#### Entropy Rain and the Solar Convective Conundrum: Dilution and Compression of Individual Convective Downflows Evan H. Anders, Daniel Lecoanet, Benjamin P. Brown





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#### **Heat Transfer of Stars**

#### > 1.5 solar masses



0.5 - 1.5 solar masses



< 0.5 solar masses





[Wikipedia]

#### The classic view of envelope convection



#### [Image courtesy of Daniel Lecoanet]

## Solar Convective Flows





Granules: ~1 Mm ~5 min

Supergranules: ~30 Mm ~1 day

Visible (Photosphere)

Call (Chromosphere)

## Flows in Simple Stratified Convection



[Anders & Brown 2017; PRF 2]

## The Solar Convective Conundrum: Horizontal Surface Flows



## The Solar Convective Conundrum: Helioseismology



## The Convective Conundrum cont'd

Red = surface observations Blue = simulations

Grey = helioseismology

[Hanasoge et al 2015]



#### Extra motivation in the asteroseismic age



Many observations which rely on stellar structure models.

→ Stellar structure models rely on convection parameterizations like MLT.

 $\rightarrow$  We need to understand stellar envelope convection better.

[Astro2020 white paper: Huber et al 2019 ]

#### Takeaways

1) There is an absence of large-scale convective power in observations of the Sun.

#### Some options:

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Entropy Rain Hypothesis: Cold downflows from the solar surface carry the solar luminosity to the bottom of the CZ, *not* warm upflows which would manifest as giant cells.

[Spruit 1997; Brandenburg 2016; and results of Käpylä et al. 2017]







## Thermals as a model for Entropy Rain

Starts as a blob of buoyant fluid

Evolves into a buoyant vortex ring

Entrains fluid from environment



[Suggested by Brandenburg 2016]

#### Thermals as a model for Entropy Rain



## 3D Visualization of an Evolved Thermal



## Flows in Simple Stratified Convection



[Anders & Brown 2017; PRF 2]

## Flows in Simple Stratified Convection



[Anders & Brown 2017; PRF 2]

## Thermals as a model for Entropy Rain

Fundamental Questions:

1.) Can entropy rain survive throughout the depth of the CZ?

2.) Can entropy rain carry the stellar luminosity?



[Suggested by Brandenburg 2016]

#### Thermal Entrainment in the Incompressible Limit

growth of thermal through transport of environmental fluid into the thermal

Depth:

 $z \propto t^{1/2}$ 

#### (implies deceleration)



[Lecoanet & Jeevanjee 2019]

#### Thermal Entrainment in the Incompressible Limit

Volume: V ∝ r<sup>3</sup> (measured) V ∝ z<sup>3</sup> (spheroidal)



[Lecoanet & Jeevanjee 2019]

#### Thermal Entrainment in the Incompressible Limit

#### Volume: $V \propto r^3 \propto z^3$ , so Radius: $z \propto r \propto t^{1/2}$

Even though thermals buoyantly accelerate, they entrain enough to stall themselves.



[Lecoanet & Jeevanjee 2019]

# Fundamental problem of thermals as a model for entropy rain:

# If thermals entrain & slow down, how can they cross the solar CZ?

## Stratification breaks symmetry

#### <u>Cold thermals:</u>

- Boussinesq-like entrainment
- stratification-enduced <u>compression</u>

#### <u>Hot thermals:</u>

- Boussinesq-like entrainment
- stratification-enduced <u>expansion</u>

## Stratification effects without buoyancy

 $\begin{array}{l} \underline{\text{Compression should go as:}}\\ \bullet \mathbf{r} \propto \boldsymbol{\rho}^{-1/2} \text{ for } \boldsymbol{horizontal} \text{ compression} \\ \bullet \mathbf{r} \propto \boldsymbol{\rho}^{-1/3} \text{ for } \boldsymbol{spherical} \text{ compression} \end{array}$ 

For  $\rho$ , the atmospheric density,

and **r**, the radius:



[Brandenburg 2016]

## **Buoyantly neutral Hill vortices**



## Non-buoyant hill vortex simulations were between expected limits (slope of -0.8).

[Brandenburg 2016]

#### Takeaways

1) There is an absence of large-scale convective power in observations of the Sun.

## 2) Entropy rain (the nonlocal transport of heat by downflows) may make giant cells unnecessary.

## Experiment: Evolve thermals in atmospheres with different stratifications

Fundamental control parameter:

S = (size of thermal) / (atmospheric scale height)

Limits:

 $S \rightarrow 0$  (Boussinesq)  $S \rightarrow O(1)$  (stratification matters)



# To first order, we see what we expect: blended compression & entrainment



#### Sims soon; First, some theory

## Goal: Model vortex ring stage of thermal's evolution

#### Important quantities:

- Buoyancy of the thermal, **B**
- Circulation of the vortex ring,
- Volume of the thermal, V
- Radius of the thermal, **r**
- Impulse of the thermal, I



## Theory of stratified vortex ring entrainment Fundamental Quantity: Impulse

$$\boldsymbol{I} = \frac{1}{2} \int_{\mathcal{V}} \boldsymbol{x} \times (\nabla \times (\rho \boldsymbol{u})) dV,$$

...which is similar to the momentum

$$I_z = (1 + k) \rho \mathcal{V} w_{\text{th}},$$

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#### and which is approximately (thin core vortex ring)

$$I_z \approx \pi \rho_0 r^2 \Gamma,$$

Assume the volume-integrated buoyancy force of the thermal is ~constant in time,

$$B \equiv \int_{\mathcal{V}} \rho_0 S_1 \, \frac{g}{c_P} \, dV.$$

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...and assume buoyancy is the only force acting on the thermal.

