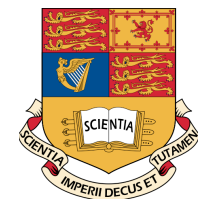


There and Back

RG Boomerangs
and
Spatially Modulated Deformations

Chris Rosen
with Donos, Gauntlett, and Sosa-Rodriguez

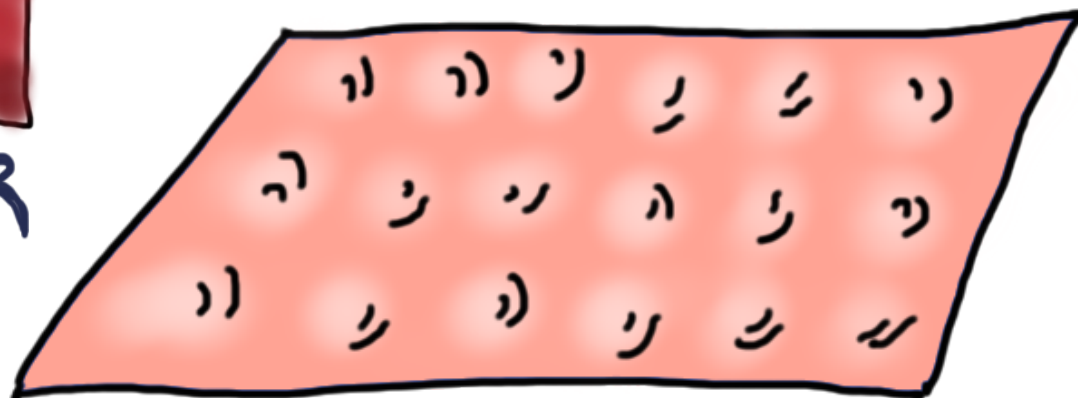
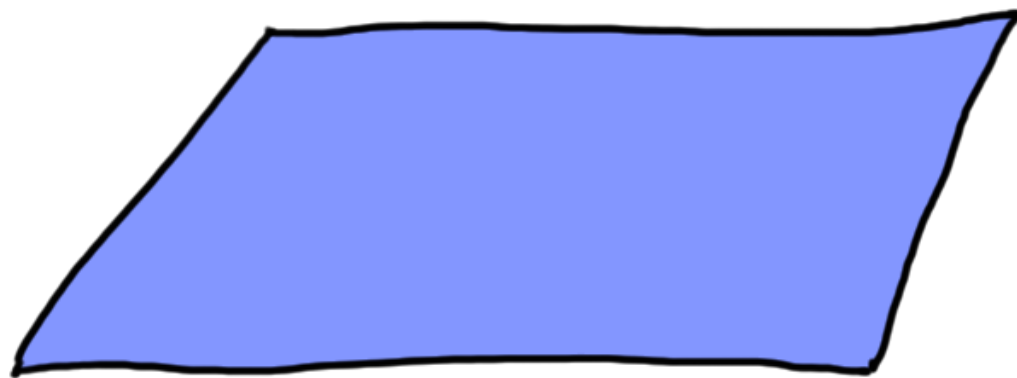


Imperial College
London

UV



IR



$$\mathcal{L}_{\text{CFT}} + m_{ij}^2 \phi^i \phi^j$$

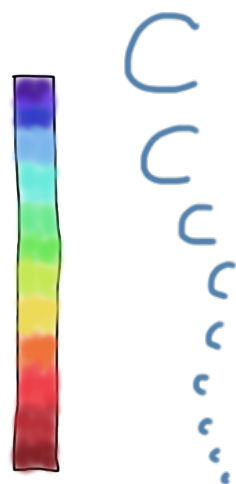


Why?



Chart the landscape of interacting fixed point theories
"Because it's there" -- Mallory

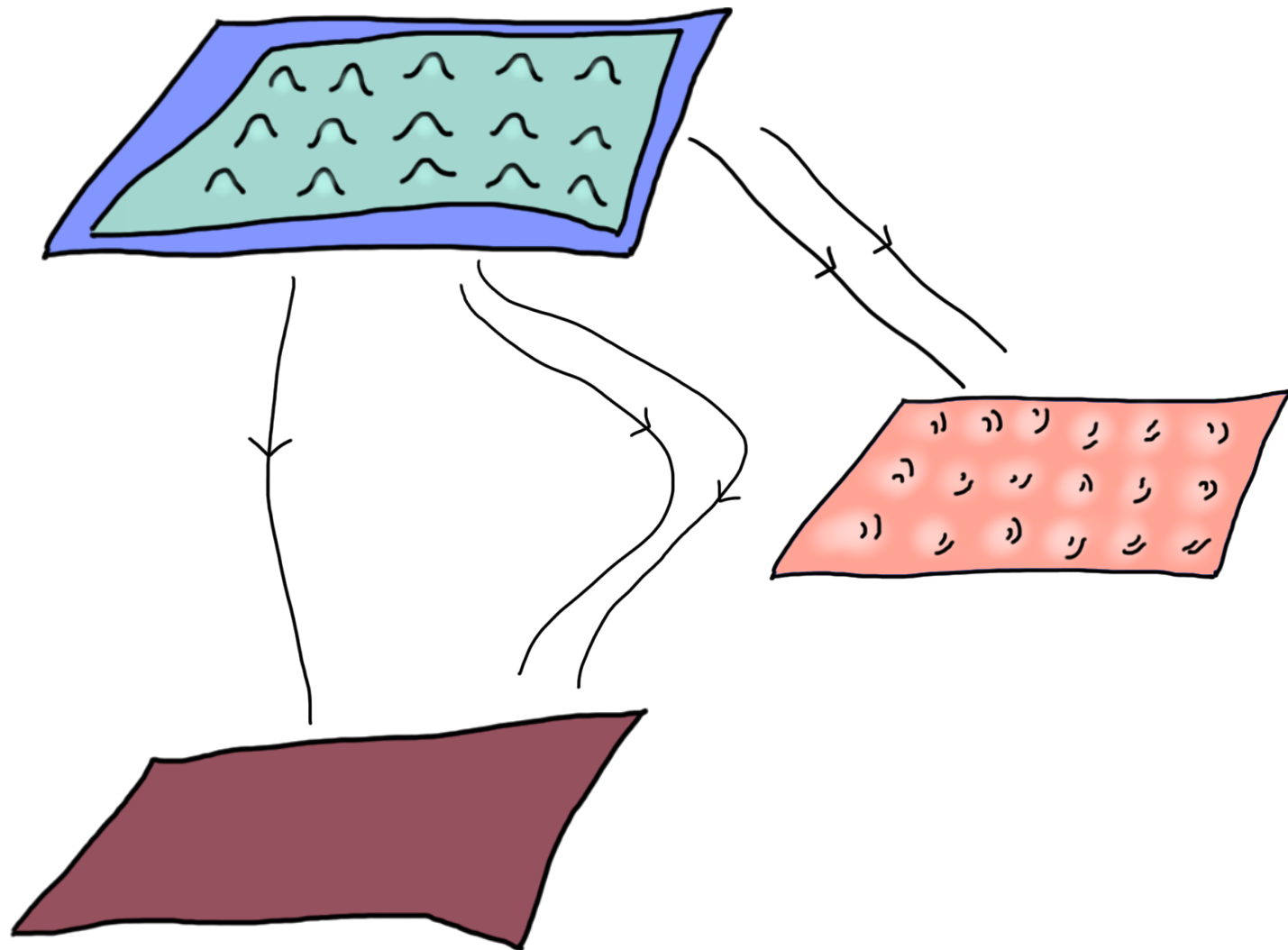
(but also to see what's possible, strengthen intuition, etc.)



Understand, criticize, develop RG results/lore

What is the validity/reach of a/F/c theorems,
entanglement measures, etc.?

