

Testing early Universe physics with upcoming observations

Emanuela Dimastrogiovanni



Nordita— July 17th, 2017



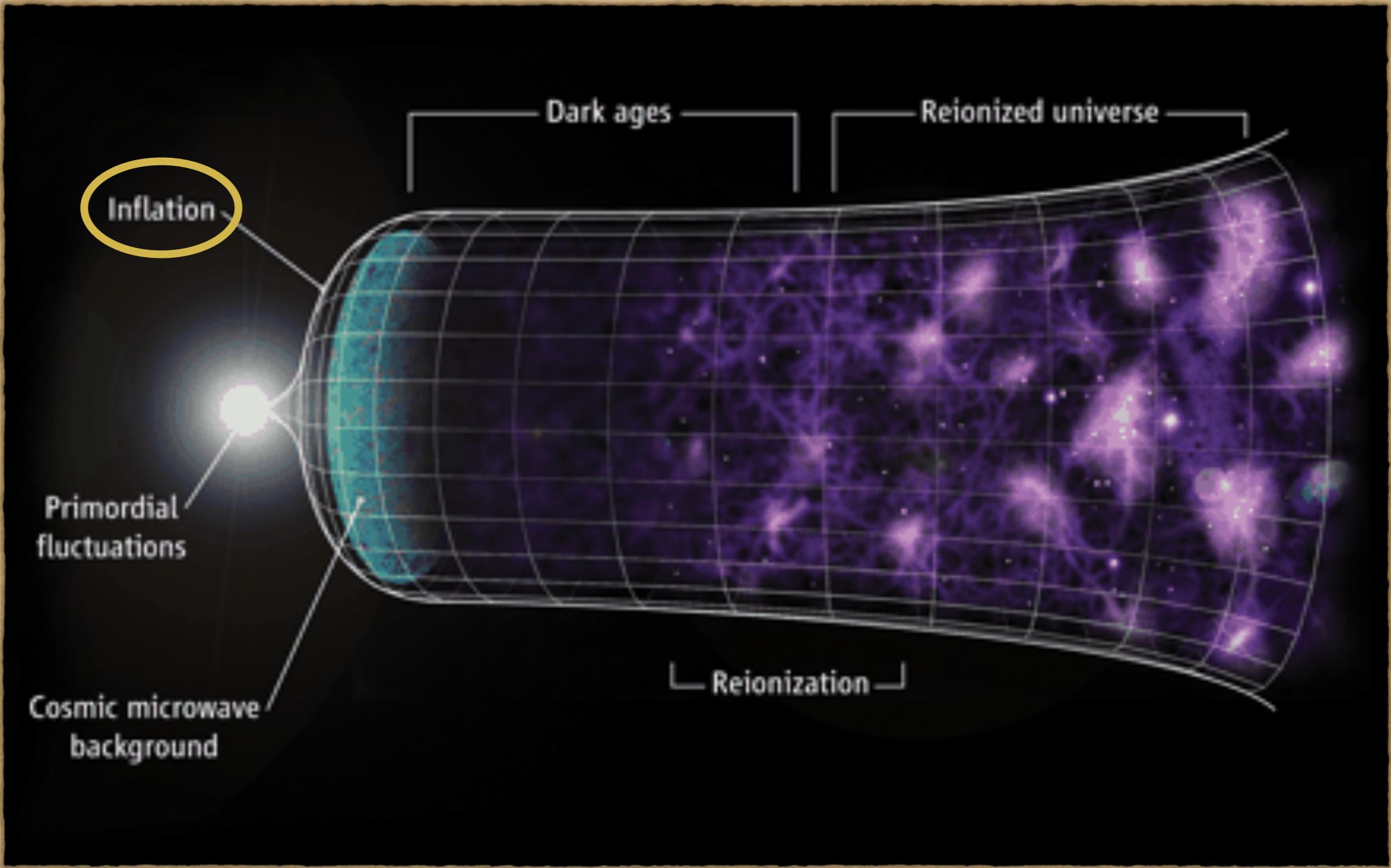
Outline:

- ED, M. Fasiello, M. Kamionkowski - 2015
- ED, M. Fasiello, D. Jeong, M. Kamionkowski - 2014

(1) Fossils from inflation
signatures of long-short mode
couplings in galaxy surveys

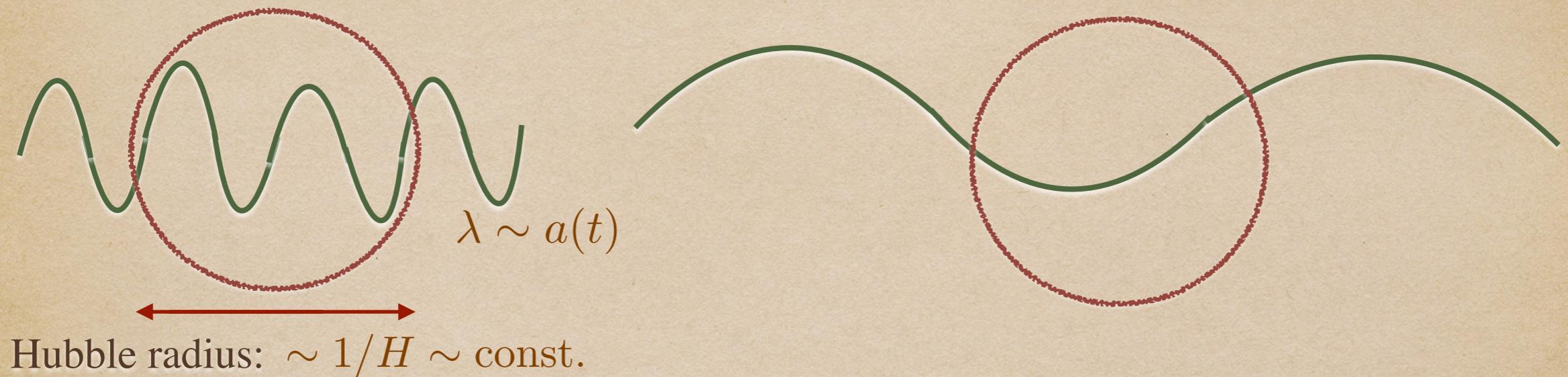
- J. Chluba, ED, M.A. Amin, M. Kamionkowski - 2017
- ED, R. Emami - 2016
- ED, L.M. Krauss, J. Chluba - 2015
- R. Emami, ED, J. Chluba, M. Kamionkowski - 2015

(2) Spectral distortions of the CMB
testing small scales primordial
perturbations



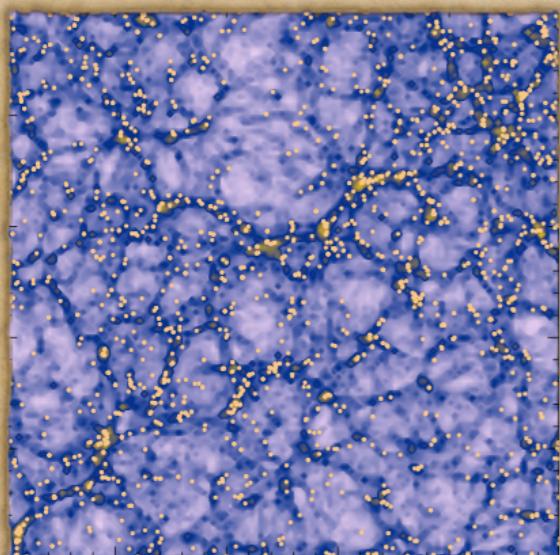
Inflation

- Era of *accelerated (nearly exponential) expansion* in the primordial Universe : $a(t) \sim e^{Ht}$



Physical scales are stretched by the expansion!

- Mechanism for the generation of cosmological fluctuations

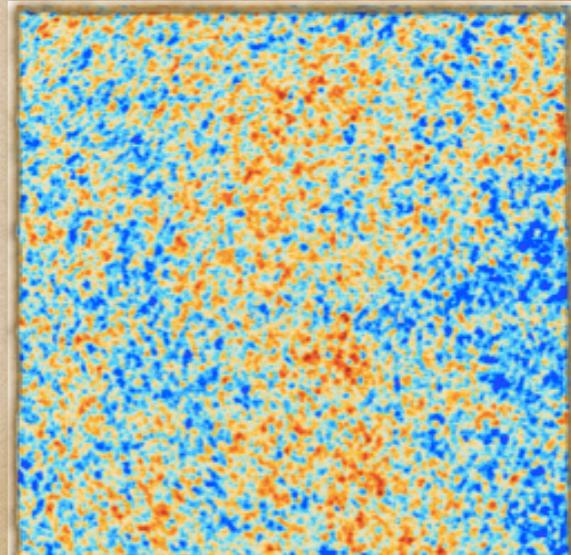


$$\phi(\vec{x}, t) = \varphi(t) + \delta\phi(\vec{x}, t)$$

Inflaton quantum fluctuations



CMB perturbations
and LSS of the Universe



Predictions: power spectra (SFSR)

Scalar fluctuations : $\delta\phi \rightarrow (\delta T/\bar{T} - \delta\rho/\bar{\rho})$

$$P_S \sim \frac{1}{\epsilon} \left(\frac{H}{M_P} \right)^2 k^{n_s-1}$$

Tensor fluctuations / primordial gravity waves :

$$P_T \sim \left(\frac{H}{M_P} \right)^2 k^{n_T}$$

Energy-scale
of inflation !

What have we learned from observations?

