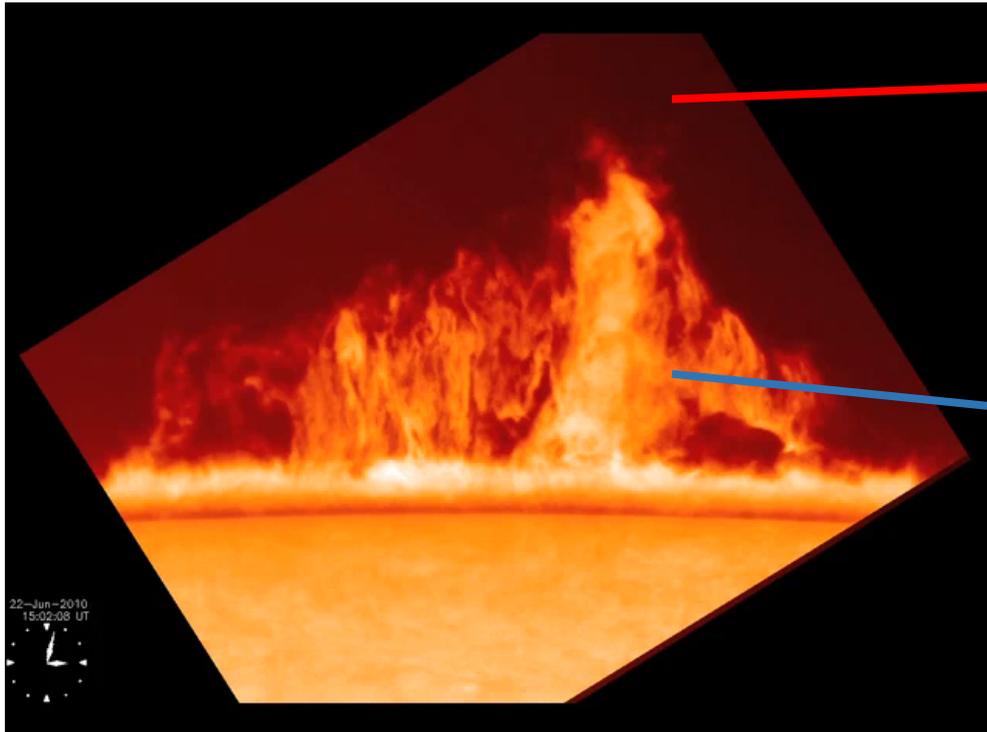


Three-dimensional MHD Simulation of Prominence Formation by Radiative Condensation

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Solar Prominence



Corona

temperature $\sim 10^6$ K
density $10^8 \sim 10^9 \text{ cm}^{-3}$

Prominence

Cool dense plasma cloud
temperature $< 10^4$ K
density $10^9 \sim 10^{11} \text{ cm}^{-3}$

Berger et al., 2011, Hinode/SOT, $H\alpha$

Mechanism of Prominence Formation

What is the origin of cool dense plasmas ?

Radiative condensation/Thermal nonequilibrium (instability):

Coronal plasmas are cooled down and condensed by radiative cooling. (Karpen et al., 2007; Luna et al., 2012; Xia et al., 2012; Kaneko & Yokoyama, 2015; Xia & Keppens, 2016)

Injection, Levitation:

Chromospheric plasmas are lifted up to coronal height by jet or emerging flux.

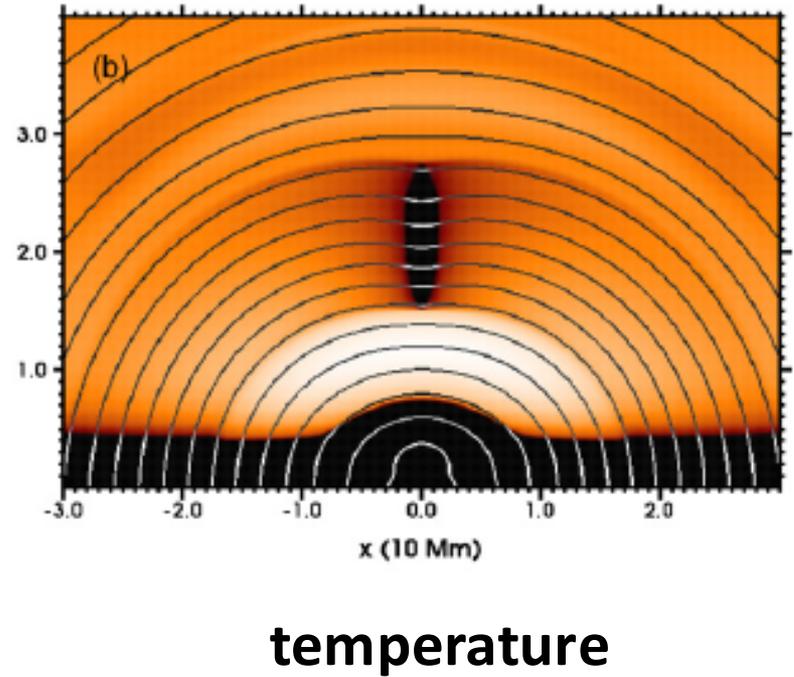
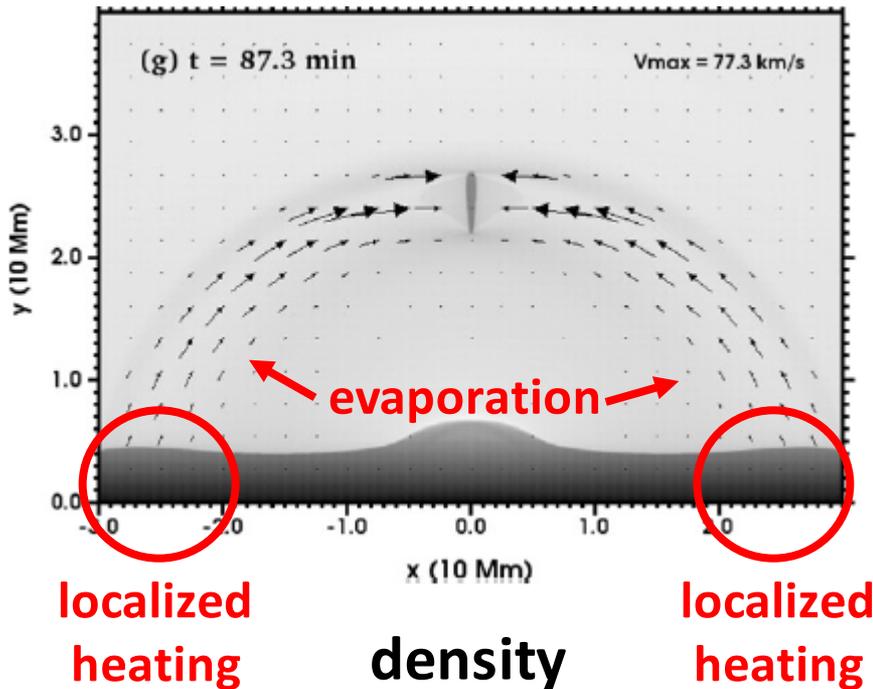
(Chae et al., 2003; Okamoto et al., 2007, 2008; Deng et al., 2000)

Evaporation-condensation model

Chromospheric evaporation by Localized footpoint heating

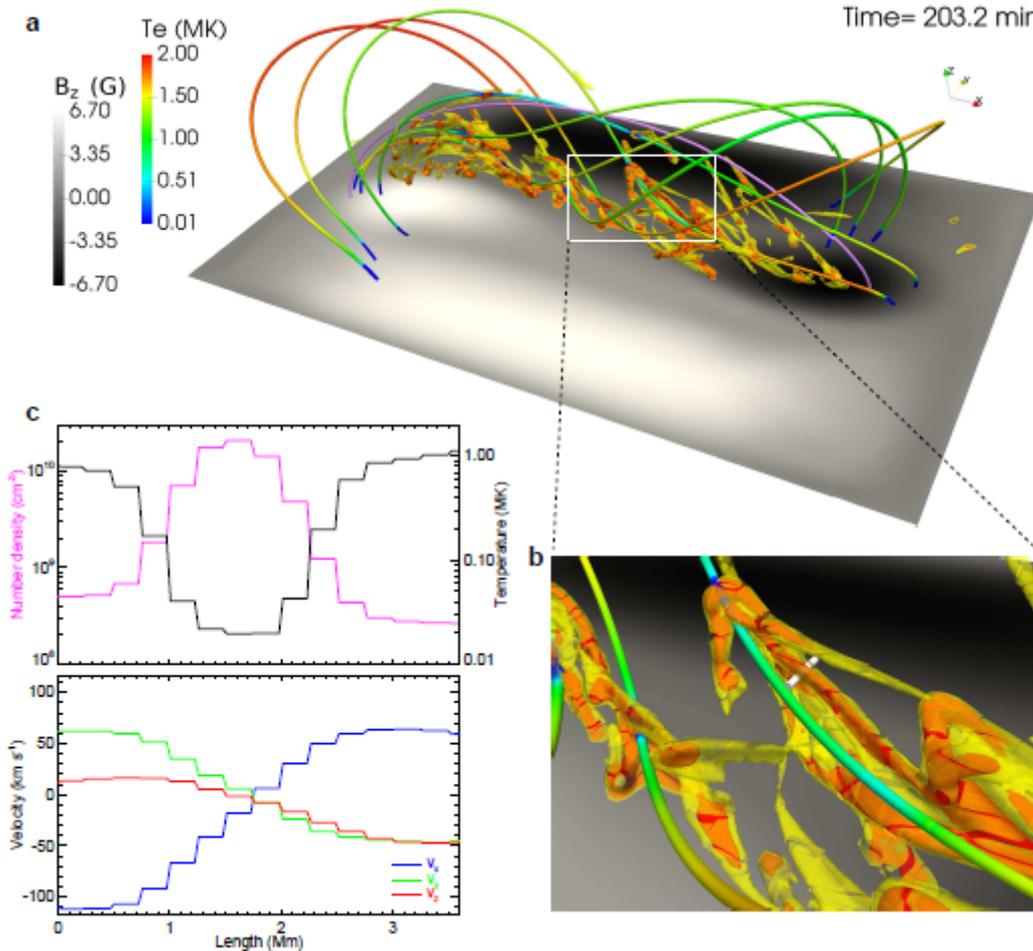
➔ Radiative condensation

Xia et al. (2012)