$$\mathcal{L}_{\text{ABJM}} + m(\textcolor{red}{x})_{ij}^2 G^i G^j$$

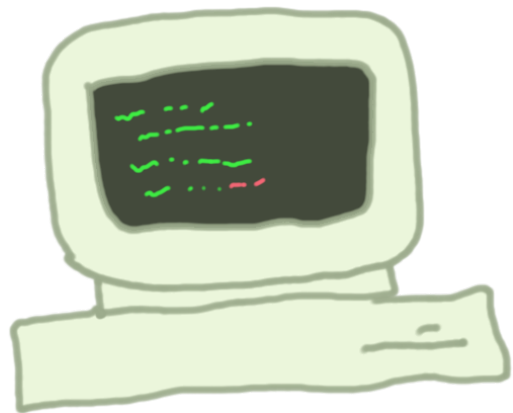


ABJM Theory deformed by a “lattice” as a QFT toy

“Conventional” Methods



- Perturbative Methods



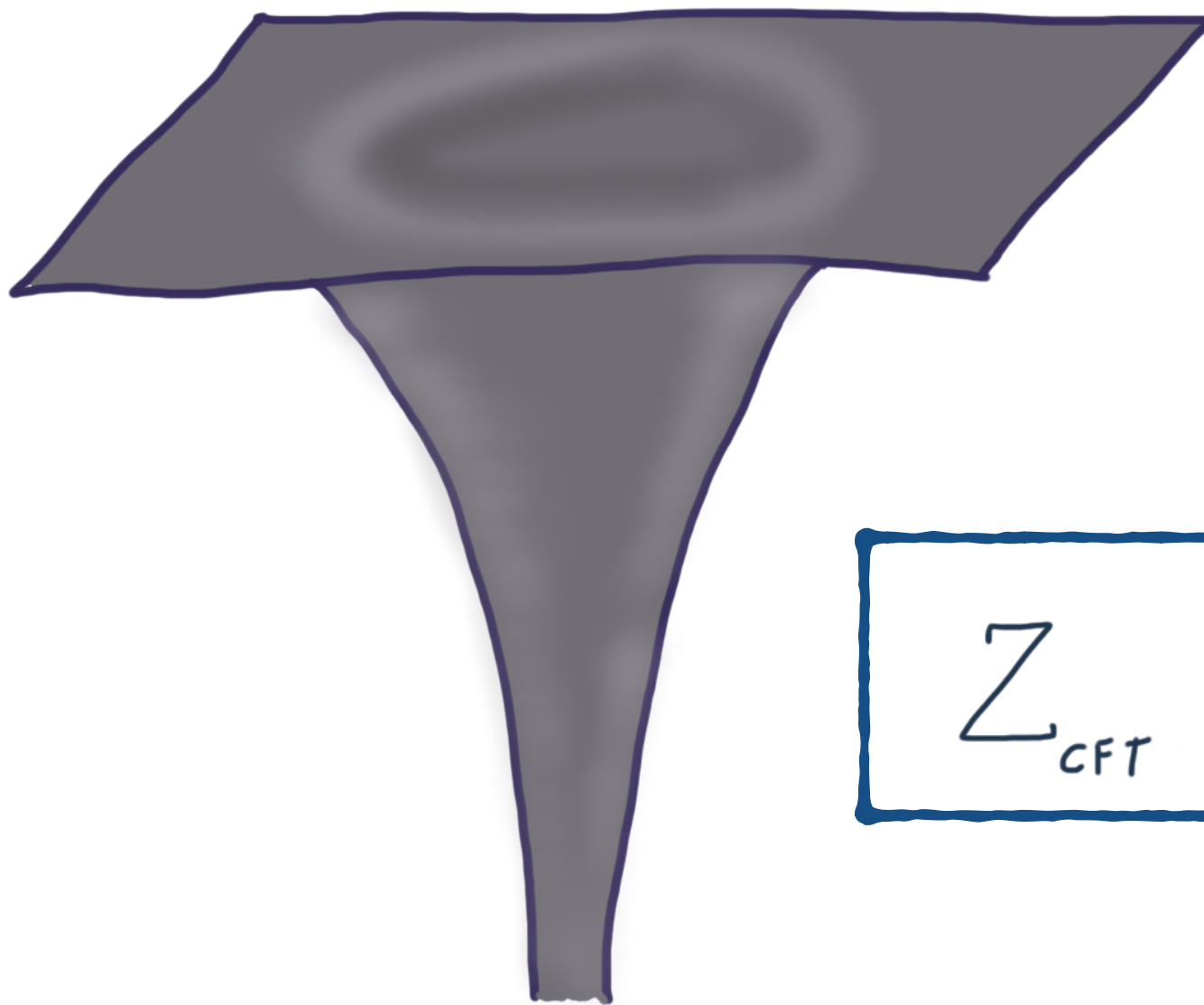
- Lattice Techniques

$$\{Q, \bar{Q}\} = P$$

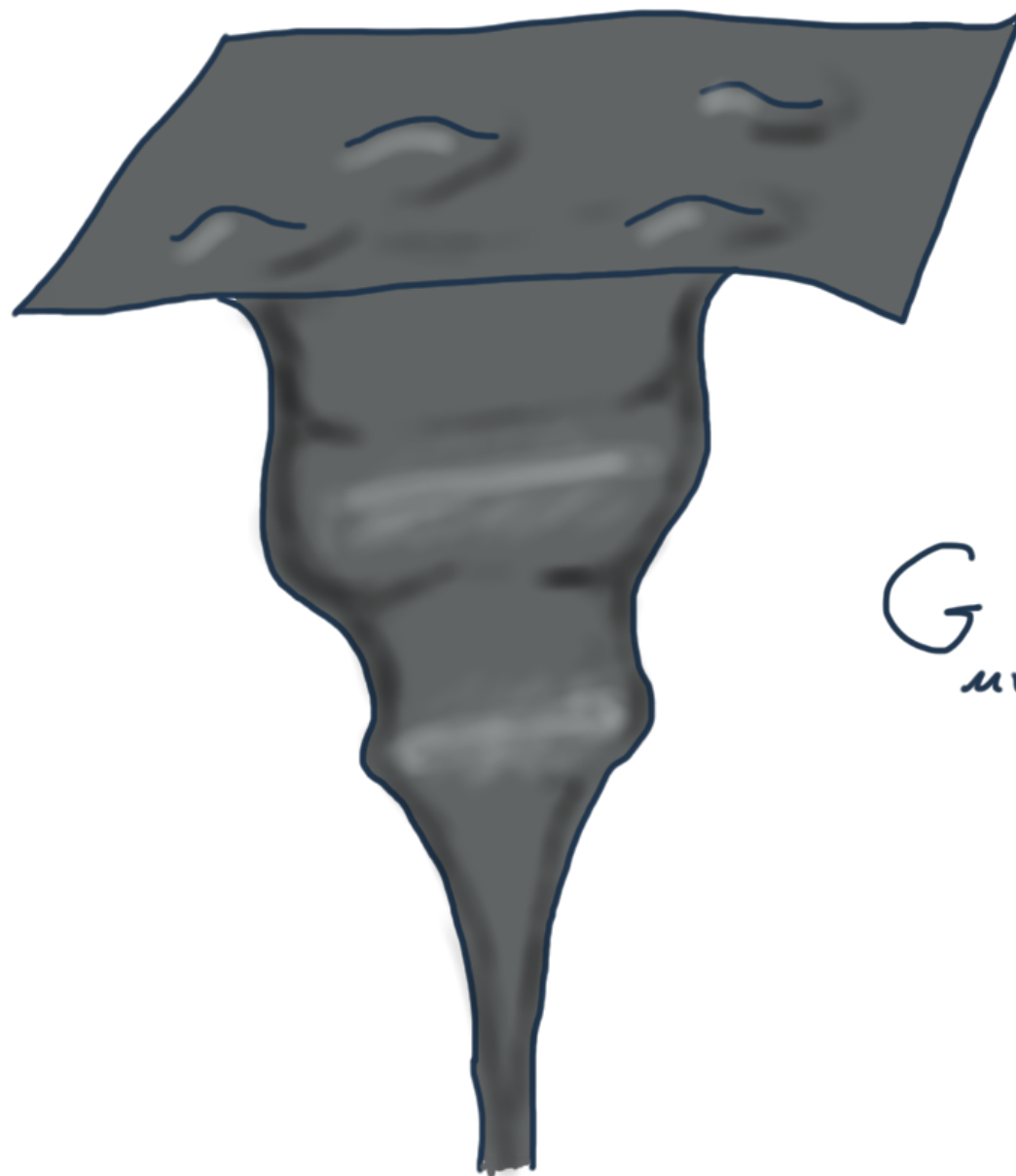
- SUSY Analysis

(... the list goes on ...)

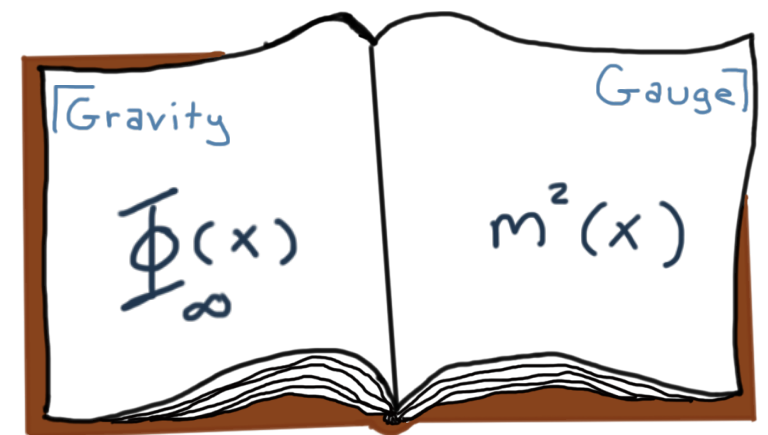
Use better variables...



$$Z_{\text{CFT}} \simeq e^{-S_{\text{SUGRA}}}$$



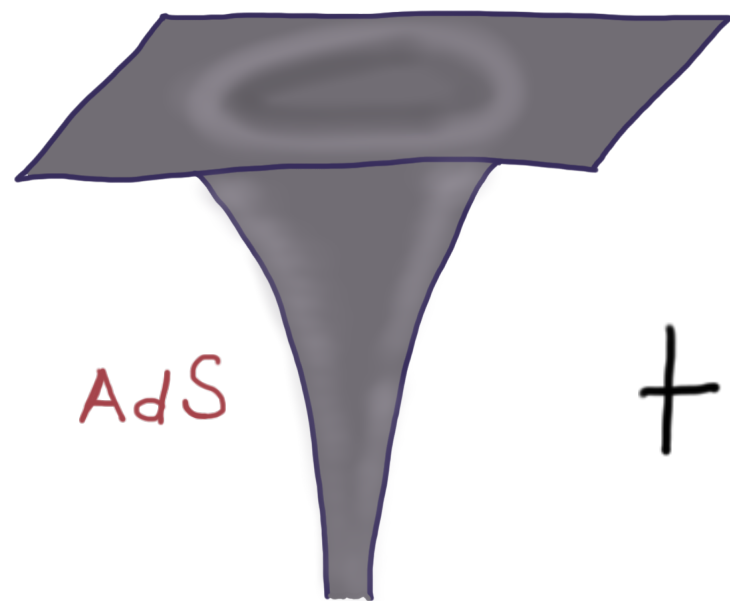
$$G_{\mu\nu} = T_{\mu\nu} + \Phi_{\infty}^i(x)$$



RG Geometrized

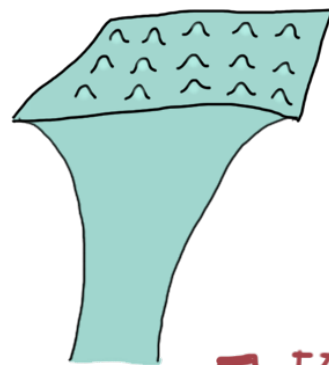
Ex. $S = \frac{1}{2k^2} \int d^4x \sqrt{-g} (R + 6 + \mathcal{L}_2) + \dots$ | $\alpha = 1, 2$
 $\mathcal{L}_2 = -G_{\alpha\bar{\beta}} \partial_\mu z^\alpha \partial^\mu \bar{z}^\beta - V(|z|)$ | $\mathcal{G} = U(1)^2$

$$z^\alpha = \gamma(r) e^{ikx^\alpha} \quad ds^2 = -U e^{-\chi} dt^2 + r^2 dx^\alpha dx^\alpha + \frac{dr^2}{U}$$



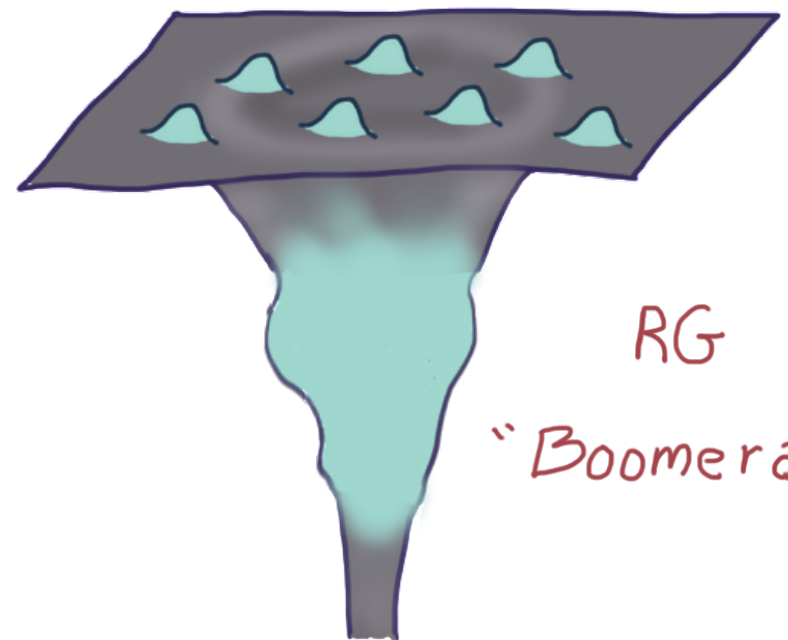
AdS

+



$$\Gamma_0 K_\nu e^{ikx}$$

=

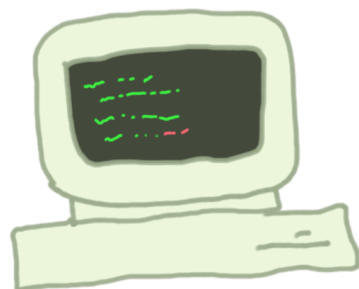
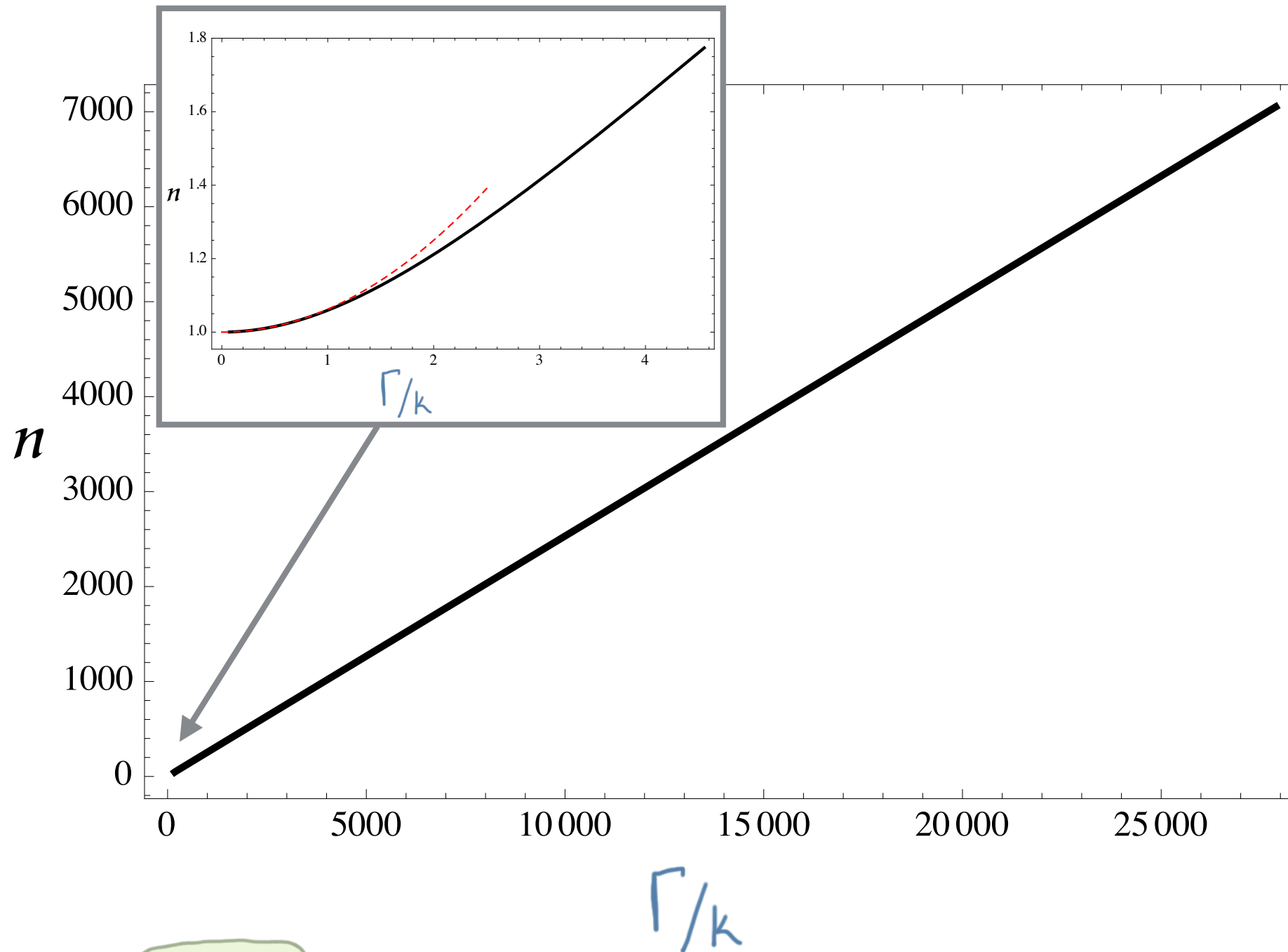