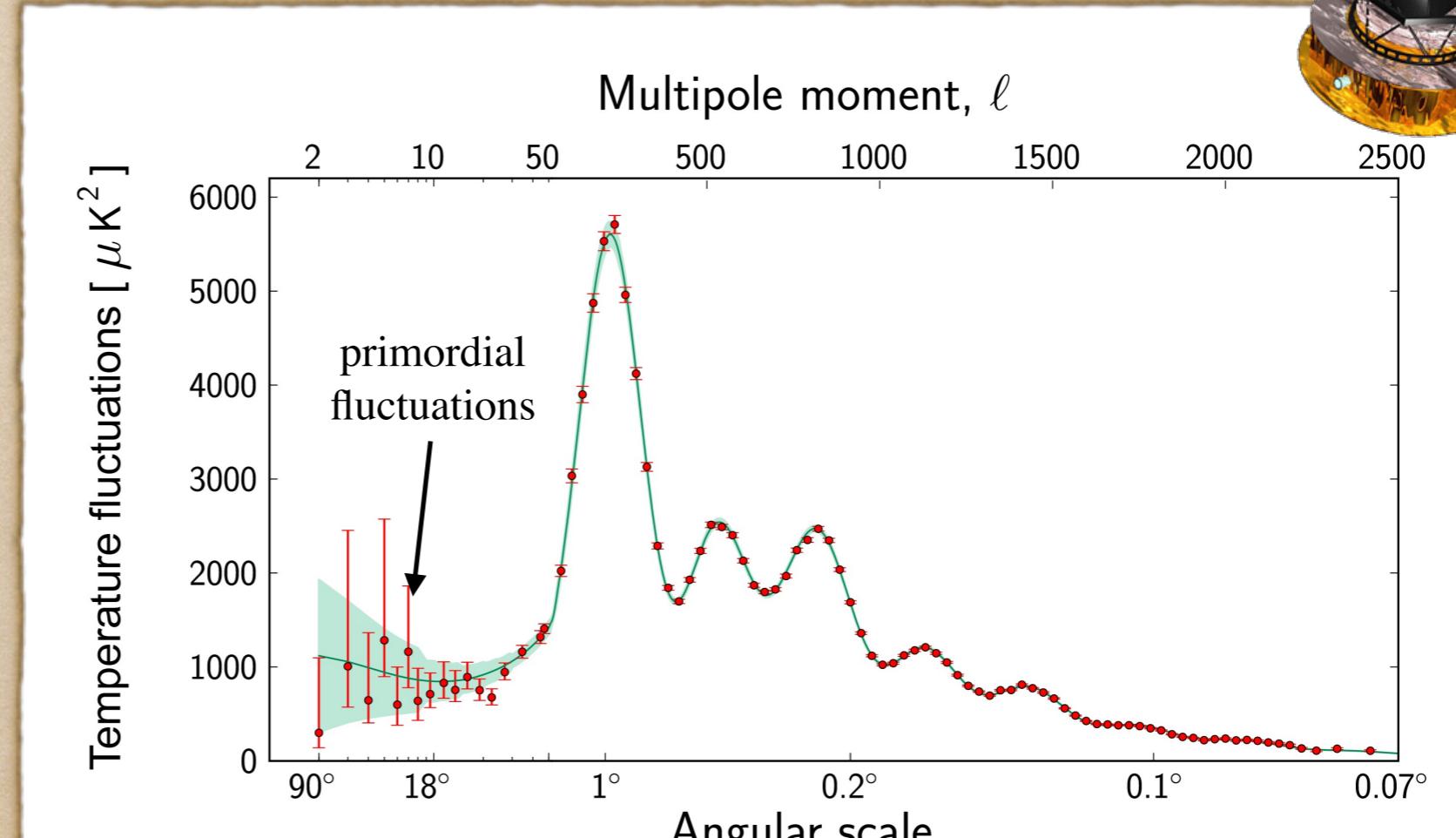
Scalar fluct. :

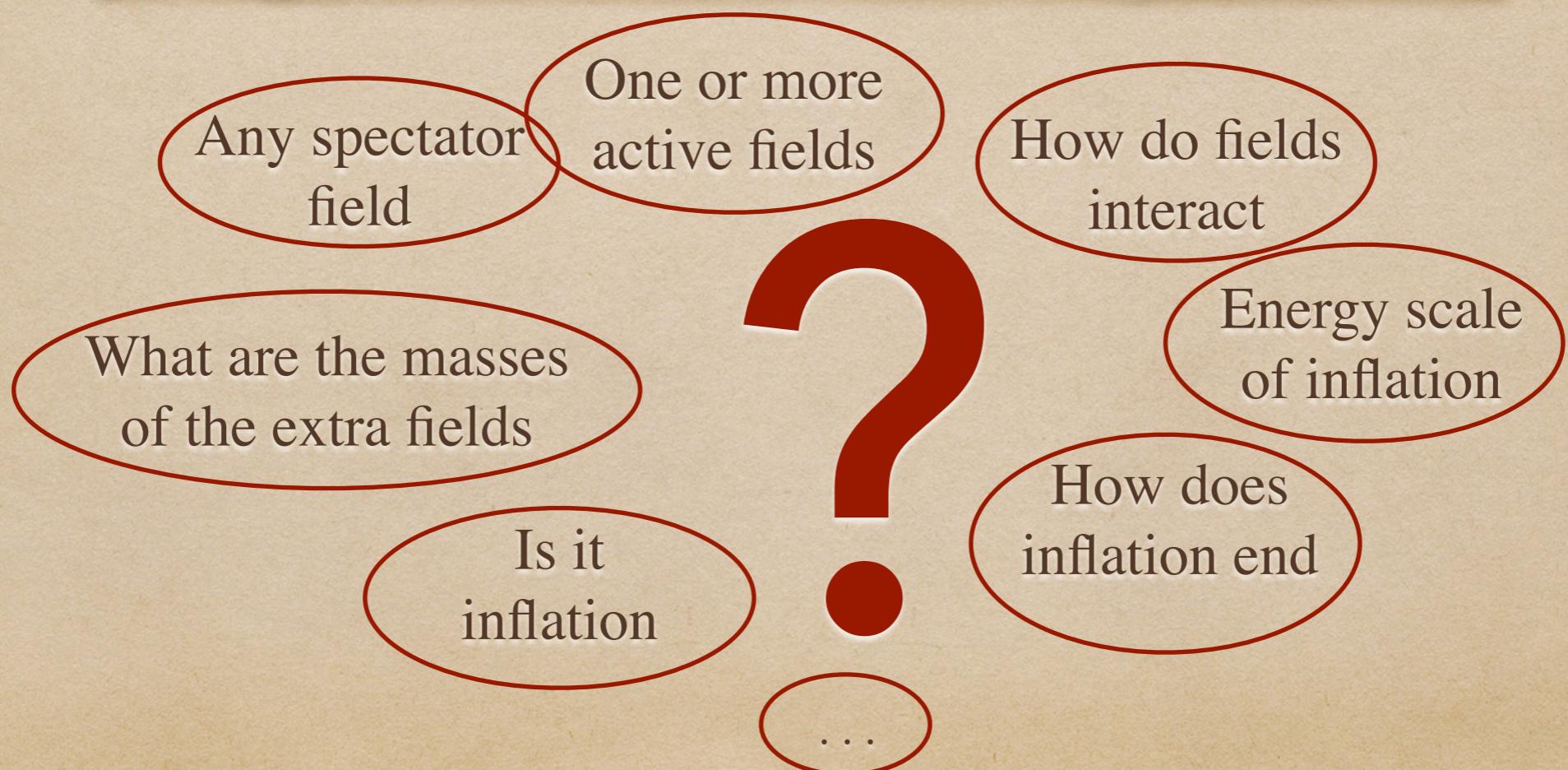
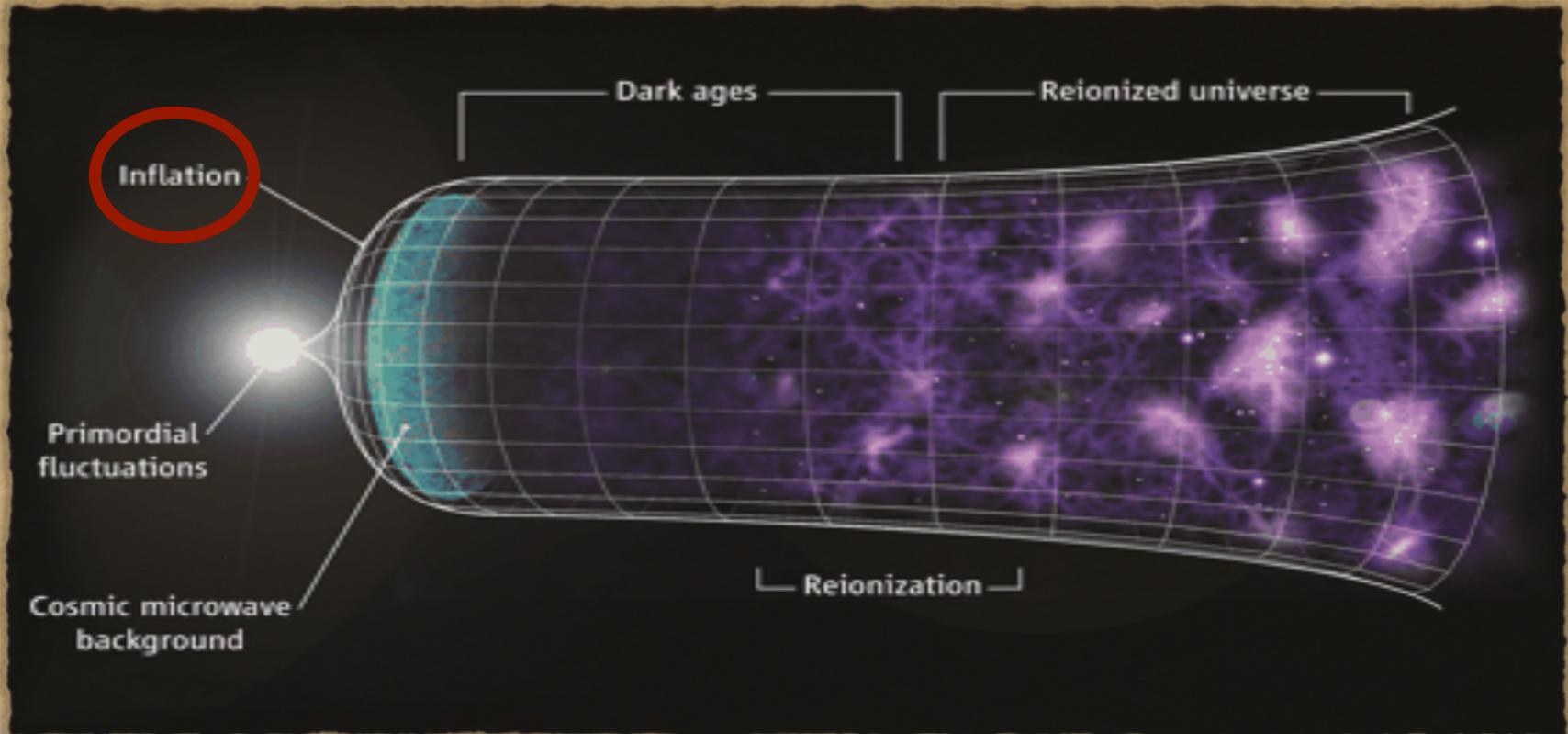
- nearly-scale invariant
 $n_s = 0.968 \pm 0.006$

