



Neutron-proton correlation phenomena in N=Z nuclei around ^{100}Sn

- An experimental perspective

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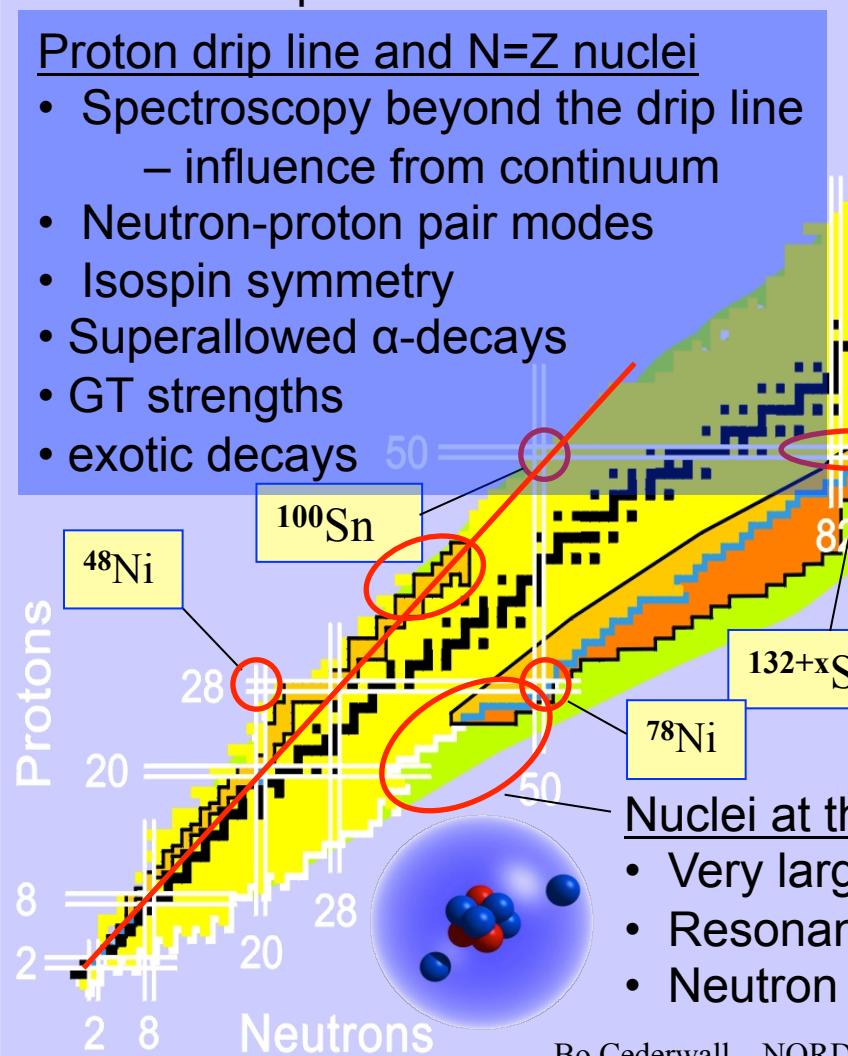
Challenges in Nuclear Structure (selection!)

Shell structure in nuclei (general)

- Structure of doubly magic nuclei
- Changes in the (effective) interactions
- Development of collective excitations

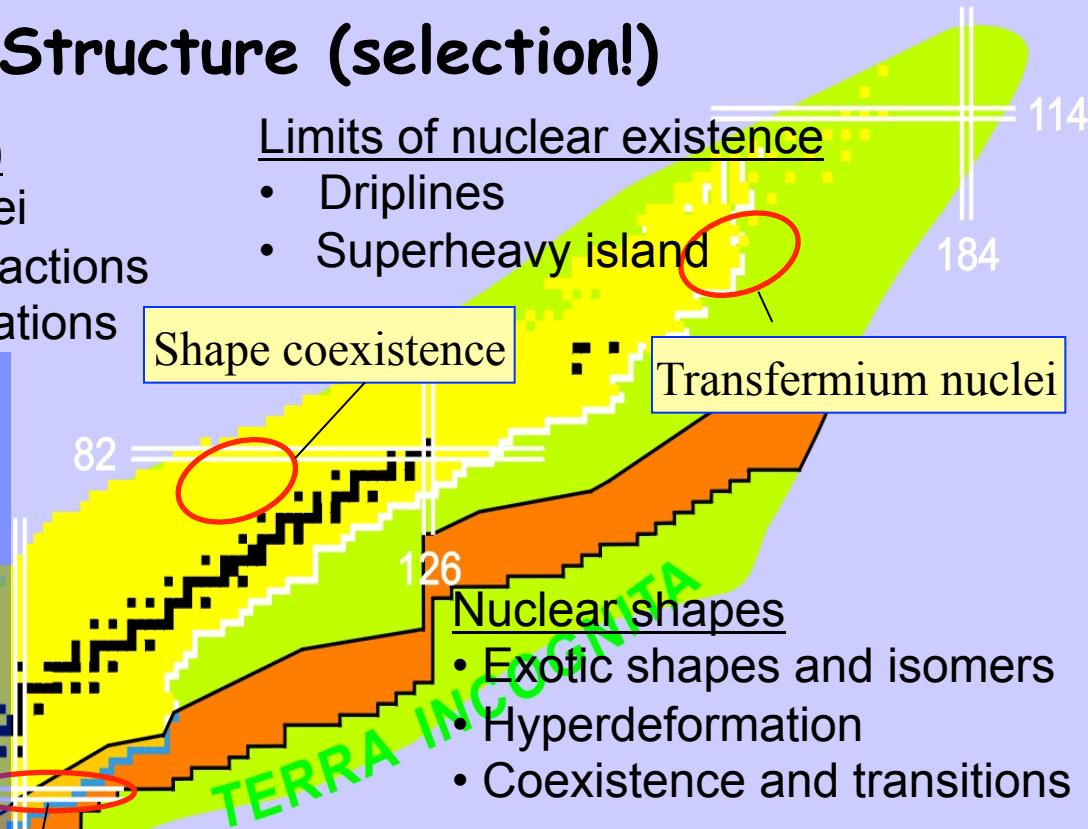
Proton drip line and N=Z nuclei

- Spectroscopy beyond the drip line
 - influence from continuum
- Neutron-proton pair modes
- Isospin symmetry
- Superallowed α -decays
- GT strengths
- exotic decays



Limits of nuclear existence

- Driplines
- Superheavy island



Neutron rich heavy nuclei ($N/Z \rightarrow 2$)

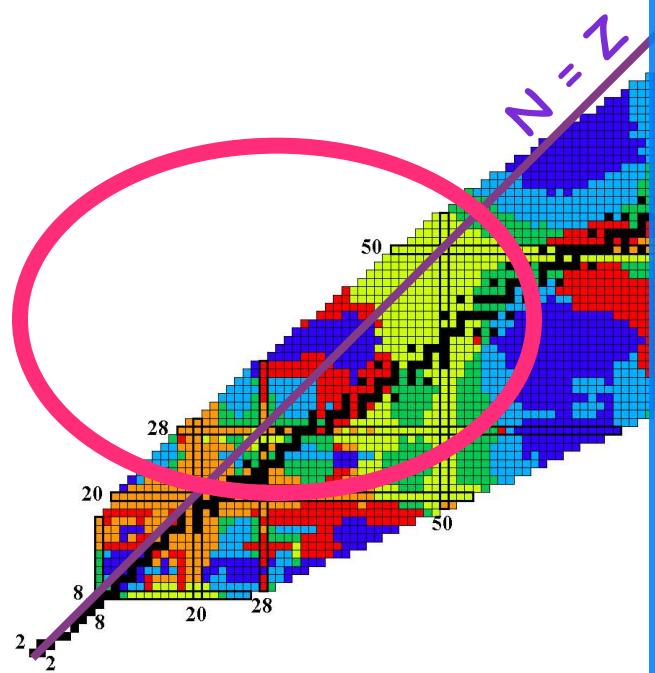
- Large neutron skins ($r_\nu - r_\pi \rightarrow 1\text{fm}$)
- New coherent excitation modes
- Shell quenching

Nuclei at the neutron drip line ($Z \rightarrow 25$)

- Very large proton-neutron asymmetries
- Resonant excitation modes
- Neutron Decay

Nuclear structure around ^{100}Sn - near the "top" of the N=Z line

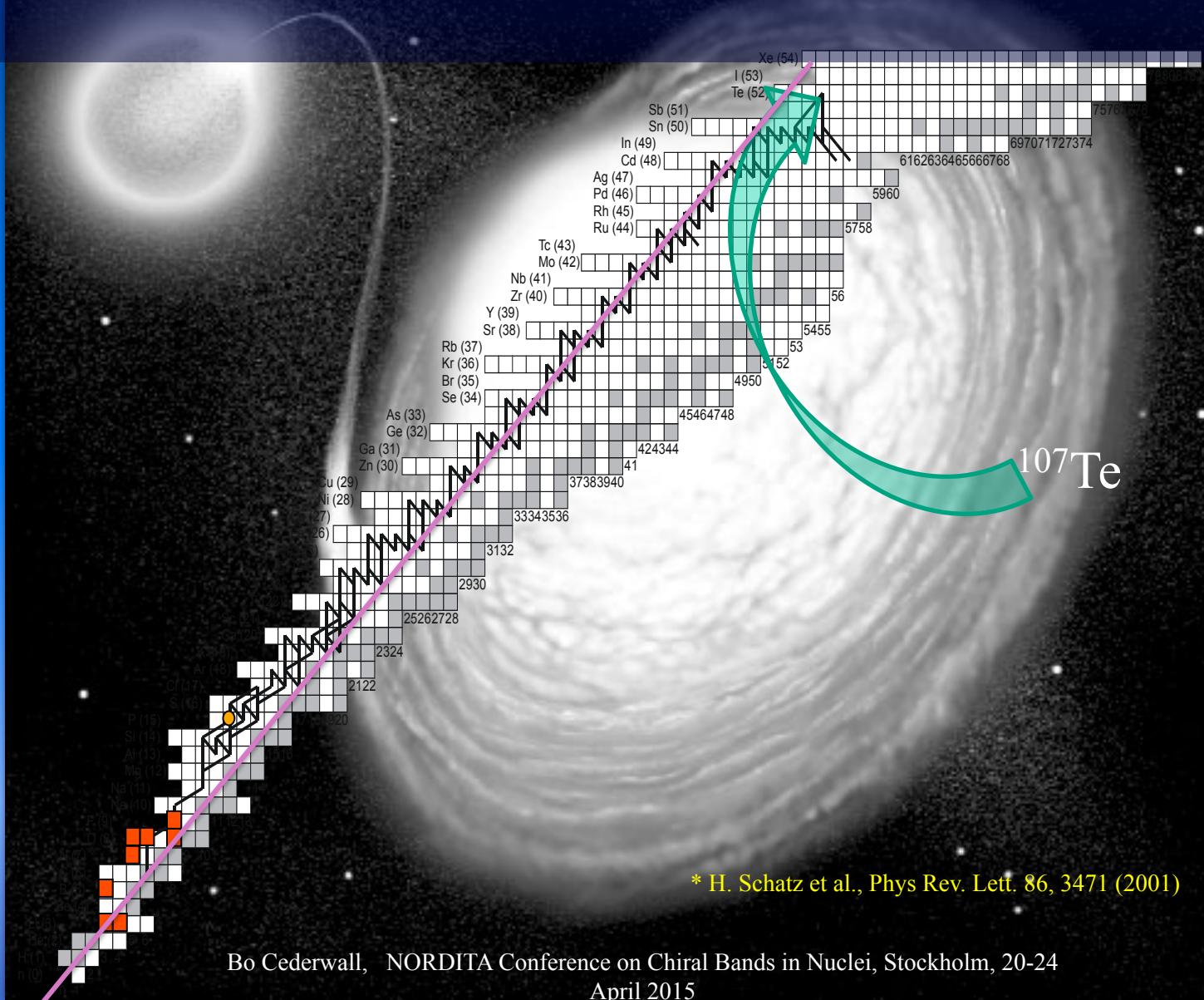
Moller Chart of Nuclides 2000
Quadrupole Deformation



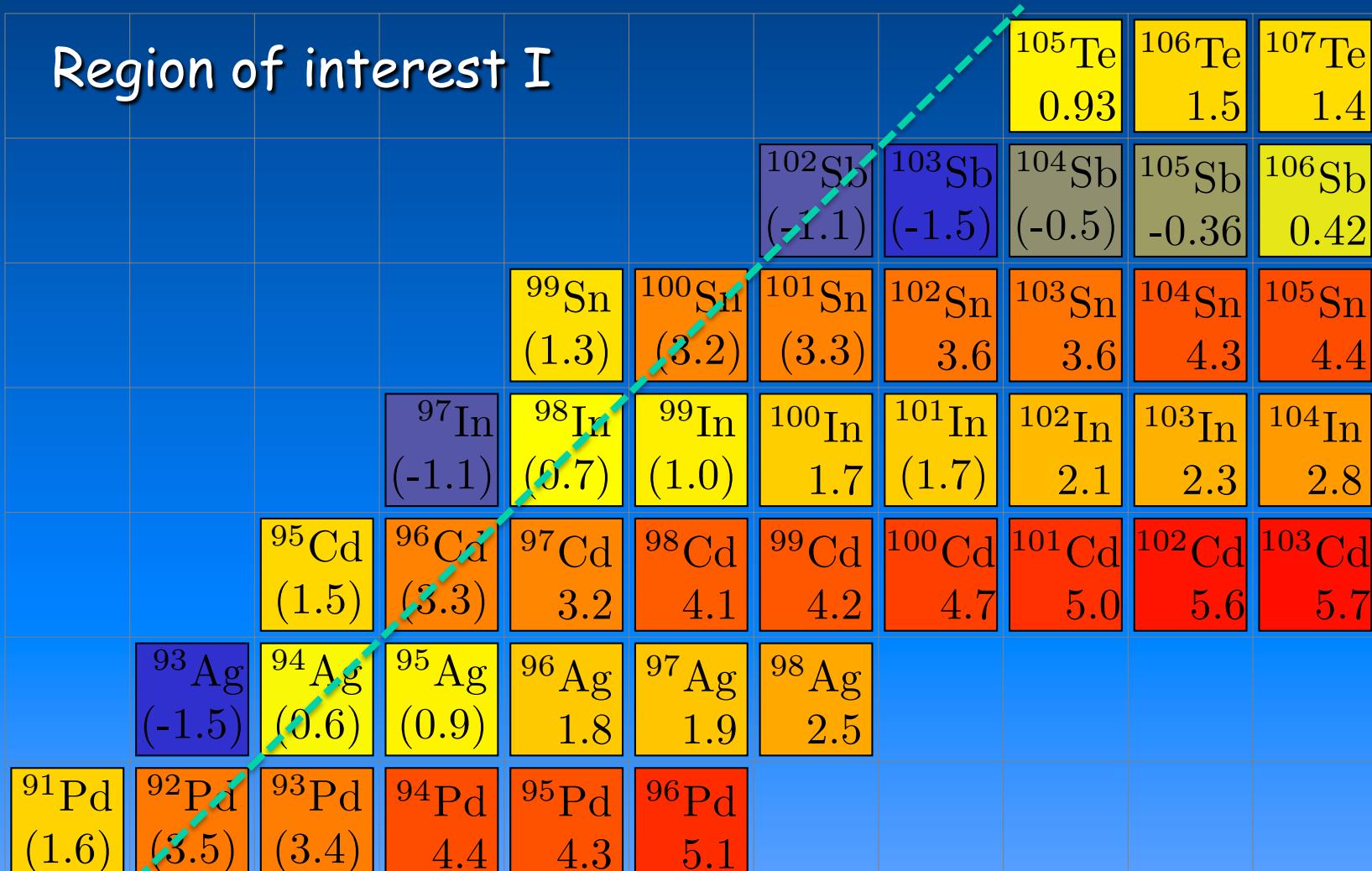
N=Z line coincides with a doubly-magic system and the proton dripline

- Neutron-proton correlations in identical orbitals
- Neutron-proton (isoscalar) pairing / Spin-aligned np coupling scheme
- Superallowed α -decays
- Emergence of "collectivity": $B(E2)$, $B(E3)$ strength
- LSSM calculations (eff. charges, interactions) can be applied and tested

Astrophysical interest:
End point of the rp-process path in X-ray bursts and
steady-state hydrogen burning on accreting neutron stars*



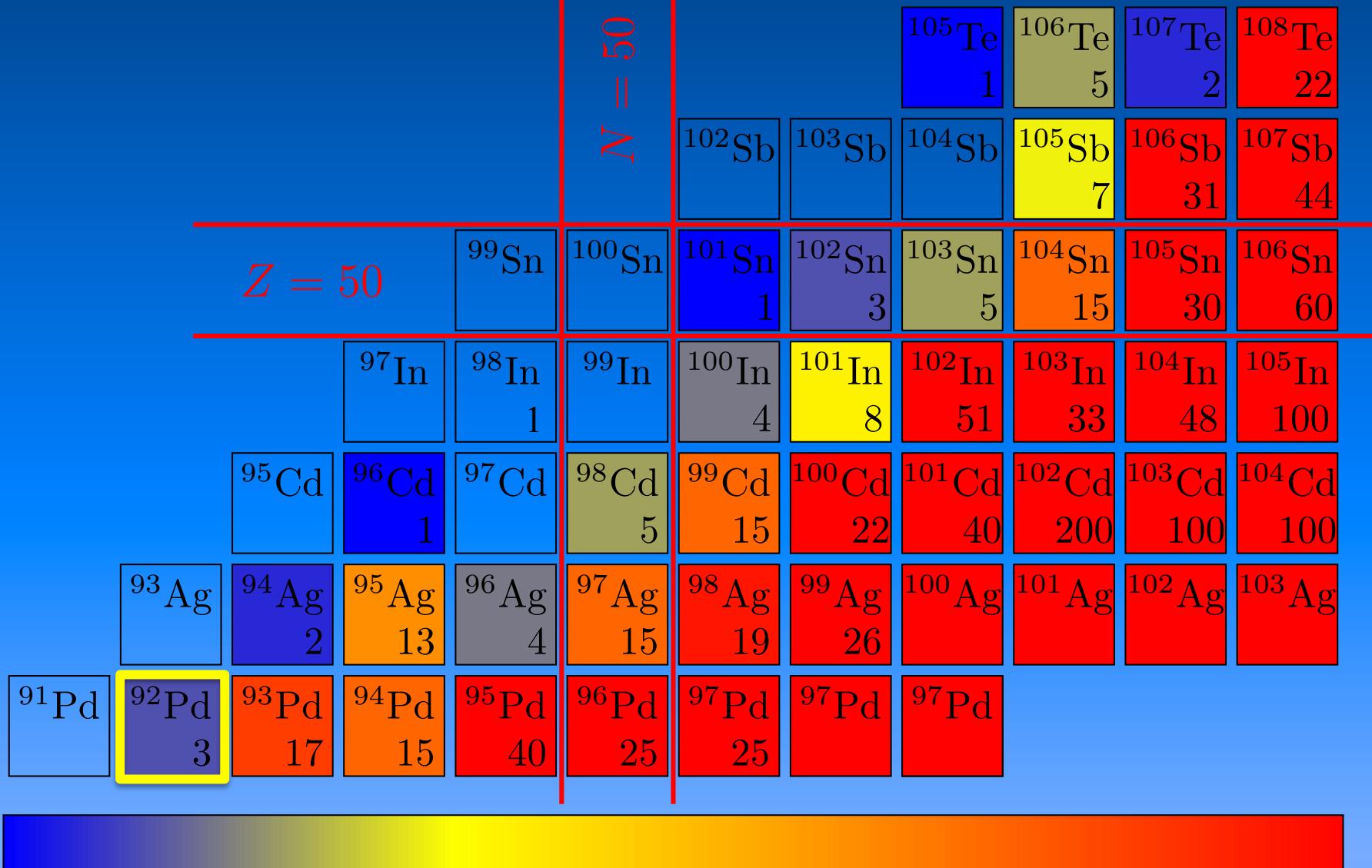
Region of interest I



Proton separation energy (MeV)

Courtesy M. Palacz

Structure data is lacking \rightarrow new generation detector arrays and facilities



Number of excited states known

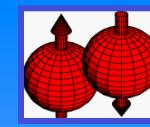
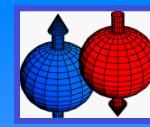
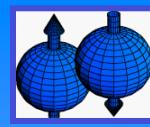
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adapted from M. Palacz

What is the nature of nuclear pair correlations near N=Z?

