

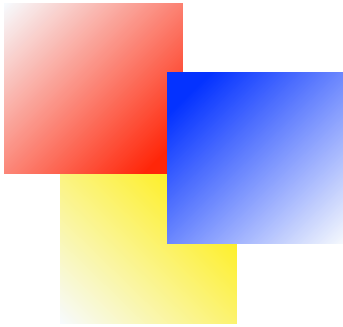


Generation of upward Alfvénic waves by the chromospheric shock waves

Munehito Shoda¹

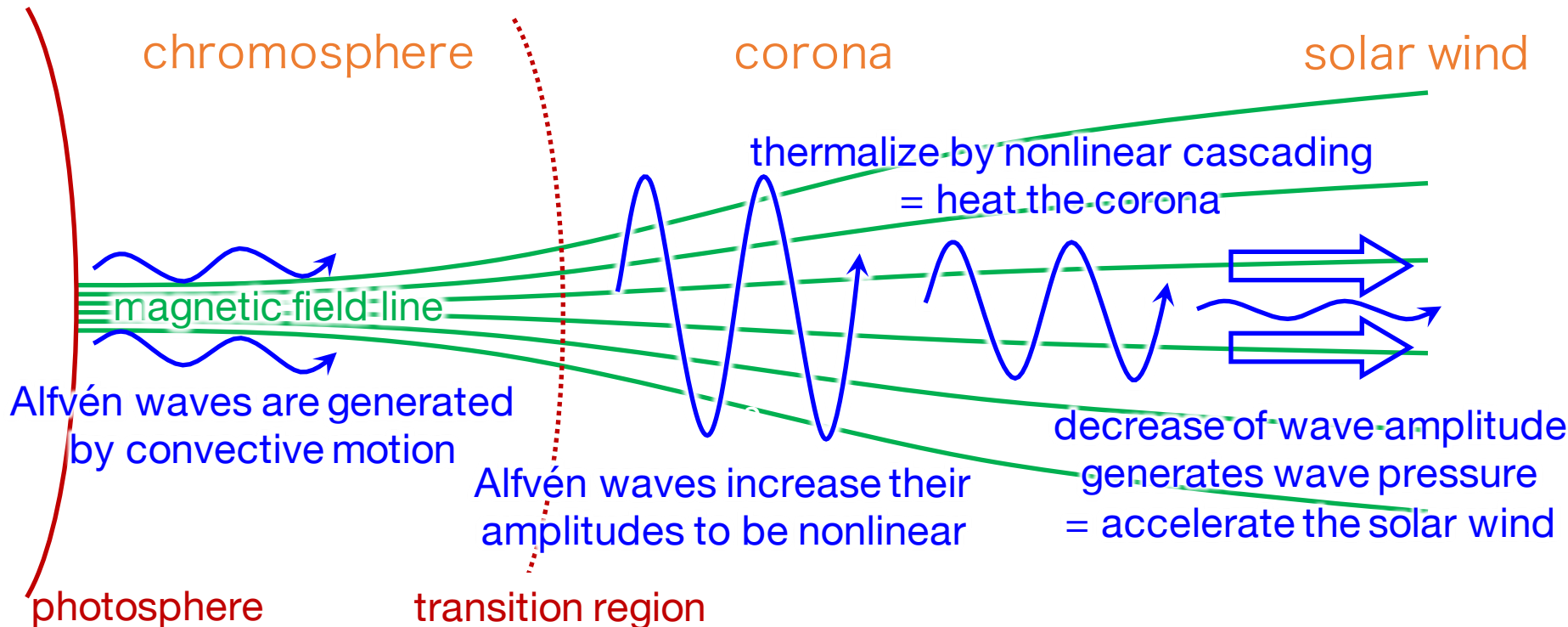
Ryuichi Kanoh², Takaaki Yokoyama¹, Toshifumi Shimizu²

1. The University of Tokyo 2. JAXA/ISAS



Introduction

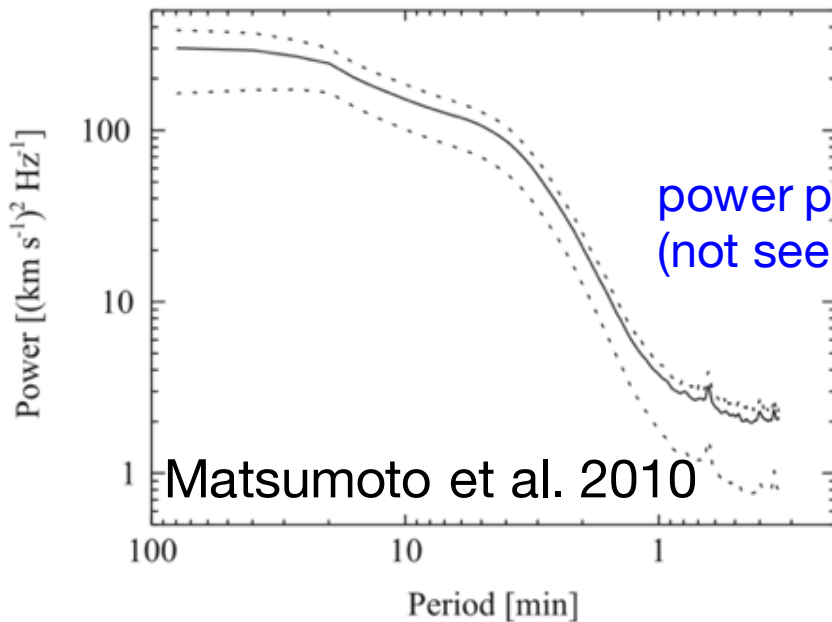
standard Alfvén-wave modeling of the corona & wind



Introduction

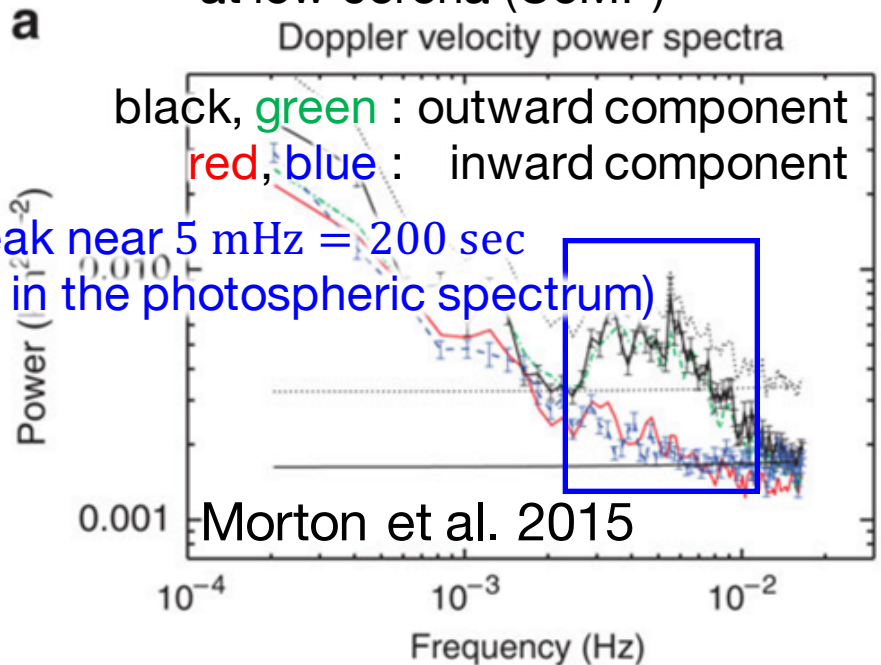
Is the standard understanding correct ?

power spectrum of transverse motion at the photosphere (*Hinode/SOT*)



power peak near 5 mHz = 200 sec
(not seen in the photospheric spectrum)

