Lecture Notes on Quantum information processing with photons and atoms

Yu-Ao Chen

CAS Center for Excellence in Quantum Information and Quantum Physics University of Science and Technology of China

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Introduction to our group

Excellence Center for Quantum Information and Quantum Physics

• Jointly supported by CAS and the Ministry of Education

Hosted by **USTC**

includes top institutes and universities on quantum physics



and excellence groups among China's universities: Tsinghua University, Peking University, Jiaotong University, etc.

Introduction to our group





China's Future National Projects

The Center is now playing a leading role in organizing

- National Science and Technology Project on Quantum Information in the next 15 years, similar to European Quantum Technologies Flagship
- National Laboratory for Quantum Information Sciences (NLQIS)







Global Quantum Communication Networks Scalable Quantum Computation and Quantum Simulation

Super-resolution Quantum Metrology







Ph. D & Postdoc position available

Lecture 1: Introduction to Quantum Physics and Quantum Information

Part 1: Quantum Foundations

Information Exchange in Human Evolution

About 100,000 years ago, Homo Neanderthalian and Homo Sapien co-existed in Europe

Homo Neanderthalian 🌩

- Stronger
- With even larger cranial capacity than modern human!



Homo Sapiens
Individually weaker than
Homo Neanderthalian, but
developed proto- symbol and
language

Why did Homo Sapien win the evolution battle and become our ancestors?

Information exchange



Coordinating groups

Social Advancement & Privacy Protection

Privacy Freedom of thoughts **Innovation and social advancement!**



Social progress makes information exchange more efficient

The information exchange has been and will continue to be accompanied by the human evolution and social development

Ever lasting questions:

- How to make information exchange more efficient?
- How to protect privacy?



Challenges in Information Security





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Ancient Greek scytale, 400 BC

《送杜少府之任蜀川》 1请弓、2请箭、3请刀、 ---王勃 5请枪旗、6请 4请甲 7请马、 风烟望五津。 锅幕 8 请衣 城阙辅三秦, 9请粮料、 赐、 10请草 与君离别意, 同是宦游人。 料、...、 39都将病、 天涯若比邻。 海内存知己, 战小胜 无为<u>在歧路,</u> 儿女共沾巾。

(D) (B) (D) (R) (22)

11th century, 北宋《武经总要》: 替换式密码

16th century, France, Vigenère cipher

Challenges in Information Security



20th Century, Switzerland



Enigma Machine



Classical encryption based on computational complexity

Challenges in Information Security



Crack via variations in the frequency of the occurrence of letters, by Al-Kindi (800-873)



Enigma machine broken by Alan Turing's Bombe machine

RSA 512: cracked in 1999

RSA 768: cracked in 2009

RSA 1024: ? shall not be used after December 31, 2013 by NIST

- The next-generation code "pairing-based cryptography" Cracked in 2012.....
- Feb. 2017, SHA-1 cracked by Google

All the classical encryption methods that depend on computational complexity, can be cracked in principle!

".....human ingenuity cannot concoct a cipher which human ingenuity cannot resolve"

-A few words on secret writing, Edgar Alan Poe (1841)

Challenges in the Computational Capacity

Classical computational bottleneck

The world's total computing power is insufficient to search a target in 280-90 database within a year

A technological limit

The Moore's law that predicts the transistor density doubles every 18 months has come to an end



0.2 nm (atomic scale) ⇒ ???



Tunneling induced leakage → The "0/1" logic in the transistors will fail

Quantum physics, after one century's development, comes to the rescue for the problems confronted in the classical information technologies



Max Planck

Albert Einstein

Niels Bohr

Erwin Schrödinger Werner Heisenberg

Paul Dirac







Classical World |*here*> or |*there*>

Quantum World |*here*> + |*there*>

A "quantum flight": from Shanghai to Stockholm, two possible routes:



When arrived

If I fell asleep during flight (do not know which route I take), I will feel : "both cold and warm"

I took both routes in one flight!?

 If I was awake during flight and checked which route I take, I will
 It confirms I can only take one of the routes!
 Feel either cold (Moscow) or warm (Singapore)

In quantum world, the state of a quantum object can be affected by measurement!

When Classical Physics Meets Life Philosophy

Newton's law precisely predicts every single movement for all objects in our daily life



- A manifest of the beauty and power of physics!
- However, does it imply determinism?
- Does it mean everything (e.g. lectures today) is already determined from Big-bang?
- Efforts meaningless?
- Fortunately, quantum mechanics tells that your act (measurement) can affect the world!

Qubit & Quantum Superposition









Qubits: Polarization of Single Photon

One bit of information per photon (encoded in polarization)

$$|H>=|"0">$$

 $|V>=|"1">$



Non-cloning theorem:

An unknown quantum state can not be copied precisely!



Single-Qubit Operations

Pauli matrix $\sigma_X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

Column vector represent of twodimensional quantum states

$$|W\rangle = \begin{pmatrix} 1\\ 0 \end{pmatrix} \qquad |V\rangle = 0$$

Pauli matrix
$$\sigma_Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$
 Pauli matrix $\sigma_Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

 0°

Two eigenstates: Two eigenstates: $|L\rangle = 1/\sqrt{2} \left(|H\rangle - i|V\rangle\right)$ $|+\rangle = 1/\sqrt{2} (|H\rangle + |V\rangle)$ $|R\rangle = 1/\sqrt{2} \left(|H\rangle + i|V\rangle\right)$ $|-\rangle = 1/\sqrt{2} (|H\rangle - |V\rangle)$ Unitary rotation: Unitary rotation: $\sigma_X |H\rangle = |V\rangle$ $\sigma_Y|H\rangle = -i|V\rangle$ $\sigma_V |V\rangle = i |H\rangle$ $\sigma_X |V\rangle = |H\rangle$

Two eigenstates: $|H\rangle$ $|V\rangle$ Unitary rotation: $\sigma_Z |H\rangle = |H\rangle$ $\sigma_Z |V\rangle = -|V\rangle$

Single-Qubit Operations





Mach-Zehnder Interferometer



$$H\Phi H \left| 0 \right\rangle = \frac{1}{2} \left(\left(e^{i\varphi} + 1 \right) \left| 0 \right\rangle + \left(e^{i\varphi} - 1 \right) \left| 1 \right\rangle \right)$$

Mach-Zehnder Interferometer



Interaction-free measurement!

Zeno Paradox



Origin of Zeno effect

Can the rabbit overtake the turtle?