- A long-standing, open question in nuclear structure physics

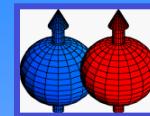
When approaching N=Z, "normal" pair correlations may remain or even be extended as neutrons and protons occupy identical quantum states:

$T=1, J=0$ ("isovector") nn, pp
as well as np Cooper pairs (neutrons and protons occupy identical orbits)



$T = 1 \quad J = 0$

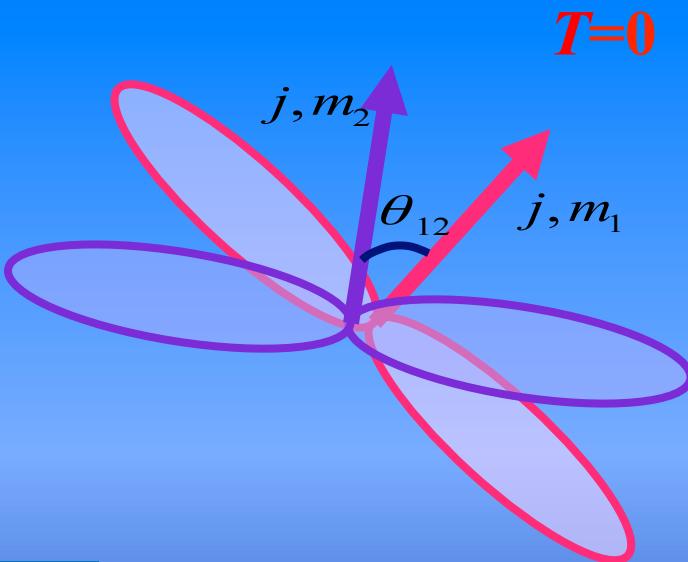
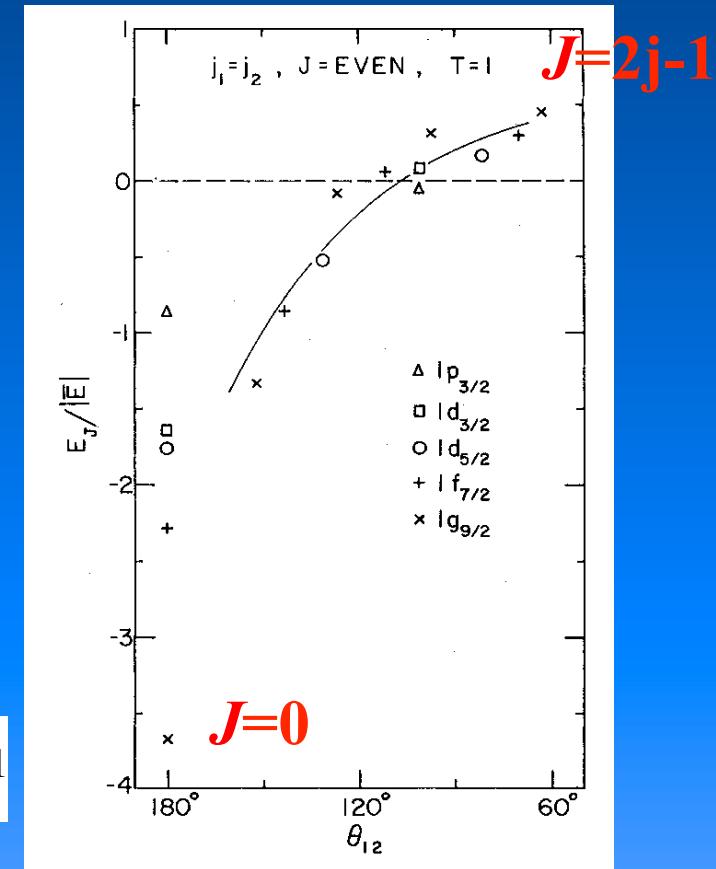
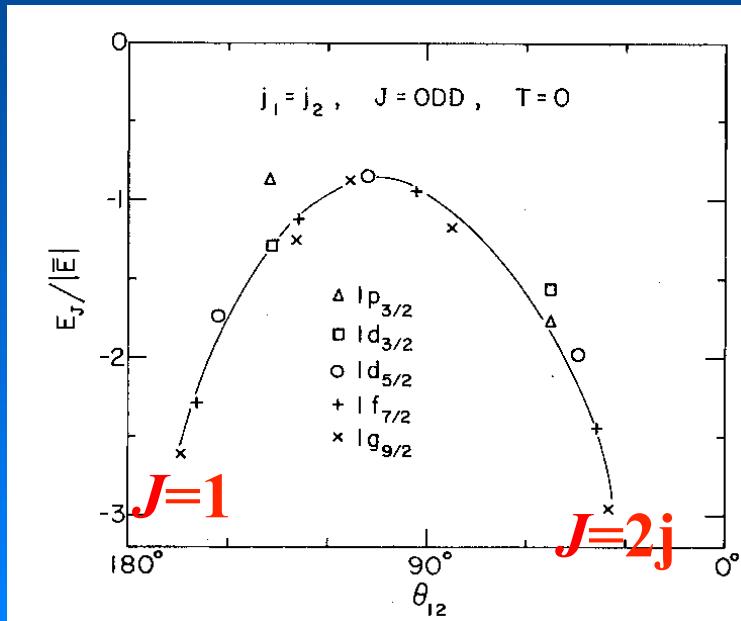
In addition: $T=0, J=1\dots$ ("isoscalar") np pairing?



$T = 0 \quad J > 0$

(We know the isoscalar effective NN interaction is strongly attractive but can it produce a correlated np pairing condensate?)

Effective (residual) interactions between nucleons in a j -shell



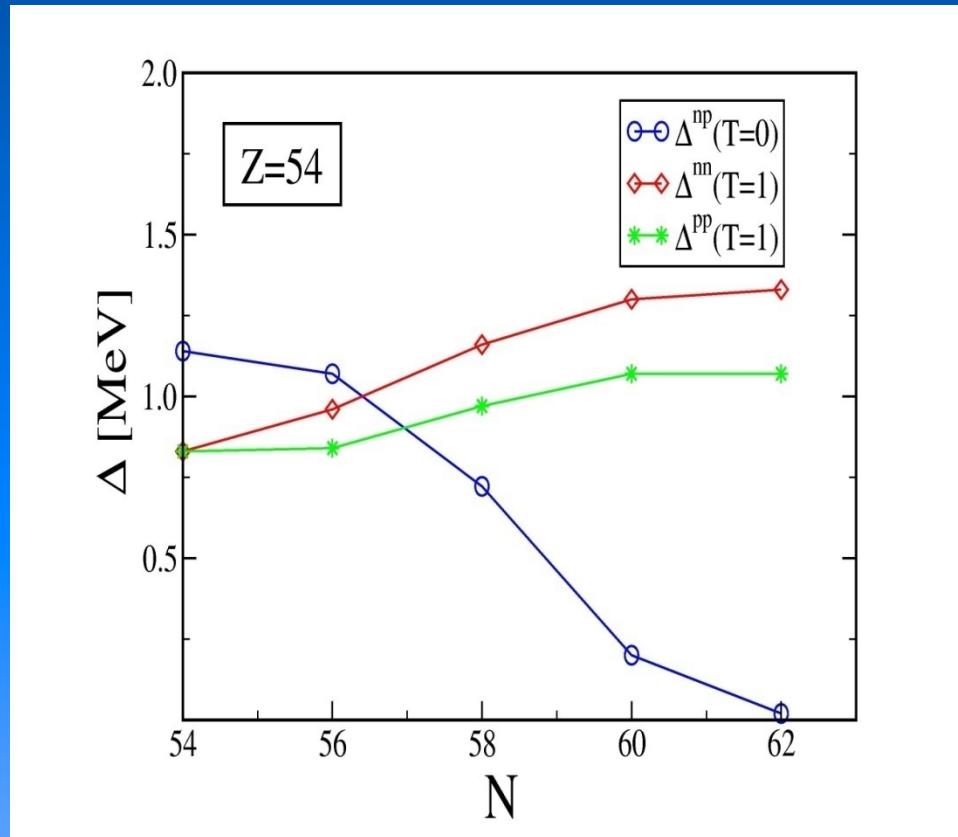
$$\cos \theta_{12} = \frac{J(J+1)}{2j(j+1)} - 1$$

J. P. Schiffer and W. W. True

Rev. Mod. Phys., Vol. 48, No. 2, Part I, April 1976

- effective force between particles, orbital overlap
- Pauli principle

The isoscalar (np) pair gap is predicted to increase sharply as $N \rightarrow Z$

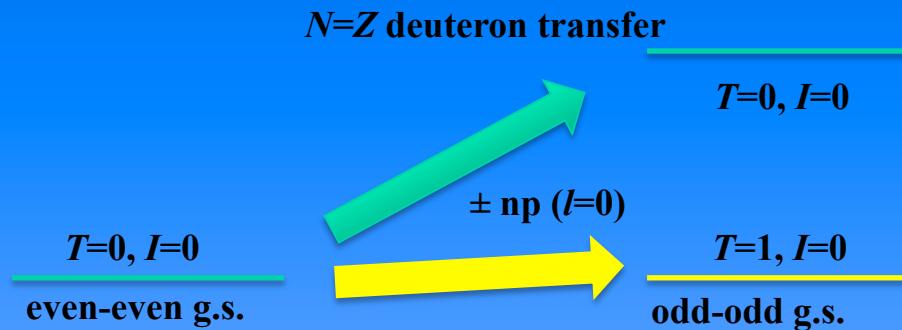


"Isospin-generalized" BCS-type calculation by W. Satula, R. Wyss
Phys. Rev. Lett. Vol. 86, 4488 (2001)

Does isoscalar (np) pairing *in the BCS sense* exist in Nature? (i.e. can we find an isoscalar pairing "deuteron" condensate somewhere?)

The experimental search for $T=0$ np pairing has focused on special features:

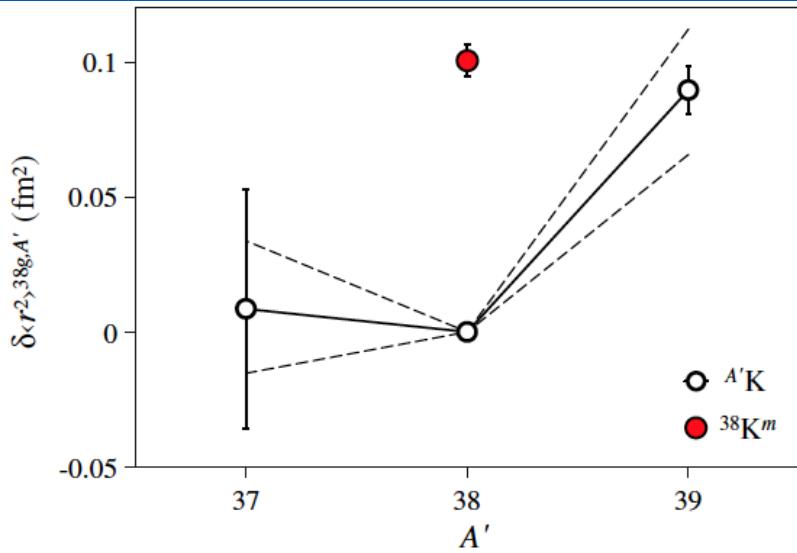
- g.s. binding energies of $N=Z$ (even-even vs odd-odd) nuclei
- high-spin properties of deformed $N=Z$ nuclei (reduced CAP, delayed alignments?)
- deuteron transfer reactions in inverse kinematics
(e.g. measure branching to $T=1$ and $T=0$ states in odd-odd $N=Z$ nuclei)



Need for (reaction) theory to develop sharp predictions

Status: No convincing evidence for so far
But much more to be done ...

Evidence for *isovector* np pairing is claimed from nuclear binding energies, rotational alignments, charge radii etc



“Proton-Neutron Pairing Correlations in the Self-Conjugate Nucleus ^{38}K Probed via a Direct Measurement of the Isomer Shift”, Bissell et al., PRL 113, 052502 (2014)

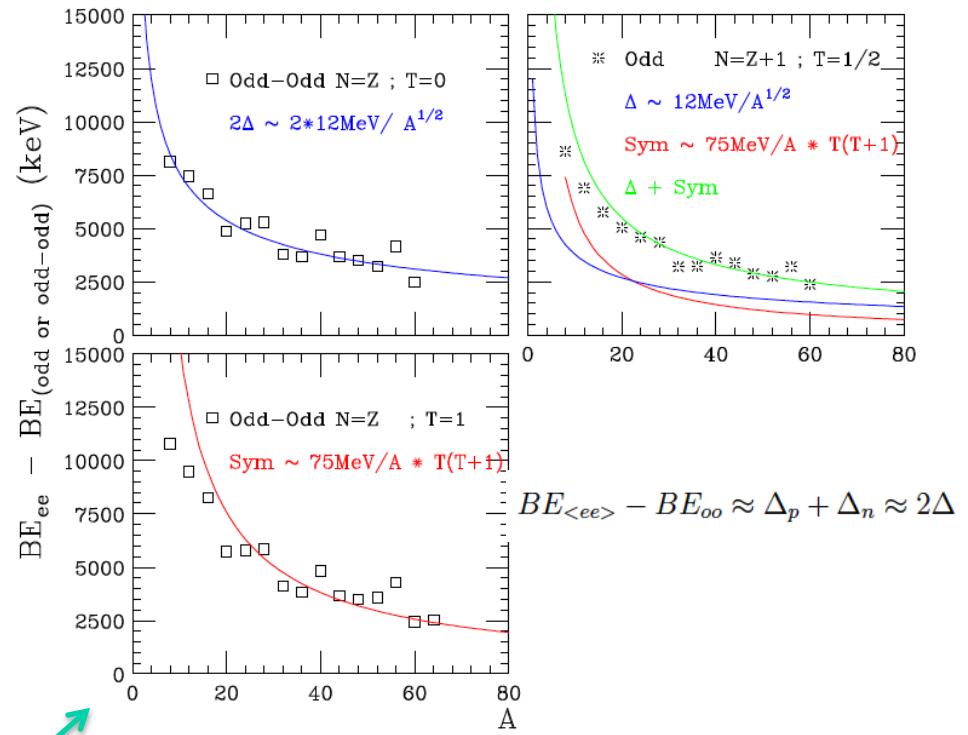
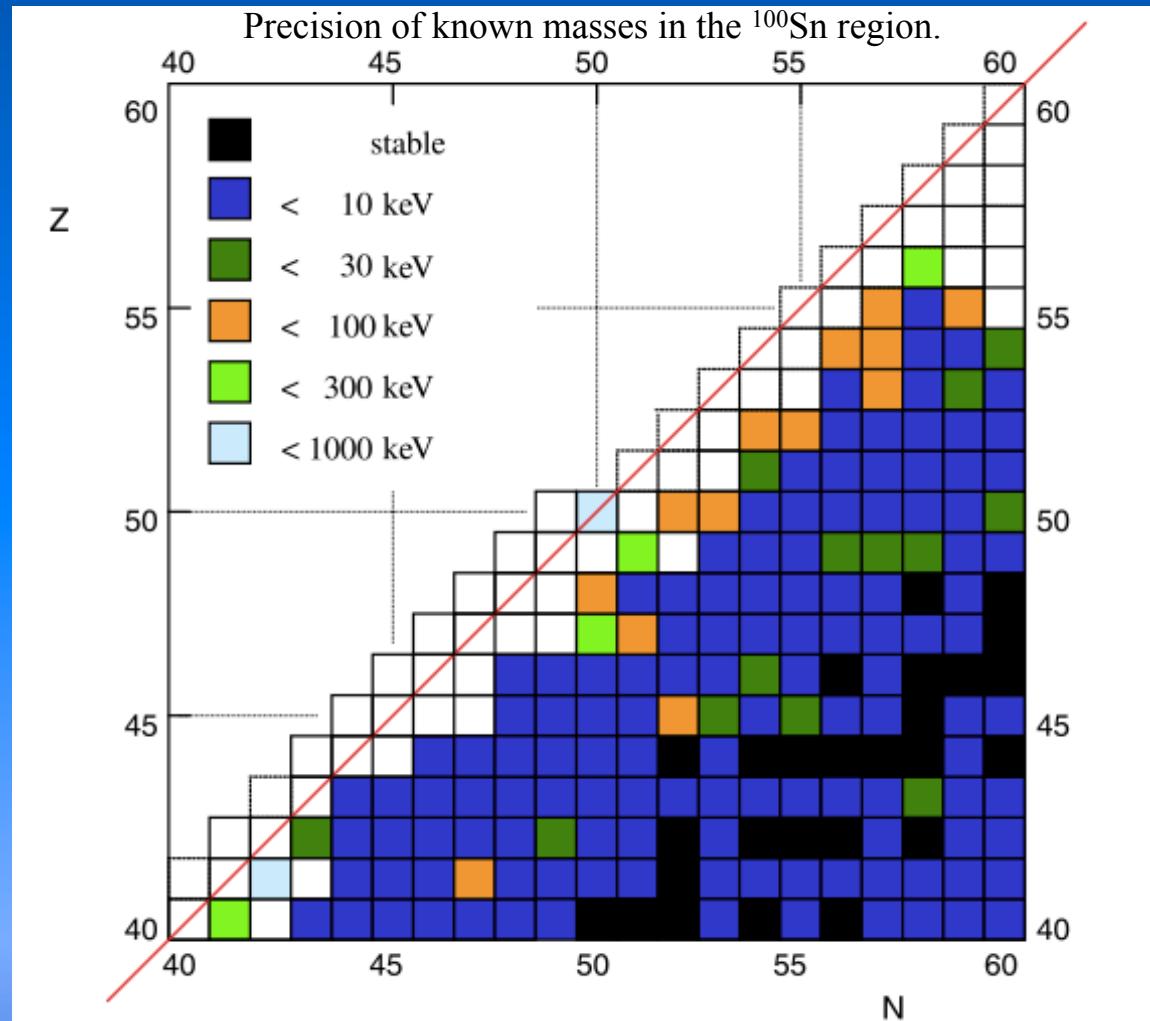


Figure 25: Summary of binding energy differences for nuclei along the $N = Z$ line.

“Overview of neutron-proton pairing”
Frauendorf, S., Macchiavelli, A.O.
Progress in Particle and Nuclear Physics
volume 78, 2014, pp. 24 - 90

“Excludes” $T=0$ pairing in g.s
Note: Data end at $N=Z \approx 30$!
Symmetry energy?

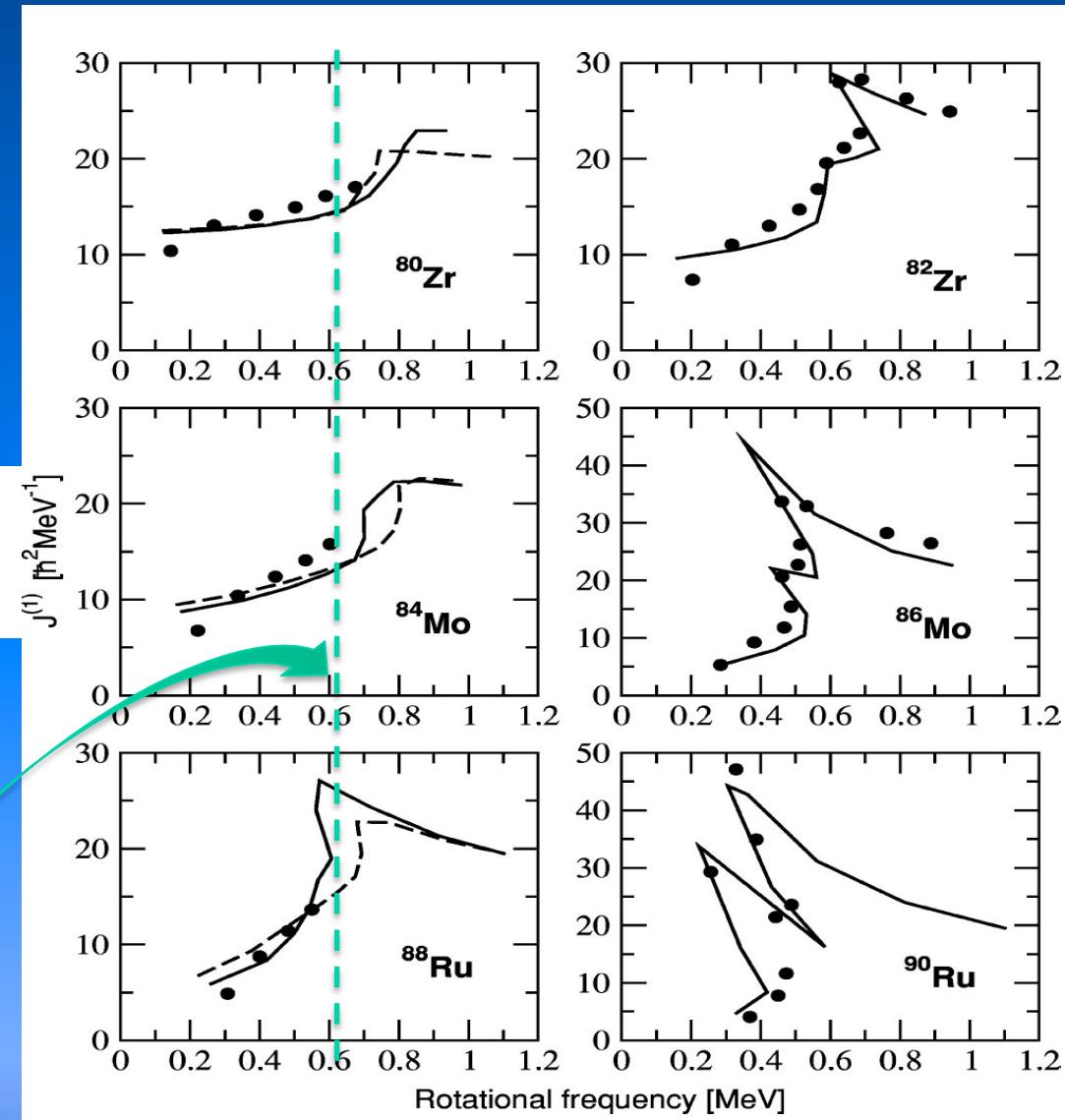
Precision mass data is crucially lacking for N=Z!



T. Faestermann, M. Górska, and H. Grawe, Prog. Part. Nucl. Phys. 69, 85 (2013)

Delayed (or absent) paired ($T=1$) bandcrossings in deformed $N=Z$ nuclei?

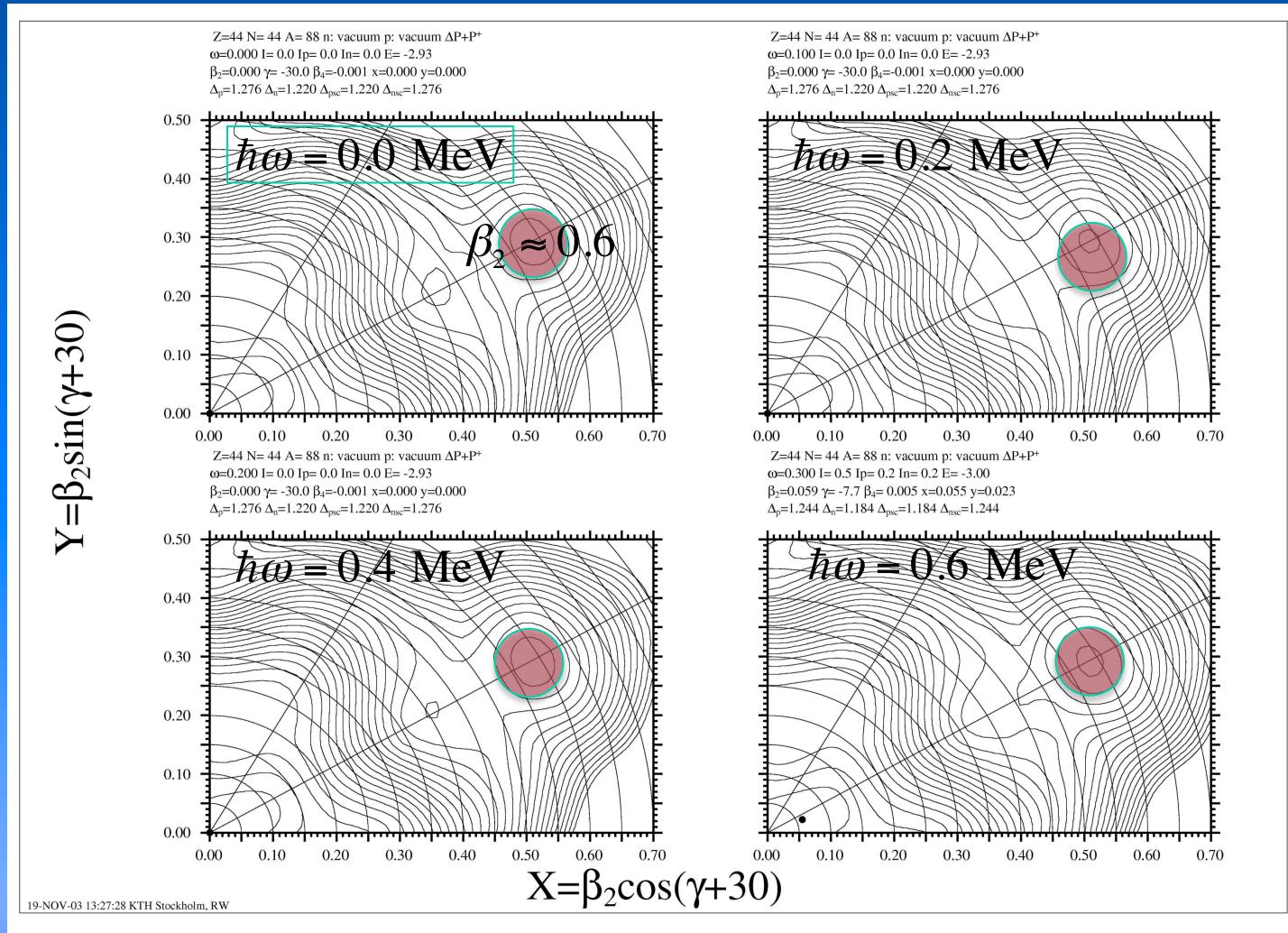
PSM _____
PSM with enhanced
np res. int. _____



Data end crucially here!

