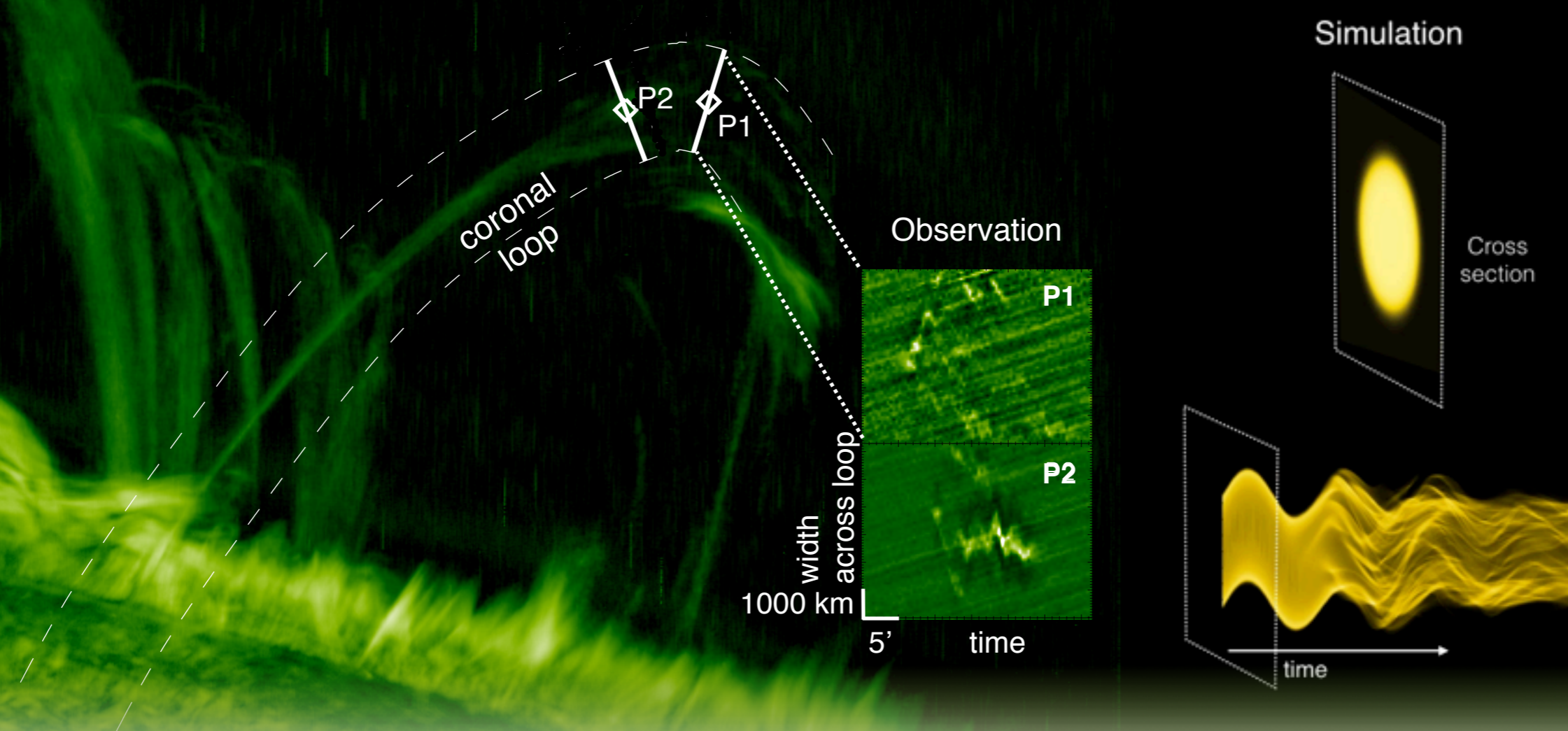


Resonant Absorption as a Feeding Mechanism for Alfvénic Turbulence and its Observable Characteristics in the Solar Atmosphere



IRIS-6@Stockholm,
20-24 June 2016



Patrick Antolin¹

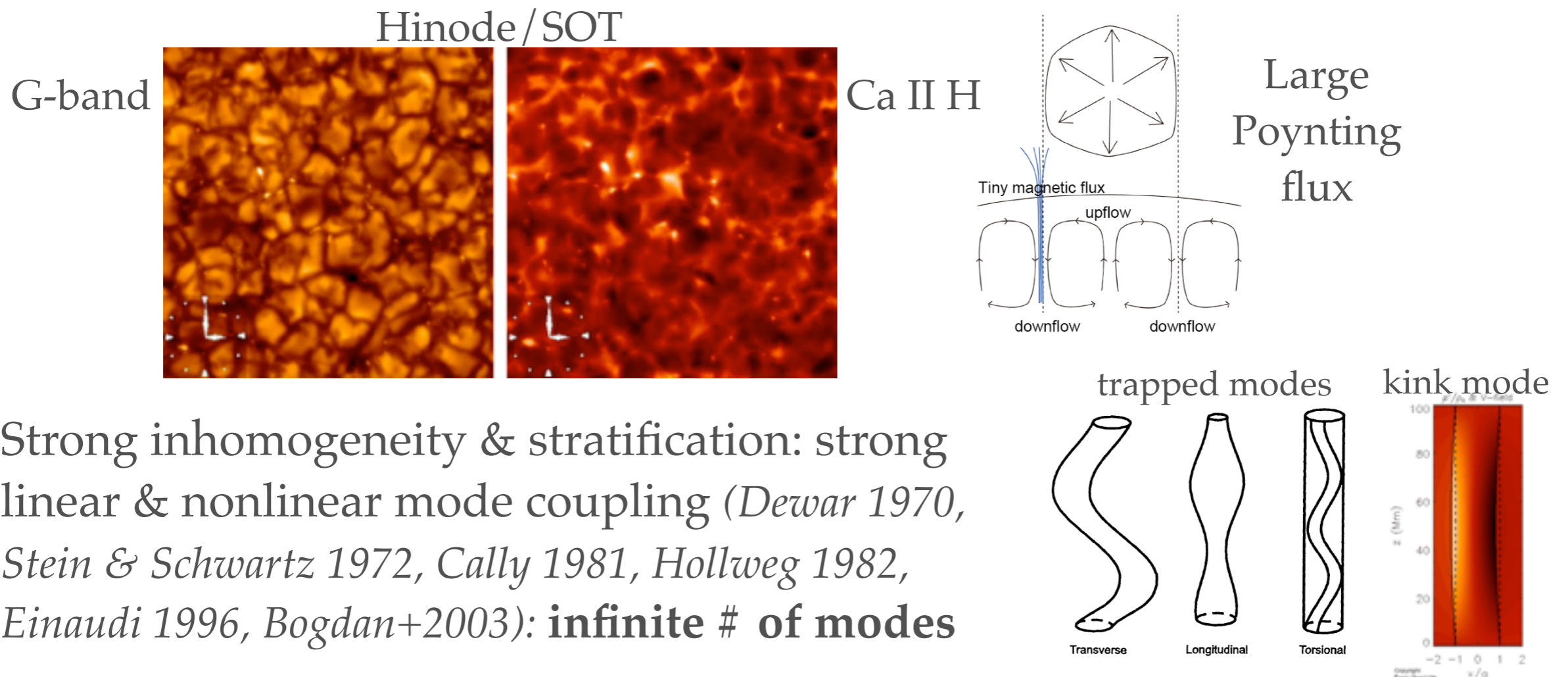


I. De Moortel¹, T. Van Doorselaere², T. Yokoyama³, T. J. Okamoto⁴, B. De Pontieu⁵

¹: U. St Andrews, ²: KU Leuven; ³: U. of Tokyo, ⁴: NAOJ, ⁵: LMSAL

Wave generation in the lower atmosphere

- Convective turbulent motion (*Biermann 1946, Schwarzschild 1948*)
- p-mode leakage ('ramp' effect) (*Michalitsanos 1973; Bel & Leroy 1977; Suematsu 1990, De Pontieu+ 2004*)
- Magnetic reconnection
- Large Poynting flux: heating candidates for chromosphere & corona (*Uchida & Kaburaki 1974, Wentzel 1974, Narain & Ulmschneider 1996*)



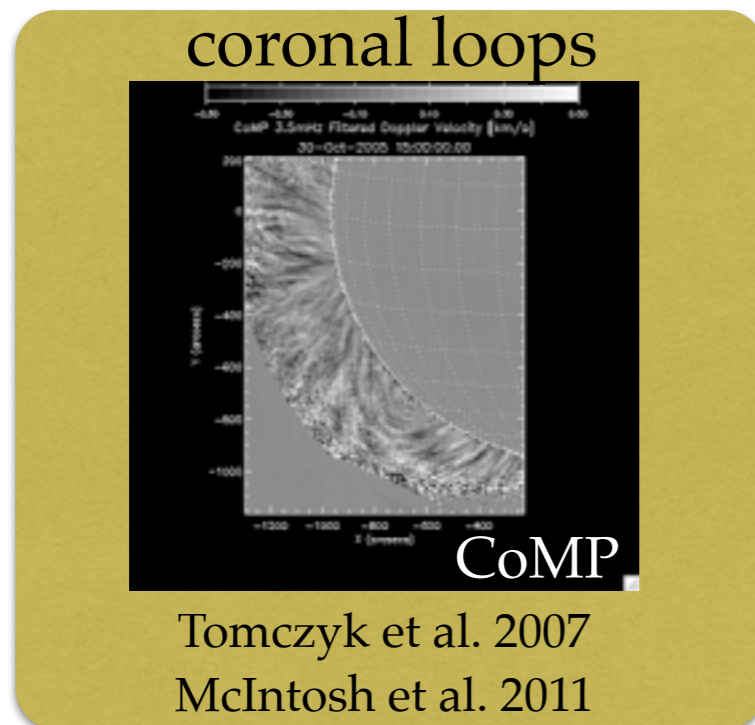
Observations of Alfvénic waves in the solar atmosphere

Transverse MHD waves ubiquitous in the solar atmosphere

What is the role of such waves in the solar atmosphere?

Are we detecting all the wave power?

Small amplitudes <
coronal heating



Large amplitudes >
chromospheric heating



(Tomczyk+ 2007, Okamoto+ 2007, De Pontieu+ 2007, Lin 2011, McIntosh+ 2011, Morton+ 2011, Antolin & Verwichte 2011, Okamoto & De Pontieu 2012, Hillier+ 2013, Schmieder+ 2013, Morton & McLaughlin 2014, De Pontieu+ 2014, Anfinogentov+ 2013, Nisticó+ 2013)

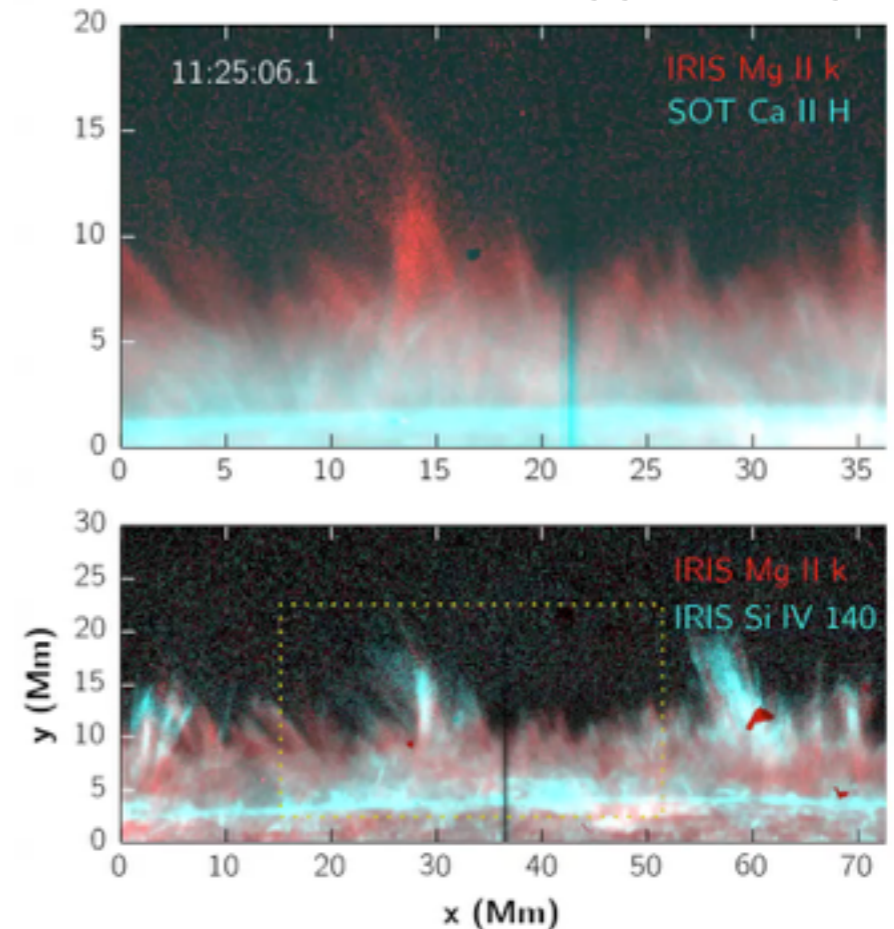
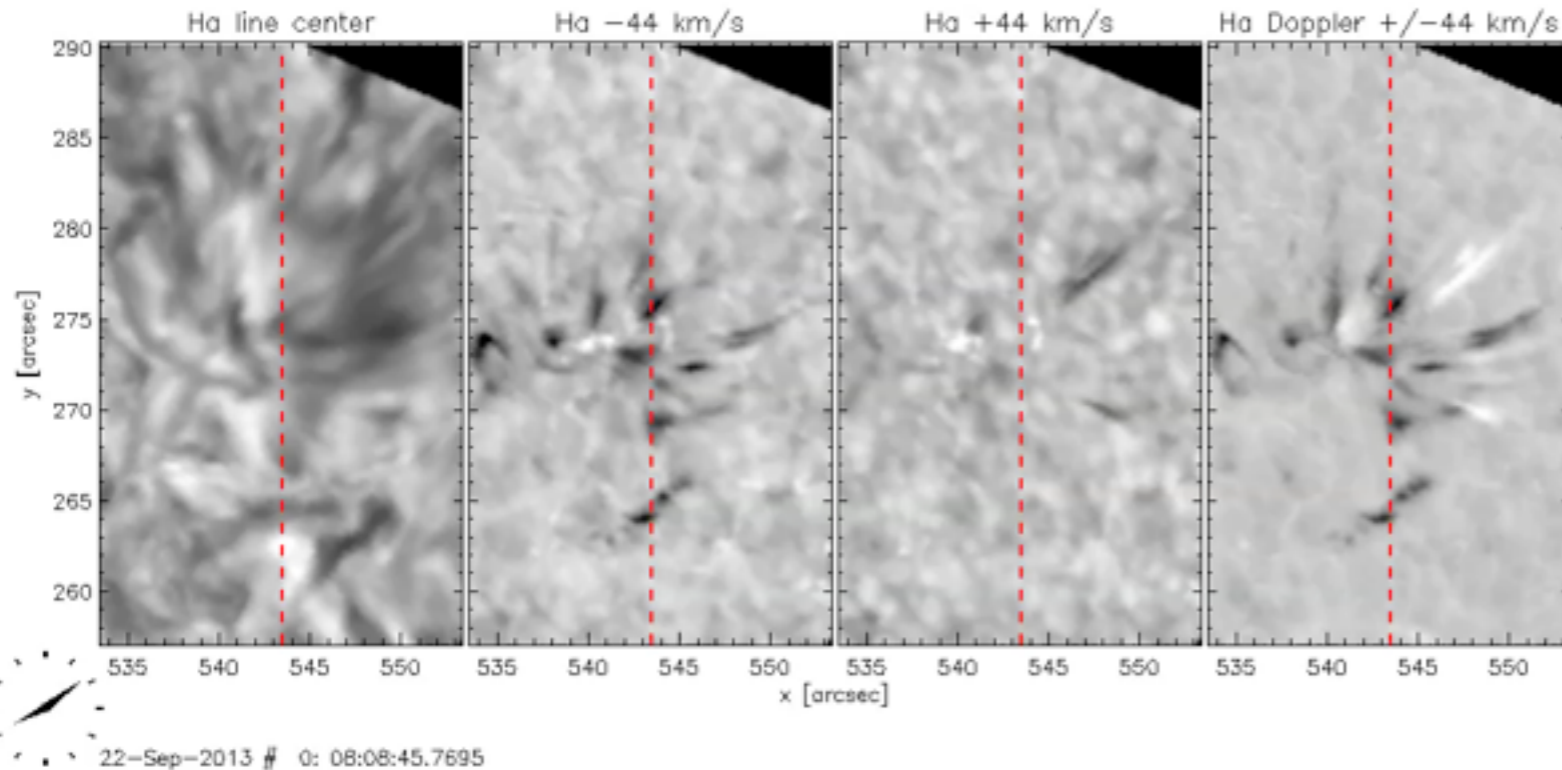
Dynamics of spicules

Alfvénic motions

SST / CRISP Observations

Roupe van der Voort+ 2015

SST + IRIS



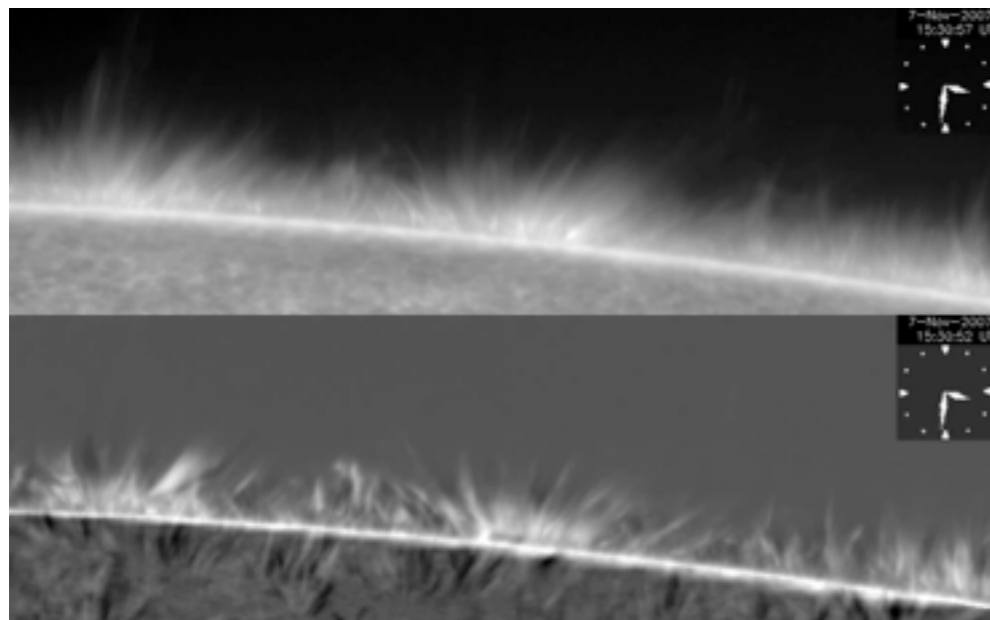
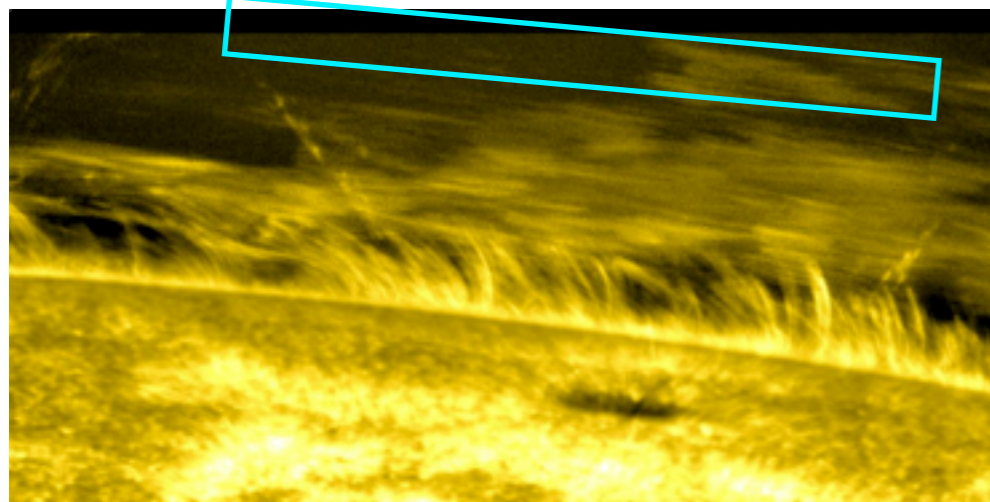
Pereira et al. 2014

- Ubiquitous jets protruding from the chromosphere into the corona
- 2 types: I - magneto-acoustic wave driven (ballistic, $v < 40$ km/s) (*Beckers 1968, Sterling 2000*)
 II - Fast disappearance in Ca II H, fast upflow ($v < 110$ km/s), mostly in QS, RBE & RBB on-disc (*De Pontieu+ 2007, Roupe vd Voort 2009, Sekse 2012, 2013*)
- New features for type II: **Multi-stranded, strong heating, swaying and torsional motions** (*Suematsu+ 2008, Pereira+2012, Skogsrud+ 2014, De Pontieu+ 2014, Roupe v.d. Voort 2015*)

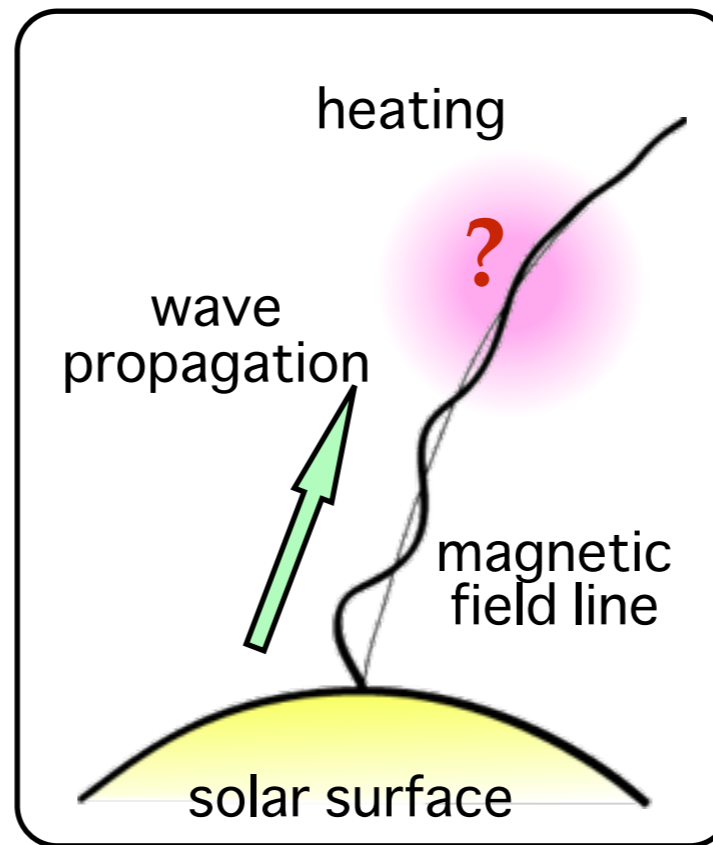
Observations of Alfvénic waves in the solar atmosphere

Transverse MHD waves ubiquitous in the solar atmosphere

(Okamoto+2007) Do they play an important role in the solar atmosphere?