Evaporation-condensation model (3D)

Xia & Keppens (2016)



Chromospheric evaporation
by footpoint heating



Enhancement of radiative
cooling inside flux rope by
injection of high density
plasmas



prominence formation by
radiative condensation



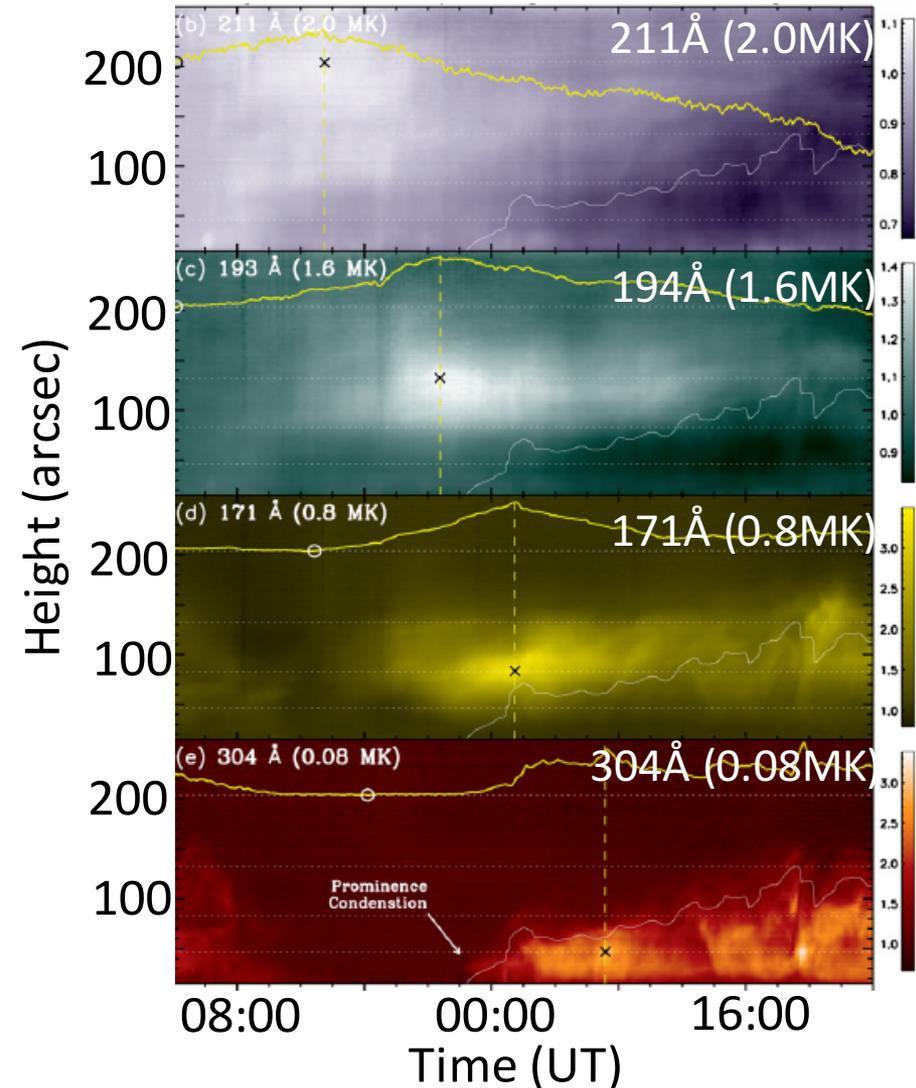
drainage of prominence to
chromosphere

Observation of in-situ condensation

Berger et al. (2012)



- Temporal intensity shift from high temperature to low temperature
-> **Radiative condensation**
- Evaporated flows had not detected.
-> **In-situ condensation** in the corona



Key steps to radiative condensation

Evaporation-condensation model

- **Enhancement of radiation** by deposition of high density plasmas
- **Limiting thermal conduction** by changing the direction of thermal flux (in long magnetic loop)

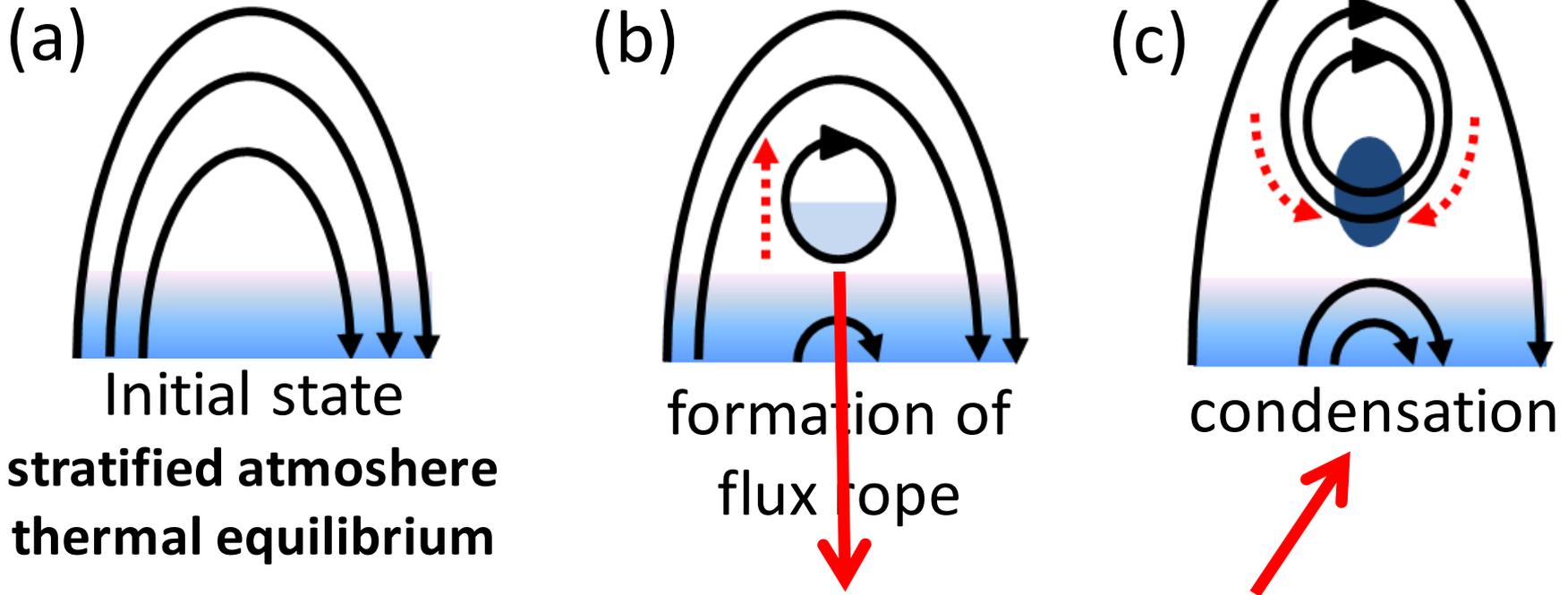


Reconnection-Condensation Model

Reconnection and subsequent **topological change** of magnetic field can also trigger radiative condensation.