Tensor fluct. :

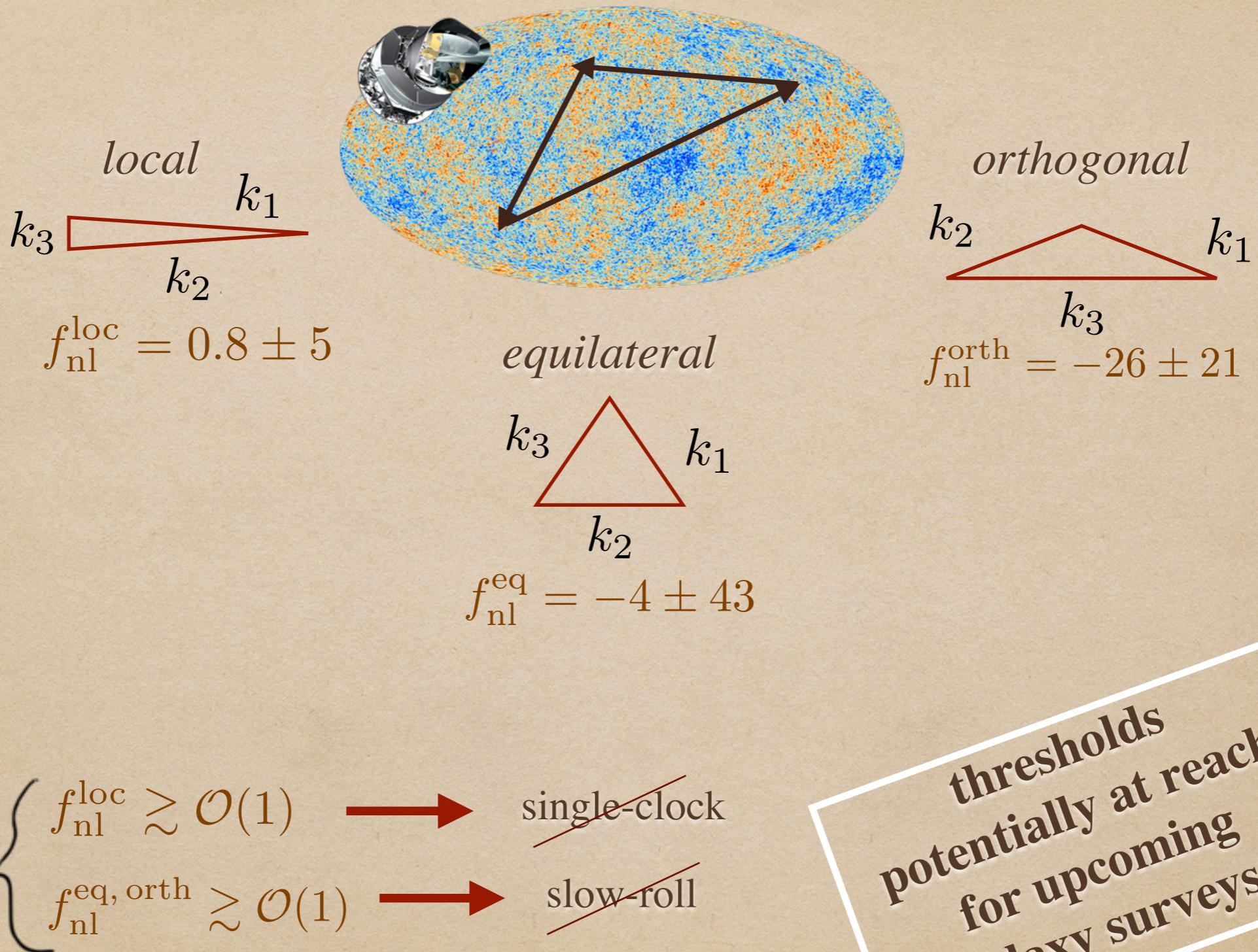
tensor-to-scalar ratio
 $r < 0.07$

BICEP2/KECK
+Planck

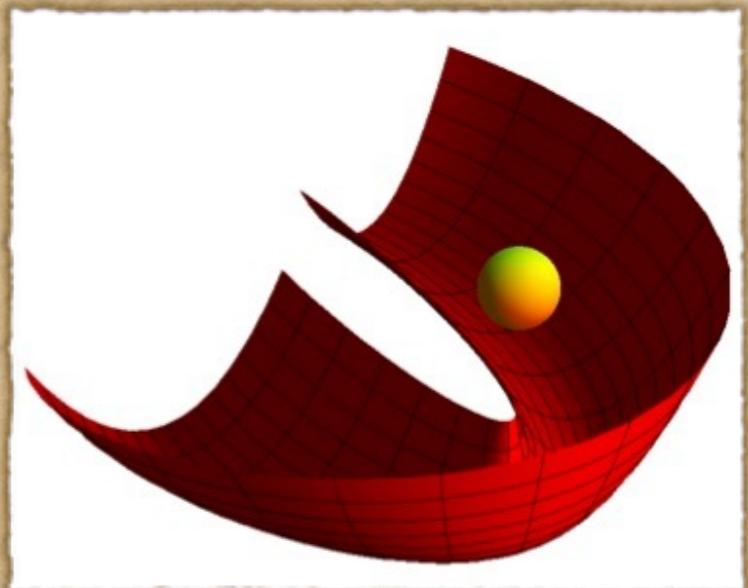




Microphysics of inflation: primordial non-Gaussianity

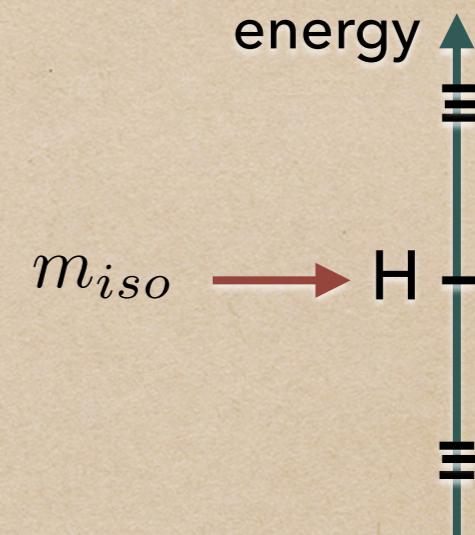
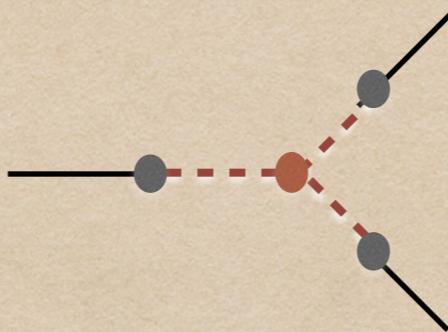


*There is more:
e.g. intermediate shapes*



Case study: Quasi-Single-Field

{ tangent: inflaton
 radial: massive isocurvaton
(also SUSY
-motivated)

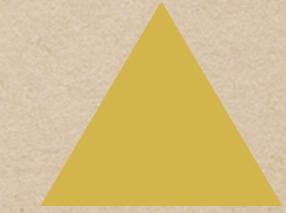


Bispectrum shape between local and equil.



$m_{iso} \rightarrow 0$

• • •



$m_{iso} \rightarrow 3H/2$

$$\mathcal{B}|_{squ.} \sim \frac{1}{k_1^3 k_3^3} \left(\frac{k_3}{k_1}\right)^{3/2 - f(m_{iso})}$$

$$f(m_{iso}) \equiv \sqrt{\frac{9}{4} - \frac{m_{iso}^2}{H^2}}$$

bispectrum shape
delivers info
on mass spectrum!

[Chen, Wang, 2010][Baumann, Green, 2011]

What happens next?

- Galaxy surveys :

photometric/spectroscopic
(DESI, LSST, Euclid,
SPHEREx, ...)
and 21cm (CHIME, ...)

{ Microphysics of inflation
(primordial non-Gaussianity,
running/features of primordial
power spectrum, ...)

Fossils from
inflation

- CMB experiments :

Spectral distortions
(PRISM, PIXIE, ...)