Corona: $\sim 3-10$ km/s
Damping often observed

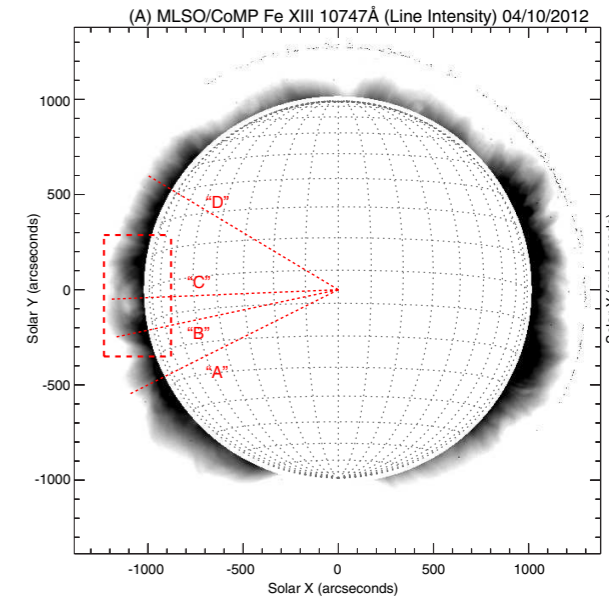
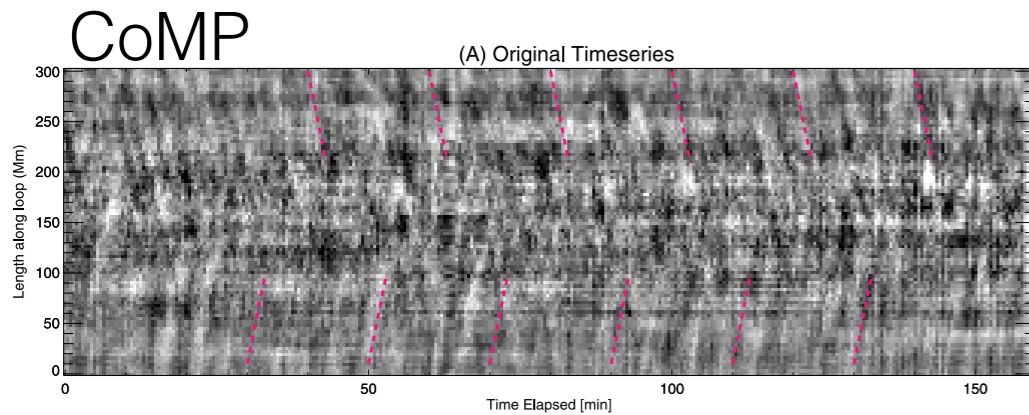


Damping can be caused by **resonant absorption (mode coupling)**: transverse waves convert into azimuthal waves

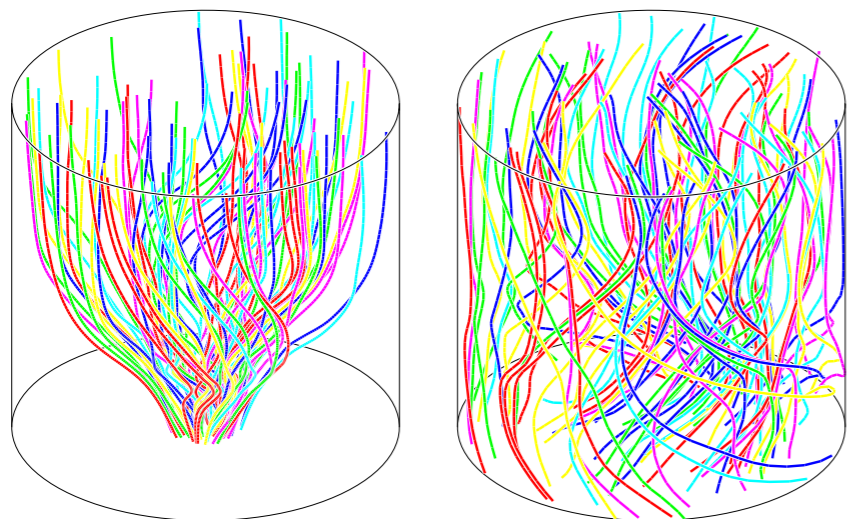
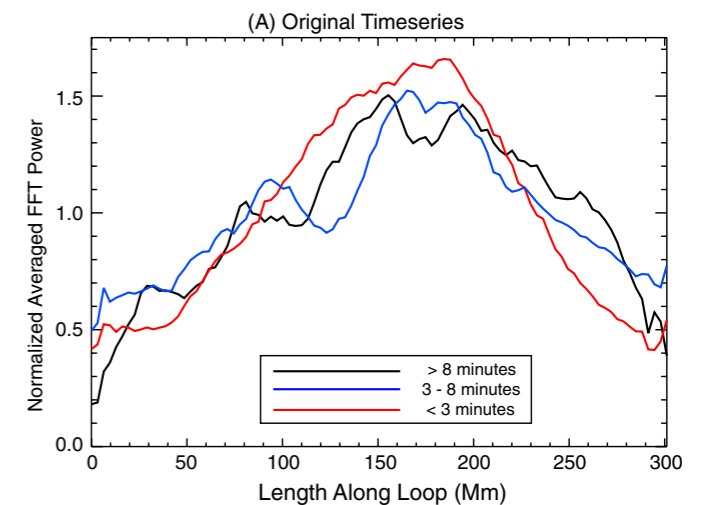
Chromosphere: ~ 20 km/s
(sufficient energy flux)

Hinode/SOT

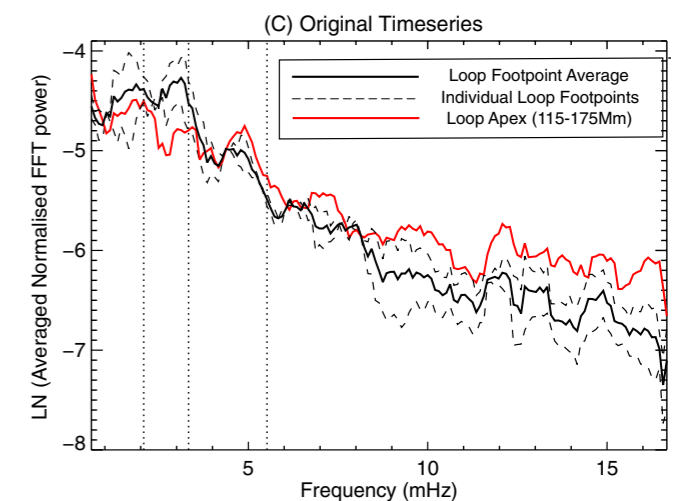
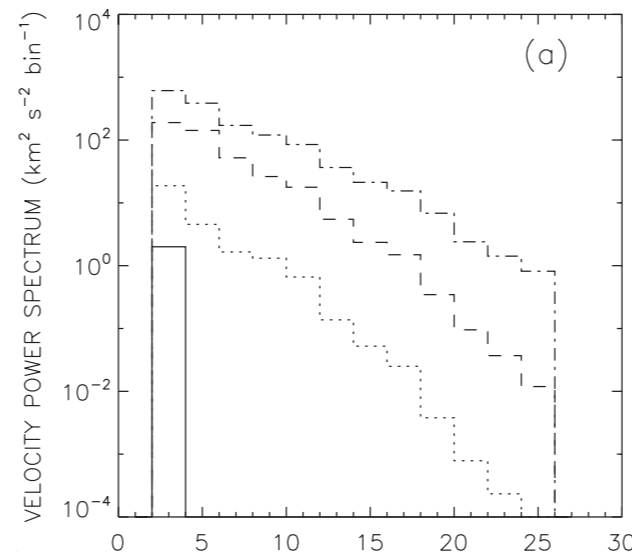
Alfvénic turbulence



- Large imbalance upward / downward wave energy flux. Increase of high frequency wave power at loop apex \rightarrow Alfvénic turbulence? (*De Moortel+ 2014, Tomczyk 2007, Tomczyk & McIntosh 2009*)
- Significant heating from Alfvénic turbulence (*Van Ballegooijen+ 2011, Matsumoto & Suzuki 2014*)
- Large non-thermal line widths may be hiding most of the wave power (*McIntosh & De Pontieu 2012*)



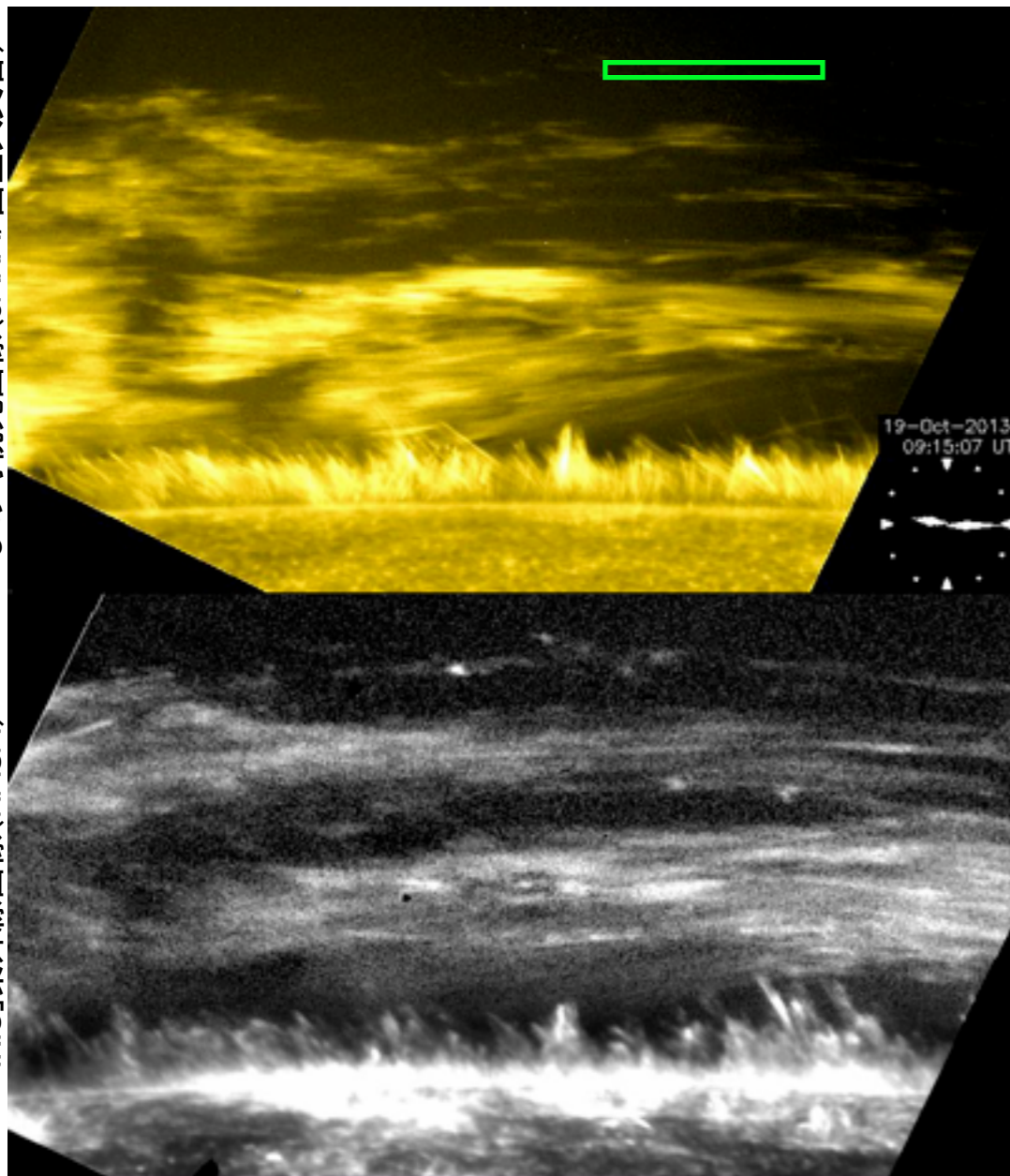
Van Ballegooijen et al. 2011



De Moortel et al. 2014

Signatures of resonant absorption

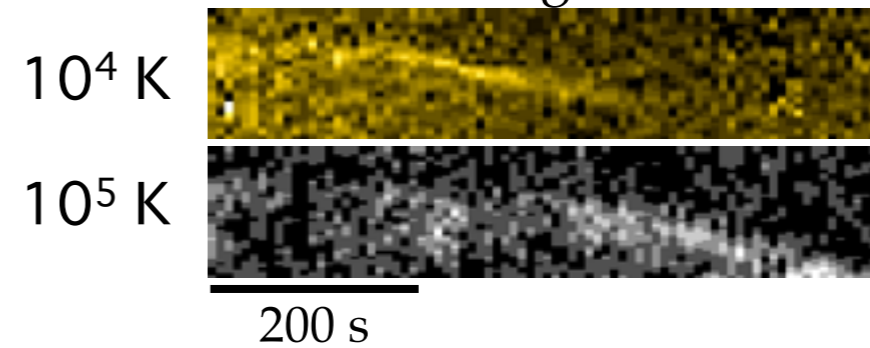
Hinode/SOT (Ca II, 10,000 K)



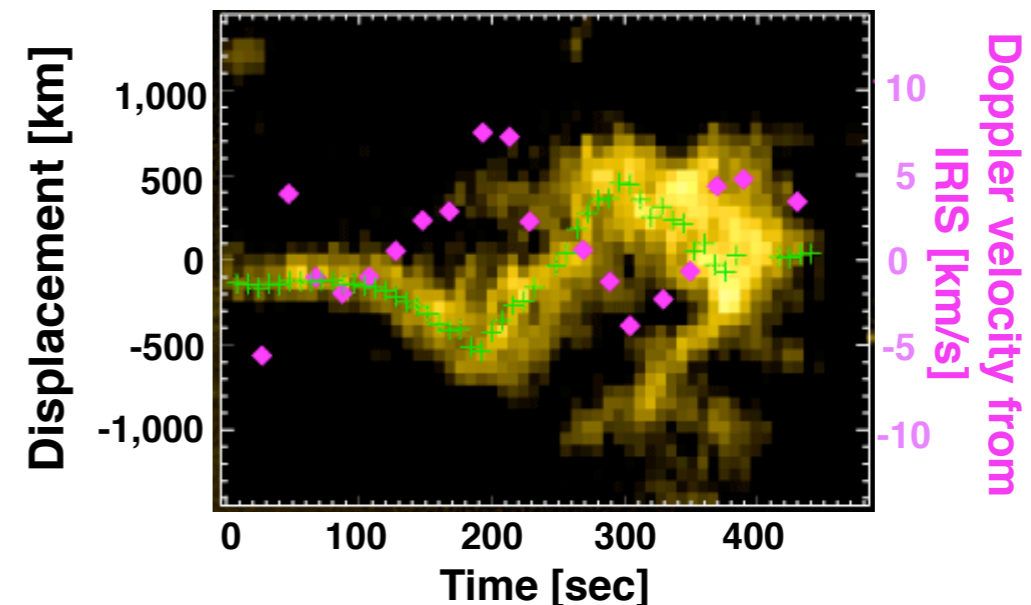
IRIS/SJI (Si IV, 100,000)

(Okamoto+2015, Antolin+2015)

x-t diagram



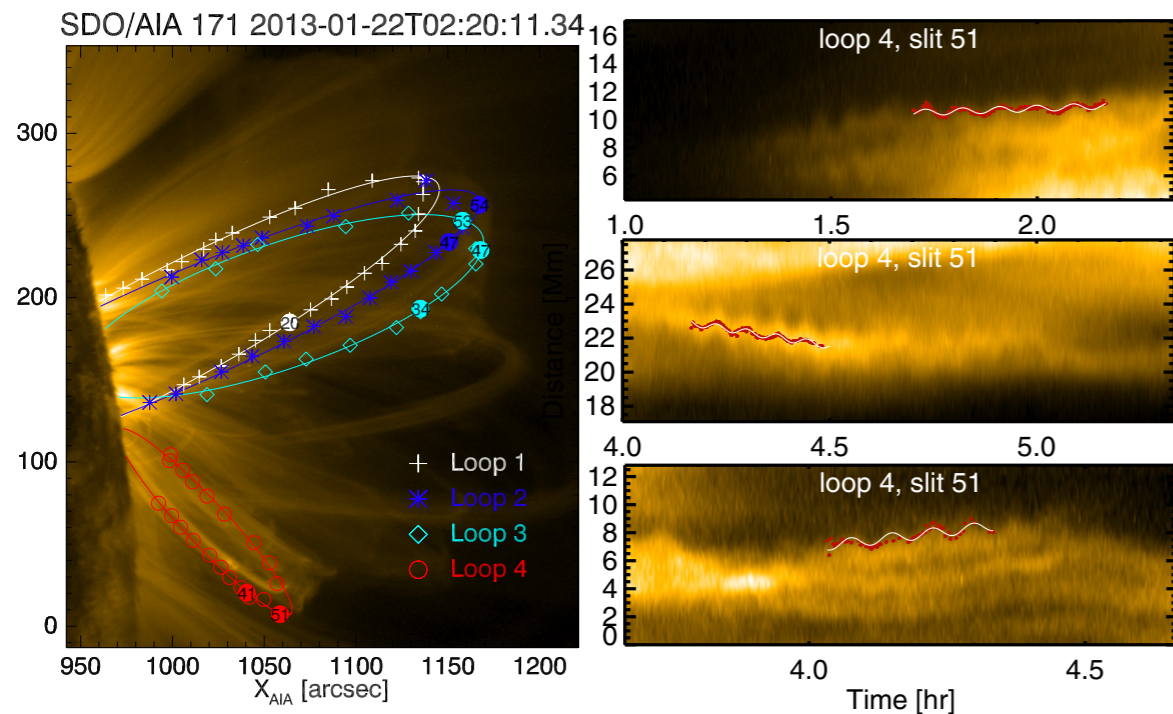
Motion of prominence plasma crossing the slit



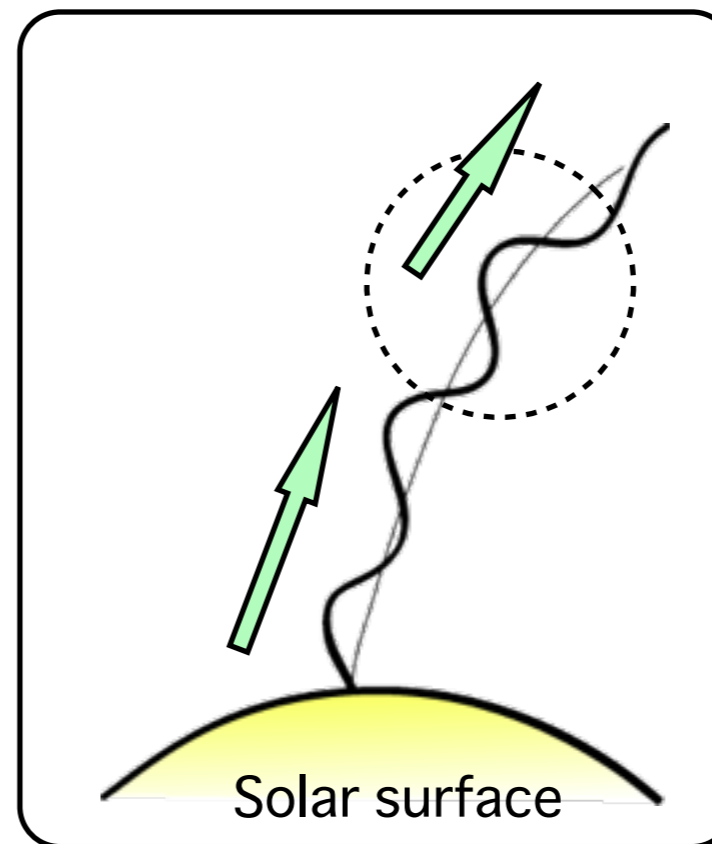
- Heating: Fading in cool line (10^4 K), subsequent appearance in hot line (10^5 K)
- POS motion out-of-phase with LOS velocity
- Thread-like structure
- ➔ Explained with 3D MHD transverse wave model: KHI + resonant absorption (current model)

Decay-less oscillations

*Nisticò+(2013), Anfinogentov+ (2013, 2015),
Goddard+(2015)*



Anfinogentov+ 2013



- Decayless
- Common
- Standing kink

Decay-less oscillations: damping+continuous low-amplitude harmonic driver?

Outline

Observations:

- Alfvénic waves are everywhere: damping & decayless
- Strong amplitudes in chromosphere > Chromospheric heating, small POS & LOS v amplitudes in corona <? coronal heating
- Wave energy expected to be in azimuthal motions -> non-thermal line widths
- Alfvénic turbulence in corona?

