

# Neutron-Mirror Stars

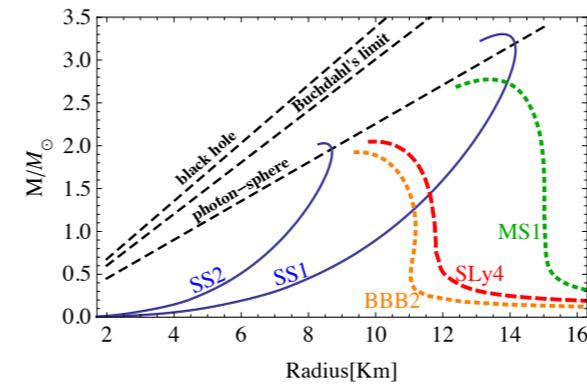
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INFN-LNGS  
[massimo@lngs.infn.it](mailto:massimo@lngs.infn.it)

work in progress with Z. Bherezhiani, R. Biondi and F. Tonelli

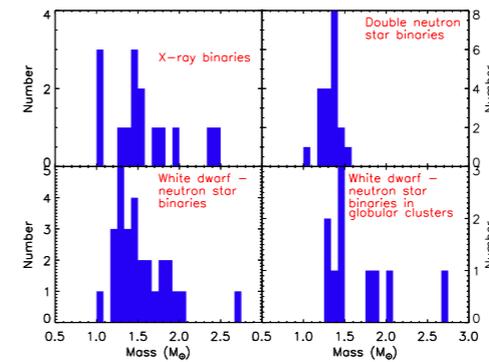
*Particle Physics with  
Neutrons at the ESS  
Stockholm, Dec 14, 2018*

# Outline

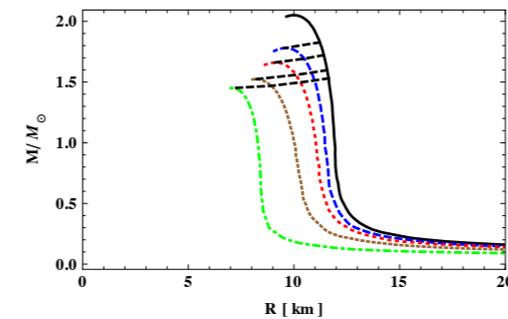
- Introduction



- Mass distribution of NSs



- Neutron-mirror stars



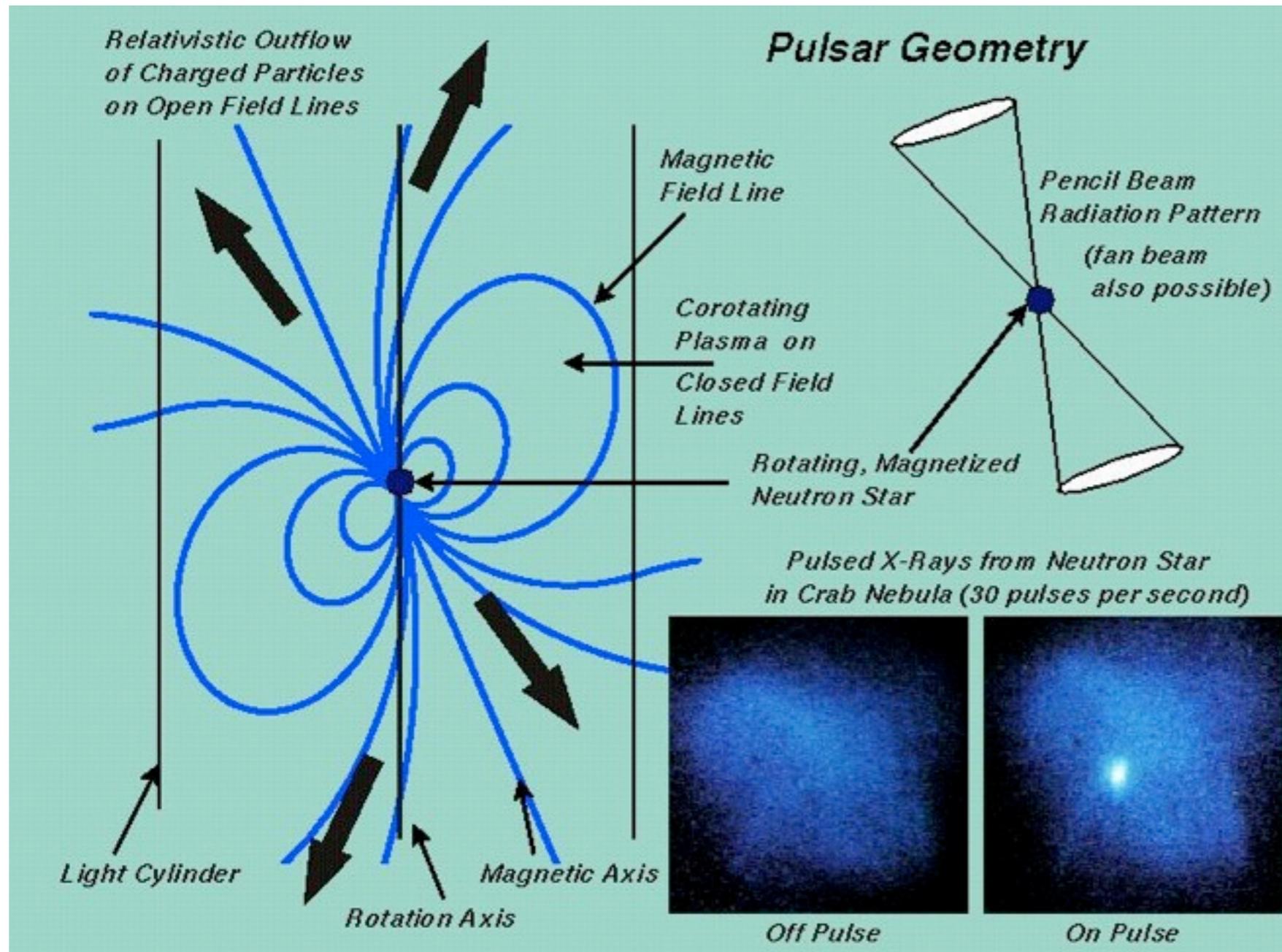
- Conclusions

# Introduction



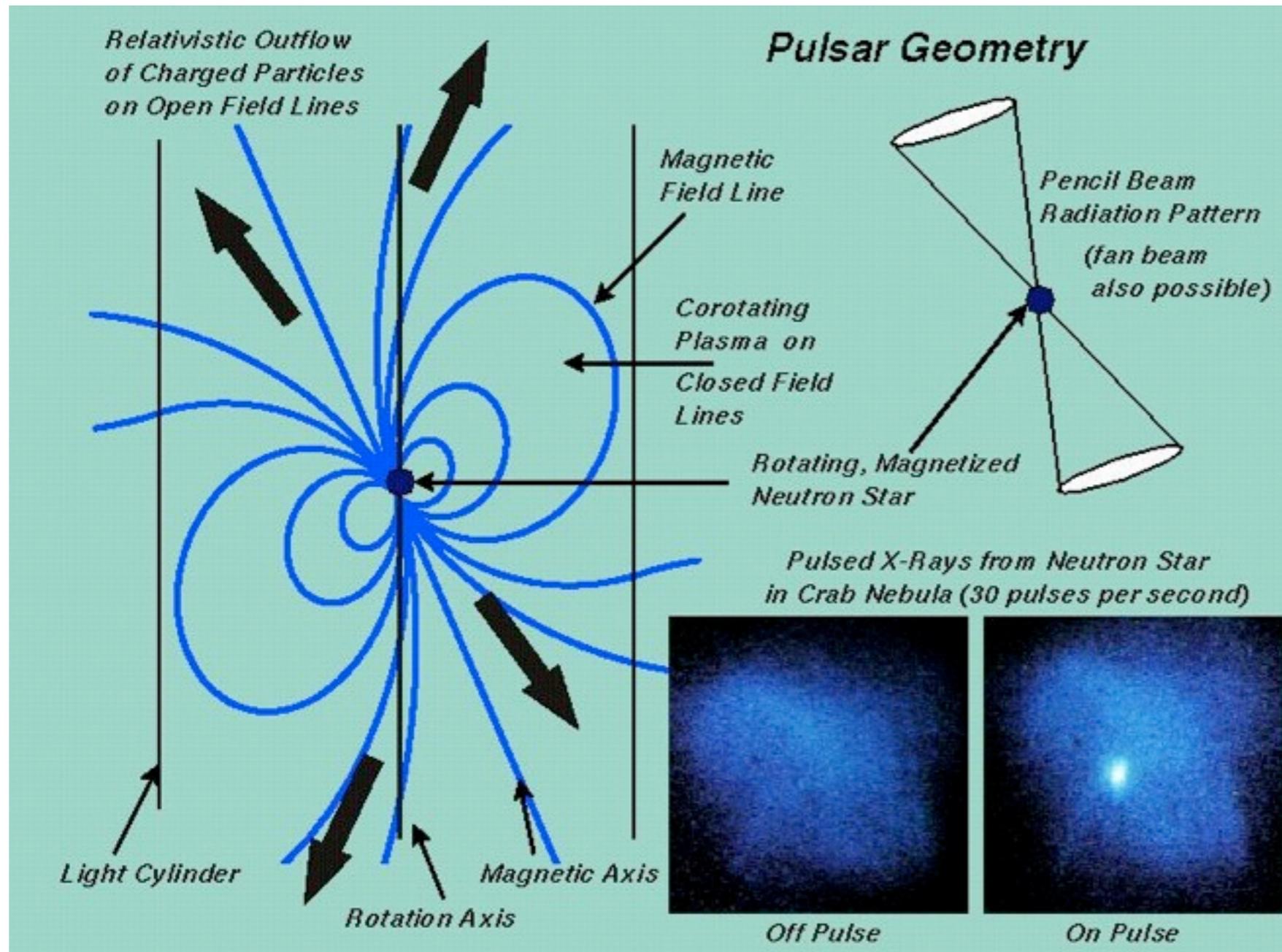
# Pulsars

Compact stars are typically observed as pulsars



# Pulsars

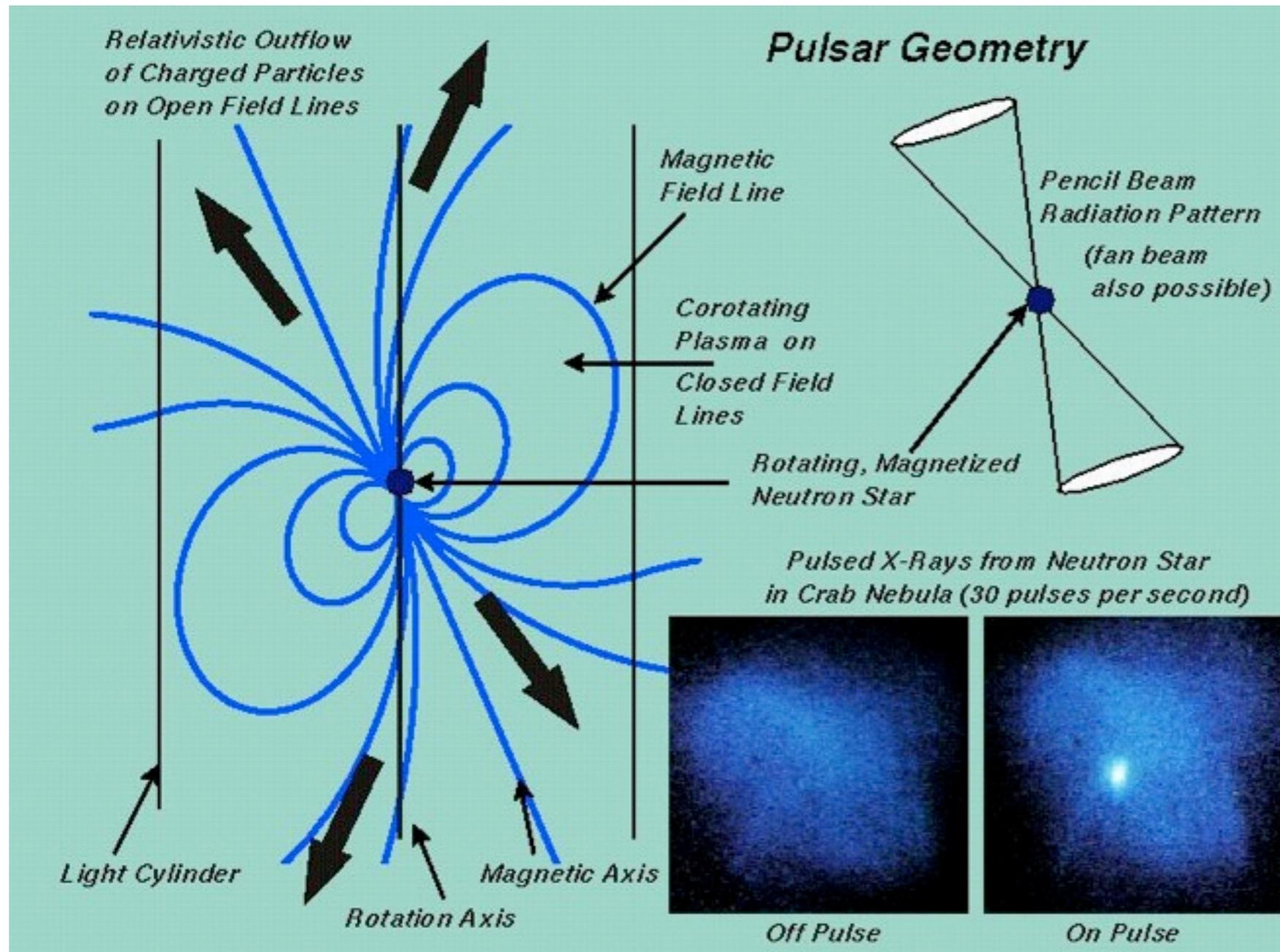
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Vela  
P~89ms