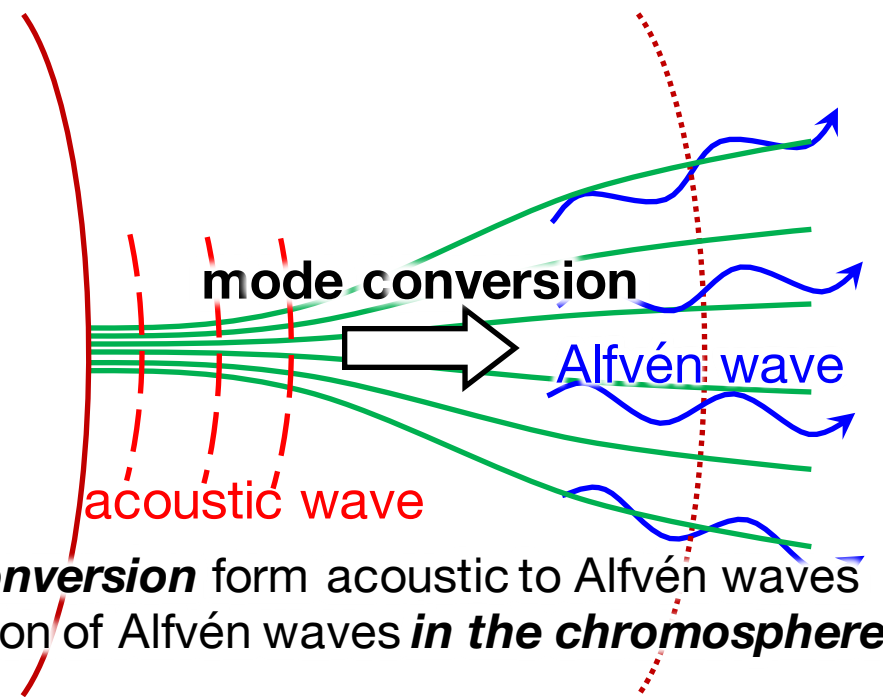
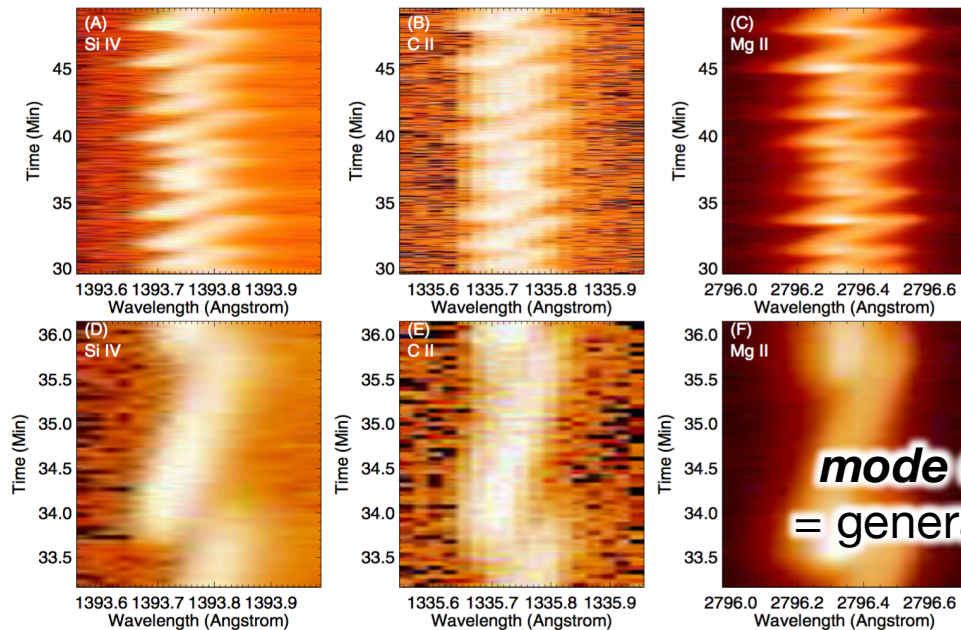
power spectra of Alfvén waves at low corona (CoMP)



Introduction

mode-conversion in the chromosphere

shock waves (N-waves) observed
in the chromosphere (Tian et al. 2014)



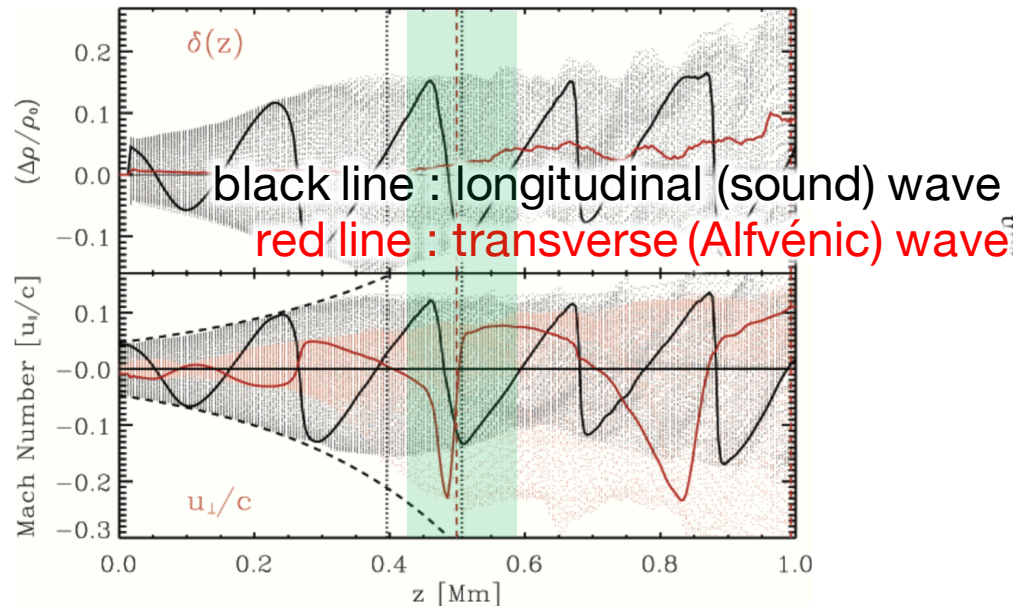
Introduction

mode-conversion in the chromosphere

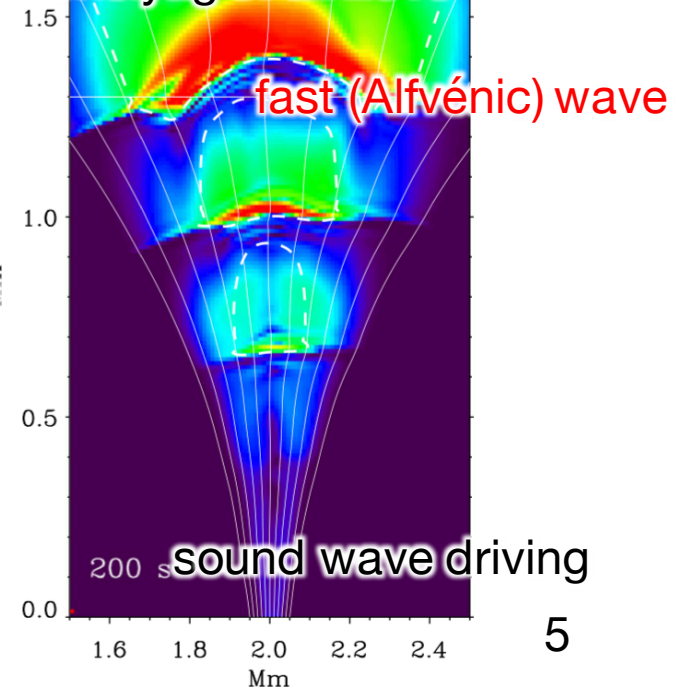
Many previous researches on mode conversion
(Bogdan+03, Cally+06, Khomenko+12, Shelyag+16 etc.)