Need for determination of observational (imaging + spectroscopic) signatures of transverse MHD waves and define their role in the solar atmosphere

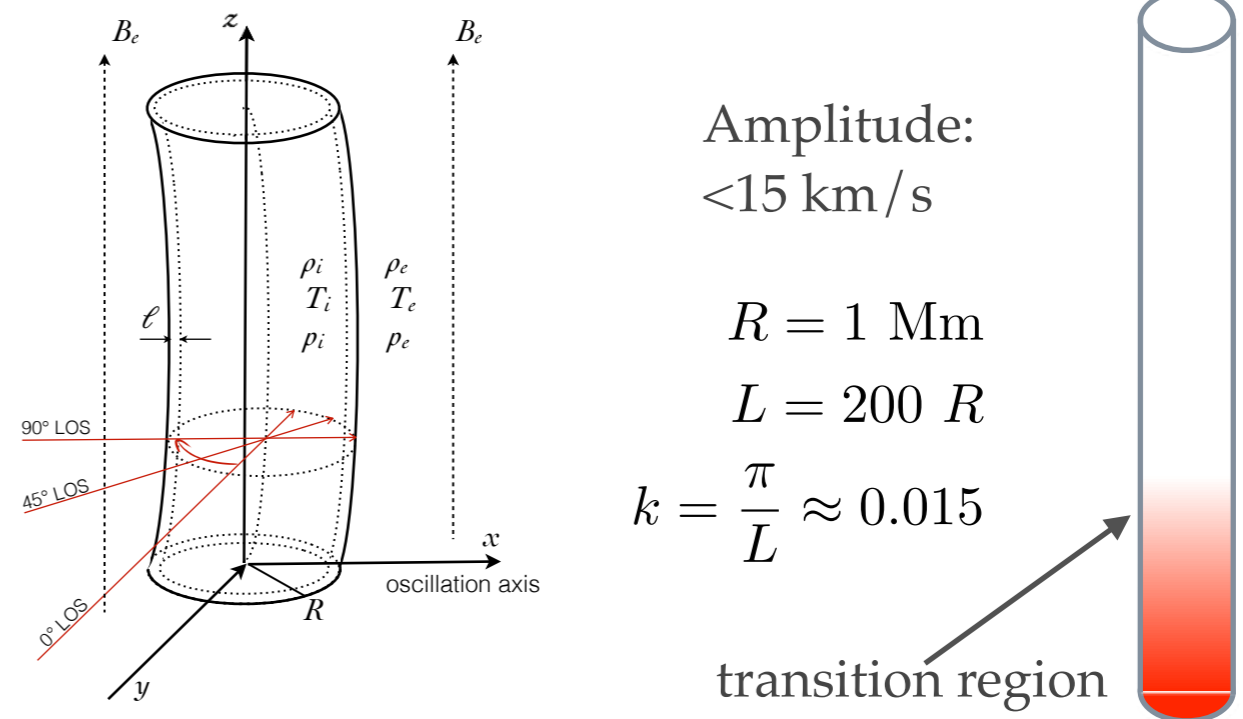
Present study:

- Modelling of transverse MHD waves (kink)
 - Prominences
 - Corona
 - Spicules

Numerical model

- 3D MHD simulations of a flux tube oscillating with the kink mode. CIP-MOCCT code (Kudoh et al. 1999) with constant resistivity and viscosity
- Grid (x,y,z): 1/4 tube = (512, 256, 100) - (1024, 512, 100) $S, R \approx 10^4 - 10^7$ spicule model
- Initial condition: sinusoidal velocity perturbation in x-direction

parameters	coronal loop	prominence	spicule
$\frac{T_i}{T_e}, T_i$ [K]	1/3, $T_i = 10^6$	1/100, $T_i = 10^4$	1/100, $T_i = 10^4$
$\frac{\rho_i}{\rho_e}, \rho_i$ [cm ⁻³]	3, $\rho_i = 3 \times 10^9$	10, $\rho_i = 10^{10}$	50, $\rho_i = 6 \times 10^{10}$
B [G]	22.8 G	18.6 G	14.5 G
c_k [km/s]	1574	776	255
$P \approx \frac{2L}{c_k}$ [s]	525	256	245
β_i	0.02	0.001	0.01
ℓ [R]	0.2-0.8	0.4	0.4



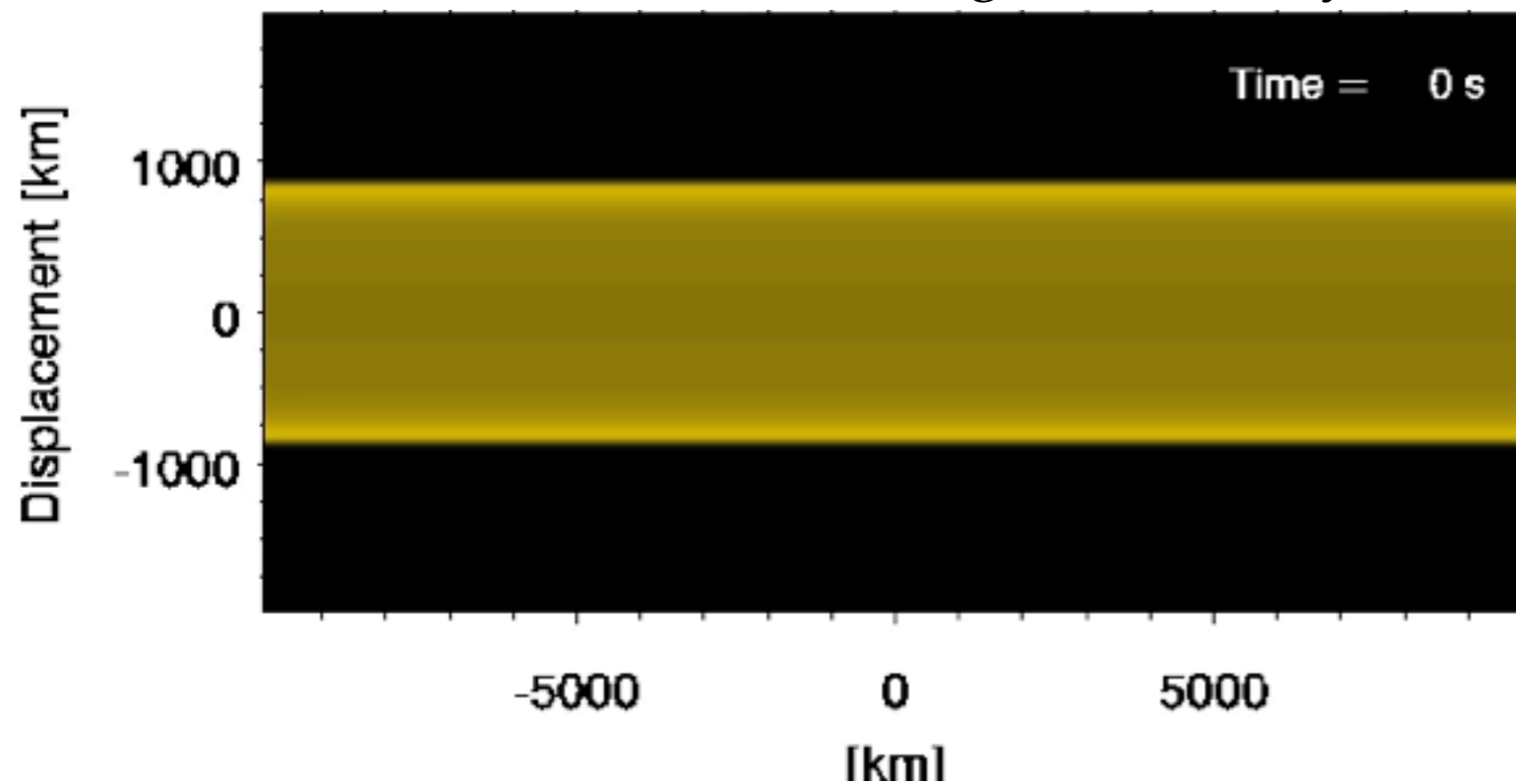
Forward modelling

- Optically thin: **FoMo** (Van Doorselaere+ 2016, Antolin & Van Doorselaere 2013)
- Optically thick: **RH** (Uitenbroek 2011)

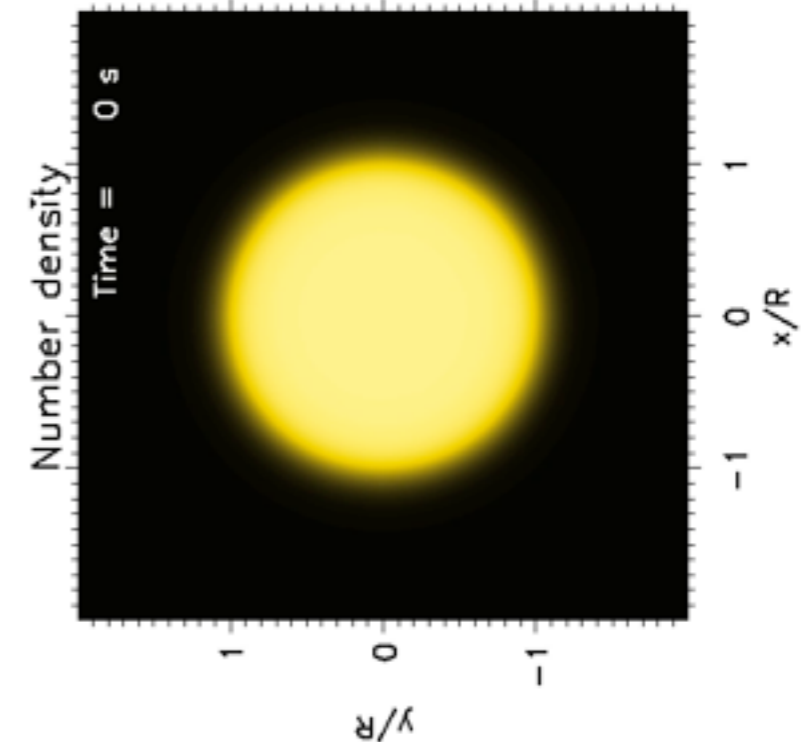
<https://wiki.esat.kuleuven.be/FoMo>

Numerical simulation

Prominence thread - Mg II k intensity

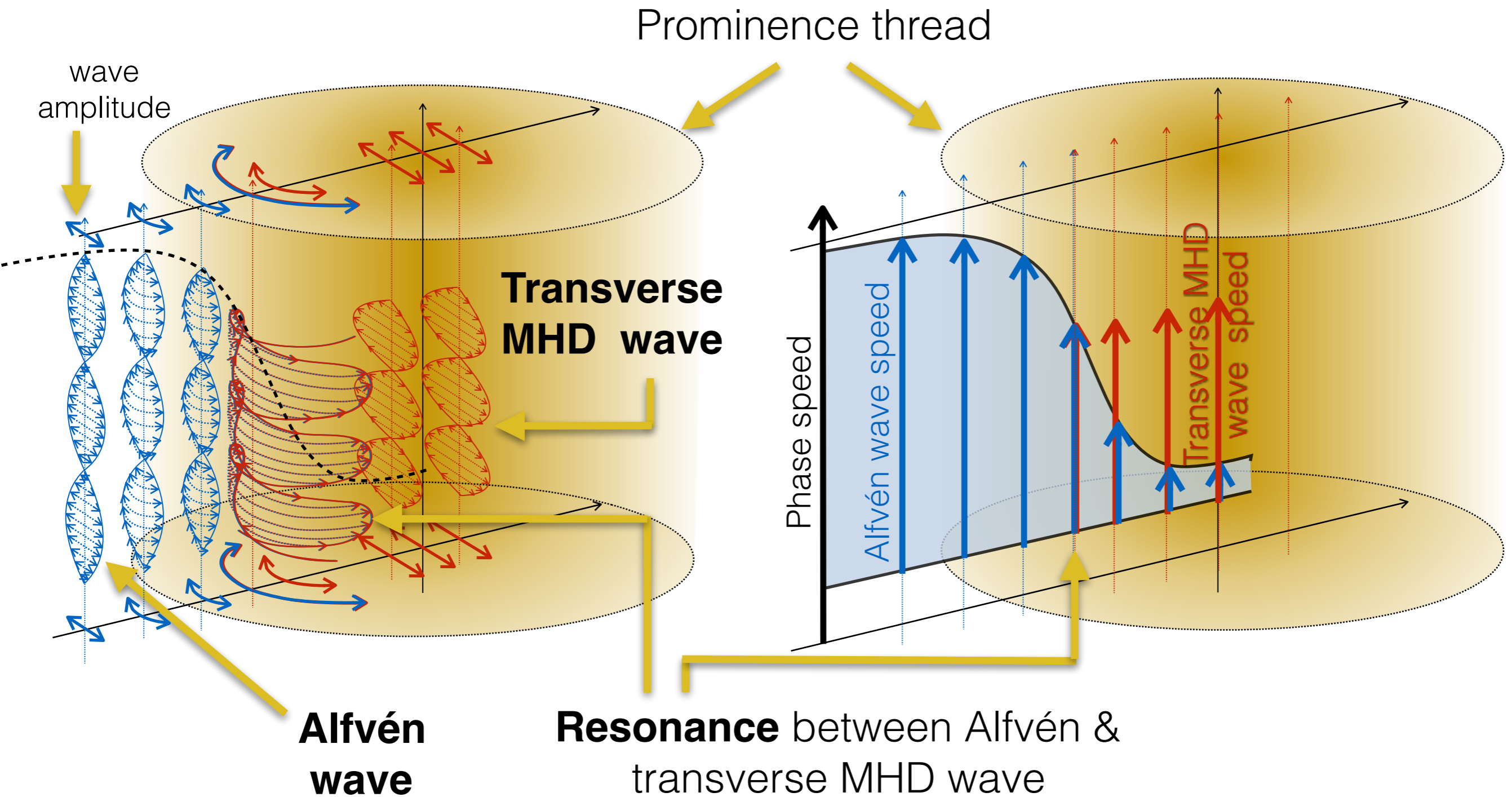


Cross-section



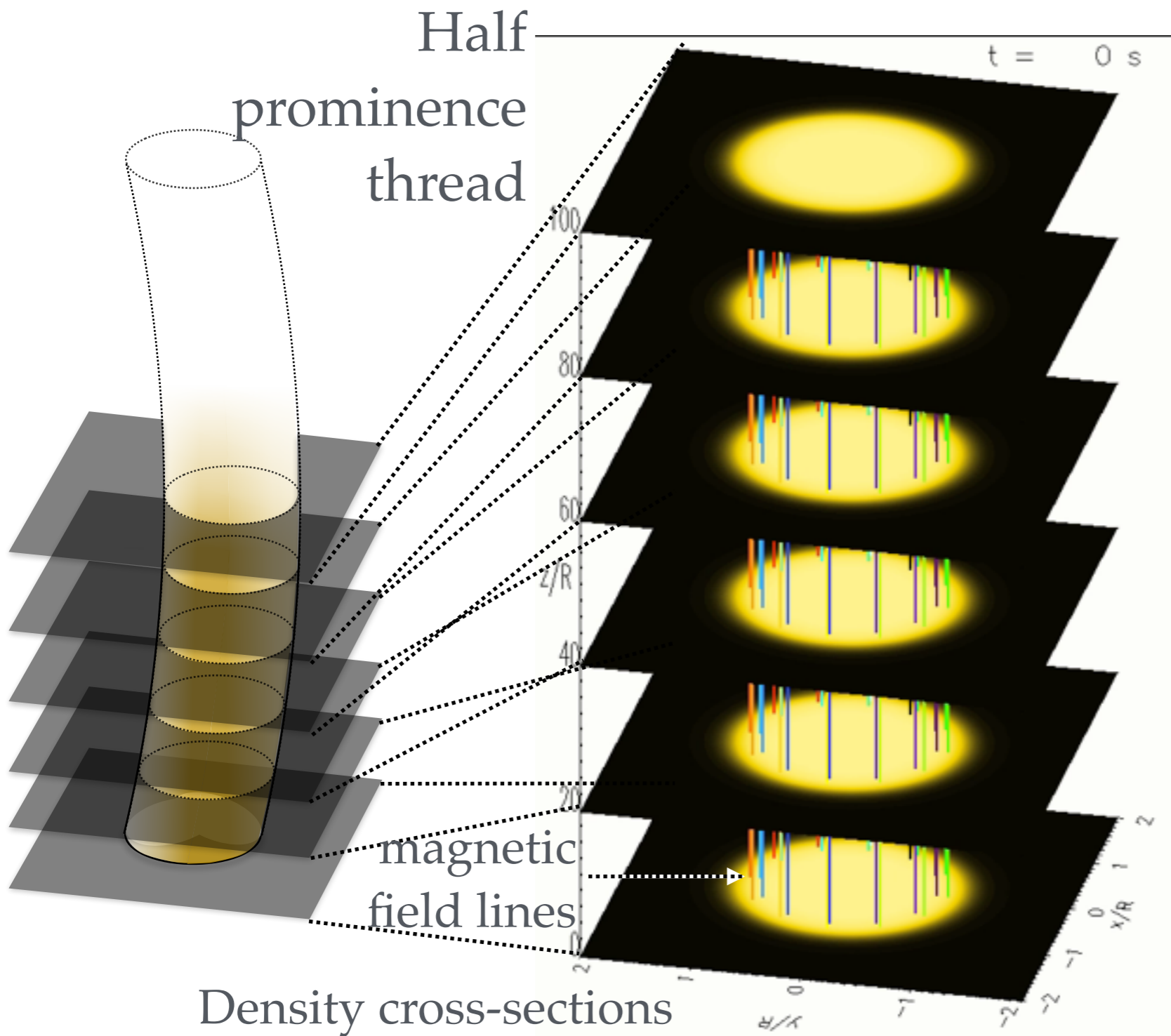
2 kinds of motion: transverse+azimuthal
Resonant absorption transfers energy from a transverse (global) wave to azimuthal (local) Alfvén waves near the boundary

Resonant absorption

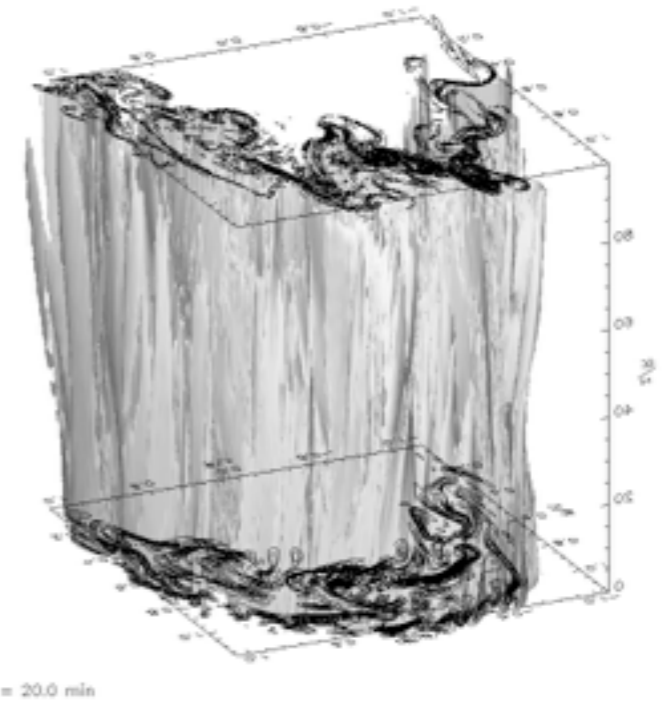


Ionson 1978, Hollweg+ 1990, Sakurai+ 1991, Goossens+ 1992, 2002, 2012; Van Doorselaere+ 2004, Arregui+ 2008, 2011; Verth+ 2010, Soler+ 2010, 2012; Pascoe+ 2010, 2012

Two mechanisms combined



Iso-contours of current density



Vortices break-up into turbulence, producing twisted current sheets

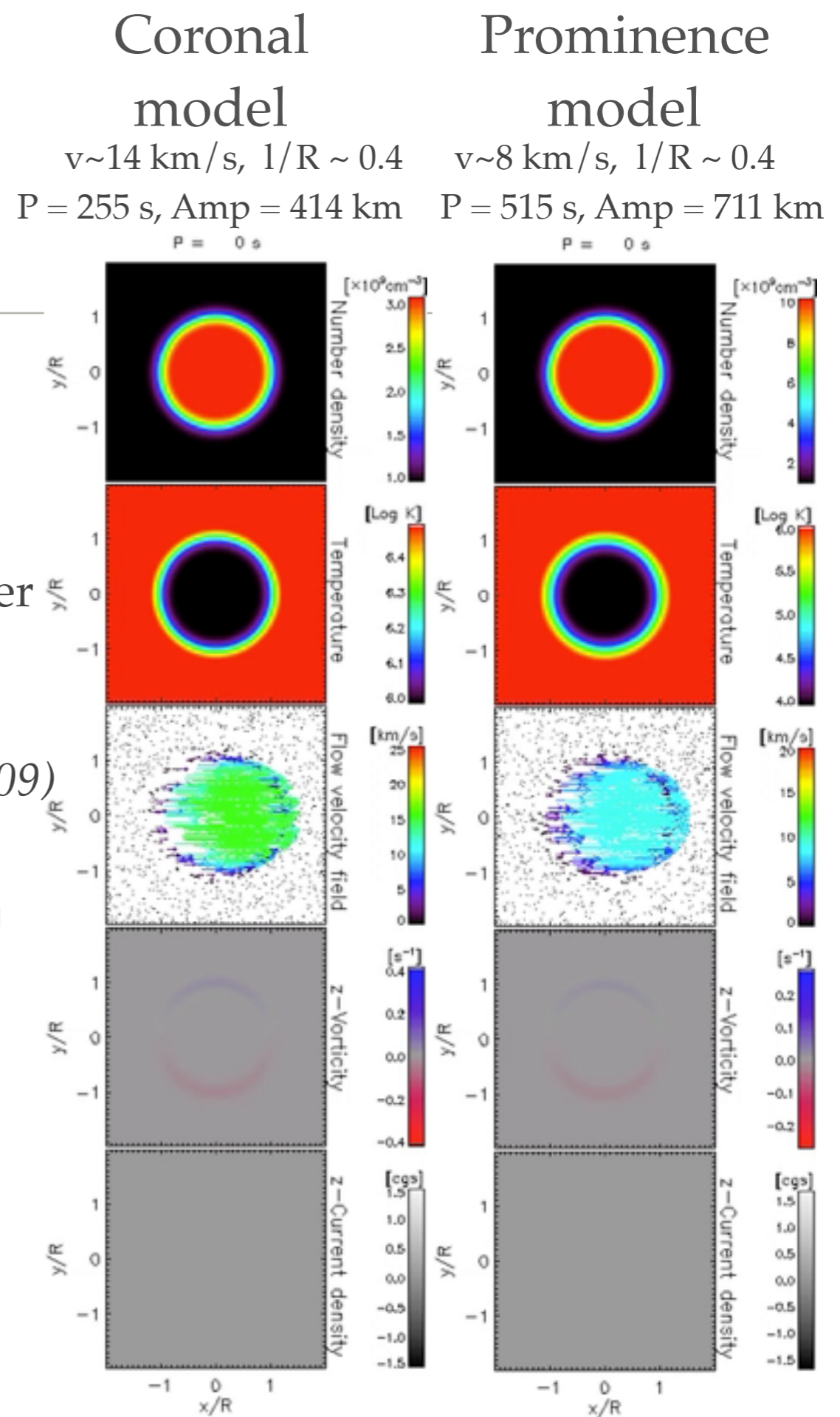
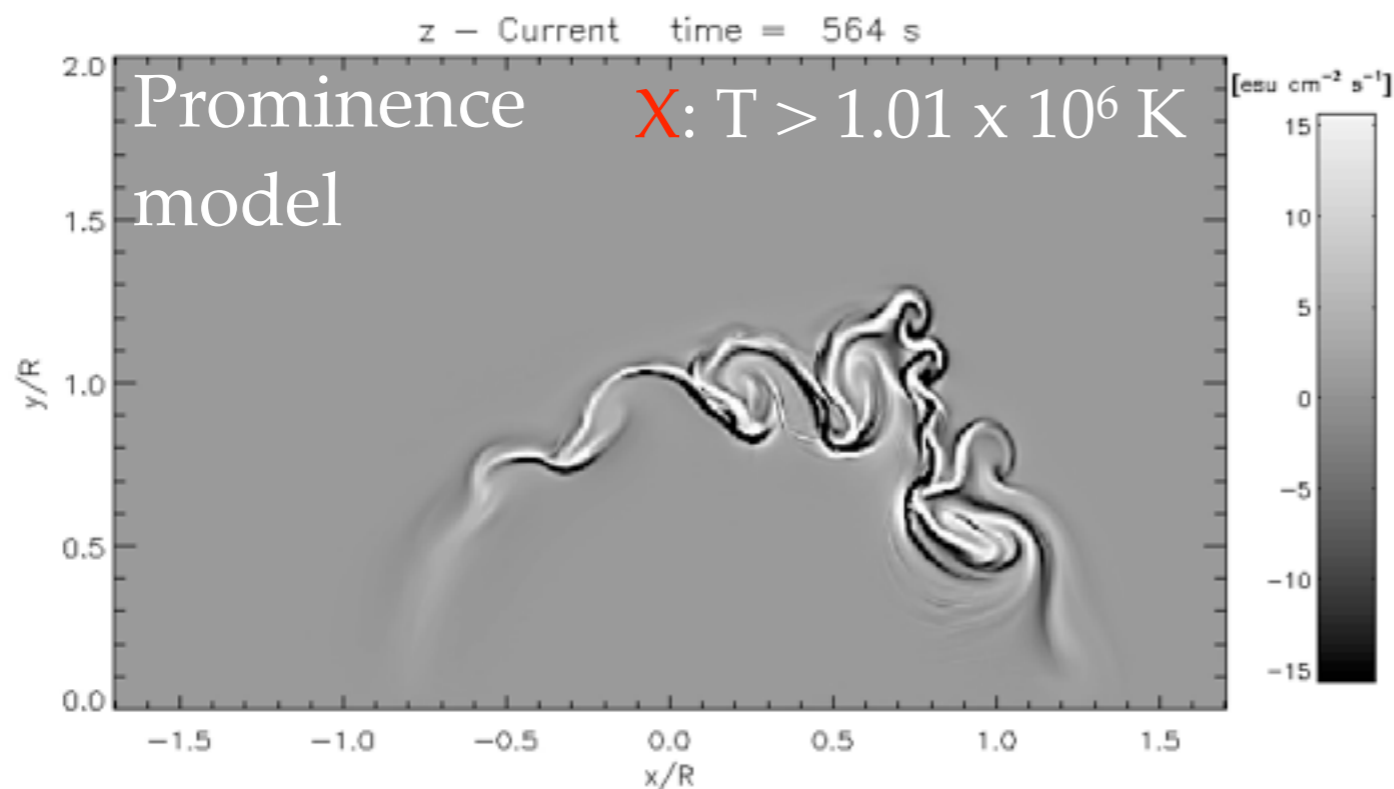
Model: Uchimoto+. 1991; Karpen+ 1993, Ofman+ 1994; Ziegler & Ulmschneider 1997; Terradas+ 2008, Soler+ 2010