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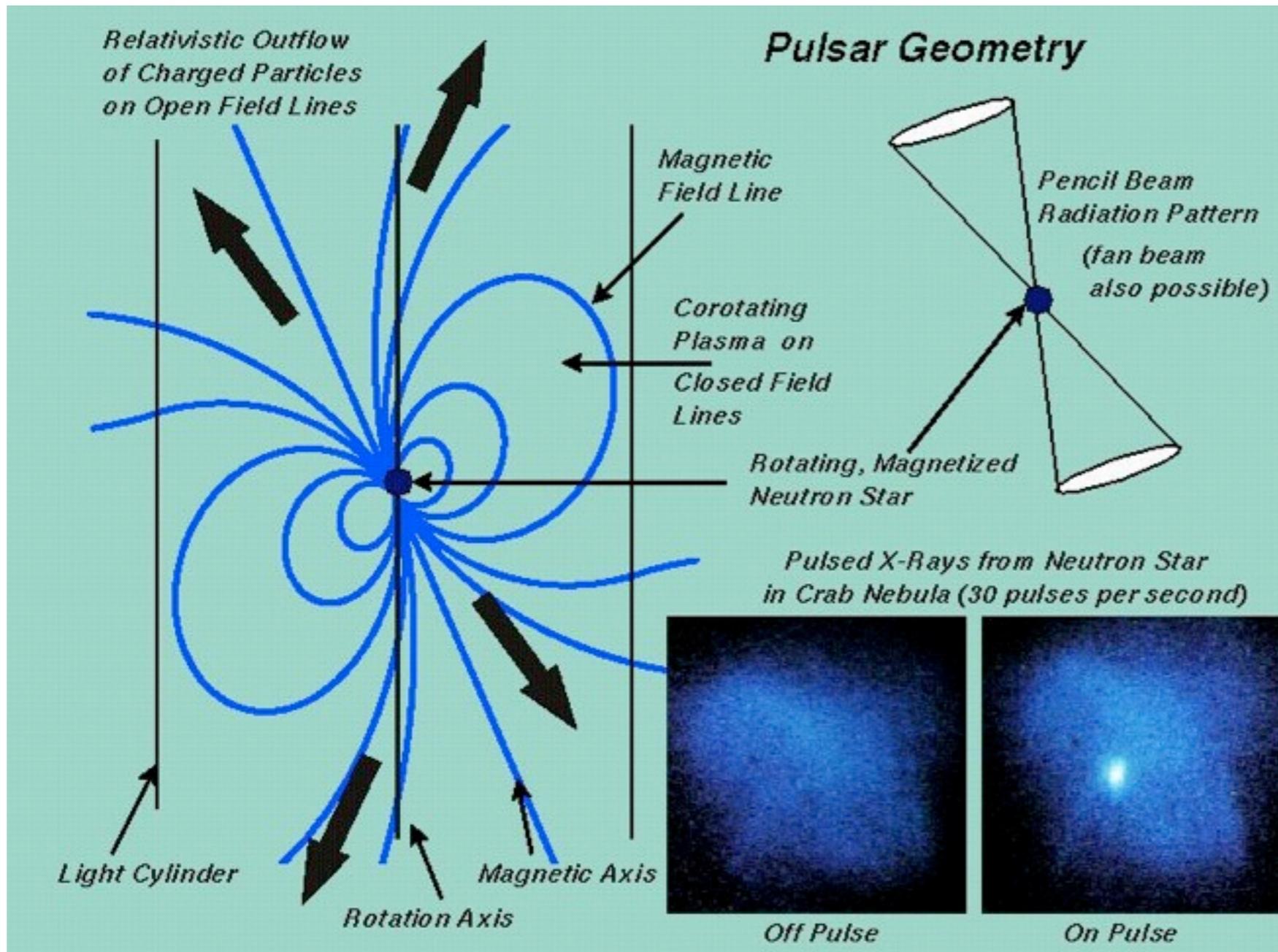


Vela  
P~89ms

Crab  
P~33ms

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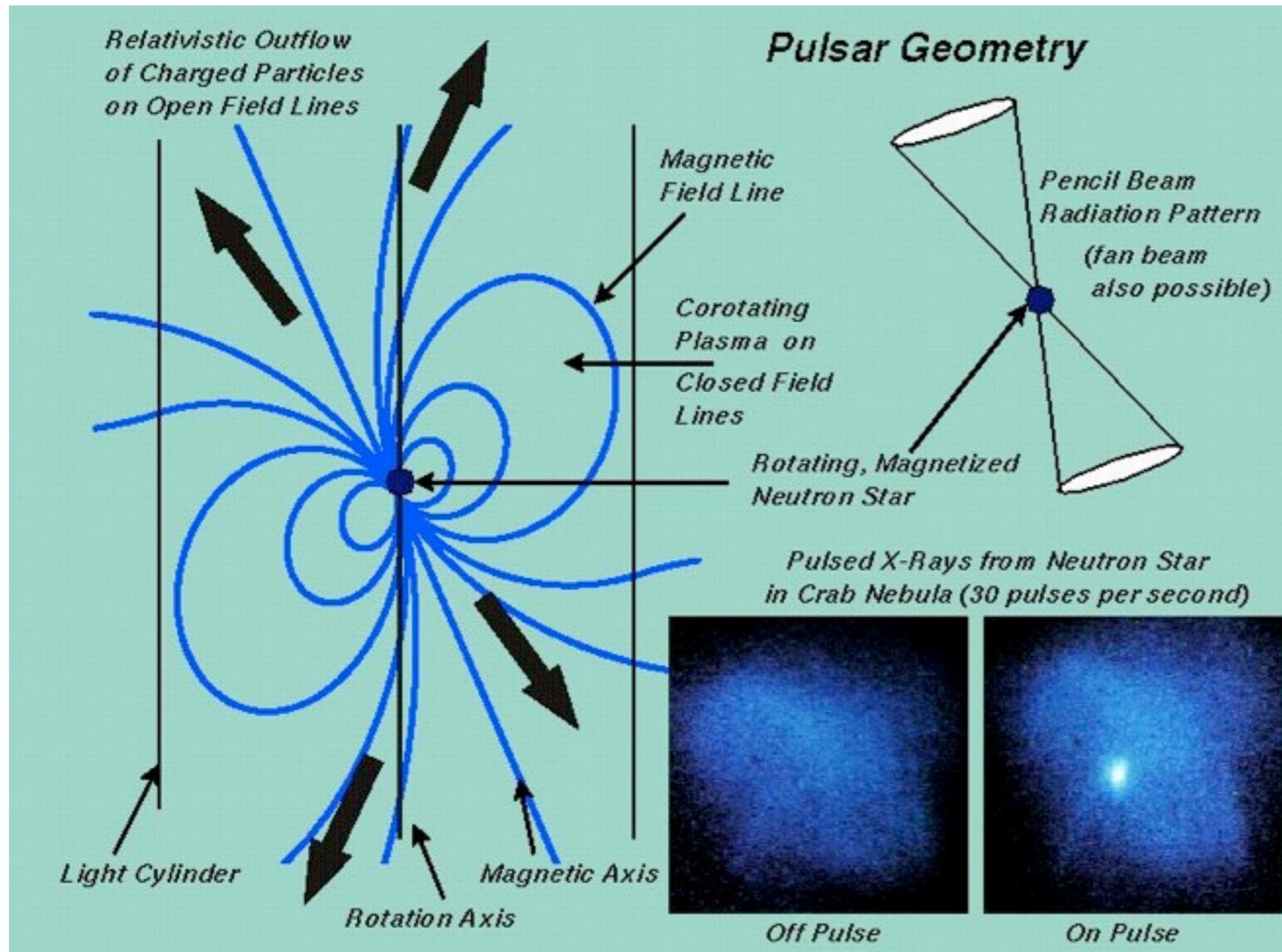
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**PSR J1748 -2446ad**  
**P~1.4ms**

