



Dark Matter Primordial Black Holes from Particle Production during Inflation

Encieh Erfani
IASBS, Iran

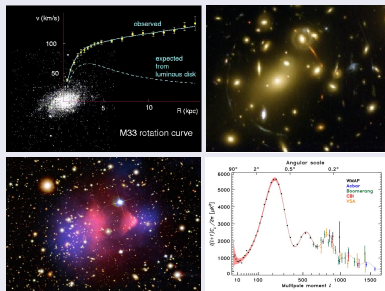
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Dark Matter

Evidences



Properties

- stable
- neutral
- weakly interacting
- right relic density

Candidates

- Axions
- Sterile neutrinos
- WIMPs
- Primordial Black Holes (PBHs)

Primordial Black Holes

Definition

A PBH is a type of black hole that is **not** formed by the gravitational collapse of a star, but by the extreme density of matter present during the Universe's early expansion.

PBHs properties

$$\text{Mass: } M_{\text{BH}} = 10^{15} \left(\frac{t}{10^{-23} \text{ s}} \right) \text{ g}$$

$$M_{\odot} \simeq 2 \times 10^{33} \text{ g}$$

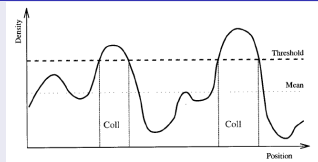
$$\text{Temperature: } T_{\text{BH}} \approx 10^{-7} \left(\frac{M}{M_{\odot}} \right)^{-1} \text{ K}$$

$$\text{Lifetime: } \tau_{\text{BH}} \approx 10^{64} \left(\frac{M}{M_{\odot}} \right)^3 \text{ y}$$

M_{BH}	τ_{BH}
A man	10^{-12} s
A building	1 s
10^{15} g	10^{10} y
The Earth	10^{49} y
The Sun	10^{66} y
The Galaxy	10^{99} y

Press-Schechter Formalism

The Press-Schechter formalism is a model for predicting the number density of bound objects of a certain mass.



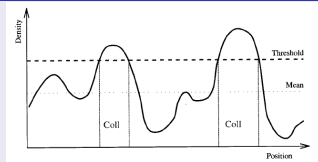
$$f(\geq M) = \gamma \int_{\zeta_{\text{th}}}^{\infty} P(\zeta; M(R))$$

$$\zeta_{\text{th}} = 0.4135$$

$$\gamma = w^{3/2} \simeq 0.2$$

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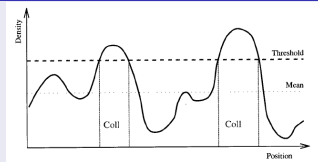
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$$\text{Gaussian PDF: } P_G(\zeta; R) = \frac{1}{\sqrt{2\pi} \sigma_\zeta(R)} \exp\left(-\frac{\zeta^2}{2\sigma_\zeta^2(R)}\right)$$

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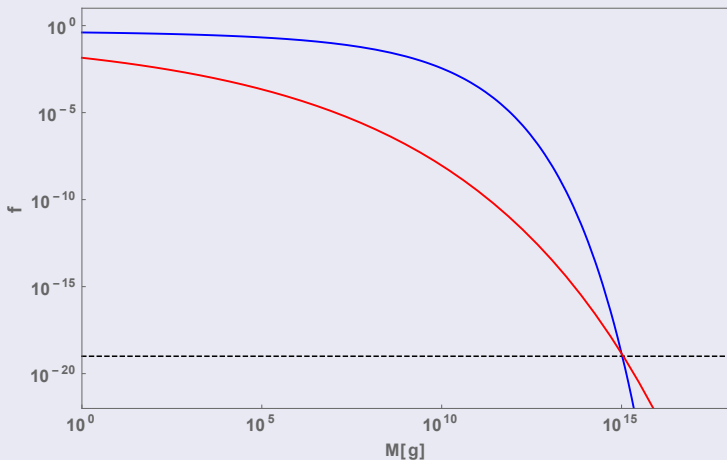
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$$\text{non-Gaussian: } P_{\text{NG}}(\zeta; R) = \frac{1}{\sqrt{2\pi} (\zeta + \sigma_g^2(R)) \sigma_g(R)} \exp\left(-\frac{\zeta + \sigma_g^2(R)}{2\sigma_g^2(R)}\right)$$

$$f_{\text{NG}} = \text{erfc}\left(\sqrt{\zeta_{\text{th}} + \sigma_g^2(R)} / \sqrt{2\sigma_g^2(R)}\right)$$

