1})_{ij}$$

$$(S_{(n)})_{ij} = \sum_{\mu,\nu} \frac{\langle \mu | Q_i | \bar{\nu} \rangle \langle \bar{\nu} | Q_j | \mu \rangle}{(E_\mu + E_\nu)^n} (u_\mu v_\nu + u_\nu v_\mu)^2$$
$$B_{q_i q_j} \to B_{\beta\beta}, B_{\beta\gamma}, B_{\gamma\gamma}$$

Moments of inertia

$$J_k = \hbar^2 \sum_{\mu,\nu} \frac{|\langle \nu | j_k | \bar{\mu} \rangle|^2 (u_\mu v_\nu - u_\nu v_\mu)^2}{(E_\mu + E_\nu)}$$

Microscopic cores, potential energy









Inertial functions, ¹²⁸Xe, Skyrme III



Odd-odd nucleus with $^{\rm 128}{\rm Xe}$ core



Odd-odd nucleus with ¹²⁸Xe core cont.

 $\Delta I = 1$ M1 and E2 transitions

In-band Inter-band M1, $g \rightarrow g$ Xe128, SIII Xe128, SLy4 M1, $s \rightarrow g$. 0 2 $\mathsf{B}(\mathsf{M1},\mathsf{I}_g\!\rightarrow\!(\mathsf{I}\!\!\cdot\!\!1)_g)\,[\mu_{\mathsf{N}}^{\ 2}]$ Xe128, NL3 $\mathsf{B}(\mathsf{M1},\mathsf{I}_{s}\!\rightarrow\!(\mathsf{I}\text{-}1)_{g})\left[\mu_{\mathsf{M}}\right]$ 2 1 1 10 14 15 Ι_g [ħ] 16 17 18 19 20 10 11 14 15 【g[九] 9 11 12 13 12 13 16 0.8 0.8 $\begin{array}{c} \mathsf{B}(\mathsf{E2},\mathsf{I}_g \to (\mathsf{I-1})_g) \; [\mathsf{e}^2 \mathsf{b}^2] \\ \mathsf{F0} \; & \mathsf{F0} \\ \mathsf{F1} \; & \mathsf{F1} \\ \mathsf{F1} \; & \mathsf{F1} \; \mathsf{F1} \\ \mathsf{F1} \; & \mathsf{F2} \; \mathsf{F1} \\ \mathsf{F1} \; & \mathsf{F2} \; \mathsf{F1} \\ \mathsf{F1} \; & \mathsf{F2} \; \mathsf{F2} \; \mathsf{F1} \\ \mathsf{F2} \; & \mathsf{F3} \; \mathsf{F2} \; \mathsf{F1} \\ \mathsf{F3} \; \mathsf{F3}$ E2, $\Delta \mathbf{I} = 1$, $\mathbf{g} \rightarrow \mathbf{g}$ E2, $\Delta I = 1$, s $\rightarrow g$ ۸ Xe128, SIII Xe128, SLy4 $\begin{array}{ll} B(E2,I_{s} \rightarrow (I\text{-}1)_{g}) \, \left[e^{2}b^{2}\right] \\ c & p \\ c & p \\ c & p \end{array}$ Δ Xe128, NL3 10 11 12 14 15 I_g[ħ] 16 17 18 10 11 14 15 9 13 19 20 9 12 13 16 17 $\int_{\sigma} [\hbar]$

Xe128, SIII Xe128, SLy4 Xe128, NL3

Xe128, SIII Xe128, SLy4

Xe128, NL3

19

18 20

18 19 20

Odd-odd nucleus with $^{\rm 128}{\rm Ba}$ core



Odd-odd nucleus with ¹²⁸Ba core cont.

 $\Delta I = 1$ M1 and E2 transitions

In-band



- Existence of the partner bands do not require rigid deformation of the core
- The S symmetry useful in model studies. Not much probable to exist in real nuclei (?)
- It is possible to obtain partner bands with microscopic cores

^{114–144}Xe. Selected energy levels *cont*.



^{114–144}Xe. Selected energy levels *cont*.



$^{114-144}\mbox{Xe}.$ Selected energy levels. Comparison with SLy4 results



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Chiral bands in nuclei, NORDITA, Stockholm, April 2015

E2 transition probabilities



E2 transition probabilities



Comparison with SLy4 results



More details on ¹²⁸Xe. Energy levels

Low energy, positive parity levels in ¹²⁸Xe



More details on ¹²⁸Xe. Energy levels

Low energy, positive parity levels in ¹²⁸Xe