{ Testing post- and
pre-recombination
physics

Primordial
perturbations
on small scales

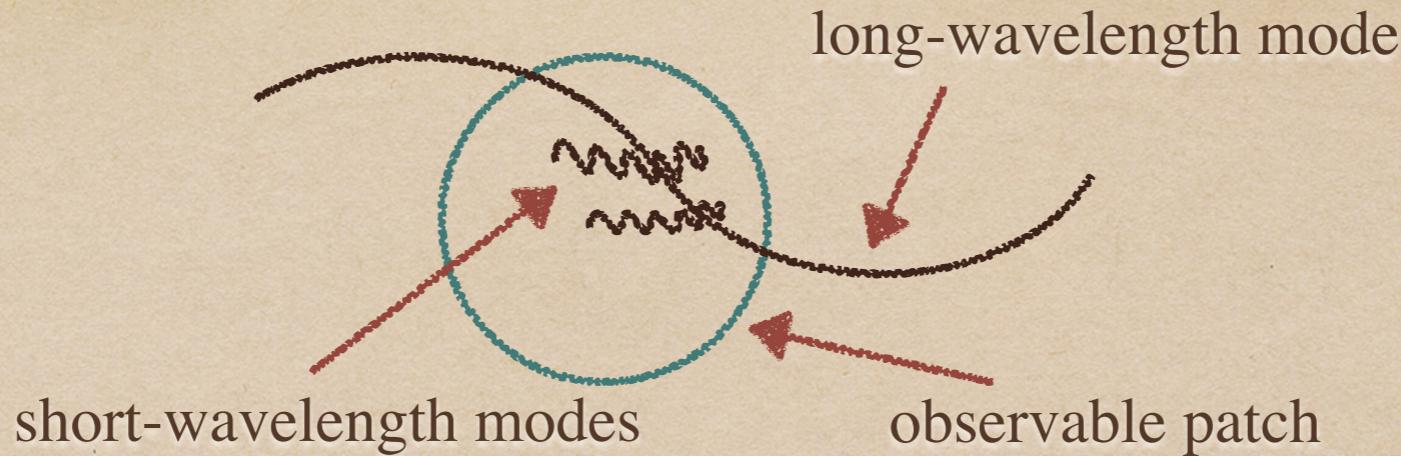
B-modes
(BICEP/KECK series,
SPTPol, SPIDER,
PIXIE, ...)

{ Origin of B-modes?
What can we learn about
the early Universe?

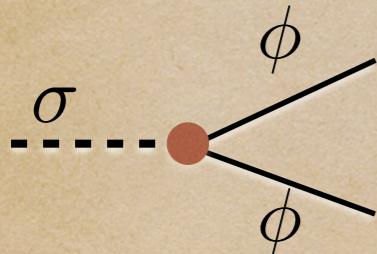
Fossils from inflation

*signatures of **multi-clock**
inflationary dynamics*

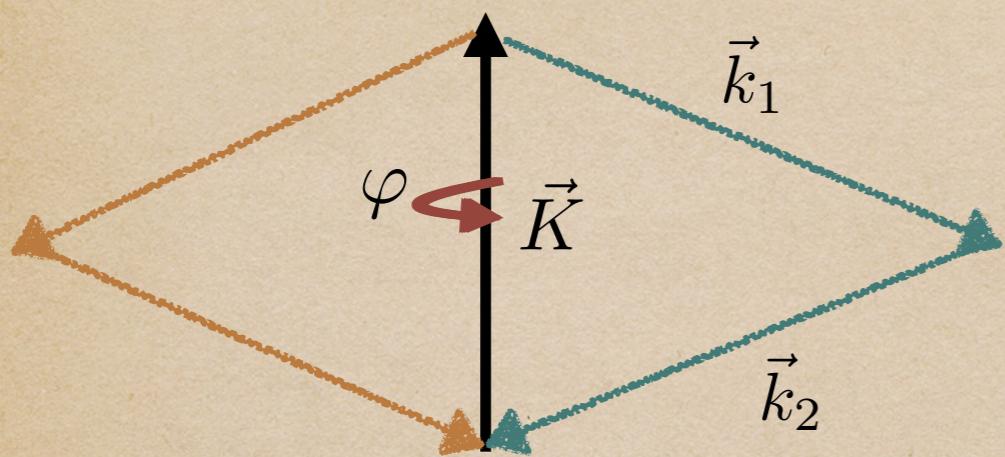
Effect of a long mode on local observables :



- Gaussian initial conditions + statistical homog./isotropy $\rightarrow \langle \delta_{\vec{k}_1} \delta_{\vec{k}_2} \rangle = \delta^{(3)}(\vec{k}_1 + \vec{k}_2) P(k_1)$ diagonal
2p corr.
- Local/squeezed non-Gaussianity $\rightarrow \vec{K}$ $\rightarrow \langle \delta_{\vec{k}_1} \delta_{\vec{k}_2} \rangle_{\vec{K}} = \delta^{(3)}(\vec{k}_1 + \vec{k}_2 + \vec{K}) \times f(\vec{k}_1, \vec{k}_2) A(K)$ off-diagonal
2p corr.!



Probing σ through off-diagonal correlation :

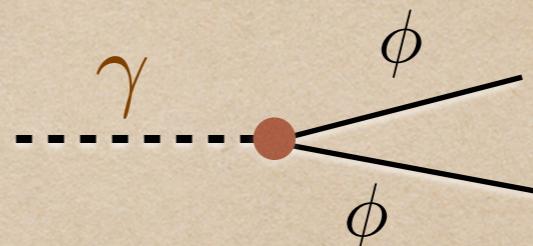


Azimuthal dependence:

Scalar	→ no φ dependence
Vector	→ $\cos \varphi, \sin \varphi$
Tensor	→ $\cos 2\varphi, \sin 2\varphi$

[Jeong, Kamiokowski, 2012]

What if $\sigma = \text{tensor mode}$ of the metric:



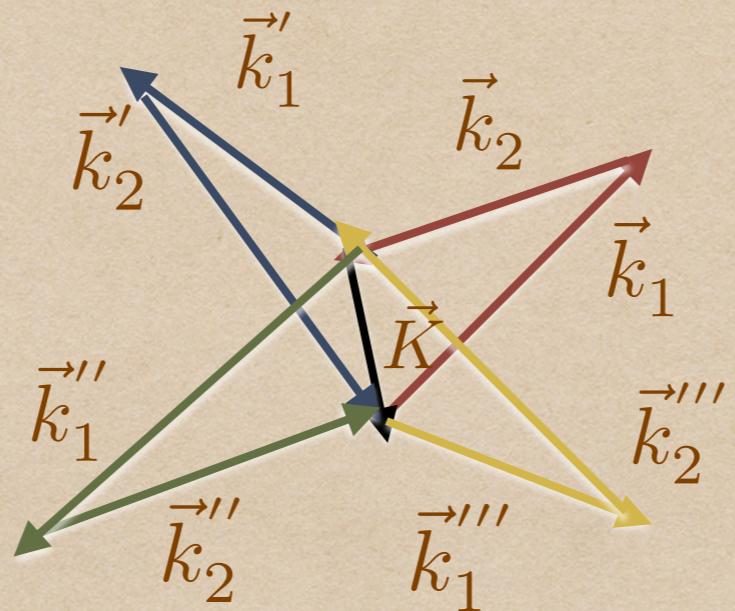
Learning about primordial tensors
through non-Gaussian effects!

[ED, Fasiello, Jeong, Kamionkowski - 2014]

[ED, Fasiello, Kamionkowski - 2015]

Estimating tensor modes amplitude :

$$\langle \delta_{\vec{k}_1} \delta_{\vec{k}_2} \rangle_{\gamma_p(\vec{K})} \sim (2\pi)^3 \delta^{(3)}(\vec{k}_1 + \vec{k}_2 + \vec{K}) \gamma_p^*(\vec{K}) \frac{B_p(\vec{k}_1, \vec{k}_2, \vec{K})}{P_\gamma^p(K)}$$



- Naive estimator:

$$\widehat{\gamma_p}(\vec{K}) = \sum_{\vec{k}_1 + \vec{k}_2 = -\vec{K}} \frac{\delta_{\vec{k}_1} \delta_{\vec{k}_2}}{B_p(\vec{k}_1, \vec{k}_2, \vec{K}) / P_\gamma^p(K)}$$

- Optimal estimator for a single mode : inverse variance weighting

$$\widehat{\gamma_p(\vec{K})} = \sigma^2 \sum_{\vec{k}} \frac{|B_p(\vec{K}, \vec{k}, \vec{K} - \vec{k})/P_\gamma^p(K)|^2}{2 V P^{tot}(k) P^{tot}(|\vec{K} - \vec{k}|)} \delta_{\vec{k}} \delta_{\vec{K} - \vec{k}}$$

$$\sigma^2 = \left[\sum_{\vec{k}} \frac{|B_p(\vec{K}, \vec{k}, \vec{K} - \vec{k})/P_\gamma(K)|^2}{2 V P^{tot}(k) P^{tot}(|\vec{K} - \vec{k}|)} \right]^{-1}$$

random variable
 variance
 $\{y_i, \sigma_i^2\}$

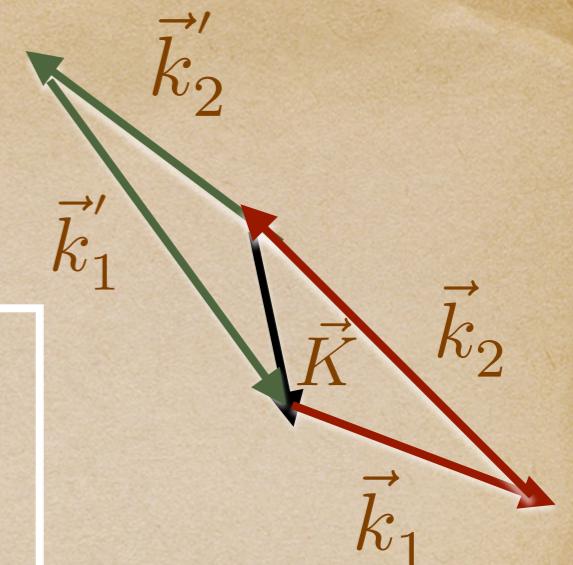
$$\hat{y} = \frac{\sum_i (y_i / \sigma_i^2)}{\sum_j (1 / \sigma_j^2)} \quad \sigma^2 = \left(\sum_i \frac{1}{\sigma_i^2} \right)^{-1}$$