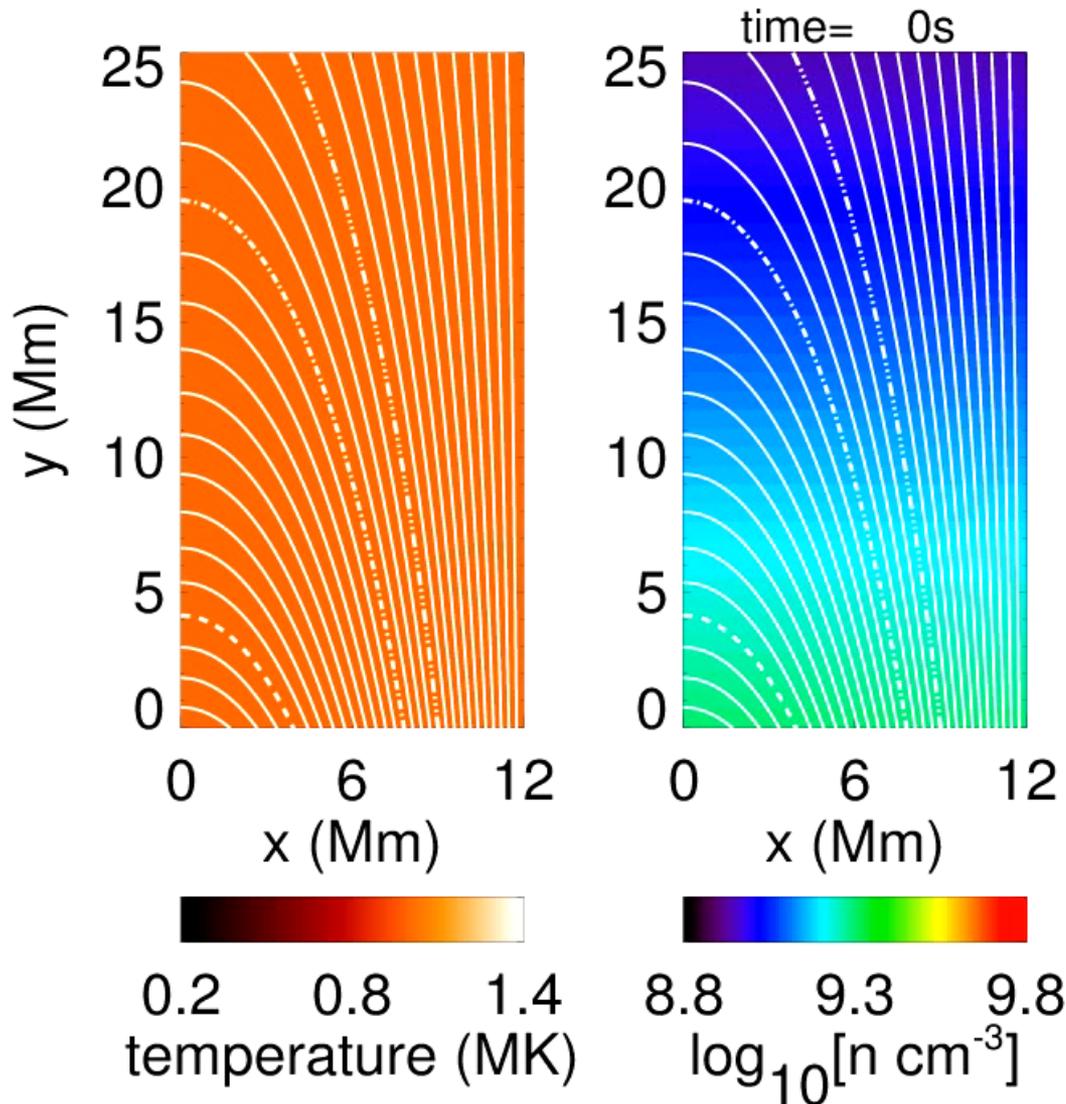
Reconnection-condensation model (2D)

Kaneko & Yokoyama (2015)



- relatively dense plasmas at the bottom (**strong radiation**)
- closed field line (**limiting thermal conduction**)

Demonstration by 2D simulation



flux rope formation



**thermal imbalance
in thermally isolated
closed loops**



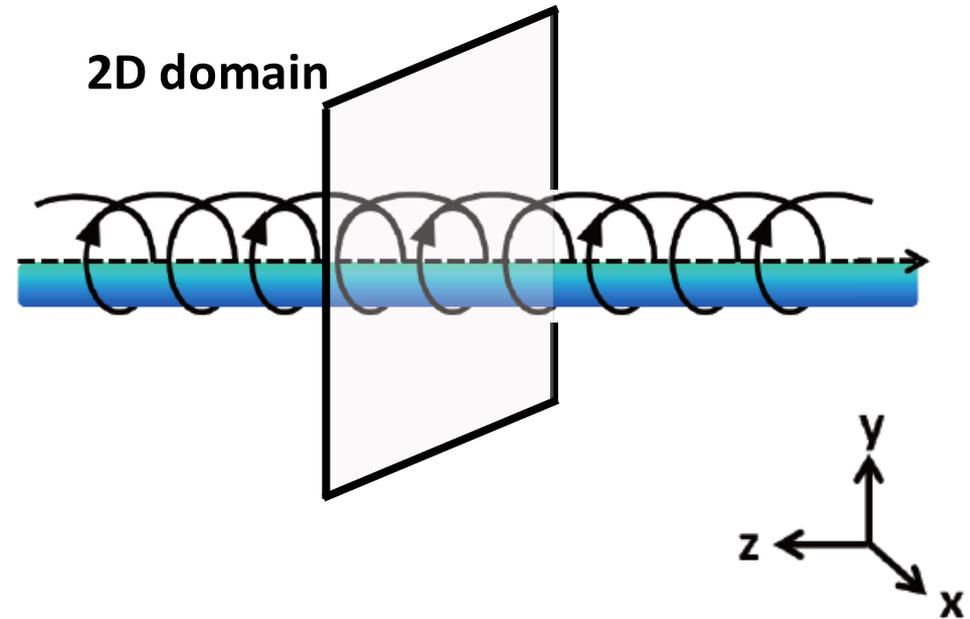
radiative condensation

Kaneko & Yokoyama (2015)

Feasibility of the model in 3D

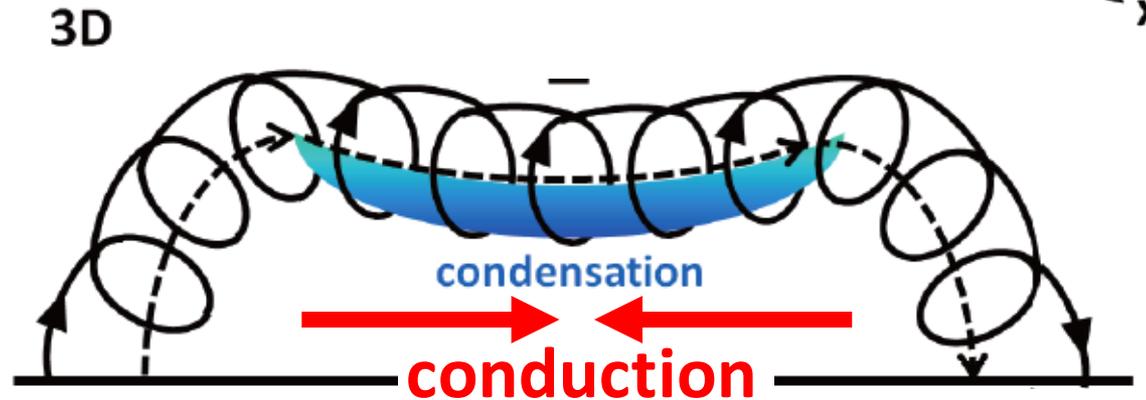
2.5D

->Temperature is uniform
perpendicular to 2D plane
=>Conduction along toroidal
components is ineffective.



3D

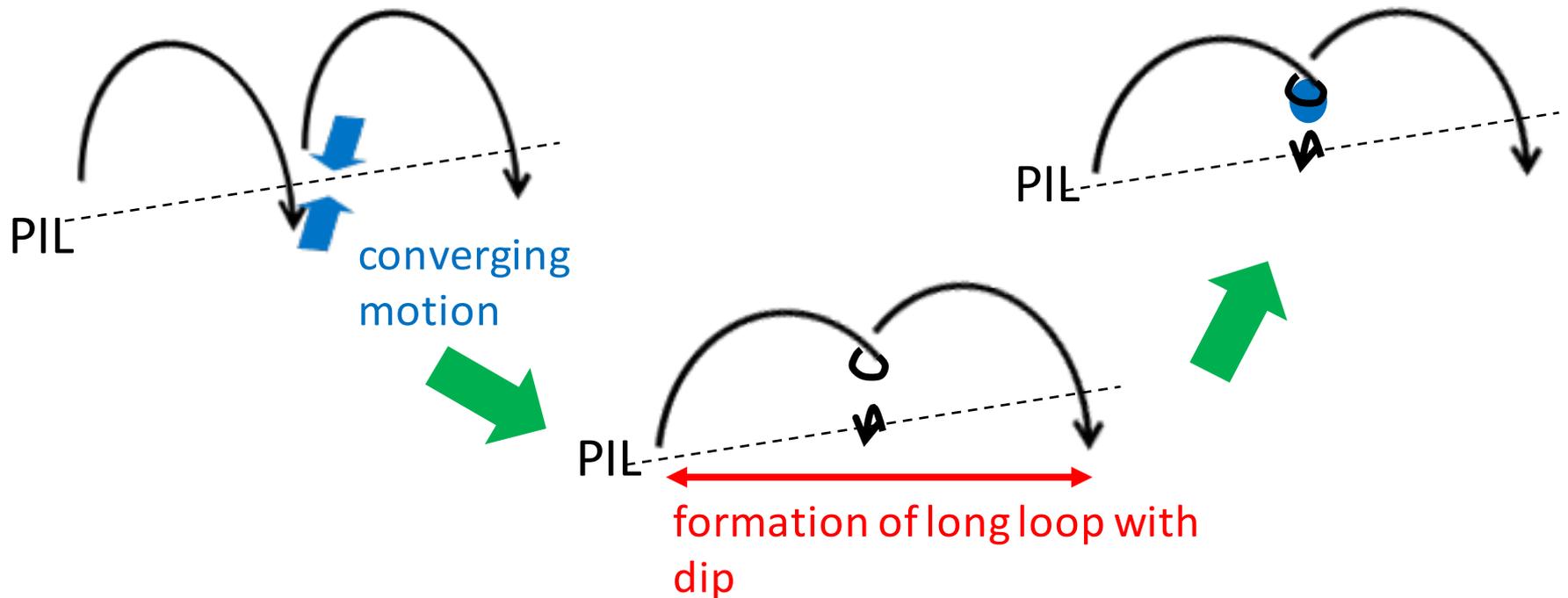
→Conduction along
toroidal magnetic field
may suppress thermal
instability.