Bogdan et al. 2003

$\beta \sim 1$



Shelyag et al. 2016

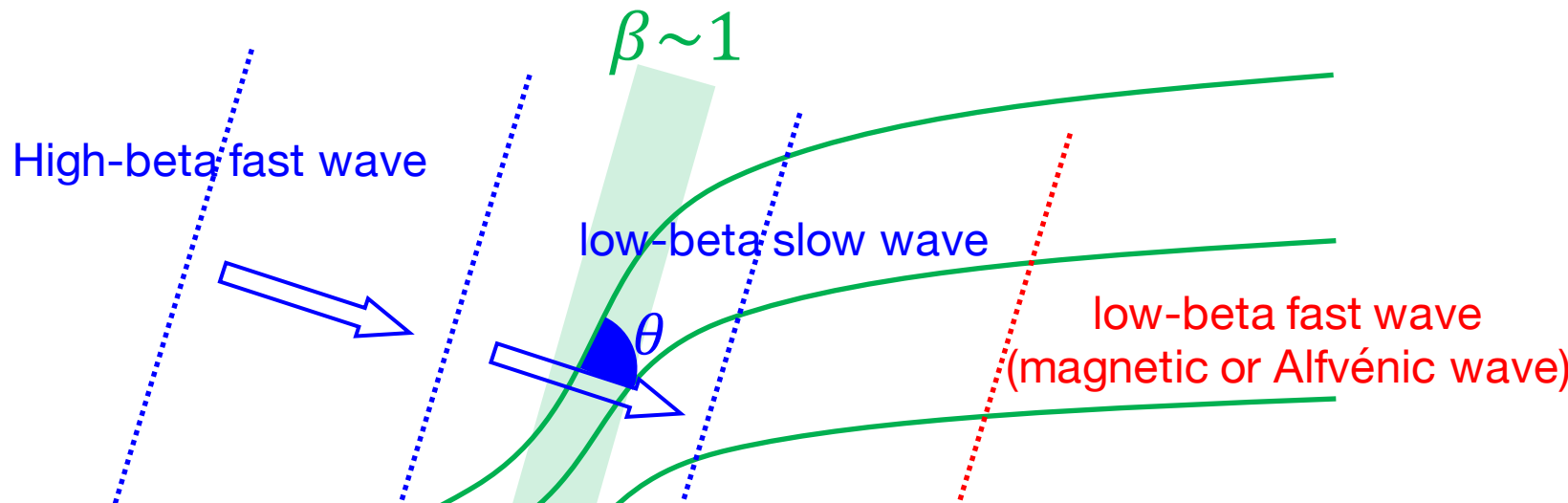




Introduction

mode-conversion: elemental process

Critical parameter for mode conversion is the *attacking angle* θ , the angle between the propagation direction & magnetic field

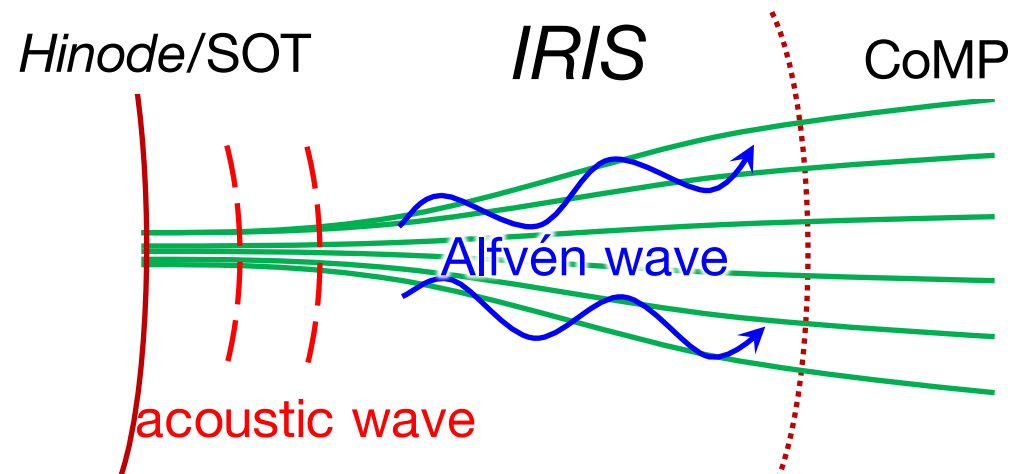


Acoustic wave – Alfvén wave interaction **inside the flux tube**
triggers similar situation

Introduction

motivation of this study

We study the mode conversion process ***inside the flux tube***
all waves propagate along the background flux tube (1D system)
impose both longitudinal & transverse waves at the photosphere



Try to find an evidence of mode conversion ***using IRIS***



Simulation method

Basic equations to be solved

1D *isothermal* MHD equations including gravity & flux tube expansion

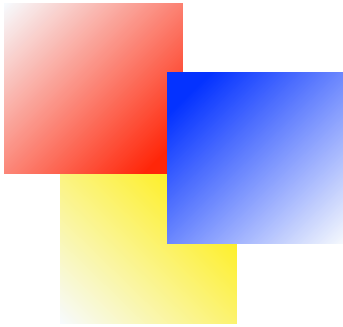
$$\frac{\partial}{\partial t} (\rho A) + \frac{\partial}{\partial x} (\rho v_x A) = 0,$$

$$\frac{\partial}{\partial t} \left(\rho v_x A \right) + \frac{\partial}{\partial x} \left[\left(\rho v_x v_x + \rho v_x v_{\perp} + \frac{B_{\perp}^2}{4\pi} \right) A \right] - \rho g A + dA \left(1 - \frac{v_{\perp}^2}{a^2} + \rho a^2 \right),$$

We do NOT solve energy equations / radiative transfer to make the physics simple.

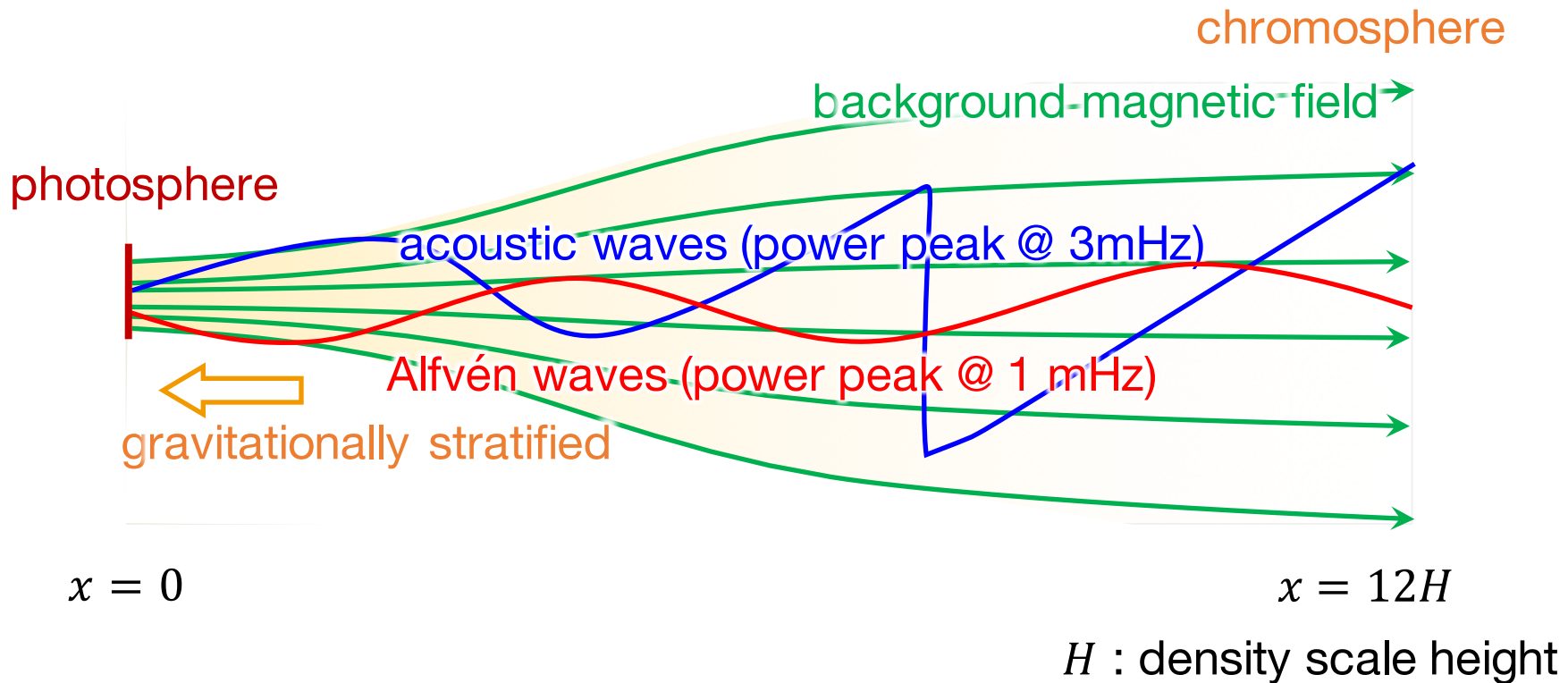
$$\frac{\partial}{\partial t} \left(\rho v_{\perp} A \right) + \frac{\partial}{\partial x} \left[\left(\rho v_x v_{\perp} + \frac{B_{\perp} v_x}{4\pi} \right) A \right] = 0,$$

$$\frac{\partial}{\partial t} \left(\mathbf{B}_{\perp} \sqrt{A} \right) + \frac{\partial}{\partial x} \left[\left(\mathbf{B}_{\perp} v_x - B_x \mathbf{v}_{\perp} \right) \sqrt{A} \right] = 0,$$