- Optimal estimator for power amplitude:
stochastic GW background with $P_p(K) = A_\gamma P_\gamma^f(K)$

$$\widehat{A}_\gamma = \sigma_\gamma^2 \sum_{p, \vec{K}} \frac{(P_\gamma^f(K))^2}{(P_p^n(K))^2} \left(\frac{\widehat{|\gamma_p(\vec{K})|^2}}{V} - P_p^n(K) \right)$$

optimal sum over different K-modes

$$\begin{cases} \sigma_\gamma^{-2} = \sum_{p, \vec{K}} \frac{(P_\gamma^f(K))^2}{2(P_p^n(K))^2} \\ P_p^n \equiv \left[\sum_{\vec{k}} \frac{|B_p(\vec{K}, \vec{k}, \vec{K} - \vec{k})/P_\gamma(K)|^2}{2V P^{tot}(k) P^{tot}(|\vec{K} - \vec{k}|)} \right]^{-1} \end{cases}$$



$$\mathcal{B}_{\gamma\zeta\zeta} \simeq \beta P_\gamma P_\zeta$$

$$\sigma_\gamma \propto \frac{1}{\beta^2} \left(\frac{k_{max}}{k_{min}} \right)^{-3}$$

Which classes of models predict these signatures at observable levels?

~~single-clock ccs~~

- (n+1)-point function fixed in terms of n-point functions
(soft limit for one of the modes)
- Apply if super-horizon modes freeze + standard initial conditions
(rescaling of background for short modes)
- Derived from symmetries of the action (invariance under space diffs)
[Maldacena 2003, Creminelli-Zaldarriaga 2004, Goldberger et al 2013, ...]

- *Isocurvature* modes (multi-field)
- *Non-Bunch Davies* initial states
[Holman - Tolley 2007, Brahma-Nelson-Shandera 2013, ...]
- *Broken space diffs* (e.g. solid inflation/space-dependent background)
[Endlich et al., 2013, Bartolo et al, 2015]

Fossil signatures at reach
for upcoming observations
(e.g. Euclid or 21cm)!

[ED, Fasiello, Jeong, Kamionkowski - 2014]
[ED, Fasiello, Kamionkowski - 2015]

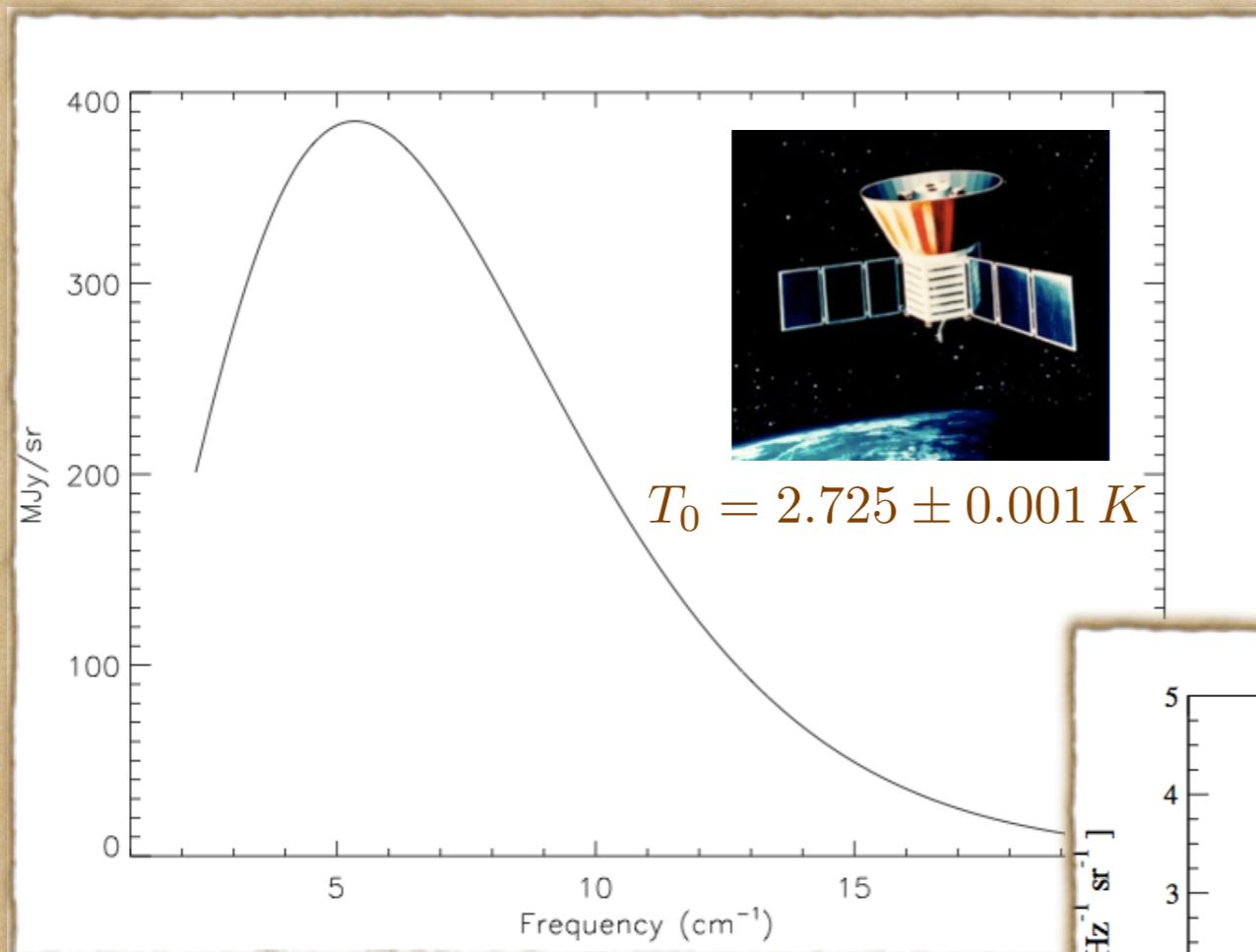
Summary

- *tss* correlation in the squeezed limit affects density power spectrum (*off-diagonal/quadrupolar anisotropy*)
- signatures observable if arising from models that evade ccs: another *test for single-clock inflation* besides the scalar (squeezed) bispectrum
- Primordial *GWs amplitude* may be estimated from off-diagonal correlation!

Spectral distortions of the CMB

*A window into primordial
perturbations on **small scales***

Cosmic Microwave spectrum from COBE/FIRAS



Deviations allowed from FIRAS:

$$|\mu| \lesssim 9 \times 10^{-5}$$

$$|y| \lesssim 1.5 \times 10^{-5}$$

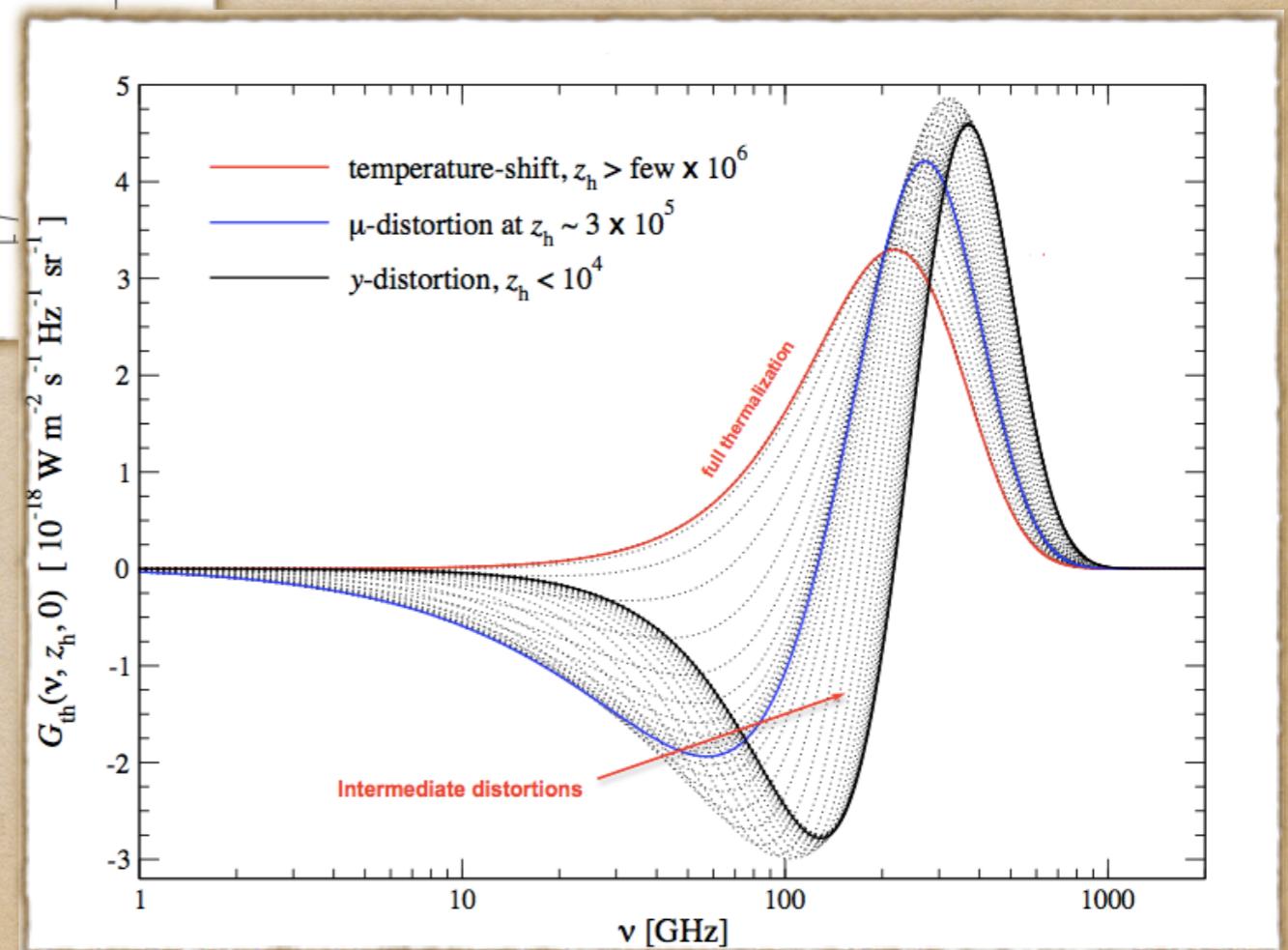
[Fixsen et al., 1996]