Extension to 3D model

Aim:

Demonstration of reconnection-condensation model by 3D MHD simulation including radiative cooling & thermal conduction

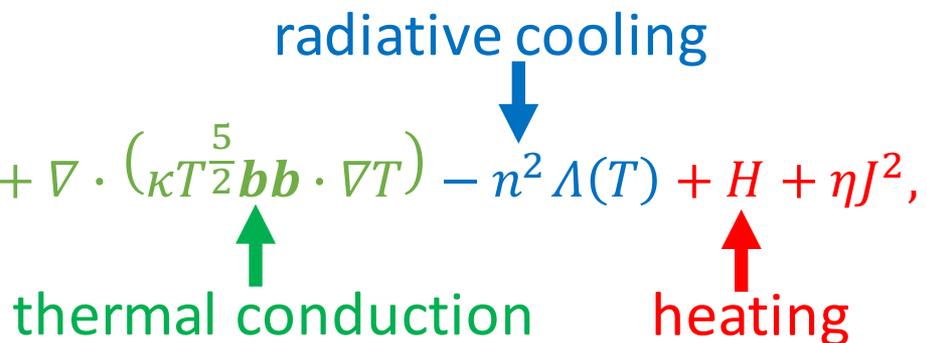


Numerical setting 1/5

Basic equations:

$$\frac{\partial \rho}{\partial t} + \mathbf{v} \cdot \nabla \rho = -\rho \nabla \cdot \mathbf{v},$$

$$\frac{\partial e}{\partial t} + \mathbf{v} \cdot \nabla e = -(e + p) \nabla \cdot \mathbf{v} + \nabla \cdot (\kappa T^{\frac{5}{2}} \mathbf{b} \mathbf{b} \cdot \nabla T) - n^2 \Lambda(T) + H + \eta J^2,$$



$$e = \frac{p}{\gamma - 1}, \quad T = \frac{m p}{k_B \rho},$$

$$\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} = -\frac{1}{\rho} \nabla p + \frac{1}{4\pi\rho} (\nabla \times \mathbf{B}) \times \mathbf{B} + \mathbf{g},$$

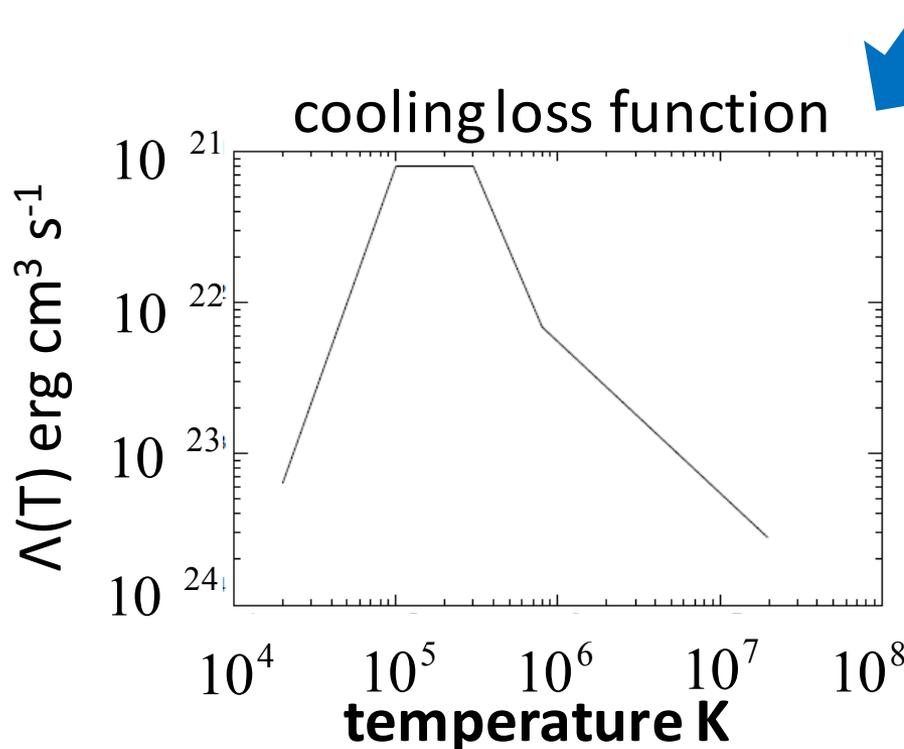
$$\frac{\partial \mathbf{B}}{\partial t} = -c \nabla \times \mathbf{E},$$

$$\mathbf{E} = -\frac{1}{c} \mathbf{v} \times \mathbf{B} + \frac{4\pi\eta}{c^2} \mathbf{J}, \quad \mathbf{J} = -\frac{c}{4\pi} \nabla \times \mathbf{B}.$$

Numerical setting 2/5

Energy equation:

$$\frac{\partial e}{\partial t} + \mathbf{v} \cdot \nabla e = -(e + p) \nabla \cdot \mathbf{v} - \underbrace{n^2 \Lambda(T)}_{\text{radiative cooling}} + \underbrace{H}_{\text{background heating}} + \eta J^2 + \nabla \cdot (\kappa T^{\frac{5}{2}} \mathbf{b} \mathbf{b} \cdot \nabla T)$$



$$H \propto P_m$$

coronal heating related to magnetic energy density (e.g. Parker, 1983)

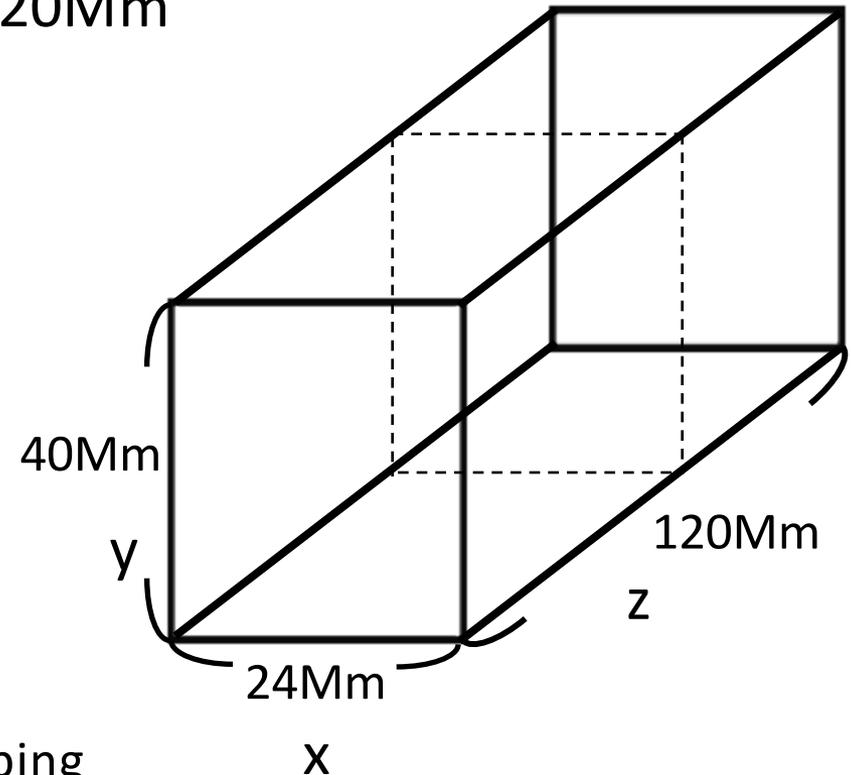
Numerical setting 3/5

Simulation box: 24Mm x 40Mm x 120Mm
(point-symmetry at $z=60\text{Mm}$)

Grid size: 120km (uniform)

Scheme

- 4th order Runge-Kutta
- 4th order central difference
- Artificial viscosity (Rempel, 2012)
- Hyperbolic divergence cleaning (Dedner, 2002)
- Thermal conduction: Super Time Stepping (RKL 2nd order, Meyer et al.2012,2014)

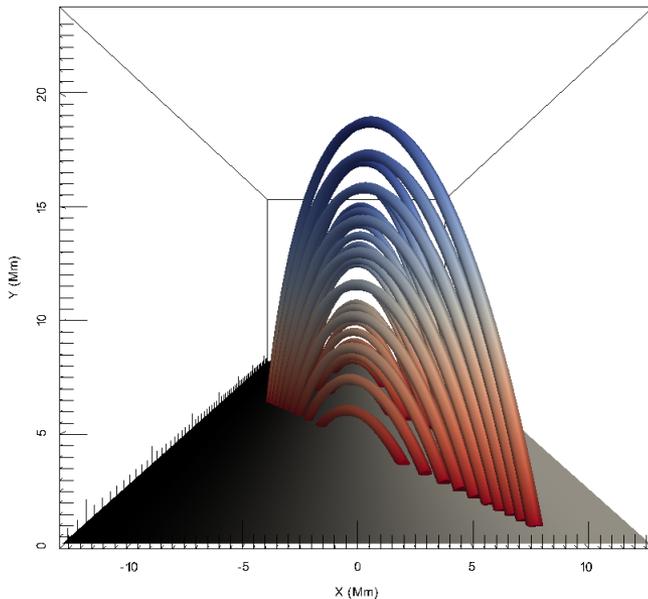


Numerical setting 4/5

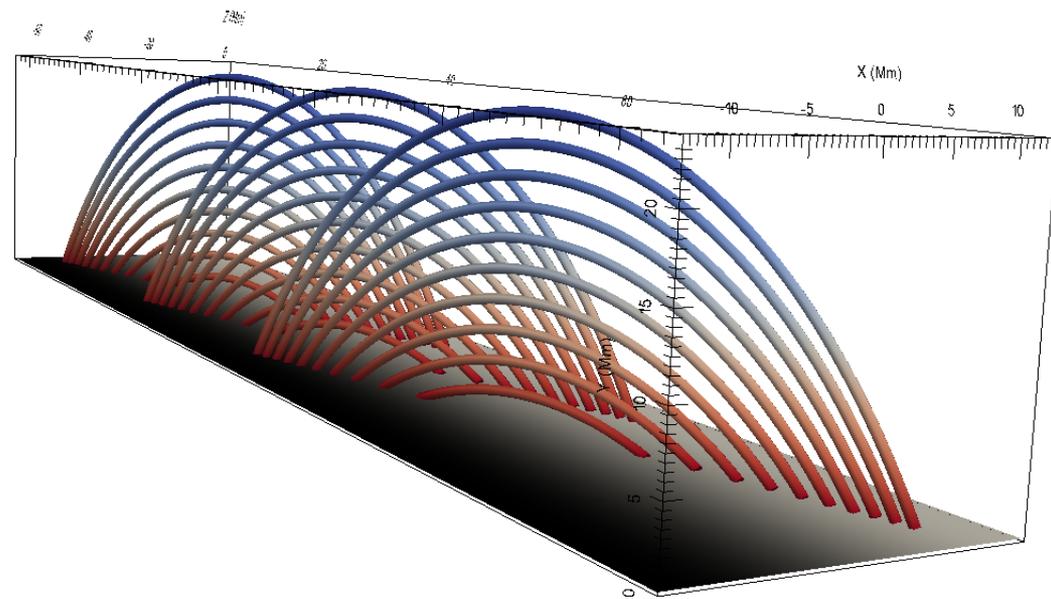
Initial condition

- magnetic field: linear force-free arcade field (< 6 G)
- stratified under uniform temperature & gravity
($T = 1\text{MK}$, $n < 2 \times 10^9 \text{ cm}^{-3}$)