Simulation method

Numerical Setup : overview of simulation



Simulation result

Numerical Result 1 : case I (uniform flux tube)

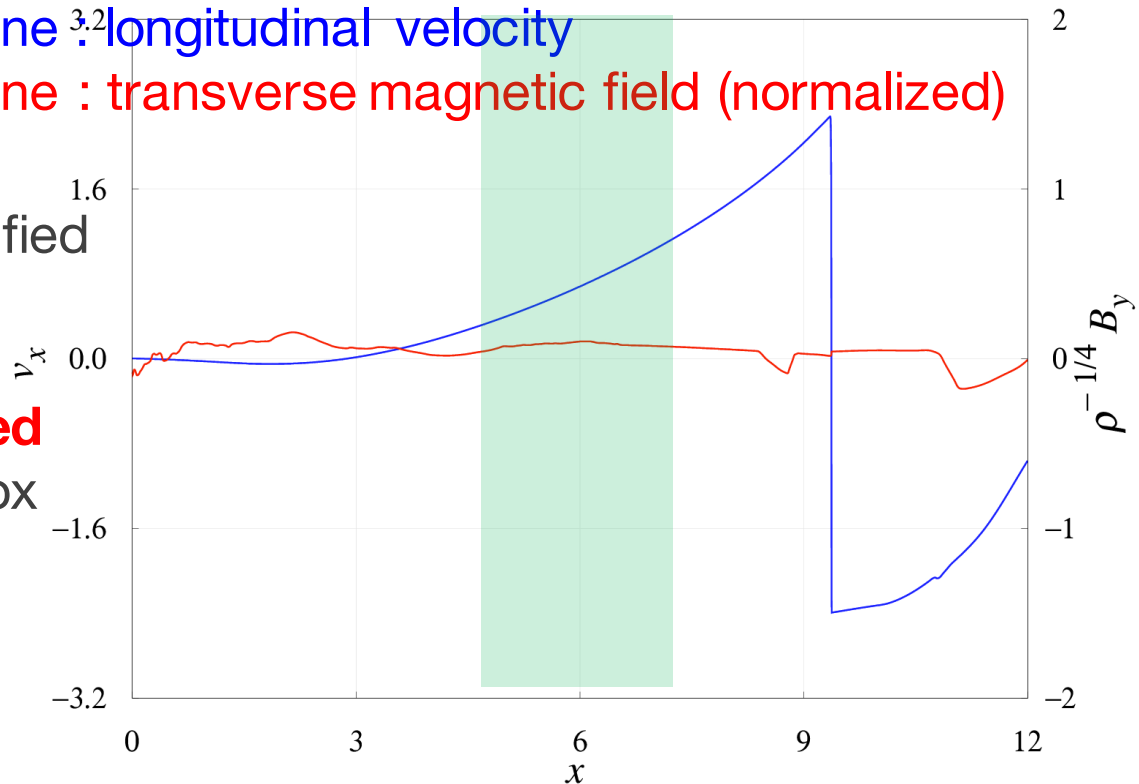
$$\beta \approx 1$$

blue line : longitudinal velocity

red line : transverse magnetic field (normalized)

Acoustic waves are amplified due to the stratification.

Alfvén waves are **amplified near the center** of the box (~magnetic canopy).



Simulation result

Numerical Result 2 : case II (expanding tube)

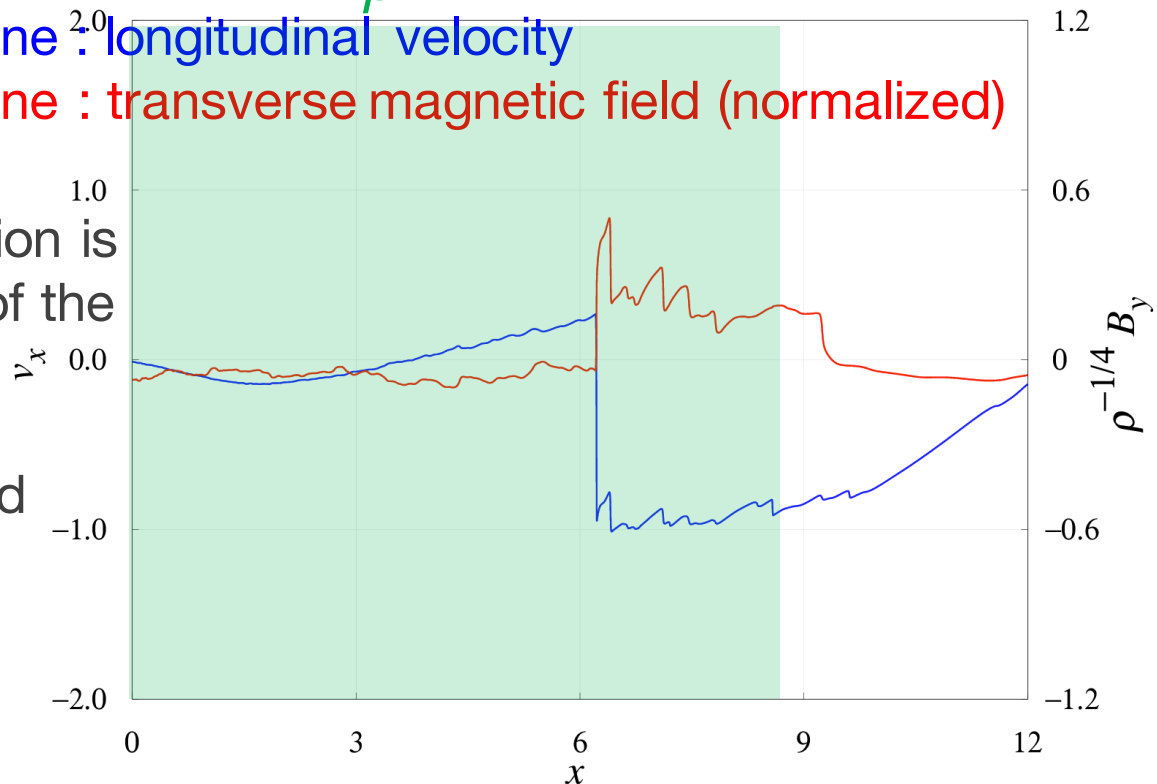
$$\beta \approx 1$$

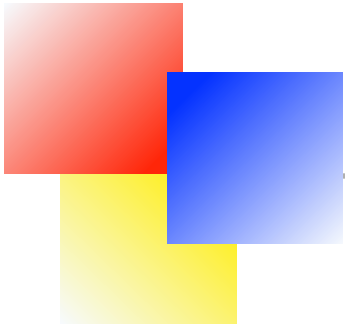
blue line : longitudinal velocity

red line : transverse magnetic field (normalized)

Acoustic-wave amplification is less significant because of the tube expansion.

Alfvén waves are amplified also in this case.