Quantum Zeno Effect



Considering neutron spin evolving in magnetic field, the probability to find it still in spin up state after time T is

$$P = \cos^2\left(\frac{\omega T}{2}\right)$$

where ω is the Larmor frequency

Quantum Zeno Effect



If we cut the bad part of the cake at time $T=\pi/2$, then at $T=\pi$ we have $G=1/4\times G_0$
Experiment



Kwiat et al., PRL 74, 4763 (1995)

Kwiat, et al., PRL 83 4725 (1999)

Quantum Zeno Effect



More complex strcture: A nested and chained version of MZI

All-Pass:

$$|100\rangle \rightarrow \cos^{m-1} q_{M} \left(\cos q_{M} |100\rangle + \sin q_{M} |010\rangle \right)$$

$$\xrightarrow{m=M} |100\rangle$$
All-Block:
$$|100\rangle \rightarrow \cos m q_{M} |100\rangle + \sin m q_{M} |010\rangle$$

 $\xrightarrow{m=M} |010\rangle$

Salih, et al., PRL 110, 170502 (2013)

Quantum entanglement

Quantum entanglement $|1\rangle$ $|1\rangle$ + $|1\rangle$

Bell states maximally entangled states:

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$$|\mathsf{F}^{\pm}\rangle_{12} = \frac{1}{\sqrt{2}} \left(|H\rangle_{1} |H\rangle_{2} \pm |V\rangle_{1} |V\rangle_{2} \right)$$
$$|\mathsf{Y}^{\pm}\rangle_{12} = \frac{1}{\sqrt{2}} \left(|H\rangle_{1} |V\rangle_{2} \pm |V\rangle_{1} |H\rangle_{2} \right)$$



Two dices are entangled

Spooky action at a distance --Albert Einstein

Quantum entanglement

GHZ states: three-photon maximally entangled states

$$\begin{split} |\Phi^{\pm}\rangle_{123} &= \frac{1}{\sqrt{2}} (|H\rangle_1 |H\rangle_2 |H\rangle_3 \pm |V\rangle_1 |V\rangle_2 |V\rangle_3) \\ |\Psi^{\pm}\rangle_{123} &= \frac{1}{\sqrt{2}} (|H\rangle_1 |H\rangle_2 |V\rangle_3 \pm |V\rangle_1 |V\rangle_2 |H\rangle_3) \\ |\Xi^{\pm}\rangle_{123} &= \frac{1}{\sqrt{2}} (|H\rangle_1 |V\rangle_2 |H\rangle_3 \pm |V\rangle_1 |H\rangle_2 |V\rangle_3) \\ |\Theta^{\pm}\rangle_{123} &= \frac{1}{\sqrt{2}} (|H\rangle_1 |V\rangle_2 |V\rangle_3 \pm |V\rangle_1 |H\rangle_2 |H\rangle_3) \end{split}$$

Manipulation of Entanglement



 $00 \to (0+1)0 = 00 + 10 \to (00 + 11)$ $\to (00 + 10) = (0 + 1)0 \to 00$

Manipulation of Entanglement



 $\begin{array}{l} 000 \rightarrow (0+1)00 = (00+10)0 \rightarrow (00+11)0 = 000+110 \\ \rightarrow 000 + 111 \rightarrow 000 + 110 = (00+11)0 \rightarrow (00+10)0 \\ = (0+1)00 \rightarrow 000 \end{array}$

Spooky Action at a Distance?



Quantum entanglement:



Quantum Non-locality

Measurement on particle A will cause instant collapse on particle B

Spooky Action at a Distance?





Quantum mechanics is certainly imposing. But an inner voice tells me that it is not yet the real thing. The theory says a lot, but does not really bring us any closer to the secret of the 'old one'. I, at any rate, am convinced that He does not throw dice.

> Einstein, stop telling God what to do!



Bell's Inequality: Testing This Battle

Experimental testable inequality: Bell, Physics 1, 195 (1964) Clauser et al., PRL 23, 880 (1969)





- $S = |E(\phi_A \phi_A) E(\phi_A \phi_B') + E(\phi_A' \phi_B) + E(\phi_A' \phi_B')|$
- Einstein's local realism: $S_{\text{max}} \leq 2$
- Quantum mechanics: $S_{\text{max}} = 2\sqrt{2}$

Bell's Inequality: Testing This Battle

A simplified case: Sakurai's Bell Inequality



Singlet state: anti-correlation of measurement results of two sides

$$|\Psi^{-}\rangle_{12} = \frac{1}{\sqrt{2}} (|H\rangle_{1}|V\rangle_{2} - |V\rangle_{1}|H\rangle_{2})$$

Pick three arbitrary directions
a, **b**, and **c**:
 $P(a0, b0) = P_{3} + P_{4}$
 $P(a0, c0) = P_{2} + P_{4}$
 $P(c0, b0) = P_{3} + P_{7}$

 $P_3 + P_4 \le P_3 + P_4 + P_2 + P_7$

Bell's Inequality: Testing This Battle

Local realism requires: $P(a0, b0) \le P(a0, c0) + P(c0, b0)$ Quantum-mechanical prediction: $P(a0, b0) = \frac{1}{2} \sin\left(\frac{a-b}{2}\right)^2$

For example $a = 90^{\circ}$, $b = 45^{\circ}$, $c = 0^{\circ}$, the inequality would require

$$\frac{1}{2}\sin^2 45^\circ \le \frac{1}{2}\sin^2 22.5^\circ + \frac{1}{2}\sin^2 22.5^\circ \implies 0.2500 \le 0.1464!$$

An unsatisfactory feature

In the derivation of BI such a local realistic and thus classical picture can explain perfect correlations and is only in conflict with statistical prediction of quantum mechanics Conflict with Local Realism

Consider a three-photon GHZ state written in σ_Z basis

$$|\Psi_{123}\rangle = \frac{1}{\sqrt{2}} \left(|H_1\rangle|H_2\rangle|H_3\rangle + |V_1\rangle|V_2\rangle|V_3\rangle\right)$$

Linear polarization basis Cir

Circular polarization basis

$$\sigma_{x}: |H'\rangle = \frac{1}{\sqrt{2}} (|H\rangle + |V\rangle),$$
$$|V'\rangle = \frac{1}{\sqrt{2}} (|H\rangle - |V\rangle).$$

$$\begin{split} \sigma_{y} : & \left| R \right\rangle = \frac{1}{\sqrt{2}} \left(\left| H \right\rangle + i \left| V \right\rangle \right), \\ & \left| L \right\rangle = \frac{1}{\sqrt{2}} \left(\left| H \right\rangle - i \left| V \right\rangle \right). \end{split}$$

Conflict with Local Realism

$$\sigma_{1y}\sigma_{2y}\sigma_{3x}: |\psi_{123}\rangle = \frac{1}{2} \Big(R_1 L_2 H_3 + L_1 R_2 H_3 + R_1 R_2 V_3 + L_1 L_2 V_3 \Big)$$

$$\sigma_{1y}\sigma_{2x}\sigma_{3y}: |\psi_{123}\rangle = \frac{1}{2} \Big(R_1 H_2 L_3 + L_1 H_2 R_3 + R_1 V_2 R_3 + L_1 V_2 L_3 \Big)$$

$$\sigma_{1x}\sigma_{2y}\sigma_{3y}: |\psi_{123}\rangle = \frac{1}{2} \Big(H_1 R_2 L_3 + H_1 L_2 R_3 + V_1 R_2 R_3 + V_1 L_2 L_3 \Big)$$

Therefor state $|\Psi_{123}\rangle$ is the eigenstate of operators $\sigma_{1y}\sigma_{2y}\sigma_{3x}, \sigma_{1y}\sigma_{2x}\sigma_{3y}, \sigma_{1x}\sigma_{2y}\sigma_{3y}$ with value -1

Conflict with Local Realism

- EPR reality criterion: operator is predeterminethe individual value of any local d
- There exists an element of local reality S_{ix} corresponding to operator

$$\sigma_{ix}(i=1,2,3).$$

All six of the elements of reality S_{ix} and S_{iy} have to be there, each with the values +1 and -1!

$$S_{1y}S_{2y}S_{3x} = -1,$$

$$S_{1y}S_{2x}S_{3y} = -1,$$

$$S_{1x}S_{2y}S_{3y} = -1.$$

What Outcomes Are Possible?