front view



side view



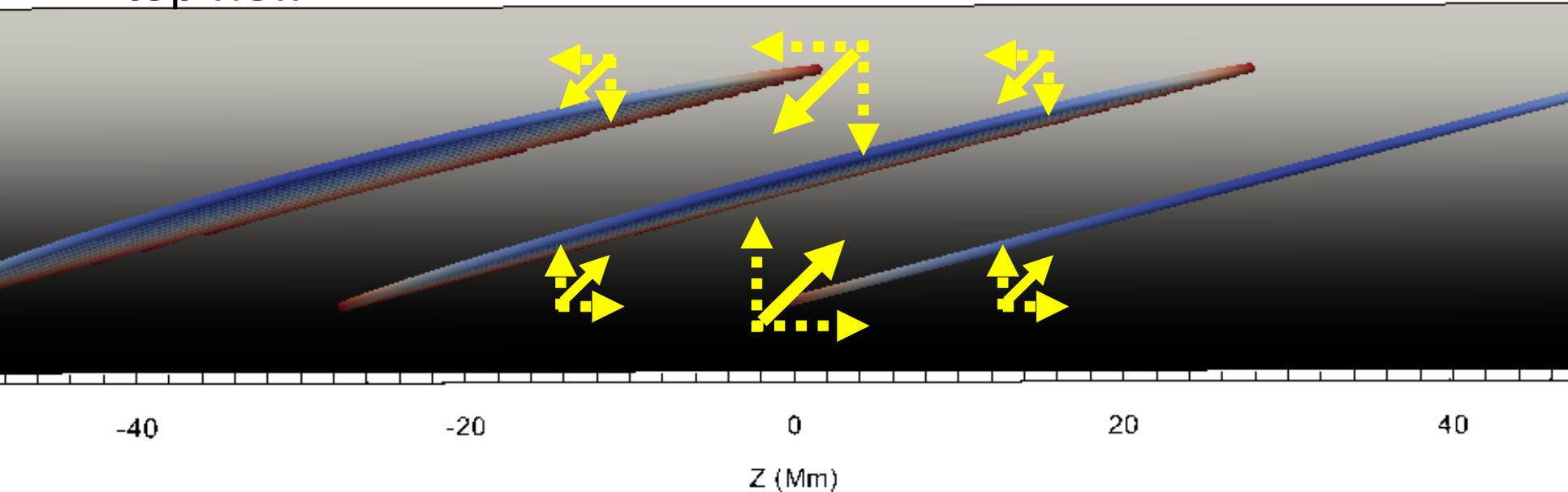
Numerical setting 5/5

Footpoint motion

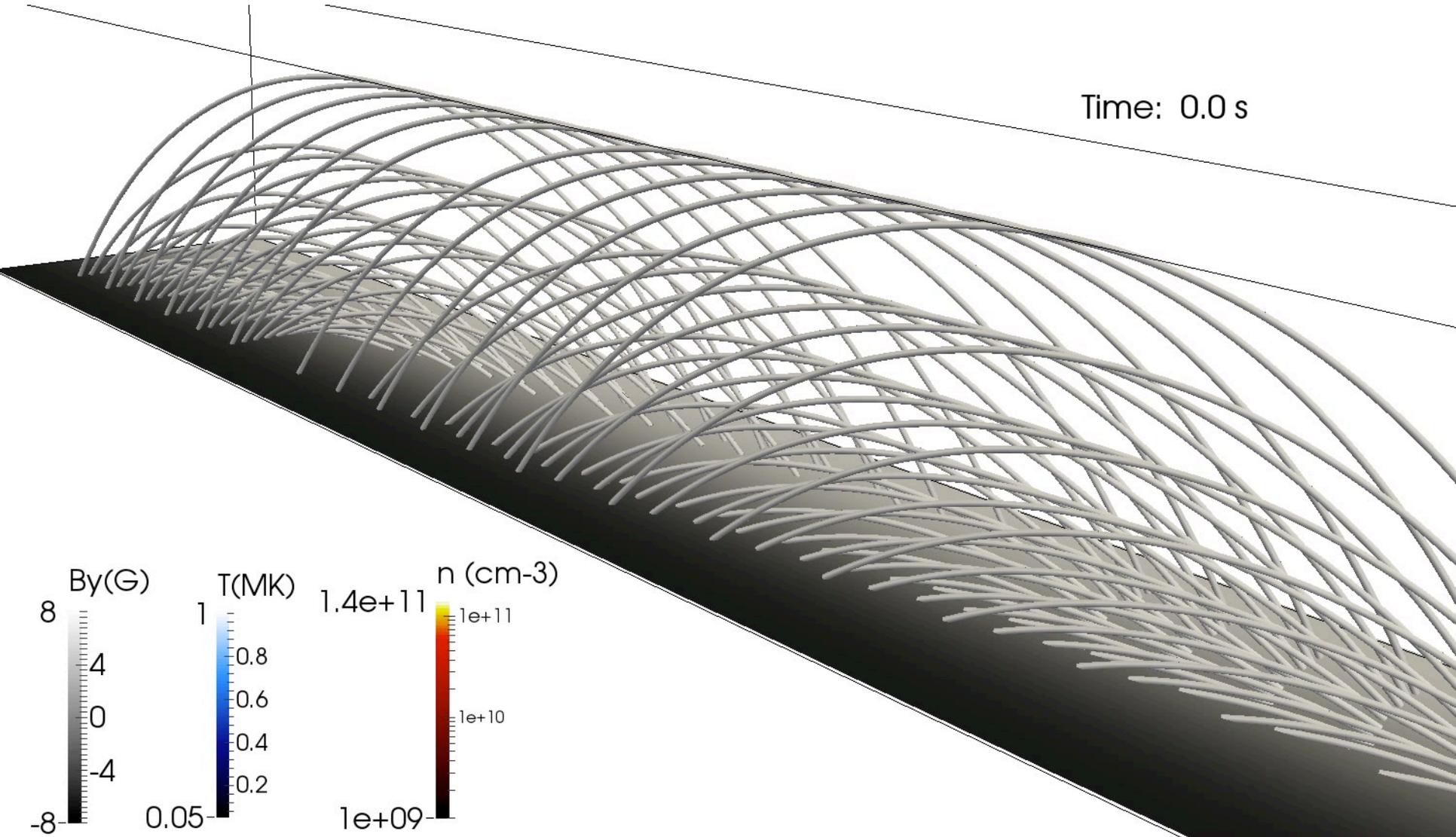
converging & anti-shearing motion

the direction in which magnetic shear is reduced
(toroidal component is reduced)