Resonant absorption & onset of K-H instability

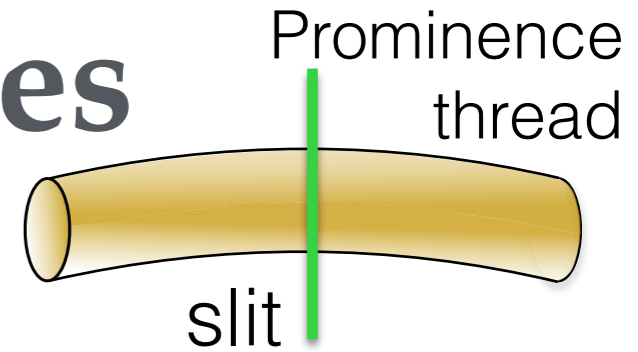
- Onset of instability (Zaqarashvili+ 2015)

$$\frac{v_0}{v_{A_i}} > \frac{\pi}{\sqrt{|m|-1}} \frac{R}{L} \sqrt{\left(1 + \left(\frac{B_e}{B_{z_i}}\right)^2\right) \left(1 - \frac{|m|-1}{\frac{\rho_i}{\rho_e} + |m|}\right)}$$

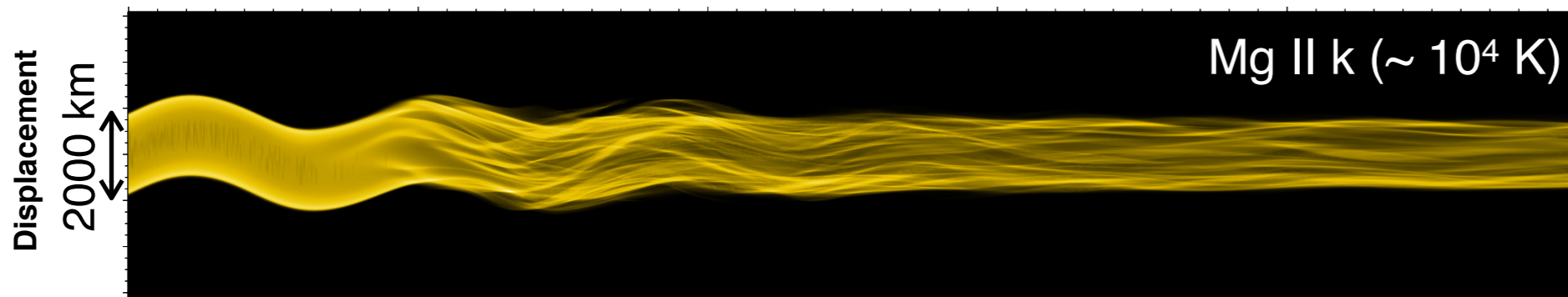
- KHI vortices obtain momentum from resonant layer
- Non-uniform boundary layer widens, mixing of plasma (Fujimoto & Terasawa 1994)
- Multiple vortices & current sheets (Ofman 1994, 2009)



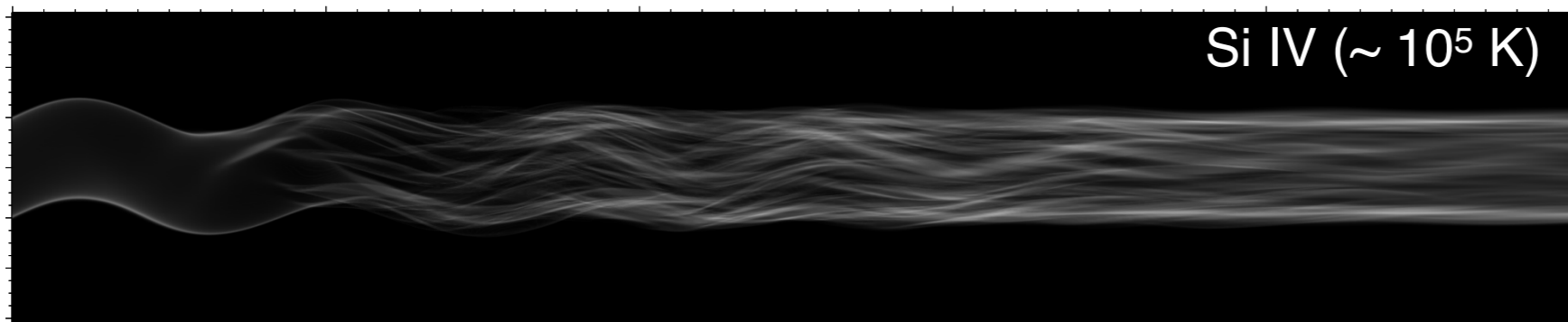
Observational signatures



45° LOS plane

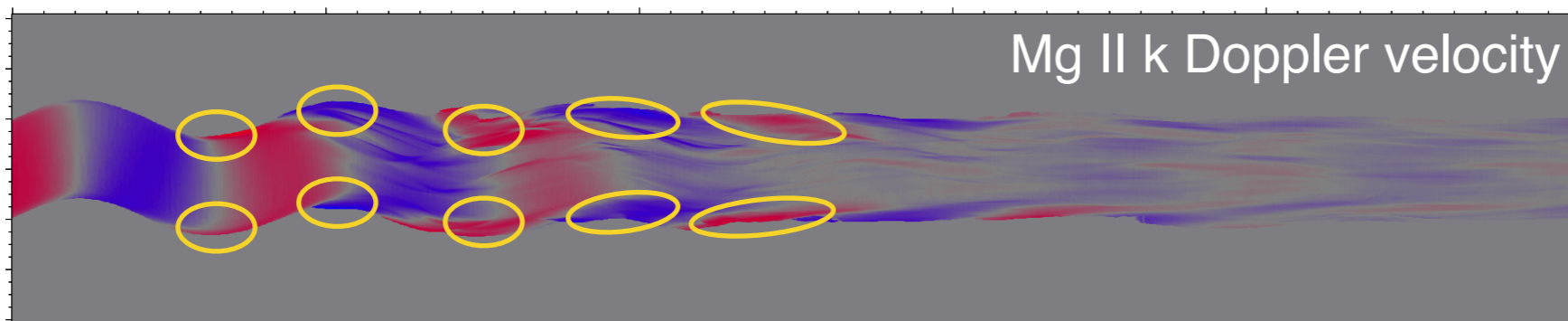


fading & thinning in chromospheric lines



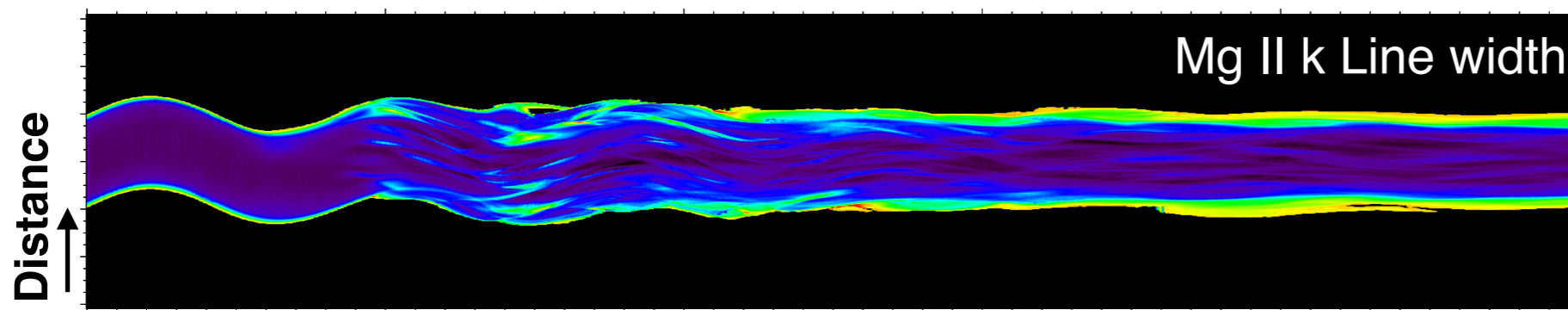
appears in TR lines, broadened

→ mixing & heating



LOS velocity out-of-phase with POS motion

→ RA+phase mixing

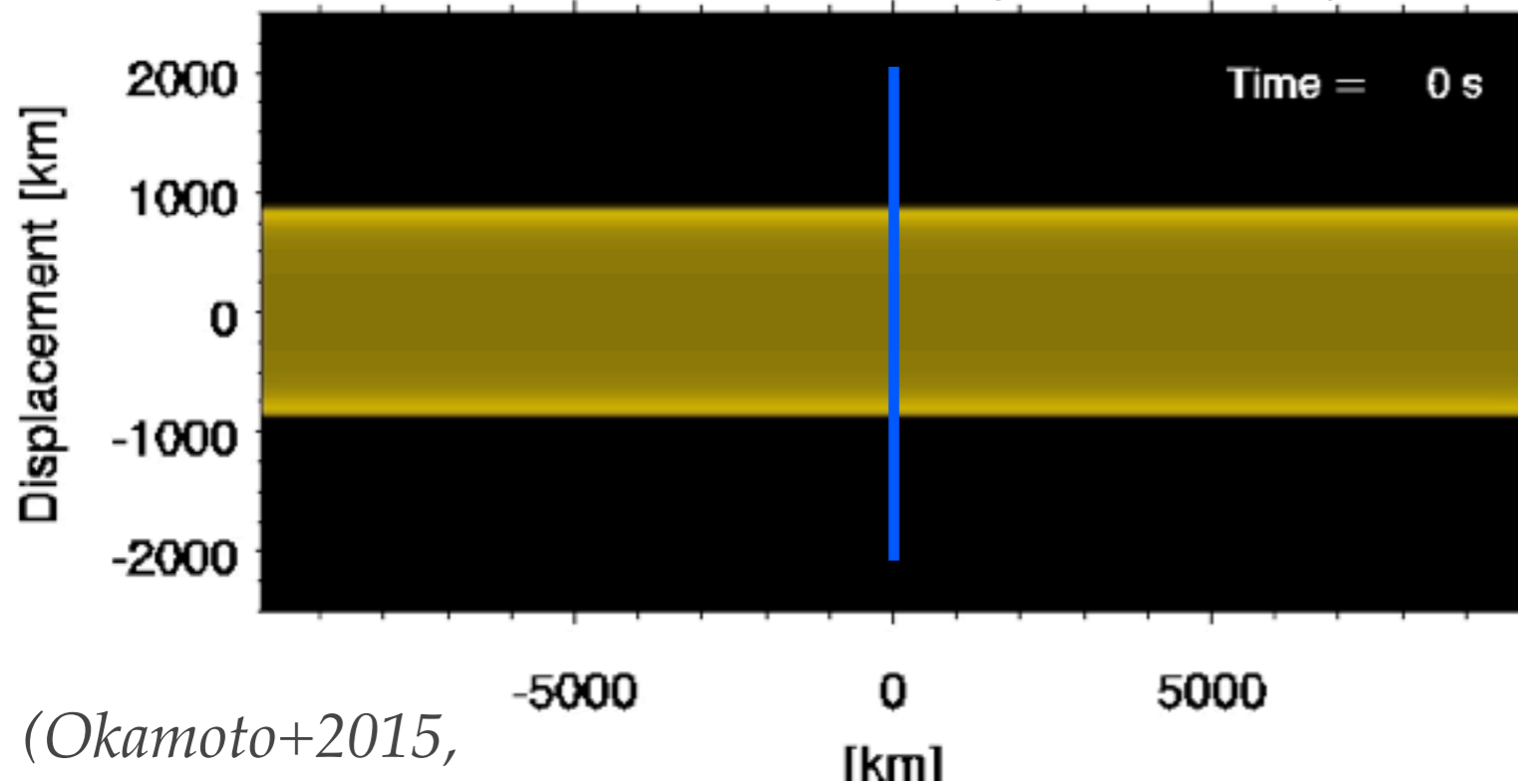


Line broadening at edges

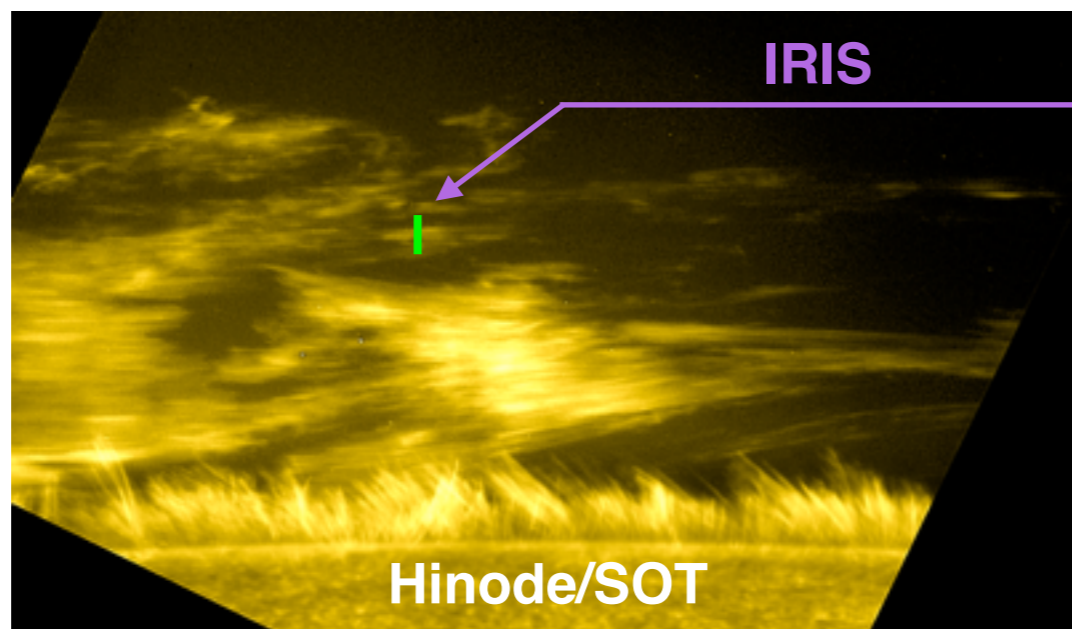
Time →
← 200 s

Comparing with Hinode & IRIS observations

Prominence thread - Mg II k intensity



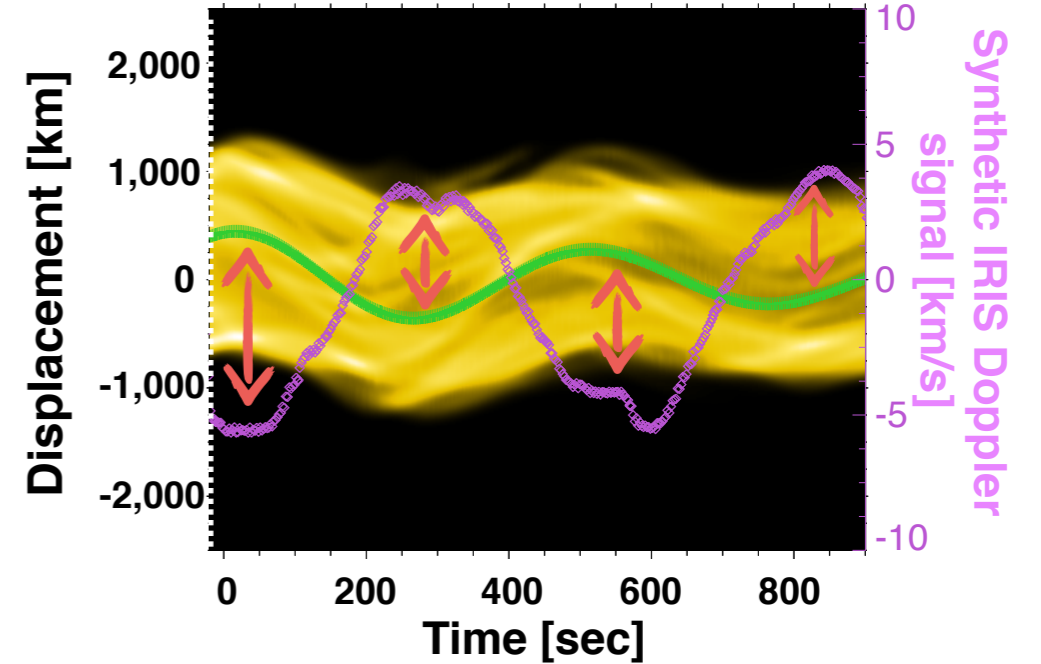
(Okamoto+2015, Antolin+2015)



