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# Neutron stars/pulsars with THEIA

M Ángeles Pérez-García

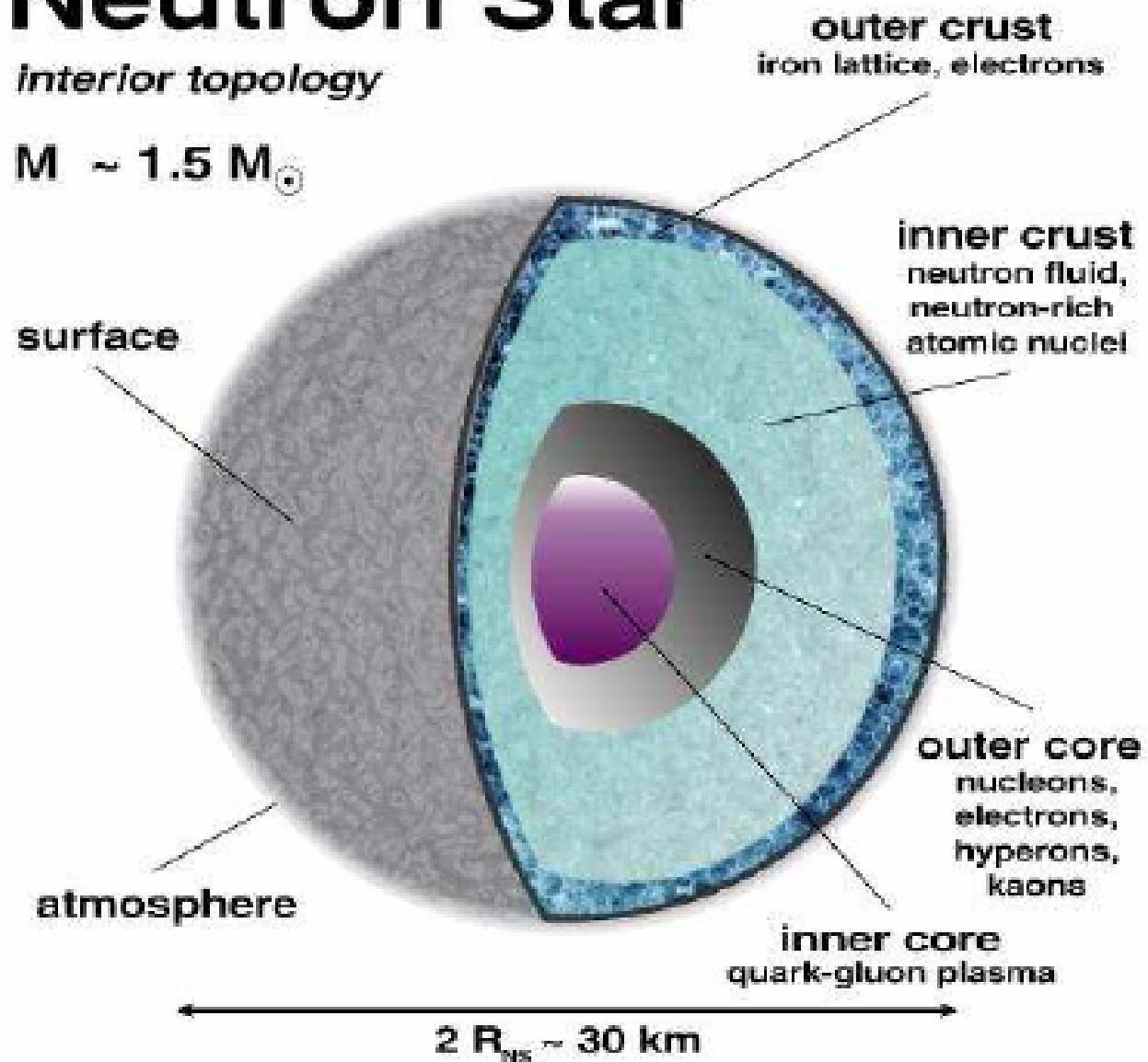
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University of Salamanca, Spain