$f(\geq M)$ diagram for the mass range $10^0 - 10^{20}$ g



Result

$$\begin{aligned} n_s(k_{\text{PBH}}) \geq 1.418 &\Rightarrow \mathcal{P}_\zeta \simeq 2 \times 10^{-2} && \text{for Gaussian PDF} \\ n_s(k_{\text{PBH}}) \geq 1.322 &\Rightarrow \mathcal{P}_\zeta \simeq 4 \times 10^{-4} && \text{for non-Gaussian PDF} \end{aligned}$$

Inflation parameters

$$\mathcal{P}_{\zeta, \text{vac.}}(k) = \mathcal{P}_{\zeta, \text{vac.}}(k_0) \left(\frac{k}{k_0} \right)^{n_s(k)-1}$$

$$n_s(k_0) - 1 \equiv \frac{d \ln \mathcal{P}_{\zeta, \text{vac.}}(k)}{d \ln k}$$

$$r = \frac{\mathcal{P}_t(k)}{\mathcal{P}_{\zeta}(k)}, \quad \mathcal{P}_{t, \text{vac.}}(k) = \frac{2}{\pi^2} \left(\frac{H}{M_{\text{P}}} \right)^2 \left(\frac{k}{k_0} \right)^{n_t}$$

$$B_{\zeta}(k_1, k_2, k_3) = f_{\text{NL}} F(k_1, k_2, k_3)$$

Observation

Planck XX, arXiv: 1502.01592

$$\ln(10^{10} \mathcal{P}_{\zeta, \text{vac.}}(k_0)) = 3.094 \pm 0.034$$

$$k_0 = 0.05 \text{ Mpc}^{-1}$$

$$n_s = 0.9645 \pm 0.0049$$

$$r_{0.002} < 0.10 \quad (95\% \text{ CL})$$

$$f_{\text{NL}} = 22.7 \pm 25.5$$

PBHs formation from Particle Production

direct or gravitational coupling of the inflaton (ϕ) to another field (χ)

$$\mathcal{L}(\phi, \chi) = -\frac{1}{2}\partial_\mu\phi\partial^\mu\phi - V(\phi) - \frac{1}{2}\partial_\mu\chi\partial^\mu\chi - U(\chi) + \mathcal{L}_{\text{int}}(\phi, \chi)$$

The equations of motion for the inflaton field:

$$H^2 = \frac{1}{3M_{\text{P}}^2} \left(\frac{1}{2}\dot{\phi}^2 + V(\phi) + \rho_\chi \right)$$

$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = \frac{\partial\mathcal{L}_{\text{int}}}{\partial\phi}$$

The inflaton fluctuations satisfy

$$\delta\ddot{\phi} + 3H\delta\dot{\phi} - \frac{\nabla^2}{a^2}\delta\phi + V''(\phi)\delta\phi = \delta\left(\frac{\partial\mathcal{L}_{\text{int}}}{\partial\phi}\right)$$

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Result

$$\mathcal{P}_\zeta(k) = \mathcal{P}_{\zeta, \text{vac.}}(k) + \mathcal{P}_{\zeta, \text{src.}}(k)$$

$$\mathcal{P}_t(k) = \mathcal{P}_{t, \text{vac.}}(k) + \mathcal{P}_{t, \text{src.}}(k)$$

$$\mathcal{L}_{\text{int}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{\alpha}{4f}\Phi F_{\mu\nu}\tilde{F}^{\mu\nu}$$