Consider measurement of 45° linear polarization basis

Local realism:

$$S_{1x}S_{2x}S_{3x} = S_{1x}(S_{1y})^2 S_{2x}(S_{2y})^2 S_{3x}(S_{3y})^2$$

= $(S_{1x}S_{2y}S_{3y})(S_{1y}S_{2x}S_{3y})(S_{1y}S_{2y}S_{3y})$
= -1

Possible outcomes:

$$V_1'V_2V_3', H_1'H_2'V_3', H_1'V_2'H_3', V_1'H_2'H_3'$$

What Outcomes Are Possible?

Quantum physics

$$|\psi_{123}\rangle = \frac{1}{2} \left(H_1' H_2' H_3' + H_1' V_2' V_3' + V_1' H_2' V_3' + V_1' V_2' H_3' \right)$$
$$\implies S_{1x} S_{2x} S_{3x} = 1!$$

Possible outcomes:

$$H_{1}'H_{2}'H_{3}', H_{1}'V_{2}'V_{3}', V_{1}'H_{2}'V_{3}', V_{1}'V_{2}'H_{3}'$$

Whenever local realism predicts a specific result definitely to occur for a measurement for one of the photons based on the results for the other two, quantum physics definitely predicts the opposite result

Bell's Inequality: Testing the Battle



Chien-Shiung Wu

First observation of quantum entanglement

The Angular Correlation of Scattered Annihilation Radiation*

C. S. WU AND I. SHAKNOV Pupin Physics Laboratories, Columbia University, New York, New York November 21, 1949

Phys. Rev. 27, 136 (1950)





- Freedman & Clauser, PRL 28, 938 (1972)
- Fry & Thompson, PRL 37, 465 (1976)

Two measurement sites are not space-like separated



 $S_{exp} = 0.101 \pm 0.020$ violates a generalized inequality $S \le 0$ by 5 standard deviations

> Drawbacks: 1. locality loophole 2. detection loophole

- Measurement devices may "tell" the EPR source their basis choices + the source may "select" according events to violate Bell inequality
- ✓ Solution: basis choice and emission of EPR source must be also space-like separated (i. e., fast and random switch of measurement basis)

Weihs et al., PRL 81, 5039 (1998)



 $S_{\rm exp} = 2.73 \pm 0.02$ violates CHSH inequality $S \le 2$ by 30 standard deviations

Drawback: detection loophole

Example 2 Detection efficiency of single photon detectors is not unity \Rightarrow some events cannot contribute to S > 2 were not detected?

✓ Solution: high detection efficiency (>83%)

Pearle, PRD 2, 1418 (1970)

Garg & Mermin, PRD 35, 3831 (1987)

Hensen et al., Nature 526, 682 (2015)



Close both detection loophole and locality loophole

- Detection efficiency>95%
- Switch time: 480ns<493m/c

But still with loopholes...

✓ Freedom of choice loophole: random number generators (RNGs) could be prior correlated ⇒ the choice of measurement bases are not truly random Brunner et al., RMP 86, 419 (2014)



Schrödinger's cat

Collapse locality loophole: measurement outcome is not defined until it is registered by a human consciousness Realized "events" have never been space-like separated Kent, PRA 72, 012107 (2005) Leggett, Compendium of Quantum Physics (Springer, 2009) Solution for both loopholes: Bell-test experiment with human-observer!

✓ Basis choice by free will

Measurement outcomes defined by consciousness Leggett, Compendium of Quantum Physics (Springer, 2009)



Requirement:

Quantum signal transit time exceeds human reaction 100ms = entanglement distribution at a distance on the order of one light-second

Quantum Information Processing (QIP)

Test of quantum nonlocality

Coherent manipulation of quantum systems

Enabling encode and process information in quantum states, outperform classical information systems in terms of



Part 2: Quantum Communication

Single-photon-based key distribution: [Bennett & Brassard 1984 protocol]



> Entanglement-based key distribution: [Ekert, PRL 67, 661 (1991)]

Sender



BB84 Protocol



If Eve is present, the probability that Alice and Bob can not tell is $(0.25)^{N}$ after they compare N raw key's value!

BB84 Security

	one-way communication	two-way communication
Upper bound	14.6%	25%
Lower bound	11.0%	18.9%

All the error rates are brought by the eavesdropping
When the error rate is lower than the lower bound, we can utilize some classical cryptography method to let Eve know nothing about the key
If the error rate is higher than the upper bound, the key is insecure

Gottesman and Lo, IEEE TIT 49, 457 (2003)

Perfect Cipher in Principle

QKD ⇒ Secure key + One-time pad



Unconditional security!

Dense Coding

Transmit two bits of information by sending one photon



Bennett & Wiesner, PRL 69, 2881 (1992)



- 1. Alice and Bob share an entangled photon pair in the state of $|\Phi^+\rangle_{12}$
- Bob chooses one of the four unitary transformation on his photon.
 The information of which choice is 2 bit.

e. g. 00:
$$I$$
 01: σ_Z 10: σ_X 11: $-i\sigma_Y$

- 3. Bob sends his photon to Alice
- 4. Alice does a joint Bell-state measurement (BSM) on the photon from Bob and her photon.
- 5. With the measurement result, she can know Bob's unitary transformation and achieve the 2 bit information.

Quantum Teleportation



Classical physics

Scanning and reconstructing

> Quantum physics

Principle of quantum measurement forbidden extracting all the information from an unknown quantum state!