top view

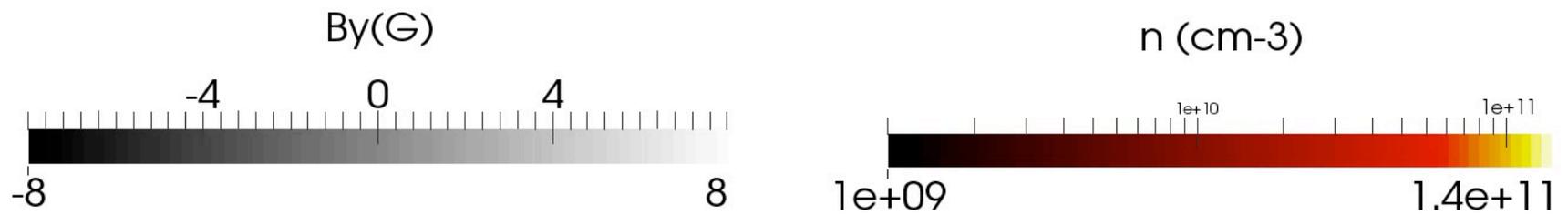


Result: side view

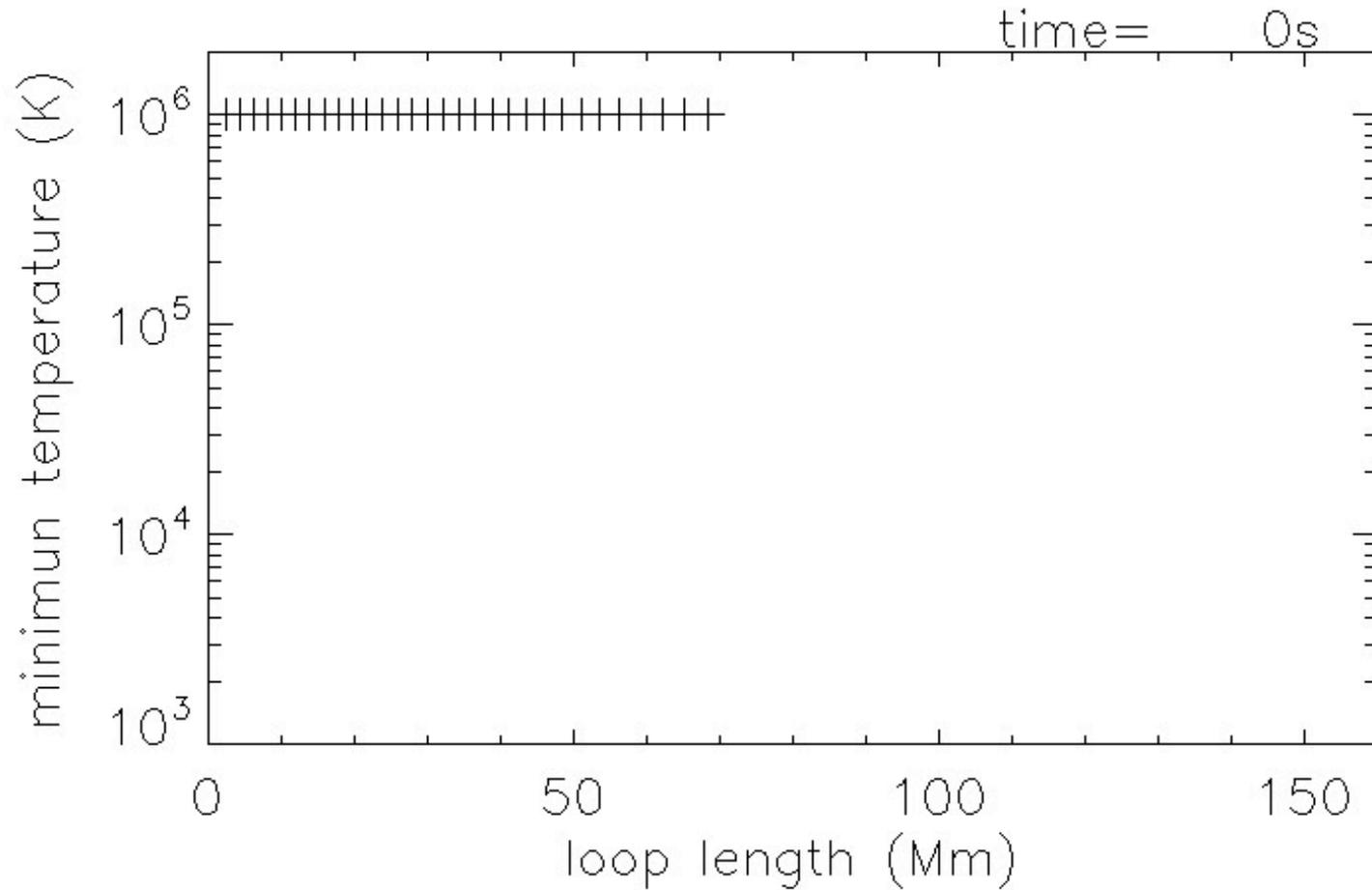


Result: top view

Time: 0.0 s

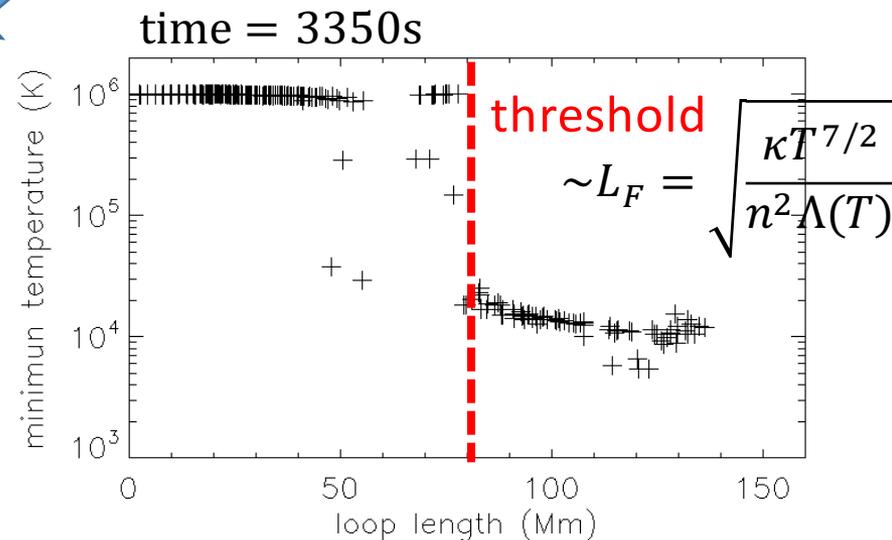
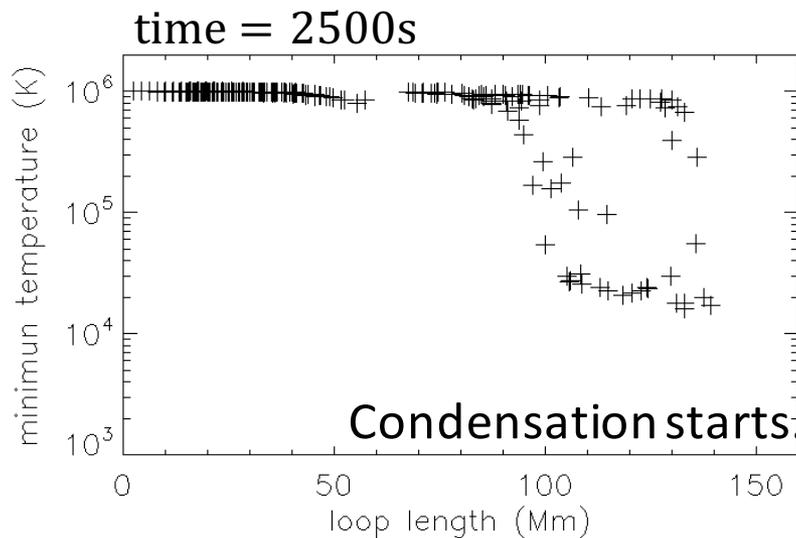
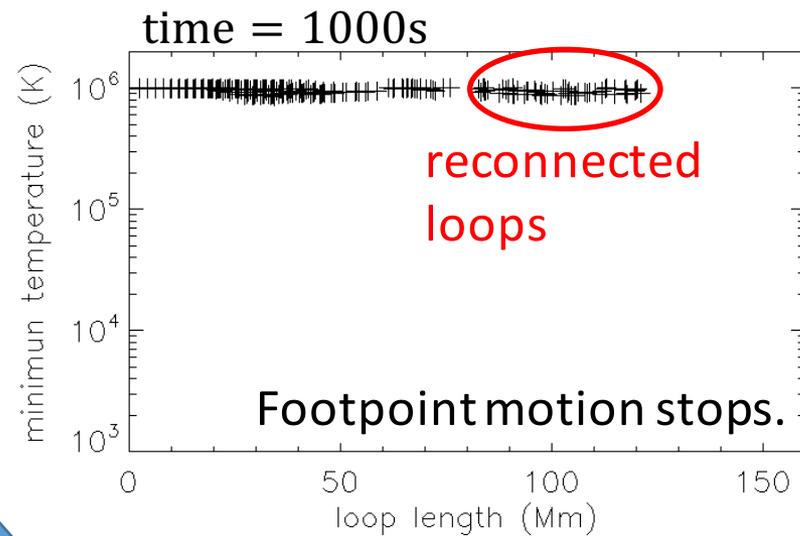
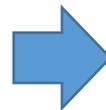
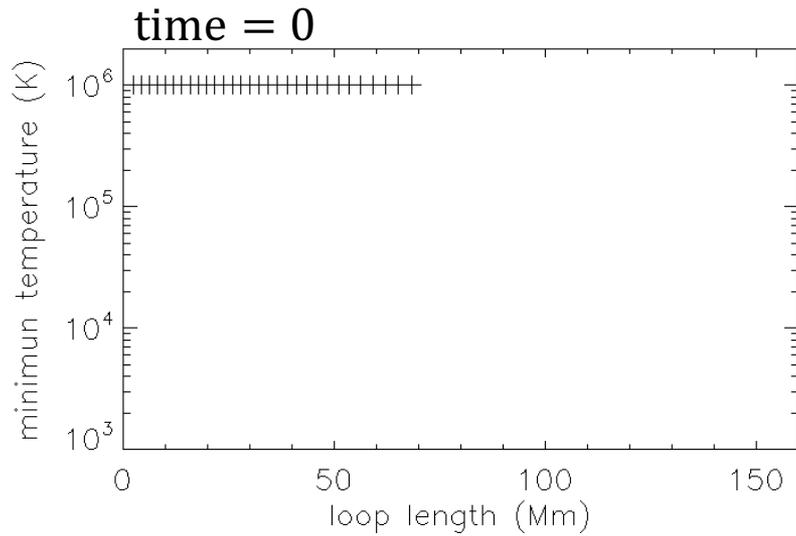


Loop length & temperature



Sufficiently long loop  Radiative cooling is not compensated by conduction.