# Neutron Star

*interior topology*

$M \sim 1.5 M_{\odot}$



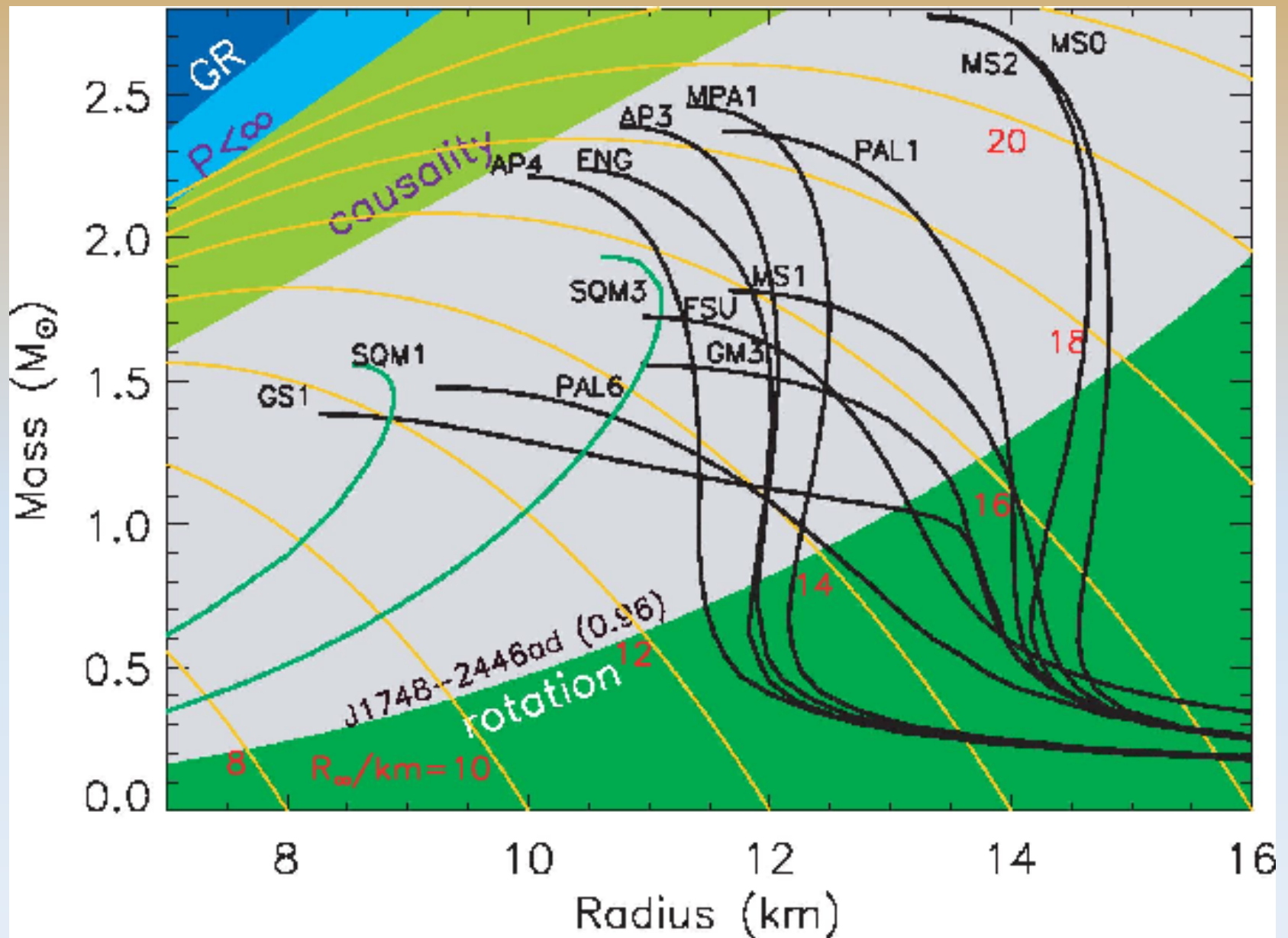
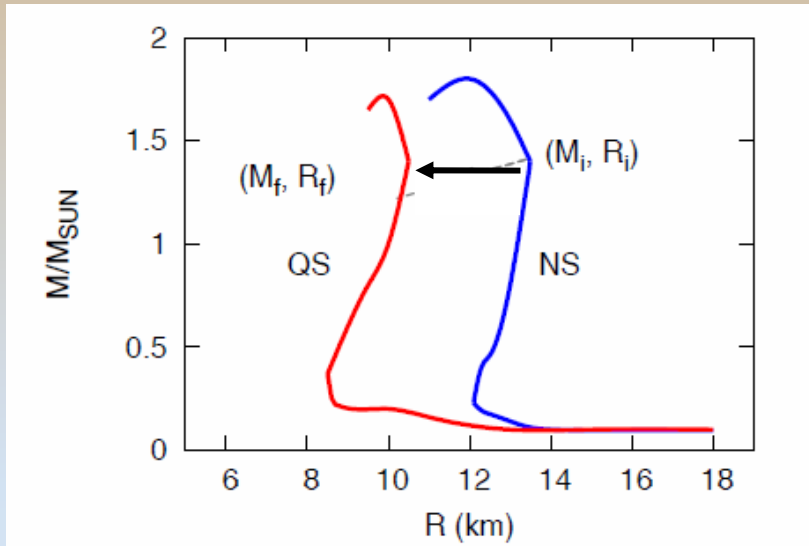


Fig courtesy of J. Lattimer

# Mass and radius determination



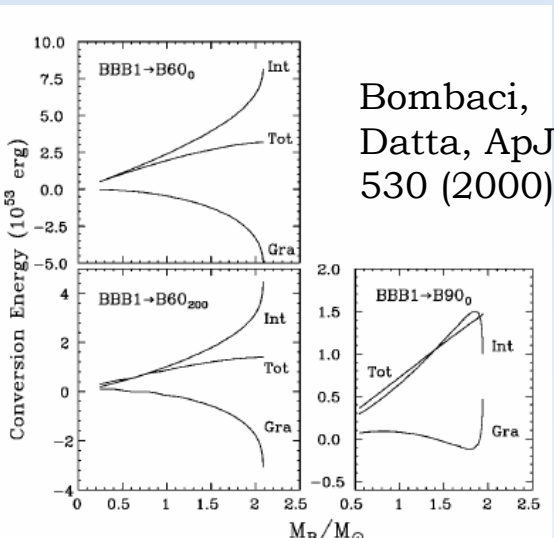
Perez-Garcia, Daigne, Silk, ApJ (2012)

- Major goal:

- Simultaneous Mass and Radius determination

- I.e. Hybrid vs pure neutron stars

- Transition to partial Quark content, Perez-Garcia, Silk, Stone PRL 105, 141101 (2010)

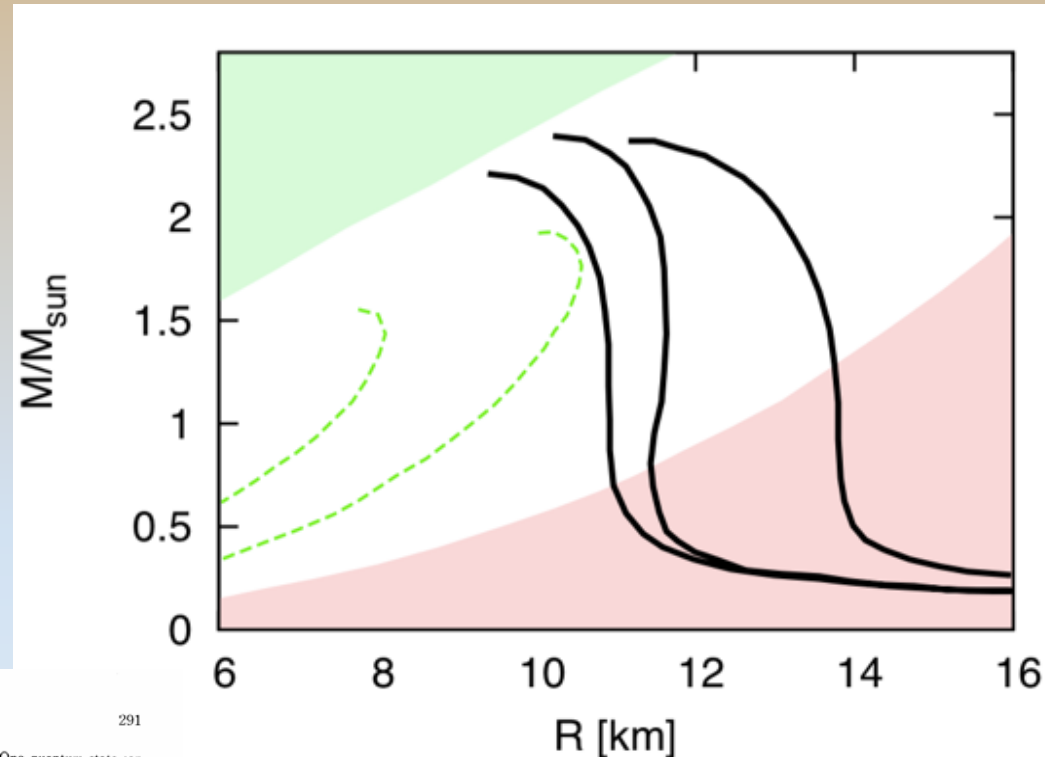


Bombaci,  
Datta, ApJ  
530 (2000)

$$\Delta E \approx GM_{NS}^2 \left( \frac{1}{R_{QS}} - \frac{1}{R_{NS}} \right) \approx 10^{53} \text{ erg}$$

# NS vs QS

- Baryons vs deconfined quark content
- Finite strangeness content
- Energetically possible (Itoh' 1970)
- Characteristic M vs R relation (dashed curve)
- Smaller radius



*Letters to the Editor* 291

Prog. Theor. Phys. Vol. 44 (1970), No. 1

**Hydrostatic Equilibrium  
of Hypothetical Quark Stars**

Naoki ITOH  
*Department of Physics, Kyoto University  
Kyoto*

March 24, 1970

Ambartsumyan and Saakyan<sup>1)</sup> initiated the study of the degenerate superdense gas of elementary particles taking into account various hyperons. Afterwards many authors<sup>2)</sup> have investigated hyperon stars further either by adding newly discovered elementary particles or by assuming some interactions between these elementary particles. A primitive and straightforward question then arises: What state occurs at the density higher than hyperon stars? No one can answer this question now, as our knowledge of strong interaction physics is very incomplete.

If a baryon consists of some fundamental particles, it may be possible that unbound fundamental particles will exist in the interior of superdense stars. Ivanenko and

with the mass  $m$ . One quantum state can be occupied by less than or equal to  $q$  para-fermions. Putting  $q=1$ , we have the usual fermion case. The ratio of the limiting momentum  $p_0$  to  $mc$ ,  $x=p_0/mc$ , is related to the number density of para-fermions,  $n$ , by