Energy/photon injections
into CMB at $z \lesssim 2 \times 10^6$

Distortion of the
blackbody spectrum

$$\Delta I_\nu \approx \int G_{\text{th}}(\nu, z') \frac{d(Q/\rho_\gamma)}{dz'} dz'$$

energy-
injection
rate



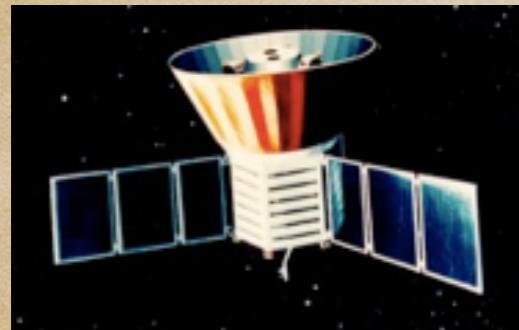
[Chluba, 2013]

Why spectral distortions matter?

CMB anisotropies:
(~20 yrs)

Spectral distortions:

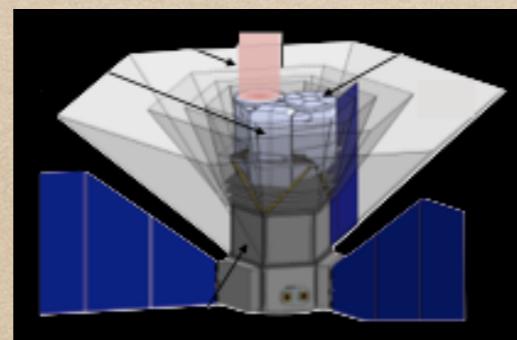
COBE/FIRAS



$$|\mu| \lesssim 9 \times 10^{-5}$$
$$|y| \lesssim 1.5 \times 10^{-5}$$

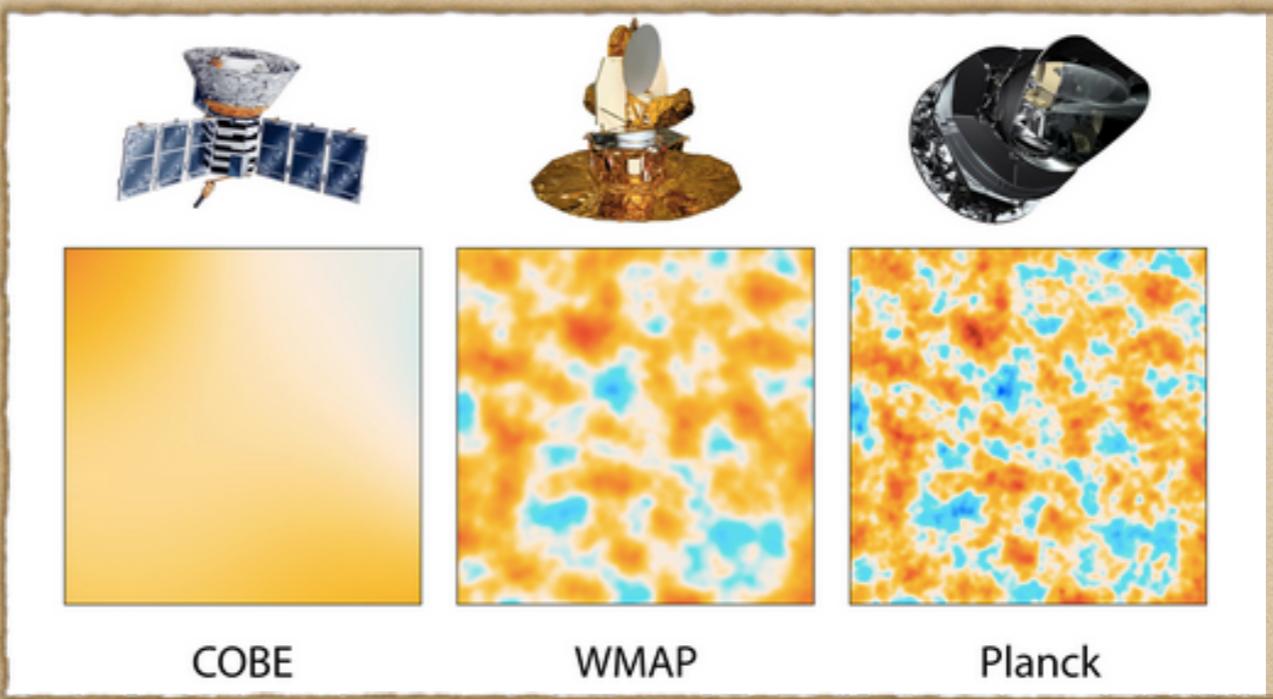
[Fixsen et al., 1996]

PIXIE



$$|\mu| \lesssim 2 \times 10^{-8}$$
$$|y| \lesssim 4 \times 10^{-9}$$

[Kogut et al., 2011]

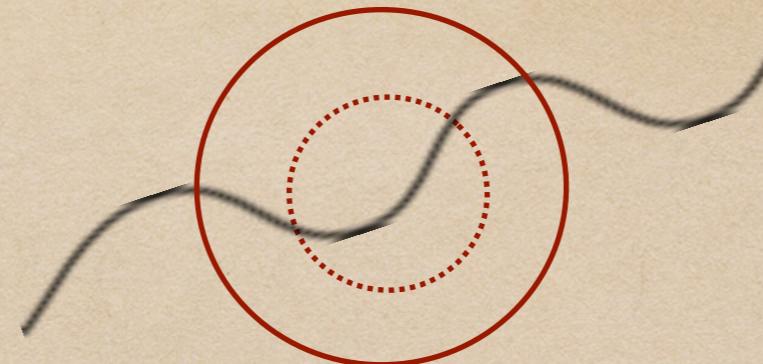


Mechanisms producing SDs include:

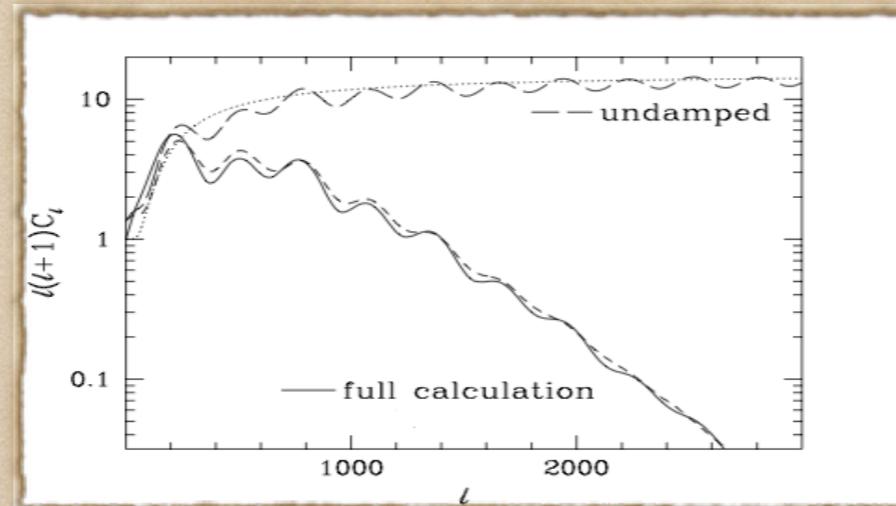
- Cosmological recombination
- Reionization, structure formation
- Dissipation of magnetic fields
- Decay/annihilation of particles
- **Dissipation of primordial fluctuations**
- ...

SDs from diffusion damping

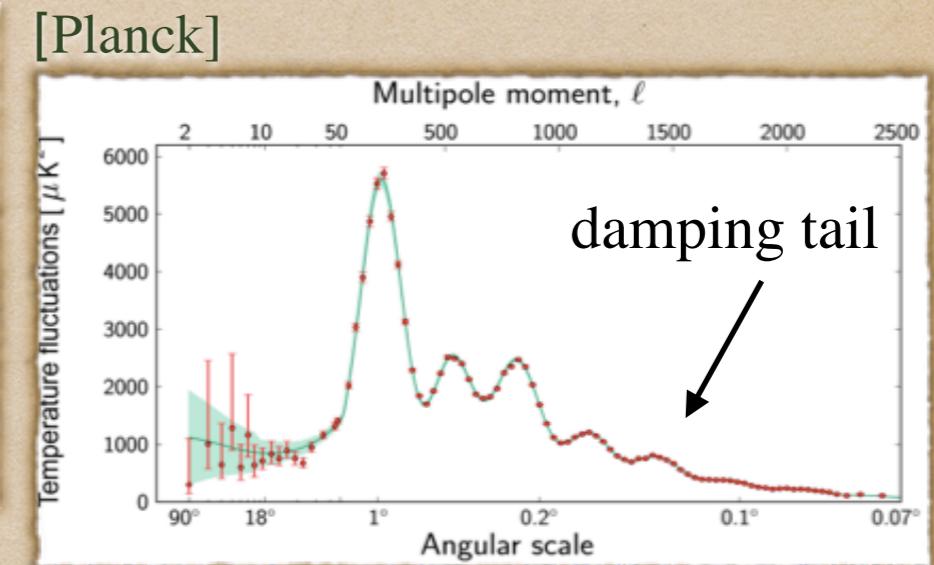
Primordial fluctuations
re-enter Hubble radius
after inflation



Isotropization
by photon
diffusion



[Hu&White, 1996]



Effective **energy**
release into CMB
[Hu et al., 1994]
[Chluba et al., 2012]

$$\frac{d(Q/\rho_\gamma)}{dz} \sim \int dk k^4 \mathcal{P}_\zeta(k) e^{-[2k^2/k_D^2(z)]} \approx \left(4 \times 10^{-6} (1+z)^{3/2} \text{Mpc}^{-1}\right)^2 \text{(damping scale)}$$

Probing scales k : $\begin{cases} [50, 10^4] \text{ Mpc}^{-1} & (\text{with } \mu \text{ distortion}) \\ [1, 50] \text{ Mpc}^{-1} & (\text{with } \gamma \text{ distortion}) \end{cases}$

probing ~ 10 extra
inflationary e-folds!