Quantum Teleportation

Initial state

$$\Phi\rangle_1 = \alpha \mid H\rangle_1 + \beta \mid V\rangle_1$$

The shared entangled pair

$$\left| \Phi^{+} \right\rangle_{23} = \frac{1}{\sqrt{2}} \left(\left| H \right\rangle_{2} \left| H \right\rangle_{3} + \left| V \right\rangle_{2} \left| V \right\rangle_{3} \right)$$

$$\begin{split} \Psi \rangle_{123} &= | \Phi \rangle_1 \otimes | \Phi^+ \rangle_{23} \\ &= | \Phi^+ \rangle_{12} \otimes \left(\alpha | H \rangle_3 + \beta | V \rangle_3 \right) + \\ &| \Phi^- \rangle_{12} \otimes \left(\alpha | H \rangle_3 - \beta | V \rangle_3 \right) + \\ &| \Psi^+ \rangle_{12} \otimes \left(\alpha | V \rangle_3 + \beta | H \rangle_3 \right) + \\ &| \Psi^- \rangle_{12} \otimes \left(\alpha | V \rangle_3 - \beta | H \rangle_3 \right) \end{split}$$



BSM results on particles 1, 2	operations on particle 3
$\left \Phi^{+} \right\rangle_{12}$	Ι
$\left \Phi^{-} \right\rangle_{12}$	σ_Z
$ \Psi^+\rangle_{12}$	σ_X
$ \Psi^{-}\rangle_{12}$	$-i\sigma_Y$

Bennett et al., PRL 73, 3801 (1993)
Quantum Teleportation



Though nowadays we can only teleport two-particle composite system..... Essential ingredient for distributed quantum information processing!

Part 3: Quantum Computation and Quantum Metrology

Quantum Computation

Quantum Parallelism



Evaluating function f(x) for many different x simultaneously

$$U\frac{1}{\sqrt{2^{N}}}\sum_{i=0}^{2^{N}-1}|i\rangle|0\rangle = \frac{1}{\sqrt{2^{N}}}\sum_{i=0}^{2^{N}-1}|i\rangle|f(i)\rangle$$
 Exponentially speedup!

This is what makes famous quantum algorithms, such as Shor's algorithm for factoring, or Grover's algorithm for searching

RSA public-key cryptosystem

Produce a large integer N

m1×m2=N, (with m1 and m2 primes)

- N is made public available and is used as a key (x) to encrypt data
- m1 and m2 are the secret keys (k) enable one to decrypt the data

$$C = E_{x}(P)$$

P = D_k(C)= D_k(E_x(P))

X: Public Key; K: Private Key

P: Plain Text; E: Encryption; C: Ciphertext; D: Decryption

Riverst, Shamir and Adleman, MIT/LCS/TR-212, Jan. 1979

- To crack a code, a code breaker needs to factorize N
- The security of RSA based on the ease with which N can be calculated from m1 and m2, and the difficulty of calculating m1 and m2 from N

- Problem: given a number, what are its prime factors ?
 e. g. a 129-digit odd number which is the product of two large primes, 11438162575788886766923577997614661201021829672124236256256184293570
 693524573389783059712363958705058989075147599290026879543541
 =3490529510847650949147849619903898133417764638493387843990820577
 x 32769132993266709549961988190834461413177642967992942539798288533
- Best factorizing algorithm requires sources that grow exponentially in the size of the number: $\exp\left(O\left(n^{1/3}\log^{2/3}n\right)\right)$, with n the length of N

Shor's Algorithm

Algorithms for quantum computation: discrete logarithms and factorizing

- E.g. factor a 300-digit number with
- Classical THz computer: 10²⁴ steps ⇒ 150,000 years
- Quantum THz computer: 10¹⁰ steps => 1 second!



Peter Shor

Foundations of Computer Science, 1994 Proceedings. 35th Annual Symposium

Code-breaking can be done in minutes, not in millennia
 Public key encryption, based on factoring, will be vulnerable!



Deutsch's problem: two types of functions f

Considering input n bits,

- Constant f: for all 2ⁿ inputs, f=0 or f=1
- Balanced f: for 2^{n-1} inputs, f=0, for another 2^{n-1} inputs, f=1

Question: given a function f, whether is it constant or balanced?

Classical deterministic algorithm: at most 2ⁿ⁻¹+1 inquiries

All outputs are the same => constant

At least 1 output is different from others => balanced

Deutsch-Jozsa Algorithm

The simplest example: (x=0 or 1)

 Constant:
 Balanced:

 f(0)=1
 f(0)=0

 f(1)=1
 f(1)=1

- Classical algorithm needs 2 inquiries
- Deutsch-Jozsa quantum algorithm: Assume f was mapped into a quantum oracle U satisfing

$$U|x\rangle|y\rangle \to |x\rangle|y \oplus f(x)\rangle \quad e. \ g., \ U_C = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \quad U_B = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

Deutsch-Jozsa Algorithm

- Prepare two qubits input state $|\psi_i\rangle = |0\rangle|1\rangle$
- Perform Hadamard operation $|\psi_i\rangle \rightarrow \frac{1}{2}(|0\rangle + |1\rangle)(|0\rangle |1\rangle)$
- Perform U

$$U|x\rangle\left(\frac{|0\rangle-|1\rangle}{\sqrt{2}}\right) = |x\rangle\left(\frac{|f(x)\rangle-|1-f(x)\rangle}{\sqrt{2}}\right) = (-1)^{f(x)}|x\rangle\left(\frac{|0\rangle-|1\rangle}{\sqrt{2}}\right) \implies$$

Output state:

$$|\psi_o\rangle = \begin{cases} \pm \left(\frac{|0\rangle + |1\rangle}{\sqrt{2}}\right) \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right), & if f(0) = f(1) \\ \pm \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right) \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right), & if f(0) \neq f(1) \end{cases}$$

• Measure the first qubit on $\{+/-\}$ basis: $|+\rangle \Rightarrow$ constant $f, |-\rangle \Rightarrow$ balanced f

Quantum algorithm only needs one inquiry

Consider a more general case with n-bit inputs x: x=0, 1, 2, ..., 2ⁿ-1

- Prepare n+1 qubits input state $|\psi_i\rangle = |0\rangle^{\otimes n}|1\rangle$ Perform Hadamard operation on all qubits $|\psi_i\rangle \rightarrow \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^{n-1}} |x\rangle|-\rangle$

The binary representation of x corresponds to values of each qubits, e. g.,

 $(|0\rangle + |1\rangle) \otimes (|0\rangle + |1\rangle) = |00\rangle + |01\rangle + |10\rangle + |11\rangle$

 $= |x = 0\rangle + |x = 1\rangle + |x = 2\rangle + |x = 3\rangle$

• Perform
$$U \rightarrow \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} (-1)^{f(x)} |x\rangle |-\rangle$$

Deutsch-Jozsa Algorithm

Measure the first n-qubit on {+/-} basis: → if and only if the output is |+>^{⊗n}, f is constant
 Only needs one inquiry!



Deutsch and Jozsa, Proc. Royal Society London A 439, 553 (1992)

Deutsch problem is not a practically important problem, but Deutsch-Jozsa algorithm firstly demonstrated the superiority of quantum computation!



Lov Grover

How quickly can you find a needle in a haystack? The simplest example:

Which one is equal to -1 in a database?