KHI takes up the resonant dynamics to the observable scale: **match**

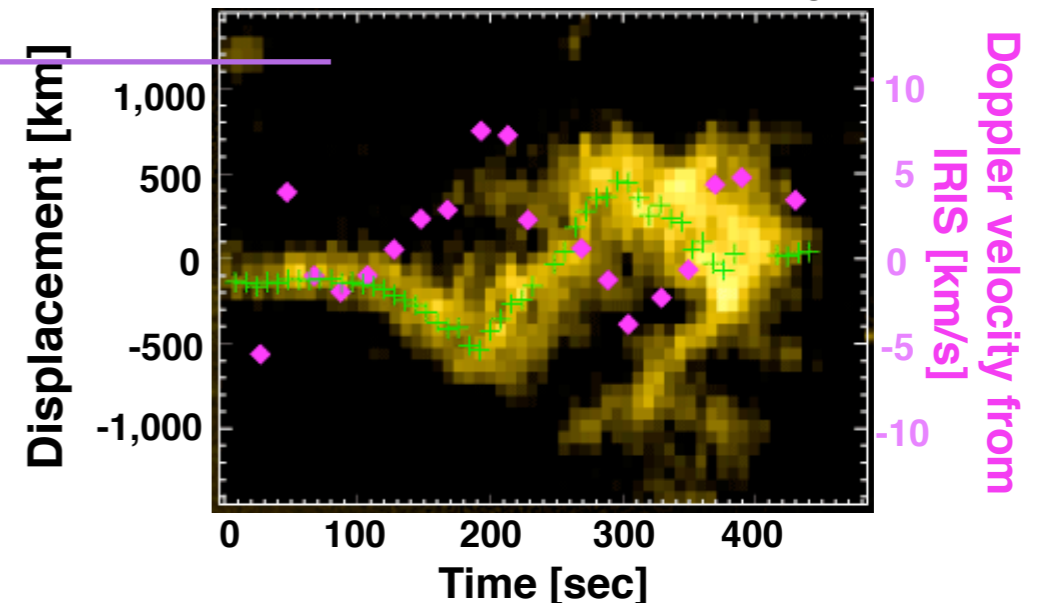
Numerical results

Motion of prominence plasma crossing the slit



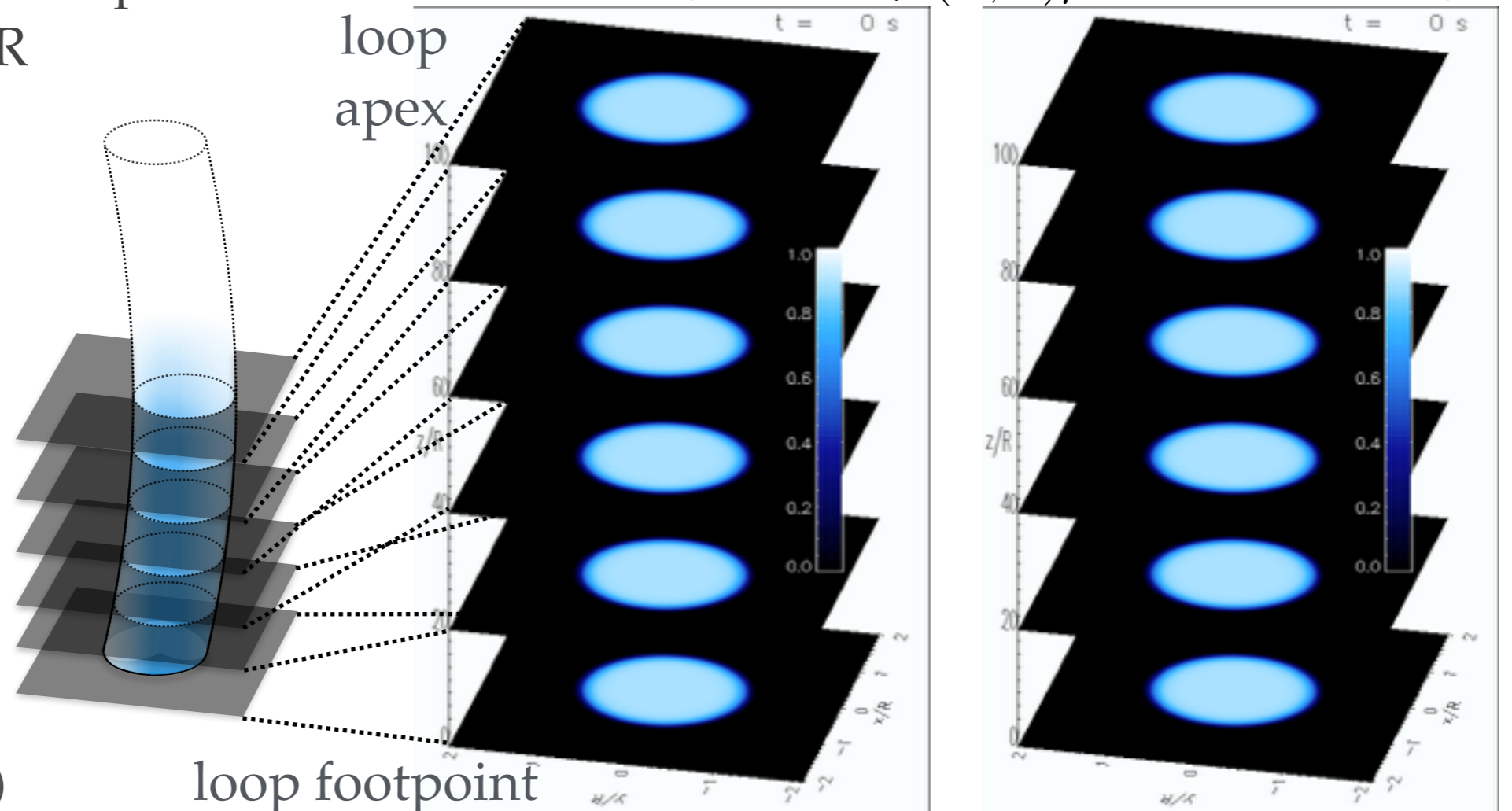
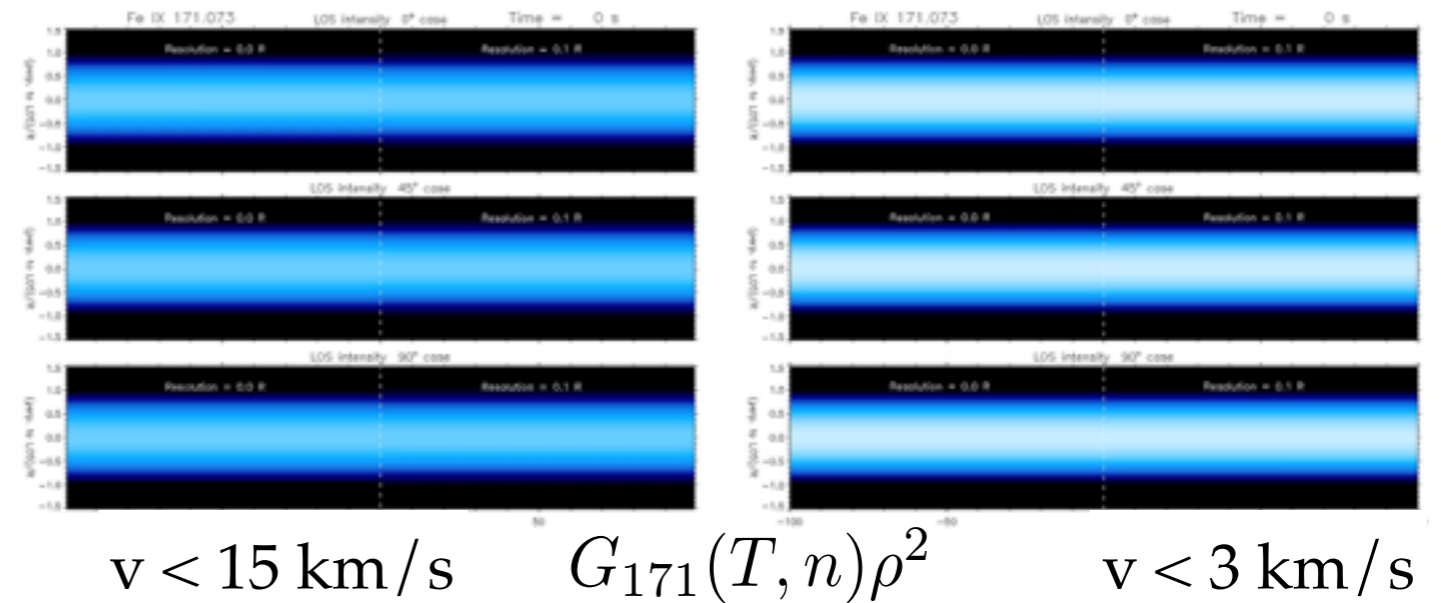
Observational results

Motion of prominence plasma crossing the slit



Strand-like structure in the corona

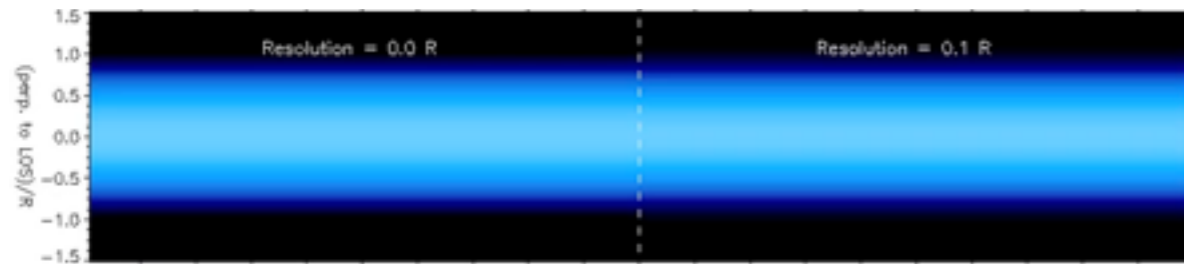
- Roll-ups (eddies) along the loop
→ strand-like structure in intensity
- Detection is strongly dependent on spatial resolution
- Lifetime for 1 strand ~ 1 period.
Widths: $0.01 R - 0.5 R$



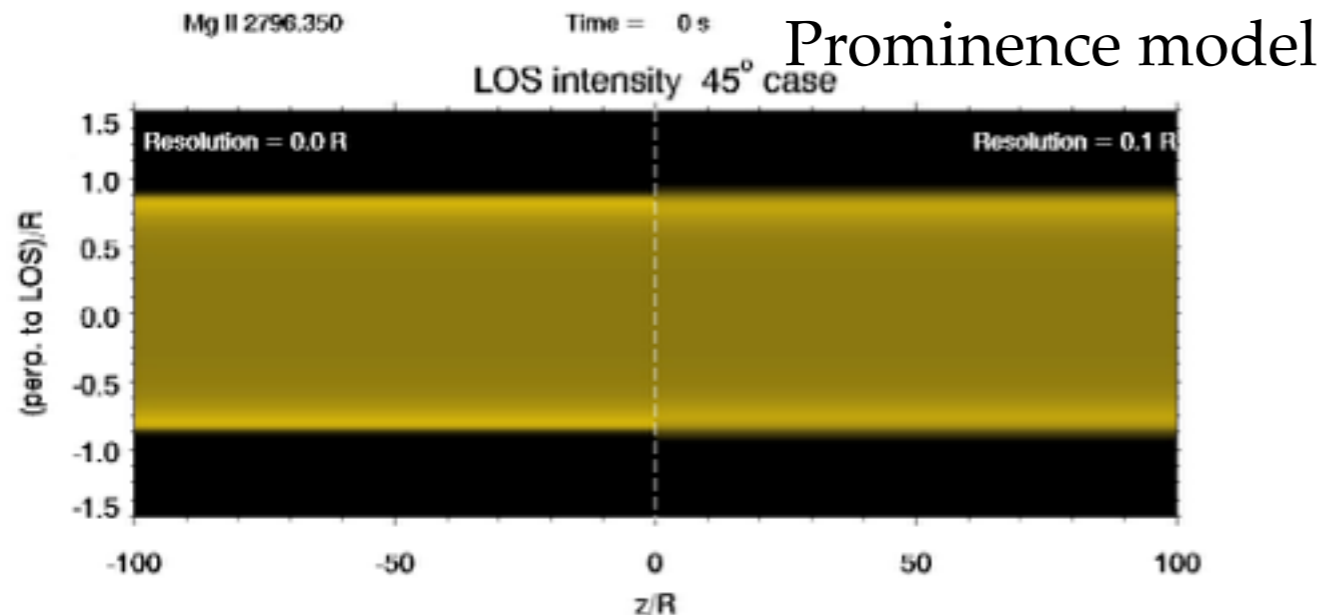
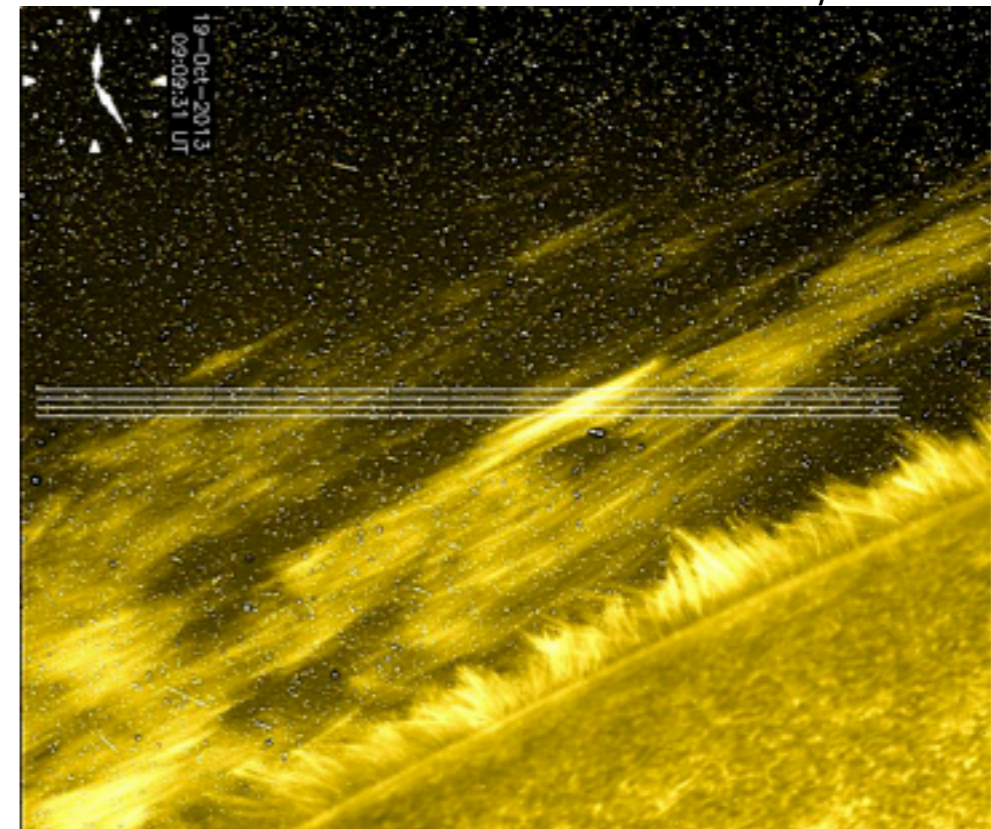
Antolin, Yokoyama &
Van Doorselaere (2014)

Strand-like structure

$v < 15 \text{ km/s}$, $1/R \sim 0.4$ Coronal model

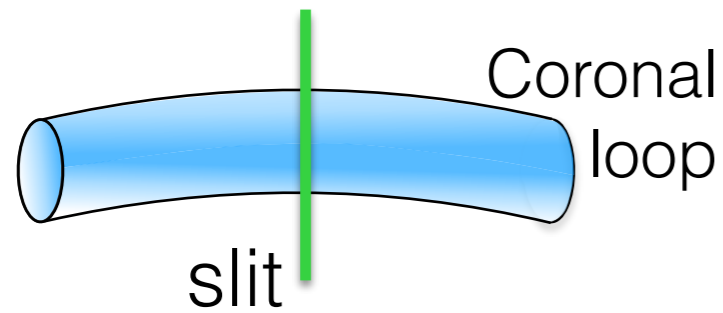


Hinode/SOT



- Roll-ups (eddies) along the loop + line-of-sight effects
 - ➔ strand-like or thread-like structure in intensity images
 - ➔ KHI vortices \longleftrightarrow part of prominence threads?
- Lifetime for 1 strand ~ 1 period. Widths: 0.01 R - 0.5 R
- Apparent crossing of strands/threads

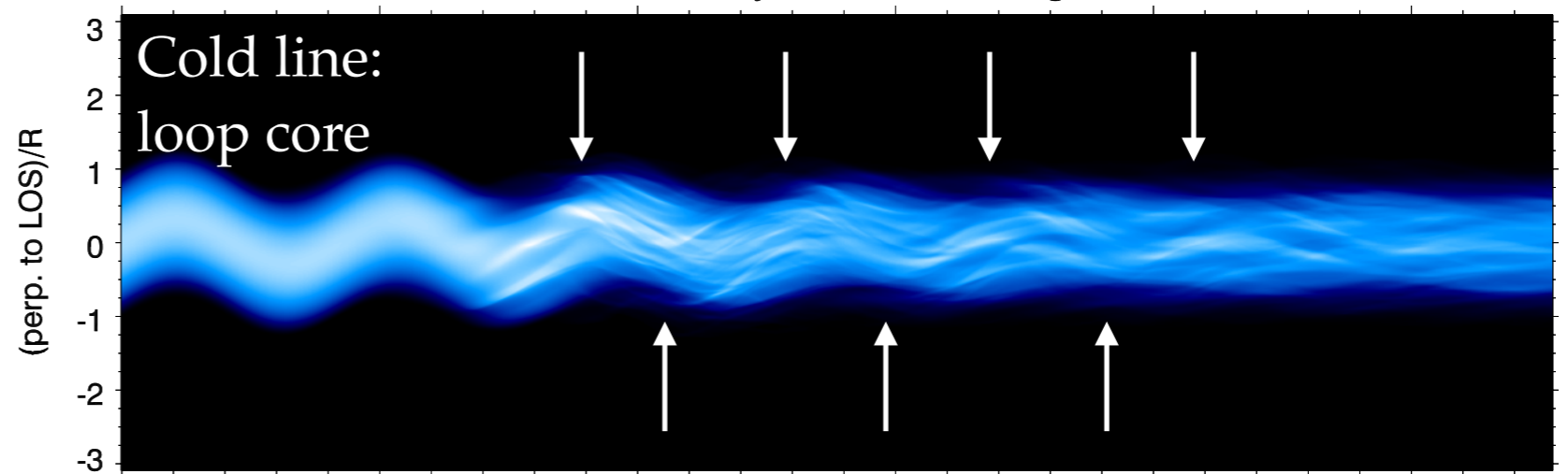
Intensity variation



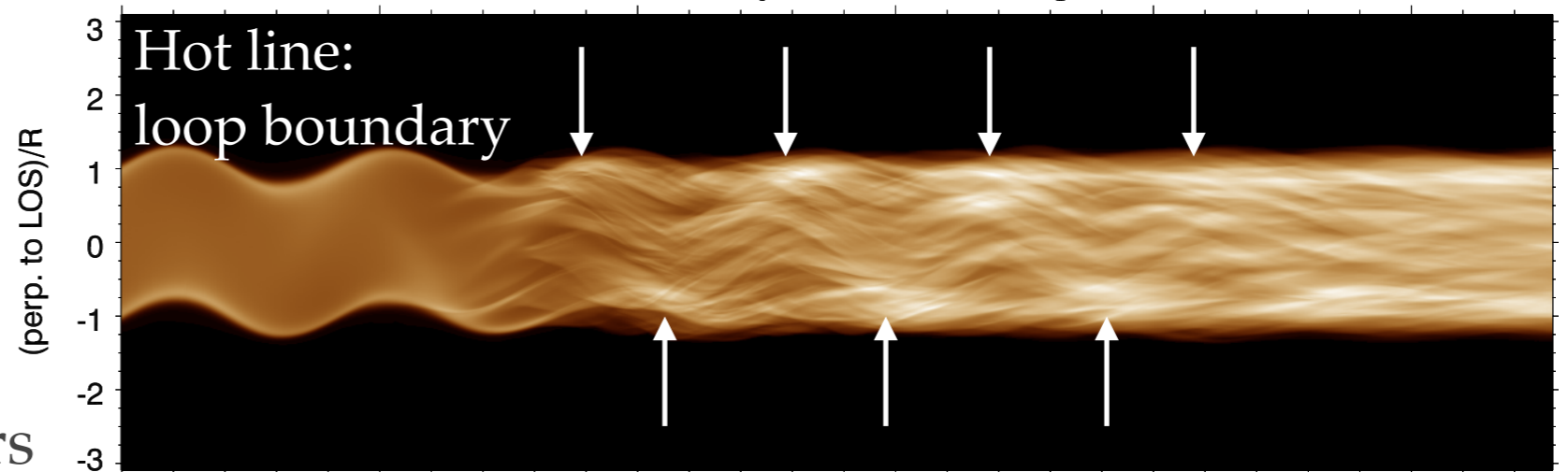
- Brightening when KHI with strand-like pattern
- Impulsive character at times of maximum displacement
- Gradual dimming in 171, intensity enhancement in 193
- Thinning in 171, enlargement in 193
- Observed damping appears different in 171 & 193

Antolin, De Moortel, Van Doorselaere, Yokoyama (2016, submitted)

Fe IX 171 Intensity - Variation along 45° LOS

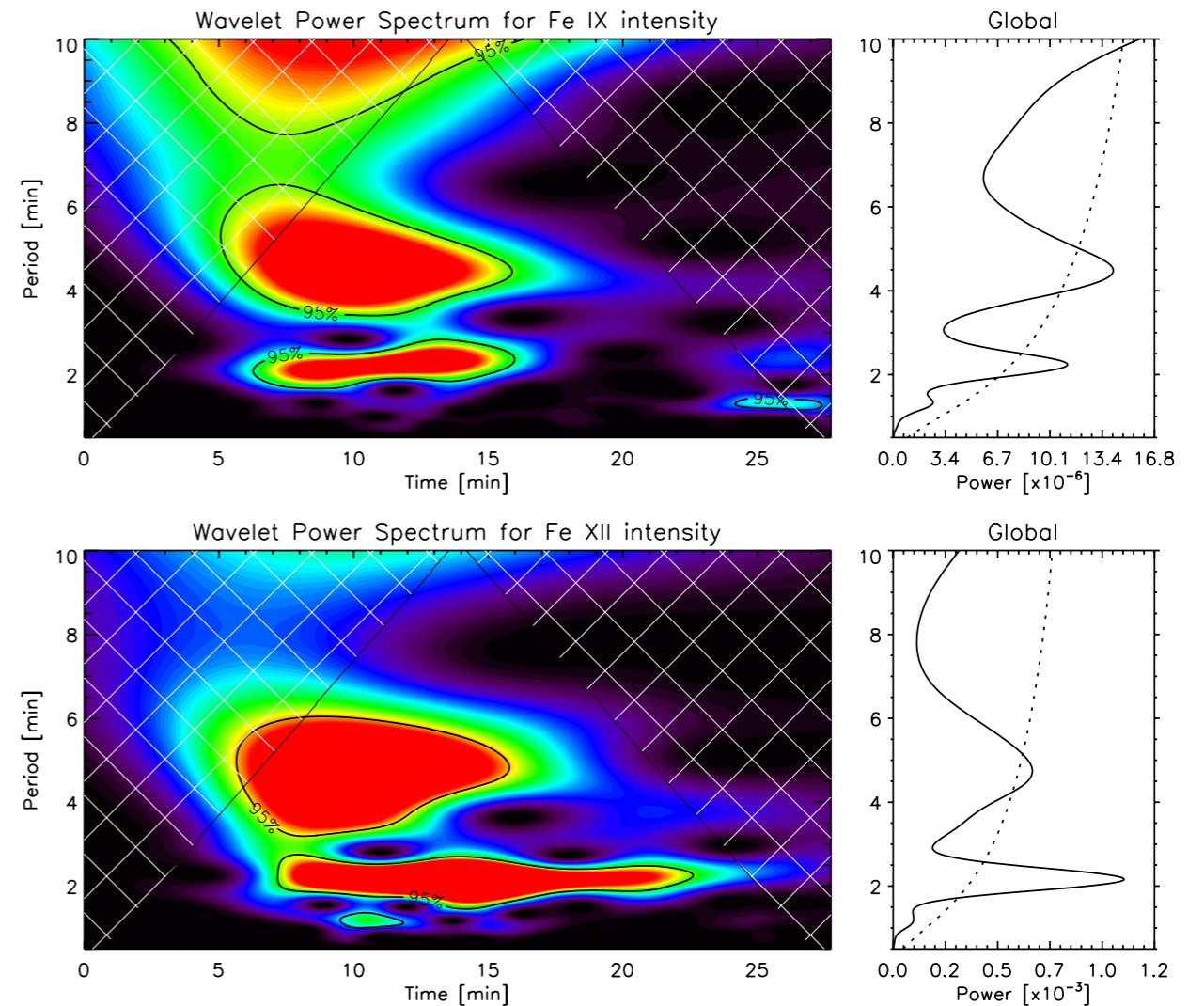
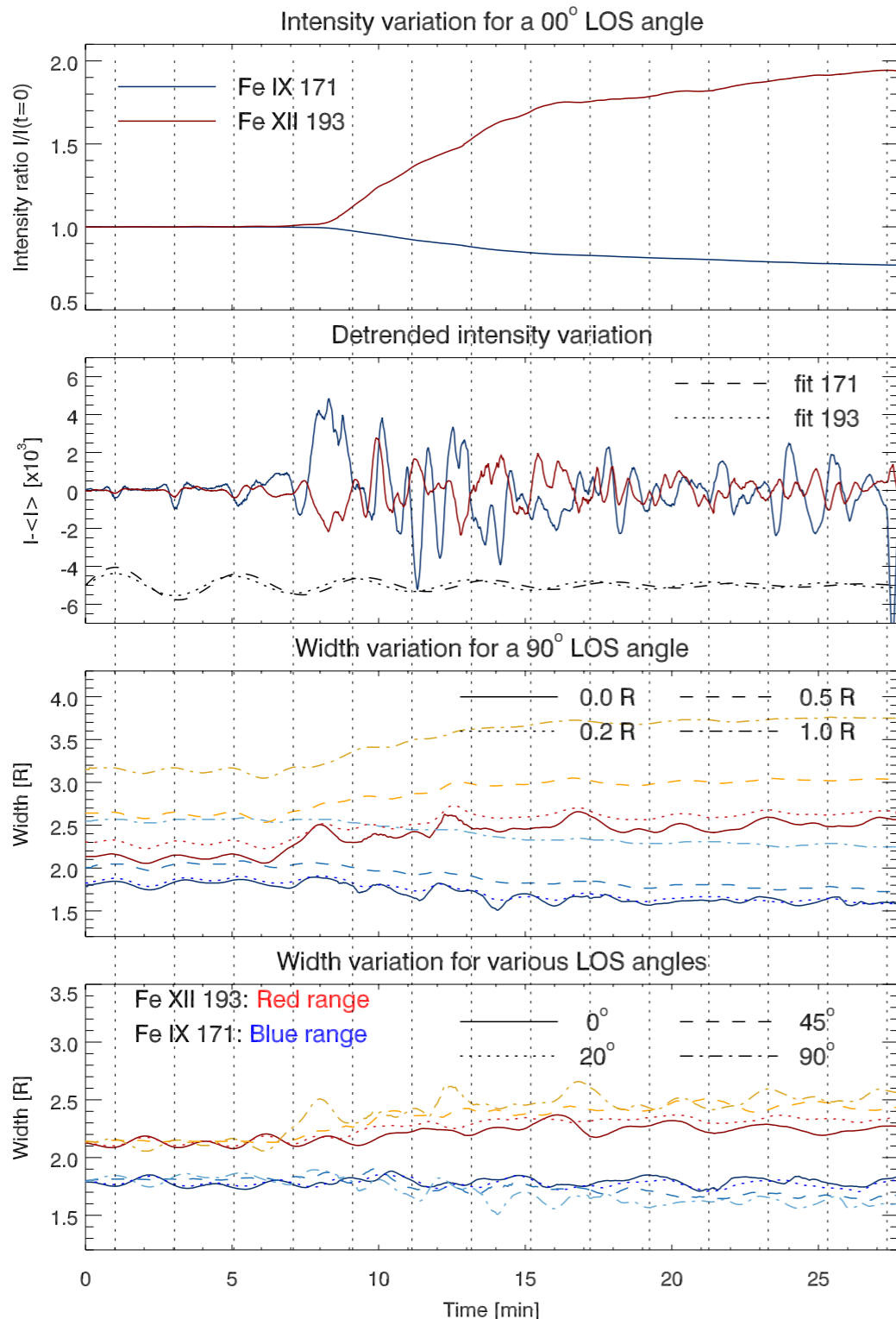