RG

"Boomerang"

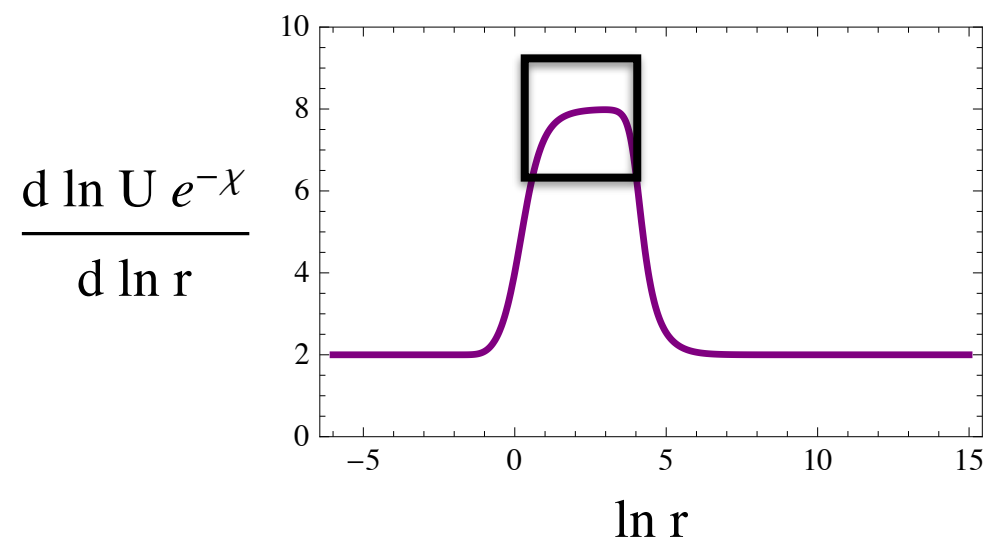
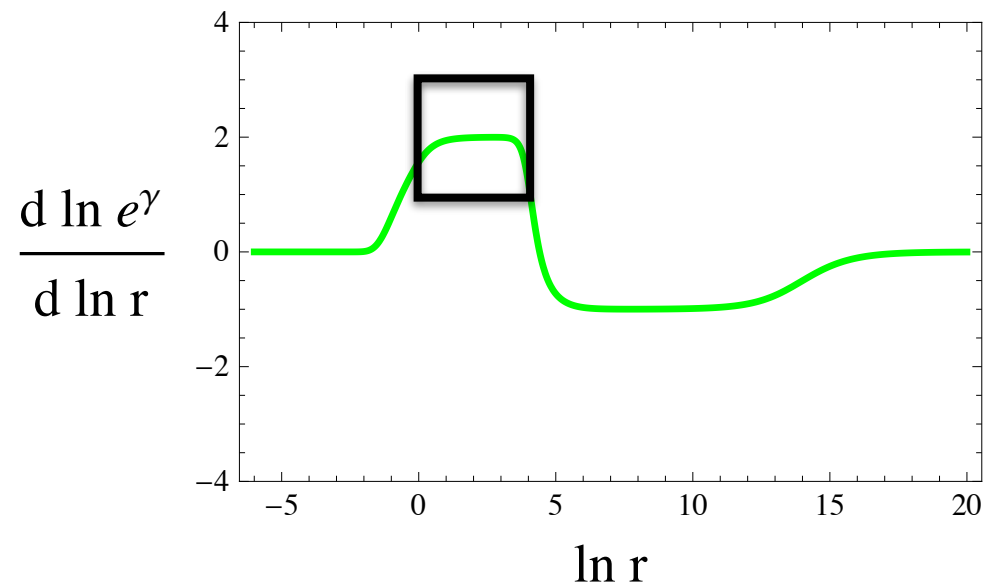
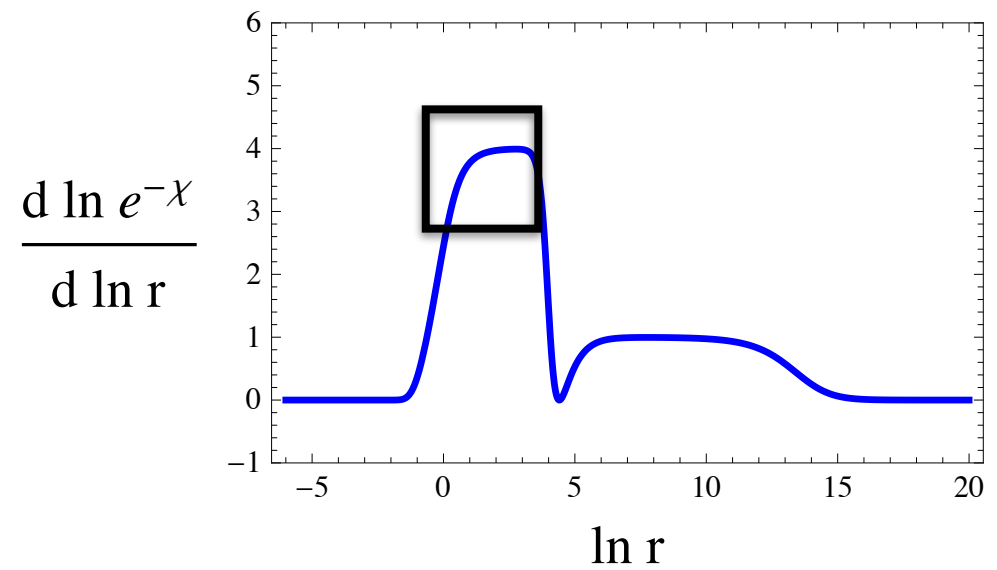
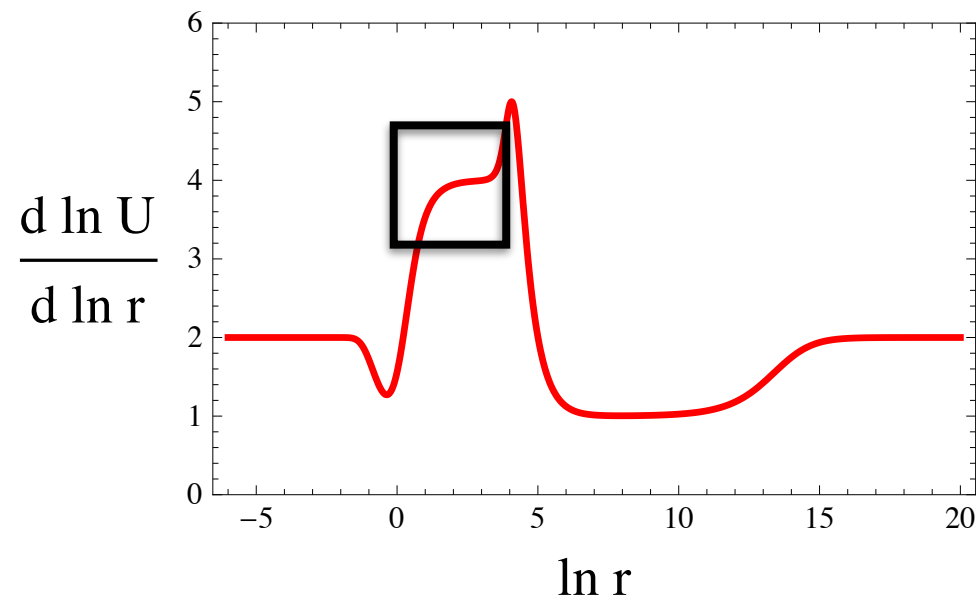
$$\eta \equiv \frac{C_{uv}}{C_{IR}} = \frac{L_{IR}}{L_{uv}} e^{\frac{1}{2}(\chi_{IR} - \chi_{uv})} = e^{\frac{1}{2}(\chi_{IR} - \chi_{uv})} \sim 1 + \# \left(\frac{\Gamma_0}{k} \right)^2 + \dots$$