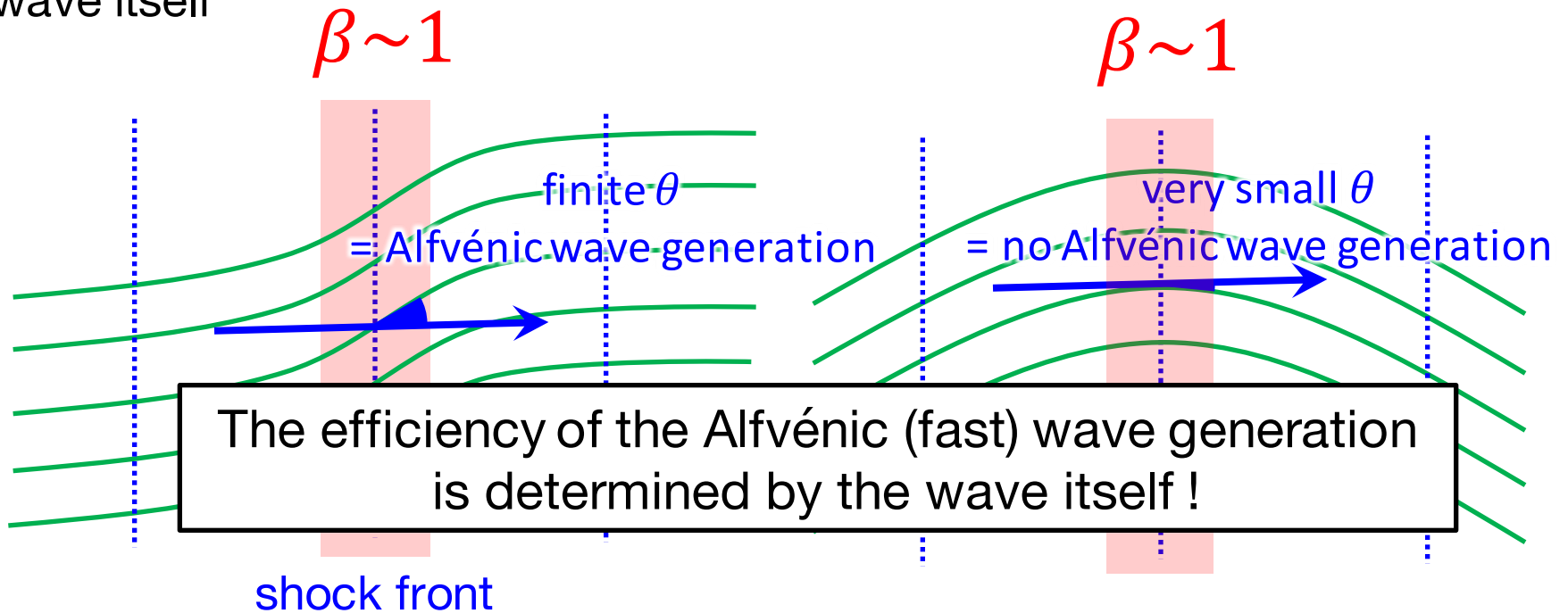


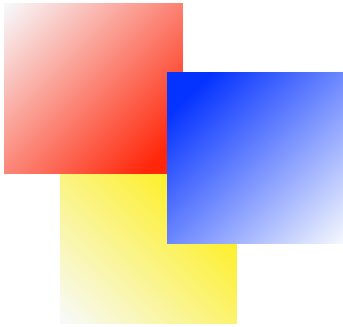


Simulation result

Role of Alfvén wave input

Attacking angle, the critical parameter is determined by Alfvénic (transverse) wave itself





Simulation result

amplification of Alfvénic flux (uniform-tube case)