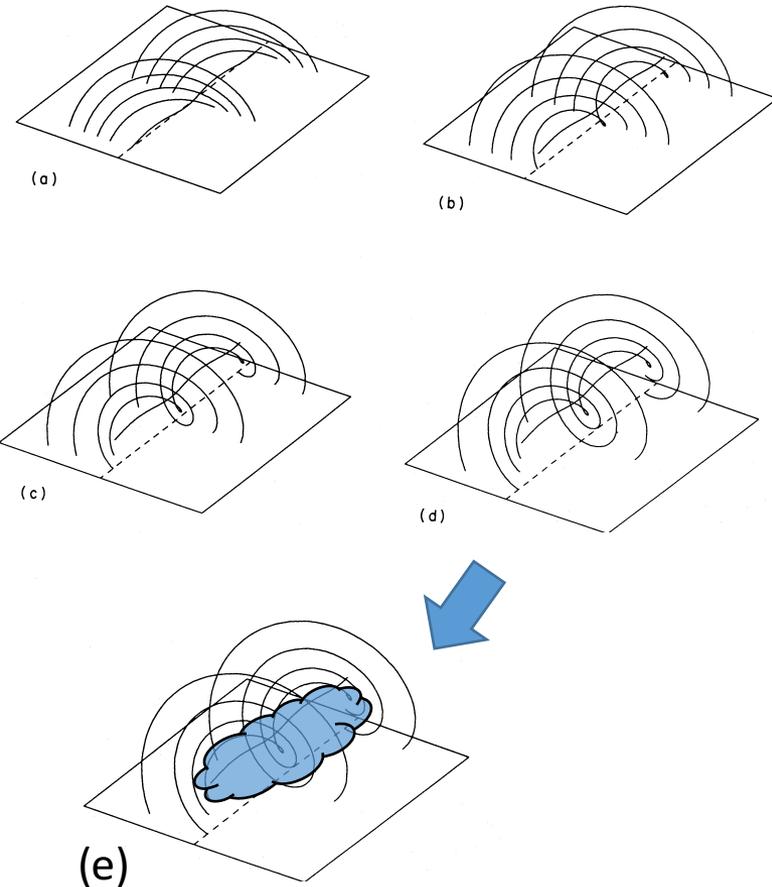
Loop length & temperature



Summary of our results

- **Reconnection-condensation scenario** has been demonstrated, in which radiative condensation is caused by the topological change of magnetic field due to reconnection, not evaporation.
- Sufficiently long reconnected loops suffer from radiative condensation.
- Converging motion can lead not only the formation of flux rope **but also the radiative condensation**.

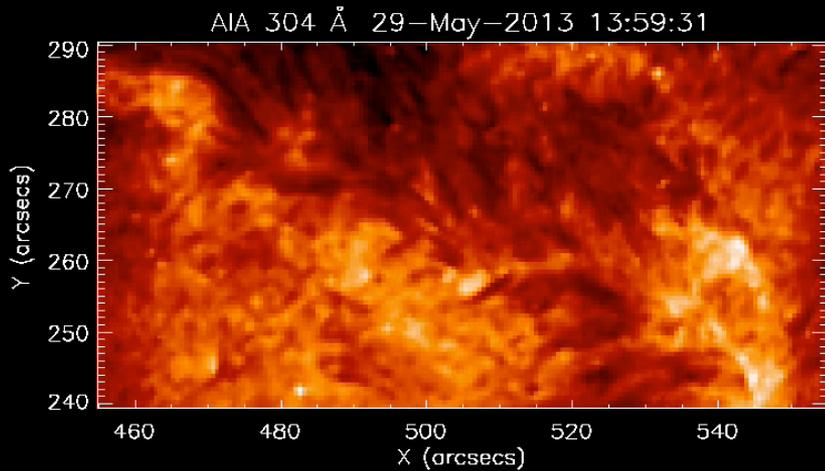
Van Ballegooijen & Martens (1989)
converging motion



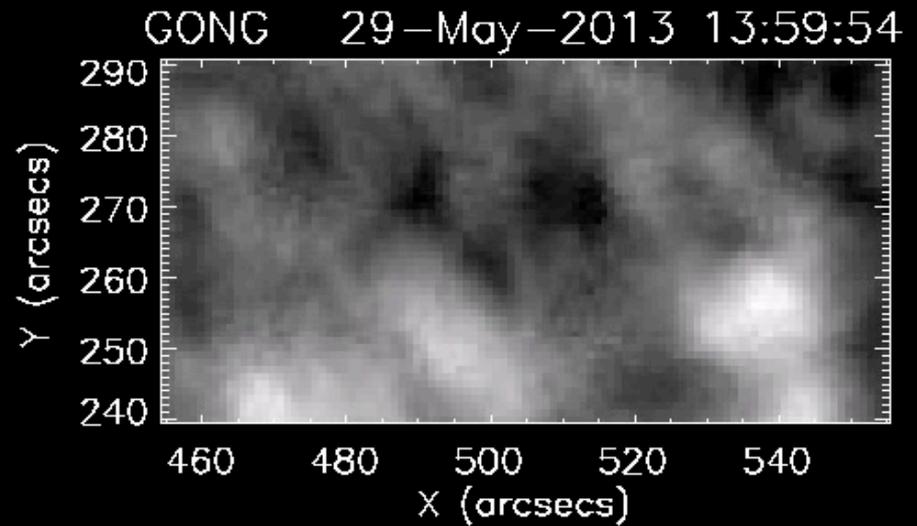
Observation of filament formation 1/2

Yang et al. (2016)