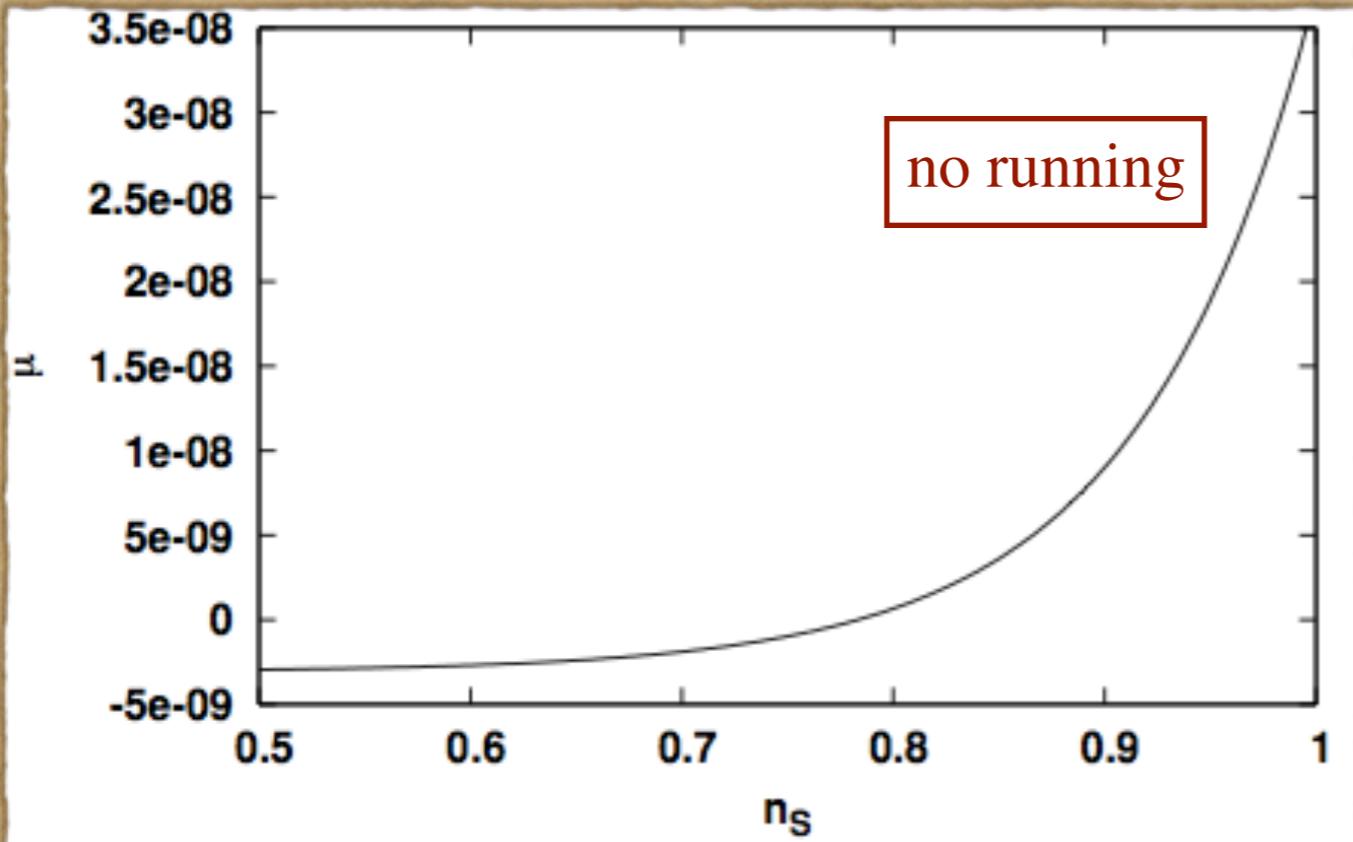
Power spectrum constraints: SFSR

$$\mathcal{P}_\zeta(k) = A_\zeta \left(\frac{k}{k_*} \right)^{n_s - 1 + \frac{1}{2} \{dn_s/d\ln k\} \ln(k/k_*)}$$

$$\begin{cases} n_s = 0.968 \pm 0.006 \\ dn_s/d\ln k = -0.006 \pm 0.007 \end{cases}$$

Planck temperature
+ polarization data
on large angular
scales (68% CL)

[Khatri, Sunyaev, Chluba, 2012]



$$\text{E.g. for } n_s = 0.96: \frac{Q_{z=2 \times 10^6}^{z=5 \times 10^4}}{\rho_\gamma} \simeq \mathcal{O}(10^{-8})$$

Sensitivity for μ distortion

$$\sigma_\mu = (1/n) \times 10^{-8}$$

E.g. : $\begin{cases} n=1 & \xrightarrow{} \text{PIXIE} \\ n=10 & \xrightarrow{} \text{PRISM} \end{cases}$

3 x PIXIE = *guaranteed discovery*

detection
(95% CL)

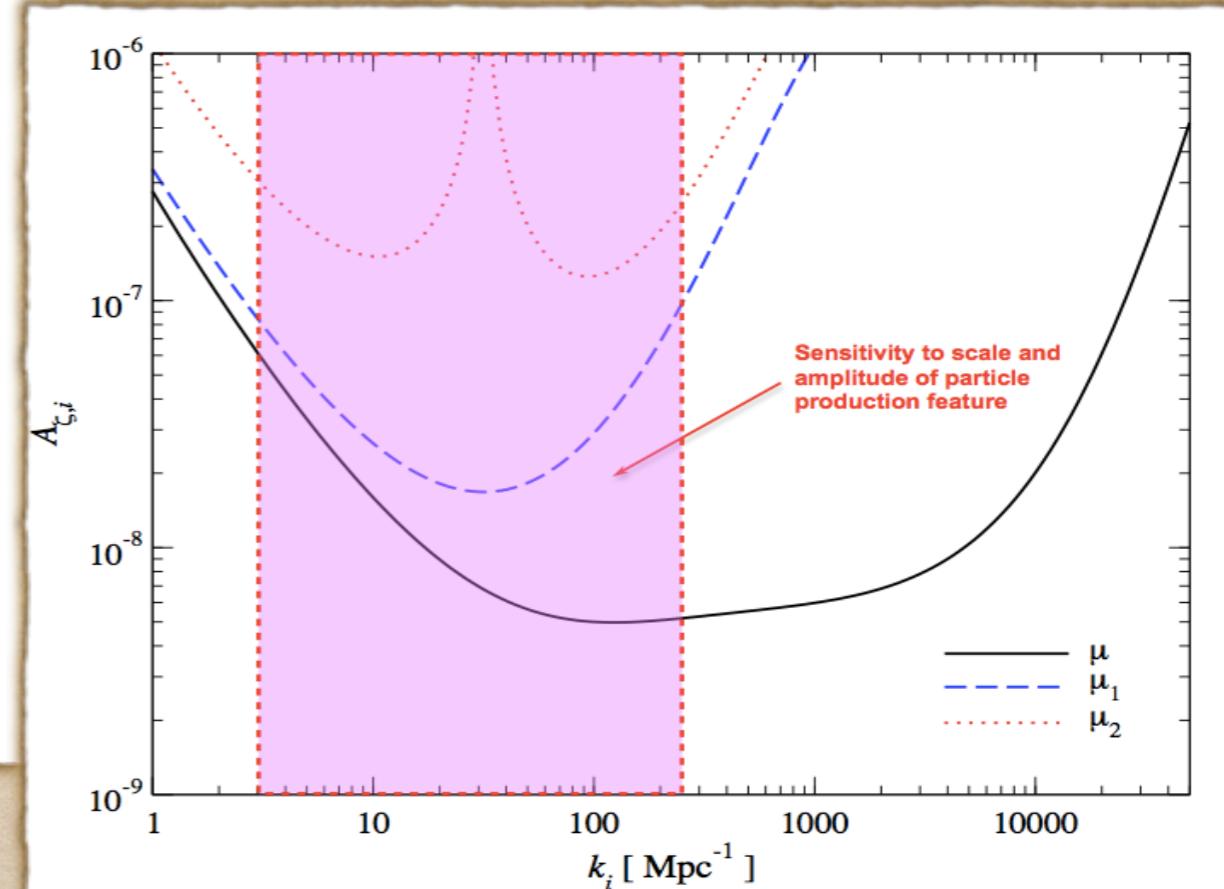
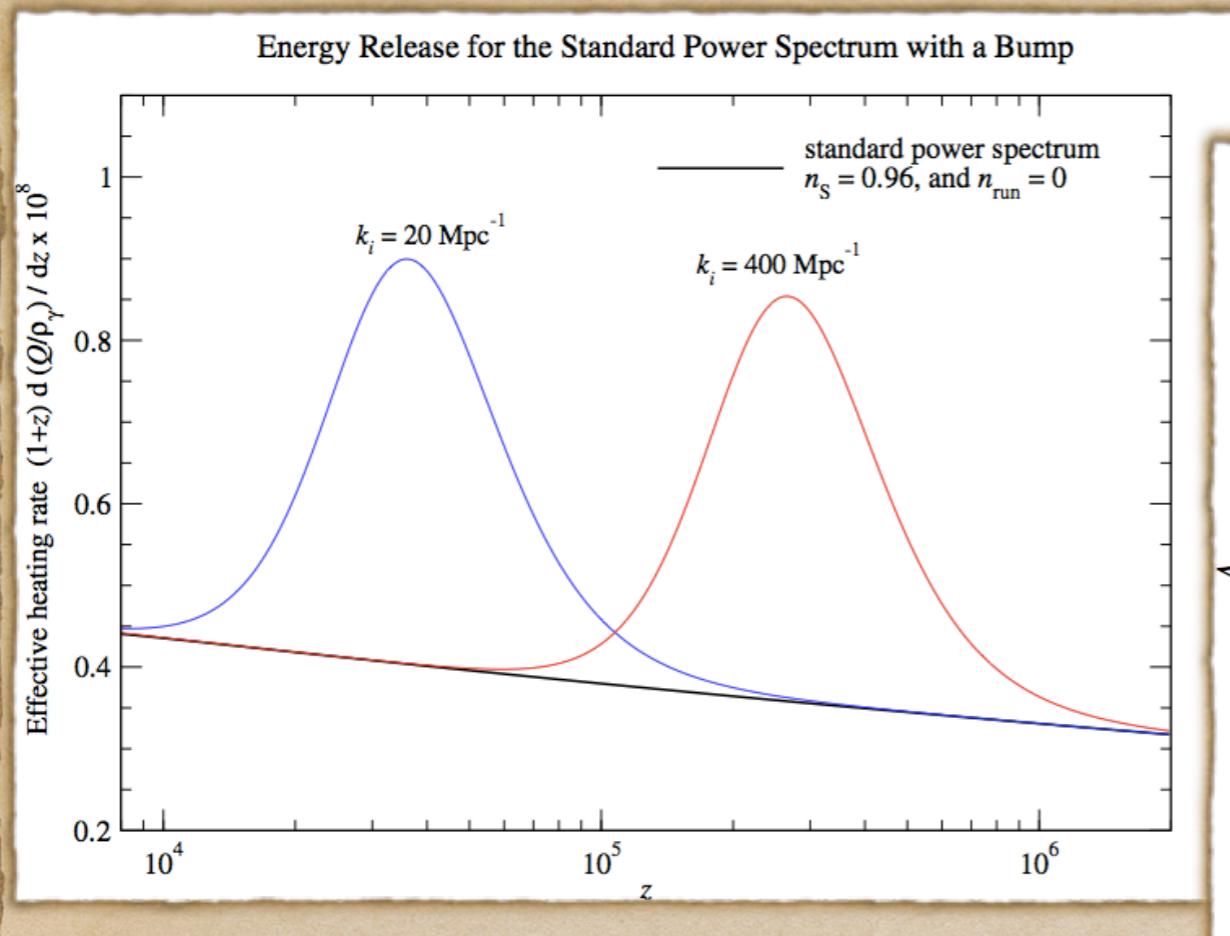
exclusion
of SFSR
(95% CL)