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**Vela**  
**P~89ms**

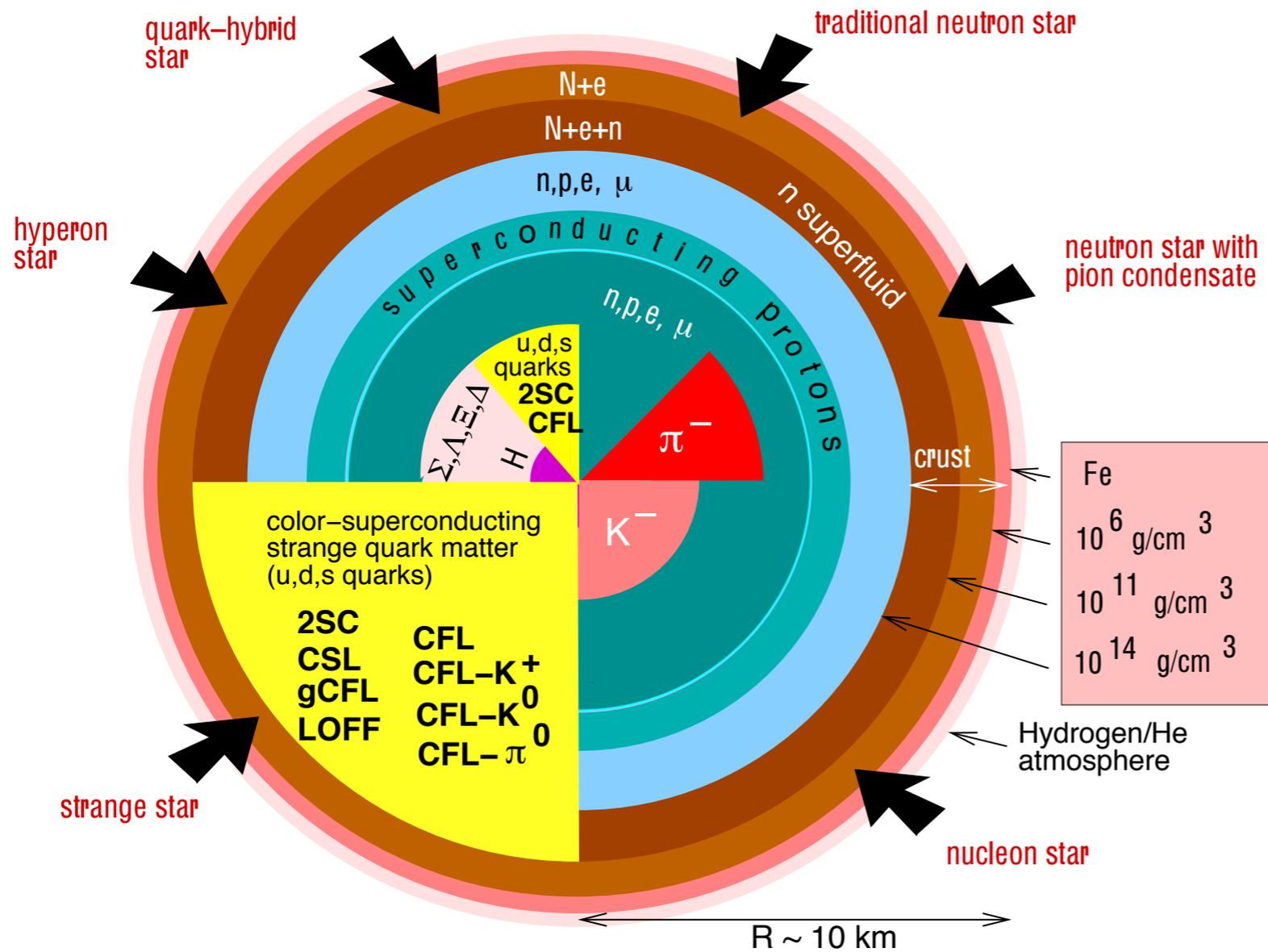
**Crab**  
**P~33ms**

**PSR J1748 -2446ad**  
**P~1.4ms**

**Modern washing machines** about 30 rounds/s thus P~33 ms as the Crab pulsar

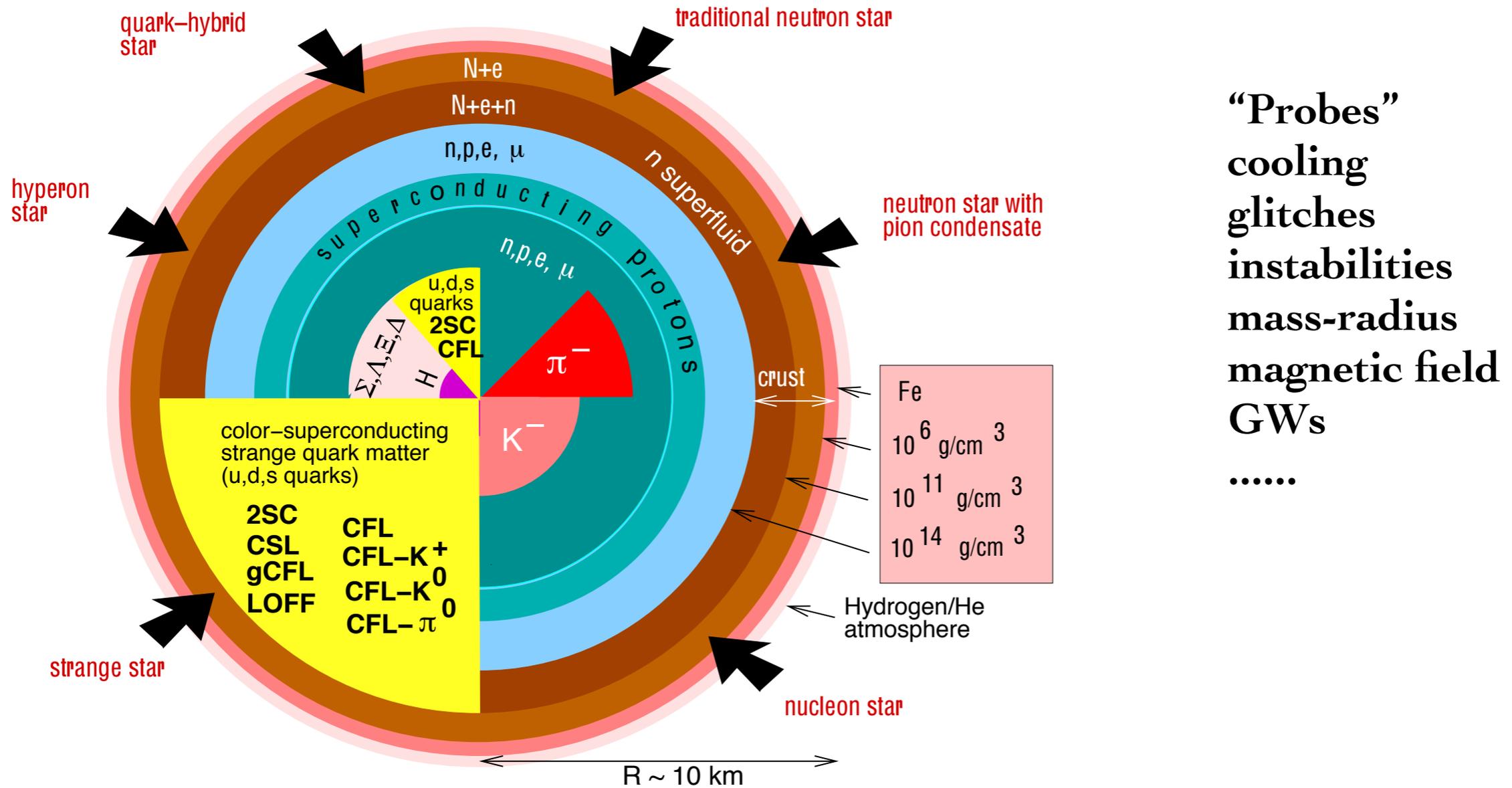
**Ferrari Engine F2004 (F1 world champion 2004)** P~3.16 ms

# Compact stellar objects



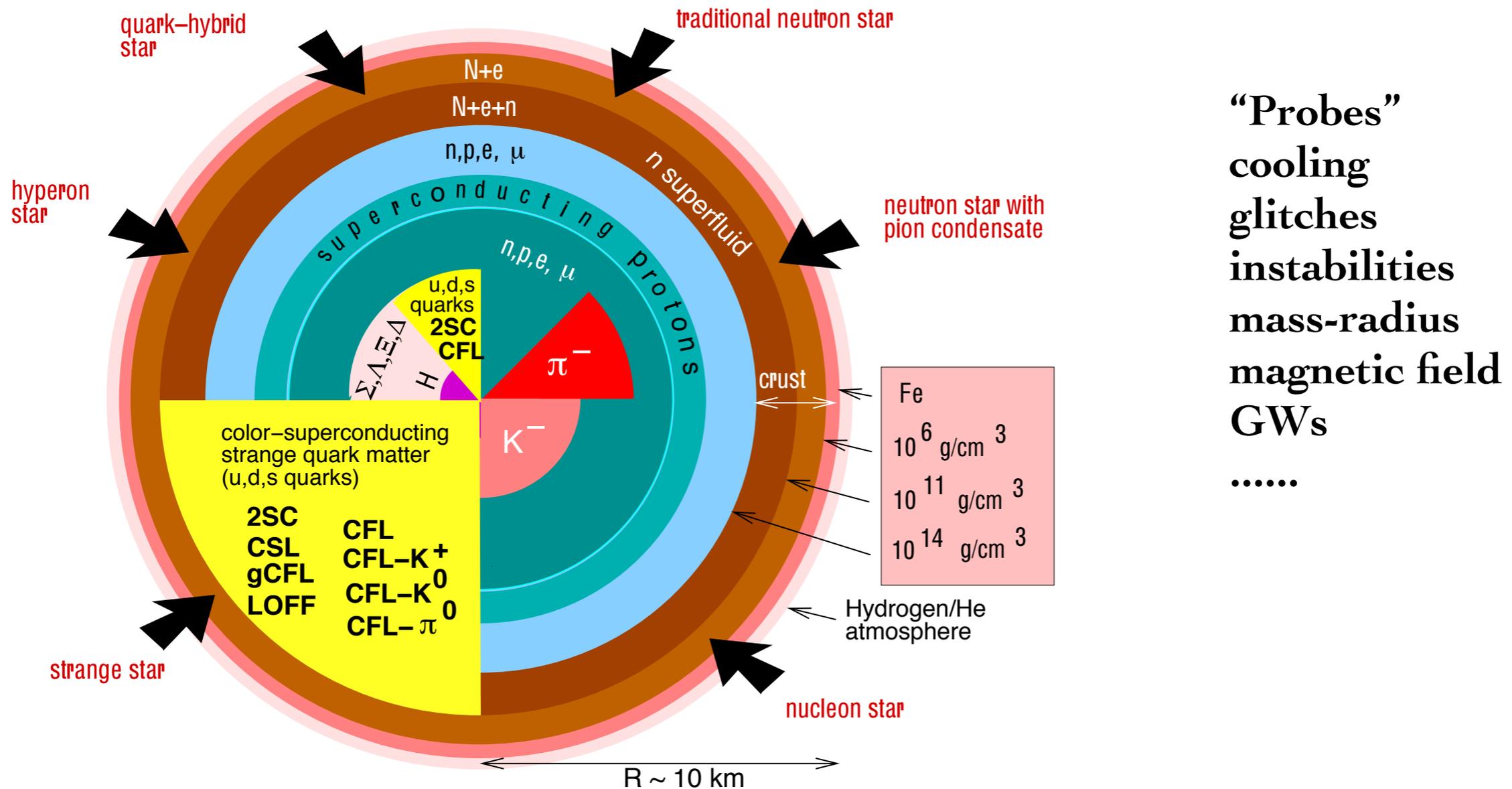
F. Weber, Prog.Part.Nucl.Phys. 54 (2005) 193

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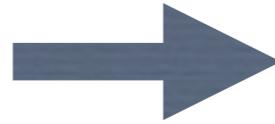
Compact stars can have an age comparable to the age of the Universe.  
 However, pulsars can be observed for a limited time... of tens of millions of years (as far as they are active)

# *Simplified Neutron Star model*

## Catalysed nuclear matter

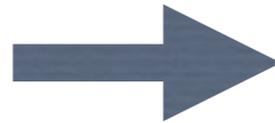
Neutron decay  $n \rightarrow p + e + \bar{\nu}_e$

Electron capture  $p + e \rightarrow n + \nu_e$



$$\mu_n = \mu_p + \mu_e$$

Electric charge neutrality



$$N_p = N_e$$

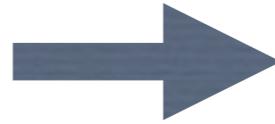
**Dynamical equilibrium: nuclear matter cools down by neutrino emission**

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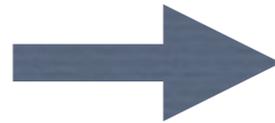
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$$N_p = N_e$$

**Dynamical equilibrium: nuclear matter cools down by neutrino emission**

**In Neutron stars, nuclear matter is mainly made of neutrons:**

$$N_p/N_n \sim 1/10$$

**For our purposes we can roughly assume that a NS is a Fermi liquid of neutrons (neglect neutron superfluidity etc.) filling a Fermi sphere**

# *Equilibrium configurations*

In general one solves the TOV's equations of hydrostatic equilibrium

$$\frac{dm}{dr} = 4\pi\rho r^2$$
$$\frac{dp}{dr} = (\rho + p) \frac{m + 4\pi p r^3}{2mr - r^2}$$

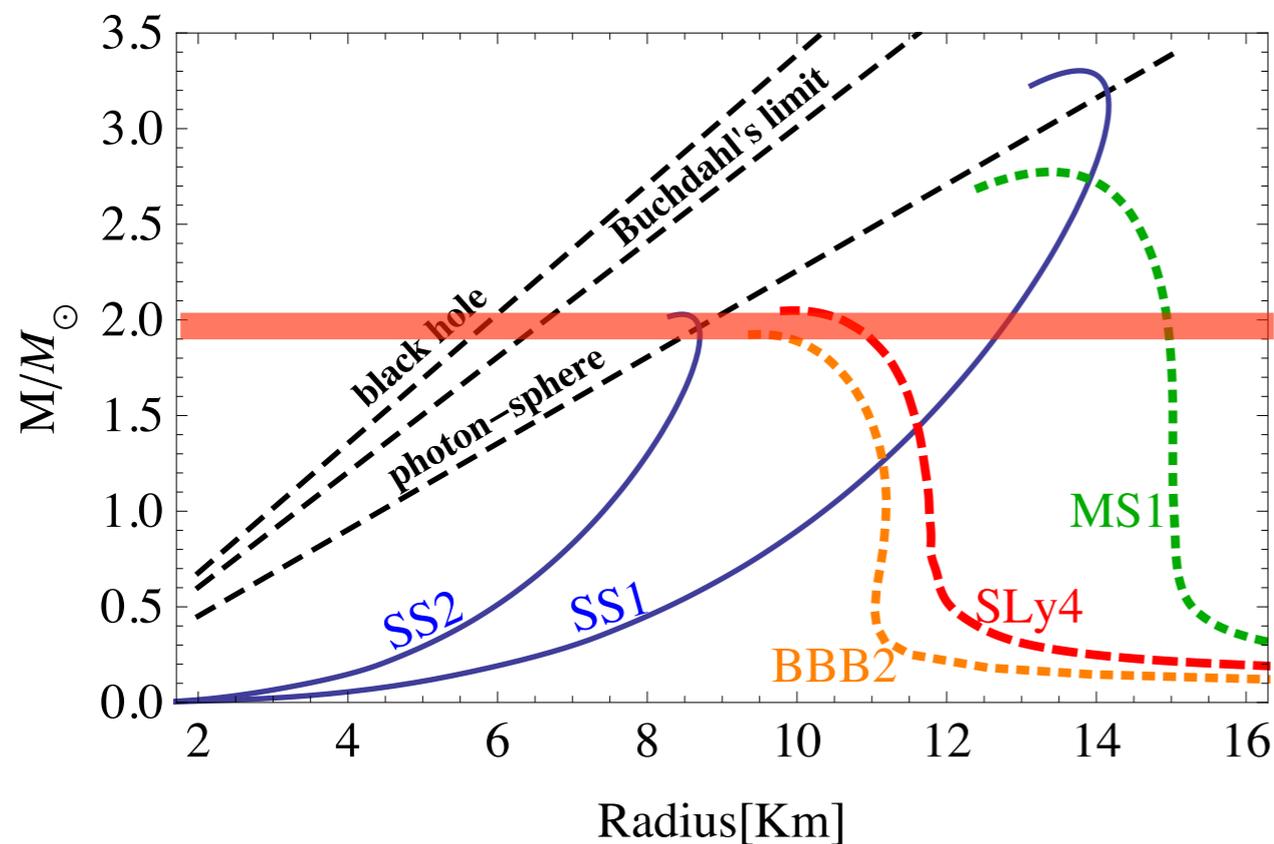
to solve these equations one needs an EoS, i.e. a relation between the pressure and the density.

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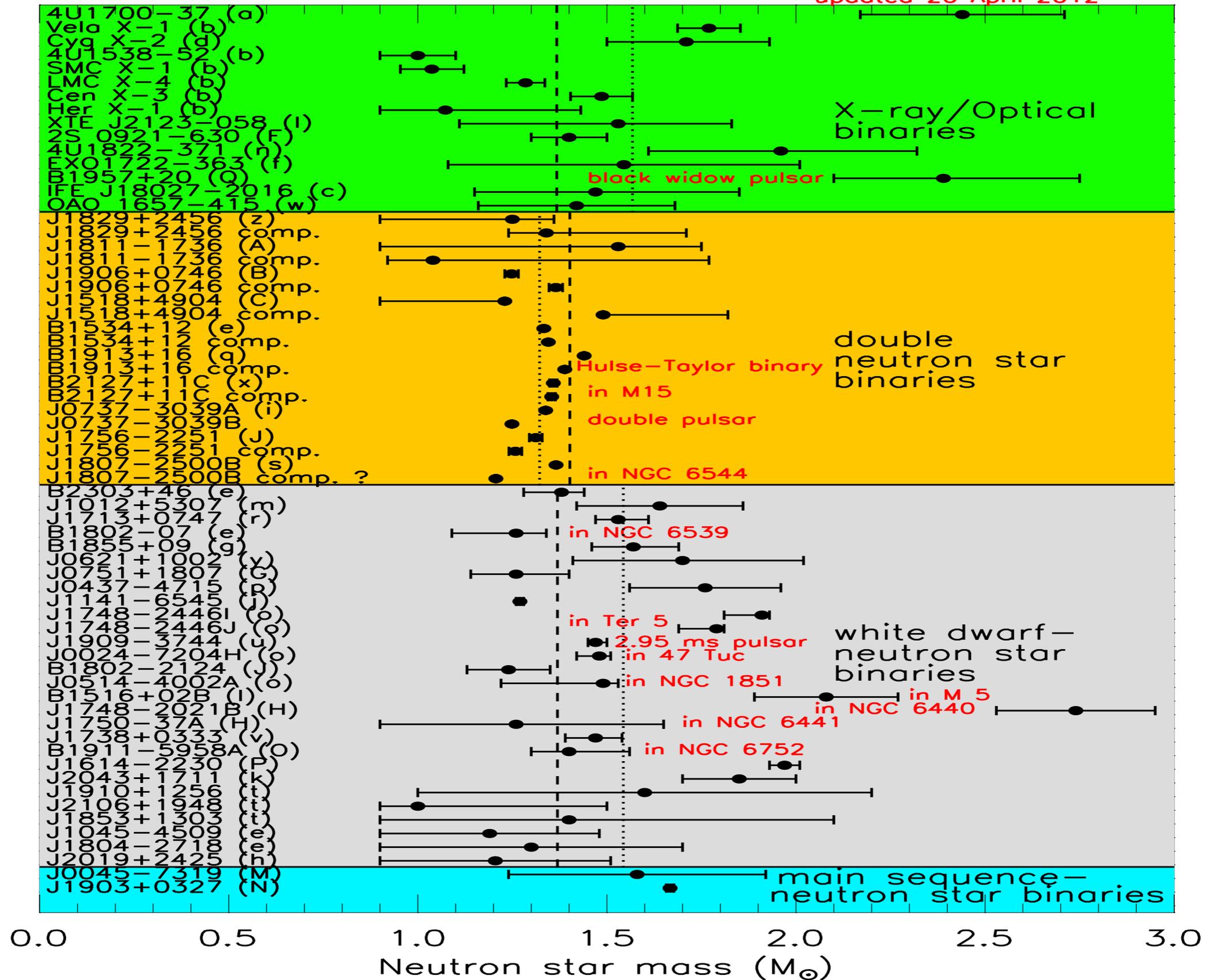
Many options for the EoS,  
leading to different mass and radii.