$$n = q \frac{8\pi m^3 c^3}{3h^3} x^3. \quad (1)$$

The pressure is given by

$$P = q \frac{\pi m^4 c^5}{3h^3} f(x) \quad (2)$$

with the function

$$f(x) = x(2x^2 - 3)(x^2 + 1)^{1/2} + 3 \sinh^{-1} x. \quad (3)$$

The internal energy of the gas  $U_{\text{kin}}$  is given by

$$U_{\text{kin}} = q \frac{\pi m^4 c^5}{3h^3} Vg(x), \quad (4)$$

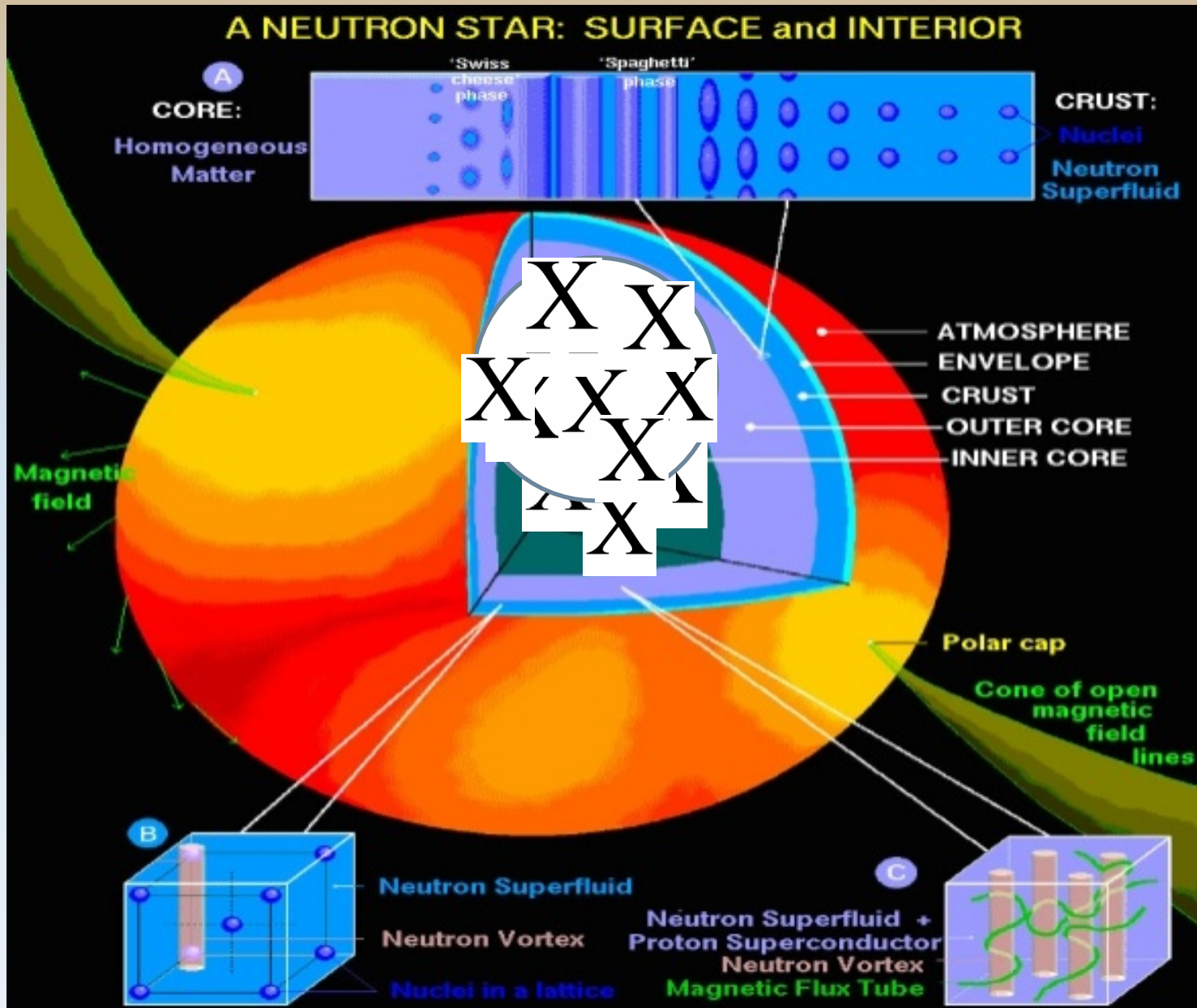
where

$$g(x) = 8x^3 \{(x^2 + 1)^{1/2} - 1\} - f(x). \quad (5)$$

We shall consider a quark star consisting of an equal amount of  $u$ ,  $d$  and  $s$ -quarks for simplicity. The condition of the charge