N. Marginenian *et al.*, Phys. Rev. C 65, 051303R (2002)

^{88}Ru - a superdeformed self-conjugate system?
 The perfect "laboratory" for investigating nuclear pairing effects?
 Deep SD minimum persists down to zero rot frequency, low excitation energy



Strong T=0 neutron-proton (np) pair correlations may lead to something different from a BCS-type of pairing condensate:

"Isoscalar spin-aligned coupling scheme"*

predicted for N=Z nuclei close below ^{100}Sn .

- Unique signature of "vibrational-like" yrast energies and "rotational-like" $B(E2)$ strengths. $B(E2;0^+\rightarrow 2^+)$ s develop differently compared with standard seniority scheme along isotopic chain as $N\rightarrow Z$
- A new manifestation of strong np-pair correlations

* 'Evidence for a spin-aligned neutron-proton paired phase from the level structure of ^{92}Pd '
B. Cederwall et al., Nature **469**, 68 (2011)

'Spin-aligned neutron-proton pair mode in atomic nuclei'
C. Qi, **J. Blomqvist**, T. Bäck, B. Cederwall, A. Johnson, R. J. Liotta, and R. Wyss
Phys. Rev. C 84, 021301 (2011)

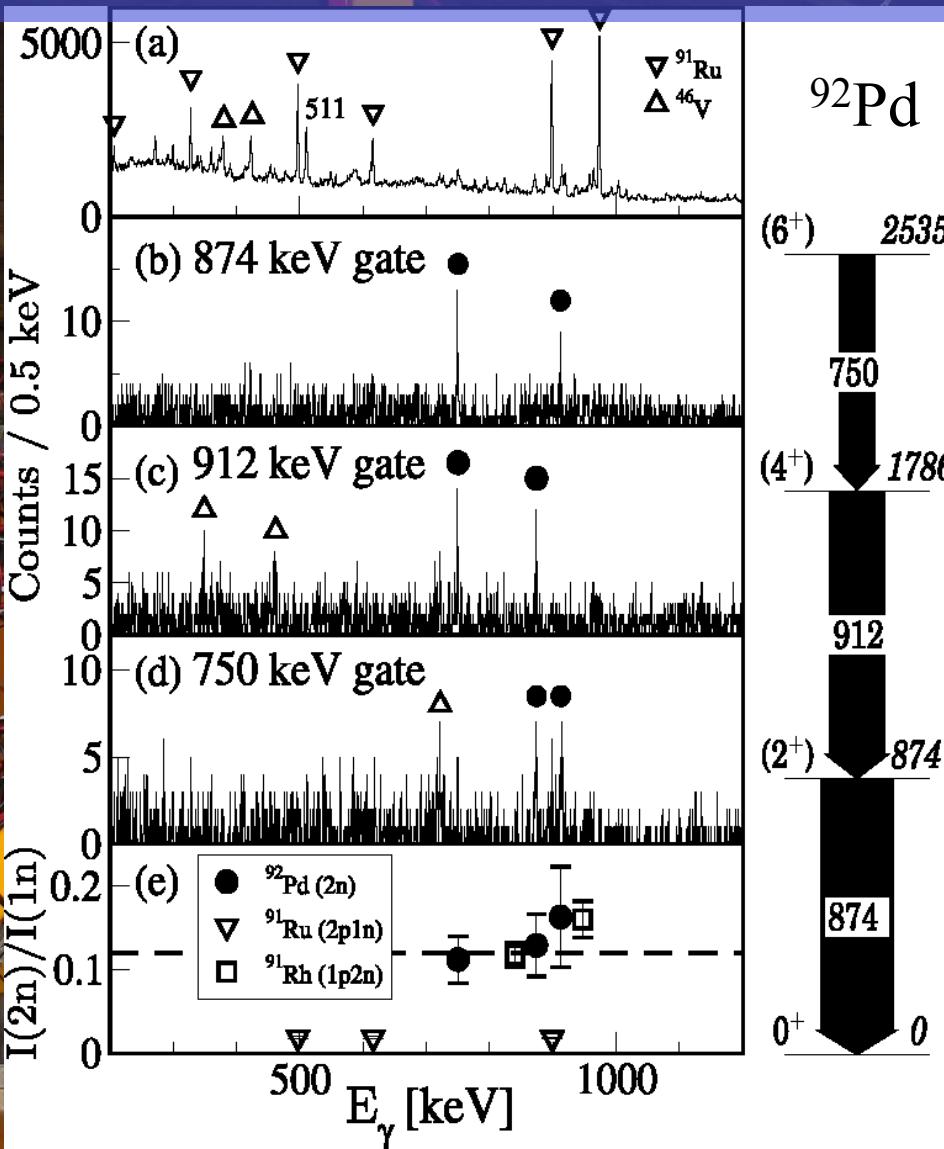


Observation of excited states in the N=Z=46 nucleus ^{92}Pd

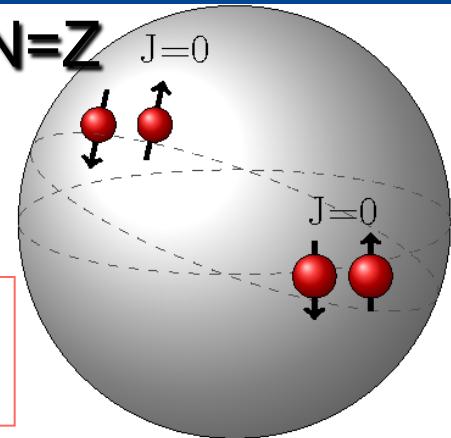
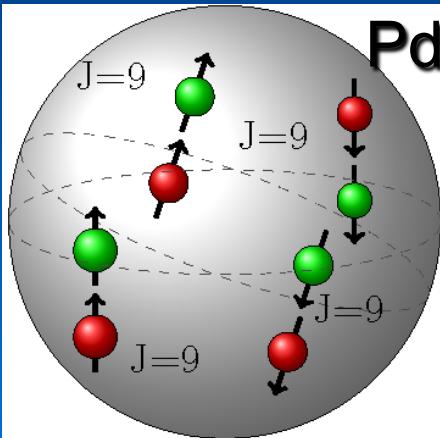
EXOGAM + Neutron Wall + Diamant experiment

laboratoire commun CEA/DSM

CNRS/IN2P3



Pd energy level systematics near N=Z - effects of np interactions



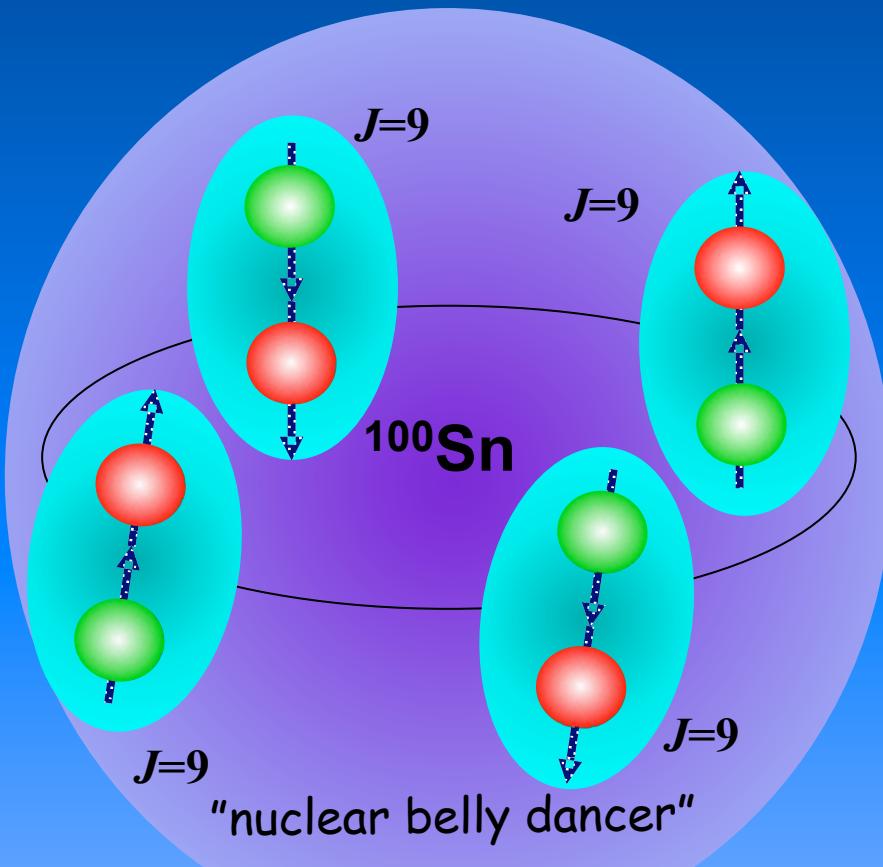
*Calculations performed in several model spaces,
i.e., $0g9/2$, $0g9/2-1p1/2$ and $0g9/2-1p1/2-0f5/2-1p3/2$
which all give very similar results.*

Int. parameters determined to reproduce exp energies in $^{94,95}Pd$, $^{93,94}Rh$

10^+ 4072	10^+ 4065	10^+ 4052	10^+ 4065	10^+ 3862	10^+ 3796	10^+ 4131	10^+ 3784
8^+ 3127	10^+ 3257			10^+ 3257			8^+ 2636
(6 ⁺) 2536	6 ⁺ 2466	8 ⁺ 2600	8 ⁺ 2749	8 ⁺ 2633	8 ⁺ 2635	8 ⁺ 2588	8 ⁺ 2792
		6 ⁺ 2110	6 ⁺ 2079	6 ⁺ 2212	6 ⁺ 2223	6 ⁺ 2128	6 ⁺ 2374
(4 ⁺) 1786	4 ⁺ 1708	4 ⁺ 1518			2 ⁺ 1405	2 ⁺ 1199	6 ⁺ 2330
					2 ⁺ 1405	2 ⁺ 1199	6 ⁺ 2380
(2 ⁺) 874	2 ⁺ 878	2 ⁺ 797	2 ⁺ 1171	2 ⁺ 1417			4 ⁺ 1709
	15						4 ⁺ 1682
0 ⁺ 0	0 ⁺ 0	0 ⁺ 0	0 ⁺ 0	0 ⁺ 0	0 ⁺ 0	0 ⁺ 0	4 ⁺ 1720
^{92}Pd exp	^{92}Pd SM	^{92}Pd T=0	^{92}Pd T=1	^{92}Pd no np	^{94}Pd no np	^{94}Pd T=1	^{94}Pd SM
							2 ⁺ 864
							2 ⁺ 861
							2 ⁺ 814
							7.5
							11
							0 ⁺ 0
							^{96}Pd exp

B. Cederwall et al., Nature **469**, 68 (2011)

Strong residual np interactions → Spin-aligned T=0 np coupling scheme for N=Z nuclei below ^{100}Sn



"Spin-aligned $T=0$ np paired phase"
Not pairing as a BCS condensate

The diagonal SM interaction matrix element that corresponds to the isoscalar $v\pi(g_{9/2})^2$ aligned np pair ($J^\pi = 9^+$),
 $V_9 = \langle g_{9/2}^2; J=9 | V | g_{9/2}^2; J=9 \rangle$,
is strongly attractive, with $V_9 \sim -2$ MeV

Aligned isoscalar np coupling:

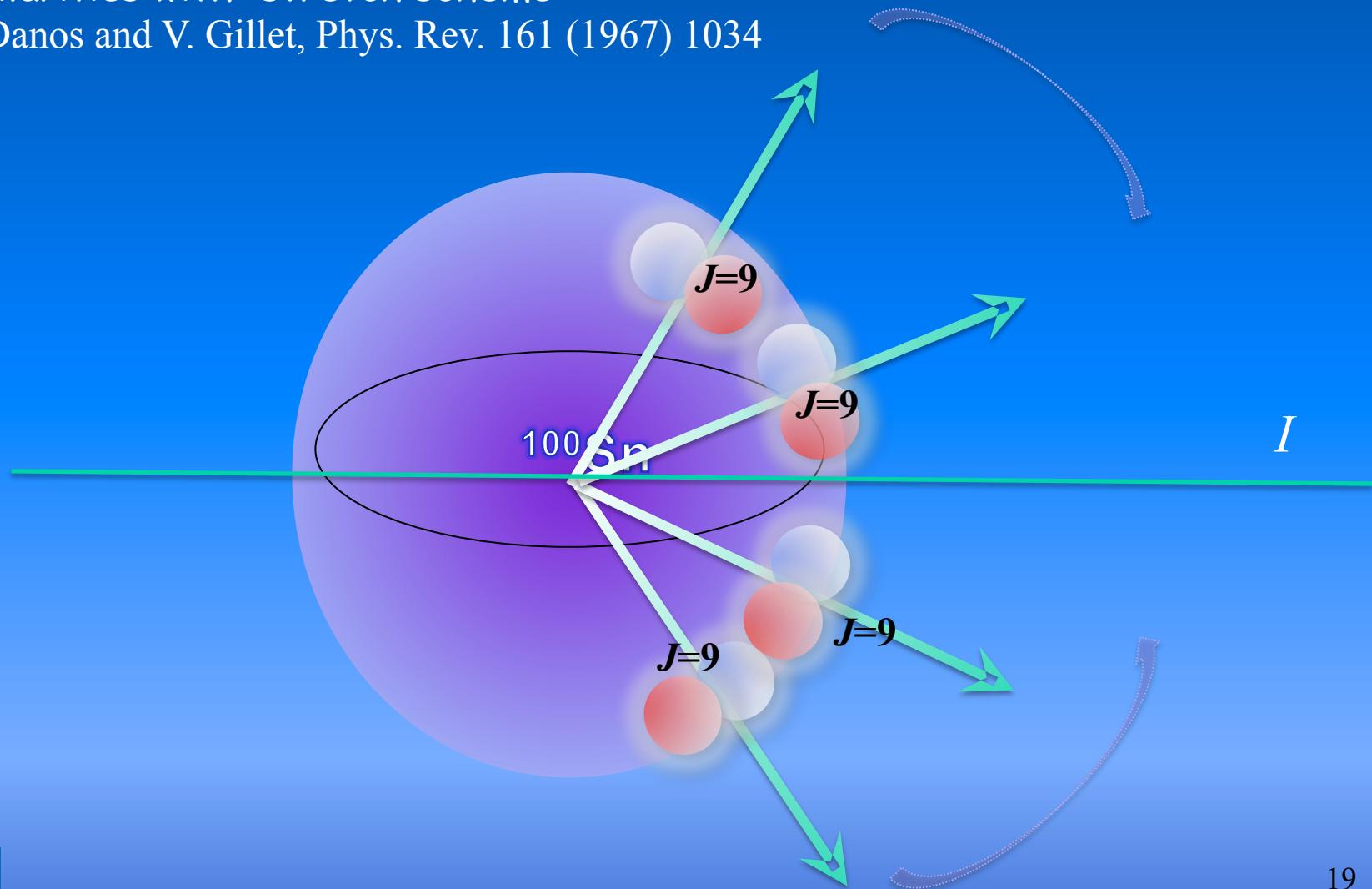
$$\Psi_{\text{G.S.}} = [(\{vg_{9/2}^{-1} \times \pi g_{9/2}^{-1}\}_{9+})^2]_{0+} \times [(\{vg_{9/2}^{-1} \times \pi g_{9/2}^{-1}\}_{7+})^2]_{0+}$$

Different from the standard textbook description of the ground states in even-even nuclei!

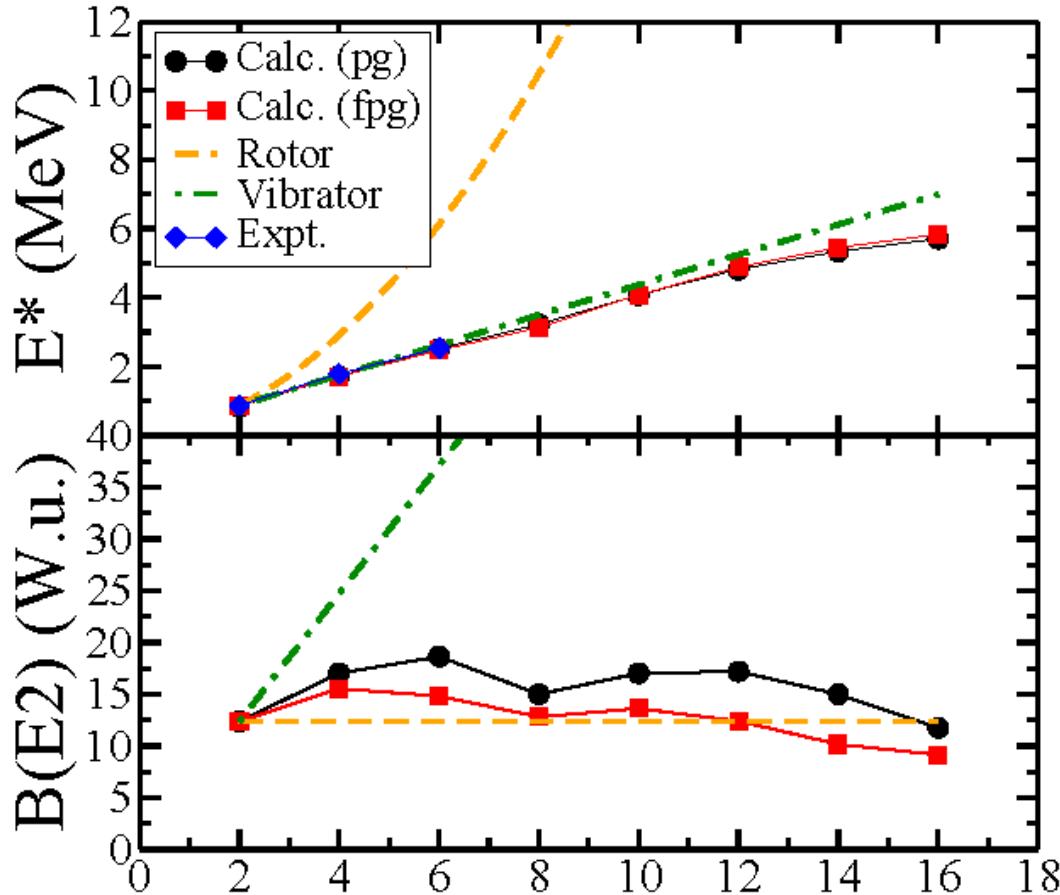
Generation of angular momentum in the isoscalar spin-aligned coupling scheme (^{92}Pd)

Similarities with "stretch scheme"