Various requirements:  
The EoS should describe matter above  
the saturation point, reproduce the  
observed NSs masses etc.

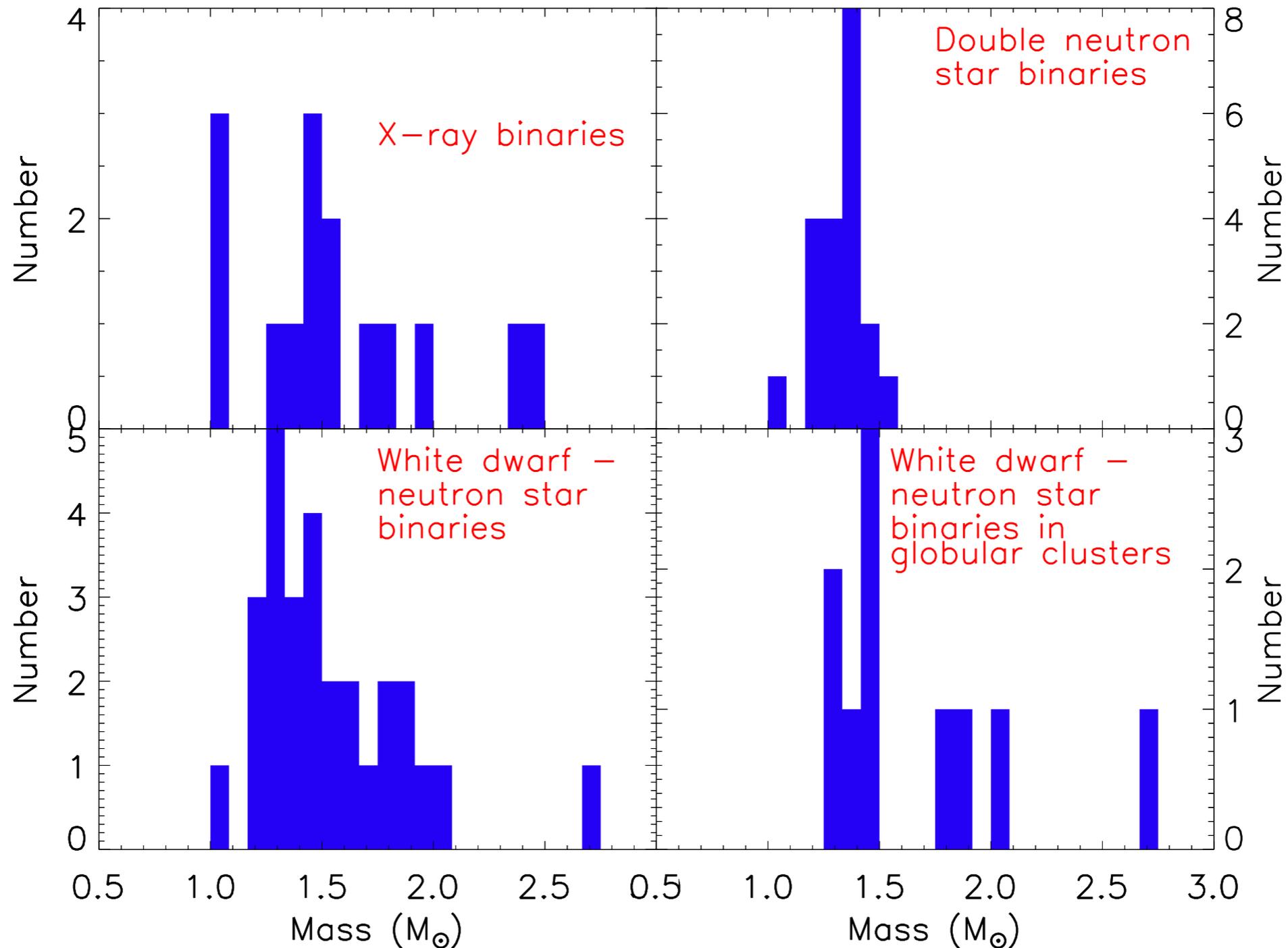
# Mass distribution of Neutron Stars

# Masses of known NSs

updated 20 April 2012



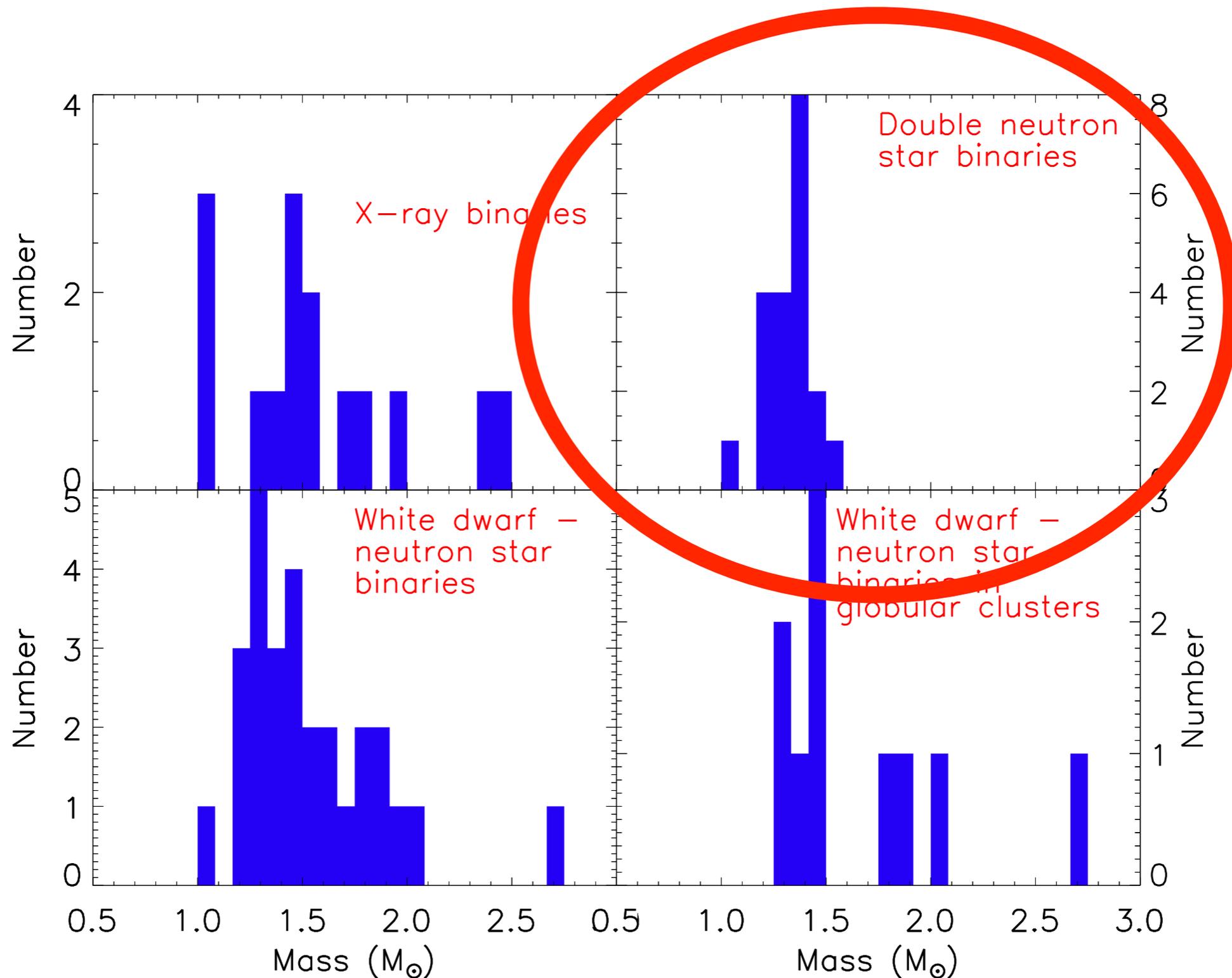
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**The reason of this peculiar distributions is unknown**

**No large mass NSs in double NS binaries!**

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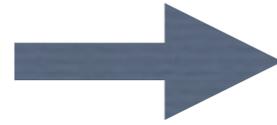


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# *Conversion of neutrons in mirror neutrons*

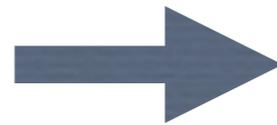
We assume that inside the NS neutrons are slowly converted in mirror neutrons



The star will (slowly) modify its equilibrium configuration.

Any small mass difference  $m_n - m_{n'}$  is irrelevant, as far as neutrons fill the Fermi sphere.

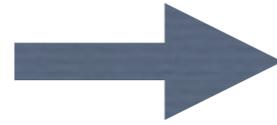
$$H \simeq \begin{pmatrix} \mathcal{E}_F & \epsilon \\ \epsilon & \mathcal{E}'_F \end{pmatrix}$$



$$\bar{P}_{nn'} = \left( \frac{\epsilon}{\mathcal{E}_F - \mathcal{E}'_F} \right)^2$$

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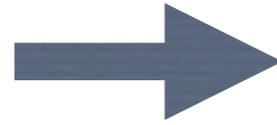
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The collision rate between neutrons is

$$\Gamma = \langle \sigma_{nn} v N \rangle \simeq 10^{24} \times (N/N_{\text{nuc}})^{4/3} \text{s}^{-1}$$

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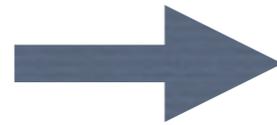
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In each collision a neutron is turned in a mirror neutron with probability  $\bar{P}_{nn'}$

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The conversion rate is therefore  $\Gamma_{n'} \simeq \left( \frac{\epsilon}{10^{-11} \text{eV}} \right)^2 \times 10^{-6} \text{yr}^{-1}$

# *Conversion of neutrons in mirror neutrons*

For the stellar evolution we also assume

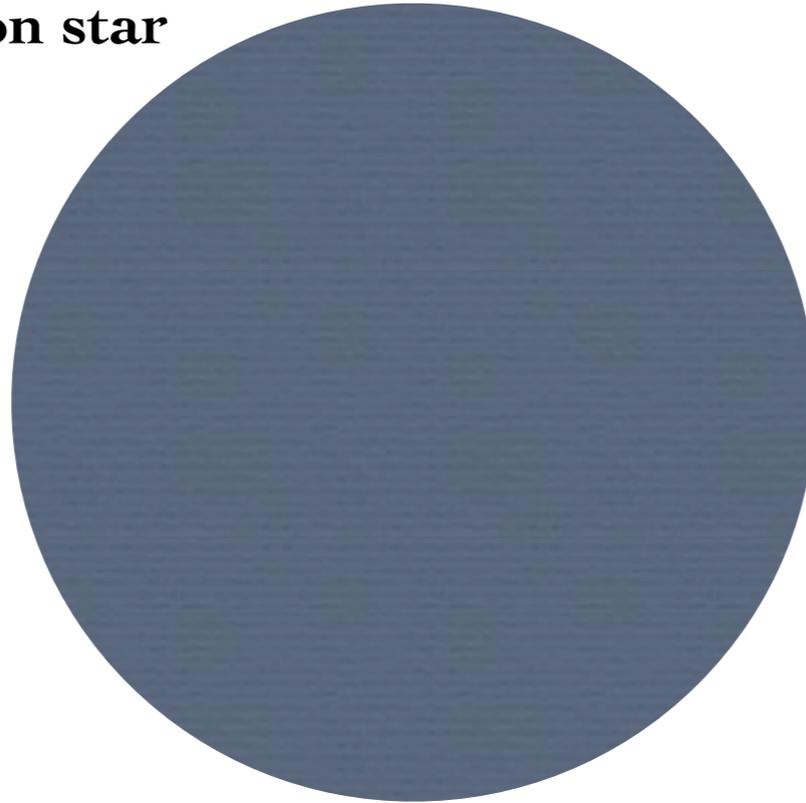
- NM and MM behave as independent fluids only coupled by gravity
- NM and MM are not, in general, in chemical equilibrium
- NM and MM have the same equation of state (EoS)
- NM and MM form two concentric stars

In other words:

MM is gravitationally trapped and the system is approximated as a two fluid noninteracting system