#### Assume circulation is constant (no baroclinic torques)

$$\frac{d\Gamma}{dt} = \oint_{\mathcal{C}} g \frac{S_1}{c_P} \hat{z} \cdot d\boldsymbol{x},$$



#### We can solve impulse for radius vs time,



→  $\mathbf{r} \propto \mathbf{\rho}^{-1/2}$  for purely horizontal compression, and →  $\mathbf{r} \propto \mathbf{t}^{1/2}$  for boussinesq regime.

#### Assume spheroidal volume

$$\mathcal{V}~pprox~mr^3$$

Substitute in the definition of velocity,

$$w_{\rm th} \equiv dz_{\rm th}/dt$$

#### Step 6: Solve for depth vs time

$$\frac{dz}{\rho(z)^{1/2}} = C \frac{dt}{(t+t_{\text{off}})^{1/2}}$$

(Simple ODE which can be solved for a prescribed density stratification)

## Experiment: two types of simulations

#### 1) 2D anelastic, cylindrically symmetric

- → Computationally cheap
- → Limit in which theory was derived

## **2) 3D fully compressible, cartesian boxes** $\rightarrow$ The "true" solution













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1) There is an absence of large-scale convective power in observations of the Sun.

- 2) Entropy rain (the nonlocal transport of heat by downflows) may make giant cells unnecessary
- 3) Buoyancy makes vortex rings expand & entrain, while stratification compresses them. Entropy rain feels both of these effects.
  4) We develop a theory that describes the size & depth of entropy rain vs. time and verify that theory with simulations. There are two regimes: "stalling" and "falling".

Stratification	
increases	

Simulation Output Parameterization						
$\overline{N_{ ho}}$		$B_{ m th}$	$\Gamma_{\rm th}$	т	$\eta$	k
2D /	Anelastic Simulations	(Cylindrical)				
0.1	Changes	-0.548	-2.17	8.05	1.04	0.732
0.5		-0.569	-2.12	8.34	0.977	0.715
1		-0.602	-2.05	8.65	0.915	0.703
2		-0.713	-1.89	9.23	0.842	0.682
3	stratification	-0.947	-1.73	9.81	0.807	0.654
4		-1.47	-1.59	10.2	0.794	0.642
5		-2.70	-1.49	10.7	0.781	0.609
6		-5.73	-1.43	10.8	0.787	0.616
3D Fully Compressible Simulations (Cartesian)						
0.1	Changes	-0.547	-2.17	8.98	1.06	0.636
0.5	with	-0.568	-2.12	8.79	0.978	0.678
1	incroscing	-0.601	-2.05	8.87	0.907	0.689
2	increasing	-0.711	-1.89	9.31	0.828	0.680
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Roughly constant as stratification increases

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Roughly constant as stratification increases



We calculated what **B** would be for a solar downflow....

S	imulation Outpu	t Par	We	e es	timate	e that
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Behavior of range of estimated thermals

Thermals travel this way (from surface to RZ)







A single thermal carries 0.01% of the solar luminosity at base of CZ.



Only a small fraction of the area of the CZ is required to carry the solar luminosity using thermals





## Takeaways

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- 2) Entropy rain (the nonlocal transport of heat by downflows) may make giant cells unnecessary
- 3) Buoyancy makes vortex rings expand & entrain, while stratification compresses them. Entropy rain feels both of these effects.
- 4) We develop a theory that describes the size & depth of entropy rain vs. time and verify that theory with simulations. There are two regimes: "stalling" and "falling"
- 5) The enthalpy fluxes carried by solar downflows is sufficient to carry the solar luminosity in the deep CZ

For more specifics, see the paper: [Anders, Lecoanet & Brown 2019; ApJ 884; arXiv: 1906.02342]

#### Preliminary Turbulent Results

- Entrainment behavior ~the same
- Acceleration & velocity behavior is quite different than the laminar cases
- Hunch: baroclinic torques

$$\frac{d\Gamma}{dt} = \oint_{\mathcal{C}} g \frac{S_1}{c_P} \hat{z} \cdot d\boldsymbol{x},$$

#### + viscous heating

$$\rho \frac{\partial S_1}{\partial t} \sim \rho \frac{\nu}{T} \nabla^2 |\boldsymbol{u}|^2; \boldsymbol{\mu} = \rho \nu$$