M. Danos and V. Gillet, Phys. Rev. 161 (1967) 1034



"Unique" signature of spin-aligned T=0 coupling scheme

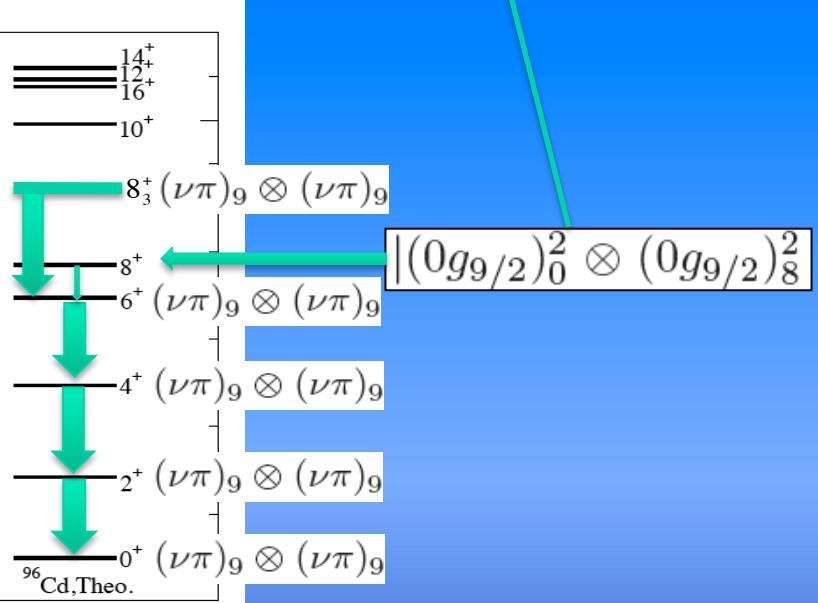
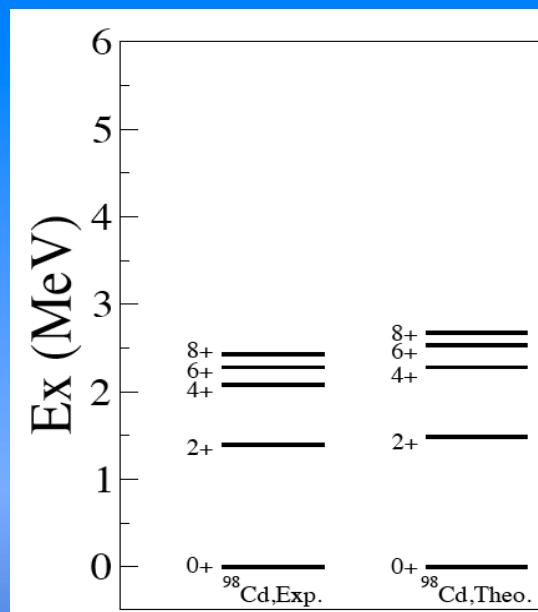
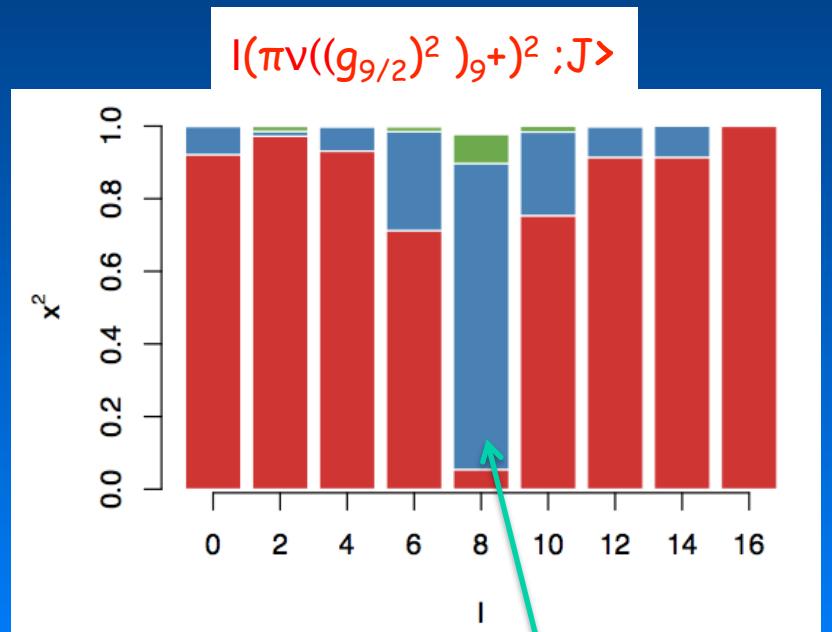
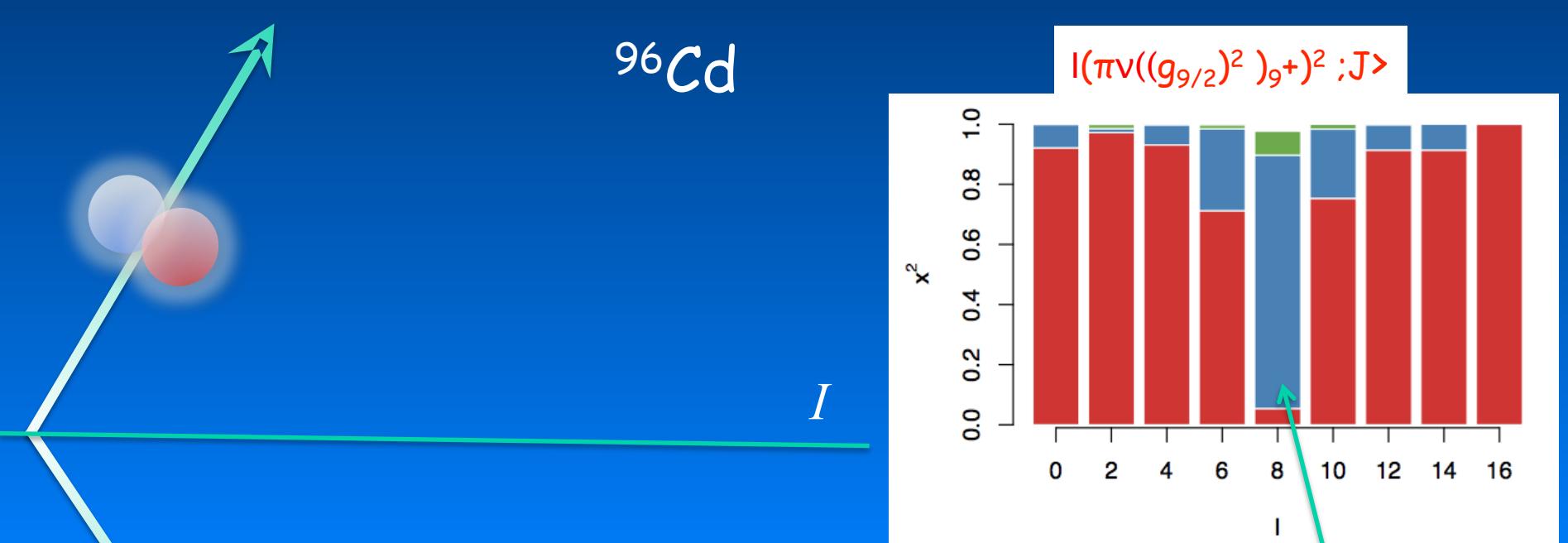


Upper: Shell model spectra of ^{92}Pd calculated within the $1\text{p}_{3/2}0\text{f}_{5/2}1\text{p}_{1/2}0\text{g}_{9/2}$ space [10] (fpg) and the $1\text{p}_{1/2}0\text{g}_{9/2}$ space (pg). Lower: $B(\text{E}2; I \rightarrow I - 2)$ values in ^{92}Pd calculated within the fpg and pg spaces. The two dashed lines show the predictions of the geometric collective model normalized to the 2^+_1 state C.Qi, J.Bломqvist, T.Bäck, B. Cederwall, A. Johnson, R. J. Liotta, and R. Wyss, PRC 84, 021301(R) (2011)

LSSM calculation, fpg (f5/2, p3/2,p1/2, g9/2)
(all shells between 28 and 50) C. Qi

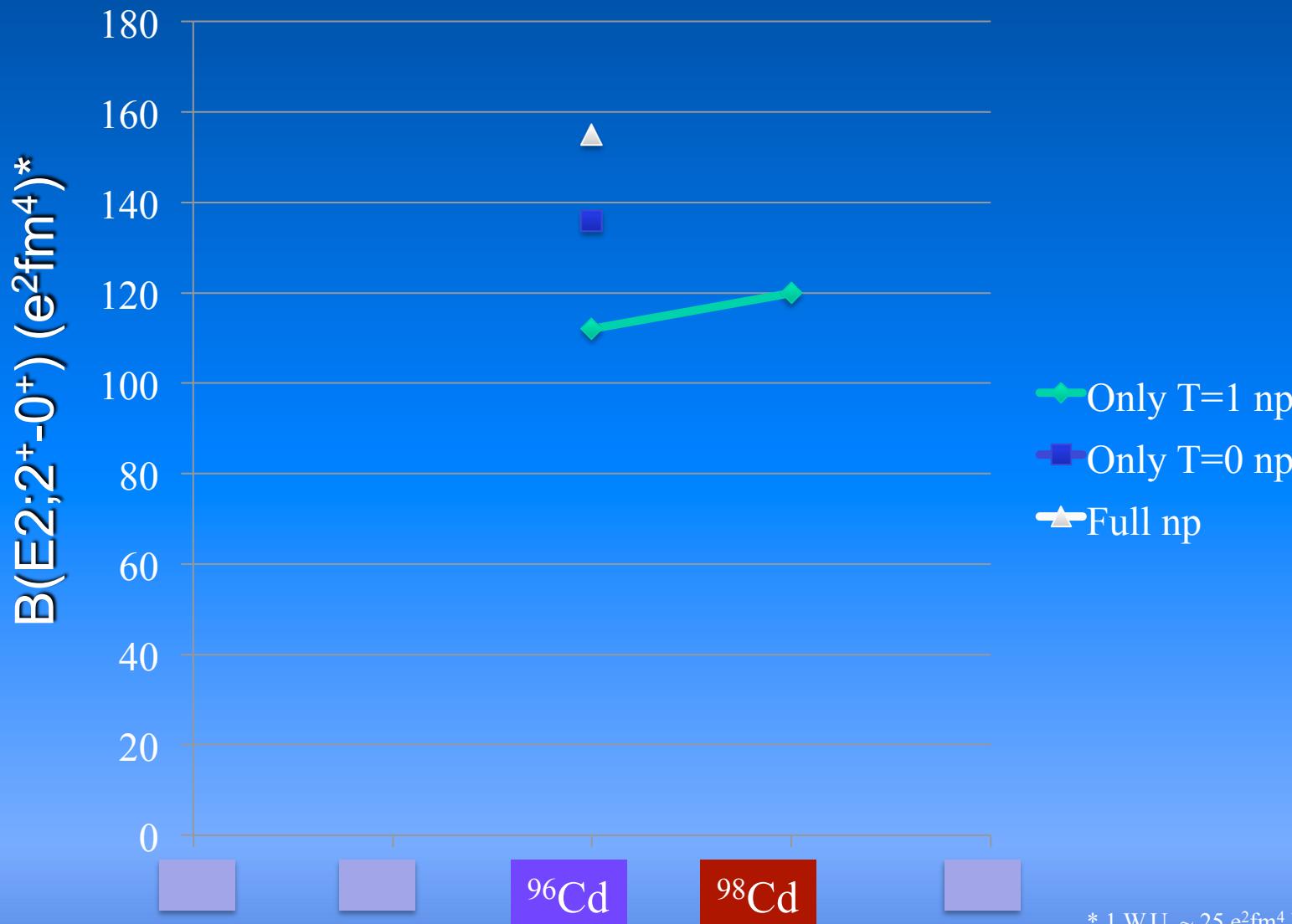


* 1 W.U. $\sim 25 e^2 \text{fm}^4$ here



C. Qi, priv. comm., Z.X. Xu et al., Nuclear Physics A 877 (2012) 51–58

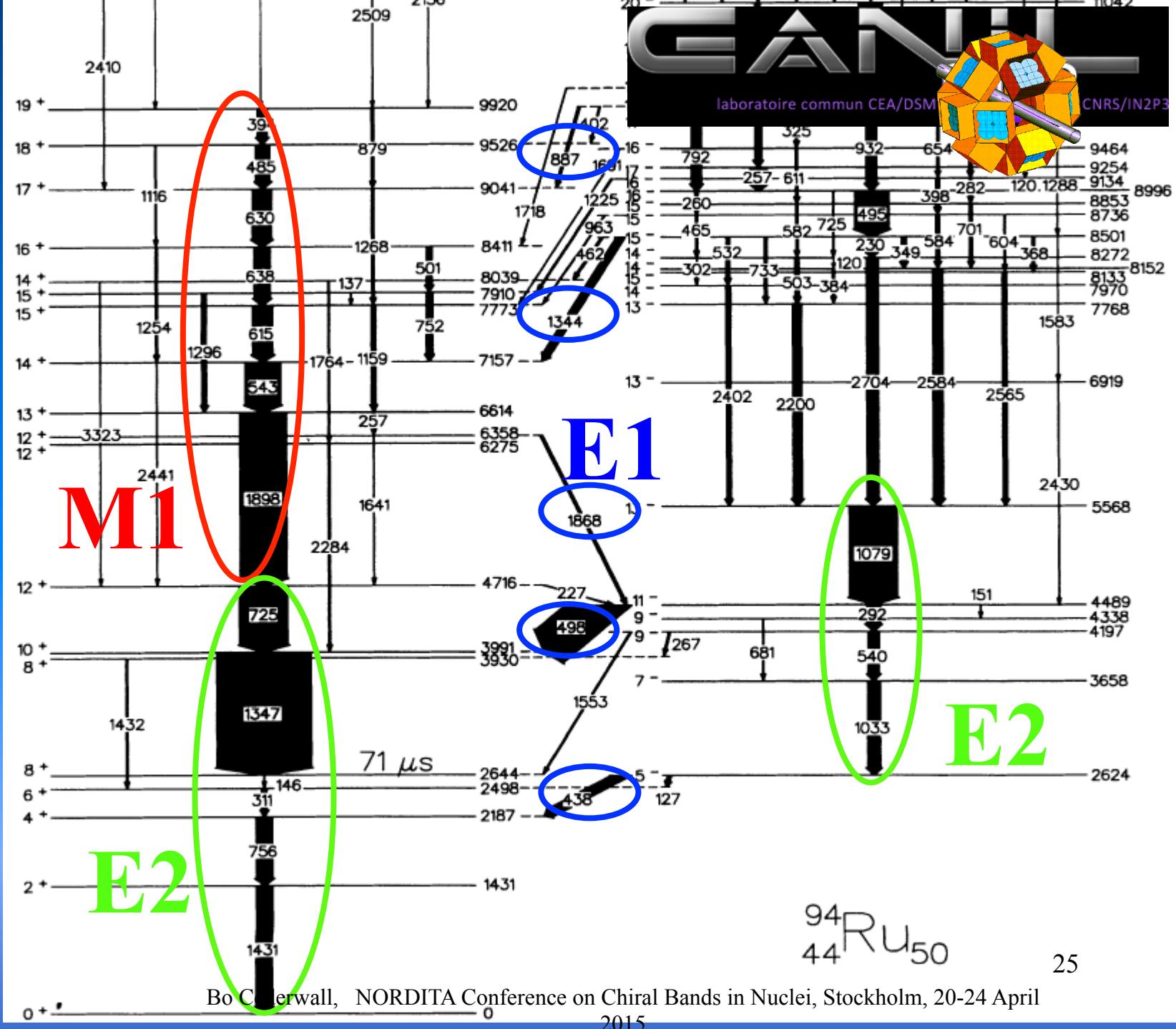
LSSM calculation, fpg (f5/2, p3/2,p1/2, g9/2)
(all shells between 28 and 50) C. Qi



* 1 W.U. $\sim 25 e^2 \text{fm}^4$ here

Critical test of LSSM interactions: Precision spectroscopy of E1 transitions in semimagic nuclei

- All low-lying states in this region are well described within the $f_{5/2}$, $p_{3/2}$, $p_{1/2}$ and $g_{9/2}$ model space.
- $E1$ transitions are “forbidden” within this space since the matrix elements $\langle f | E1 | i \rangle$ vanish for all possible combinations of initial states i and final states f .
- Presence of $E1$ transitions → other (higher or deeper lying) single-particle states are active.
- A sensitive probe of SM parameters: Even a minute admixture of such configurations in the wave function may greatly increase the probability of $E1$ decay since the $E1$ single-particle matrix element is very large in comparison with any other multipole mode.
- The observed $E1$ transition strengths serve as a critical test of the shell-model wave function with respect to the model subspace from which it is constructed. In particular for core excited states.



E1 hindrance factors in N=50 nucleus ^{94}Ru

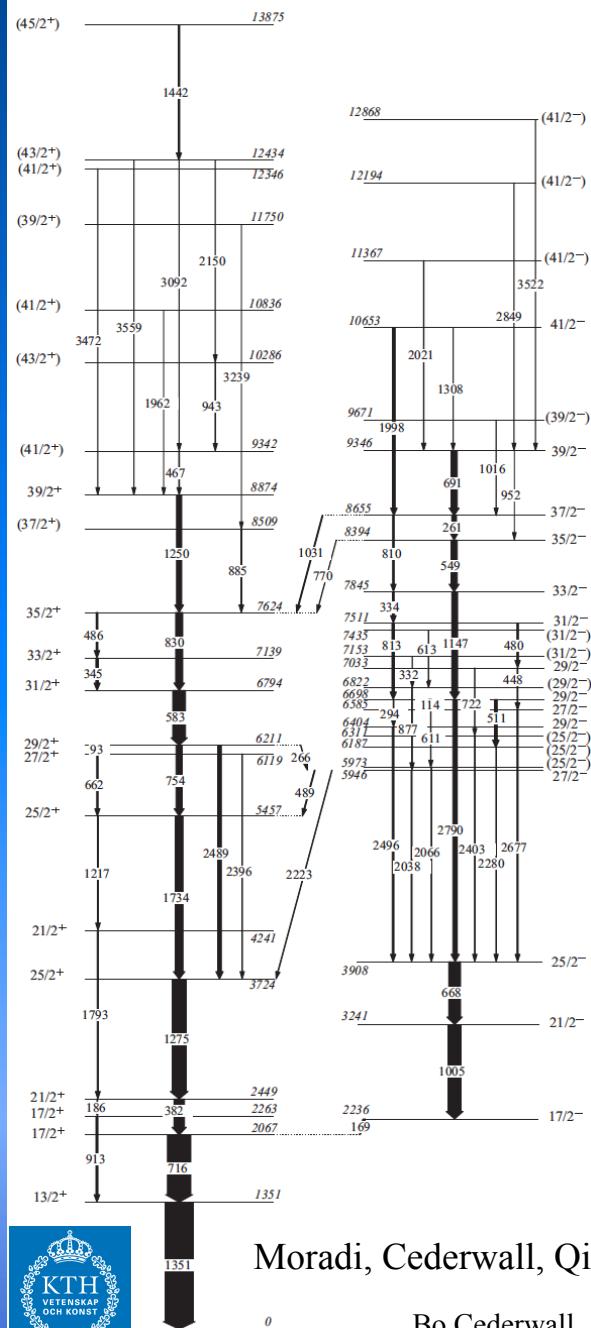
TABLE II. The hindrance factors^a H for the observed $E1$ γ -ray transitions as deduced from the branching ratios and lifetimes [30] of the initial states. Uncertainties are given in parentheses. $\Delta E_\gamma = \pm 0.5$ keV.

E_γ (keV)	$J_i^\pi \rightarrow J_f^\pi$	$H \times 10^5$ (W.u.) $^{-1}$
257	$13^+ \rightarrow 12_1^-$	0.006 (1)
462	$15_2^- \rightarrow 14_2^+$	0.051 (5)
402	$18_1^- \rightarrow 18^+$	0.188 (25)
1344	$15_2^- \rightarrow 14_1^+$	0.451 (32)
227	$12_1^+ \rightarrow 11^-$	0.57 (27)
887	$18_1^- \rightarrow 17^+$	1.09 (12)
438	$5^- \rightarrow 4^+$	1.90 (17)
498	$11^- \rightarrow 10^+$	4.27 (19)

$$^a H = \frac{A^{2/3}}{15.5 \times B(E1)}.$$

Moradi, Qi, Cederwall et al., PRC, 014301 (2014)

E1 hindrance in $^{95}\text{Rh}_{50}$



E_γ (keV)	$J_i^\pi \rightarrow J_f^\pi$	$H \times 10^5$ (W.u.) $^{-1}$
266	$29/2^+ \rightarrow 27/2^-$	0.11(1)
1031	$37/2^- \rightarrow 35/2^+$	0.29(1)
770	$35/2^- \rightarrow 35/2^+$	0.35(1)
169	$17/2^- \rightarrow 17/2_1^+$	2.8(2)

$$^a H = \frac{A^{2/3}}{15.5 \times B(E1)}$$

Table II of this work, there are a number of weak E1 transitions present, even at low excitation energies (below 5 MeV). We evaluated all possible E1 transitions among states within the expanded shell model space. Considering first the E1 decay from the first $17/2^-$ state to the yrast $17/2_1^+$ state, these states are predominantly of $\pi(1p_{1/2}^{-1}0g_{9/2}^{-4})$ and $\pi(0g_{9/2}^{-5})$ character, respectively, as discussed above; i.e., without possibility of E1 decay. This transition has the largest hindrance, 2.8×10^5 W.u. $^{-1}$ among the E1 decays observed in ^{95}Rh which is reflected by the long (~ 19 ns) half-life of the $17/2^-$ state [2]. The core-excitation components in these two states are mainly of a one-neutron character. In our calculation the contribution to the transition in terms of occupation probability from the high lying shells $1d_{5/2}$ and $0g_{7/2}$ is 0.02, while the corresponding contribution from the deep-lying shells $1p_{3/2}$ and $0f_{5/2}$ is approximately 10^{-4} . Therefore the E1 hindrance factor is of the order 5×10^5 W.u. $^{-1}$, which is consistent with the value given in Table II. The absence of E1 transitions depopulating the following negative-parity states up to $25/2^-$ indicates that the influence of the core excited configurations is limited in these states as predicted in Ref. [8]. E1 decays observed from the higher-lying negative-parity states as well as from the $29/2_1^+$ state signal significant contributions from core-excited configurations with one neutron being excited from below the $N = 50$ shell closure to the $1d_{5/2}$ or $0g_{7/2}$ orbits, in agreement with the calculations presented in Ref. [20].