red : without acoustic waves

weak acoustic wave

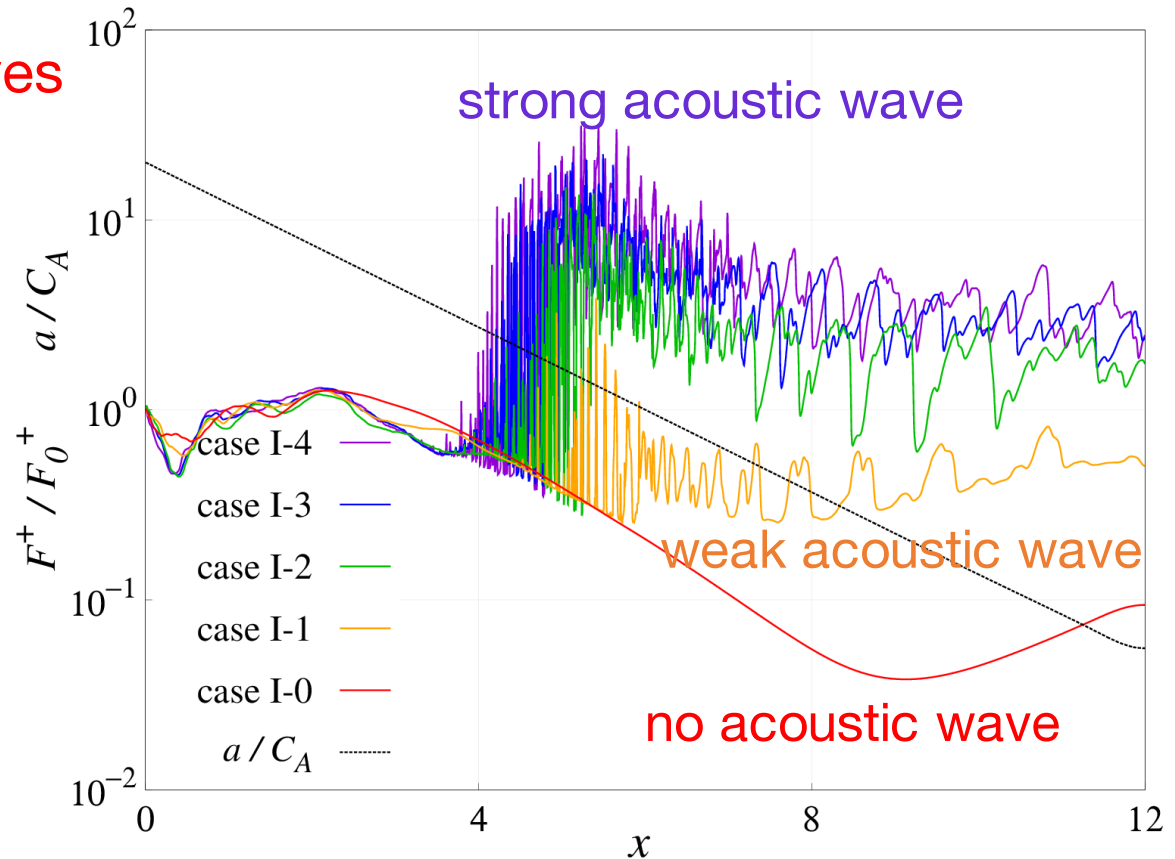
orange

green

blue

violet

strong acoustic wave



Observation method

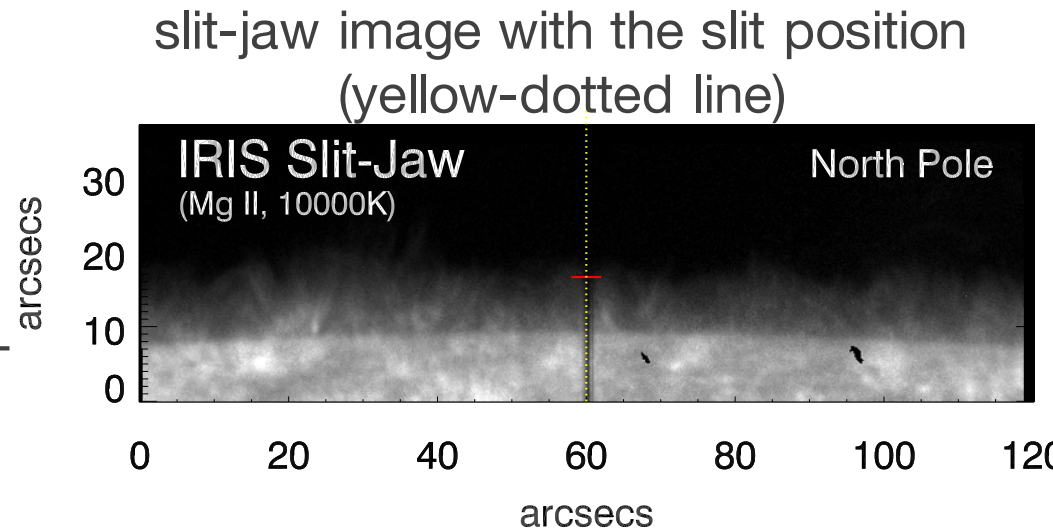
Observation by IRIS

observation time:
2014. 4. 4. 00:20UT~01:32UT

cadence: 9 seconds

sit-and-stare observation of the northern polar region

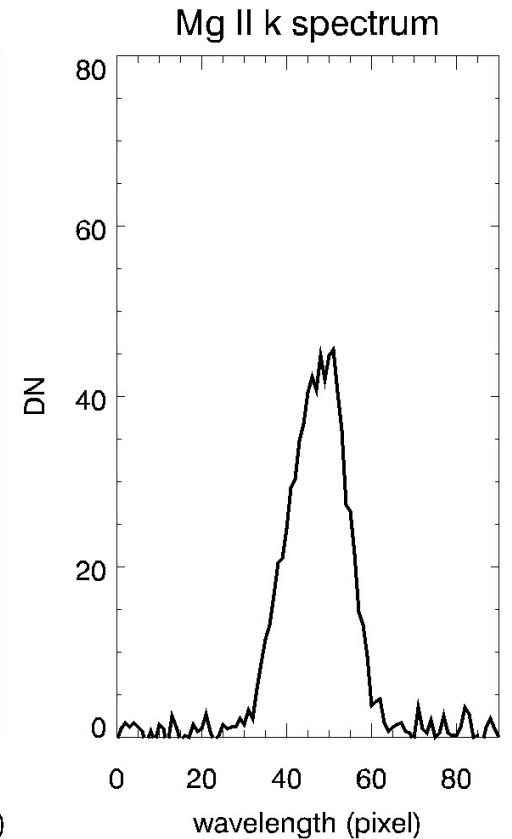
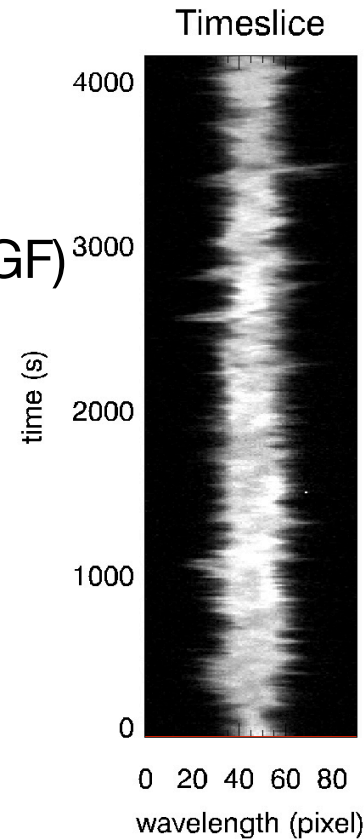
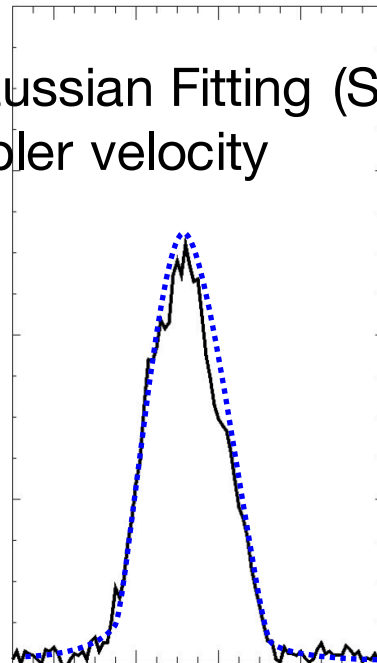
Mg II k line is used for analysis



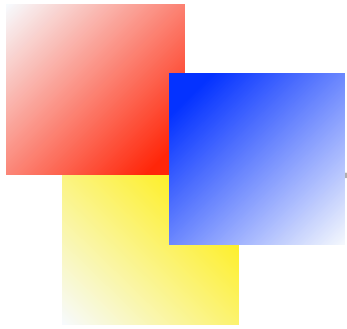
Observation result

Line profile & timeslice at one pixel

Applied Single Gaussian Fitting (SGF)
to derive the Doppler velocity



2016/6/23



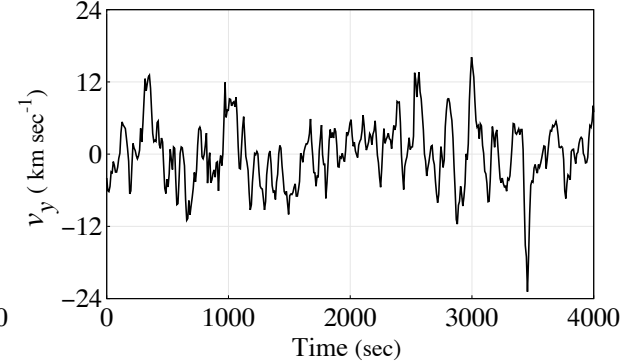
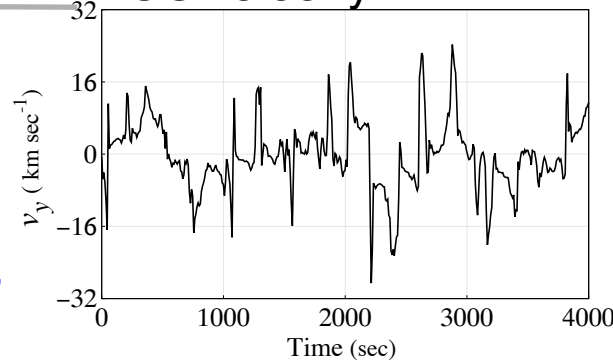
Result

*Observational Result
: comparison with simulation*

simulation

observation

LOS velocity



From top to bottom

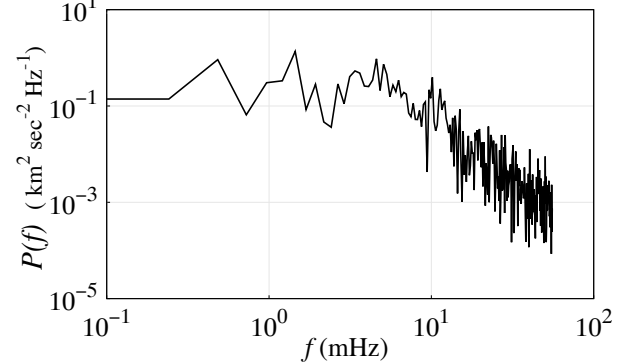
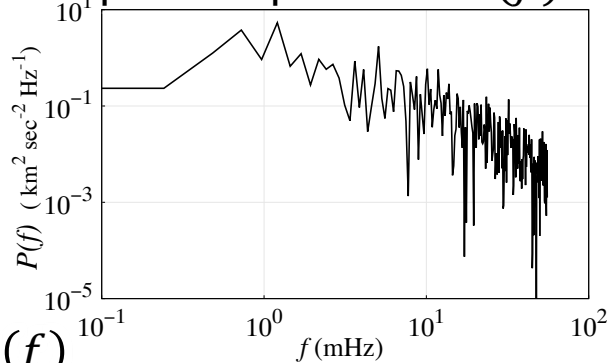
line-of-sight velocity v_y

power spectrum $P(f)$

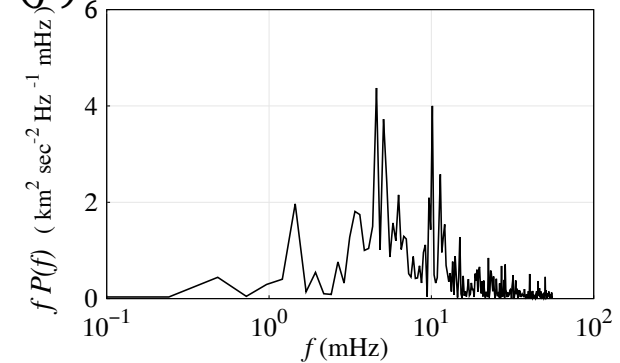
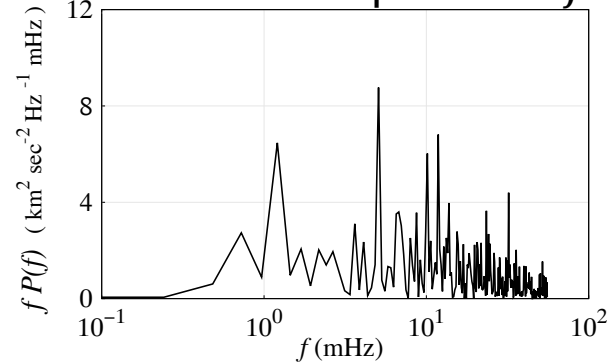
normalized power spectrum $fP(f)$