Serial	0	1	2	3	4	5	
Value	1	1	-1	1	1	1	

Classically search

- Sequentially try all N possibilities
- Average search takes N/2 steps

➤Quantum search

- Simultaneously try all possibilities
- Refining process reveals answer
- Average search takes N^{1/2} steps

> A databased is encoded with a N×N diagonal matrix R (rotate phase)

$$R = \begin{bmatrix} 1 & 0 & 0 & & 0 \\ 0 & 1 & 0 & \cdots & 0 \\ 0 & 0 & -1 & & 0 \\ \vdots & & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & 1 \end{bmatrix} \qquad R|i\rangle = \begin{cases} -|i\rangle, i = x \\ |i\rangle, \text{ otherwise} \end{cases}$$
 The task is to find x

> Take a *m*-qubit register $(2^m=N)$, and prepare the registers in an equal superposition state of all the states

$$|\varphi\rangle = \frac{1}{\sqrt{N}} \sum_{i=0}^{N} |i\rangle$$

Perform rotate phase matrix R on the register



Then perform diffusion operator D



> Measure the register to get the specific state $|x\rangle$ with nearly unity probability

Formulas



Formulas • After n iteration: $|\varphi_n\rangle = (DR)^n |\varphi\rangle \approx \sin\left(\frac{2n}{\sqrt{N}}\right) |x\rangle + \frac{\cos\left(\frac{2n}{\sqrt{N}}\right)}{\sqrt{N}} \sum_{i=1}^{N-1} |i\rangle$ $(i \neq x)$ The probability to collapse into the x
 Choose iteration steps $P = \left| \sin\left(\frac{2n}{\sqrt{N}}\right) \right|^2$ $n = \left| \frac{\pi}{4} \sqrt{N} \right|$ $1 - P \le \cos^2\left(\frac{\pi}{2} - \frac{1}{\sqrt{N}}\right) \xrightarrow{N \to \infty} 0$ The probability of failure:

Grover, PRL 79, 325 (1997)

Quantum Metrology

Super-resolution with multi-particle entanglement



Part 4: Quantum Repeaters and Quantum Error Correction



Unavoidable interaction with environment and decoherence will happen

 $|0\rangle|E\rangle \xrightarrow{U(t)} |0\rangle|E_0(t)\rangle \qquad |1\rangle|E\rangle \xrightarrow{U(t)} |1\rangle|E_1(t)\rangle$

- $|0\rangle$, $|1\rangle$ represents the qubit state and $|E\rangle$ represents the environment initial state, U(t) is the joint unitary time evolution operator
- For arbitrary qubit state:

$$(\alpha_{0}|0\rangle + \alpha_{1}|1\rangle)|E\rangle \xrightarrow{U(t)} \alpha_{0}|0\rangle|E_{0}(t)\rangle + \alpha_{1}|1\rangle|E_{1}(t)\rangle \qquad \rho_{q}(t) = Tr_{E}\rho_{q+E} = \begin{bmatrix} |\alpha_{0}|^{2} & \alpha_{0}\alpha_{1}^{*}\langle E_{1}|E_{0}\rangle \\ \alpha_{1}\alpha_{0}^{*}\langle E_{0}|E_{1}\rangle & |\alpha_{1}|^{2} \end{bmatrix}$$

The off-diagonal element of the qubit density matrix will drop down with the rate $\langle E_0(t)|E_1(t)\rangle=e^{-\Gamma t}$

The maximally entangled state will be in some mixed state with a certain entanglement fidelity due to the process × Photon loss increases exponentially with channel length: A ∝ e^{-ΓL} (e. g., in commercial fiber Γ = 0.2dB/km)
 × For 1000 km commercial fiber, even with a perfect 10 GHz single-photon source and ideal detectors, only 0.3 photon can be transmitted on average per century!

Solutions in Quantum Communication

Quantum repeater



Solution to photon loss:
 Entanglement swapping

Solution to decoherence:
 Entanglement purification

Briegel et al., PRL 81, 5932 (1998)

Entanglement Swapping

Entangling the remote particles which never interacted!



$$\begin{split} \left|\Psi\right\rangle_{1234} &= \left|\Phi^{+}\right\rangle_{14} \otimes \left|\Phi^{+}\right\rangle_{23} \\ &= \left|\Phi^{+}\right\rangle_{12} \otimes \left|\Phi^{+}\right\rangle_{34} + \left|\Phi^{-}\right\rangle_{12} \otimes \left|\Phi^{-}\right\rangle_{34} + \left|\Psi^{+}\right\rangle_{12} \otimes \left|\Psi^{+}\right\rangle_{34} + \left|\Psi^{-}\right\rangle_{12} \otimes \left|\Psi^{-}\right\rangle_{34} \end{split}$$

Zukowski et al., PRL 71, 4287 (1993)

Entanglement Swapping



Without entanglement swapping, the total cost in multi-stage is ~1/P^{2N}
 With entanglement swapping, the total cost is ~1/P²
 (assume that the emission probability of ERP sources is unity)

Entanglement Purification



Fidelity: $F = \langle \Psi^- | M | \Psi^- \rangle$

Goal: to extract from a large ensemble of low-fidelity M a small subensemble with sufficiently high fidelity

Scheme for Entanglement Purification

Random bilateral Pauli rotation on each photon in the states M = change arbitrary mixed state into Werner state: [Werner, PRA 40, 4277 (1989)]

$$W_F = F|\Psi^-\rangle\langle\Psi^-| + \frac{1-F}{3}|\Psi^+\rangle\langle\Psi^+| + \frac{1-F}{3}|\Phi^+\rangle\langle\Phi^+| + \frac{1-F}{3}|\Phi^-\rangle\langle\Phi^-| + \frac{1-F}{3}|\Psi^-\rangle\langle\Phi^-| + \frac{1-F}{3}|\Psi^-| + \frac{1-F}{3}|\Psi^$$

- For two same pairs of Werner states, we consider them as source pair and target pair respectively
- > A unilateral σ_Y is performed on each of the two pairs: $|\Psi^{\pm}\rangle \leftrightarrow |\Phi^{\mp}\rangle$



i. e., states with a large component (F > 1/2) of $|\Phi^+\rangle$, and equal components of the other three Bell states

Scheme for Entanglement Purification

> Perform CNOT operation on source and target pairs:



Measure target pair in {H/V} basis, keep the unmeasured source pair when measuring results are same