Fe XII 193 Intensity - Variation along 45° LOS



➔ Hotter channel more sensitive to vortex dynamics (and to the resonant flow)

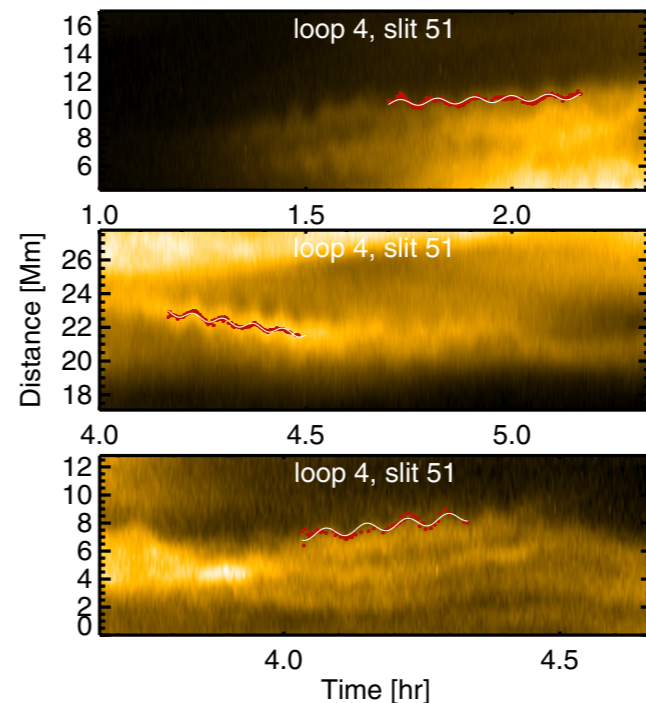
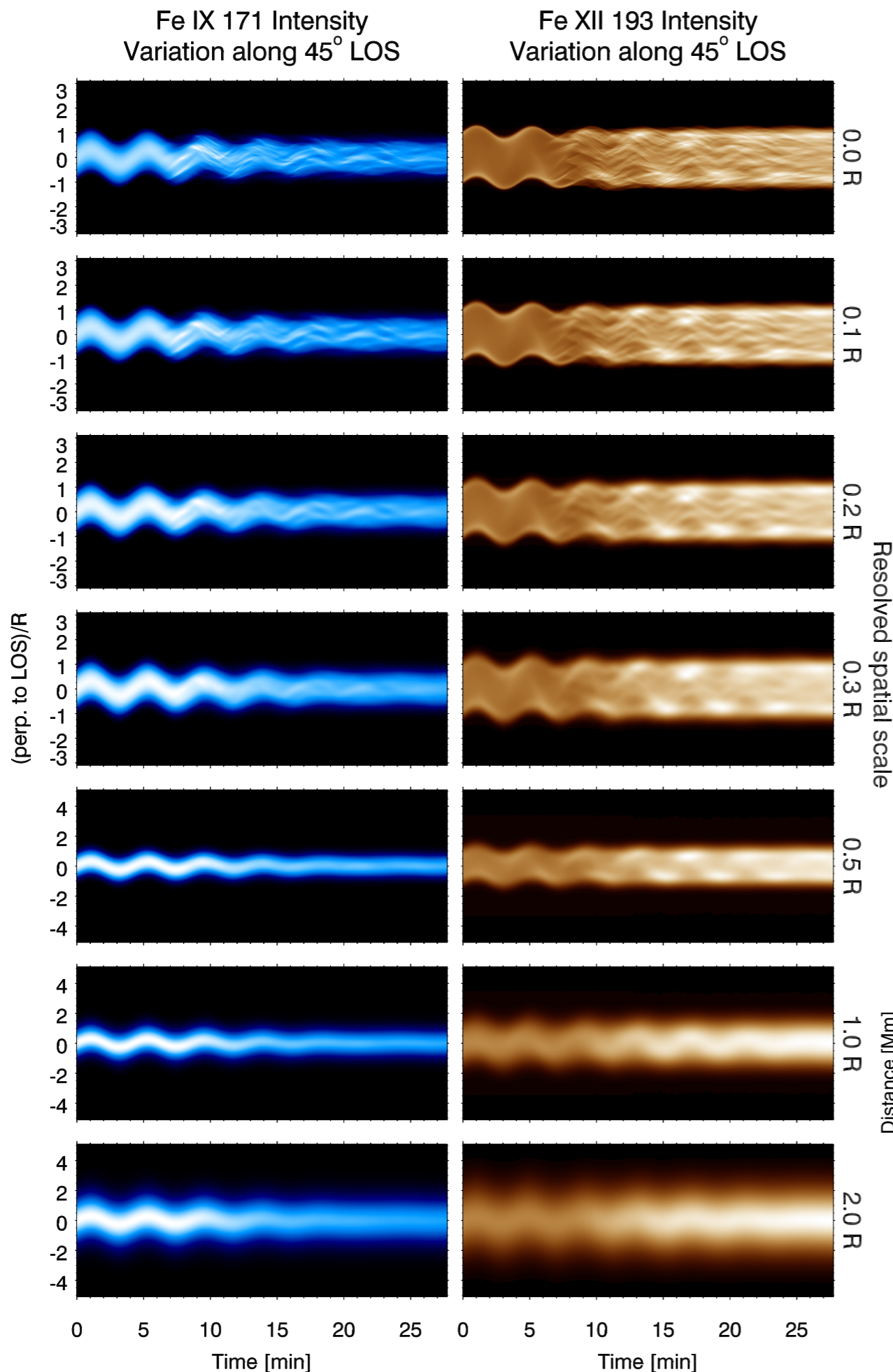
Intensity & loop width variation



- Double periodicity in intensity linked to KHI vortex generation
- Width variation linked to centrifugal force or flute modes

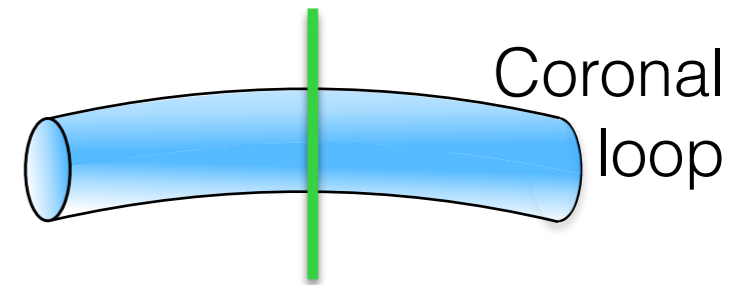
Decay-less oscillations: Alfvén waves

- Damping observed in 171 is not observed in 193 at low resolution
- Due to ensemble of correlated vortices (which show much less damping due to resonance) & brightening at maximum displacement

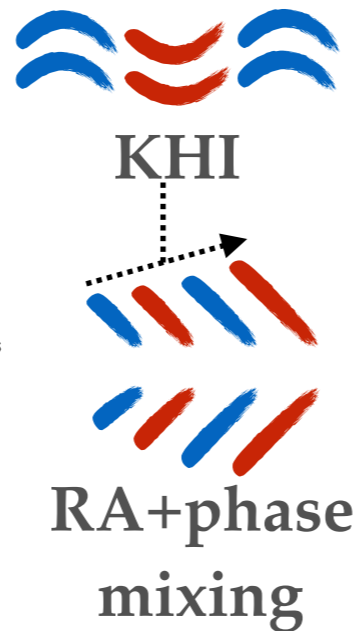
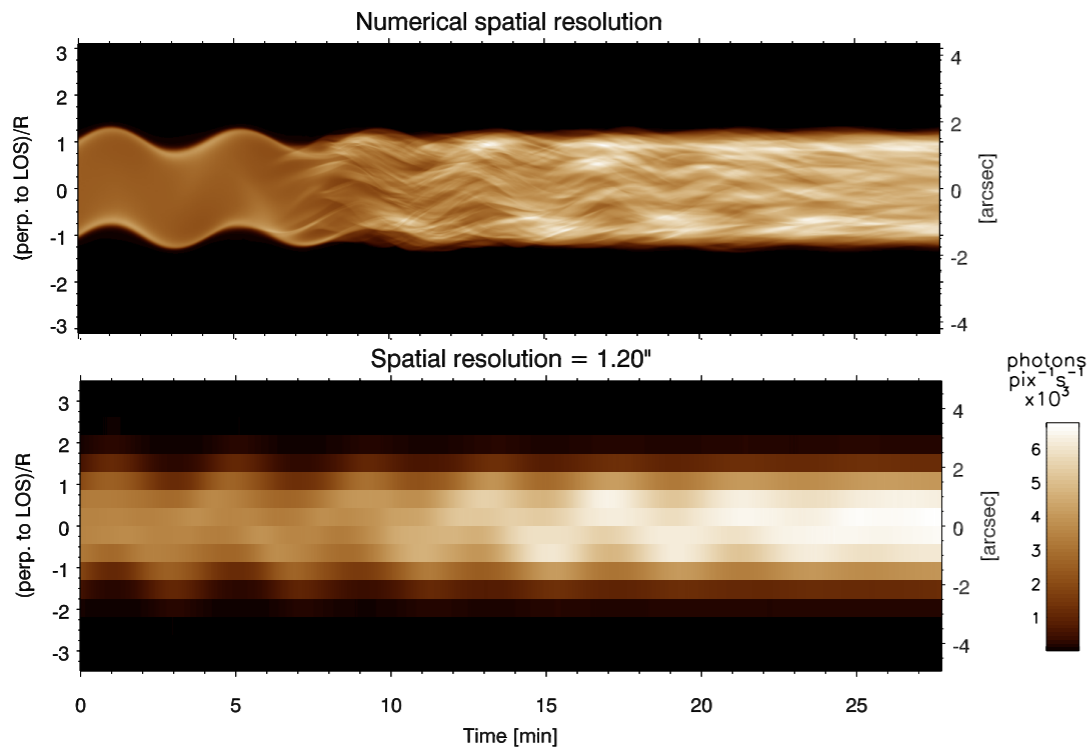


Nisticò+(2013),
Anfinogentov+ (2013,
2015),
Goddard+(2015)

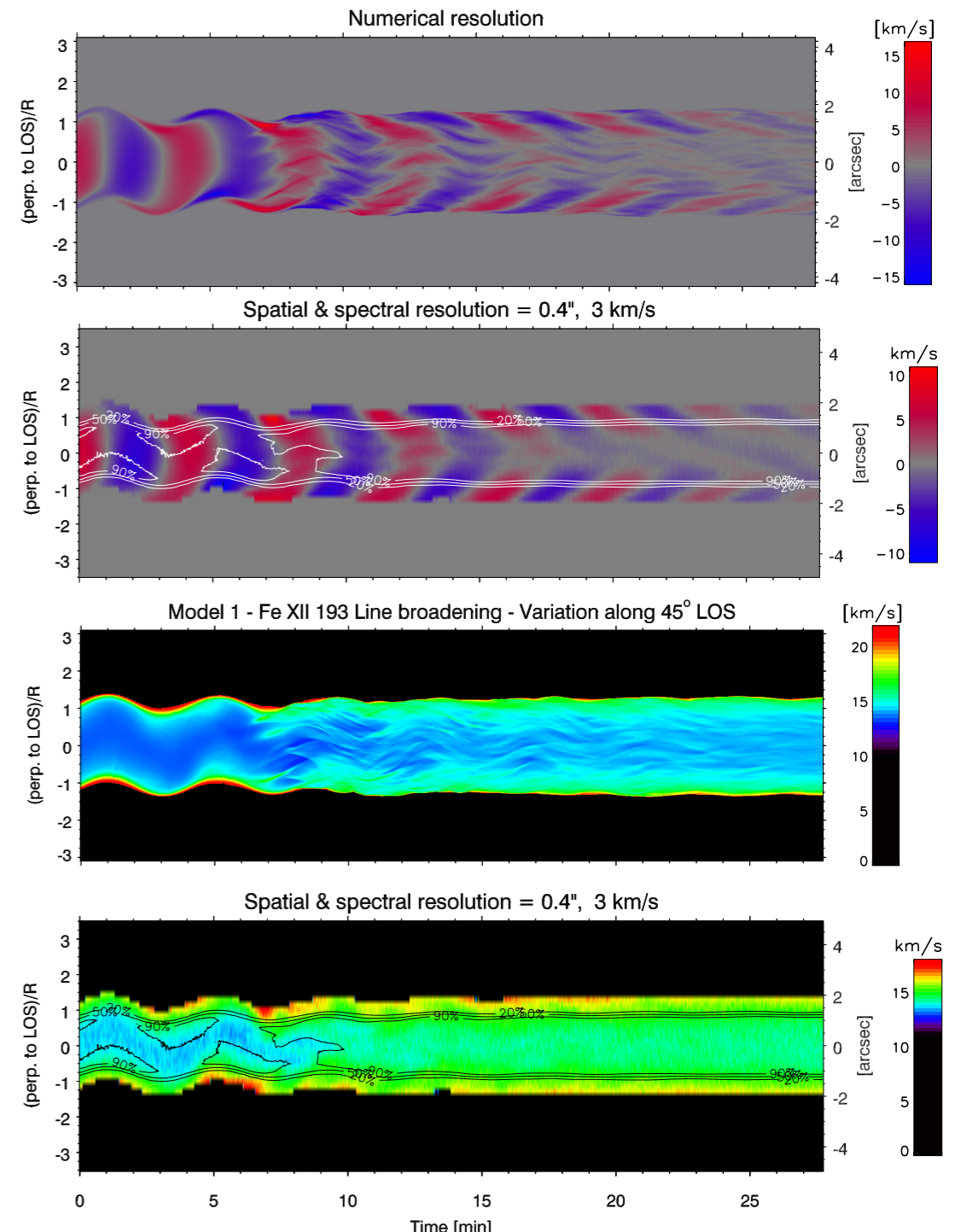
Doppler velocities & line widths



Fe XII 193 Intensity - Variation along 45° LOS



Fe XII 193 Doppler velocity - Variation along 45° LOS



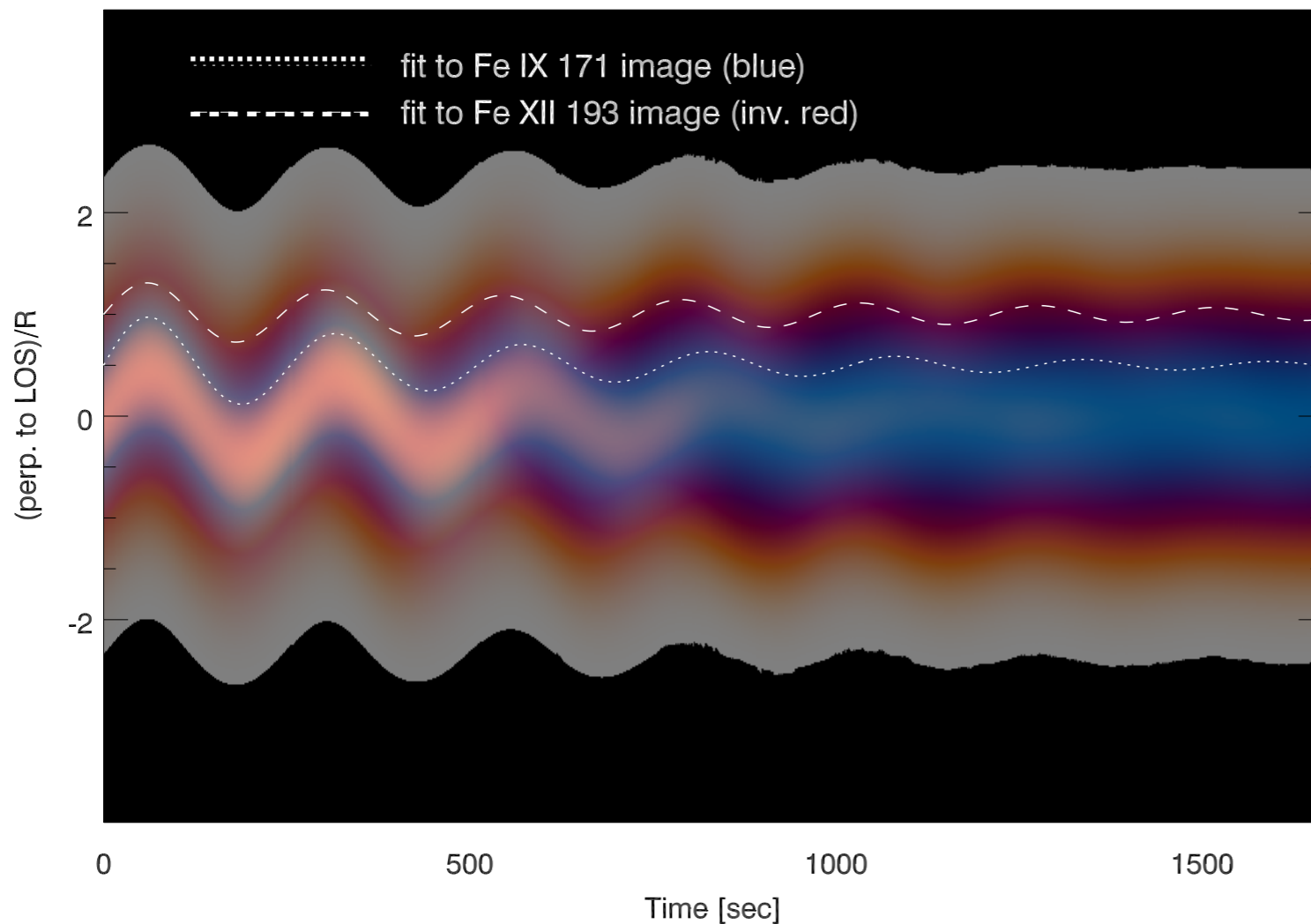
- Apparently decay-less
- Strand-like structure
- Broad arrow-shaped Doppler maps
- Broad line widths at edges

(Antolin+2016, submitted)

Decay-less oscillations: Alfvén waves

Antolin+ (2016, submitted)