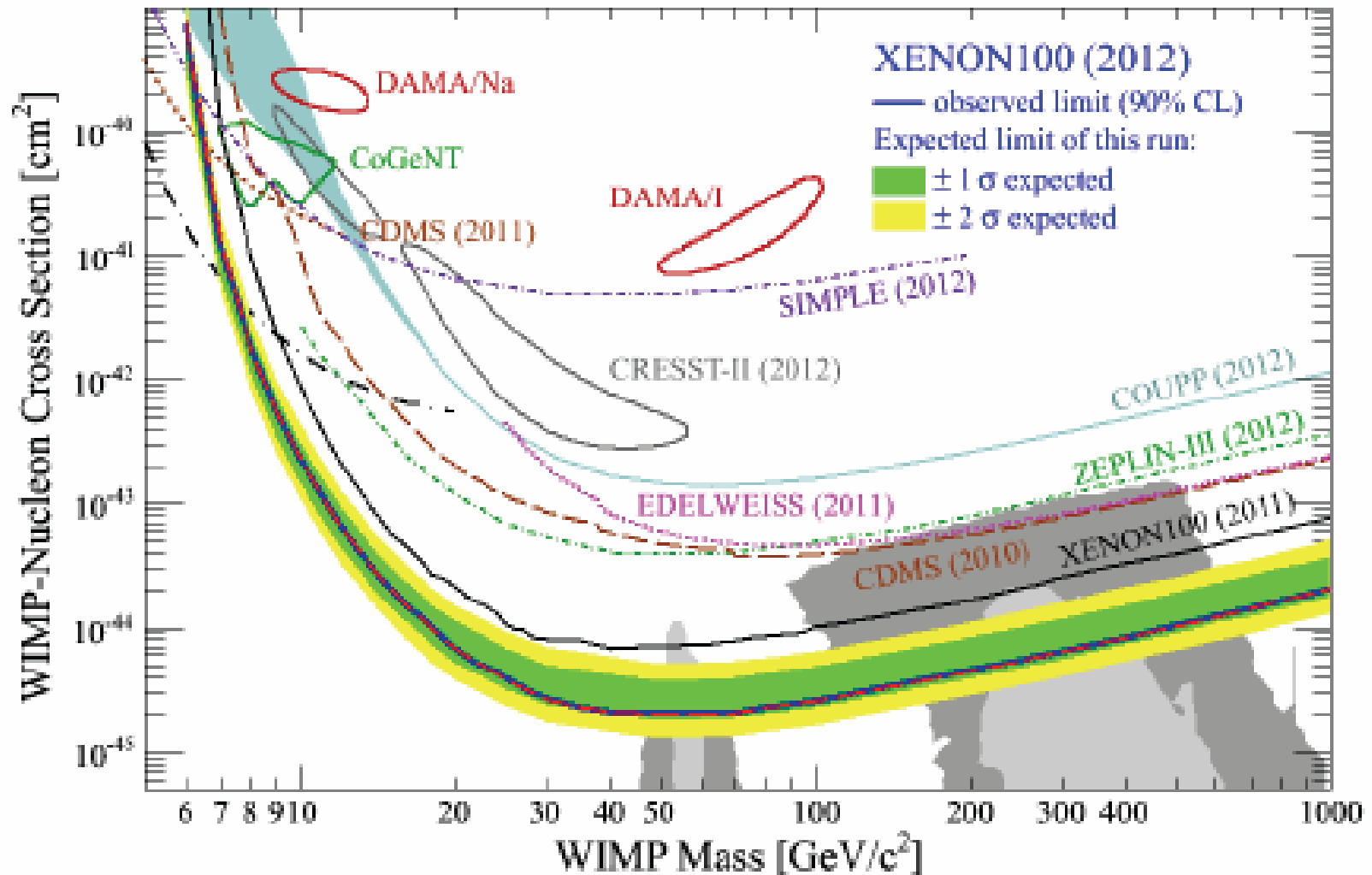
Fig 1-Previous proposal  
Theia

# UCO and Dark Matter sections are related..





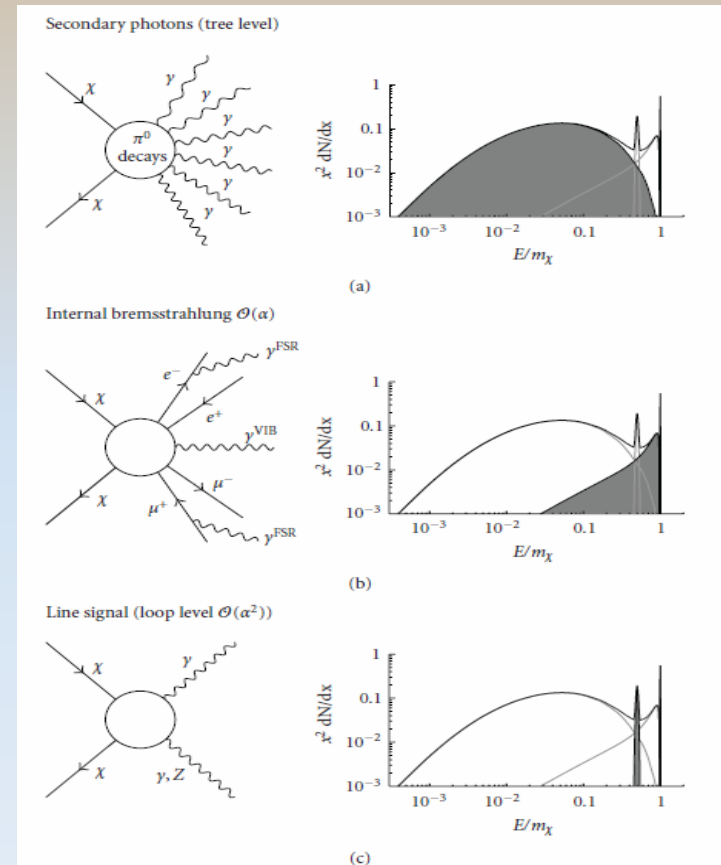
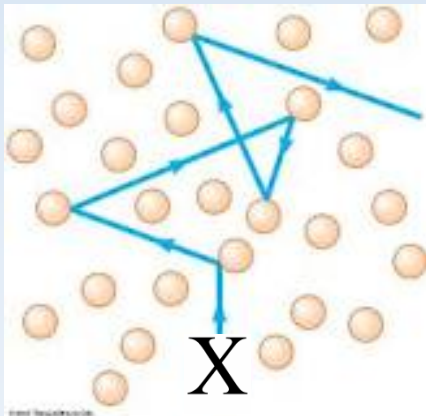
# Dark matter direct search



# NS opaque to Dark Matter

The **efficiency** of NS to capture DM is **much larger** than for the sun since:

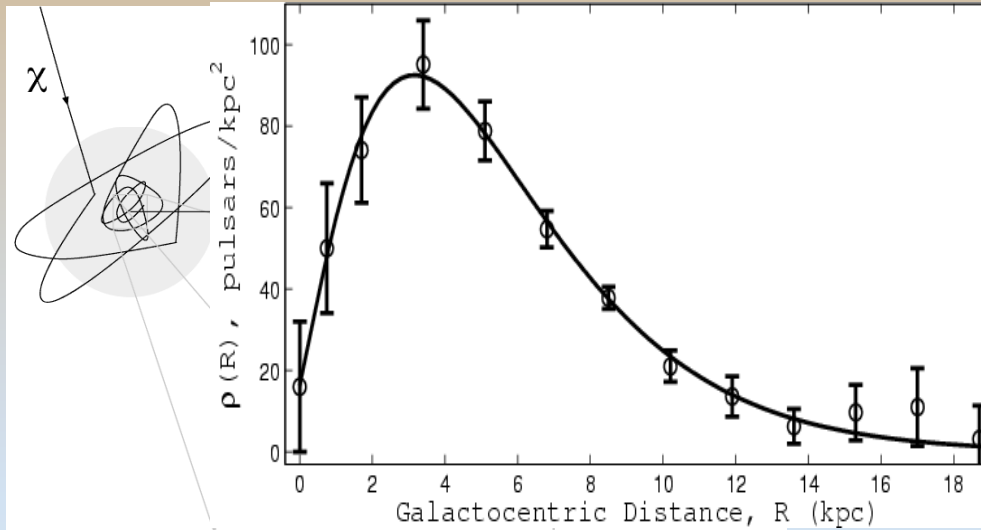
Magnitude	Sun	Neutron star
Central mass density [g/cc]	$10^2$	$10^{14}$
Mean free path [ $1/\sigma n$ ] cm	$10^{14}$	100
Capture rate [ $s^{-1}$ ]	$10^{23}$	$10^{25}$



•DM can be accreted from galactic profile by many massive astrophysical objects [Goldman, Nussinov, Press, Spergel, Kouvaris, Lavallaz, Fairbairn, Silk, Stone, Bramante, Zurek, Freese, Perez-Garcia..]

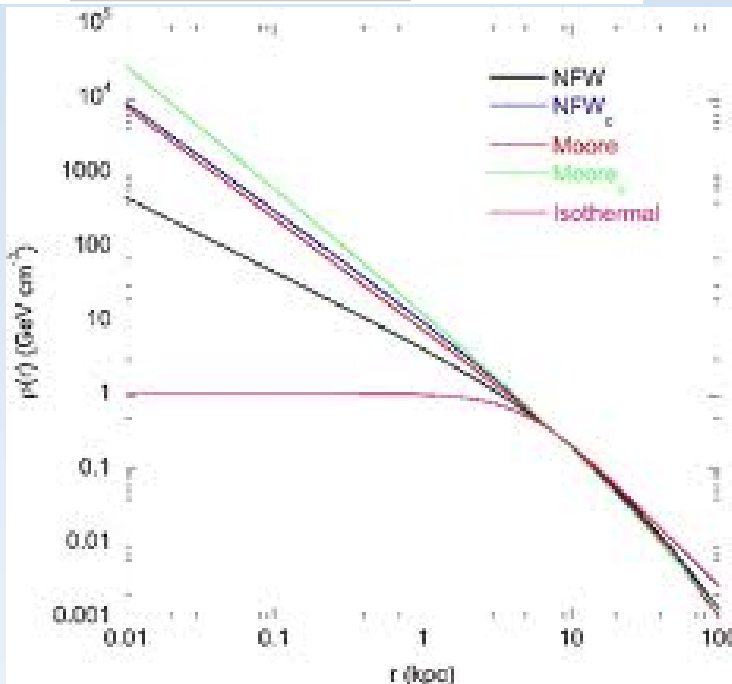


# Engine model: NS accretor of DM9



- Pulsar profile peaks at  $\sim 3$  Kpc
- DM particle could be accreted from the galactic halo
- A number of DM profiles can provide a density of local solar ( $\sim 8$  Kpc)

$$\rho_{DM,0} \approx 0.3 \text{ GeV}/c^2$$



$$F = \frac{3.042 \cdot 10^{25}}{m_\chi (\text{GeV})} \frac{\rho_{DM}}{\rho_{DM,0}} (s^{-1})$$

# Theia proposal: amplitudes

A system at distance  $D$  made of two objects of 1 Solar mass each, one being a main sequence star and the other a compact object, the astrometric amplitude is:

$$A = 150 \mu\text{as} \left( \frac{P}{1 \text{ yr}} \right)^{2/3} \left( \frac{D}{3 \text{ Kpc}} \right)^{-1} \quad (1)$$

The HMXBs usually have longer orbital periods,  $P$ :  
Theia previous proposal  
days to hundreds of days (about 1 yr).

The LMXBs are more compact,  $P$  range from tens of minutes to several hours. Y. Chou, arXiv:1408.6638v2., Q. Liu A&A 469, 807(2007)  
[XMM, Chandra, INTEGRAL data]

# Theia proposal: accuracy Mass

A Sun-like star at 3 Kpc has an apparent magnitude  $R = 17$ .