Non-Perturbative Solutions

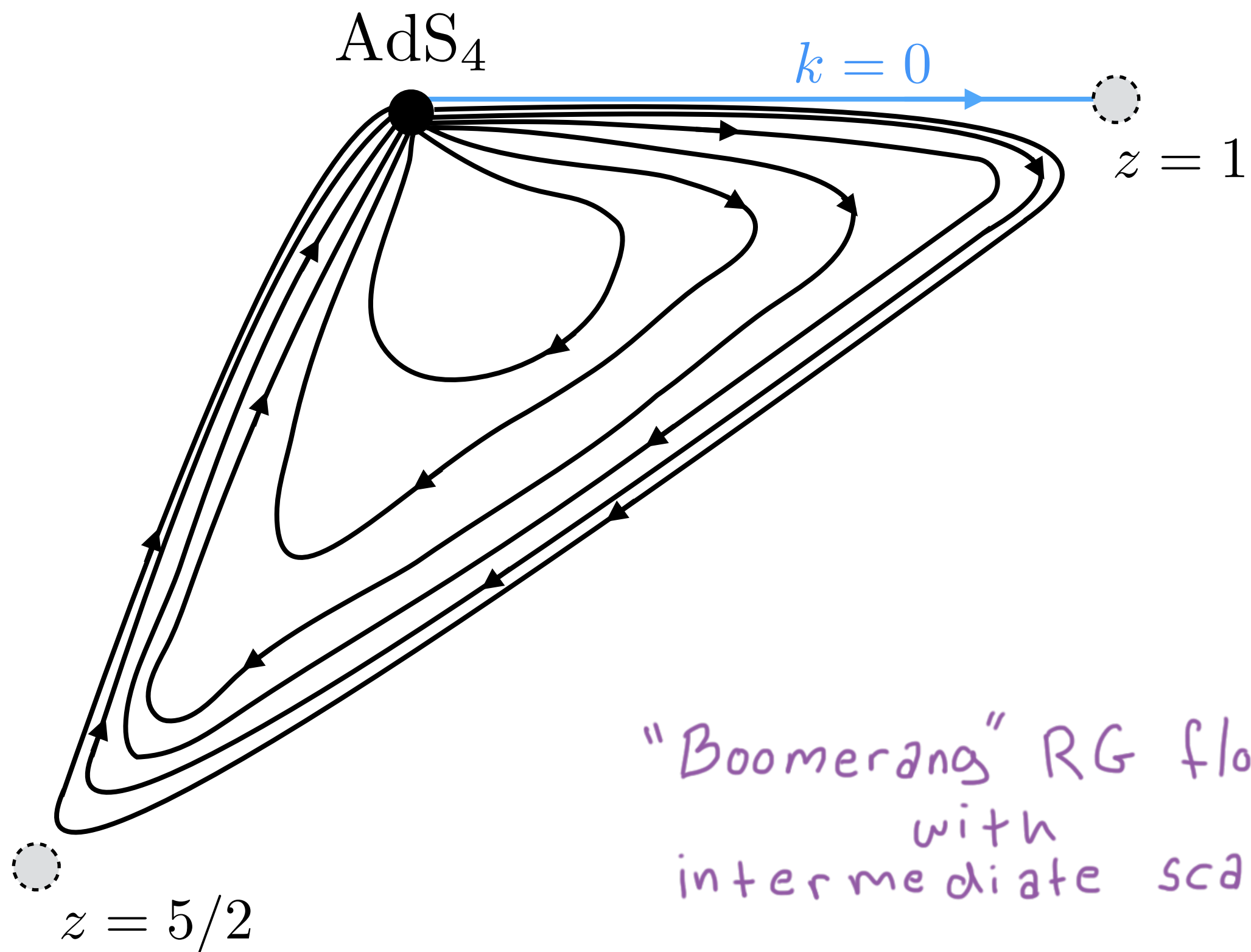


Boomerangs persist to BIG deformations

Non-Perturbative Solutions

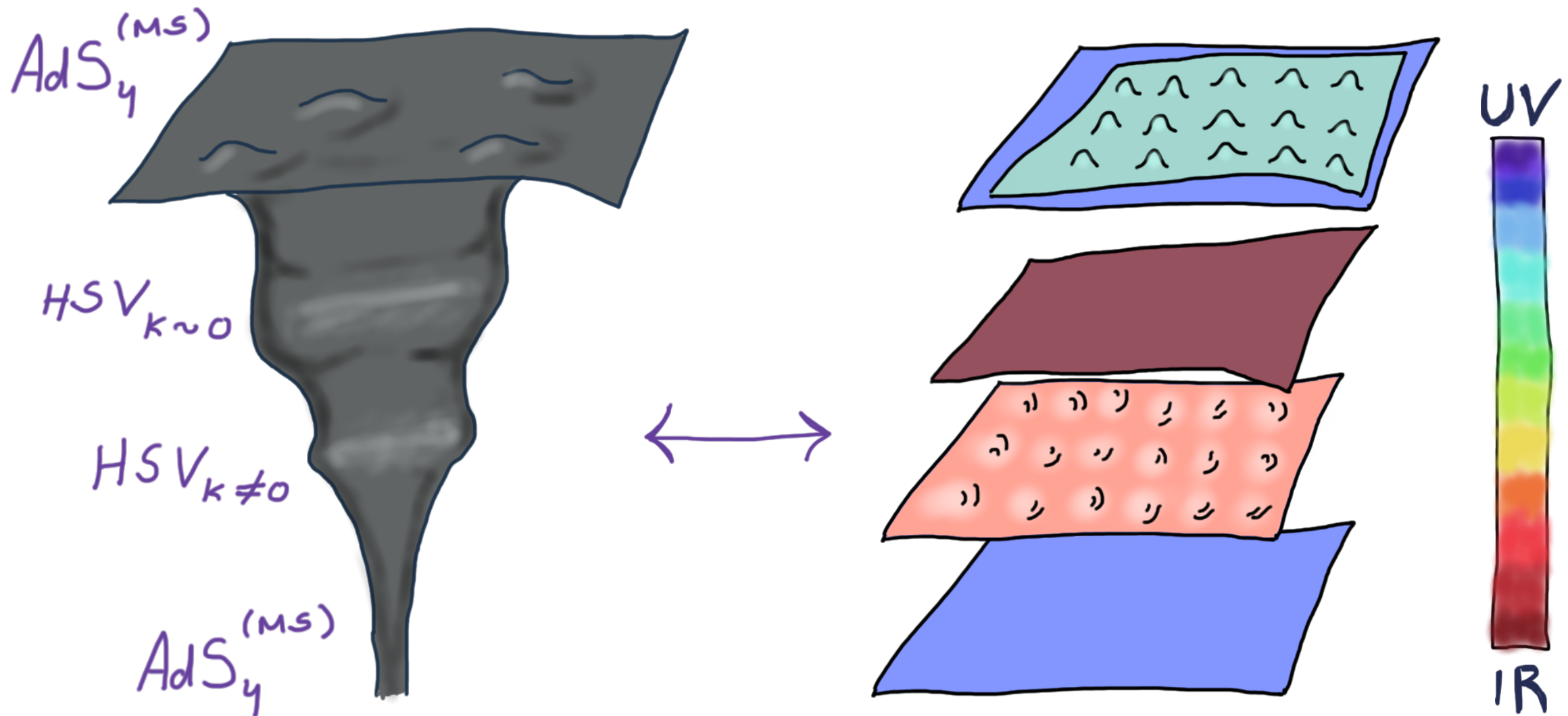


$$ds^2 = \frac{1}{\rho} \left(-\frac{1}{\rho^3} dt^2 + d\vec{x}^2 + d\rho^2 \right) \quad \text{HSV: } \phi=1, z=5/2$$



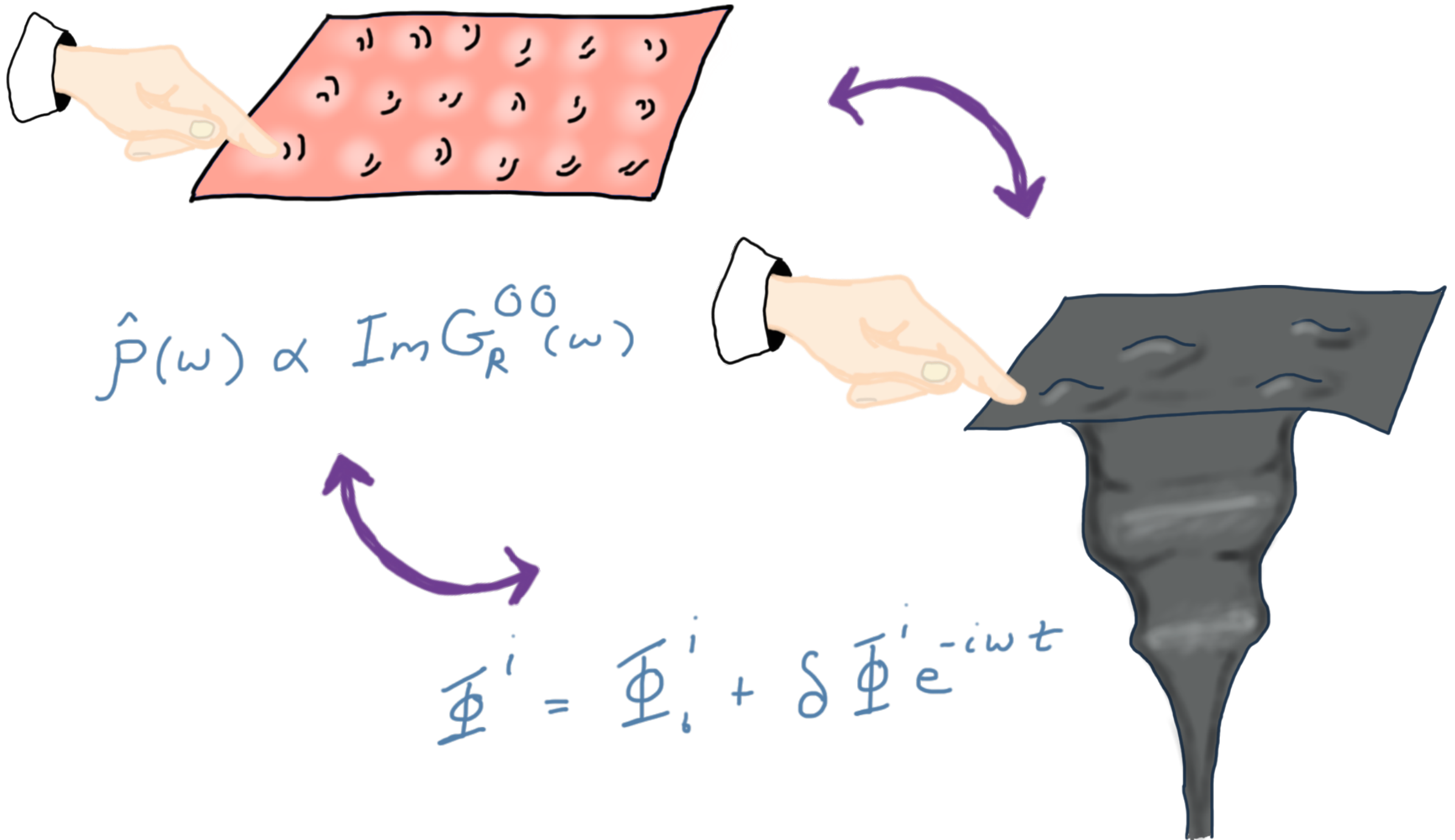
"Boomerang" RG flows
with
intermediate scaling

The Big Picture

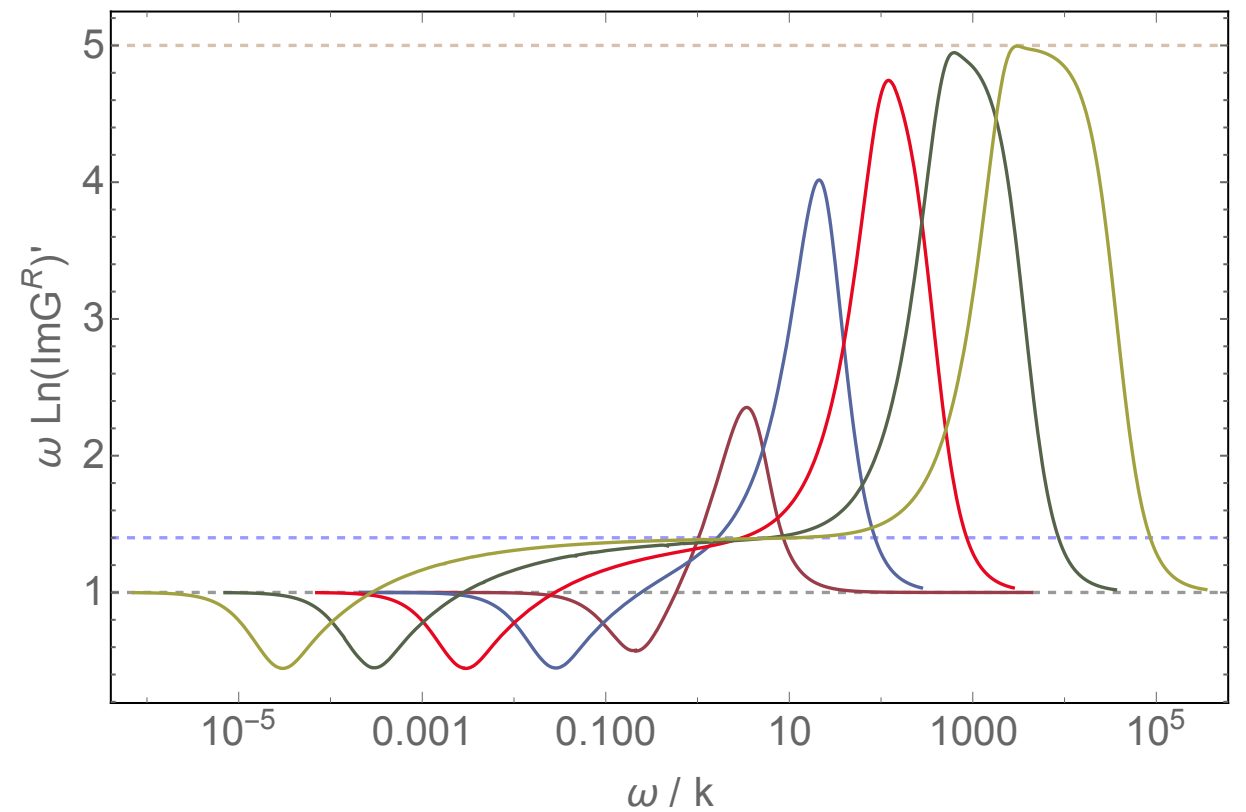
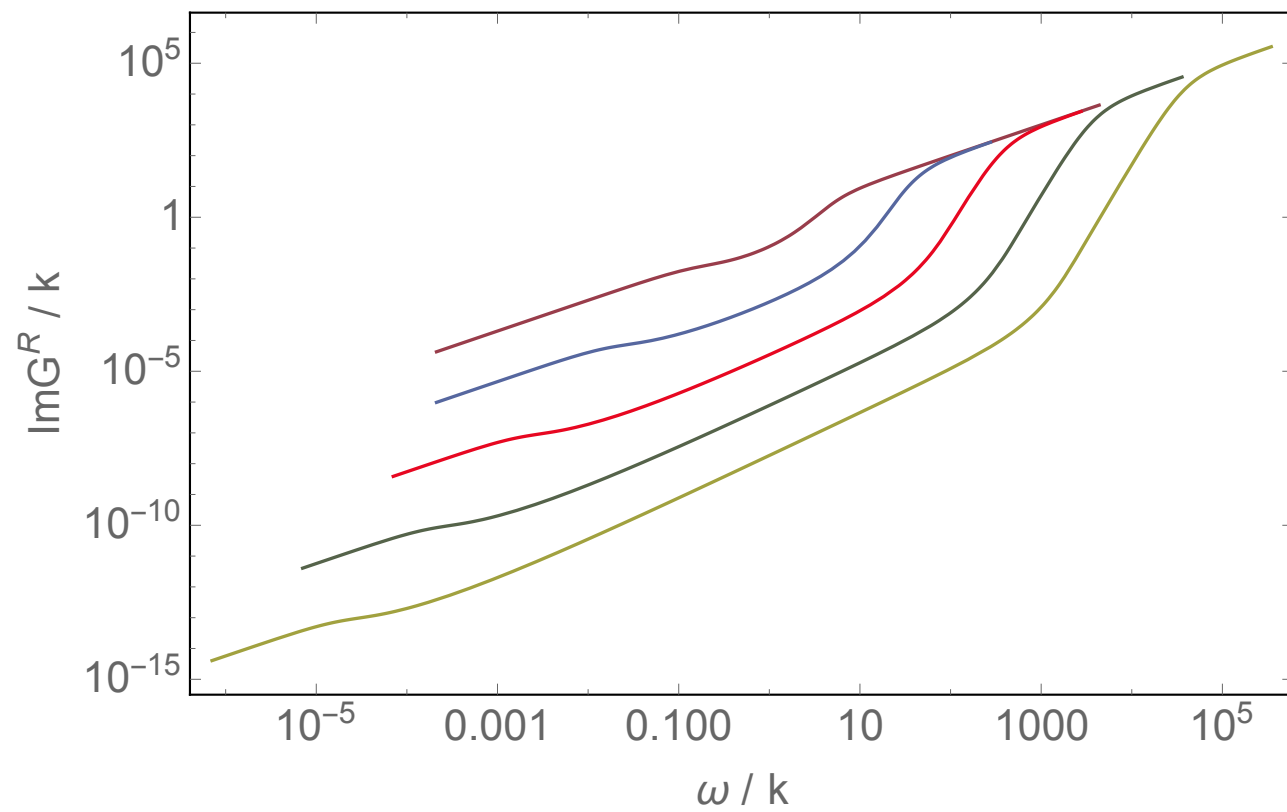
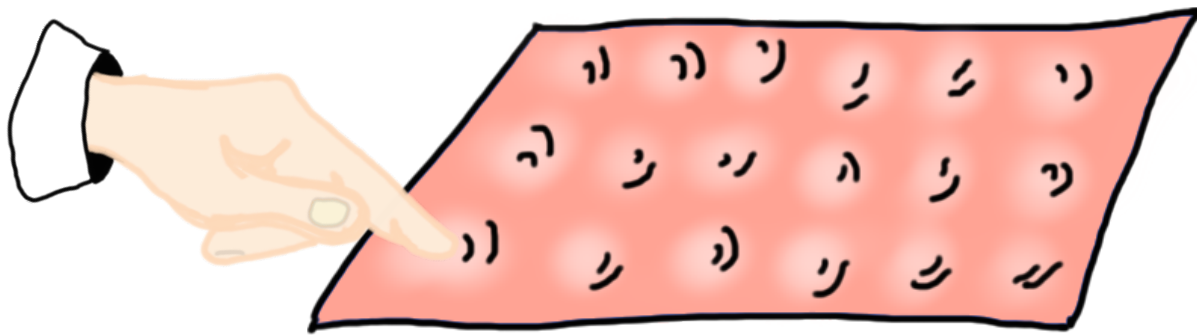