Region of interest II: Island of α and p radioactivity "NE" of ^{100}Sn

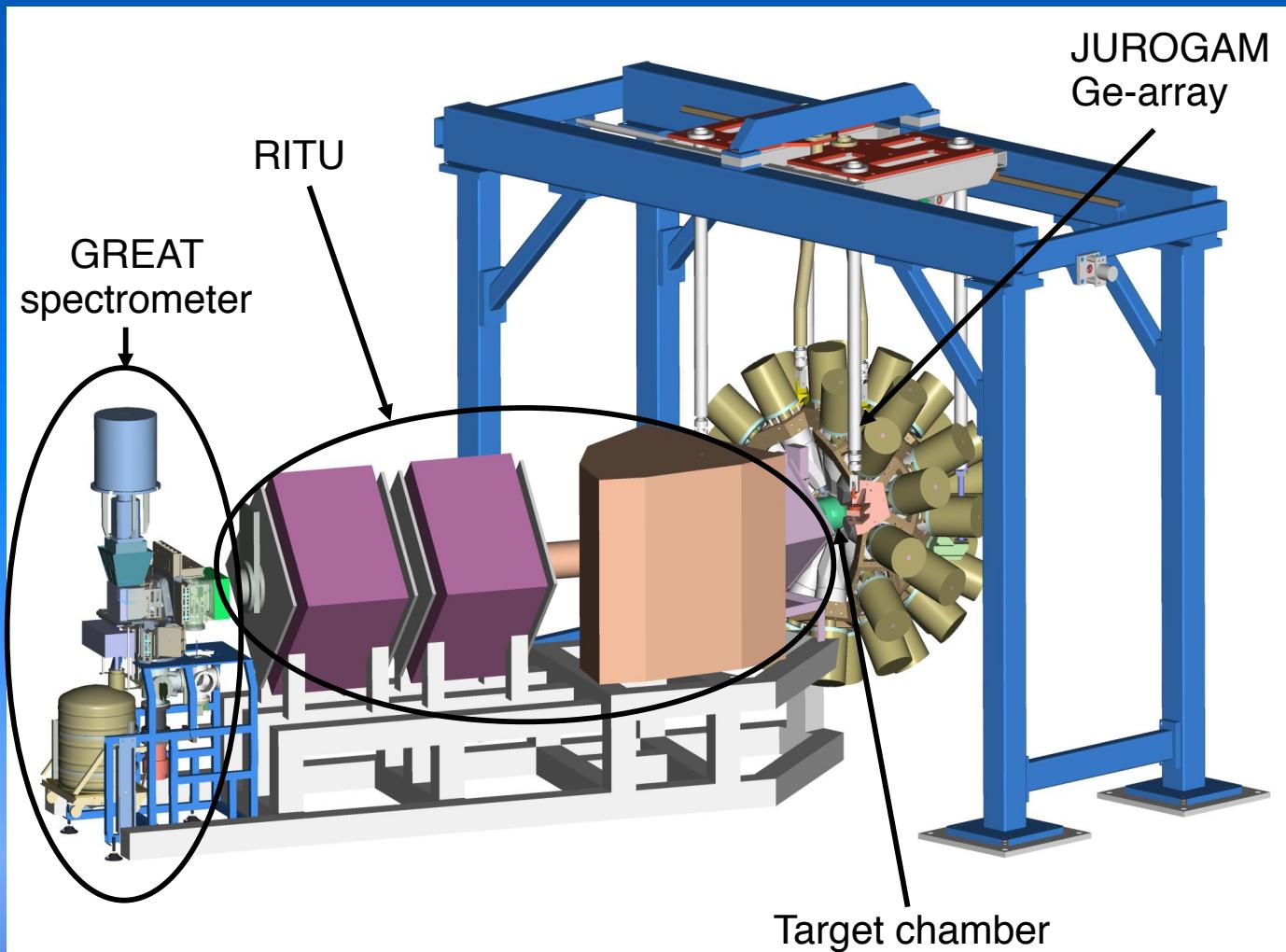
				x.s. ~ 50 nb			Ba 113 100 ms	Ba 114 0.43 s	Ba 115 0.45 s	Ba 116 1.3 s
							Cs 112 500 μs	Cs 113 17 μs	Cs 114 0.57 s	Cs 115 1.4 s
		x.s. ~ 25 nb								
	x.s. ~ 10 nb						Xe 109 13 ms	Xe 110 105 ms	Xe 112 2.7 s	Xe 113 2.8 s
			N = Z				I 108 36 ms	I 109 100 μs	I 111 0.65 s	I 112 3.42 s
				Te 105 0.70 μs	Te 106 70 μs	Te 107 3.1 ms	Te 108 2.1 s	Te 109 4.6 s	Te 110 18.6 s	Te 111 26.2 s
				Sb 103 1.5 μs	Sb 104 0.44 s	Sb 105 1.12 s	Sb 106 1.1 s	Sb 107 4.6 s	Sb 108 7.6 s	Sb 109 16.7 s
									Sb 110 24.0 s	Sb 111 75 s
	Sn 100 0.94 s	Sn 101 3 s	Sn 102 3.8 s	Sn 103 7.0 s	Sn 104 20.8 s	Sn 105 34 s	Sn 106 2.1 m	Sn 107 2.9 m	Sn 108 10.3 m	Sn 109 18.0 m
										Sn 110 4.11 h

**Recoil-decay tagging (RDT) *) has become a crucial tool
for structural studies of heavy, proton rich nuclei**

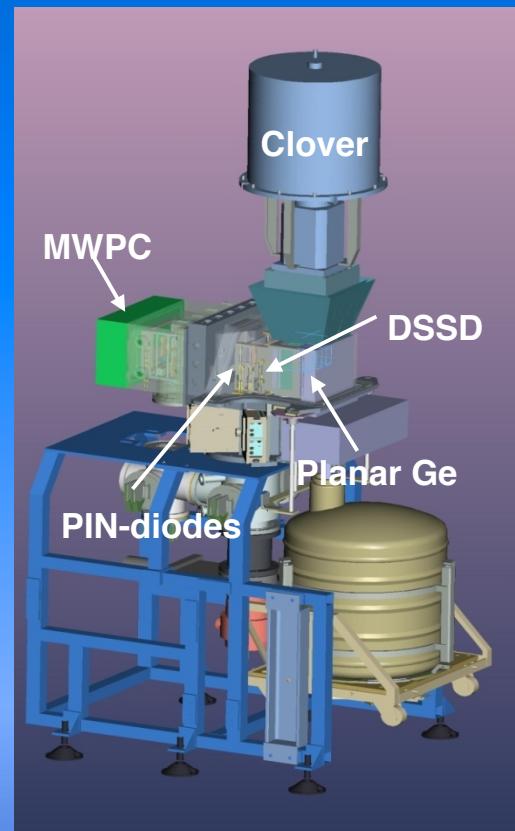
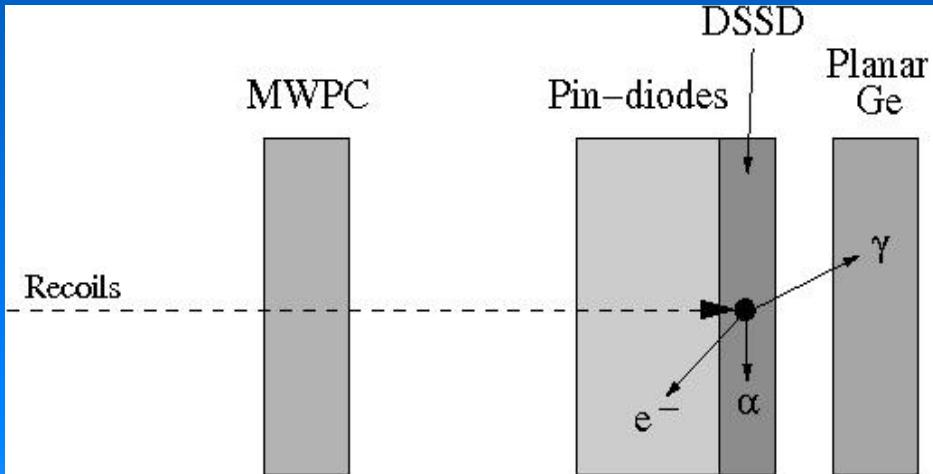
- **Recoil-decay tagging spectroscopy (as we know it nowadays)
started in the $A \sim 100$ ($^{108,109}\text{Te}$) region E.S. Paul et al.**
- **Extremely low production cross sections prevented further exploration**
- **Technical advances (RITU + GREAT, TDR ...) were needed to proceed further**
- **Technique now getting ready for high-intensity beams**

***) R.S. Simon *et al.*, Z.P.A. 325, 197 (1986): NaI + SHIP @ GSI
E.S. Paul *et al.*, P.R.C. 51, 78 (1995): Eurogam (45 HPGe) + DRS @ Daresbury**

Experimental set-up at JYFL; Univ. of Jyväskylä Cyclotron Laboratory



Gamma Recoil Electron Alpha Tagging - GREAT



MWPC - Multi Wire Proportional Counter: Recoil discriminator

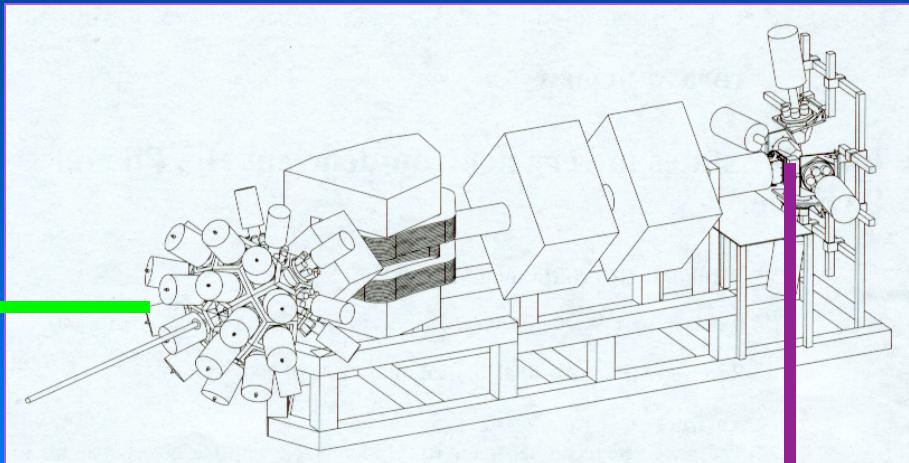
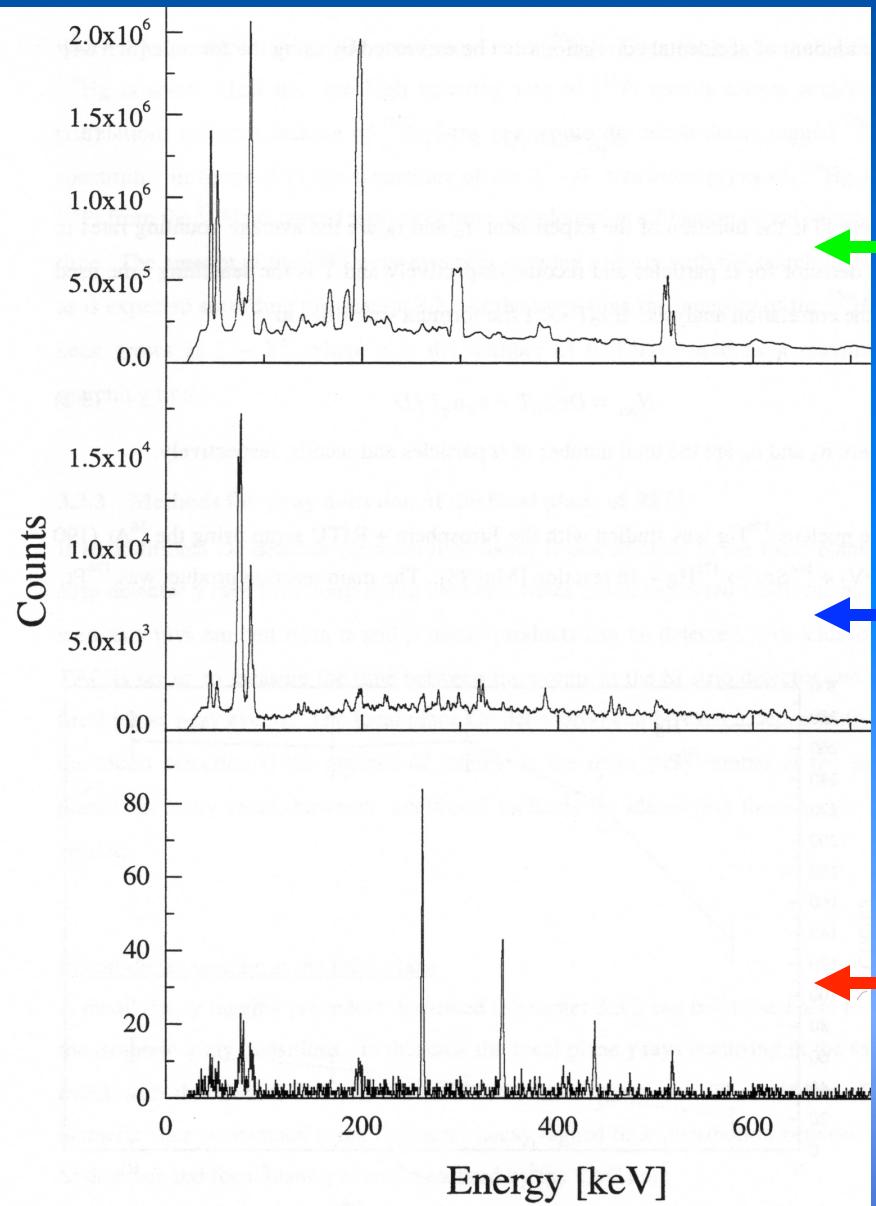
PIN-diodes: Detection of β -particles and conversion e^-

DSSD - Double-sided Silicon Strip Detector:

Charged particle detection (alpha, proton)

Planar Ge and Clover: Detection of delayed gamma rays
following radioactive decays or from isomeric states

RDT selectivity

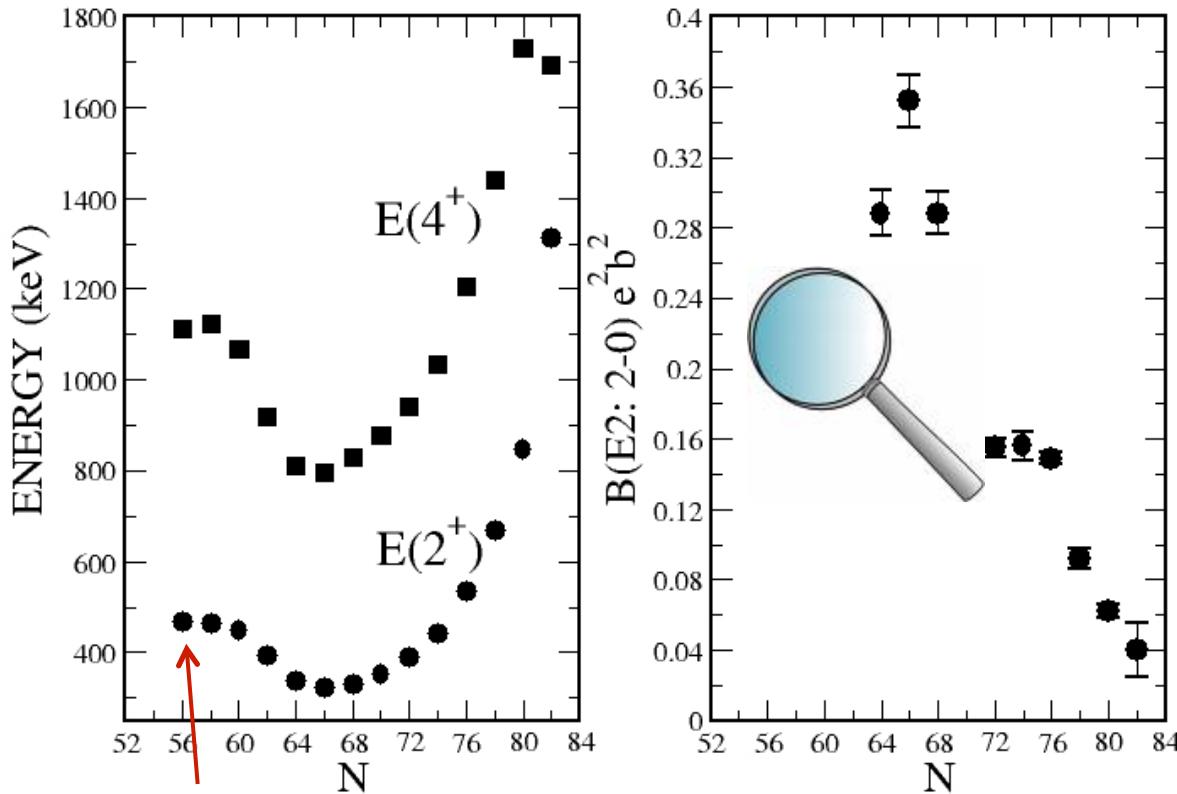


$\gamma - recoil$

$\gamma - recoil - \alpha$

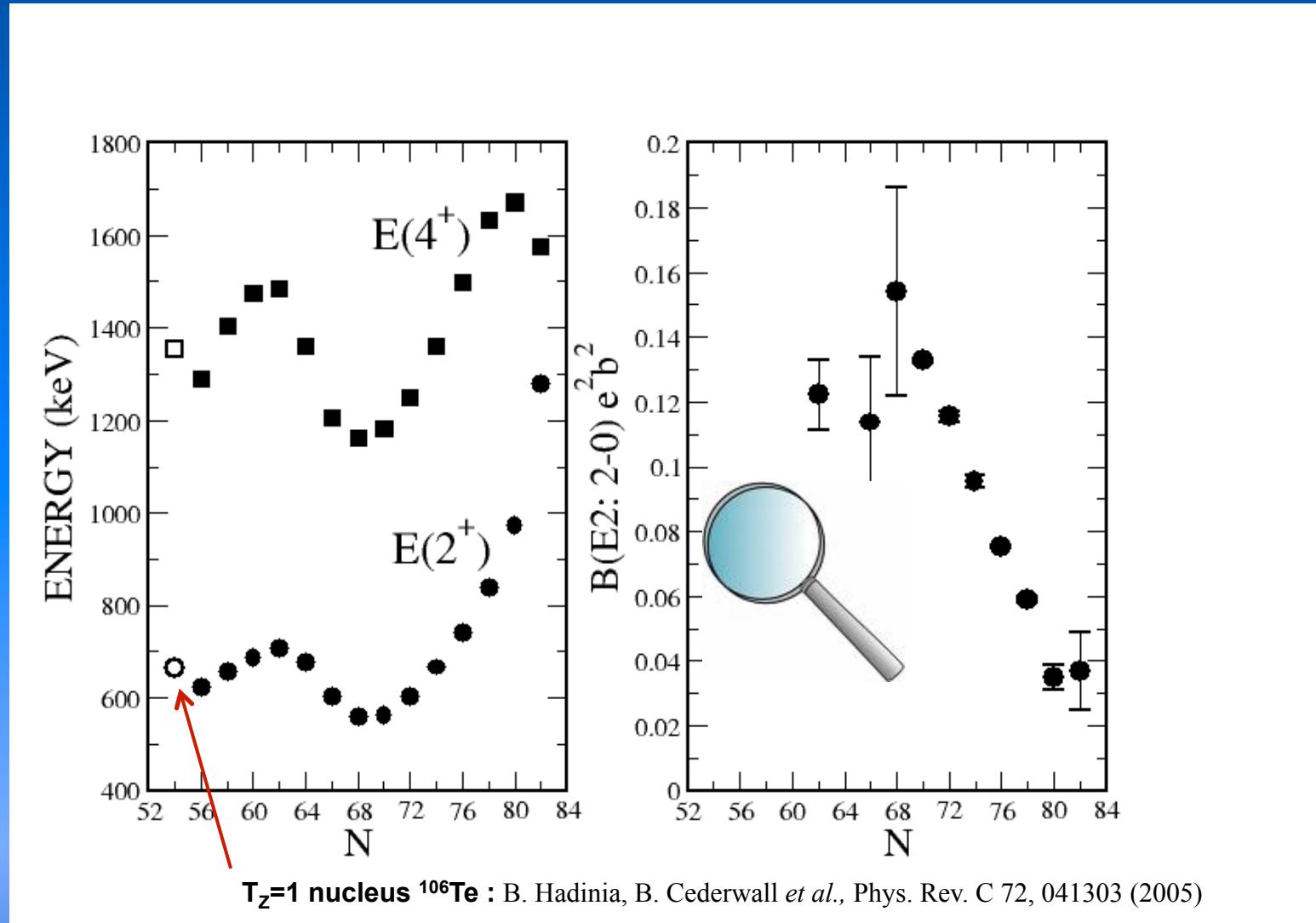
Evidence for enhanced collective strength in Te and Xe nuclei approaching N=Z?

Xe experimental $E(2^+)$ and $B(E2; 2^+ \rightarrow 0^+)$ systematics

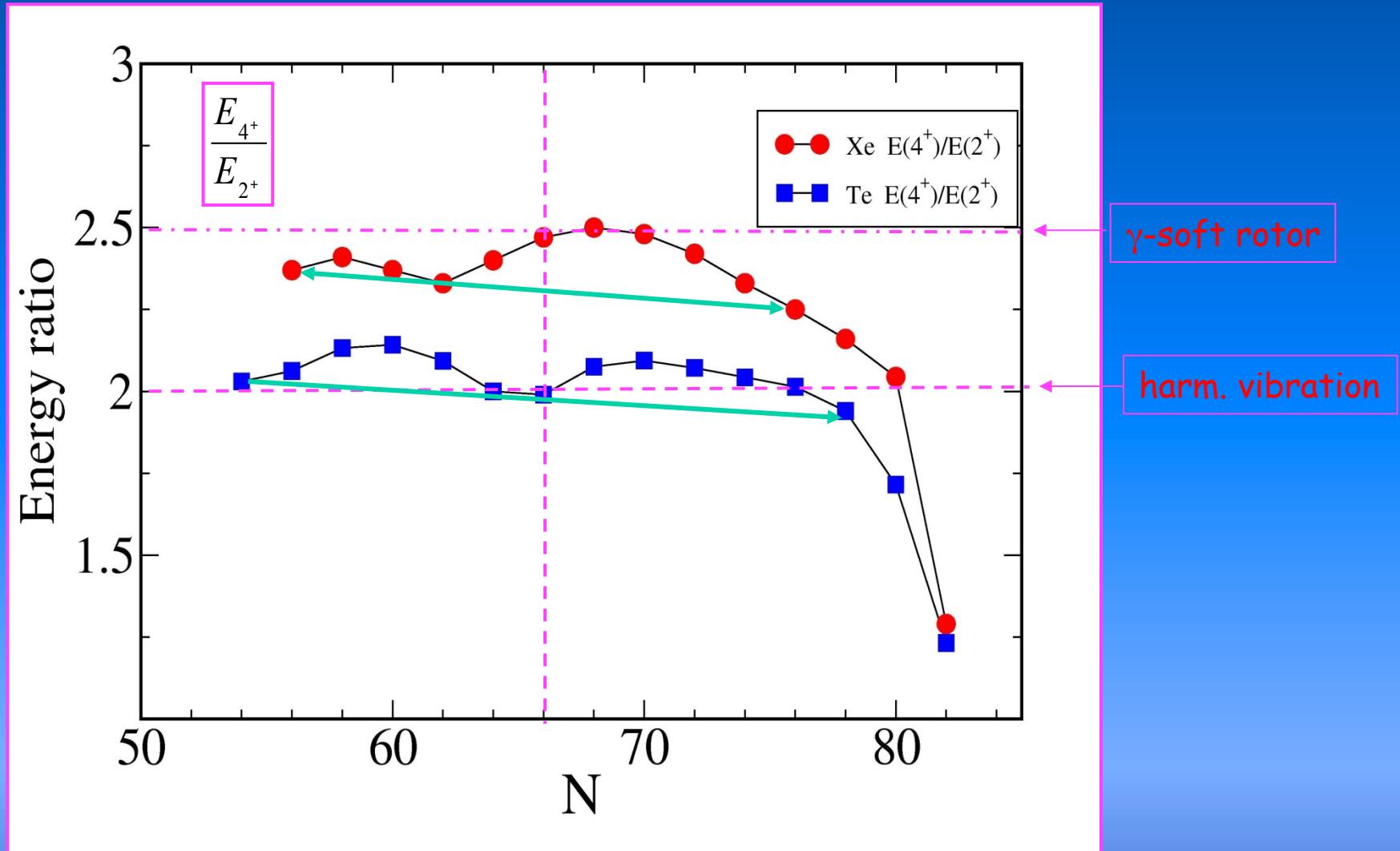


$T_z=1$ nucleus ^{110}Xe ; M. Sandzelius, B. Hadinia,, B. Cederwall *et al.*, Phys. Rev. Lett. 99, 022501 (2007)