# *Qualitative expectations*

Neutron star



Nuclear  
matter



Mirror  
matter



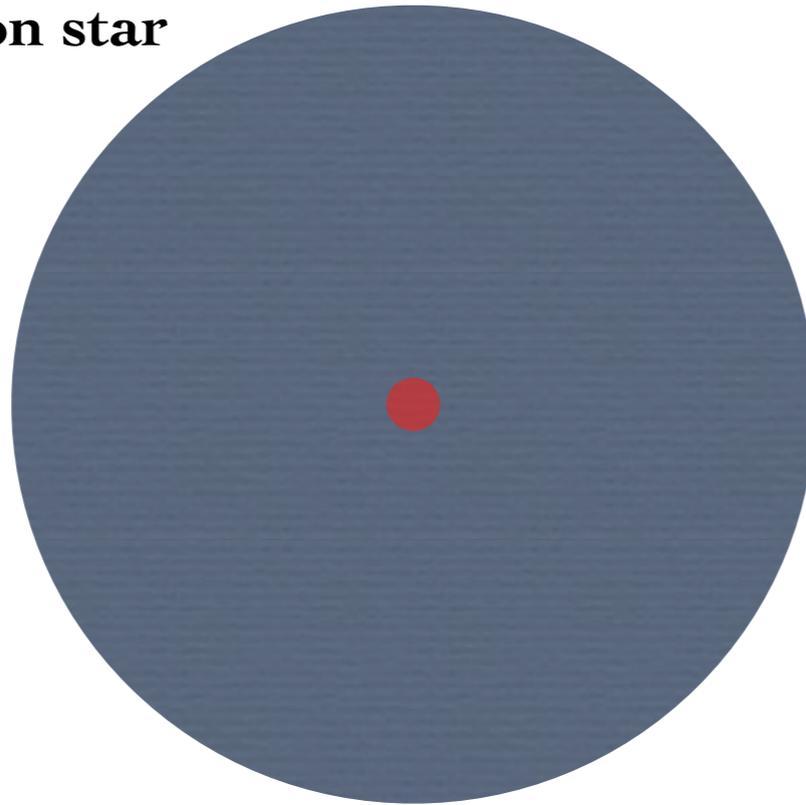
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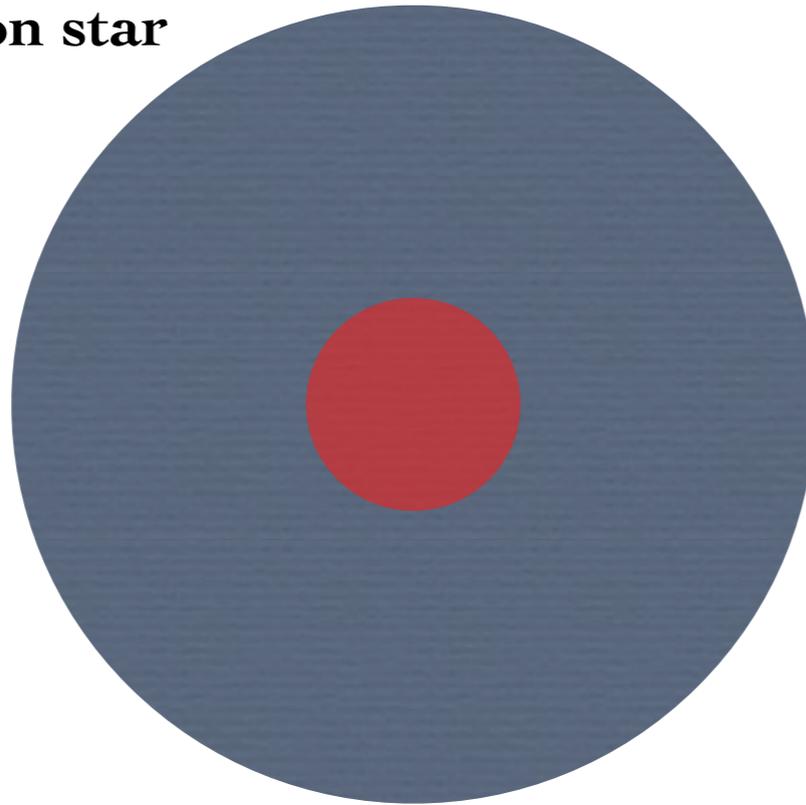
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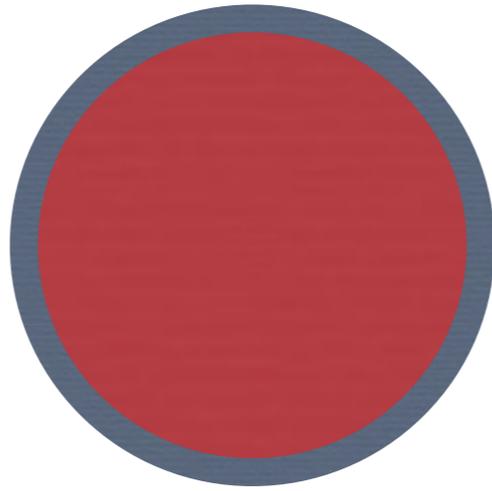
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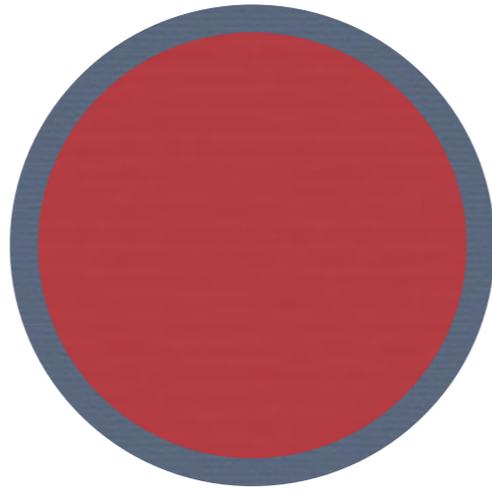
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The conversion speed depends  
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We expect to be slow, but  
comparable with typical NS  
life-time

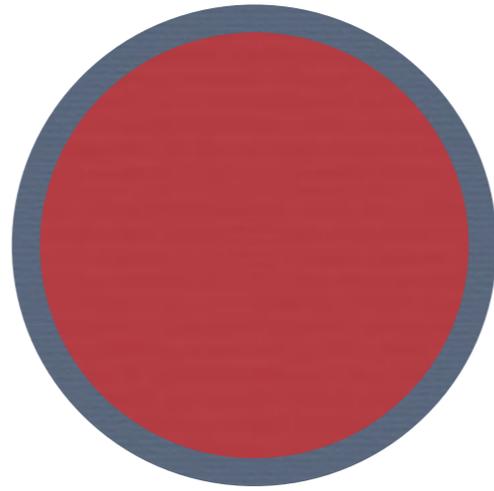
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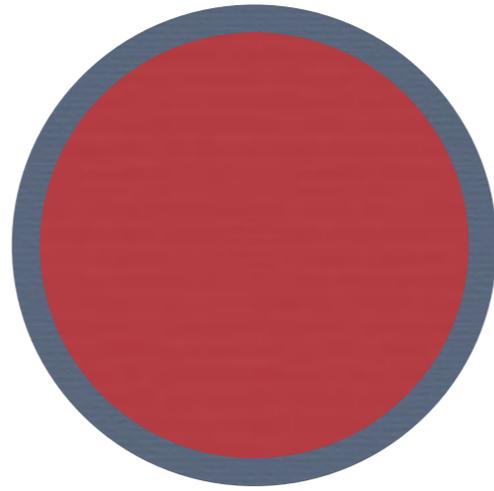
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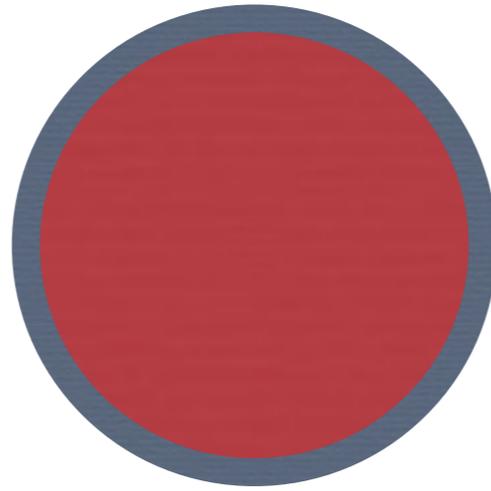
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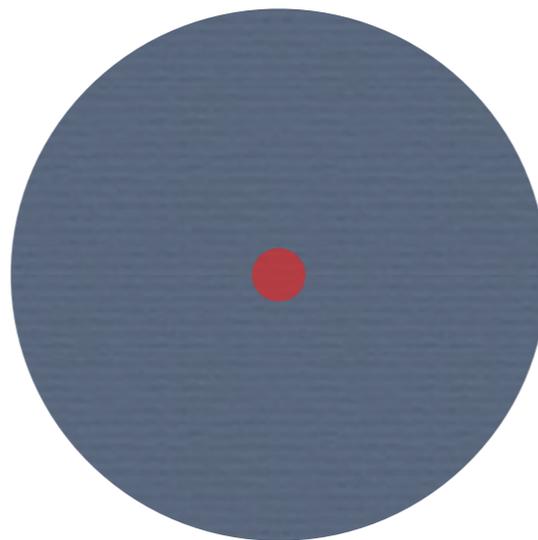
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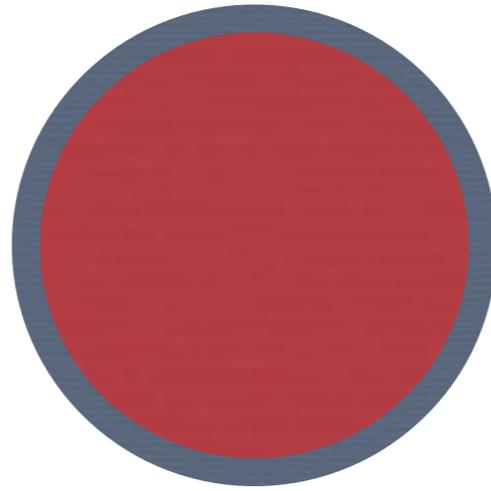
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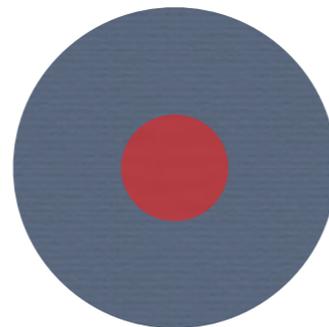
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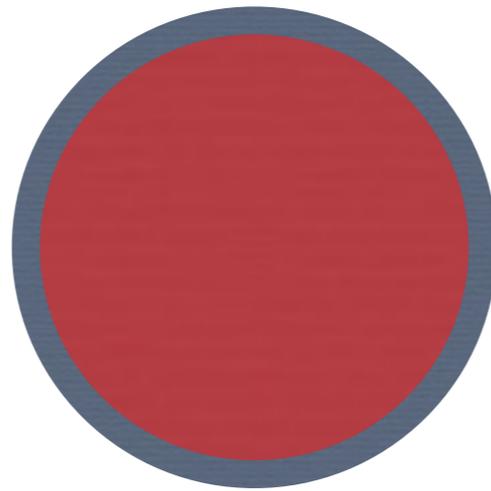
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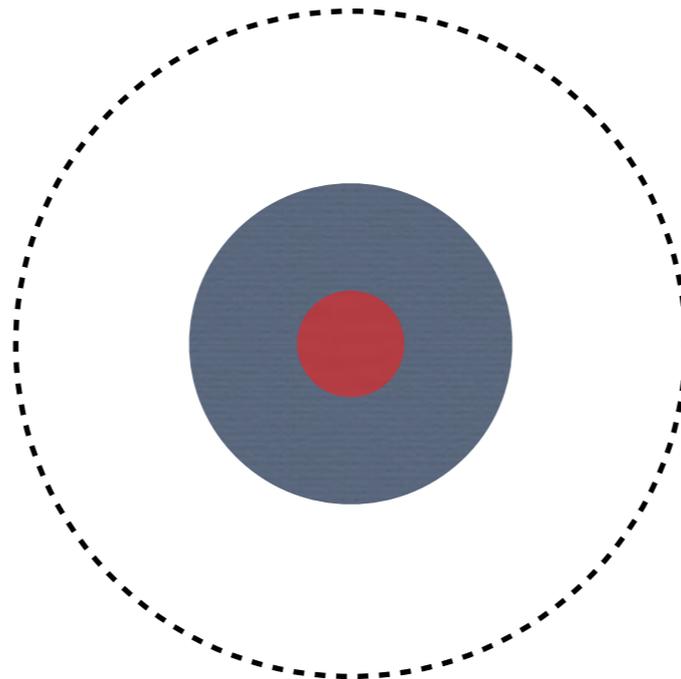
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The final configuration is a NS with smaller radius and a smaller mass, or...