How is intermediate energy scaling realized in ABJM?

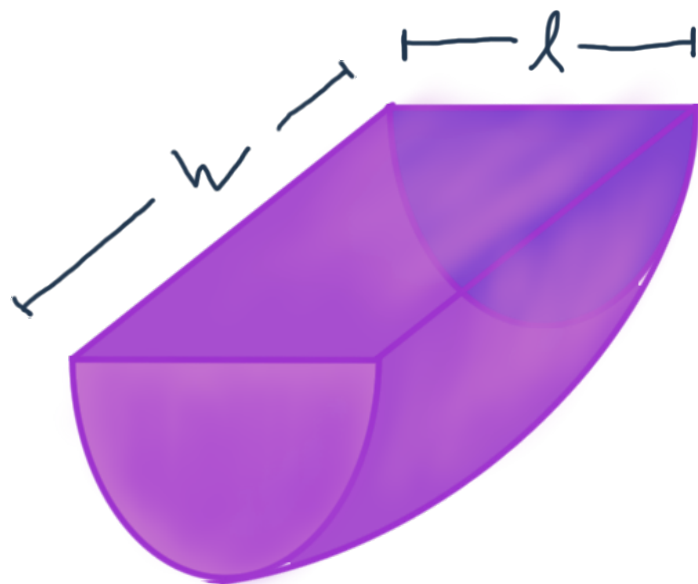
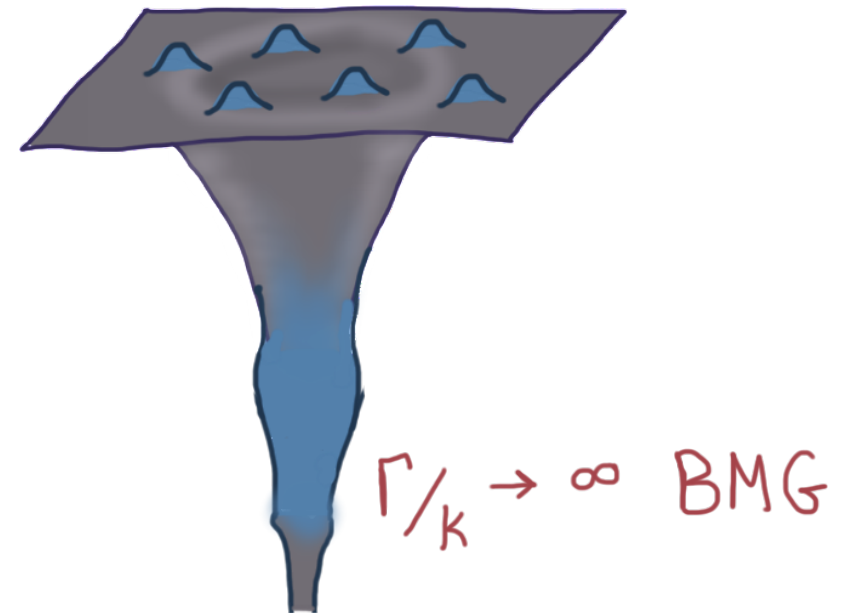
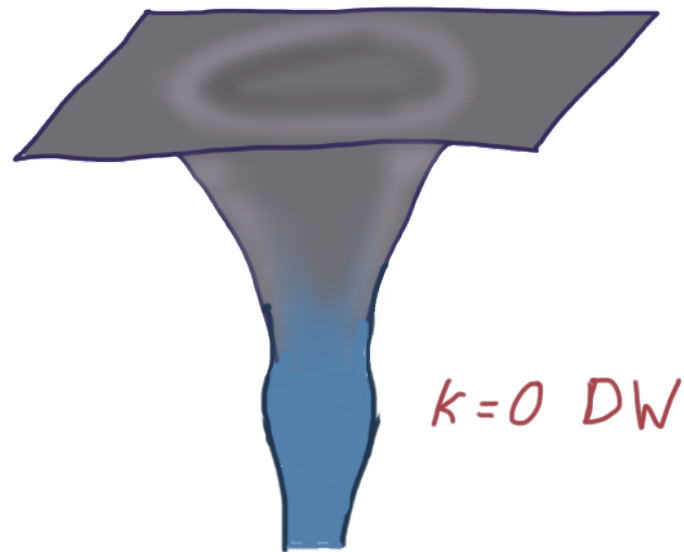
Scaling Signatures: Linear Response



Scaling Signatures: Linear Response



Scaling Signatures: Entropic c-Functions

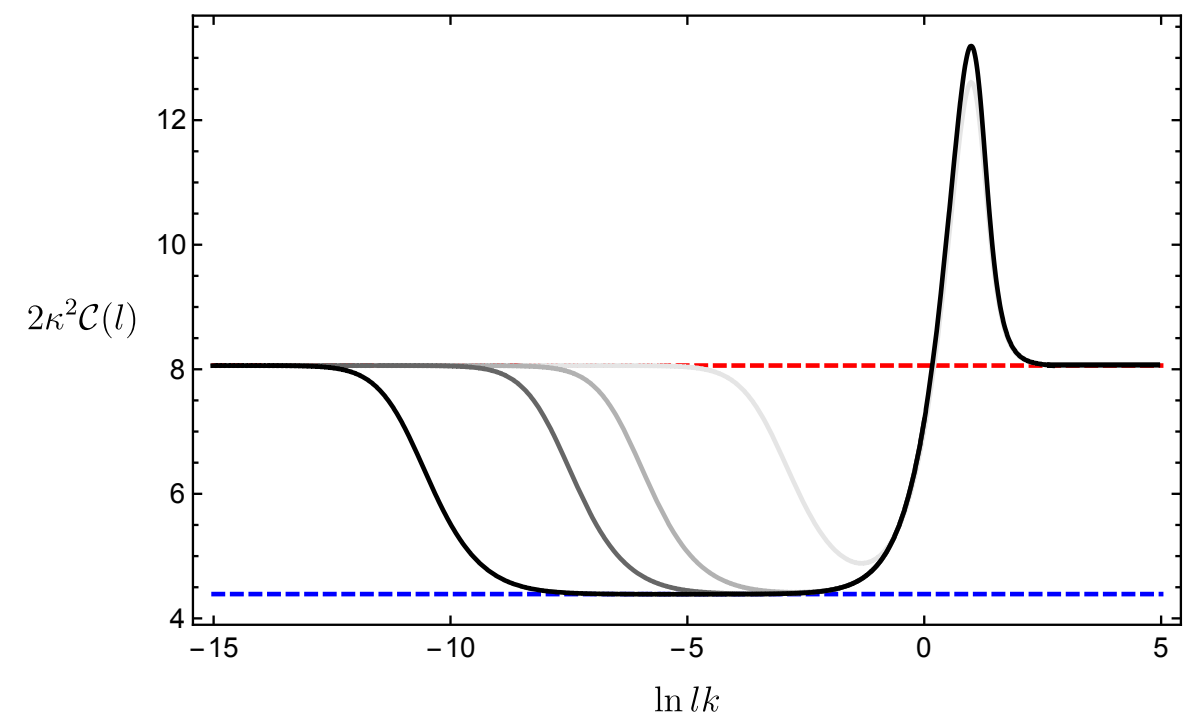
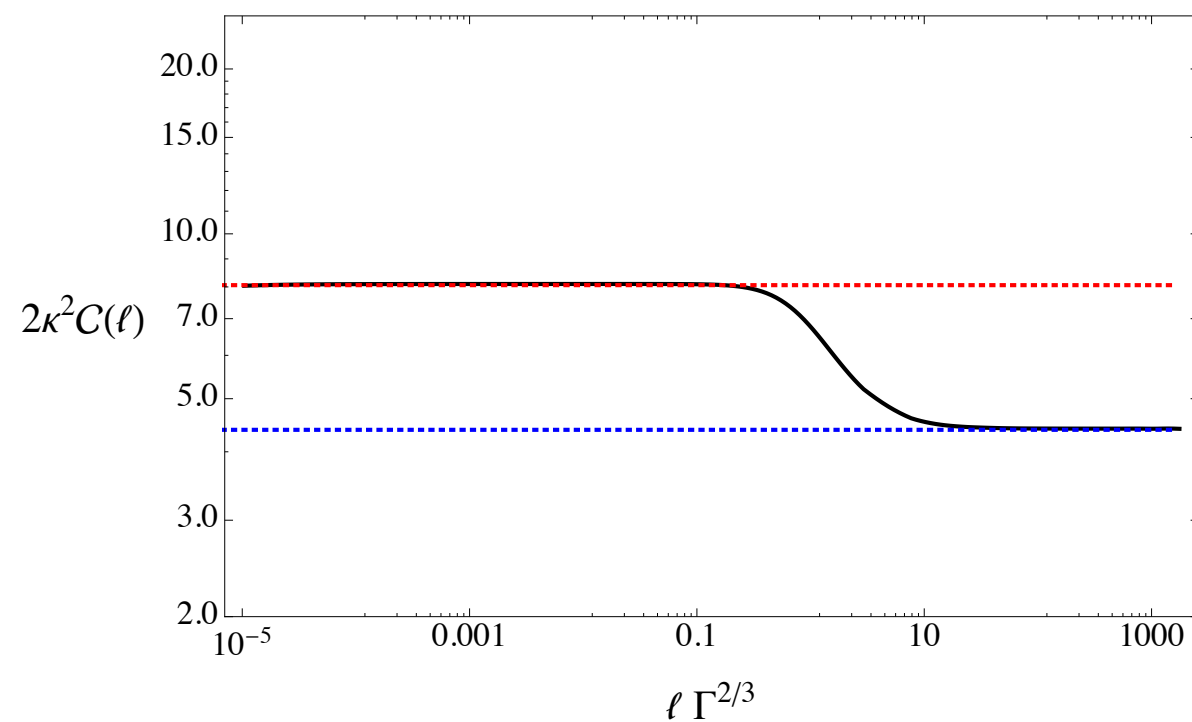
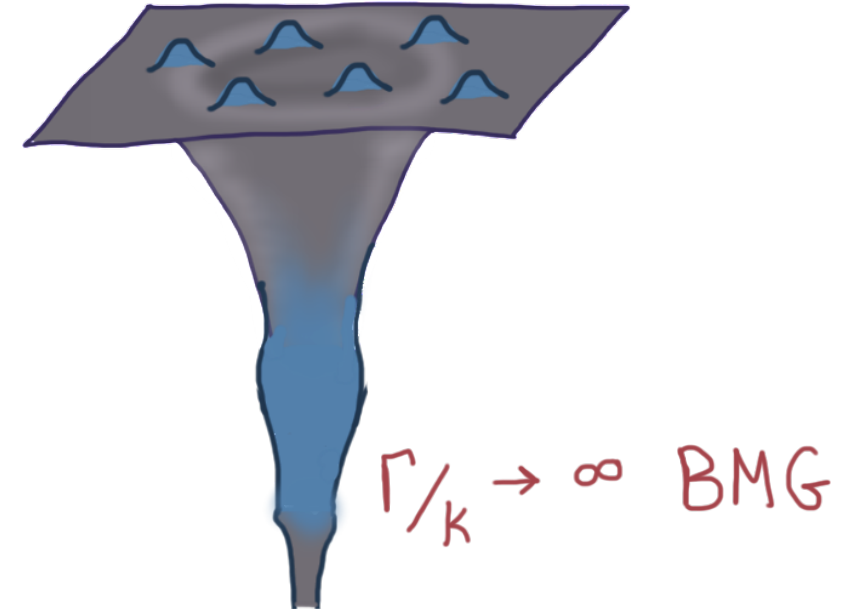
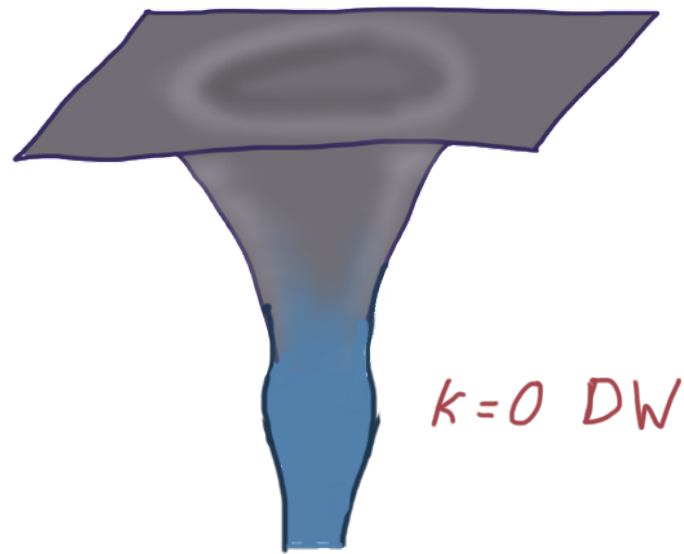