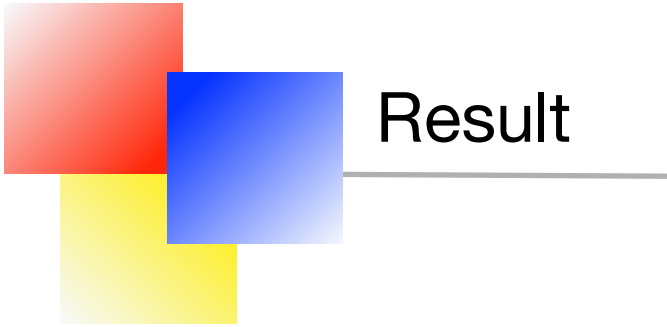
= representing the energy-containing scale

power spectrum $P(f)$



normalized spectrum $fP(f)$



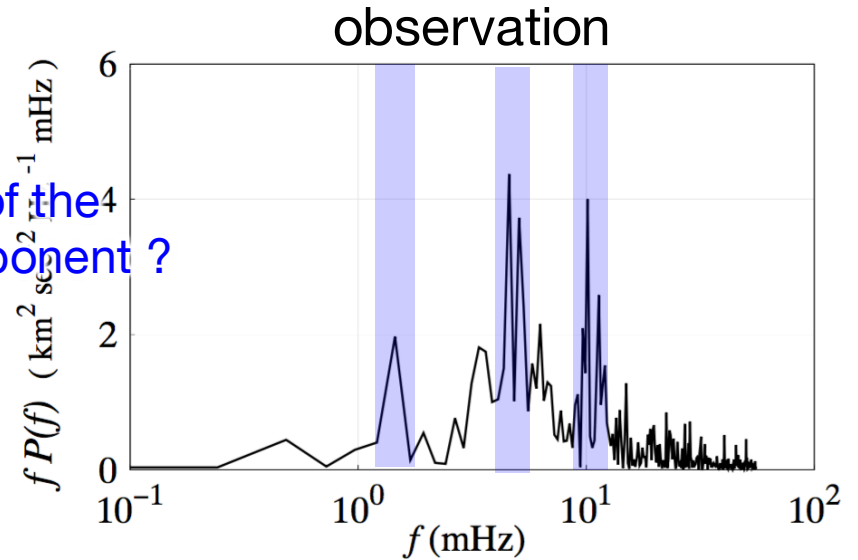
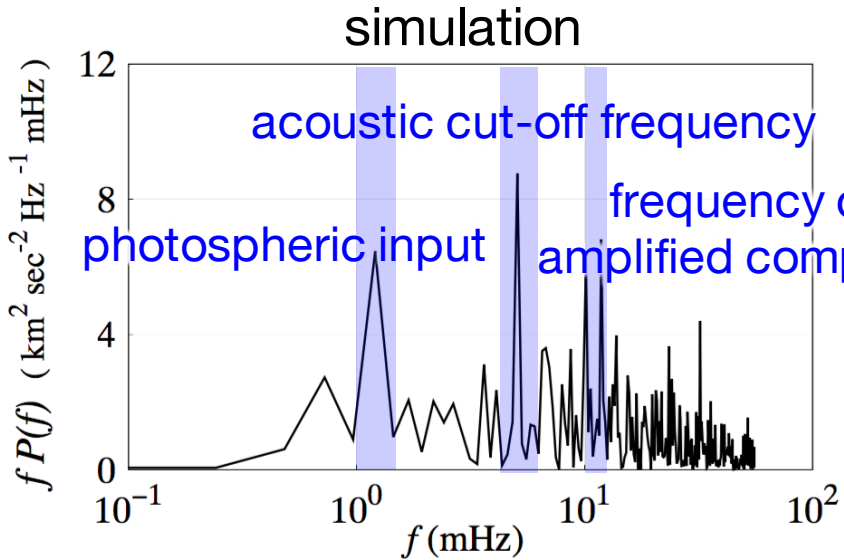
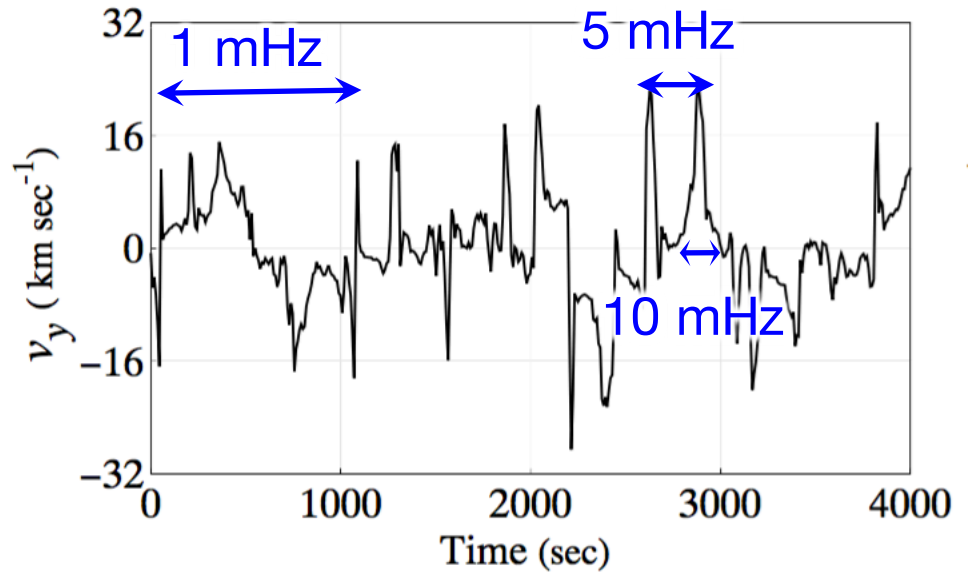


Result

Observational Result : cor

Typical three peaks (1mHz, 5n simulation and observation)

mode conversion occurs under the formation height of Mg II k line



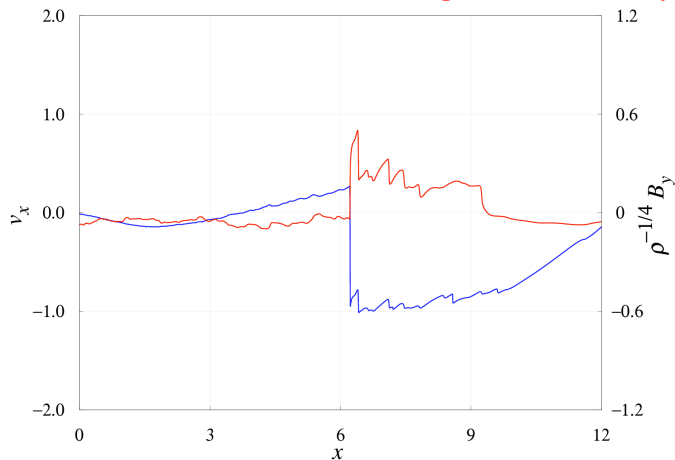
Summary

Using 1D simulation, we have shown intermittent, impulsive Alfvénic-wave generation triggered by acoustic (shock) waves inside the flux tube.

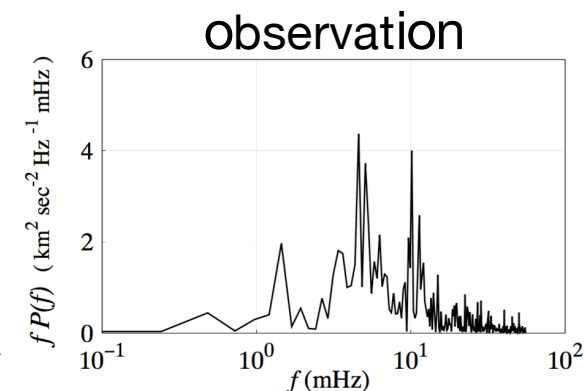
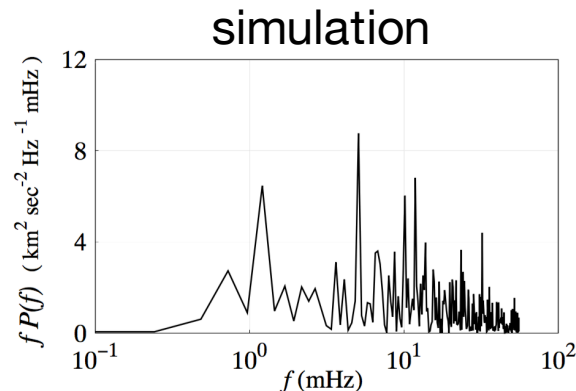
Similarity of the normalized power spectrum (energy-containing scale) shows the mode conversion occurs under the Mg II k formation height.