Te experimental $E(2^+)$ and $B(E2; 2^+ \rightarrow 0^+)$ systematics

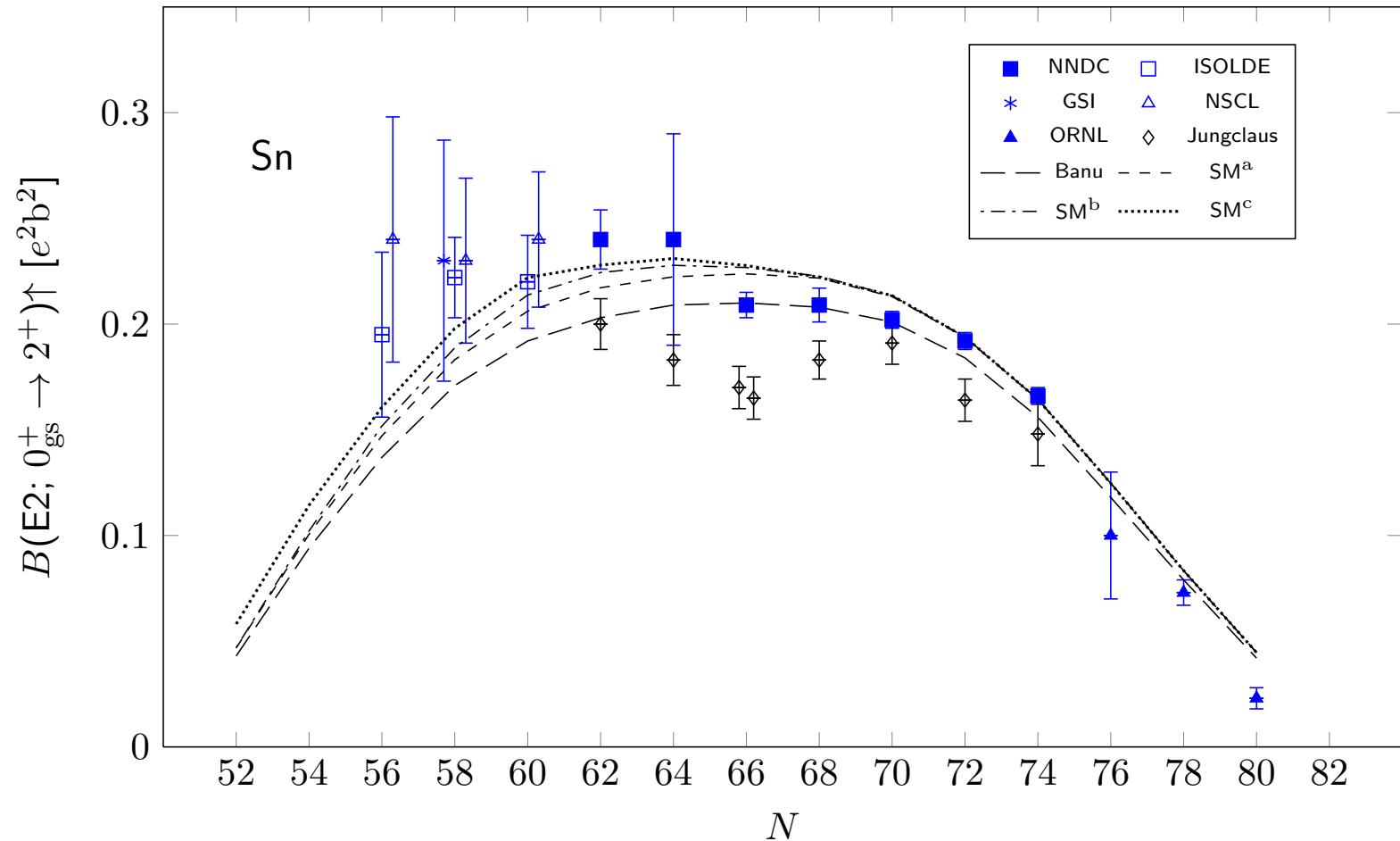


Xe and Te energy ratios

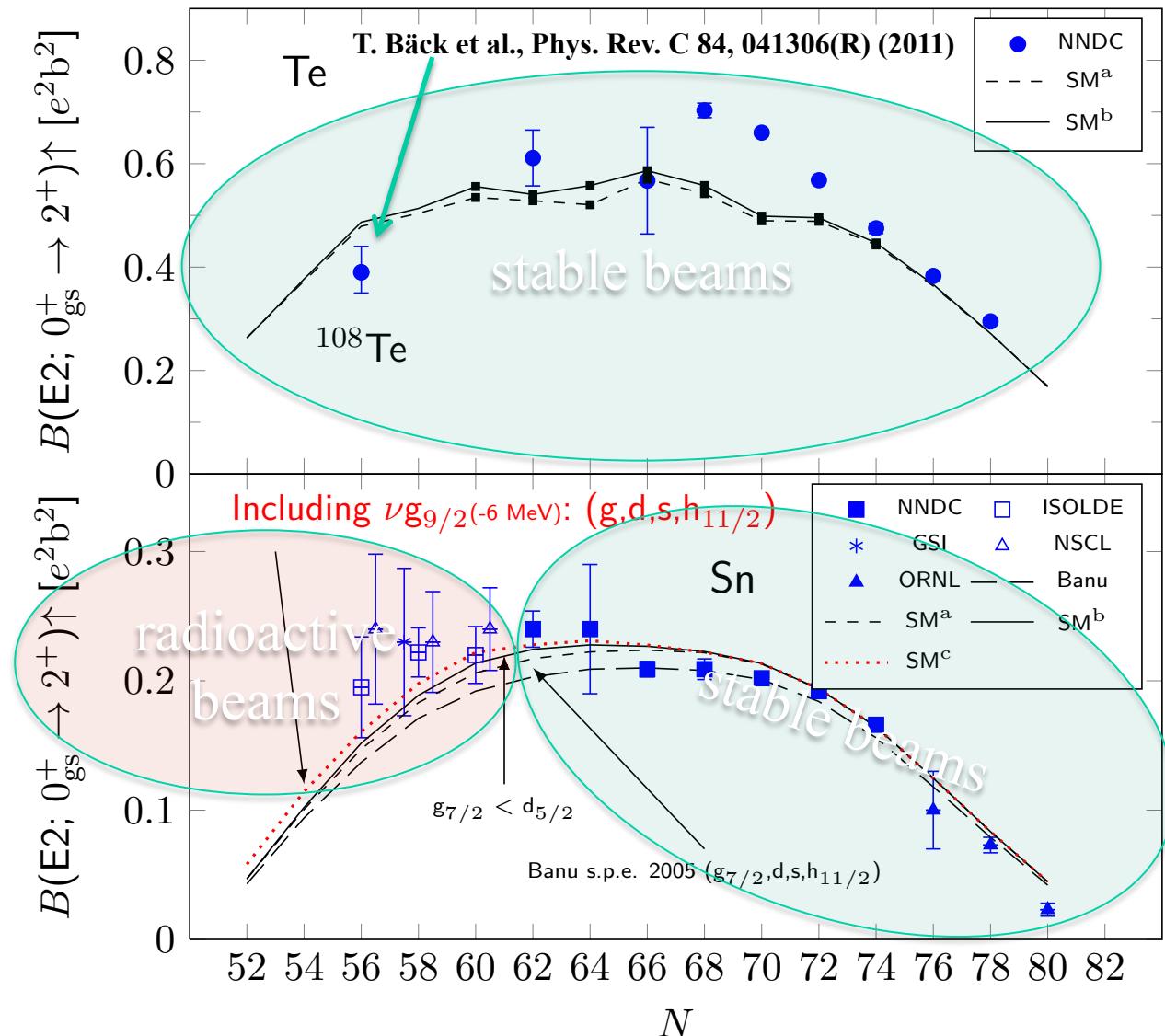


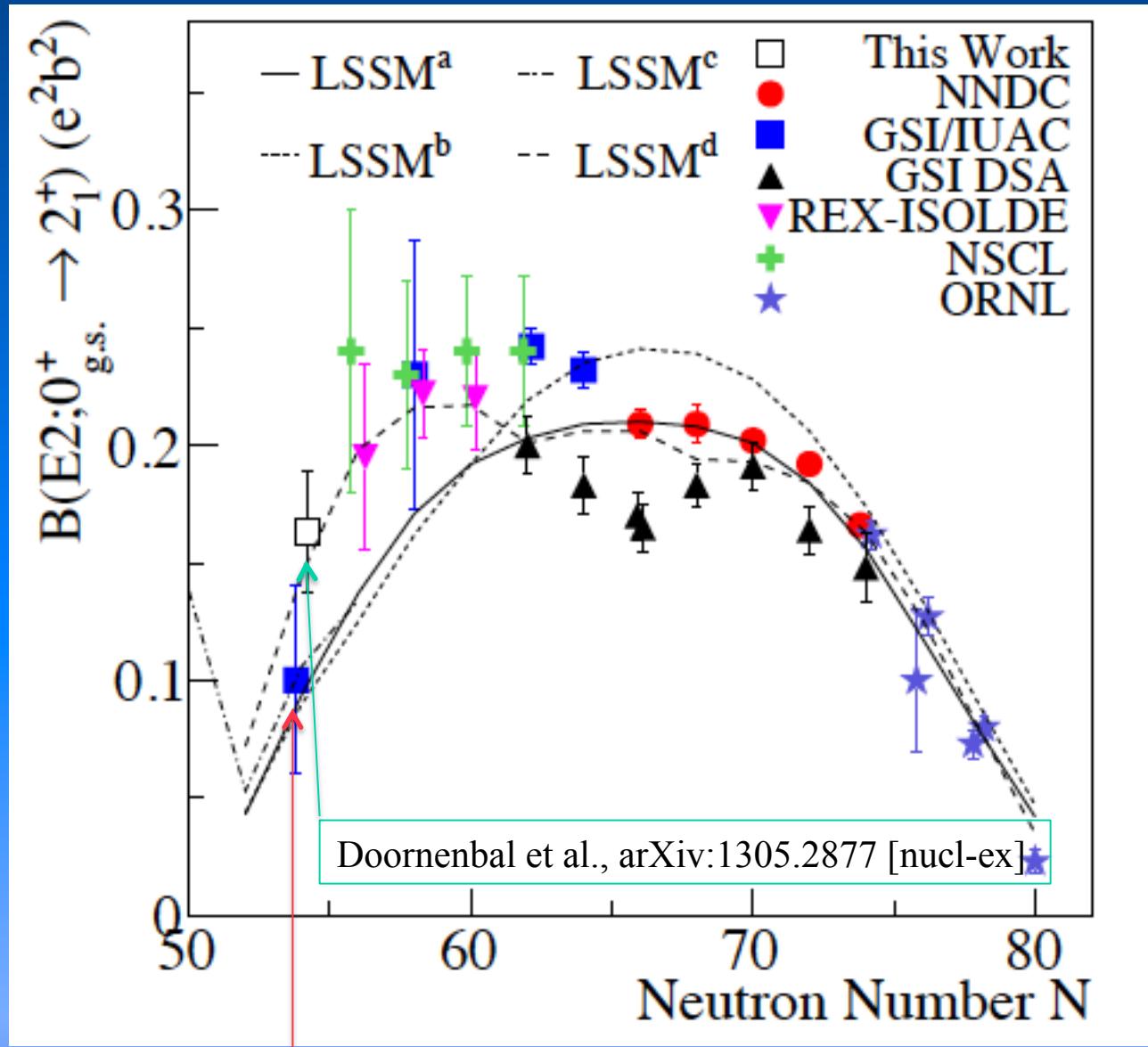
Evidence for weakening of the N=Z=50 shell closure??

Systematics of $B(E2; 0^+_1 \rightarrow 2^+)_\text{↑}$ for Sn isotopes Status pre 2013



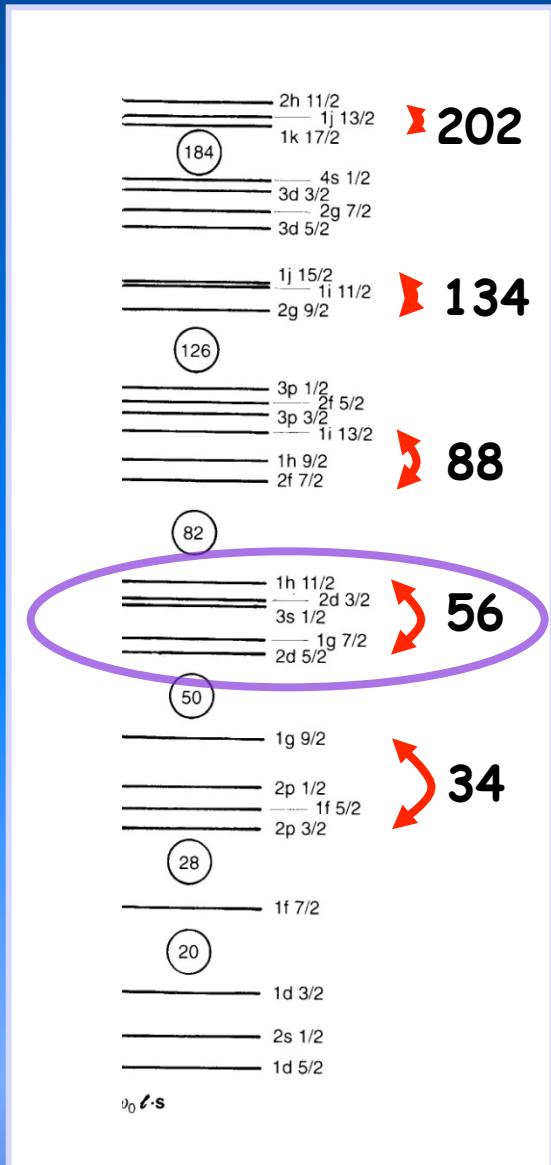
Systematics of $B(E2; 0_{gs}^+ \rightarrow 2^+) \uparrow [e^2 b^2]$ for Te and Sn isotopes





^{104}Sn , PRESPEC, Guastalla et al., PRL 110, 172501 (2013)

Enhanced octupole deformation and correlations near N=Z



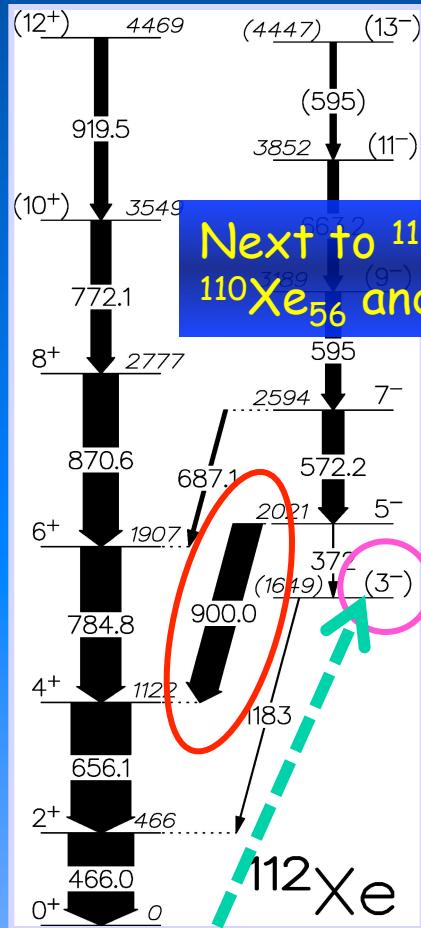
Strong octupole correlations are expected in nuclei where normal-parity single-particle states and intruder states differing by $\Delta l = \Delta j = 3$ are near the Fermi surface.

"Doubly-magic octupole-deformed" nucleus predicted by theory (J. Skalski, Phys Lett. **B**, 1990)

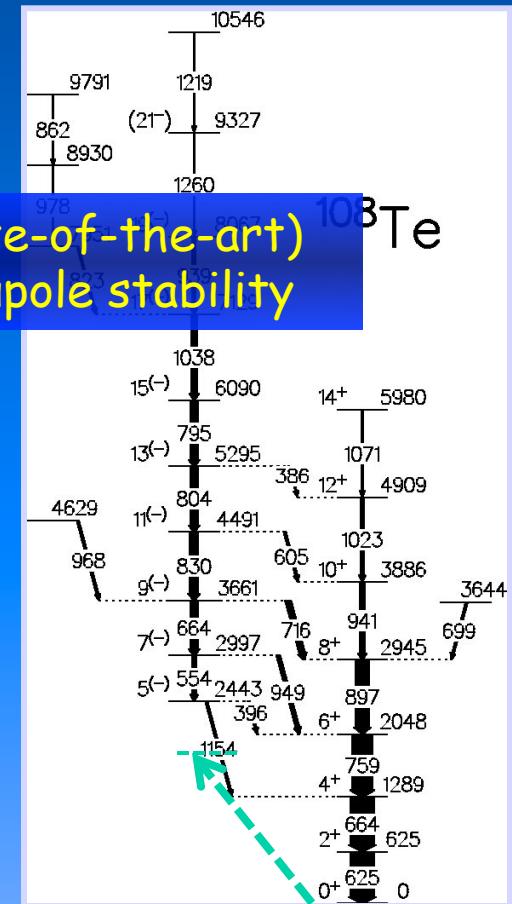
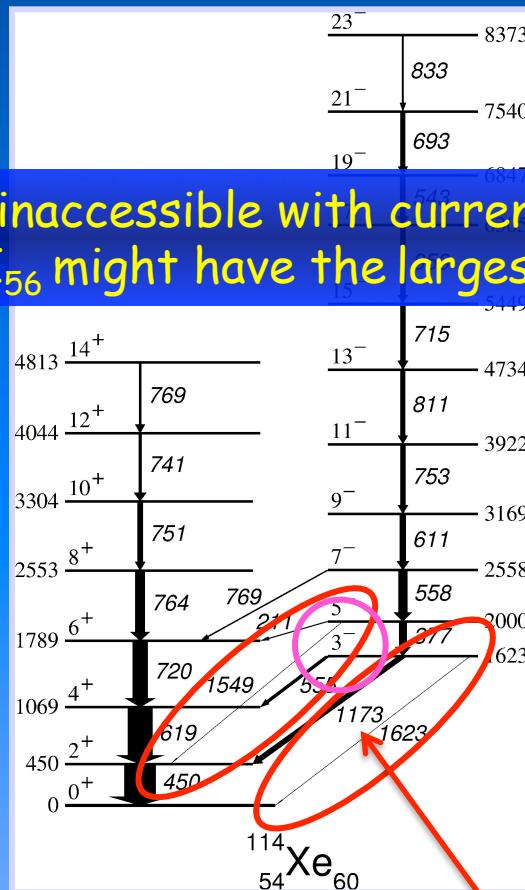


Coherent octupole correlations for neutrons and protons should occur near N=Z.
Can we observe additional enhancement due to dynamic np correlations ?

Enhanced octupole correlations in light Te-Xe nuclei due to dynamical np coupling ($\pi(v)d_{5/2}^- v(\pi)h_{11/2}$) ?