$$S_{\mathcal{A}} = 4\pi \mathcal{A}$$

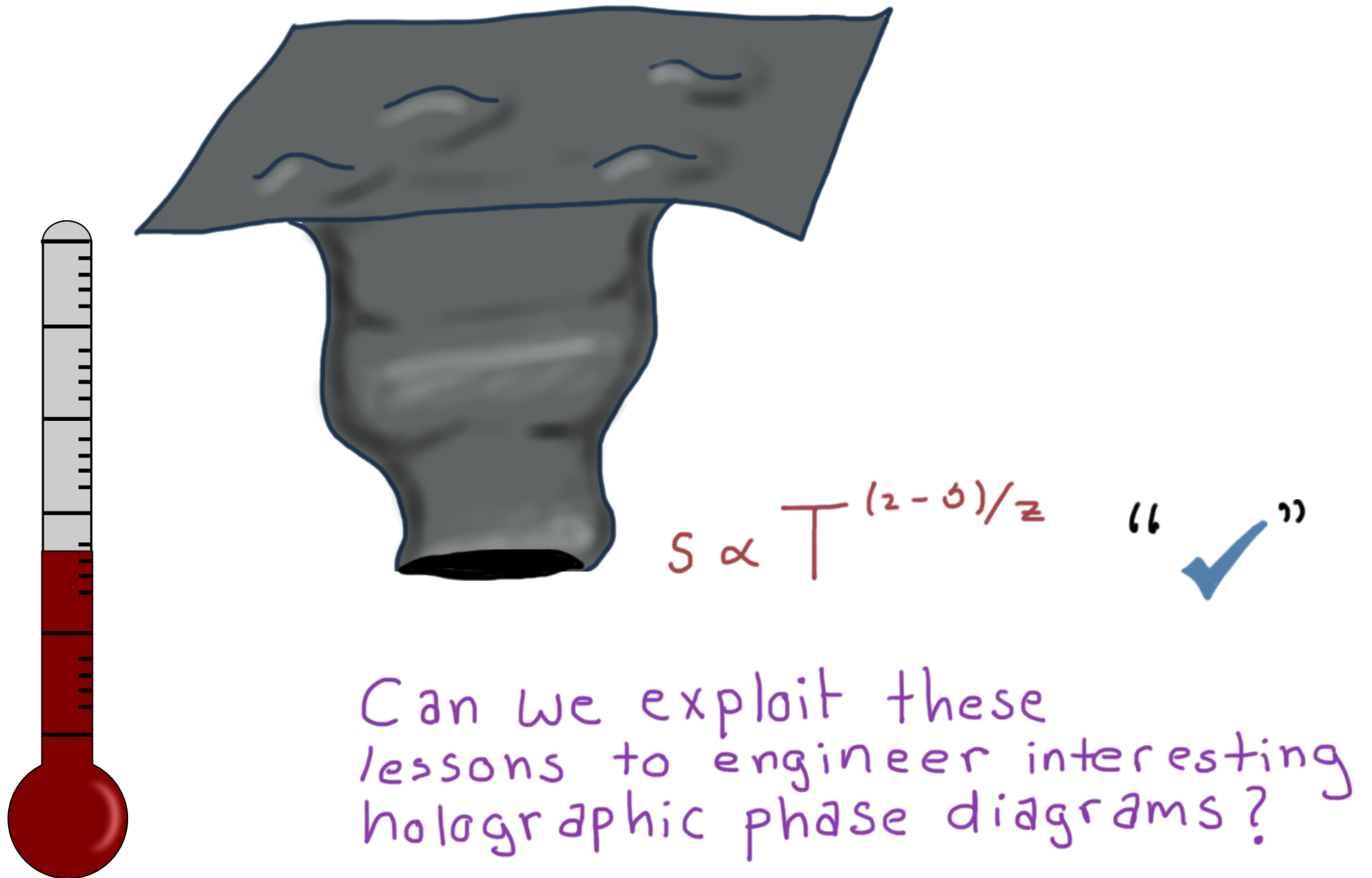
$$C(\ell) = \frac{\ell^3}{w^2} \frac{dS_{\mathcal{A}}}{d\ell}$$

RG
Monotone?

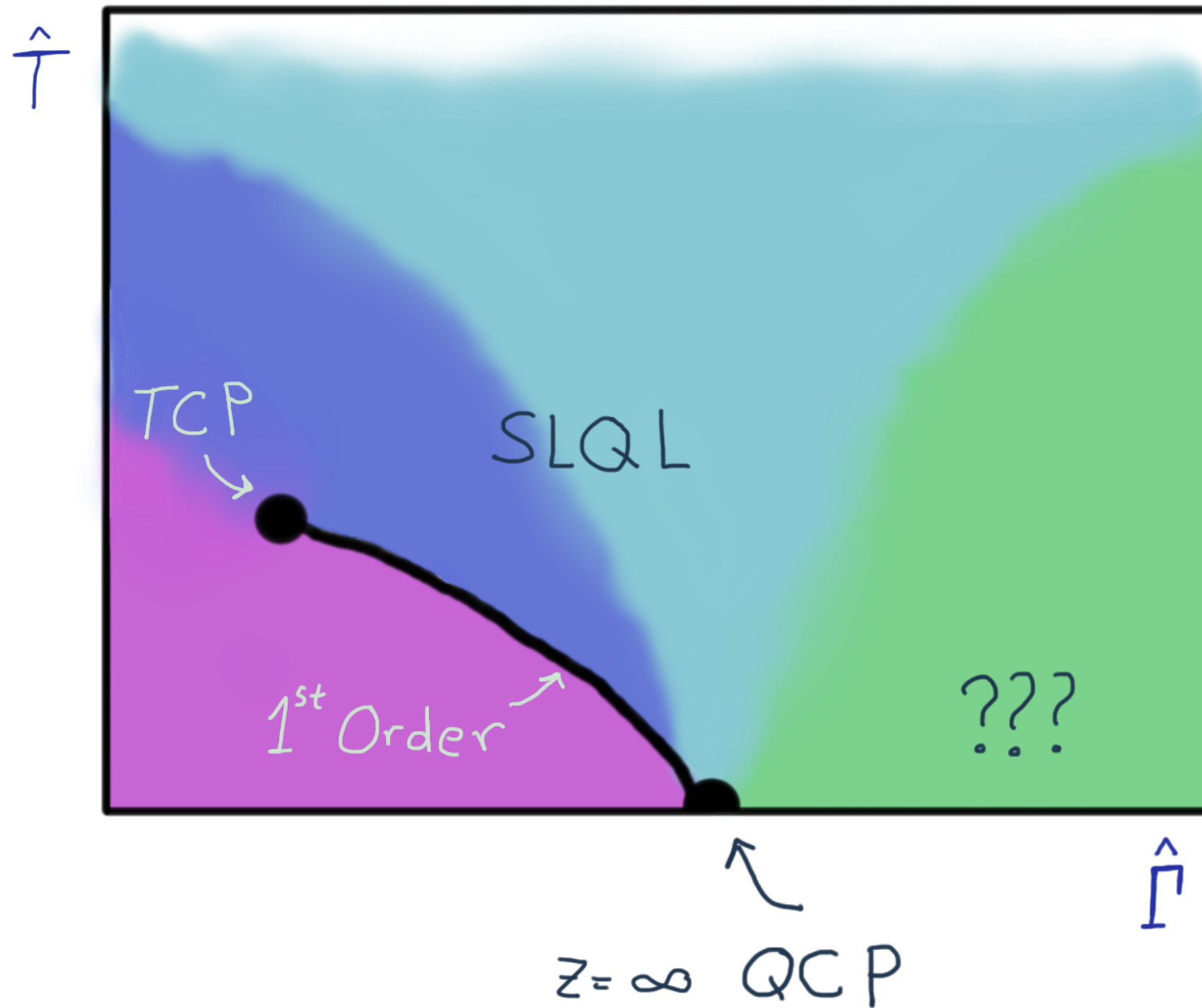
Scaling Signatures: Entropic c-Functions



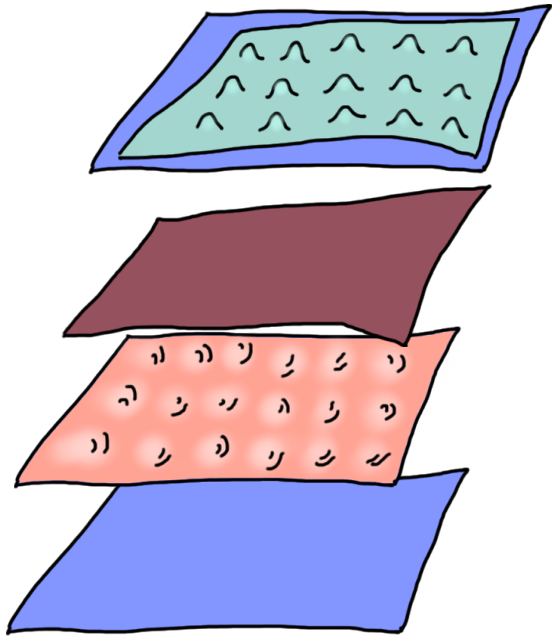
Scaling Signatures: Thermodynamics



Scaling Signatures: Quantum Critical Behavior



Some Open Questions and Outlook



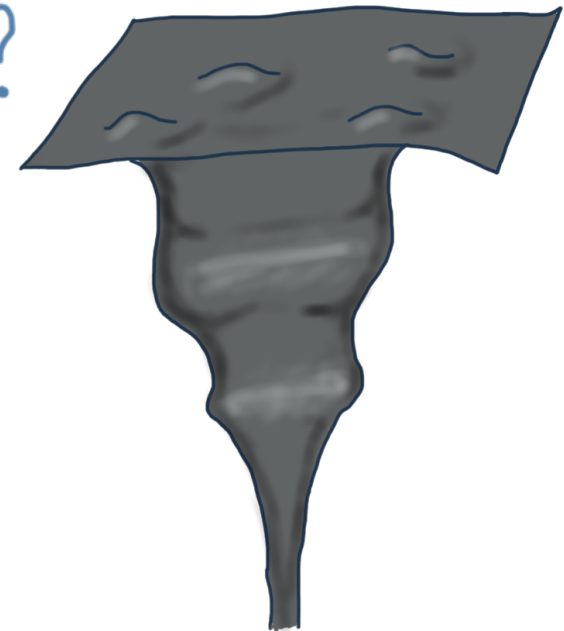
Related solutions with more/less symmetry
Options abound...