Drobobility	Bef	fore	After	
Probability	Source	Target	Source	Target
F ²	$ \Phi^+ angle$	$ \Phi^+ angle$	$ \Phi^+ angle$	$ \Phi^+ angle$
F(1-F)/3	$ \Phi^{-} angle$	$ \Phi^+ angle$	$ \Phi^{-} angle$	$ \Phi^+ angle$
F(1-F)/3	$ \Psi^+ angle$	$ \Phi^+ angle$	$ \Psi^+ angle$	$ \Psi^+ angle$
F(1-F)/3	$ \Psi^{-} angle$	$ \Phi^+ angle$	$ \Psi^{-} angle$	$ \Psi^+ angle$
$(1-F)^2/9$	$ \Psi^+ angle$	$ \Psi^+ angle$	$ \Psi^+ angle$	$ \Phi^+ angle$
$(1-F)^2/9$	$ \Psi^{-} angle$	$ \Psi^+ angle$	$ \Psi^{-} angle$	$ \Phi^+ angle$
F(1-F)/3	$ \Phi^+ angle$	$ \Psi^+ angle$	$ \Phi^+ angle$	$ \Psi^+ angle$
$(1-F)^2/9$	$ \Phi^{-} angle$	$ \Psi^+ angle$	$ \Phi^{-} angle$	$ \Psi^+ angle$
F(1-F)/3	$ \Phi^+ angle$	$ \Phi^{-} angle$	$ \Phi^{-} angle$	$ \Phi^{-} angle$
$(1-F)^2/9$	$ \Phi^{-} angle$	$ \Phi^{-} angle$	$ \Phi^+ angle$	$ \Phi^{-} angle$
$(1-F)^2/9$	$ \Psi^+ angle$	$ \Phi^{-} angle$	$ \Psi^{-} angle$	$ \Psi^{-} angle$
$(1-F)^2/9$	$ \Psi^{-} angle$	$ \Phi^{-} angle$	$ \Psi^+ angle$	$ \Psi^{-} angle$
$(1-F)^2/9$	$ \Psi^+ angle$	$ \Psi^{-} angle$	$ \Psi^{-} angle$	$ \Phi^{-} angle$
$(1-F)^2/9$	$ \Psi^{-} angle$	$ \Psi^{-} angle$	$ \Psi^+ angle$	$ \Phi^{-}\rangle$
F(1-F)/3	$ \Phi^+ angle$	$ \Psi^{-} angle$	$ \Phi^{-} angle$	$ \Psi^{-} angle$
$(1-F)^2/9$	$ \Phi^{-}\rangle$	$ \Psi^{-}\rangle$	$ \Phi^+ angle$	$ \Psi^{-}\rangle$

Scheme for Entanglement Purification

> After that, the component of $|\Phi^+\rangle\langle\Phi^+|$ of the target pair will be

$$F' = \frac{F^2 + \frac{1}{9}(1 - F)^2}{F^2 + \frac{2}{3}F(1 - F) + \frac{5}{9}(1 - F)^2} > F, \quad \text{when } F > \frac{1}{2}$$

- > Equivalently, the fidelity $\langle \Psi^- | M | \Psi^- \rangle$ is equal to F'
- Via several this kind processes, we can purify a general mixed state into a highly entangled state



Bennett et al., PRL 76, 722 (1996)

Solutions in Quantum Computation

Quantum error correction

Analogy between classical error correction:

Goal: store an unknown single bit for a time t

Errors:

- In a time interval au one error occurs with the probability $P_{ au}$
- Only one type of error: bit flips $0 \rightarrow 1, 1 \rightarrow 0$
- Suppose errors cause each physical bit to be flipped independently

Classical Error Correction Code

> Correct the errors by using a "redundant coding", e.g. :

Physical bits $000 \rightarrow \text{Logical bit } 0$ Physical bits $111 \rightarrow \text{Logical bit } 1$





Network for decoding

After the errors occur

- Probability of no errors: $(1 P_{\tau})^3$
- Probability of error in one bit: $3P_{\tau}(1-P_{\tau})^2$
- Probability of error in two bits: $3P_{\tau}^2(1-P_{\tau})$
- Probability of error in three bits: P_{τ}^3

Correction for a long time

To keep the state for a very long time t:

Correct errors as frequently as possible

- Consider $P_{\tau} = c\tau$ for time τ sufficiently short
- Divide t in N intervals of duration $\tau = t/N$
- After the time t:

$$P_t^c = \left[1 - 3\left(\frac{ct}{N}\right)^2 + 2\left(\frac{ct}{N}\right)^3\right] \to 1, when N \gg 3(ct)^3$$

Zeno effect!

Classical Error Correcting Codes

Decode the logical bits by taking the majority answer of the three bits and correct the encoded bits

$000 \rightarrow 000$	$111 \rightarrow 111$
$001 \rightarrow 000$	<mark>0</mark> 11 → 111
$010 \rightarrow 000$	1 <mark>0</mark> 1 → 111
$100 \rightarrow 000$	11 <mark>0</mark> → 111

The correct state with a probability $P_{\tau}^{c} = (1 - P_{\tau})^{3} + 3P_{\tau}(1 - P_{\tau})^{2} = 1 - 3P_{\tau}^{2} + 2P_{\tau}^{3}$ Measurement of error destroys superpositions
 No-cloning theorem prevents repetition
 Multiple types of errors

- Bit flip (σ_X) : $\sigma_X(\alpha|0\rangle + \beta|1\rangle) \rightarrow \alpha|1\rangle + \beta|0\rangle$
- Phase flip (σ_z) : $\sigma_Z(\alpha|0\rangle + \beta|1\rangle) \rightarrow \alpha|0\rangle \beta|1\rangle$
- Mixed (σ_Y) : $-i\sigma_Y(\alpha|0\rangle + \beta|1\rangle) \rightarrow \alpha|1\rangle \beta|0\rangle$

A 3-bit quantum error correction scheme uses an encoder and a decoder circuit



Any correction must be done without looking at the output

3-qubit Error Correction

Similar to classical error correction $|0\rangle \rightarrow |000\rangle$ $|1\rangle \rightarrow |111\rangle$

Superposition $\alpha |0\rangle + \beta |1\rangle \rightarrow \alpha |000\rangle + \beta |111\rangle$



Quantum circuit for encoding

Decoder looks just like the encoder
3-qubit Error Correction

All the possible error conditions

Error	Decoded	Correction	
$\alpha 000\rangle+\beta 111\rangle$	$\alpha 000\rangle + \beta 100\rangle = (\alpha 0\rangle + \beta 1\rangle)$	00>	
$\alpha 100\rangle + \beta 011\rangle$	$\alpha 111\rangle + \beta 011\rangle = (\alpha 1\rangle + \beta 0\rangle)$	11>	Flip the top qubit
$\alpha 010\rangle + \beta 101\rangle$	$\alpha 010\rangle + \beta 110\rangle = (\alpha 0\rangle + \beta 1\rangle)$	10>	
$\alpha 001\rangle+\beta 110\rangle$	$\alpha 001\rangle + \beta 101\rangle = (\alpha 0\rangle + \beta 1\rangle)$	01>	

After decoding, the states of syndrome qubits are orthogonal enabling to distinguish which qubit is the error occurred on

Phase-flip Error Correction

Phase flip: $\alpha |0\rangle + \beta |1\rangle \rightarrow \alpha |0\rangle - \beta |1\rangle$ Represent with $\{+/-\}$ basis $|\pm\rangle = \frac{1}{\sqrt{2}}(|0\rangle \pm |1\rangle)$

 $(\alpha + \beta)|+\rangle + (\alpha - \beta)|-\rangle \rightarrow (\alpha + \beta)|-\rangle + (\alpha - \beta)|+\rangle$ Bit flip error in {+/-} basis

Similar to bit flip error correction, a logical qubit is encoded with



Concatenated code

▶ Bit flip error correction: $|0\rangle \rightarrow |000\rangle |1\rangle \rightarrow |111\rangle$

Consider $|000\rangle$ as $|0'\rangle$, $|111\rangle$ as $|1'\rangle$

> Phase flip error correction: $|0'\rangle \rightarrow |+'+'+'\rangle \quad |1'\rangle \rightarrow |-'-'-'\rangle$

Corrects both bit flip and phase flip errors!

$$|0\rangle \rightarrow \left(\frac{|000\rangle + |111\rangle}{\sqrt{2}}\right)^{\otimes 3} \qquad |1\rangle \rightarrow \left(\frac{|000\rangle - |111\rangle}{\sqrt{2}}\right)^{\otimes 3}$$

Shor, PRA 52, 2493 (1995)

Shor's 9 Qubits Error Correcting Code



- General single-qubit error: $E = c_0 I + c_1 \sigma_Z + c_1 \sigma_X + c_3 \sigma_X \sigma_Z$
- Each term will be represented with orthogonal states of syndrome bits
 If a code can correct both bit flip and phase flip errors, it can correct arbitrary single-qubit error!