SDO/AIA 304 Å



GONG H α

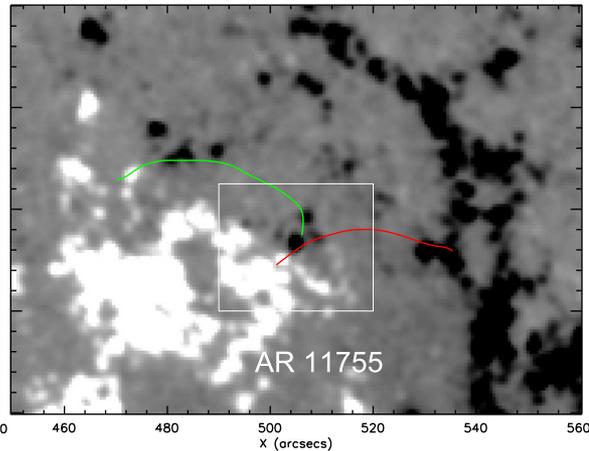


Reconnection \rightarrow Condensation

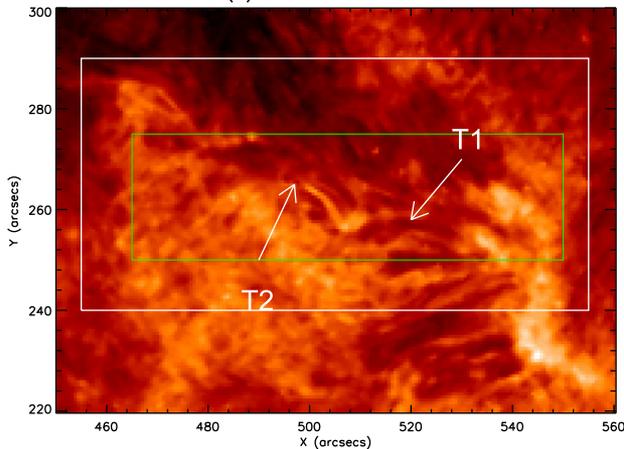
Observation of filament formation 2/2

Yang et al. (2016) SDO/AIA 304 & HMI

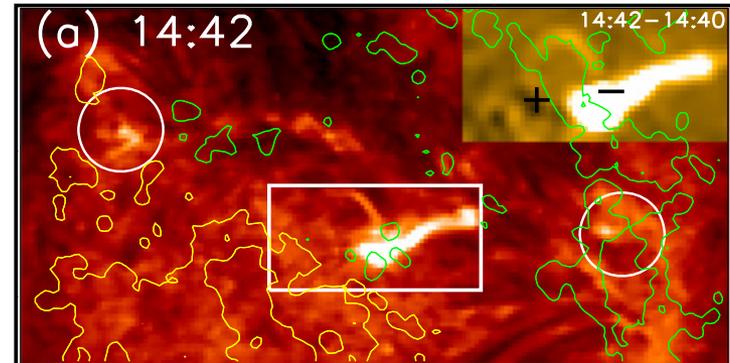
Cancellation



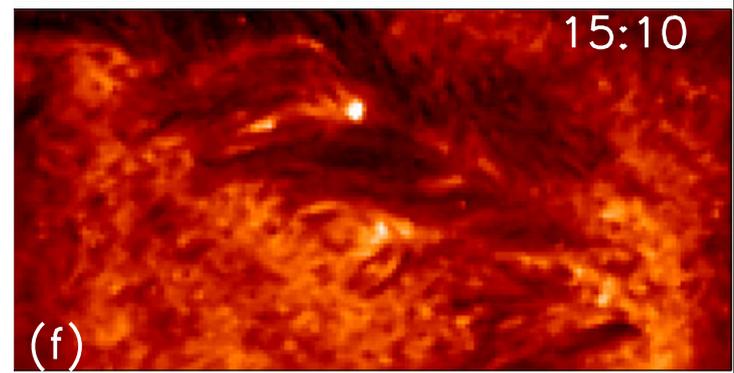
(a) AIA 304 Å 14:34 UT



Reconnection

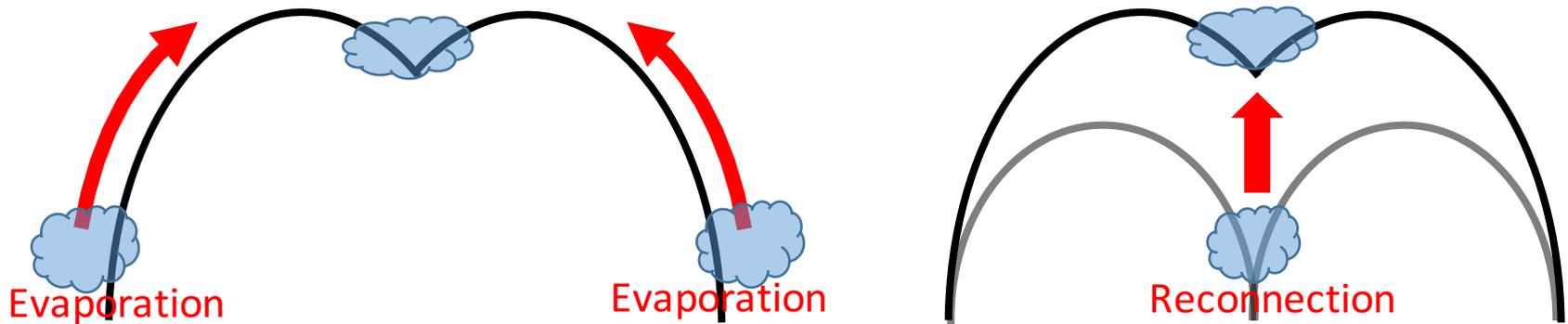
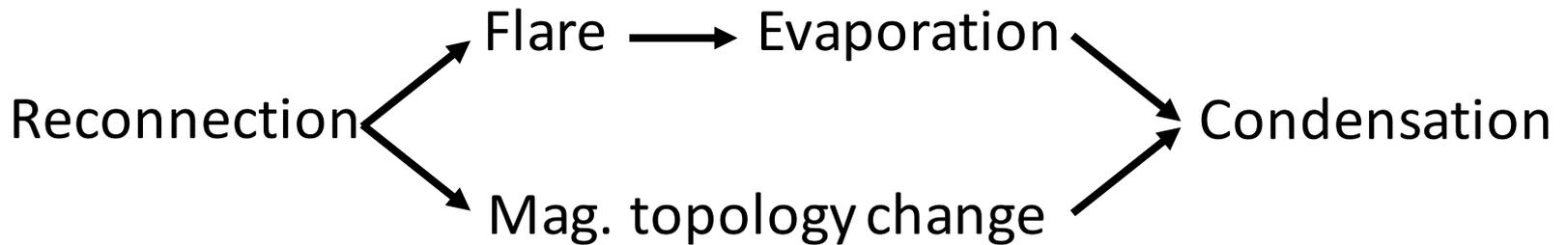


Filament formation (condensation)



Discrimination of models by observation

- From which route does the high density plasmas come ?
- How is the thermal evolution ?



Phase-mixing in Prominence 1/2

Cross-field waves

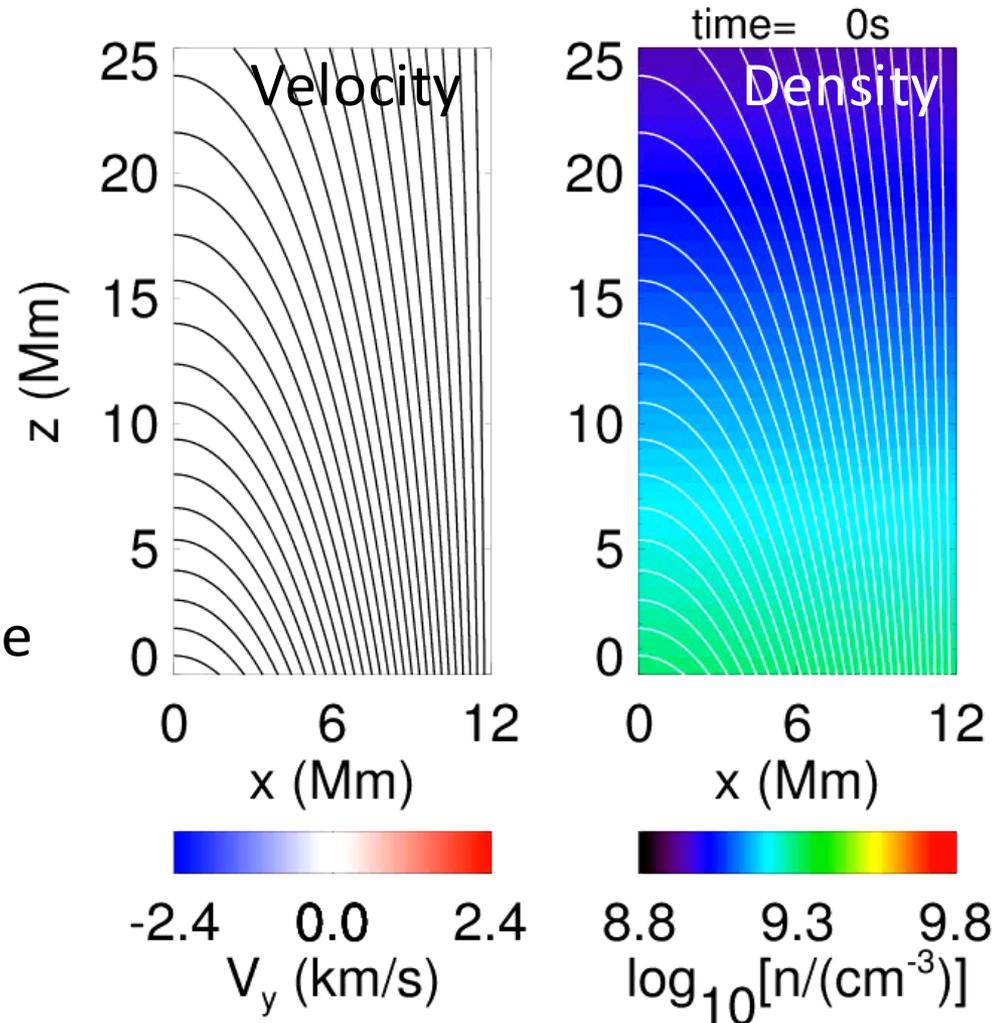


~~Fast mode ?~~

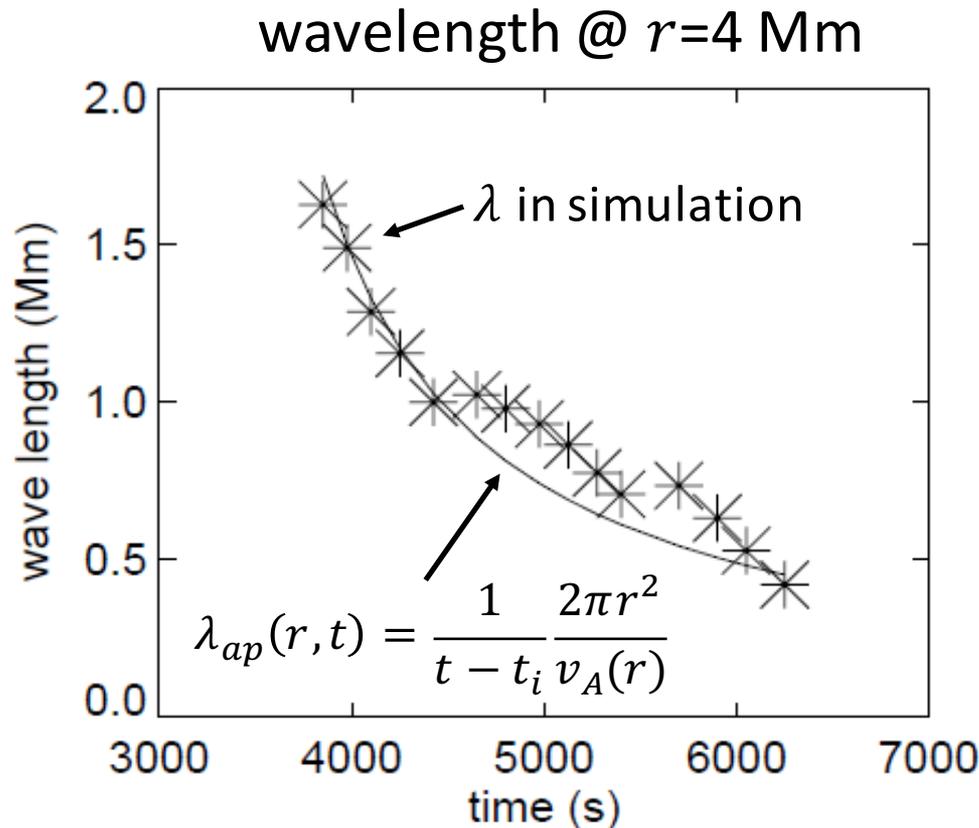


Phase mixing of standing
Alfven waves in the flux rope

Kaneko et al. (2015)



Phase-mixing in Prominence



Kaneko et al. (2015)

Phase-mixing



- Cross-field **wavelength decreases with time**.
- Eigen frequencies vary across magnetic field.

High spatiotemporal resolution of IRIS has advantage to detect phase-mixing.

Conclusion

- We demonstrate **reconnection-condensation model** by 2D & 3D MHD simulation including radiative cooling & thermal conduction
- In 3D simulation, although thermal conduction along toroidal magnetic components are effective, radiative condensation is triggered **in the dip of sufficiently long helical magnetic field**.
- IRIS multiwavelength observation with high spatiotemporal resolution has advantage for detect the flows & waves associated with prominence condensations.

Papers:

- ✓ T. Kaneko & T. Yokoyama (2015), ApJ, 806, 115
- ✓ T. Kaneko, M. Goossens, R. Soler, J. Terradas, T. Van Doorselaere, T. Yokoyama & A. N. Wright (2015), ApJ, 812, 121