- Solⁿs in different dimensions/theories
- Solⁿs dual to systems @ finite- n
- Solⁿs away from the Q-lattice
- SUSY higher curvature corrections
- More, more, more

Flows are (perturbatively) stable in IR ? UV

The HSV solutions likely are not -- what about full flow?

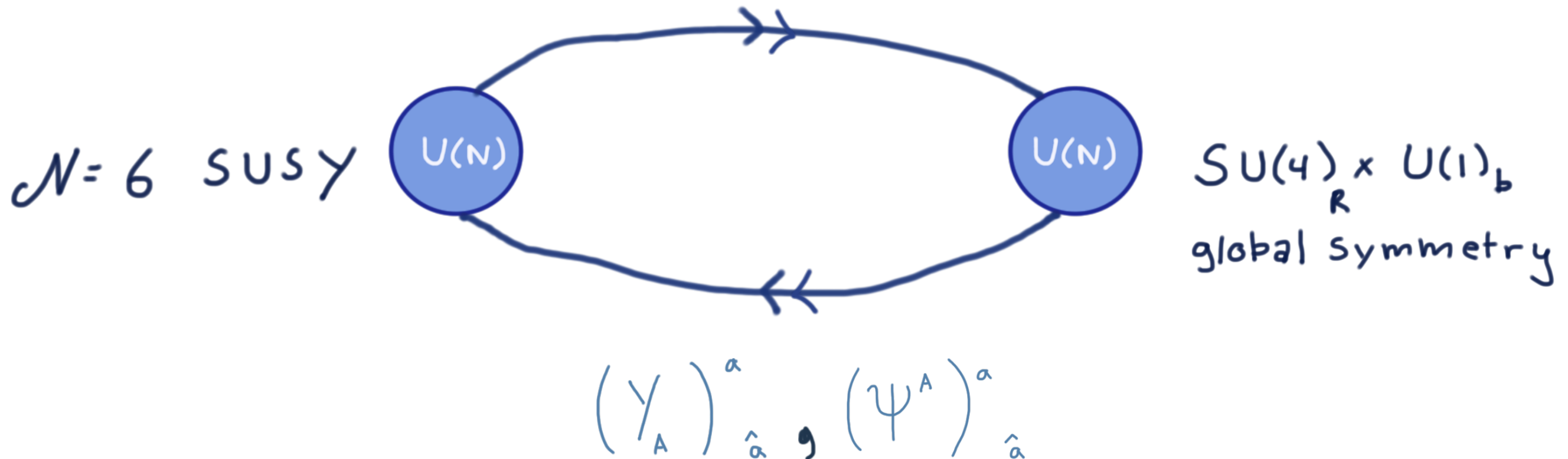
Other probes may yield further surprises
Top down fermionic spectral functions a natural candidate



Thank You

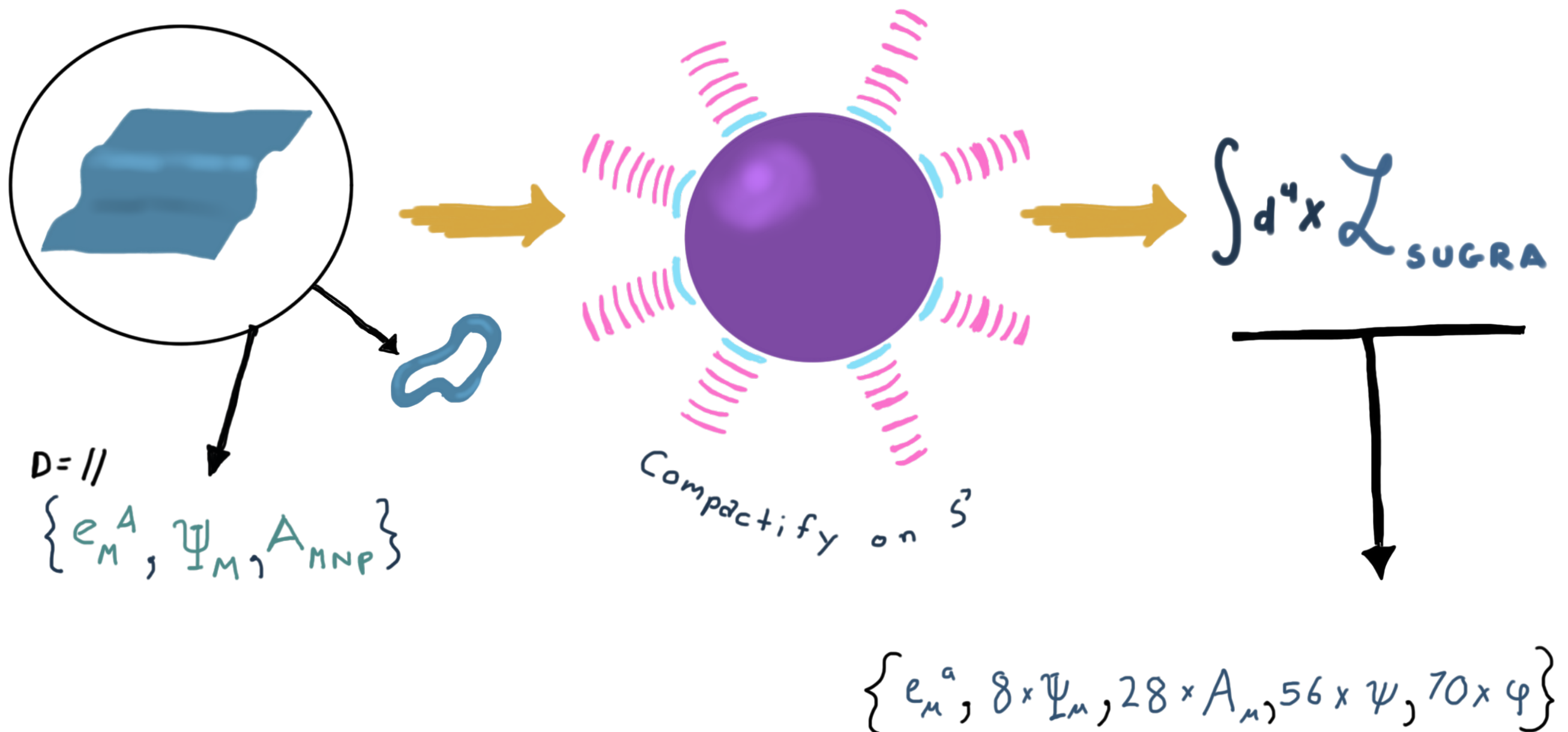
* DETAILS AND ASSORTED EXTRAS BEYOND THIS POINT

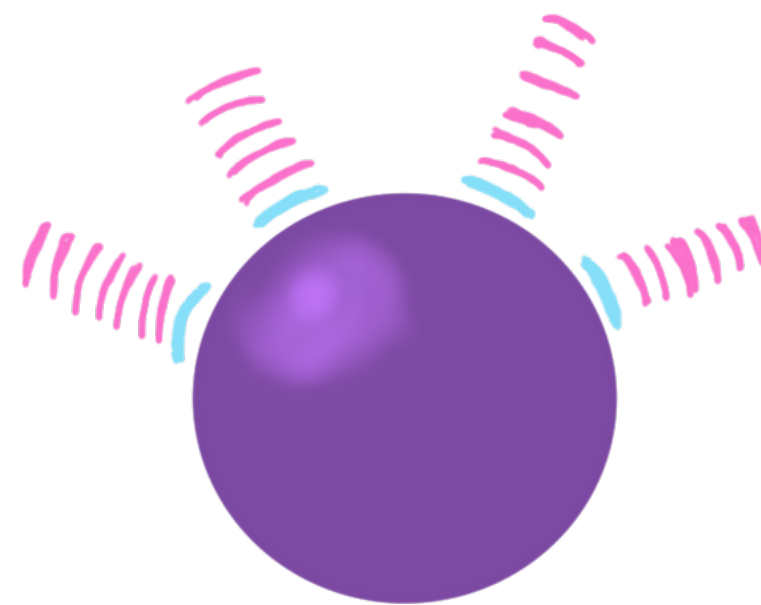
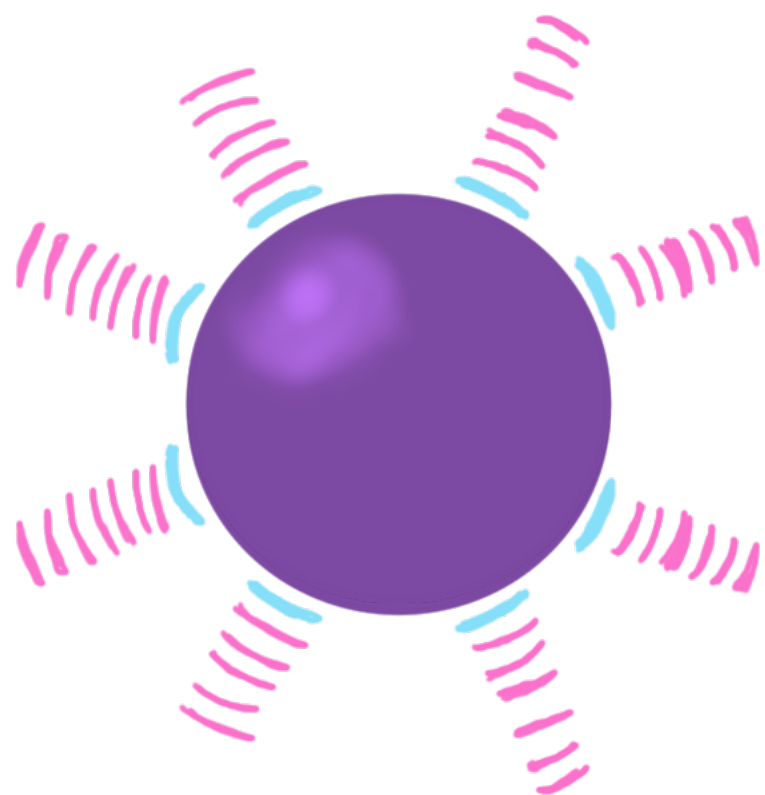
ABJM as a toy model of d=2+1 strongly interacting matter



$$S_{\text{ABJM}} \sim \int A \wedge dA + \frac{2}{3} A^3 - \hat{A} \wedge d\hat{A} - \frac{2}{3} \hat{A}^3 - (D\gamma)^2 + \bar{\psi} \not{D} \psi + \gamma^2 \psi^2 - V(\gamma^6)$$

Truncations we know and love





$$U(1)^4 \subset SO(8)$$

$$\Phi_i = \lambda_i e^{i\sigma_i}$$

$$e^{-1} \mathcal{L}_{\text{SUGRA}} = R - \frac{1}{2} \sum_i^3 \left((\partial \lambda_i)^2 + \sinh^2 \lambda_i (\partial \sigma_i)^2 \right) - \text{Re} \left(F^{+a} \mathcal{M}_{ab} F^{+b} \right) + 8g^2 \sum_i^3 \cosh \lambda_i$$

($\mathcal{N}=2$ w/ gravity + 3x Vector) ("STU" model)

A SUGRA Q-lattice Ansatz

$$e' \mathcal{L}_{\text{SUGRA}} = R - \frac{1}{2} \sum_i^3 \left((\partial \lambda_i)^2 + \sinh^2 \lambda_i (\partial \sigma_i)^2 \right) + g^2 \sum_i^3 \cosh \lambda_i$$

Periodic Spatial Modulation



$$\lambda_1 = \sigma_1 = 0$$

$$\lambda_2 = \lambda_3 = \gamma(r)$$

$$\sigma_2 = kx \quad \sigma_3 = ky$$

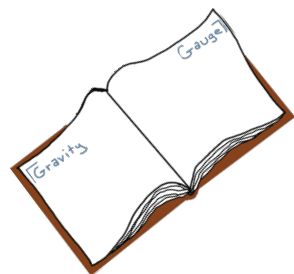
$$F_1 = F_2 = F_3 = F_4 = 0$$

$$ds^2 = \left[-U e^{-\gamma/2} dt^2 + r^2 d\vec{x}^2 \right] + \frac{dr^2}{U}$$

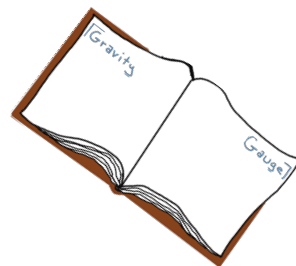


Broken Poincare Invariance

ODEs!!