More Efficient Code

The Steane (CSS) code

 $|0\rangle \rightarrow |0_{L}\rangle \equiv \frac{1}{\sqrt{8}} [|000000\rangle + |101010\rangle + |0110011\rangle + |100110\rangle + |100110\rangle + |100110\rangle + |100110\rangle + |101001\rangle]$ $|1\rangle \rightarrow |1_{L}\rangle \equiv \frac{1}{\sqrt{8}} [|111111\rangle + |0101010\rangle + |1001100\rangle + |0011001\rangle + |100011\rangle + |001010\rangle]$

Calderbank and Shor, PRA 54, 1098 (1996) Steane, Proc. R. Soc. Lodon A 452, 2551 (1996)

Argument

Encode 1 logical qubit using n physical qubits
 n-1 syndrome bits => them can represent 2ⁿ⁻¹ states at most
 n qubits => 3n possible errors and also the case of no errors



$$|0\rangle \xrightarrow{encode} \frac{1}{\sqrt{8}} \left(|00000\rangle - |01111\rangle - |10011\rangle + |11100\rangle \right)$$
$$+ |00110\rangle + |01001\rangle + |10101\rangle + |11010\rangle$$

$$|1\rangle \xrightarrow{encode} \frac{1}{\sqrt{8}} (|11111\rangle - |10000\rangle + |01100\rangle - |00011\rangle + |11001\rangle + |10010\rangle - |01010\rangle - |00101\rangle)$$

Laflamme et al., PRL 77, 198 (1996)

5 Qubits Error Correcting Code

$$\begin{aligned} \alpha |0\rangle + \beta |1\rangle &\xrightarrow{encode} \rightarrow \frac{\alpha}{\sqrt{8}} \left(|0_1 0_2 0_3 0_4 0_5\rangle - |0_1 1_2 1_3 1_4 1_5\rangle - |1_1 0_2 0_3 1_4 1_5\rangle + |1_1 1_2 1_3 0_4 0_5\rangle \right) \\ &+ |0_1 0_2 1_3 1_4 0_5\rangle + |0_1 1_2 0_3 0_4 1_5\rangle + |1_1 0_2 1_3 0_4 1_5\rangle + |1_1 1_2 0_3 1_4 0_5\rangle \right) \\ &+ \frac{\beta}{\sqrt{8}} \left(|1_1 1_2 1_3 1_4 1_5\rangle - |1_1 0_2 0_3 0_4 0_5\rangle + |0_1 1_2 1_3 0_4 0_5\rangle - |0_1 0_2 0_3 1_4 1_5\rangle \right) \\ &+ |1_1 1_2 0_3 0_4 1_5\rangle + |1_1 0_2 1_3 1_4 0_5\rangle - |0_1 1_2 0_3 1_4 0_5\rangle - |0_1 0_2 1_3 0_4 1_5\rangle \right) \end{aligned}$$



Rules of Shifting Phase and Flip Bit



Signal after Hadamards



Step-by-step Analysis of Encoding Circuit



Step-by-step Analysis of Encoding Circuit



Step-by-step Analysis of Encoding Circuit



- Assuming at most 1 qubit error and the error is just as likely to affect any qubit
- The decoding circuit is the encoding circuit in reverse:



Assume encoded qubit damaged such that:

$$\begin{aligned} \frac{a}{\sqrt{8}} \left(\left| 0_{1} 0_{2} 0_{3} 0_{4} 0_{5} \right\rangle - \left| 0_{1} 1_{2} 1_{3} 1_{4} 1_{5} \right\rangle - \left| 1_{1} 0_{2} 0_{3} 1_{4} 1_{5} \right\rangle + \left| 1_{1} 1_{2} 1_{3} 0_{4} 0_{5} \right\rangle \\ + \left| 0_{1} 0_{2} 1_{3} 1_{4} 0_{5} \right\rangle + \left| 0_{1} 1_{2} 0_{3} 0_{4} 1_{5} \right\rangle + \left| 1_{1} 0_{2} 1_{3} 0_{4} 1_{5} \right\rangle + \left| 1_{1} 1_{2} 0_{3} 1_{4} 0_{5} \right\rangle \right) \\ + \frac{b}{\sqrt{8}} \left(\left| 1_{1} 1_{2} 1_{3} 1_{4} 1_{5} \right\rangle - \left| 1_{1} 0_{2} 0_{3} 0_{4} 0_{5} \right\rangle + \left| 0_{1} 1_{2} 1_{3} 0_{4} 0_{5} \right\rangle - \left| 0_{1} 0_{2} 0_{3} 1_{4} 1_{5} \right\rangle \\ + \left| 1_{1} 1_{2} 0_{3} 0_{4} 1_{5} \right\rangle + \left| 1_{1} 0_{2} 1_{3} 1_{4} 0_{5} \right\rangle - \left| 0_{1} 1_{2} 0_{3} 1_{4} 0_{5} \right\rangle - \left| 0_{1} 0_{2} 1_{3} 0_{4} 1_{5} \right\rangle \end{aligned}$$

