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Topological solitons in confined system

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INTRODUCTION: BELAVIN-POLYAKOV SOLITONS

The two-component fields are represented by vector

$$\psi = \begin{pmatrix} \psi^1 \\ \psi^2 \end{pmatrix} \quad (1)$$

The components $\psi^{1,2}$ may be written as $\psi^{1,2} = \sqrt{(N^{1,2})}z^{1,2}$, where $N^{1,2} = \psi^{1,2\dagger}\psi^{1,2}$ and $z^{1,2} = z^{1,2}(x, y, t)$ are complex numbers. In what follows we consider $N^{1,2} = N = \text{const.}$ Then we have $\psi = \sqrt{N}\zeta$, where

$$\zeta = \begin{pmatrix} z^1 \\ z^2 \end{pmatrix}, \quad (2)$$

with the normalization condition $z^1 z^1 + z^2 z^2 = 1$.

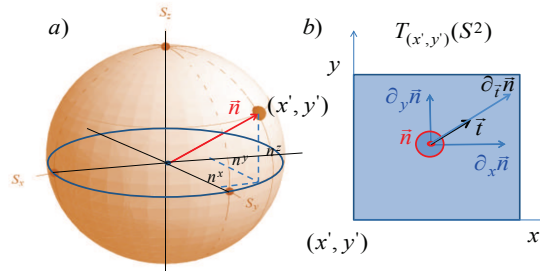


FIG. 1: a) Poincaré sphere with $\vec{n} = (n_x, n_y, n_z)$ in the point (x', y') representing the quantities defined in eq. (3); b) Tangent plane $T_{(x', y')}(S^2)$ from which the condition eq. (8) between a vector \vec{n} and their derivatives may be read off.

We next make transformation

$$n^{s_i} = n^{s_i}(x, y) = \zeta^\dagger \tau^{s_i} \zeta, \quad (3)$$

where τ_{s_i} , with $s_i = x, y, z$, are the usual Pauli matrices. For the components of the vector $\vec{n} = (n^x, n^y, n^z)$ hold $\sum_{s_i} (n^{s_i})^2 = 1$. The transformation is a mapping from a complex plane to a 2-sphere, that is, $\vec{n} : \mathbb{R}^2 \rightarrow S^2$. The sphere described by the tip of the vector \vec{n} can be identified with the Poincaré (or Bloch) sphere. The points in the Poincaré sphere may be interpreted as pseudo-spins [8]. Figure 1 describes the above statements.

Now, we follow the discussion in [9] and take following boundary condition

$$\lim_{\rho \rightarrow \infty} n(x, y) = (0, 0, 1). \quad (4)$$

for $\rho = \sqrt{x^2 + y^2}$. Direct consequence of (4) is possibility of compactification from \mathbb{R}^2 to S^2 (we consider only one chart of the atlas)

$$\vec{n}_i : S^2 \rightarrow S^2. \quad (5)$$

All mappings defined by 5 form configuration space: $Q = \{\vec{n}_i\}$. As homotopy group of configuration space $\pi_1(Q)$ is isomorphic to homotopy group $\pi_2(S^2)$

$$\pi_1(Q) = \pi_2(S^2) = \mathbb{Z} \quad (6)$$

it is itself union of homotopy classes Q_{N_i}

$$Q = \bigcup_{N_i} Q_{N_i}. \quad (7)$$

A class is determined by $N_i \in \mathbb{Z}$ which is degree of the mapping $\vec{n}_i : N_i = \text{deg}(\vec{n}_i)$.

At a fixed point on the Poincaré sphere (x', y') ([1]) (see figure 1b and expression (9))

$$\partial_\alpha n^{s_1}(x, y) = \varepsilon_{\alpha\beta} \varepsilon_{s_1 s_2 s_3} n^{s_2}(x, y) \partial_\beta n^{s_3}(x, y) \quad (8)$$

which means fixing covariant derivative \mathcal{D}_α ($\alpha = x, y$ are now coordinates of tangent plane), i.e. demanding parallel transport on a curve on the sphere

$$\mathcal{D}_\alpha \vec{n}(x, y) \equiv \partial_\alpha \vec{n}(x, y) - \varepsilon_{\alpha\beta} \partial_\beta \vec{n}(x, y) \times \vec{n}(x, y) = 0. \quad (9)$$

The field ζ in terms of n^{s_i} reads ([6]) in circular polarization basis

$$\zeta = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{1+n^z} \\ \frac{n^x + in^y}{\sqrt{1+n^z}} \end{pmatrix}. \quad (10)$$

Comparing last vector with (1) one sees that $z^1 = \sqrt{1+n^z}$ and $z^1 z^2 = (n^x + in^y)$.

In what follows we chose polar coordinates for local field $\zeta = \zeta(x, y) = \zeta(r \cos \phi, r \sin \phi)$. From the other side $\vec{n} = (\sin \Theta \cos \Phi, \sin \Theta \sin \Phi, \cos \Theta)$ where $\Theta = \Theta(r, \phi)$ and $\Phi = \Phi(r, \phi)$ are global coordinates corresponding to polar and azimuthal angles of Poincaré sphere. Hence

$$\zeta = \begin{pmatrix} \cos \frac{\Theta}{2} \\ \sin \frac{\Theta}{2} e^{i\Phi} \end{pmatrix}. \quad (11)$$

Belavin and Polyakov showed that equality (8) implies Cauchy-Riemann conditions on complex function $u = z^2/z^1$

$$\partial_r u = -\frac{i}{r} \partial_\phi u. \quad (12)$$

In global coordinates $u = \tan \frac{\Theta}{2} e^{i\Phi}$ giving

$$\frac{d\Theta}{\sin \Theta} = -\left(\frac{N_i}{r} + i \frac{d\Phi(r)}{dr}\right) dr \quad (13)$$

where we used

$$\Phi(r, \phi) = N_i \phi + \Phi(r). \quad (14)$$

M being topological charge and $\Theta = \Theta(r)$.

1) Standard 2D Belavin-Polyakov solitons ([9]) are then solution of (13) for $\Phi(r) = \Phi_0$

$$\Theta(r) = 2 \arctan \left(\frac{r^{N_i}}{R} \right). \quad (15)$$

where R is constant describing soliton radius. One can check that $\lim_{r \rightarrow 0} \Theta(r) = 0$ and $\lim_{r \rightarrow \infty} \Theta(r) = \pi$ (a meridian on Poincaré sphere) which are necessary conditions for Belavin-Polyakov soliton ([12]).

2) Another kind of solution exists if $\Phi(r)$ varies with r .

$$\Phi(r, \phi) = N_i \phi + \Phi(r). \quad (16)$$

and this topic will be considered in the next section.

SOLITONS ON TORUS: NLSE(GP) EQUATION

The fundamental group of the torus $S^1 \times S^1$ is $\mathbb{Z} \oplus \mathbb{Z}$, which is abelian, so its covering spaces are in 1 - 1 correspondence with the subgroups of this group, namely with $p\mathbb{Z} \oplus q\mathbb{Z}$,

$p, q = 0, 1, 2, 3, \dots$. These covering spaces are tori, if $p, q > 0$, cylinders, if $p = 0$ or $q = 0$ but not both, or \mathbb{R}^2 if $p = q = 0$. The solutions of the form (16) will be appropriate for the second case.

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