[Cabass, Melchiorri, Pajer, 2016]

Power spectrum constraints: ~~SESR~~

$$\mu \approx 2.2 \int \mathcal{P}_\zeta \left[e^{-\left(\frac{k \cdot \text{Mpc}}{5400}\right)^2} - e^{-\left(\frac{k \cdot \text{Mpc}}{31.6}\right)^2} \right] d \ln k$$

$$y \approx 0.4 \int \mathcal{P}_\zeta(k) e^{-\left(\frac{k \cdot \text{Mpc}}{31.6}\right)^2} d \ln k$$



- e.g. from particle production
- bump of amplitude $\mathcal{A}_{\zeta,i}$, localized around k_i
- intermediate distortions to remove degeneracies

[Chluba et al., 2015]

General note:

bounds from SDs competitive
w.r.t. those from PBHs and UCMHs!

Bispectrum constraints

- Bispectrum in the squeezed limit from distortion-temperature correlation
[Pajer - Zaldarriaga, 2012]

Example: local ansatz

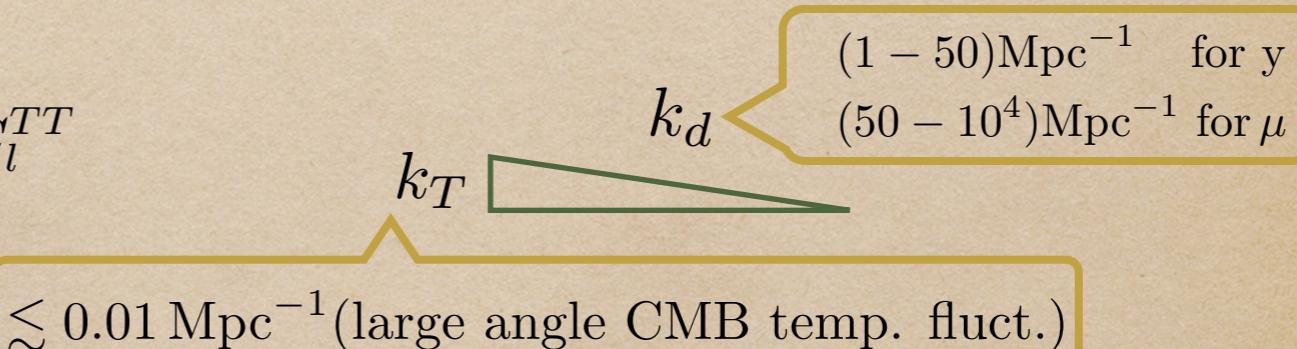
$$\mathcal{R}(\vec{x}) = r(\vec{x}) + \frac{3}{5} f_{nl} r^2(\vec{x})$$

Long-short mode decomposition: $\mathcal{R} = \mathcal{R}_L + \mathcal{R}_s$

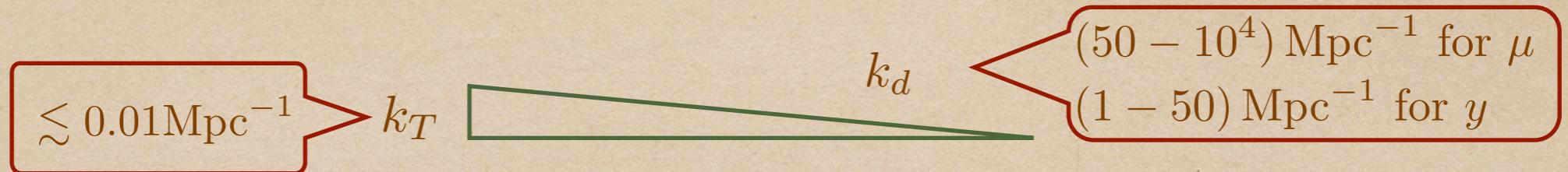


$$\mathcal{R}_s(\vec{x}) \approx r_s(\vec{x}) \left[1 + \frac{6}{5} f_{nl} \mathcal{R}_L(\vec{x}) \right]$$

$$\begin{aligned} \frac{\Delta T}{T} &\sim \frac{\mathcal{R}}{5} \rightarrow C_l^{(T,d)} \sim f_{nl}(k_d) C_l^{TT} \\ \frac{\Delta d}{d} &\sim \frac{\delta \langle \mathcal{R}^2 \rangle}{\langle \mathcal{R}^2 \rangle} \end{aligned}$$



Scale-dependence of non-Gaussianity



Smallest detectable μ -T and y -T correlations ($1-\sigma$):

- template: $B_{\mathcal{R}} \simeq \frac{12}{5} f_{nl}(k_s) P_{\mathcal{R}}(k_s) P_{\mathcal{R}}(k_L)$
- uniform f_{nl} signal on μ and y scales

smallest detectable monopoles for a given experiment

$$f_{nl}^{(\mu)} \simeq 10^2 \left(\frac{\mu_{\min}}{10^{-9}} \right) \left(\frac{\langle \mu \rangle}{2 \times 10^{-8}} \right)^{-1}$$

$$f_{nl}^{(y)} \simeq 10^2 \left(\frac{y_{\min}}{2 \times 10^{-10}} \right) \left(\frac{\langle y \rangle}{4 \times 10^{-9}} \right)^{-1}$$

$$\langle \mu \rangle \approx \int d \log k \Delta_{\mathcal{R}}^2(k) W_{\mu}(k)$$

$$\langle y \rangle \approx \int d \log k \Delta_{\mathcal{R}}^2(k) W_y(k)$$

[Emami, ED, Chluba, Kamionkowski - 2015]

Summary

SDs from diffusion damping gives access to ~ 10 extra inflationary e-folds!

→ **Vast range of modes** to uncover statistics of **primordial fluctuations**

	k/Mpc	fnl
CMB anisot./galaxy surveys	$(10^{-4} - 1)$	bispectrum / off-diagonal correlations, ...
y distortion	$(1 - 50)$	y -T correlation
μ distortion	$(50 - 10^4)$	μ -T correlation

Next :

- More non-Gaussian *observables*, e.g. :
 - Including *residual (r-type) distortions* ($10^4 \lesssim z \lesssim \text{few} \times 10^5$)
- Evaluating *non-primordial* distortions:
e.g. from black-holes, (BSM/DM)particles decays/annihilations, ...

Thank you