Combined Intensities Fe IX & Fe XII - Resolution = 1.0 R - LOS angle = 00°



Period in 171 ~ 255 s

Period in 193 ~ 243 s

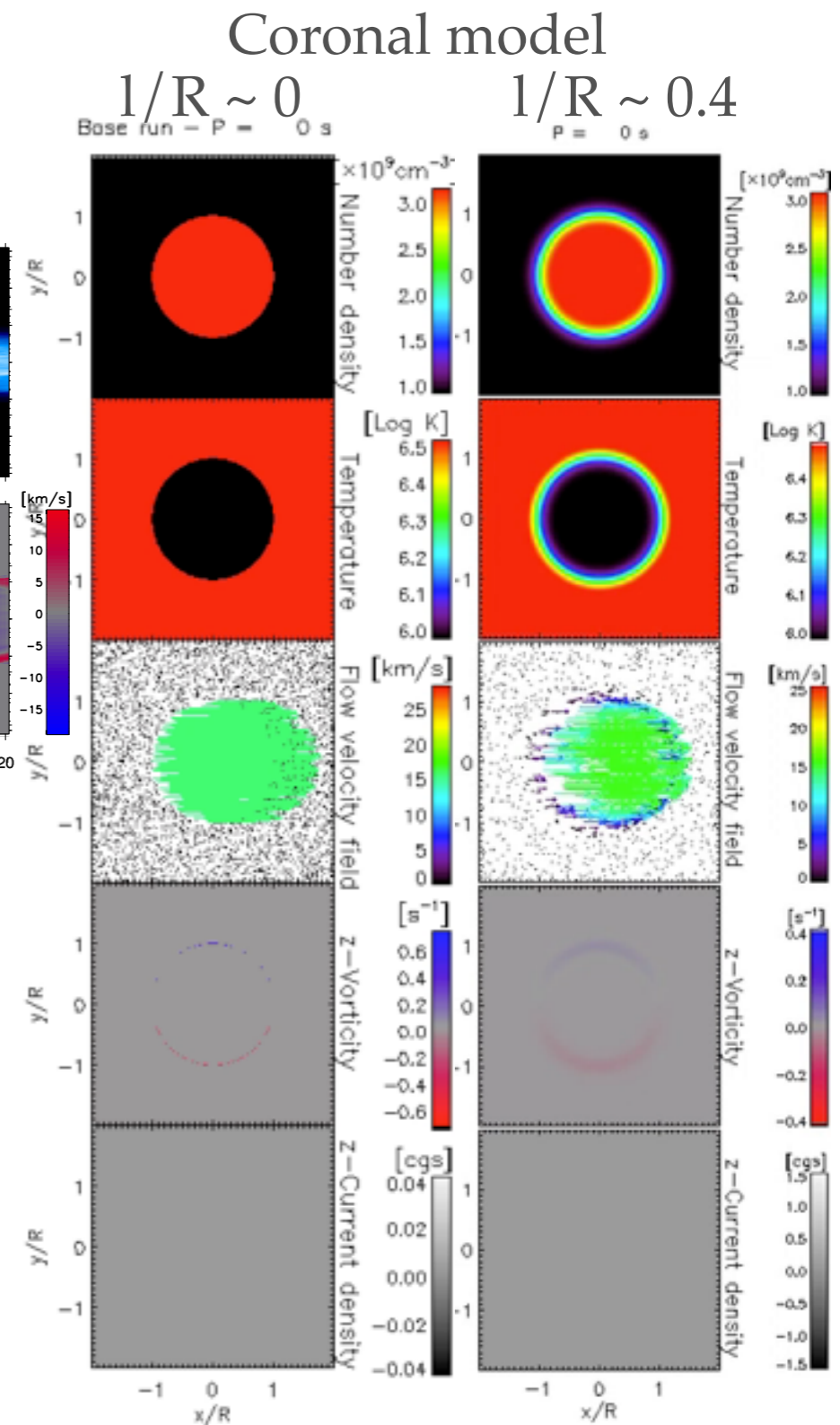
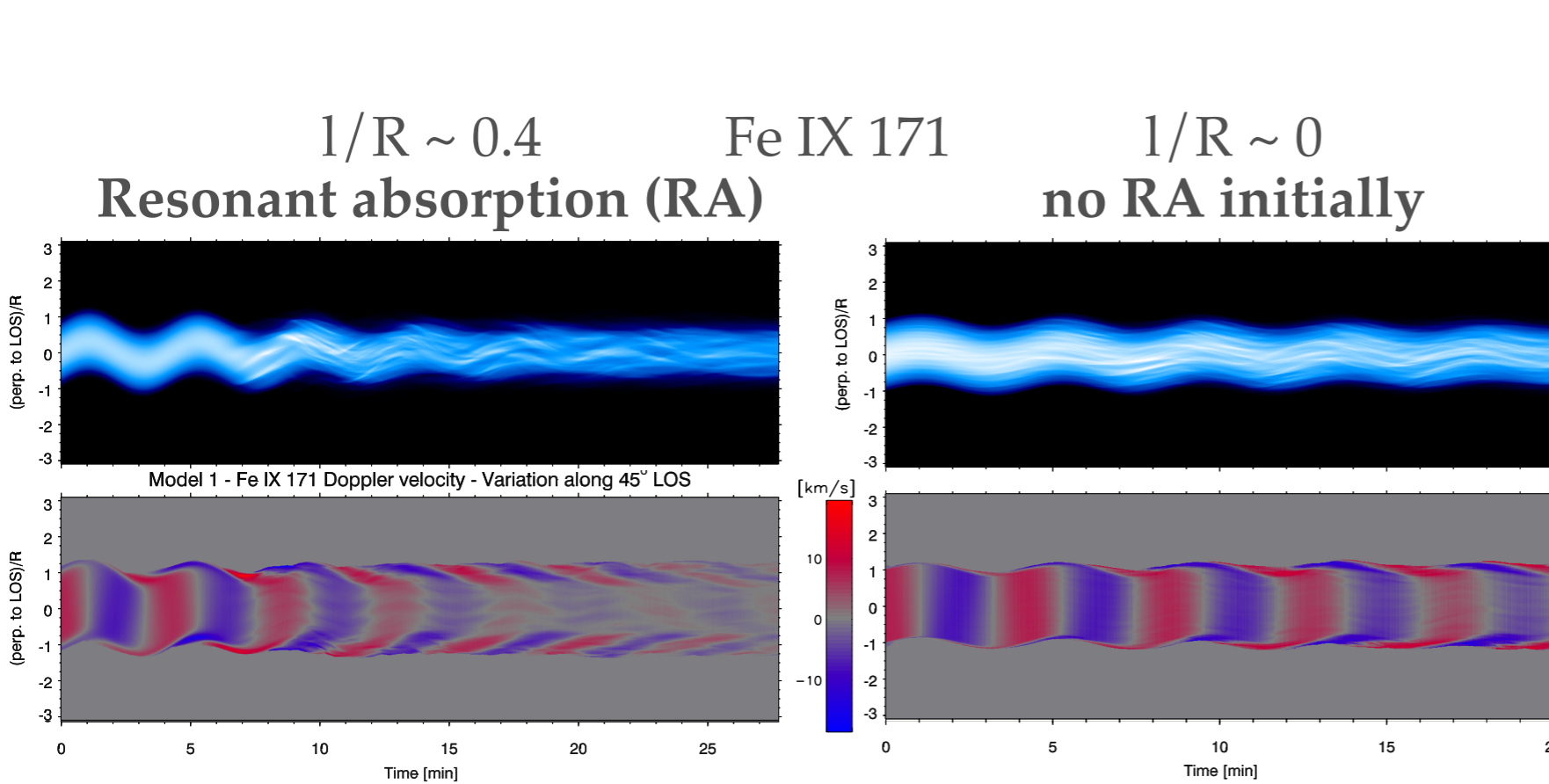
- Oscillation in 193 and 171 go out-of-phase: effect of phase mixing
- Oscillation in 193 reflects azimuthal Alfvén wave oscillation at boundary

$$\frac{\rho_b}{\rho_i} = \frac{1}{2} \left(1 + \frac{\rho_e}{\rho_i} \right) \left(\frac{P_b}{P_k} \right)^2$$

$$\rightarrow n = 1.8 \times 10^9 \text{ cm}^{-3}$$

➔ Combining channels sensitive to different temperatures we can perform high resolution seismology

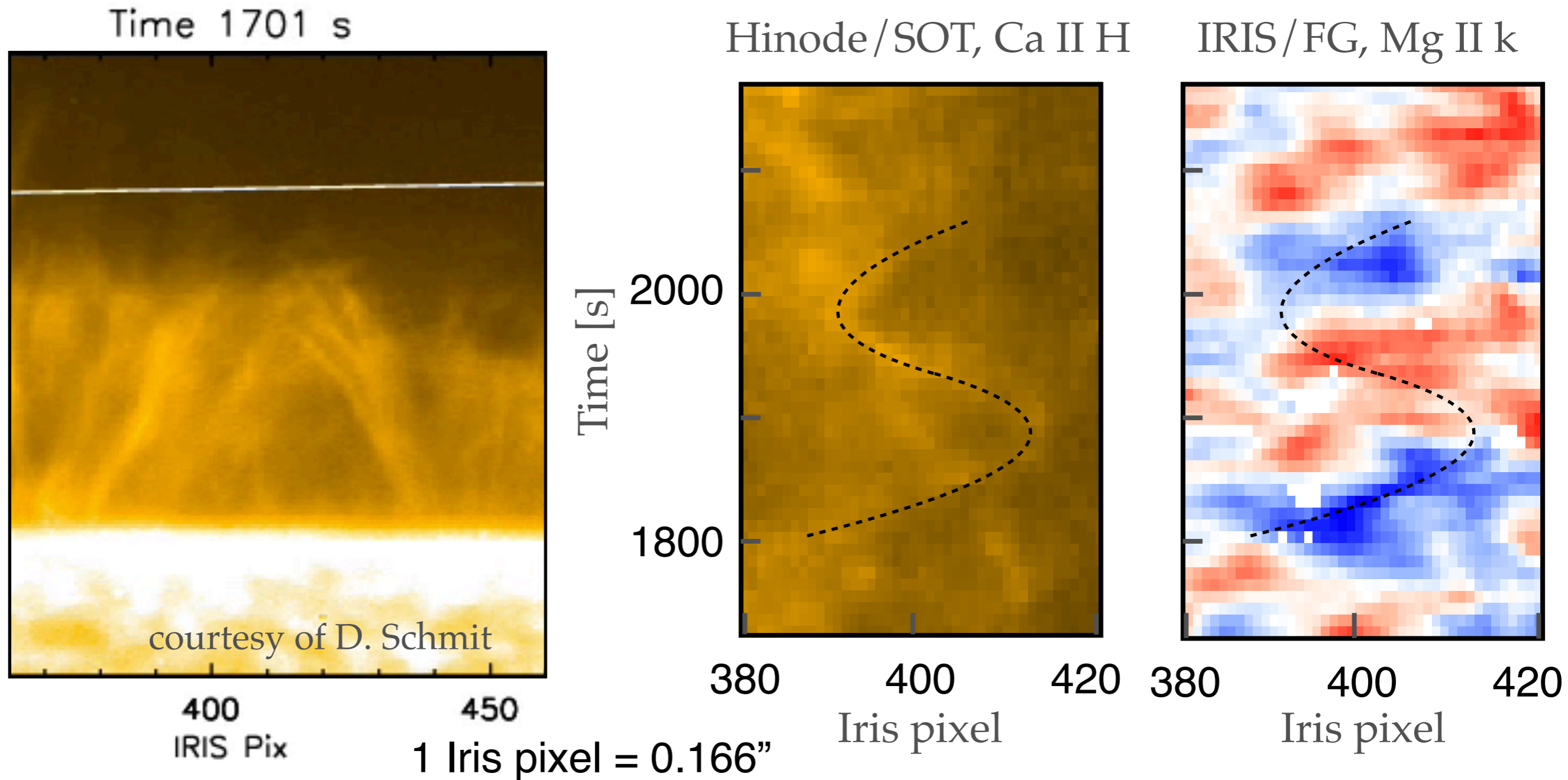
Effect of resonant absorption



- Larger wavenumbers dominate first for thinner boundaries (higher growth rates). “Inverse” cascade.
- No RA in boundary layer \rightarrow less kinetic energy in boundary layer. Small wavenumbers have lower growth rates (no large vortices)
- No boundary layer: less damping
- **Significant difference in dynamics of core & boundary layer**

IRIS/Hinode observations: spicules

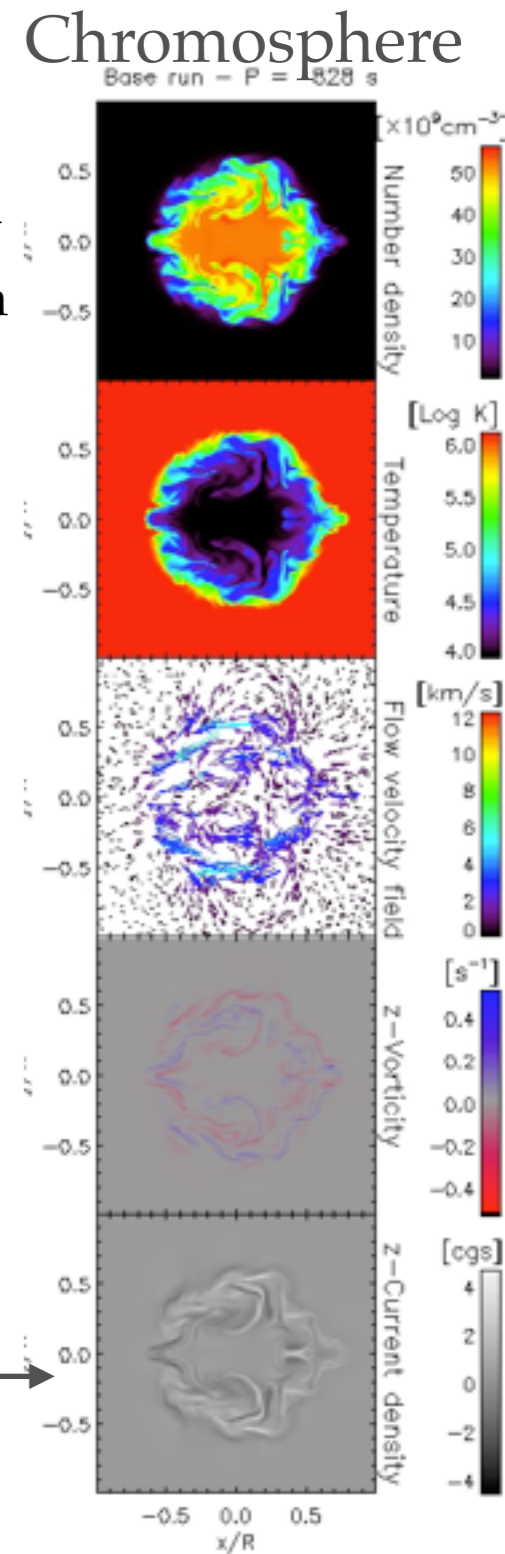
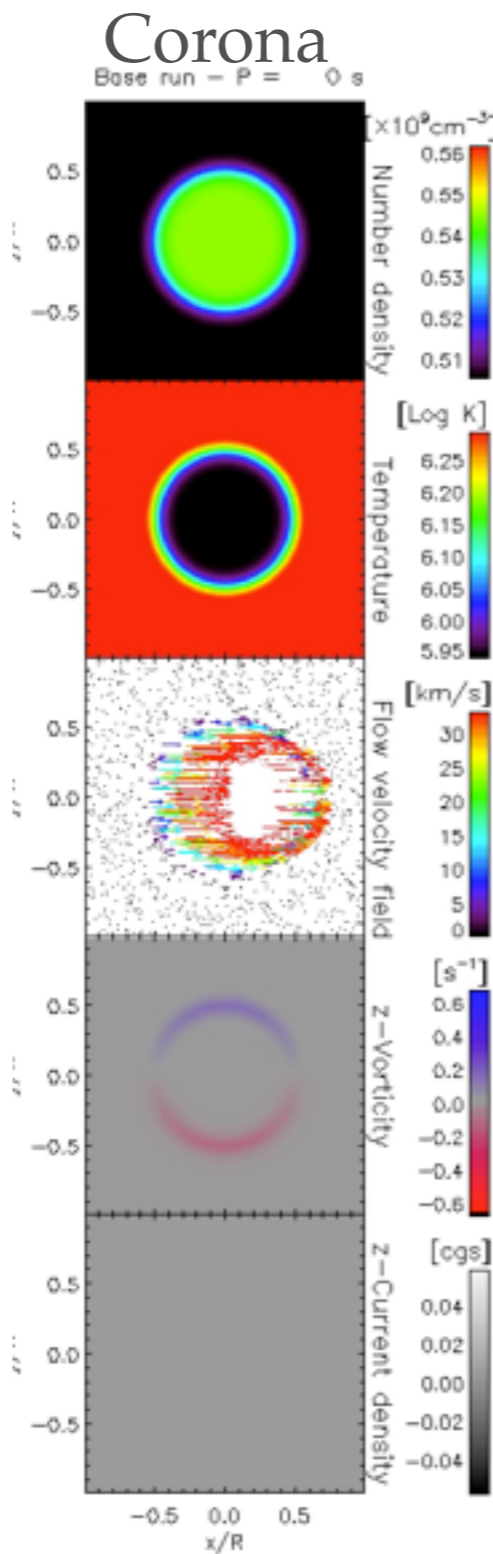
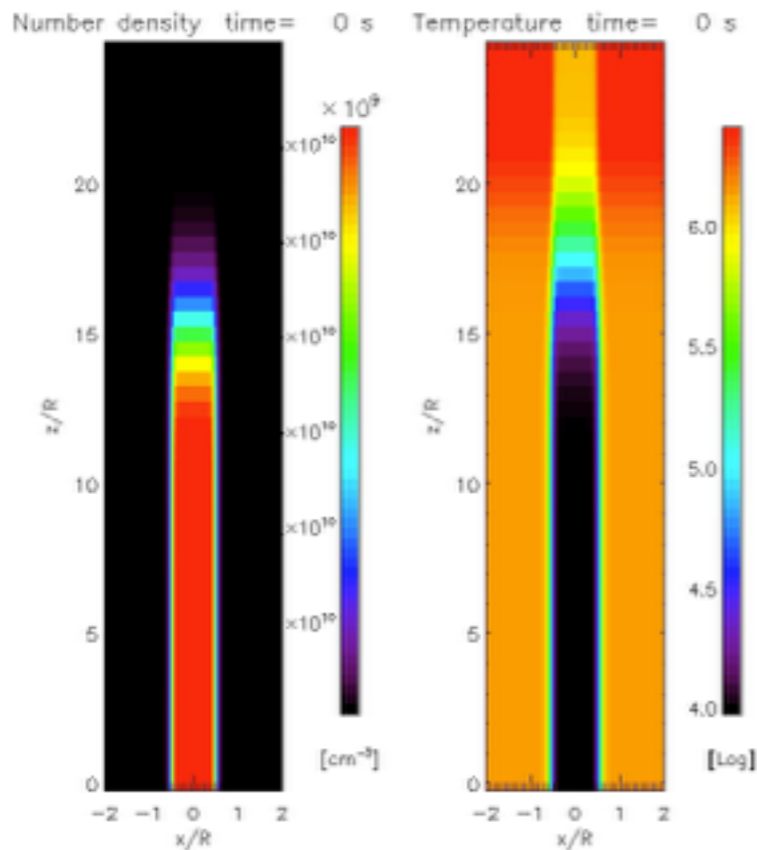
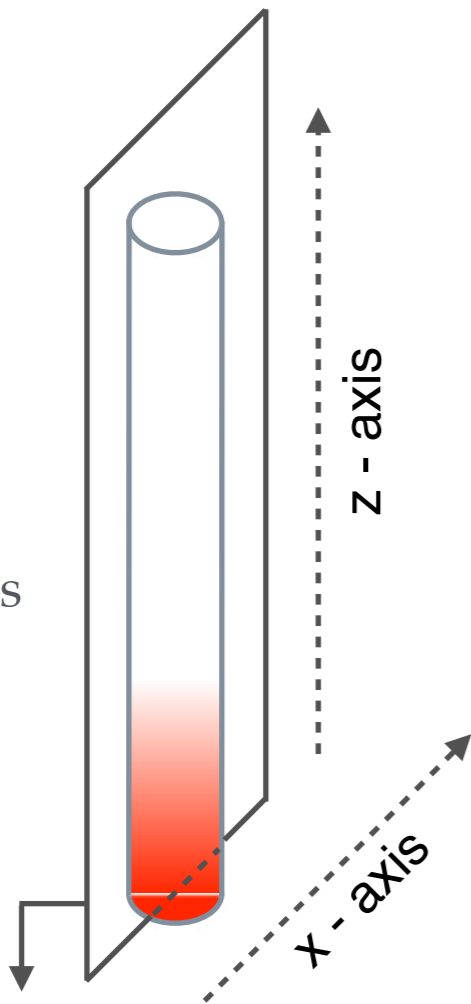
Antolin, Schmit, De Pontieu, Pereira (2016, in prep.)



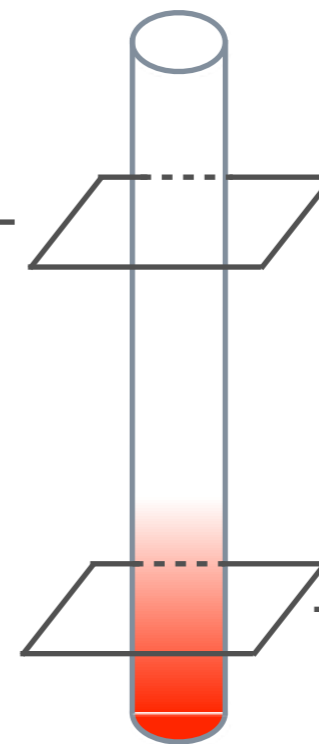
Spicule model

standing

Amp. ~ 12 km/s
 Period ~ 245 s



Strong vorticity
 destroys loop in
 corona
 (Magyar & Van
 Doorselaere
 2016)

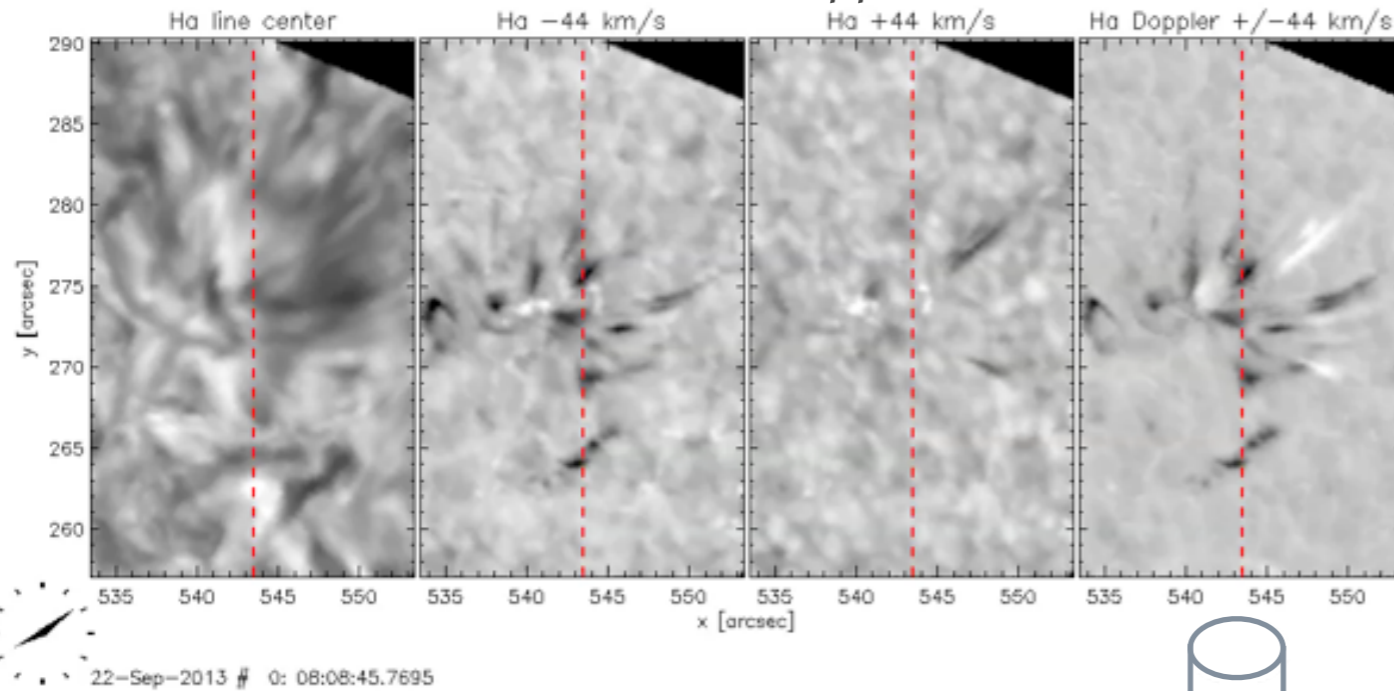


Spicule dynamics

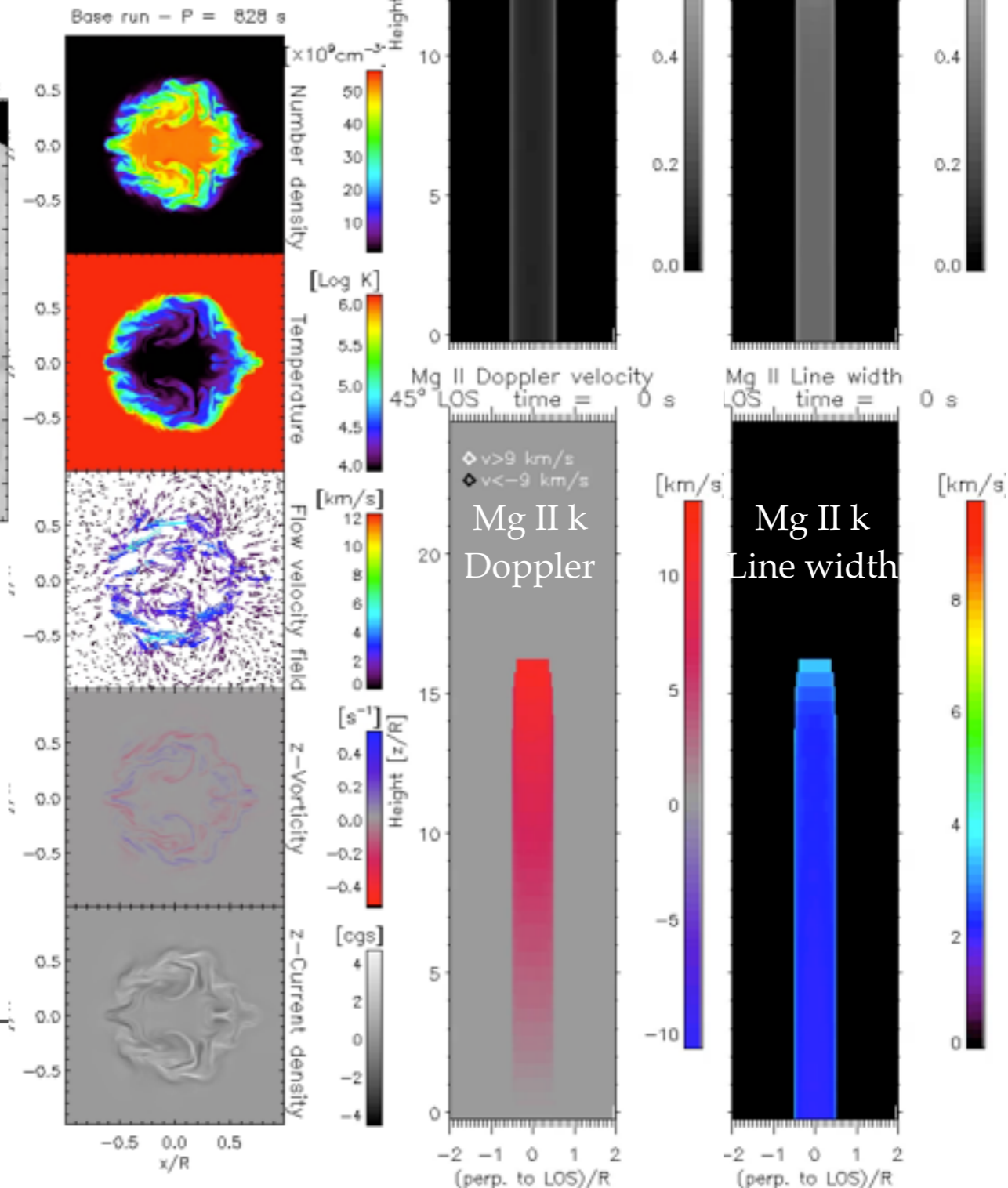
standing kink

SST/CRISP Observations

Rouppé van der Voort+ 2015

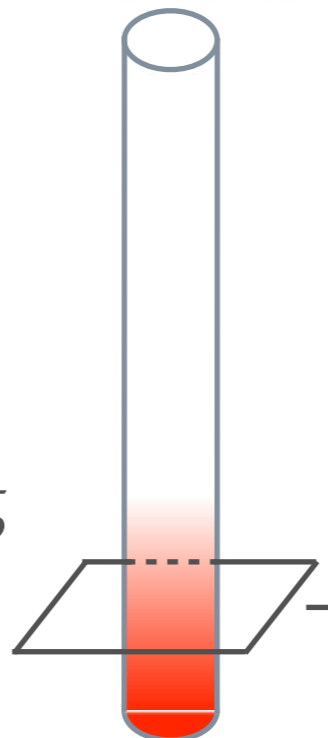


Simulated spicule cross-section



- KHI vortices +RA can explain spicular transverse dynamics
- Potential candidate for heating signatures

(Antolin+2016, in prep.)