blue line : longitudinal velocity

red line : transverse magnetic field (normalized)



Normalized power spectrum
of the transverse velocity

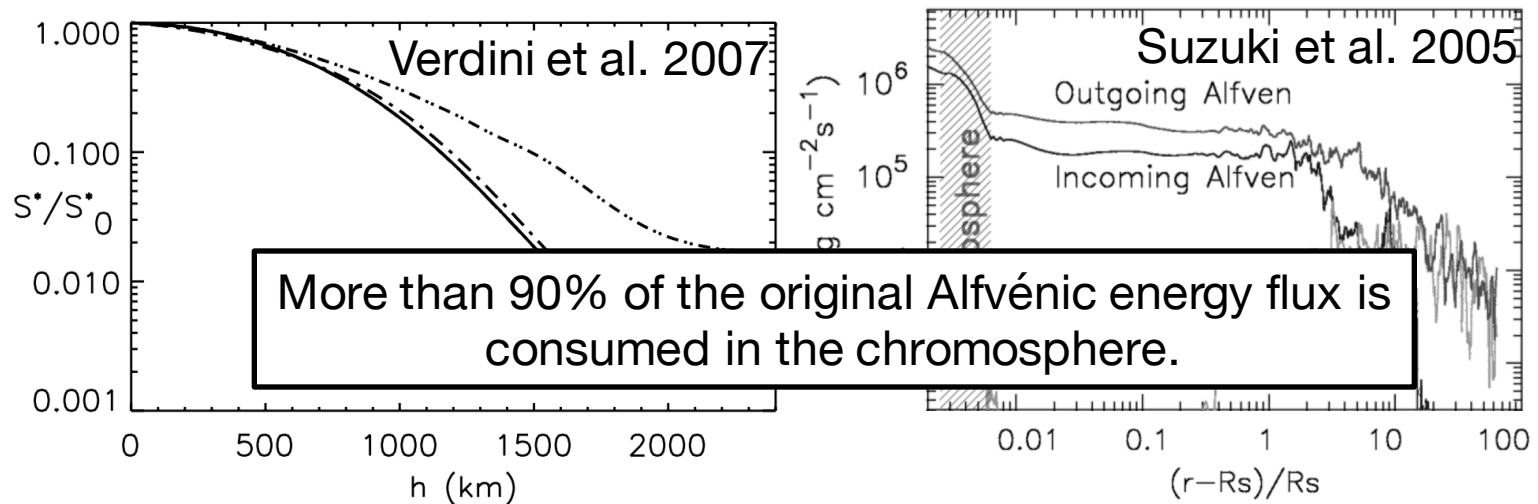


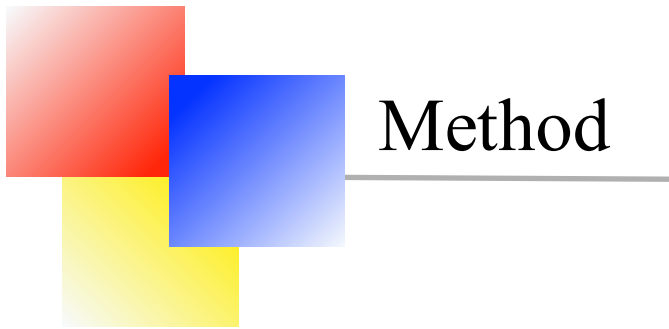
Introduction

On the role of the chromosphere in coronal heating

The chromosphere was thought to be just a *dissipative medium connecting the wave source & the corona*.

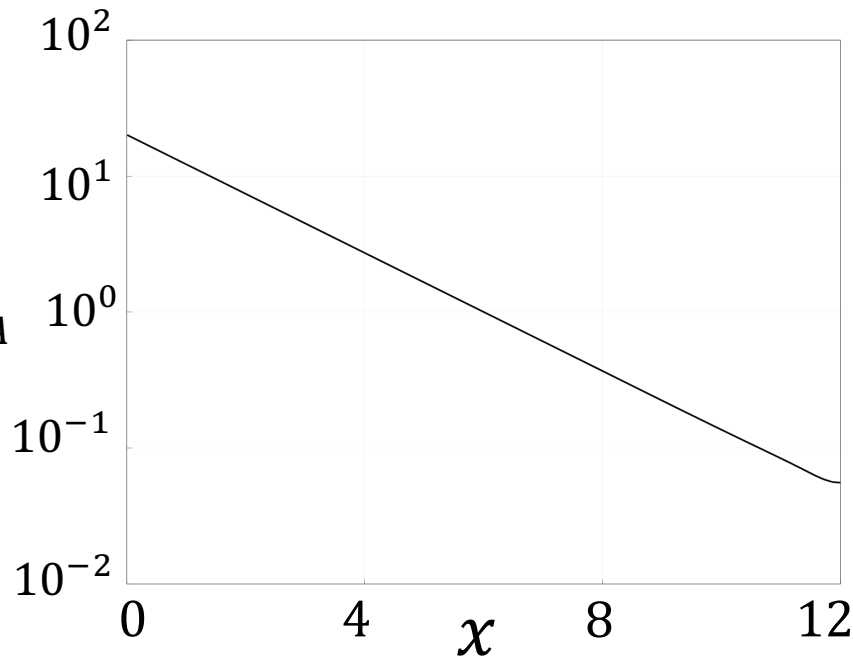
Alfvén-wave energy flux as a function of height (two different models).



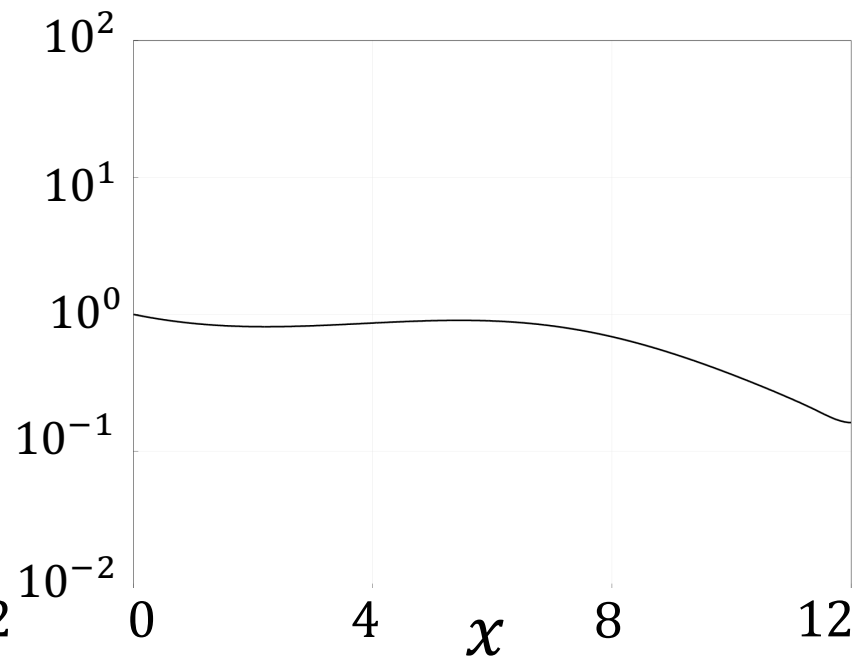


Numerical Setup : background plasma beta

case I : uniform flux tube



case II : expanding flux tube





Results

Numerical Result : amplification of Alfvénic flux (case I)

Upward energy flux of Alfvén(ic) wave is obtained as below.

$$F^+ = \frac{1}{T_{\text{sim}}} \int_0^{T_{\text{sim}}} \left[\frac{1}{4} \rho \mathbf{z}_+^2 C_A A \right] dt$$

T_{sim} : simulation time (needed for temporal averaging)

\mathbf{z}_+ : outward Elsässer variable $\mathbf{z}_+ = \mathbf{v}_\perp - \mathbf{B}_\perp / \sqrt{4\pi\rho}$

A : cross section of magnetic flux tube

Discussion

Effect of ambipolar diffusion

Ambipolar diffusion (diffusion caused by neutrals) is recently thought to be important for high-frequency (≥ 10 mHz) Alfvén waves.
(Goodman 2011, Khomenko+ 2012, Shelyag+ 2016)