The uncertainty of a Theia observation, 10 visits, 4 h each ( $\sim 0.2\%$  of the total mission time), for a  $R = 17$  object is  $5 \mu\text{as}$  ( $1 \sigma$ )

signal to noise  $S/N = 30$  measurement of the astrometric signal

*3% accuracy in measurement of the mass of a compact object located 3 kpc with a 1 yr orbital period.*

Minimum Mass for NS:  $0.2 M_{\text{sun}}$   $\rightarrow$  no min for QS

# Theia proposal: accuracy Radius

At the distance of 3 kpc, the parallax of objects is 300  $\mu$ as.

For a  $R = 17$  object, the distance will be measured with a 1.5% accuracy.

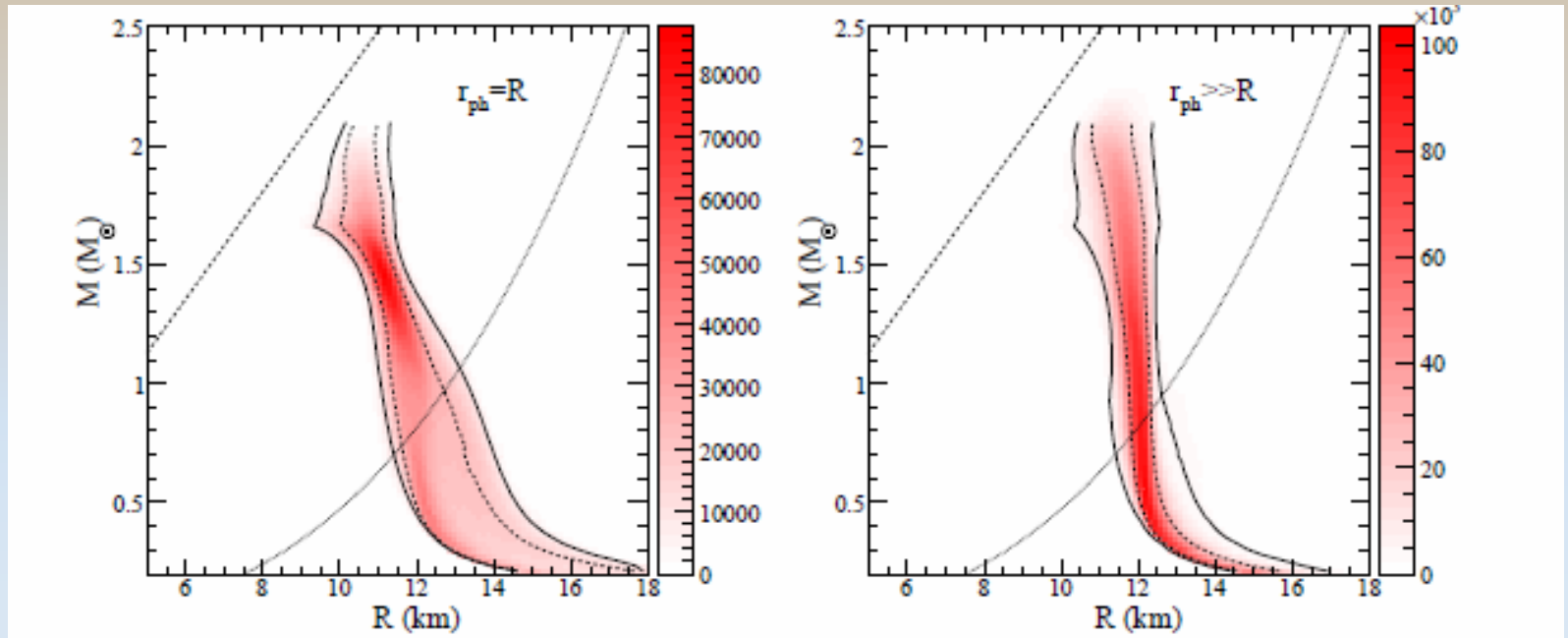
Accurate knowledge of the distance, combined with X-ray measurements of the effective temperature,  $T_{eff}^{\infty}$

$T_{eff}^{\infty}$  is gravitationally red-shifted.

Derived bolometric flux,  $F_{bol}$ , provides a measurement of the radius,

$$R^{\infty} = D \left( \frac{F_{bol}^{\infty}}{\sigma_{SB} T_{eff}^4} \right)^{1/2}$$

# M vs R and EOS



They use three Type-I X-ray bursters with photospheric radius expansion, and three transient low-mass X-ray binaries.

Steiner et al, Astrophysical Journal 722 (2010)

