New results for 5 loop massless propagators integrals



A. Georgoudis

Department of Physics and Astronomy Uppsala University

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A. Georgoudis



p-Integrals I

We want to compute massless 2-point integrals at 5-loops. We are interested in computing this two families of integrals:

- Momentum-space integrals
- Position-space integrals (i.e propagators of type $x_{ij} = x_i x_j$).



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We can map the planar sector of this two types of integrals but we have to compute the non planar part separately.





Integral-by-parts Identities

■ IBPs are integration of a total derivative:

$$0 = \int \frac{\mathrm{d}^D \ell_1}{\mathrm{i} \pi^{D/2}} \dots \frac{\mathrm{d}^D \ell_L}{\mathrm{i} \pi^{D/2}} \sum_{j=1}^L \frac{\partial}{\partial \ell_j^{\mu}} \frac{v_j^{\mu}}{D_1^{\nu_1} \dots D_m^{\nu_m}},$$



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They can be used to express a general integral as a combination of Master Integrals:

$$\mathsf{I} = \sum_{i=1}^N c_i \mathsf{I}_i,$$

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Integral-by-parts Identities: constraints

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Reducing convergent integrals!



Integral-by-parts Identities: Example

Reduction of two-loop propagator integral





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The spurious pole give a constrain on two epsilon orders. Convergence of the starting integral and the insertion of the value of the trivial integral $M_{\text{Dbubble}} = \frac{1}{\epsilon^2}$ gives:

$$M_{
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Simple algebra gave us a integral result!

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- Start from a 4-point finite conformal integral.
- We can now take the limit x₂ → x₁ = 0 which introduces two scales x₂² ≪ x₃².
- x_i is of the order of x_2 or order of x_3 . For $x_2^2 < x_i^2$ we can write:

$$\frac{1}{(x_2 - x_i)^2} = \sum_{n=0}^{\infty} \frac{(2x_2 \cdot x_i - x_2^2)^n}{(x_i^2)^{n+1}}$$

 \blacksquare 2^{*I*} possible such regions.

For finite integrals the expansion is still finite.



 $C_{1} = \int \frac{\mathrm{d}^{4} x_{5} \dots \mathrm{d}^{4} x_{9}}{x_{15}^{2} x_{16}^{2} x_{26}^{2} x_{27}^{2} x_{29}^{2} x_{34}^{2} x_{36}^{2} x_{37}^{2} x_{45}^{2} x_{48}^{2} x_{56}^{2} x_{58}^{2} x_{59}^{2} x_{78}^{2} x_{79}^{2} x_{89}^{2}}.$

Example



UPPSALA

$$C_1 = \int \frac{\mathrm{d}^4 x_5 \dots \mathrm{d}^4 x_9}{x_{15}^2 x_{16}^2 x_{26}^2 x_{27}^2 x_{29}^2 x_{34}^2 x_{36}^2 x_{37}^2 x_{45}^2 x_{48}^2 x_{56}^2 x_{58}^2 x_{59}^2 x_{78}^2 x_{79}^2 x_{89}^2}$$

$$\begin{split} C_1|_{|x_{12}|\ll|x_{13}|} &\sim x_{13}^2 \int \frac{\mathrm{d}^4 x_5 \dots \mathrm{d}^4 x_9}{x_{15}^2 x_{16}^2 x_{26}^2 x_{27}^2 x_{29}^2 x_{56}^2 x_{58}^2 x_{59}^2 x_{78}^2 x_{79}^2 x_{89}^2} \\ &\sim u^{-11+\frac{5d}{2}} (x_{13}^2)^{-10+\frac{5d}{2}} \mathcal{P}^{(5)}(d) \,. \end{split}$$

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UPPSALA

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$$P^{(5)}(d=4-2\varepsilon)=\frac{c_1}{\varepsilon^5}+\frac{c_2}{\varepsilon^4}+\frac{c_3}{\varepsilon^3}+\frac{c_4}{\varepsilon^2}+\frac{c_5}{\varepsilon}+c_6+\mathscr{O}(\varepsilon).$$

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$$P^{(5)}(d=4-2\varepsilon) = \frac{c_1}{\varepsilon^5} + \frac{c_2}{\varepsilon^4} + \frac{c_3}{\varepsilon^3} + \frac{c_4}{\varepsilon^2} + \frac{c_5}{\varepsilon} + c_6 + \mathscr{O}(\varepsilon).$$
$$c_1 = c_2 = c_3 = c_4 = c_5 = 0.$$

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We can fix the c_6 by using magic identities! Take the 4-point conformal integral with points $\{x_7, x_8, x_9\}$. Exchange $x_2 \leftrightarrow x_3, x_4 \leftrightarrow x_5$ and expand:

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$$C_2|_{|x_{12}|\ll|x_{13}|}\sim \frac{120\,\zeta(3)\zeta(5)}{u}.$$

This also fixes the finite term of the five-loop p-integral to be

$$c_6 = 120 \zeta(3)\zeta(5)$$
.

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How?





Results

With IBPs and conformal symmetry constraints we were able to fix:

- Expansions of 169 out of the 187 genuine five-loop masters.
- All but two of the 95 planar masters were fixed to transcendental weight 9 in this way.
- In order to get the 169 integrals up to trascendentality 9 we used parametric integration (8 single orders of 7 integrals needed).



Glue and Cut symmetry

If we want to compute non planar momentum space integrals we can use a different method.



Glue and Cut symmetry





In the *G*-scheme normalization:

- We find that ζ_2 is absent!
- The coefficients of ε^{-n} never have transcendental contributions of weight larger than 9-2n (where n > 0).
- We can redefine odd zeta values to eliminate even zetas in the expansion of all master integrals.
- Any dependence on multiple zeta values comes in the combination $\varpi = \zeta_{2,6} \zeta_{5,3}$.



- Presented recent result for 5-loop master integral.
- All the complication is encoded in the IBP reductions.
- Applying CaG symmetry to higher loops to check and obtain new results.
- Derive similar results for different dimensions (e.g d = 3).



Thank You