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direct coupling

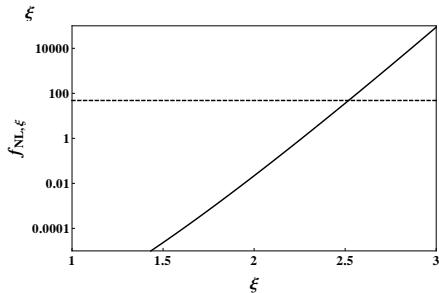
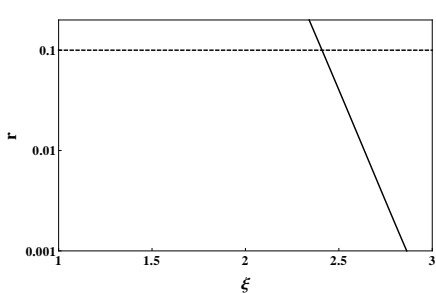
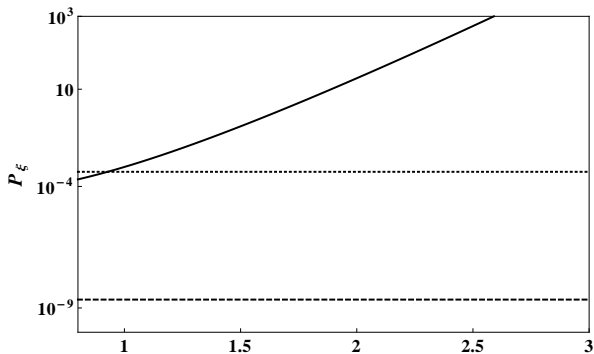
$$\mathcal{P}_{\zeta} = \mathcal{P}_{\zeta, \text{vac.}} (1 + 7.5 \times 10^{-5} \epsilon^2 \mathcal{P}_{\zeta, \text{vac.}} X^2)$$

$$r = 16\epsilon \frac{1 + 2.2 \times 10^{-7} \mathcal{P}_{t, \text{vac.}} X^2}{1 + 7.5 \times 10^{-5} \epsilon^2 \mathcal{P}_{\zeta, \text{vac.}} X^2}$$

$$f_{\text{NL}, \zeta}^{\text{equil.}} \approx 4.4 \times 10^{10} \epsilon^3 \mathcal{P}_{\zeta, \text{vac.}}^3 X^3$$

where

$$X \equiv \frac{e^{2\pi\xi}}{\xi^3} \quad \xi \equiv \frac{\alpha}{2fH} \dot{\Phi}$$

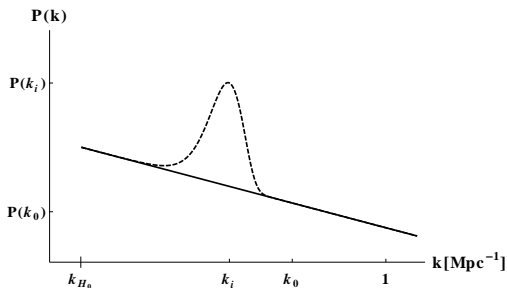


Scalar Production

$$\mathcal{L}_{\text{int}}(\phi, \chi) = -\frac{g^2}{2} (\phi - \phi_0)^2 \chi^2$$

$$\mathcal{P}_{\zeta, \text{src.}}(k) \sim A k^3 e^{-\frac{\pi}{2} \left(\frac{k}{k_i}\right)^2}$$

Result



$$A \lesssim 4 \times 10^{-4}$$

Conclusions

- The fluctuation which arise at inflation are the most likely source of PBHs.
- The spectral index at scale of PBHs formation should be at least 1.418 (1.322) for Gaussian (non-Gaussian) PDF.
- The most stringent constraints on the gauge production parameter is derived from the non-production of DM PBHs at the end of inflation and the bounds from the bispectrum and the tensor-to-scalar ratio are weaker.
- In the scenario where the inflaton field coupled to a scalar field, the model is free of DM PBHs overproduction in the CMB observational range if the amplitude of the generated bump in the scalar power spectrum, A is less than 4×10^{-4} .