Next to ^{112}Ba (inaccessible with current state-of-the-art)
 $^{110}\text{Xe}_{56}$ and $^{109}\text{I}_{56}$ might have the largest octupole stability

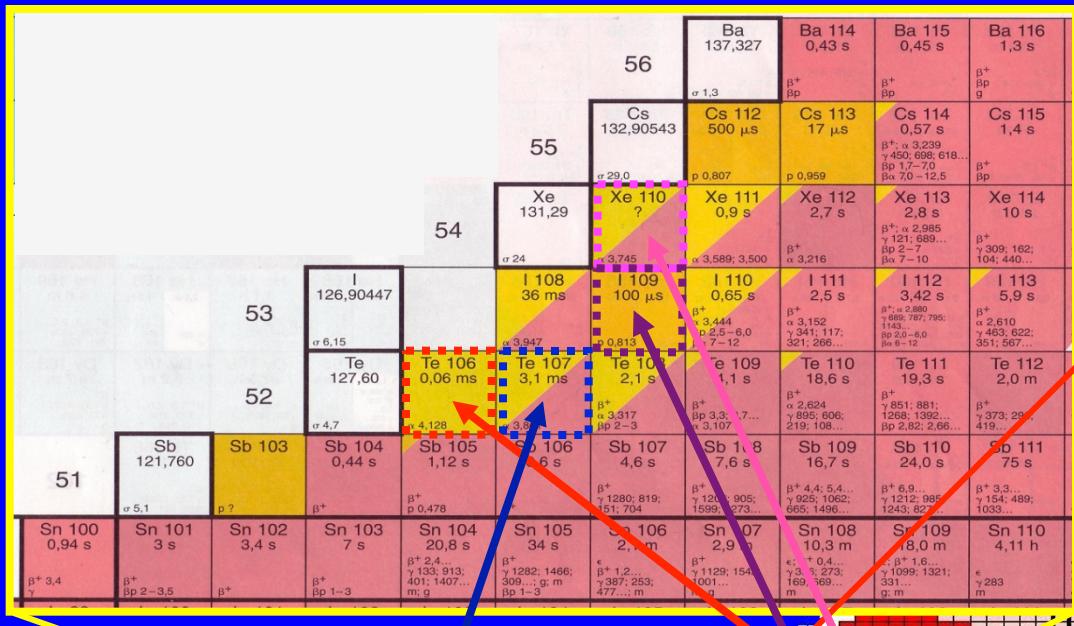


J.F. Smith et al.
Phys. Lett B523, 13 (2001)
(Gammasphere)

G. de Angelis et al.
Phys. Lett B535, 93 (2002)
(Euroball)

G.F. Lane et al.
Phys. Rev. C57, R1022 (1998)
(Gammasphere)

The “island” of alpha and proton radioactivity “NE” of ^{100}Sn → opportunity for RDT spectroscopy of exotic N~Z nuclei



N = Z

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Challenges

- Difficult separating fusion-evaporation residues:
Can only be populated in near-symmetric reactions



- Large fusion cross section, many evaporation channels open

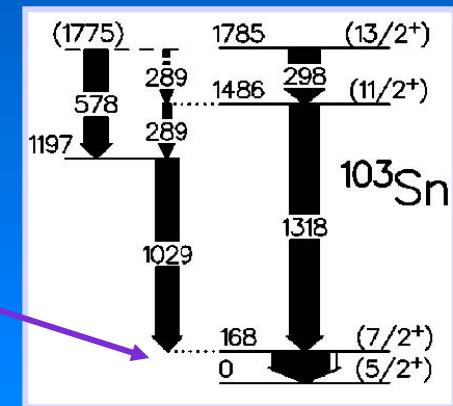
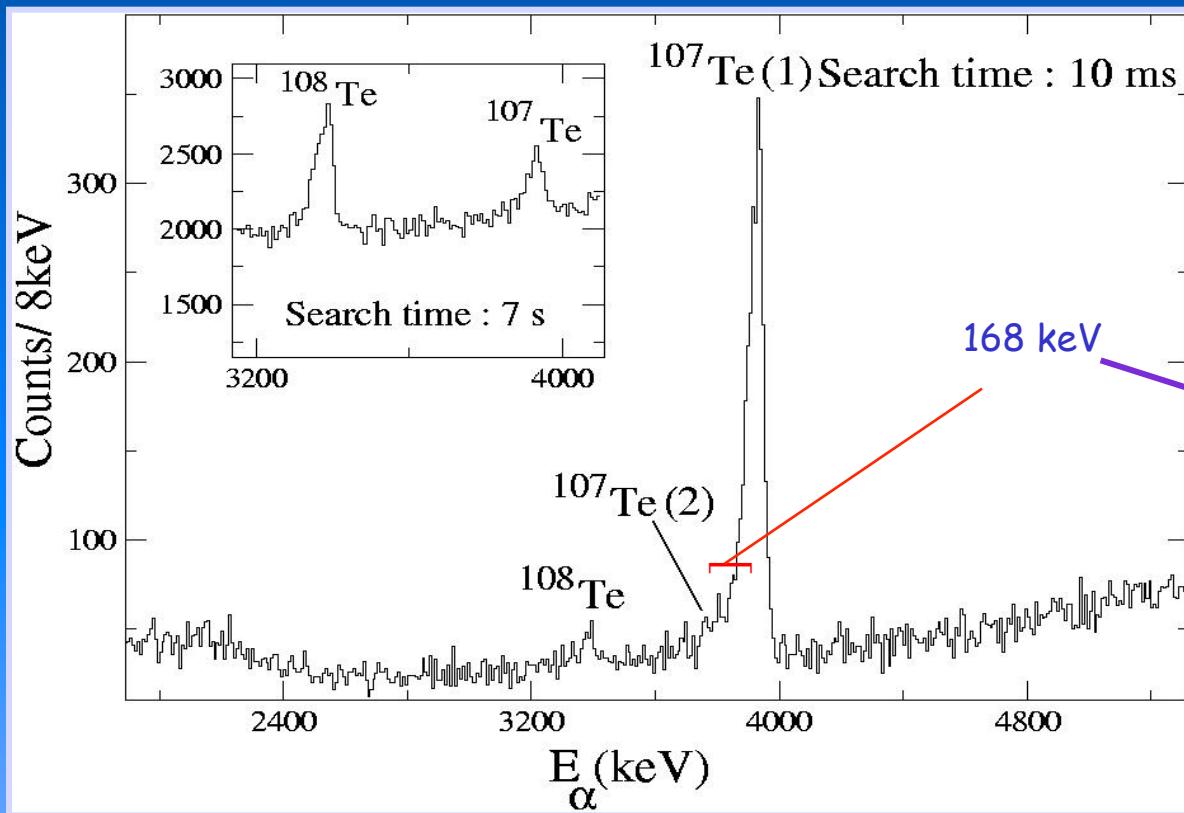


Plane detect



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^{58}Ni ($^{52}\text{Cr}, 3n$) $^{107}\text{Te}^*$ @ 187 MeV
 Recoil-correlated α decays @ RITU focal plane
 4pnA, ~5 days

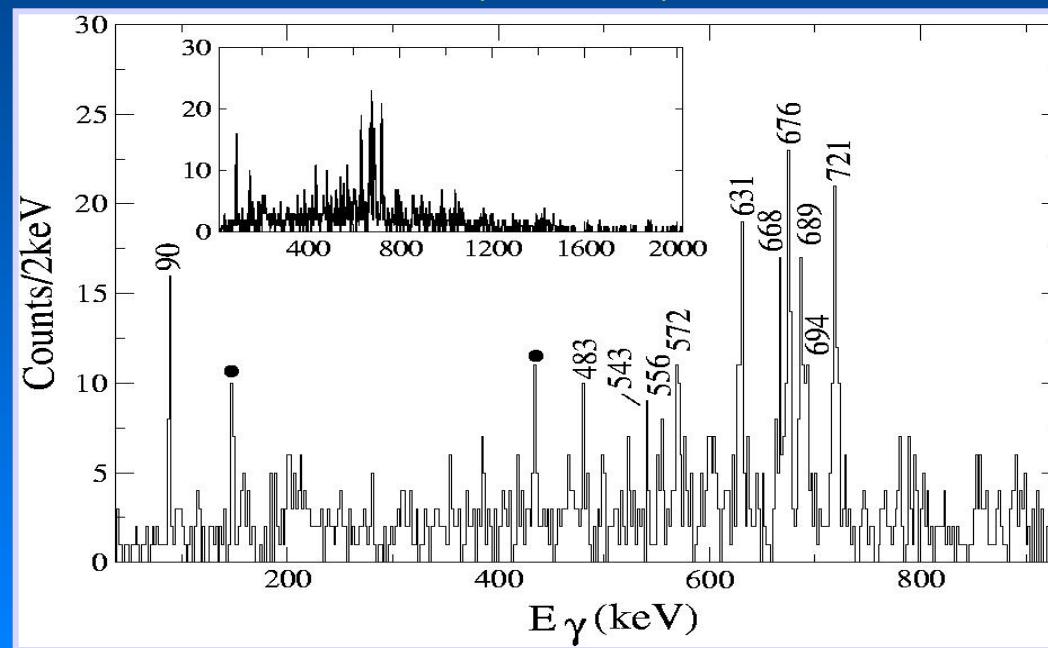


C. Fahlander et al.
[Phys. Rev. C63, 021307 \(2001\)](#)
[\(Euroball\)](#)

D. Seweryniak et al.
[Phys. Rev. C66, 051307 \(2002\)](#)
[\(Gammasphere + FMA\)](#)

Alpha-decay branching ratio : 70%
 Half life : 3.1ms
 $\sigma = 1 \mu\text{b}$

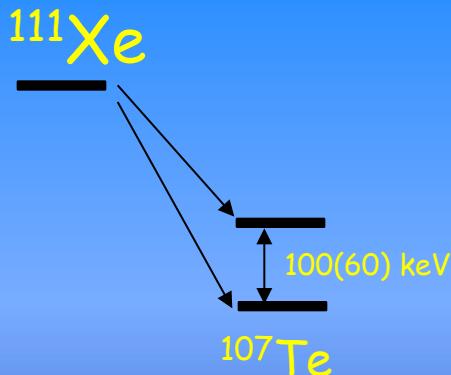
$^{58}\text{Ni} ({}^{52}\text{Cr}, 3n) {}^{107}\text{Te}^*$



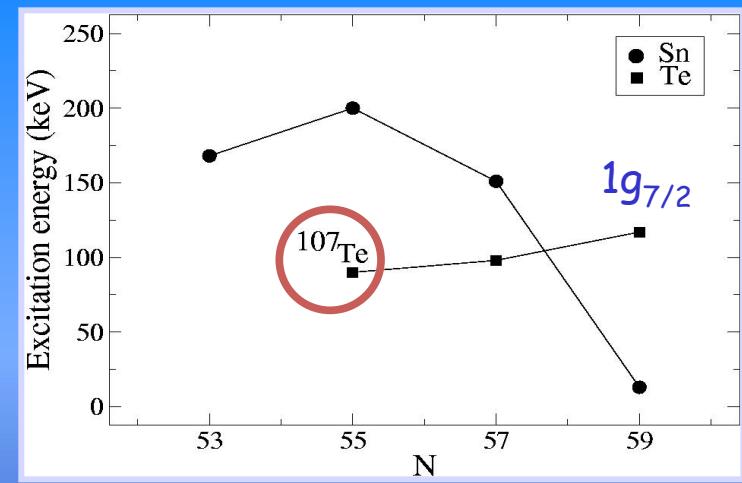
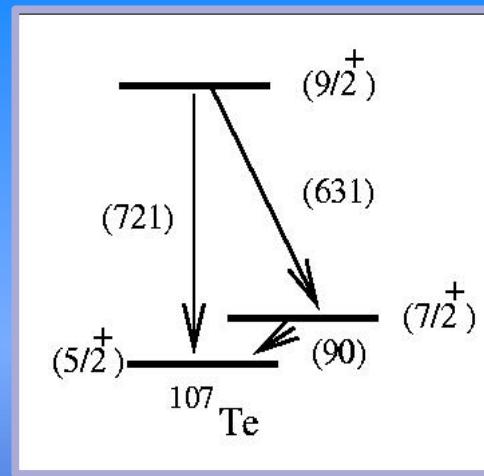
B. Hadinia et al.,
Phys. Rev. C. 2004

Recoil-decay correlated gamma-ray spectrum

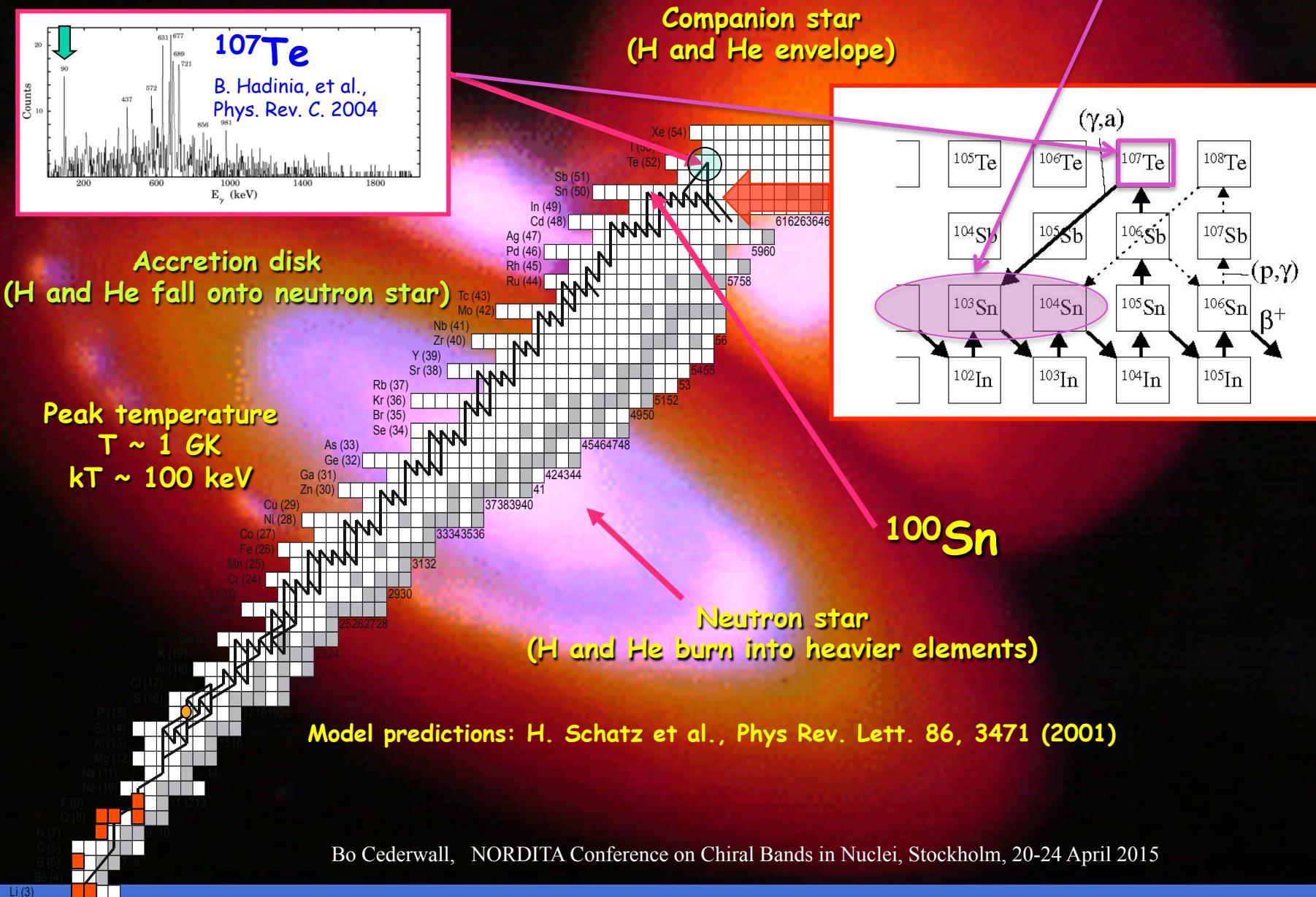
Tentative level scheme

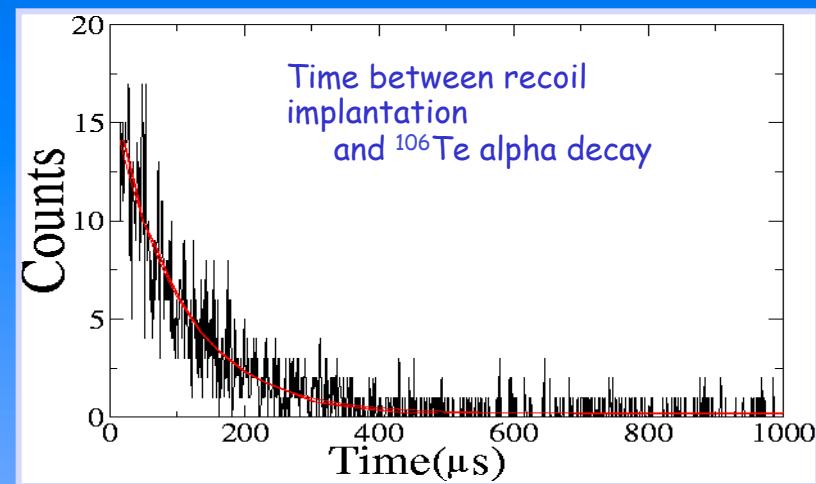
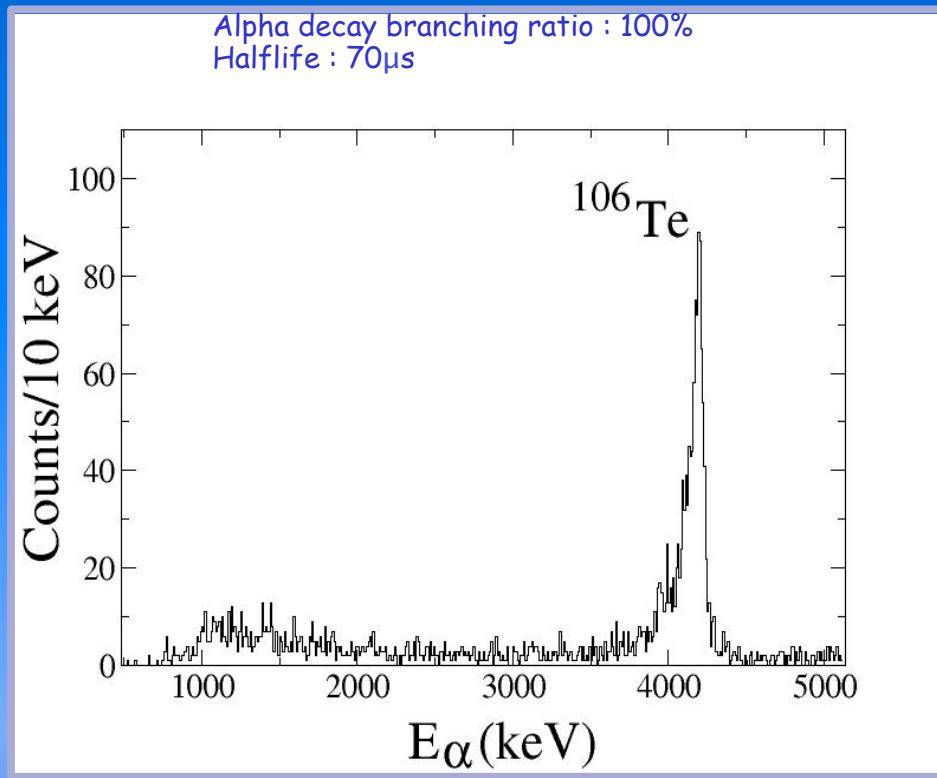
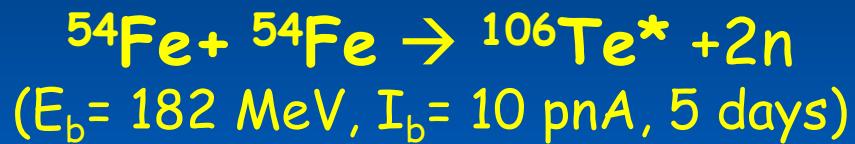


D. Schardt et al.
Nucl. Phys. A368, 153 (1981)



The predicted rp-process end point in X-ray bursters and accreting neutron stars

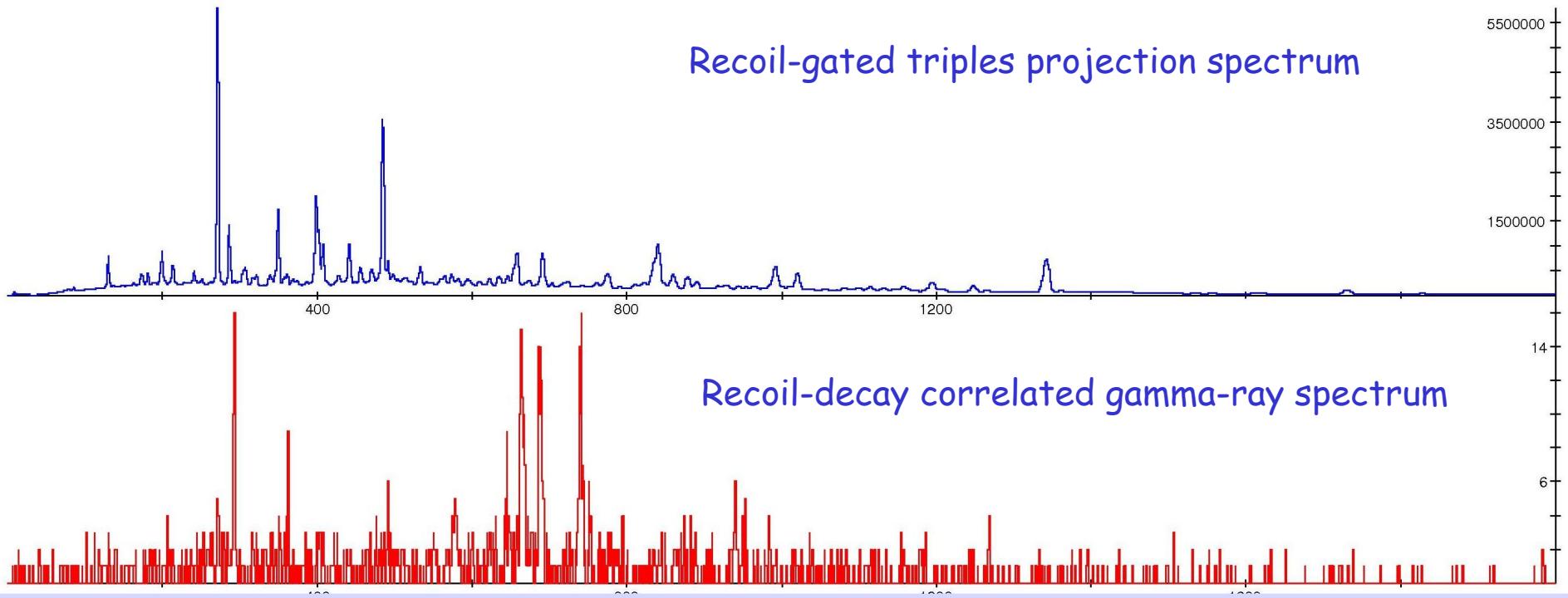




B. Hadinia, et al., Phys. Rev. C 72, 041303 (2005)
45

^{106}Te gamma rays

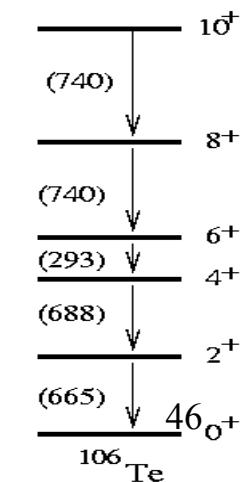
$\sigma = 25 \text{ nb}$ - (Then) a new limit for in-beam γ -ray spectroscopy!