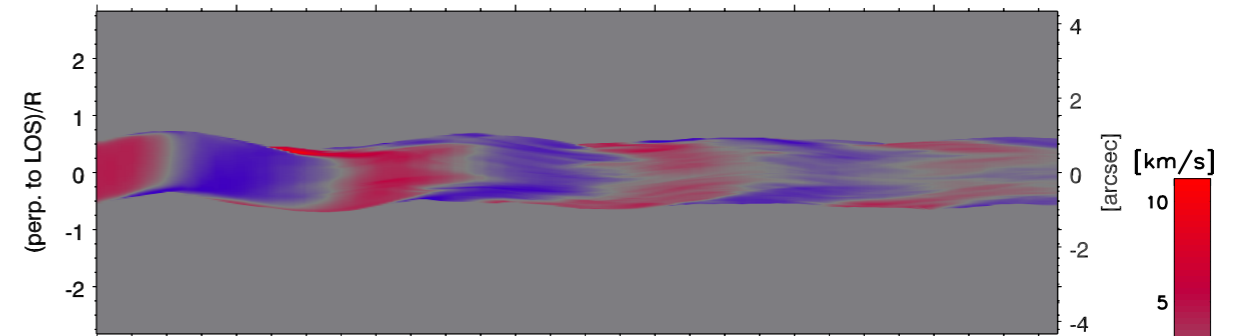


Spicule dynamics

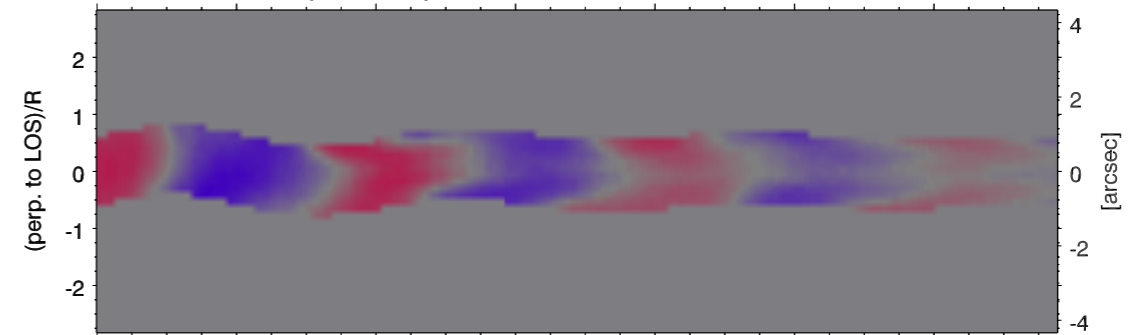
Mg II

Mg II 2796 Doppler velocity - Variation along 45° LOS

Numerical resolution

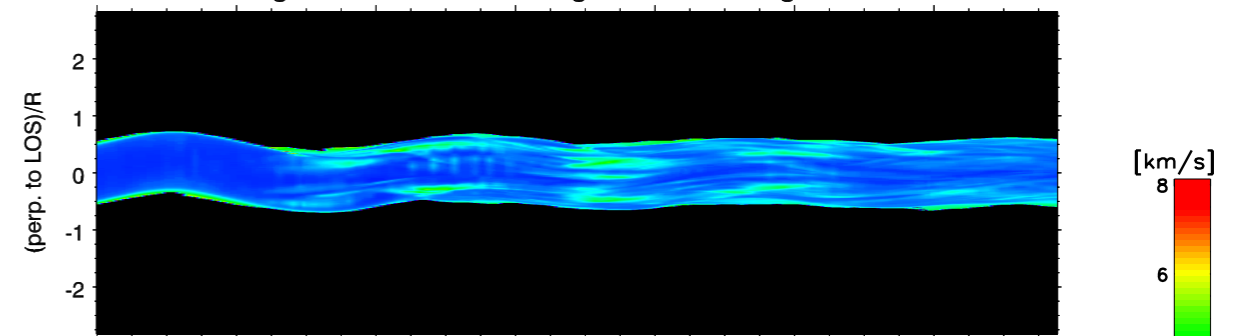


Spatial & spectral resolution = 0.332", 03 km/s

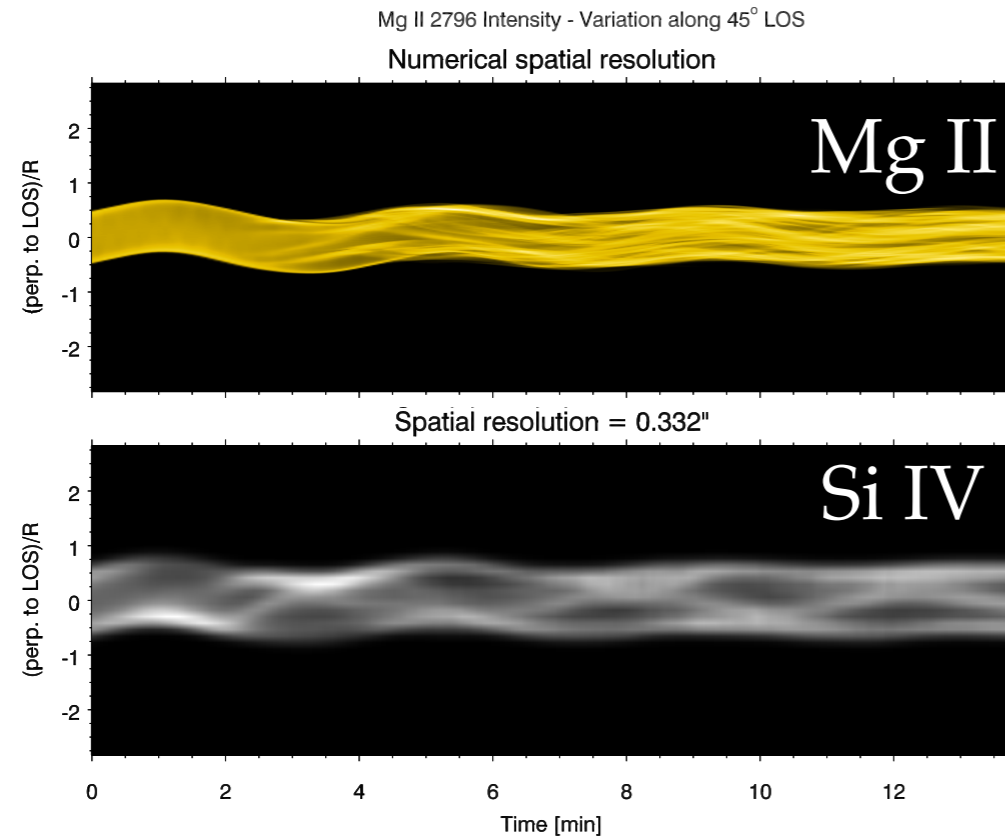
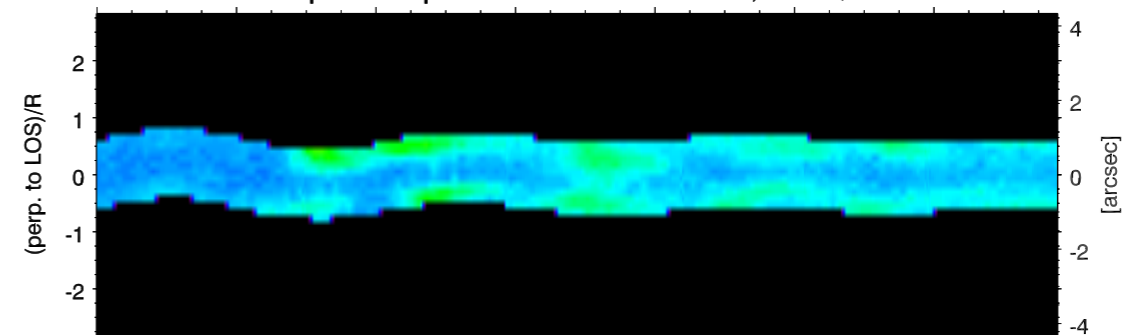


Mg II 2796 Line broadening - Variation along 45° LOS

Mg II 2796 Line broadening - Variation along 45° LOS



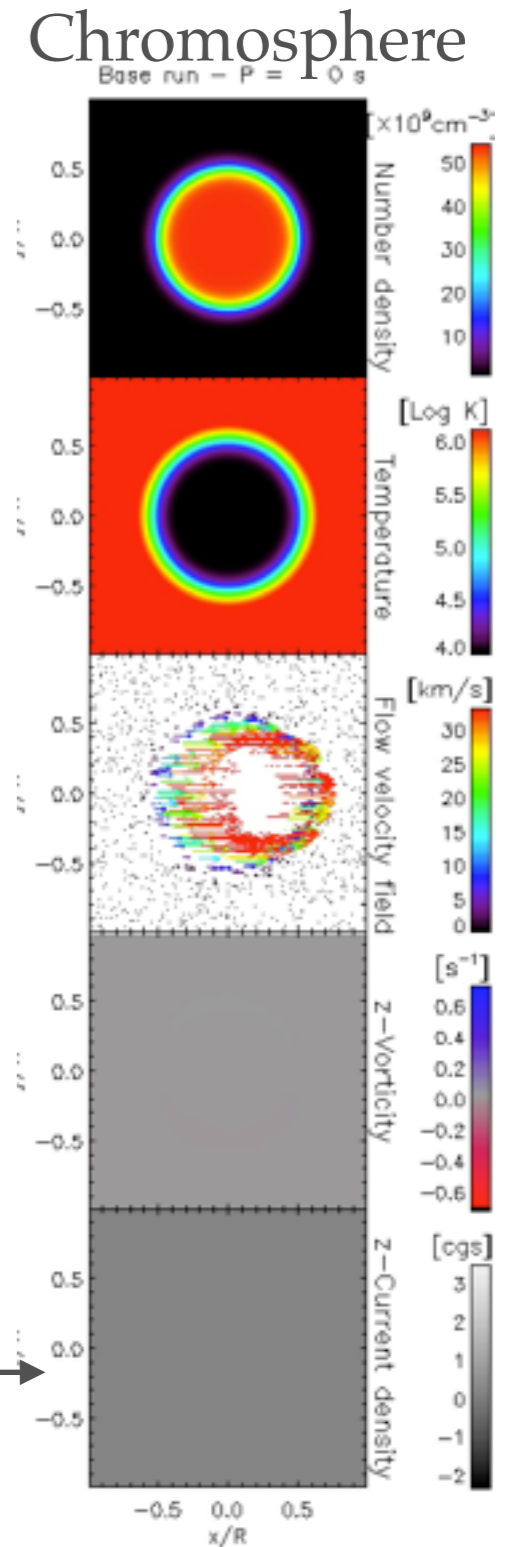
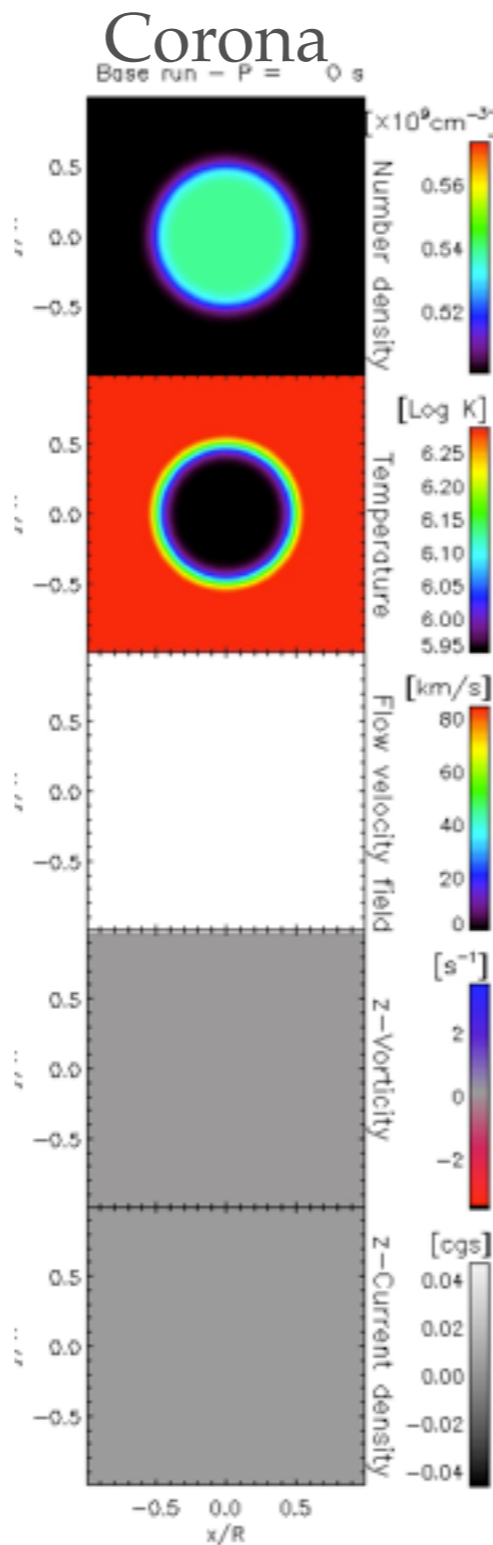
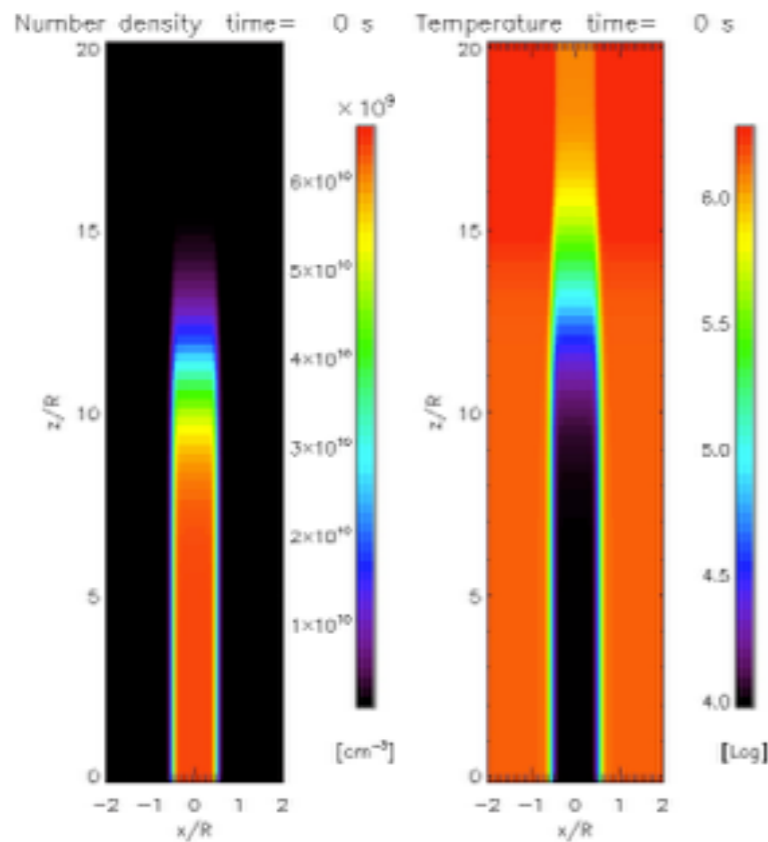
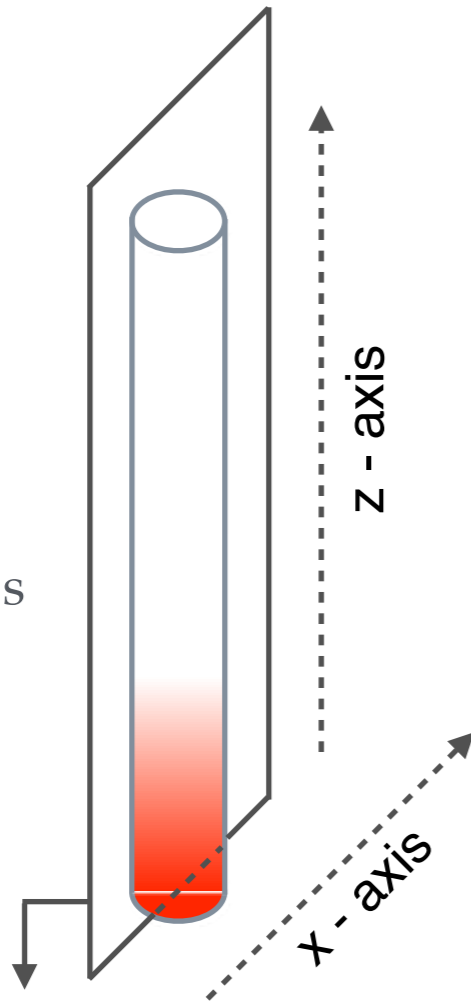
Spatial & spectral resolution = 0.332", 03 km/s



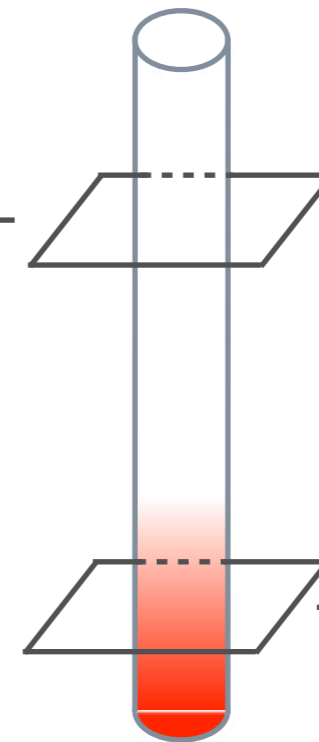
- Similar features as for the prominence and the coronal model

Spicule model *propagating*

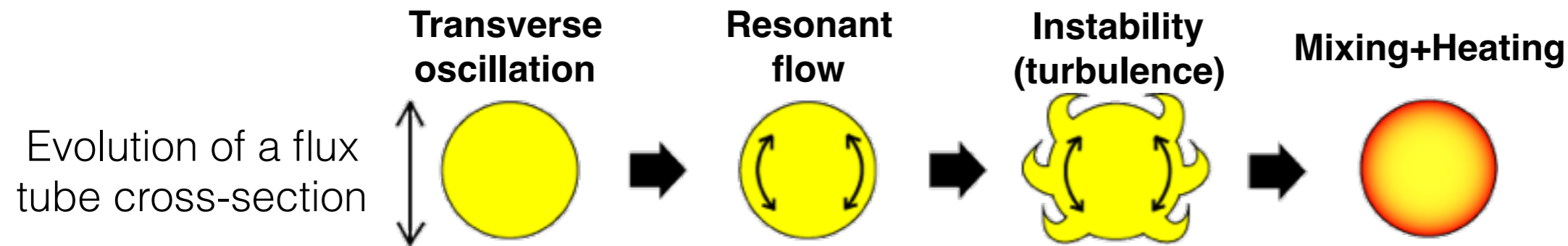
Amp. ~ 30 km/s
Period ~ 245 s



Strong vorticity
destroys loop in
corona
(Magyar & Van
Doorselaere
2016)



Conclusions



- ✿ 3D MHD model of transverse MHD wave in prominence shows 2 mechanisms at play
- ✿ Resonant absorption + dynamic (KHI) instabilities = heating
- ✿ KHI fine structure leads to **strand-like structure** in intensity images (*Antolin+ 2014*). **Thread-like structure/strand-like structure in observations = Alfvénic vortices?**
- ✿ Out-of-phase (90° - 180°) behaviour between POS motion in cool lines and hot lines: **very good match with IRIS/Hinode observations** (*Okamoto+ 2015, Antolin+ 2015*)
- ✿ Resonant absorption enhances significantly KHI dynamics: **Alfvénic turbulence**
- ✿ Damping in cool lines (probing loop core temperature). **Oscillation in hot line can be decay-less (probing boundary layer). Due to ensemble of vortices and heating.** May explain observed decay-less observations (*Nisticò+2013, Anfinogentov+2013*)
- ✿ KHI fine structure dependent on boundary layer width & amplitude. Differences in damping and phase between cool/hot emission lines: **seismology tool**
- ✿ May explain observed spicule transverse dynamics
- ✿ Model valid for **corona, prominences and spicules.**

Thank you for your attention!



European Research Council

Established by the European Commission



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