BH horizon



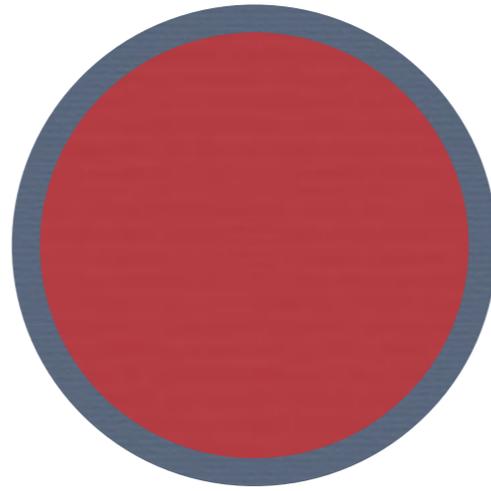
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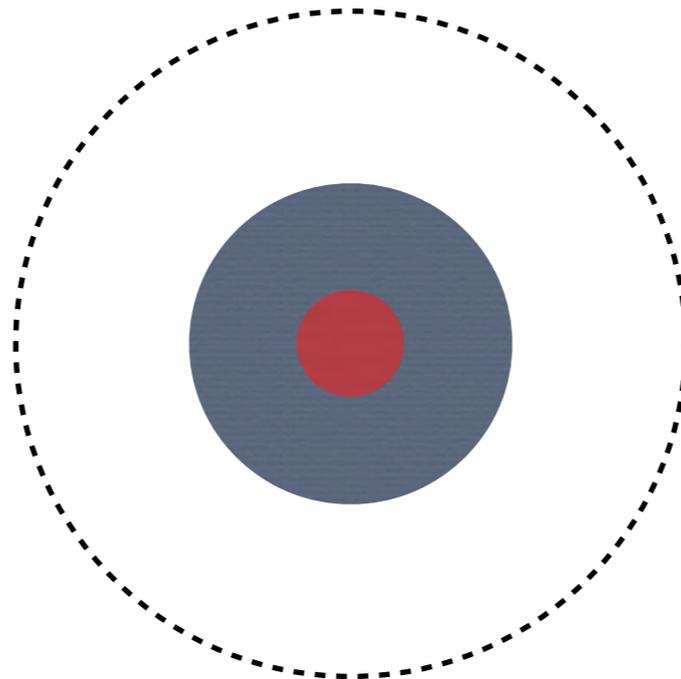


The conversion speed depends on many factors.

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The final configuration is a NS with smaller radius and a smaller mass, or...

A black hole (BH) triggered by radial unstable oscillations



BH horizon



# *Equilibrium configurations*

The TOV's equations of hydrostatic equilibrium for one fluid

$$\frac{dm}{dr} = 4\pi\rho r^2$$
$$\frac{dp}{dr} = (\rho + p) \frac{m + 4\pi p r^3}{2mr - r^2}$$

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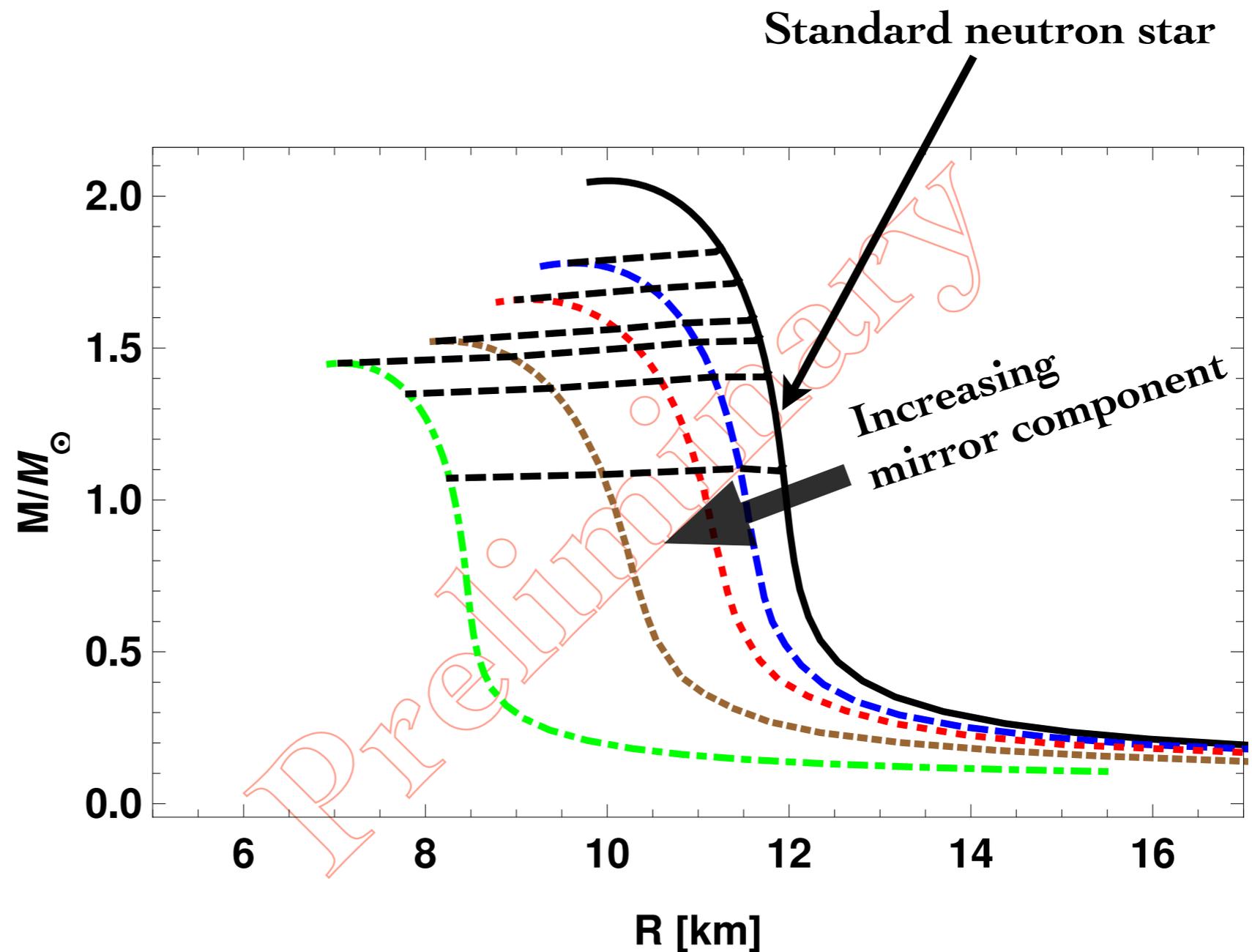
$$\frac{dm}{dr} = 4\pi\rho r^2$$
$$\frac{dp}{dr} = (\rho + p) \frac{m + 4\pi p r^3}{2mr - r^2}$$

are modified to include coupled (by gravity) TOV's equations for a two fluid system

$$\frac{dm_1}{dr} = 4\pi\rho_1 r^2 \quad \text{for } r \leq R_1$$
$$\frac{dm_2}{dr} = 4\pi\rho_2 r^2$$
$$\frac{dp_2}{dr} = (\rho_2 + p_2) \frac{m_1 + m_2 + 4\pi(p_1 + p_2)r^3}{2(m_1 + m_2)r - r^2}$$
$$\frac{dp_1}{dr} = \frac{\rho_1 + p_1}{\rho_2 + p_2} \frac{dp_2}{dr} \quad \text{for } r \leq R_1$$

# Conversion of neutrons in mirror neutrons

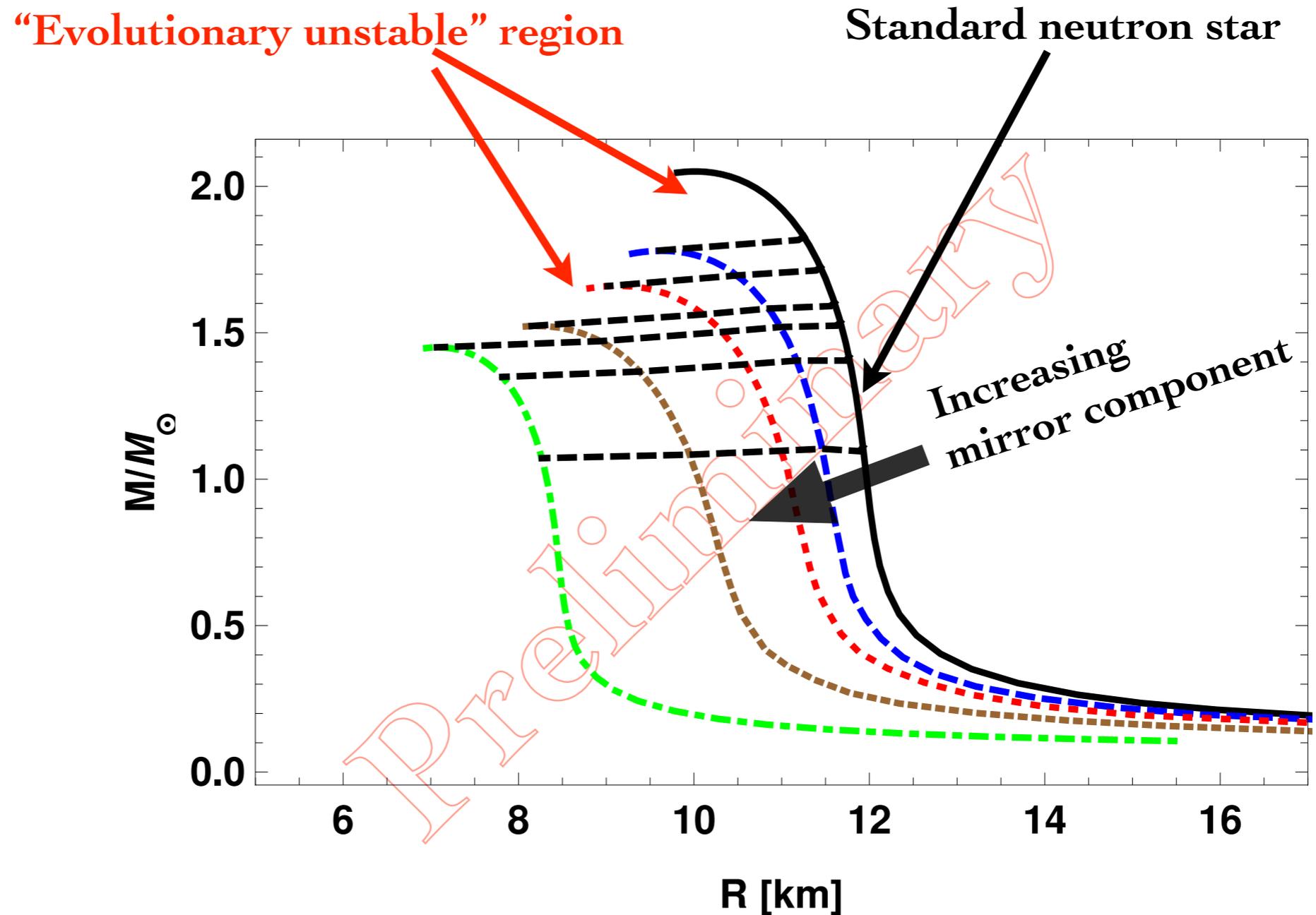
Solving the coupled TOV's equations we determine the equilibrium configurations