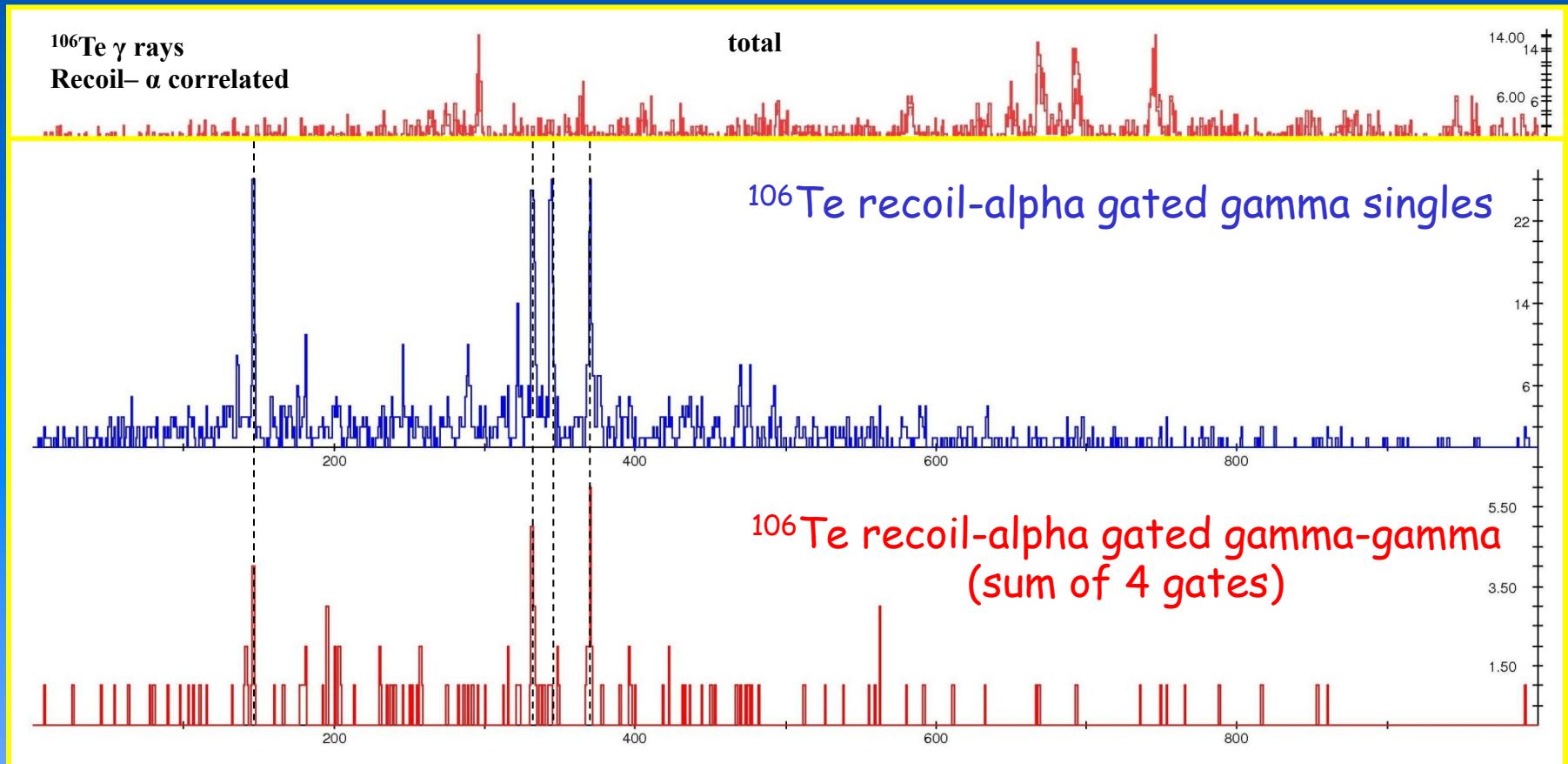
Selectivity: $\sim 10^{-7}$

Tentative level structure of ^{106}Te

B. Hadinia, et al., Phys. Rev. C 72, 041303 (2005)



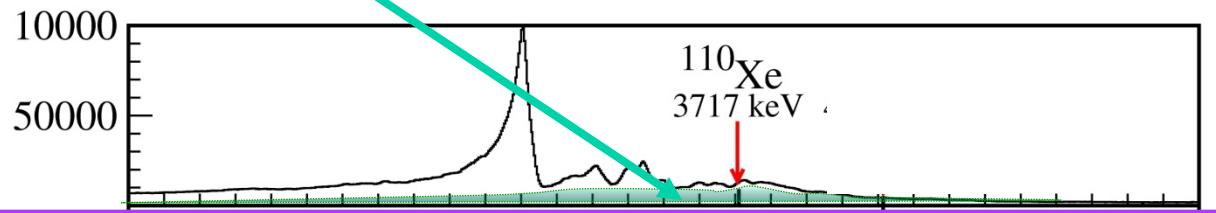
Gamma-gamma coincidences at $\sigma \sim 25$ nb



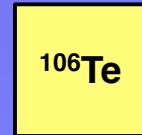
Identification of excited states in ^{110}Xe

M. Sandzelius, et al. Phys. Rev. Lett 2007

β -delayed protons



$E_\alpha = 3.72 \text{ MeV}$
 $t_{1/2} \approx 100 \text{ ms}$
 $b_\alpha \approx 70\%$



$E_\alpha = 4.13 \text{ MeV}$
 $t_{1/2} \approx 70 \mu\text{s}$
 $b_\alpha = 100\%$



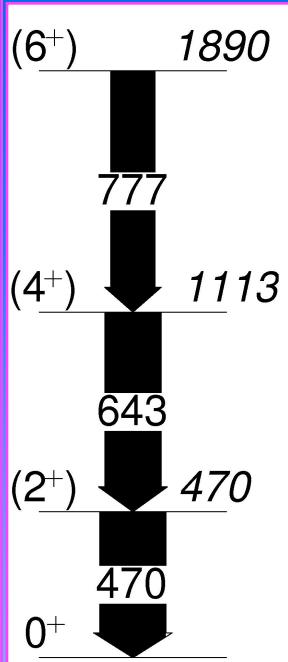
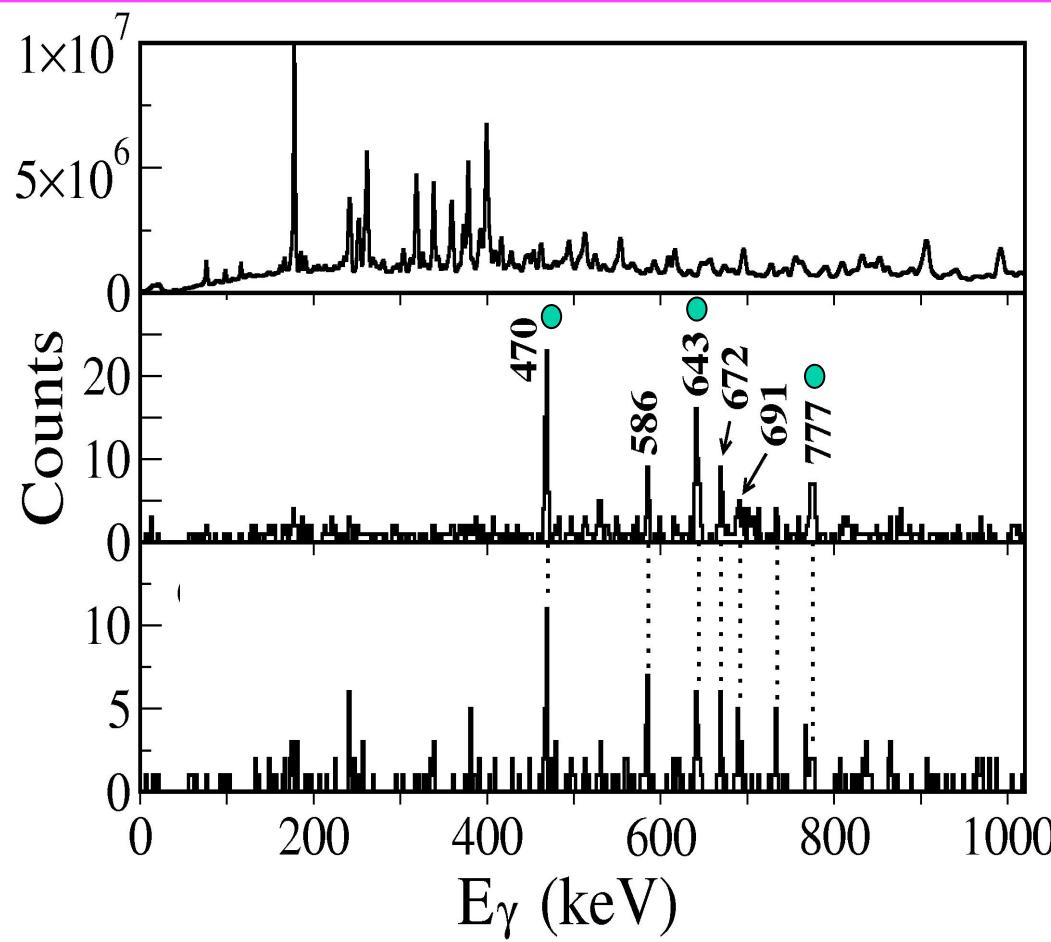
First evidence for superallowed α -decay at $T_z=1$

Clean mother-daughter correlations essential for selecting the ^{110}Xe nuclei

M. Sandzelius, et al. Phys. Rev. Lett 2007

β -delayed protons

$^{110}\text{Xe } \alpha$ decay
Recoil – α correlation time 330 ms

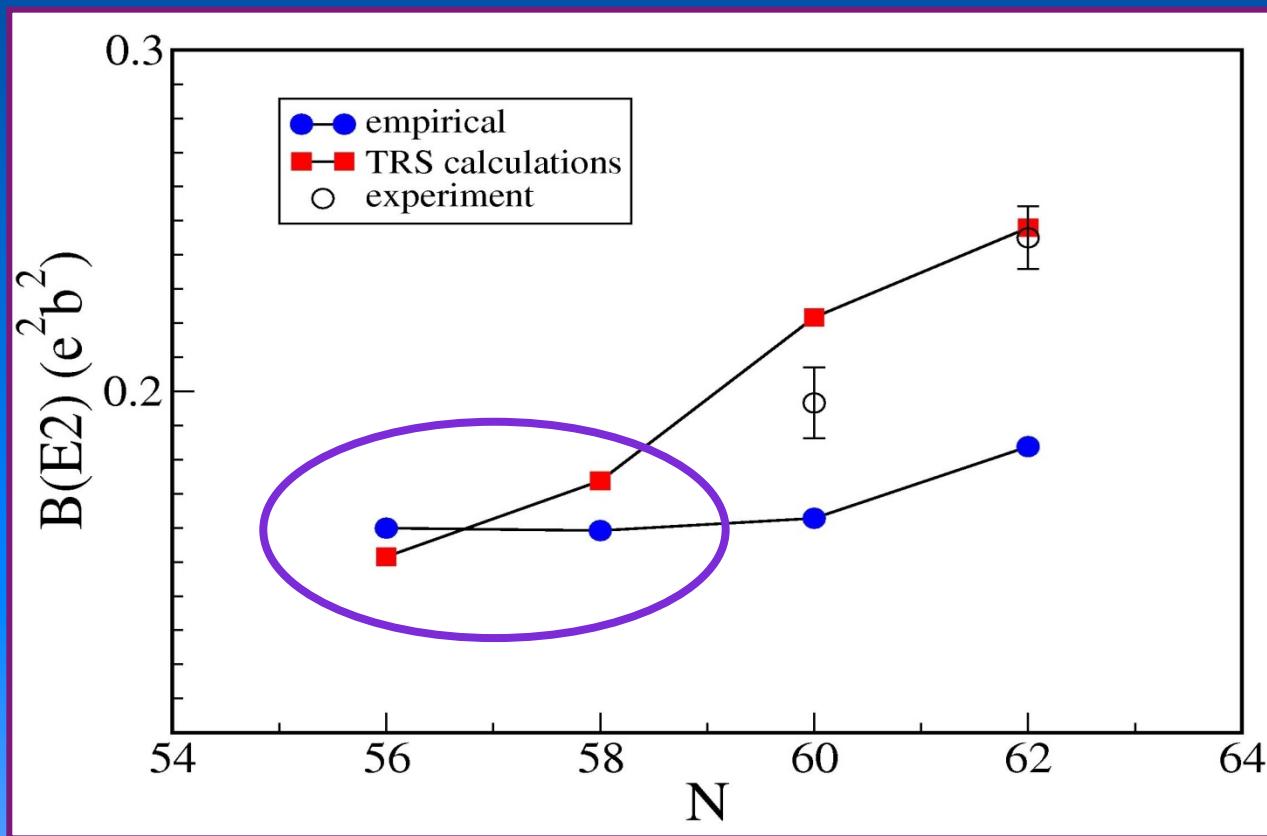


Total recoil-tagged

^{110}Xe mother-daughter correlated

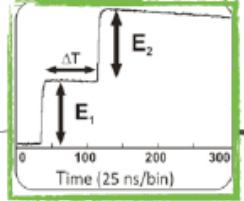
^{110}Xe mother-daughter correlated, sum $\gamma\gamma$

Comparing theory with experimental $B(E2)$ values (Raman estimates) for extremely neutron deficient Xe isotopes



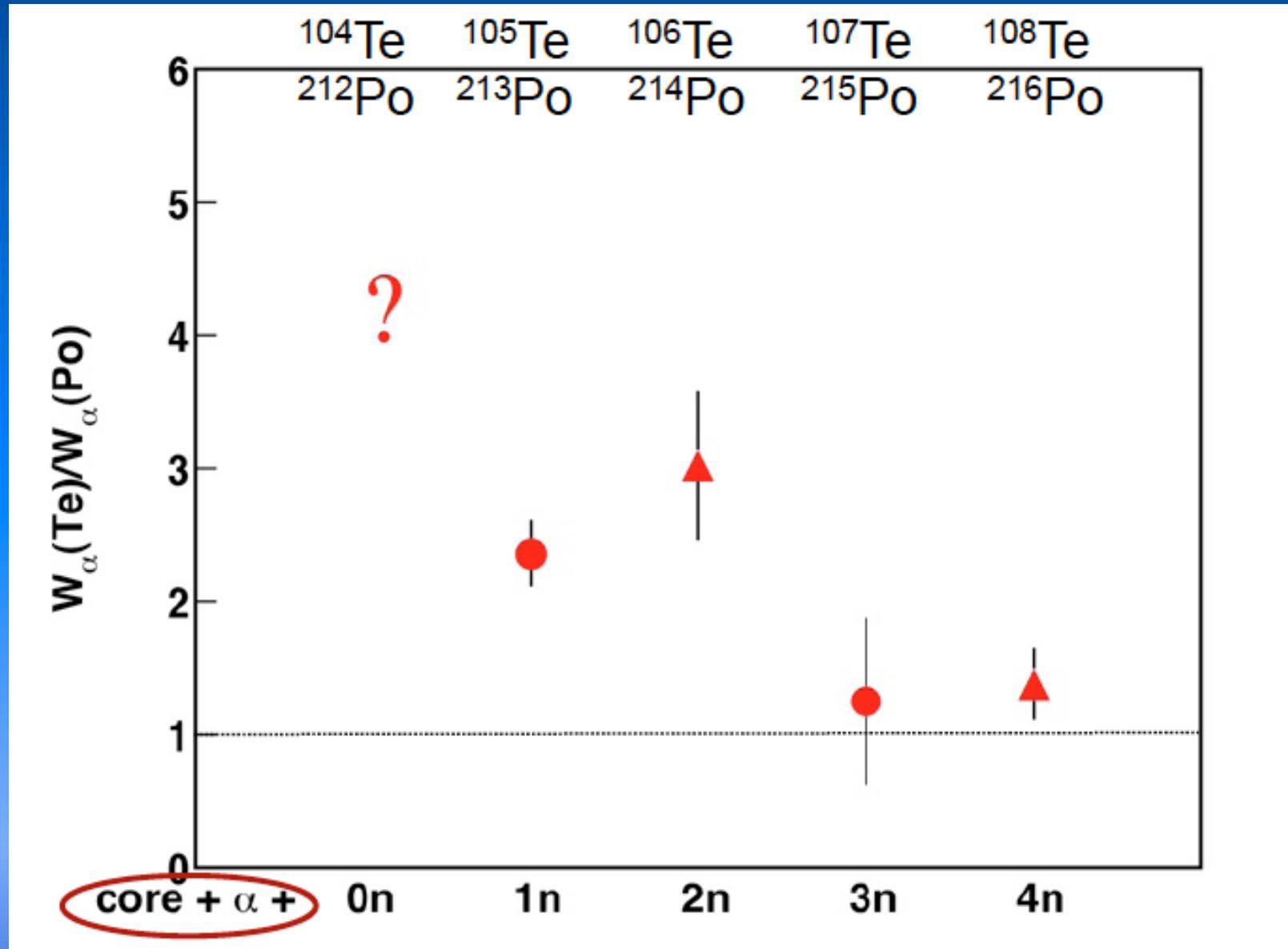
- $B(E2)$ values are a measure of nuclear collectivity
- Theoretical models predict a *decrease* in $B(E2)$ values for decreasing N
- The empirically deduced values^{*)} reveal a leveling off and a even a small *increase* of the $B(E2)$ value for ^{110}Xe

Decay spectroscopy: "High-intensity" stable beams and new detection schemes

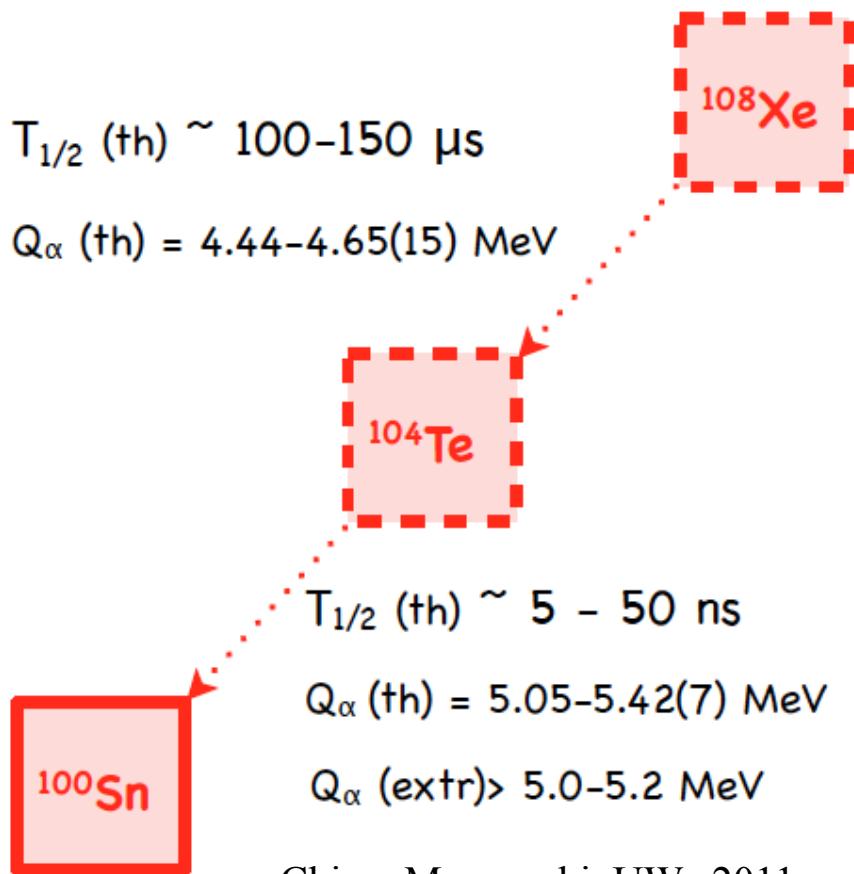
$^{109}\text{Xe} \rightarrow ^{105}\text{Te} \rightarrow ^{101}\text{Sn}$ Exp I Liddick et al., PRL97 082501 (2006)	$^{109}\text{Xe} \rightarrow ^{105}\text{Te} \rightarrow ^{101}\text{Sn}$ Exp II I.G. Darby et al, Phys. Rev. Lett. 105 , 2010
$^{54}\text{Fe} + ^{58}\text{Ni} \rightarrow ^{112}\text{Xe}^*(3n) ^{109}\text{Xe}$ ~8 part. nA standard target	$^{54}\text{Fe} + ^{58}\text{Ni} \rightarrow ^{112}\text{Xe}^*(3n) ^{109}\text{Xe}$ ~50 part. nA rotating target
DSSD (40x40, 65 μm) + veto detectors	DSSD (40x40, 65 μm) + veto detectors + HPGe (11.2% @ 123 keV)
 Digital DAQ: alpha catcher mode	

Chiara Mazzocchi, UW, 2011

Superallowed alpha decay: enhanced preformation
- Strongly connected to np correlations in N=Z systems



Chiara Mazzocchi, UW, 2011⁵²



Theory:

Xu and Ren, PRC74, 2006;
Mohr, EPJA31, 2007

Extrapolated limits:

Liddick et al., PRL97, 2006;
Seweryniak et al., PRC73, 2006

Experimental opportunities at the end of the N=Z line:

AGATA @ GANIL → SPIRAL2

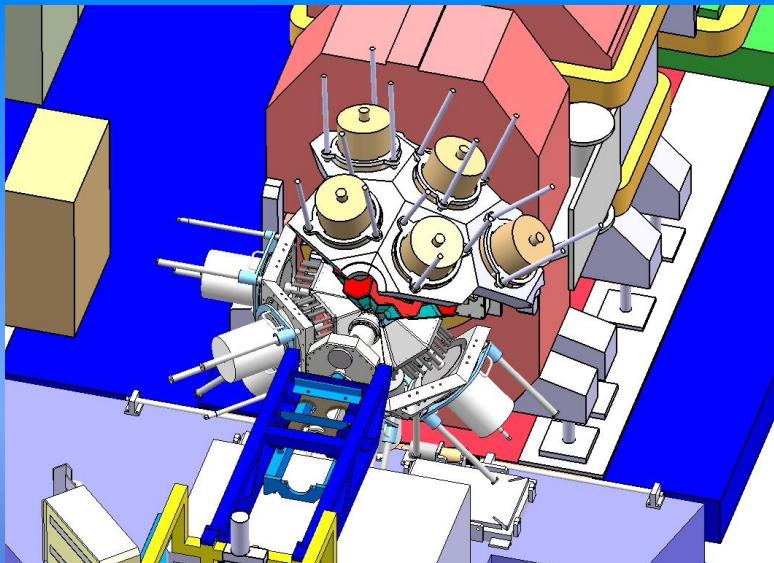
- RDT/RDDS using VAMOS in gas-filled mode
- L.E. Coulex
- Spectroscopy
- np-transfer



RIKEN (DALI2+ZDS) ?

AGATA-HISPEC @ FAIR

- H.E. Coulex, knockout



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"Conclusions"

Proton rich and N=Z nuclei - major potential for important discoveries

- Neutron-proton correlation effects
 - isoscalar pair modes
 - superallowed alpha decays
 - development of collective excitations
 - GT strengths Not discussed here
 - Isospin symmetry Not discussed here
 - Exotic particle decay modes (2p, ...) Not discussed here
 - Influence from proton state continuum on structure and correlations Not discussed here
 - High-spin physics Discussed only briefly here
 - ...
- "High-intensity" stable beams coupled to improved detector instrumentation one important route forward!

np pair modes - outlook

- ^{100}Sn region a prime “laboratory” for investigating effects of strong residual neutron-proton correlations including T=0 pairing
- Strong need for mass measurements in N=Z nuclei beyond A \approx 60 to pin down ground-state isoscalar pairing issue
- Deuteron transfer reactions with RIBs (inverse kinematics) another promising route
- Angular momentum response in (super) deformed N=Z systems
- Isoscalar spin-aligned coupling scheme is a possible, different “paired phase” with characteristic features that need further experimental verification at the new facilities
- Need sharper theoretical predictions for reactions & structure



Thank You