, Phase and bit flip on 3rd qubit $(-i\sigma_p)$

$$\frac{\alpha}{\sqrt{8}} \left(\left| \begin{array}{c} \begin{array}{c} \begin{array}{c} \alpha \\ \end{array} \\ \hline \end{array} \\ \left| \begin{array}{c} 0 \\ \end{array} \\ \begin{array}{c} 0 \\ \end{array} \\ \begin{array}{c} 0 \\ \end{array} \\ \left| \begin{array}{c} 0 \\ \end{array} \\ \begin{array}{c} 0 \\ \end{array} \\ \begin{array}{c} 0 \\ \end{array} \\ \left| \begin{array}{c} 0 \\ \end{array} \\ \left| \begin{array}{c} 0 \\ \end{array} \\ \begin{array}{c} 0 \\ \end{array} \\ \begin{array}{c} 0 \\ \end{array} \\ \left| \begin{array}{c} 0 \\ \end{array} \\ \left| \begin{array}{c} 0 \\ \end{array} \\ \begin{array}{c} 0 \\ \end{array} \\ \left| \begin{array}{c} 0 \\ \end{array} \\ \\ \left| \begin{array}{c} 0 \end{array} \\ \\ \left| \begin{array}{c} 0 \\ \end{array} \\ \\ \left| \begin{array}{c} 0 \end{array} \\ \\ \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \\ \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \\ \\ \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \end{array} \right| \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \end{array} \right| \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \end{array} \right| \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \\ \\ \\ \end{array} \right| \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \\ \\ \end{array} \right| \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \\ \\ \end{array} \right| \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \\ \end{array} \right| \left| \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \end{array} \right| \left| \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \end{array} \\ \\ \\ \\ \\ \end{array} \right| \left| \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \end{array} \\ \\ \\ \\ \\ \end{array} \right| \left| \left| \left| \begin{array}{c} 0 \end{array} \\ \\ \\ \\ \\ \\ \end{array} \right| \left| \left| \left| \left| \left| \left| \left|$$

Continuation of Error Analysis in Decoder



Continuation of Error Analysis in Decoder



$$\frac{\alpha}{\sqrt{8}} \Big(- |0_1 0_2 1_3 0_4 1_5 \rangle + |0_1 1_2 1_3 1_4 1_5 \rangle + |1_1 0_2 1_3 1_4 1_5 \rangle - |1_1 1_2 1_3 0_4 1_5 \rangle \\ - |0_1 0_2 1_3 1_4 1_5 \rangle + |0_1 1_2 1_3 0_4 1_5 \rangle + |1_1 0_2 1_3 0_4 1_5 \rangle - |1_1 1_2 1_3 1_4 1_5 \rangle \Big)$$
Shifting phase on bit
$$+ \frac{\beta}{\sqrt{8}} \Big(+ |1_1 1_2 0_3 1_4 1_5 \rangle - |1_1 0_2 0_3 0_4 1_5 \rangle - |0_1 1_2 0_3 0_4 1_5 \rangle + |0_1 0_2 0_3 1_4 1_5 \rangle \\ + |1_1 1_2 0_3 0_4 1_5 \rangle - |1_1 0_2 0_3 1_4 1_5 \rangle - |0_1 1_2 0_3 1_4 1_5 \rangle + |0_1 0_2 0_3 0_4 1_5 \rangle \Big)$$



Re-express equation to prepare for Hadamard transform:

$$\begin{split} &\frac{\alpha}{\sqrt{8}} \left(-|0_{1}0_{2}1_{3}0_{4}1_{5}\rangle + |0_{1}1_{2}1_{3}1_{4}1_{5}\rangle + |1_{1}0_{2}1_{3}1_{4}1_{5}\rangle - |1_{1}1_{2}1_{3}0_{4}1_{5}\rangle \right) \\ &\quad -|0_{1}0_{2}1_{3}1_{4}1_{5}\rangle + |0_{1}1_{2}1_{3}0_{4}1_{5}\rangle + |1_{1}0_{2}1_{3}0_{4}1_{5}\rangle - |1_{1}1_{2}1_{3}1_{4}1_{5}\rangle \right) \\ &\quad + \frac{\beta}{\sqrt{8}} \left(+ |1_{1}1_{2}0_{3}1_{4}1_{5}\rangle - |1_{1}0_{2}0_{3}0_{4}1_{5}\rangle - |0_{1}1_{2}0_{3}0_{4}1_{5}\rangle + |0_{1}0_{2}0_{3}1_{4}1_{5}\rangle \right) \\ &\quad + |1_{1}1_{2}0_{3}0_{4}1_{5}\rangle - |1_{1}0_{2}0_{3}1_{4}1_{5}\rangle - |0_{1}1_{2}0_{3}1_{4}1_{5}\rangle + |0_{1}0_{2}0_{3}0_{4}1_{5}\rangle \right) \\ &= \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) \left(\frac{\alpha}{\sqrt{4}} \left(-|0101\rangle + |111\rangle - |0111\rangle + |1101\rangle \right) + \frac{\beta}{\sqrt{4}} \left(-|1001\rangle + |0011\rangle - |1011\rangle + |0001\rangle \right) \right) \\ &= \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) \left(-\alpha|1\rangle + \beta|0\rangle \left(\frac{|01\rangle + |11\rangle}{\sqrt{2}} \right) \\ &= \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) \left(-\alpha|1\rangle + \beta|0\rangle \left(\frac{|0\rangle + |1\rangle}{\sqrt{2}} \right) (|1\rangle) \Rightarrow \text{ Input to Hadamard} \end{split}$$

Continuation of Error Analysis in Decoder



- Qubits 1,2,4 and 5 are the syndrome bits which indicate the exact error that occurred and the current state of qubit 3:
- So apply a phase shift and a bit flip on qubit 3 to obtain the protected qubit $\alpha|0\rangle + \beta|1\rangle$

Syndromes Table after Decoding

Syndrome qubits: q1, q2, q4, q5 Output qubit: q3

Syndrome states	Error on					Output qubit	Syndrome states	Error on				Output qubit	
	q1	q2	q3	q4	q5			q1	q2	q3	q4	q5	
0000						$\alpha 0\rangle + \beta 1\rangle$	1000	Z					$-\alpha 0\rangle - \beta 1\rangle$
0001		Х				$\alpha 0\rangle - \beta 1\rangle$	1001				Y		$-\alpha 1\rangle - \beta 0\rangle$
0010				Ζ		$-\alpha 0\rangle - \beta 1\rangle$	1010			Z			$\alpha 0\rangle - \beta 1\rangle$
0011					Х	$-\alpha 0\rangle - \beta 1\rangle$	1011				Х		$-\alpha 1\rangle - \beta 0\rangle$
0100		Ζ				$-\alpha 0\rangle - \beta 1\rangle$	1100					Ζ	$\alpha 0\rangle - \beta 1\rangle$
0101		Y				$\alpha 0\rangle - \beta 1\rangle$	1101			Y			$-\alpha 1\rangle + \beta 0\rangle$
0110	Х					$-\alpha 1\rangle - \beta 0\rangle$	1110	Y					$-\alpha 1\rangle - \beta 0\rangle$
0111			Х			$-\alpha 1\rangle - \beta 0\rangle$	1111					Y	$-\alpha 0\rangle + \beta 1\rangle$



Recalling that these codes require $P_{\tau} \propto \tau \ll 1$ for time τ sufficiently short In realistic devices, τ cannot be infinite small

threshold for tolerable error rate

Highest threshold: 2.02×10^{-5} Extremely hard to achieve! Spedalieri *et al.*, Quantum Inf. Comput. 9, 666 (2009)

- A possible solution: Topological error correction! Raussendorf et al., Ann. Phys. 321, 2242 (2003)
 - Topological homology of 3D cluster state (encoding one logical qubit with 180 physical qubits)
 - Relax the error threshold rate from 10⁻⁵ to 10⁻²



Thanks for your attention!