Dashed line correspond to constant baryonic + mirror matter number

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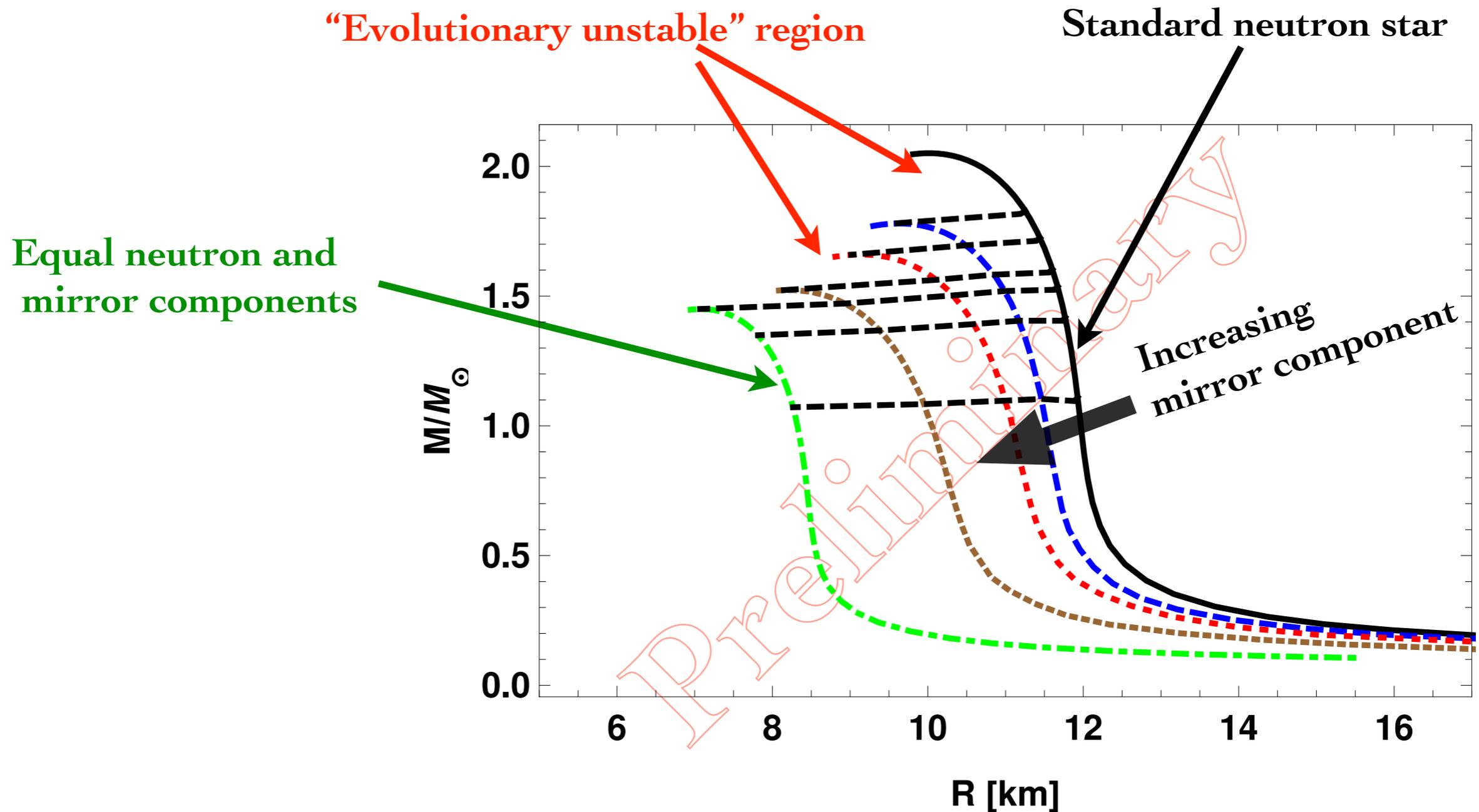
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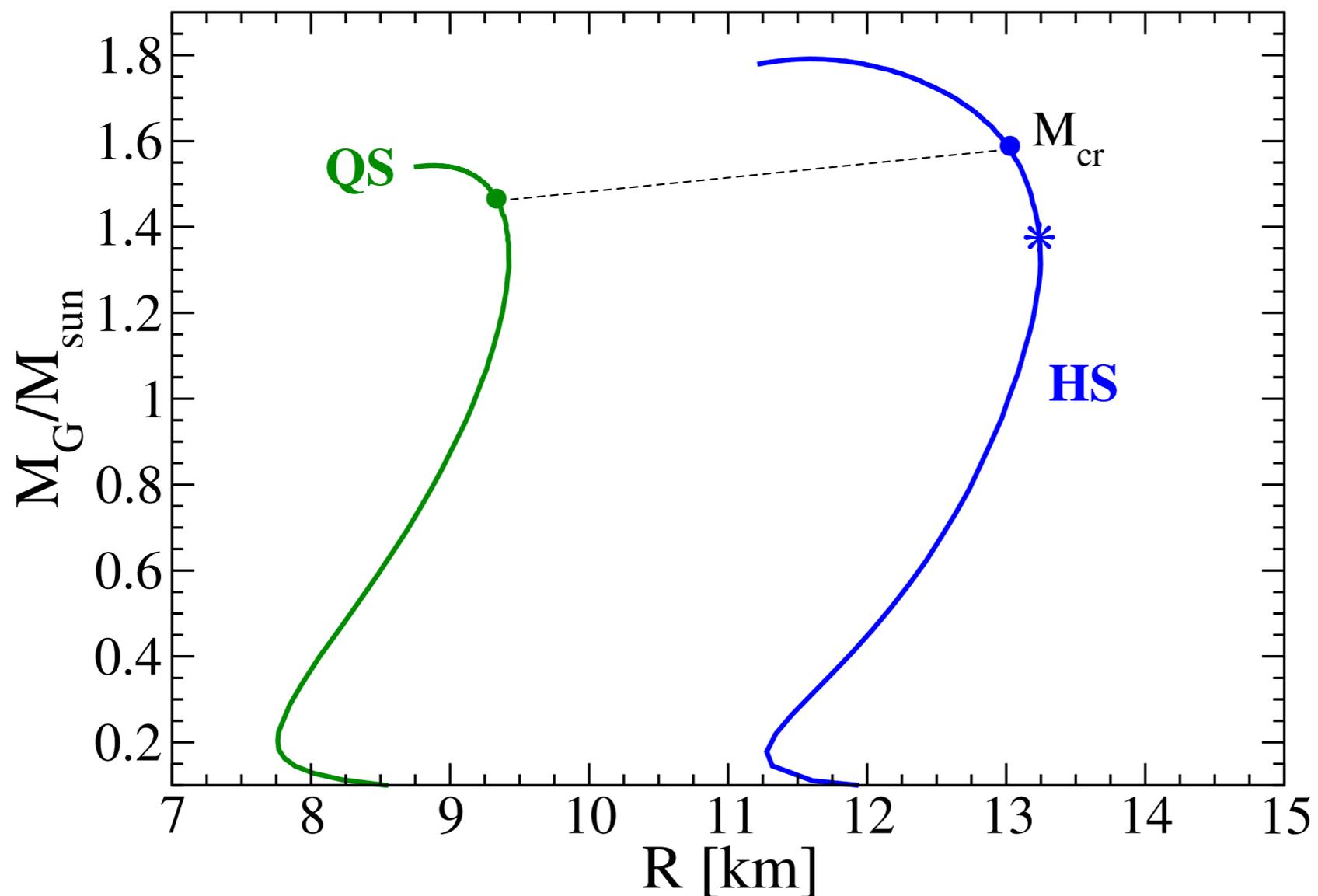


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# *Twin stars*

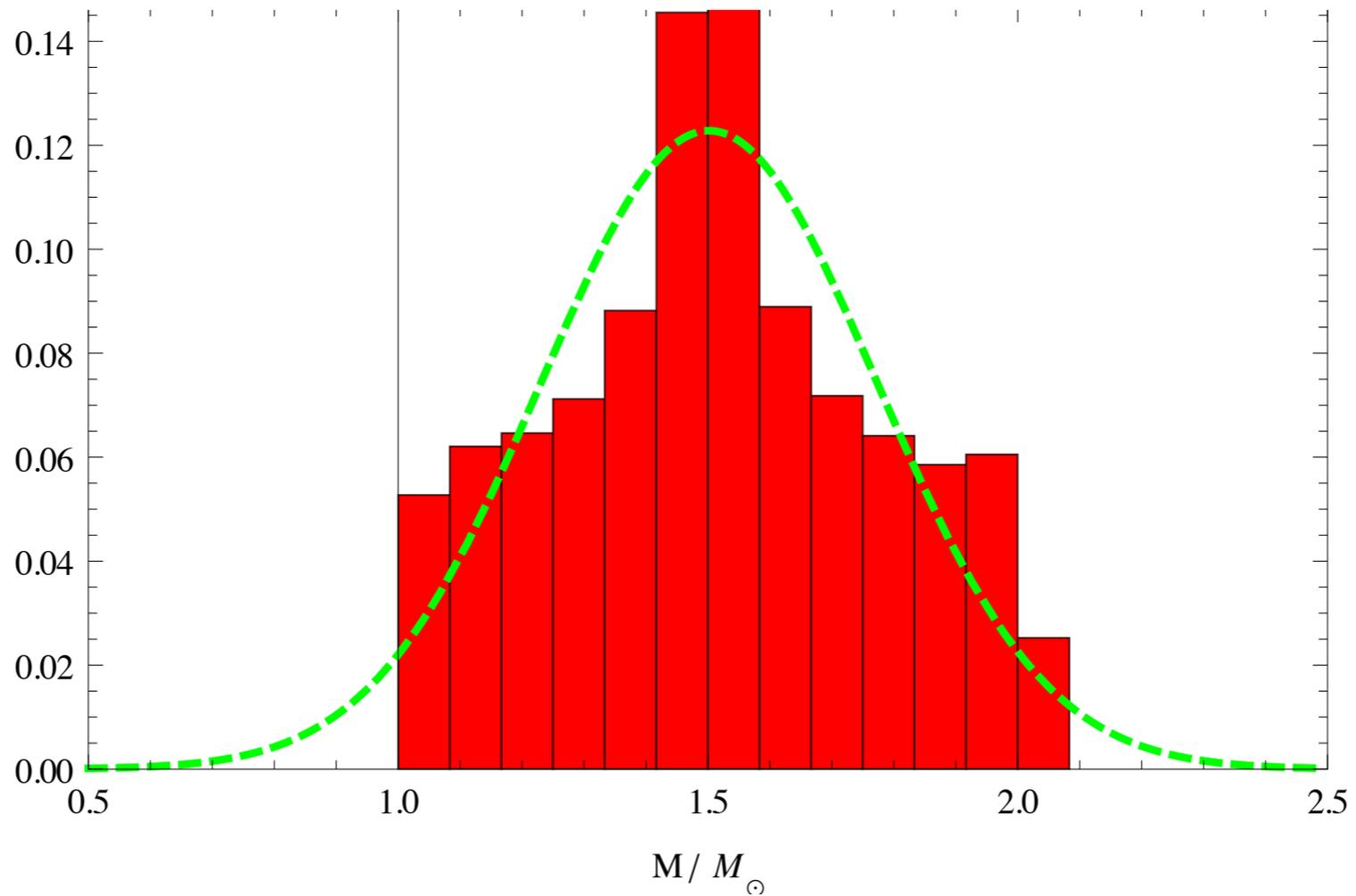
Assuming that during the stellar evolution the standard hadronic matter is converted in a different kind of matter one can obtain twin stars

Example: conversion of Neutron Stars in Hybrid Stars (with a deconfined quark core)



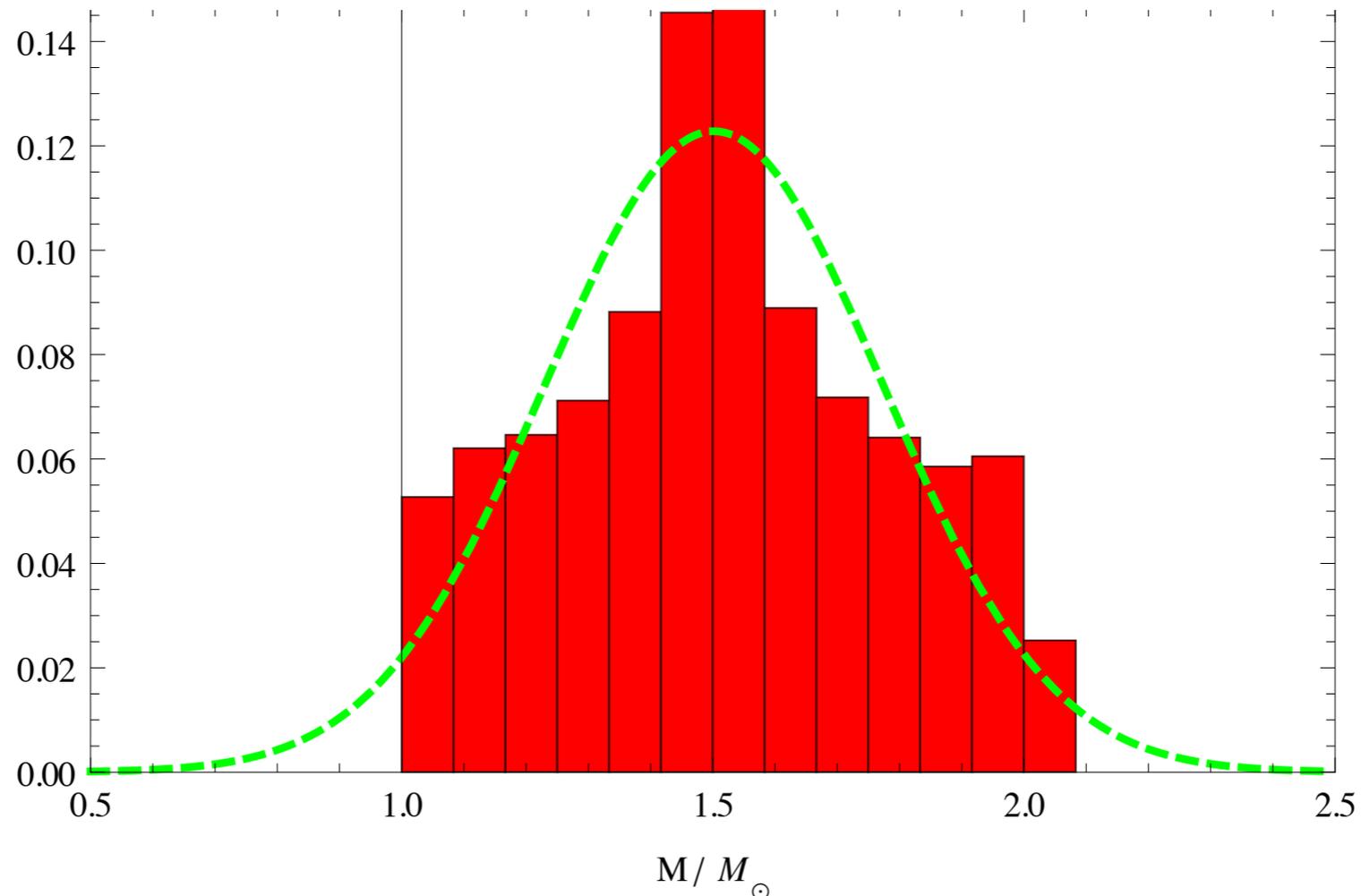
# *Procedure*

We generate NSs at a constant rate with an initial mass distribution:



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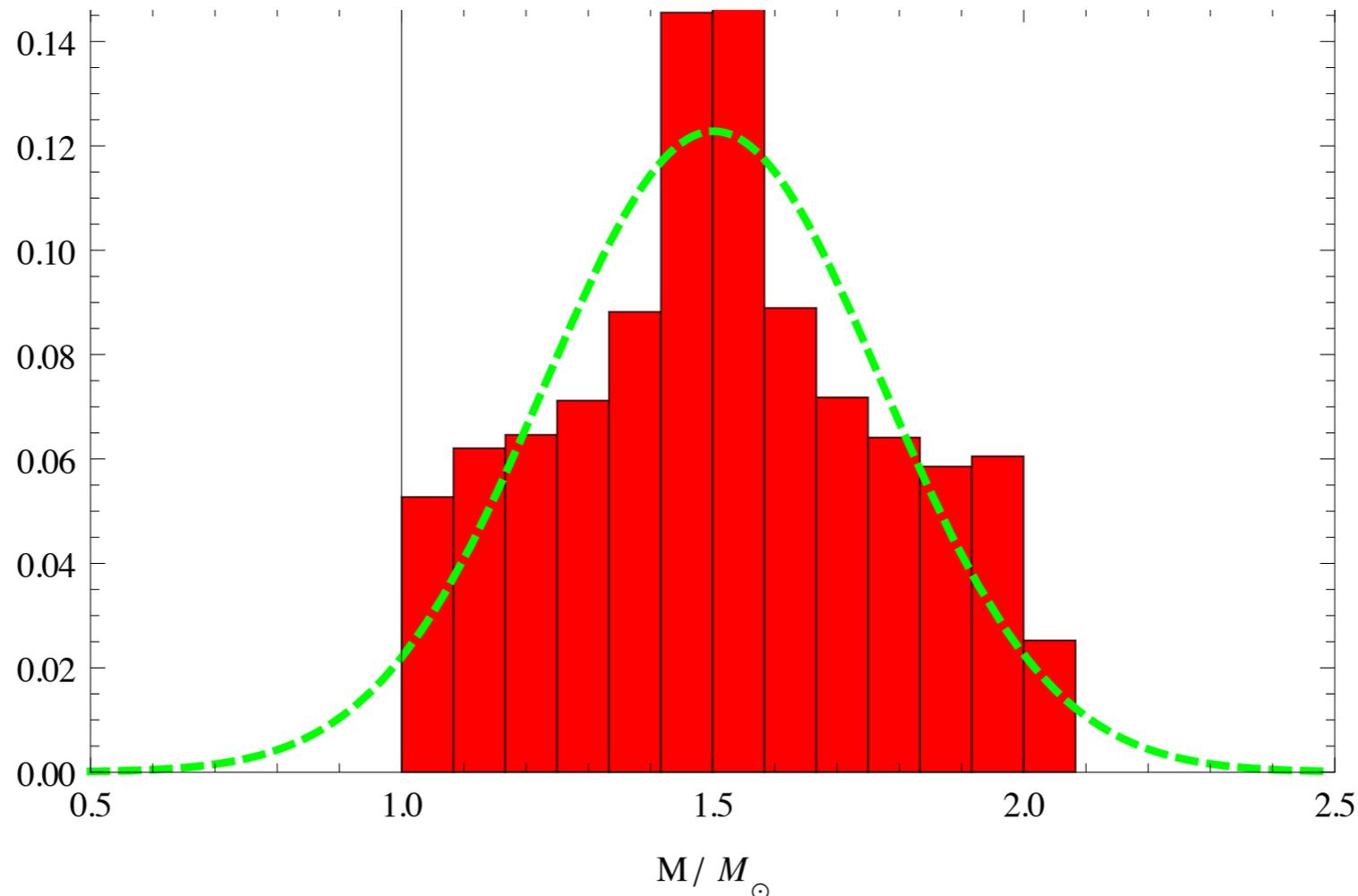


## Time evolution:

- 1) At every time step a certain number of NSs is created and evolves along a total constant baryonic number line.
- 2) We turn off the creation rate and allow the system to evolve.

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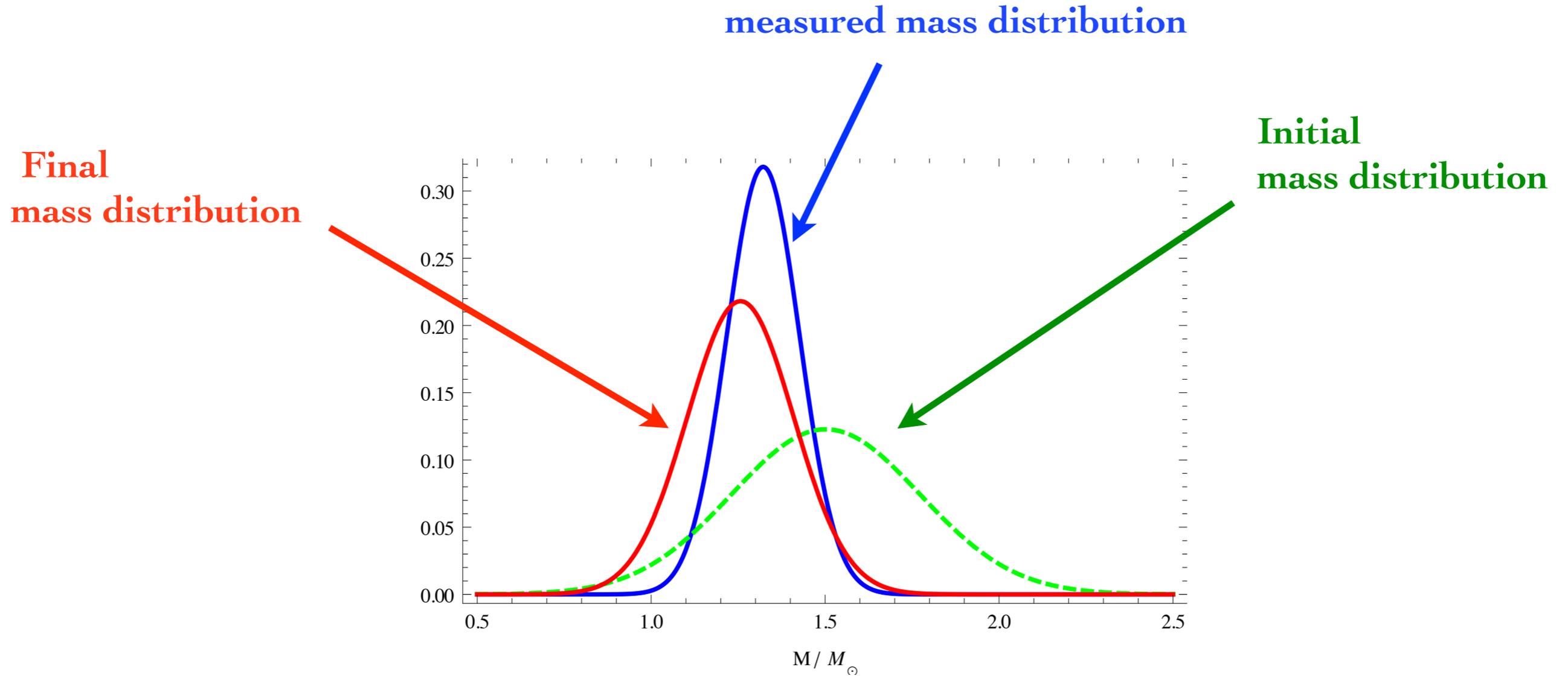


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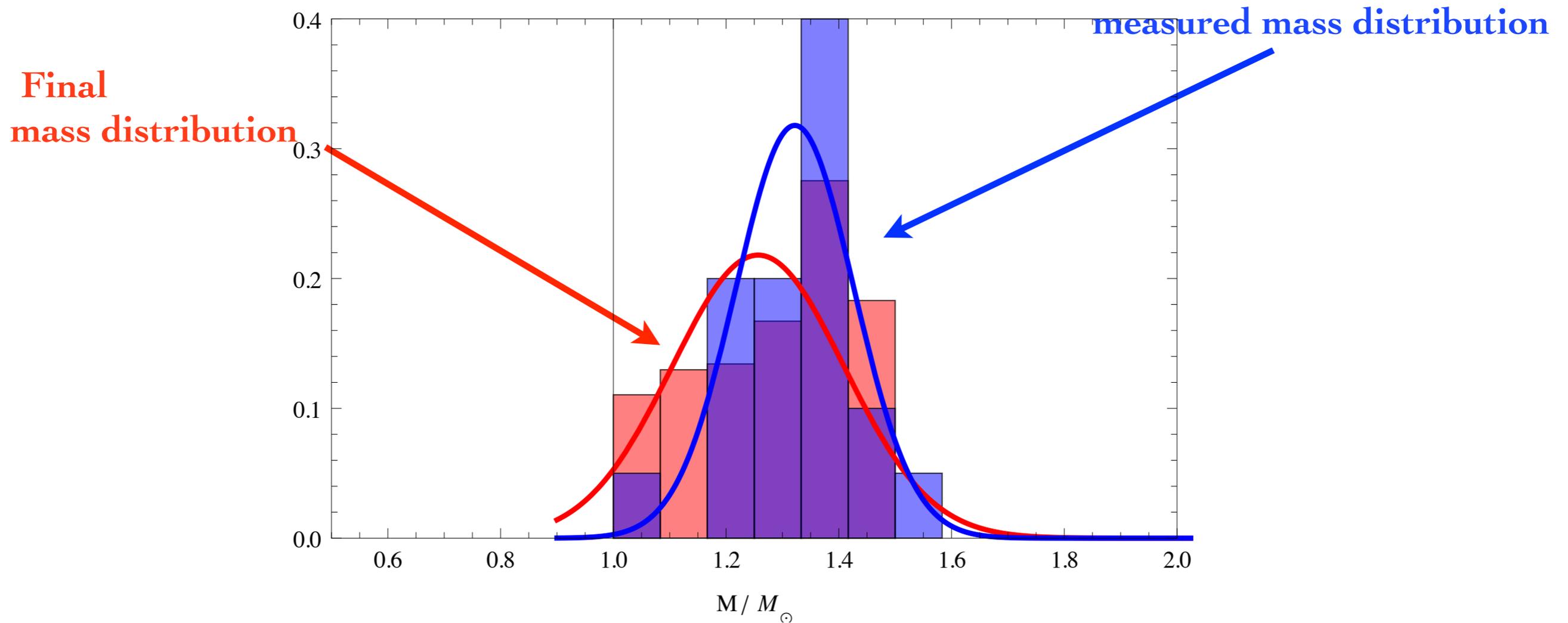
We tried different initial distributions and creation rates, obtaining similar results.

# *Mass distribution of NSs*



Thanks to the “decay” of massive NSs, the peak moves to the left and shrinks

# *Mass distribution of NSs*



One interesting aspect is that both the measured mass distribution and our Final mass distribution cannot be well fitted by a Gaussian. They have a similar asymmetry: in both distributions the high mass NSs are suppressed.

# *Conclusions*

- Neutron star masses are distributed in a peculiar way
- Assuming that standard baryonic matter inside NSs slowly evolves suppresses the large mass NSs
- The NSs mass distribution then shrinks and is asymmetric
- NSs born with a mass close to two solar masses may become unstable by the conversion of neutrons in mirror neutrons

**BACKUP**

# *Taxonomy of compact stars*

## Neutron star



$$R \sim 10 \text{ km} \quad M = 1 - 2 M_{\odot}$$

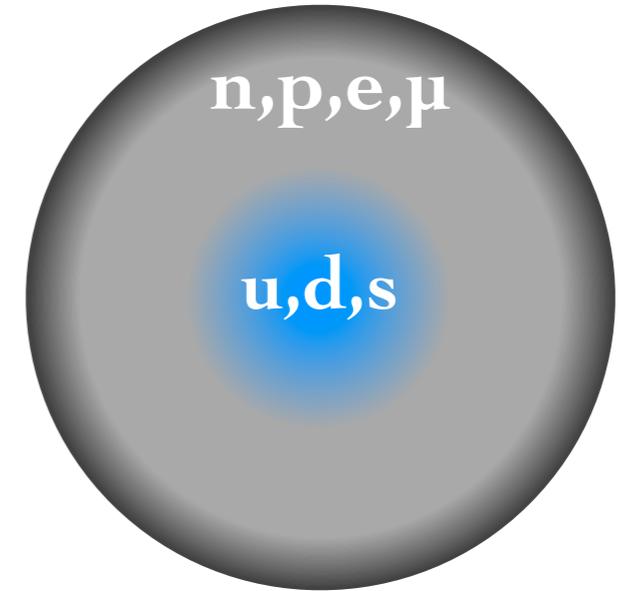
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Neutron star



$R \sim 10 \text{ km}$   $M = 1 - 2 M_{\odot}$

Hybrid star



$R \sim 10 \text{ km}$   $M = 1 - 2 M_{\odot}$

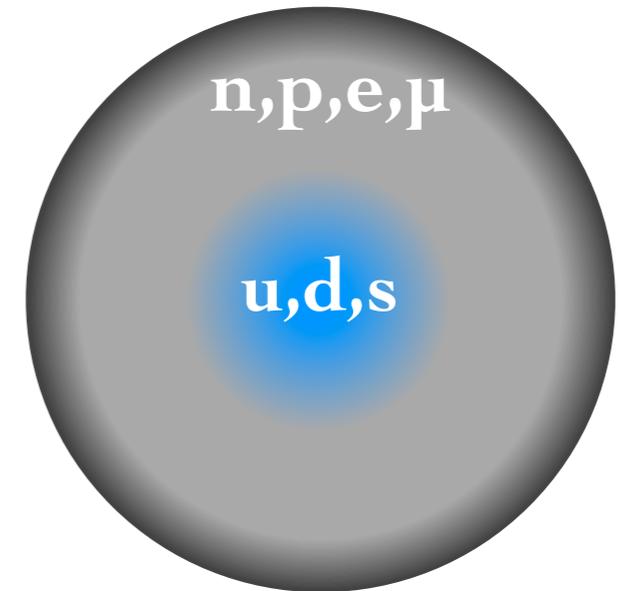
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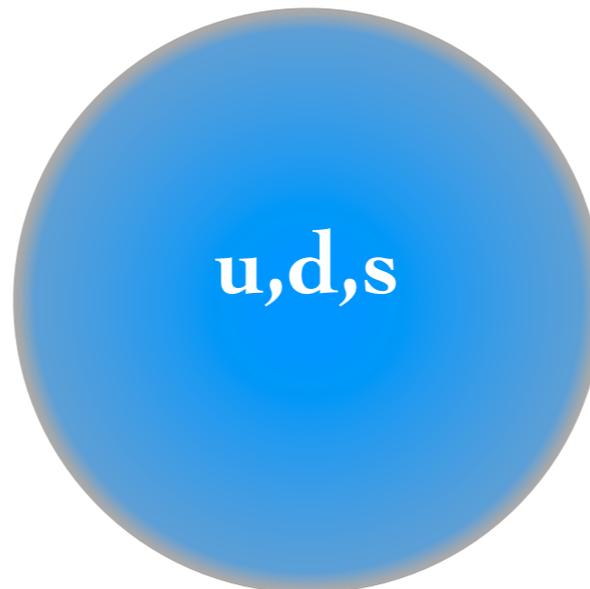
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**Hybrid star**



$$R \sim 10 \text{ km} \quad M = 1 - 2 M_{\odot}$$

**Strange star**



$$R \sim 0 - 10 \text{ km} \quad M < 3 M_{\odot}$$

# Increasing baryonic density