Operator Matching



A top-down perk...

SUGRA

$$ds_{S^7}^2 = (d\psi + \omega)^2 + ds_{\mathbb{CP}^3}^2$$

$$\lambda e^{i\sigma} = \chi^+ \uparrow$$

$$i \gamma / 35_v \rightarrow 15_0 \oplus \dots$$

$$\uparrow$$

$$35_c \rightarrow 15_0 \oplus \dots$$

ABJM

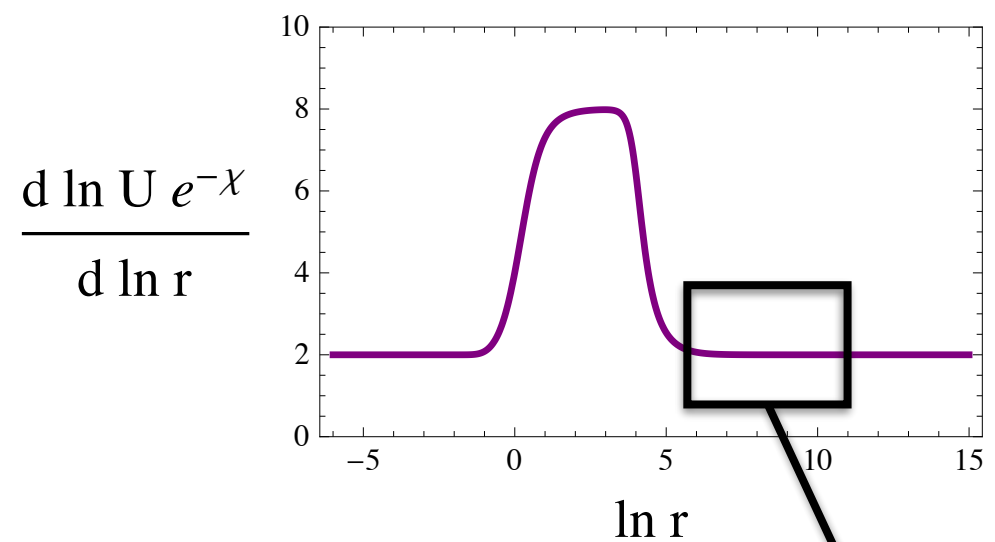
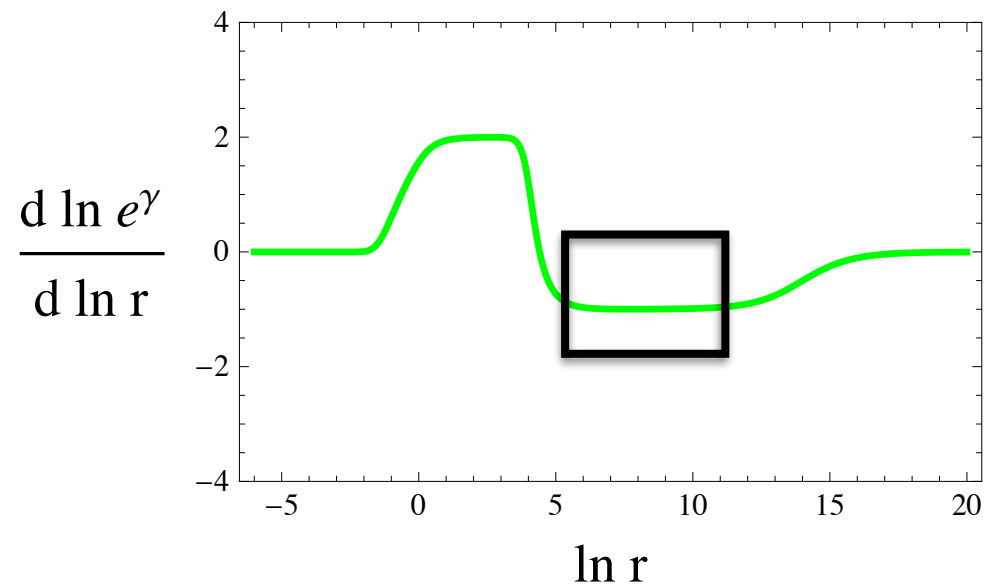
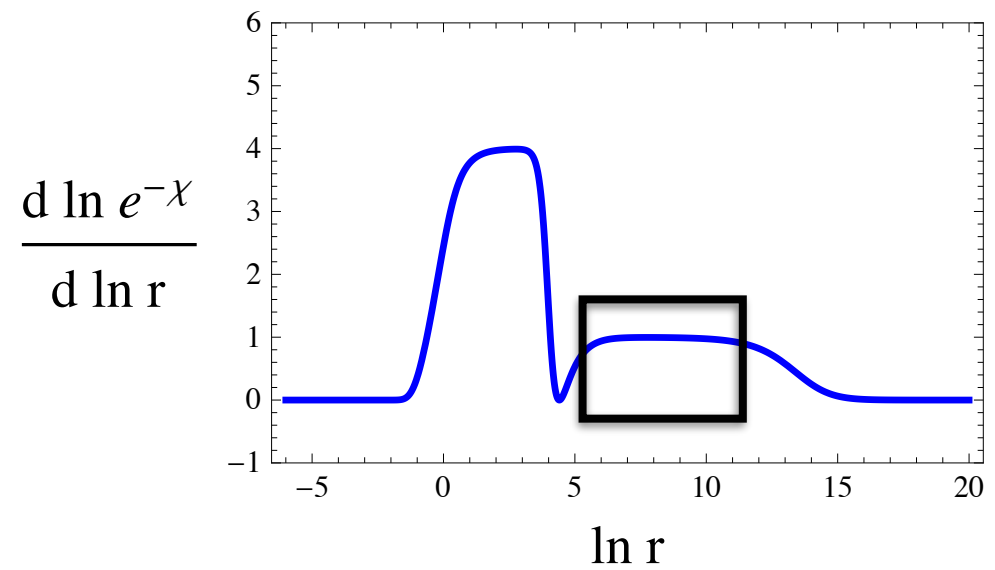
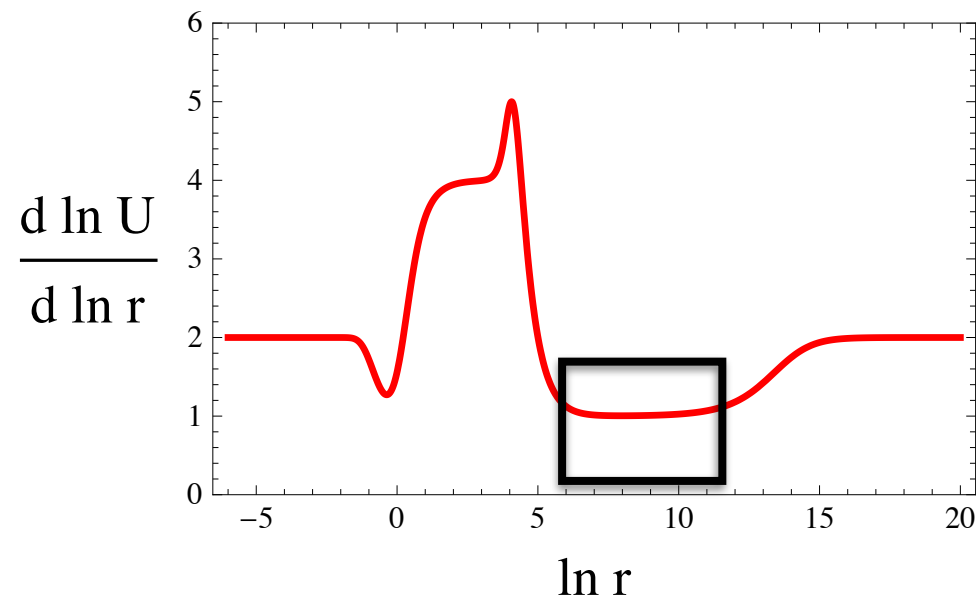
$$SU(4)_R \times U(1)_b$$

$$\mathcal{O}_{Y^2}^{\Delta=1} \sim \text{tr} \left(Y_A^\dagger Y^B - \frac{1}{4} \delta_A^B Y^2 \right)$$

$$\mathcal{O}_{\psi^2}^{\Delta=2} \sim \text{tr} \left(\psi_A^\dagger \psi^B - \frac{1}{4} \delta_A^B \psi^2 \right)$$

$$\Delta \mathcal{L} \sim \hat{\Gamma} \cos kx \mathcal{O}_{Y^2} + \hat{\Gamma} \cos ky \mathcal{O}'_{Y^2} + \Gamma \sin kx \mathcal{O}_{\psi^2} + \Gamma \sin ky \mathcal{O}'_{\psi^2}$$

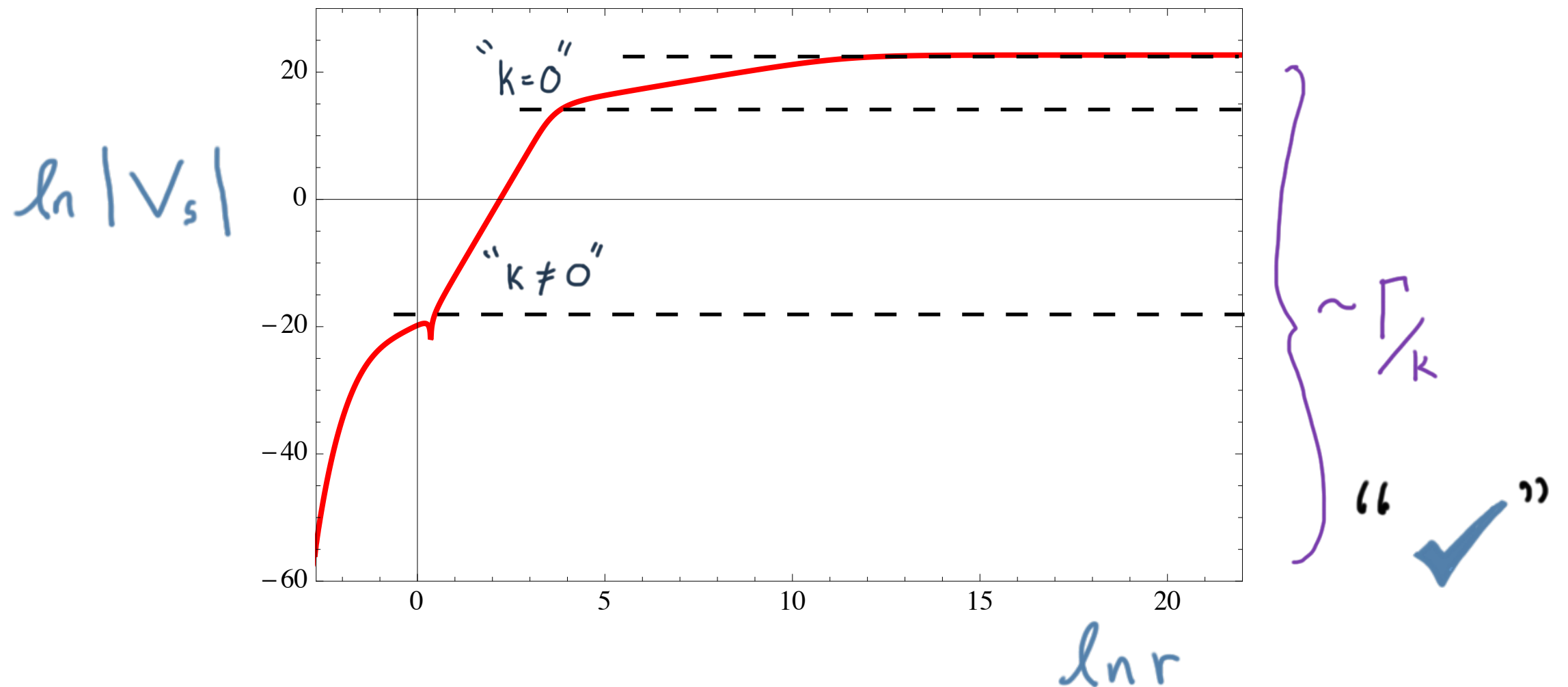
Nonlinear Solutions



$$ds^2 = \frac{1}{p^4} (-dt^2 + d\vec{x}^2 + d\rho^2)$$

HSV: $b = -2$, $z = 1$

Scaling Signatures: Linear Response



$$\left(\square + \dots \right)_{ij} \delta \Phi^j = 0 \quad \Rightarrow \quad -\psi'' + V_s \psi = \omega^2 \psi$$

but details matter...