# M vs R and bursting/quiescent LMXRBs

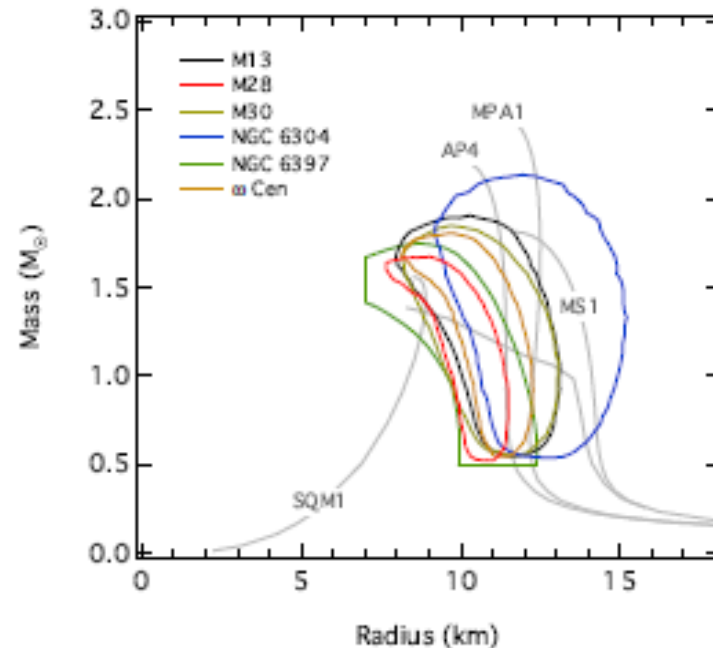
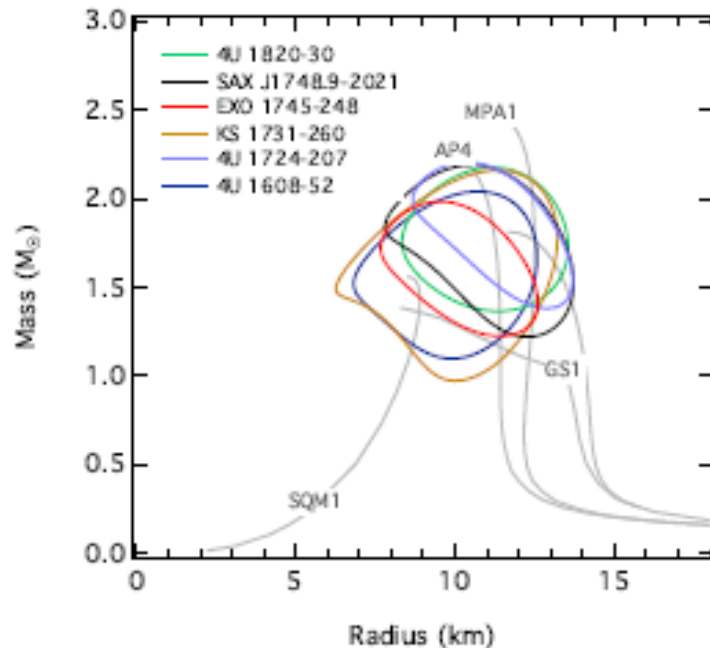


TABLE 2  
PROPERTIES OF QUIESCENT LMXBs

Source	$N_H^a$ ( $10^{22} \text{ cm}^{-2}$ )	$kT_{\text{eff}}^b$ (eV)	P.L. Norm. <sup>b</sup> ( $10^{-7} \text{ keV}^{-1} \text{ s}^{-1} \text{ cm}^{-2}$ )	Distance <sup>c</sup> (kpc)
M13	$0.01^{+0.05}_{-0.01}$	$82^{+26}_{-10}$	$3.9^{+3.8}_{-3.0p}$	$7.1 \pm 0.4^1$
M28	$0.25^{+0.03}_{-0.02}$	$123^{+41}_{-10}$	$8.3^{+4.8}_{-4.8p}$	$5.5 \pm 0.3^2$
M30	$0.01^{+0.03}_{-0.01}$	$98^{+31}_{-12}$	$8.6^{+8.2}_{-8.1p}$	$9.0 \pm 0.5^{3,4}$
$\omega$ Cen	$0.13^{+0.03}_{-0.03}$	$77^{+24}_{-8}$	$0.9^{+1.3}_{-0.7p}$	$4.59 \pm 0.08^{5,6}$
NGC 6304	$0.38^{+0.11}_{-0.10}$	$97^{+80}_{-11}$	$2.5^{+2.7}_{-2.0p}$	$6.22 \pm 0.26^7$
NGC 6397	$0.12^{+0.01}_{-0.01}$	$76^{+14}_{-7}$	$4.1^{+2.0}_{-2.0}$	$2.44 - 2.58^8$

<sup>a</sup> The hydrogen column density is fixed at the best-fit value when calculating confidence contours over the neutron star mass and radius, except for NGC 6397, see Heinke et al. (2014).

<sup>b</sup> Subscript p indicates that the posterior likelihood did not converge to zero within the hard limit of the model.

<sup>c</sup> References: 1. Harris et al. (1996, 2010 revision); 2. Servillat et al. (2012); 3. Carretta et al. (2000); 4. Lugger et al. (2007); 5. Watkins et al. (2013); 6. see also the discussion in Heinke et al. (2014); 7. Guillot et al. (2013) and references therein; 8. Heinke et al. (2014)

reason why the masses and radii they infer for 4U 1608–52 are very different between the bursts they consider (see their Figure 8) and their net result is marginally consistent with the independently measured distance to this source.

F. Özel et al  
(2015), arXiv:1505.05155v1

# Conclusions

- Theia can beat in accuracy in Mass and Radius determination with respect to GAIA.
- Typical binary systems with masses can provide mass accuracy to 3%, orbital period determination.
- Radius accuracy to 1.5%, bursting sources not necessary...but thermal emitters. X-ray missions.
- Complementary test to EOS → already important goal itself due to the difficulty of measurements.
- Best case scenario → testing QS and Quark matter EOS by selecting “small candidates”.