

# The little hierarchy problem at the LHC

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• Experimental “problems” of the SM:

- Gravity
- Dark matter
- Baryon asymmetry

• Experimental “hints” of physics beyond the SM

- Neutrino masses
- Quantum number unification

• Theoretical puzzles of the SM:

- $\langle H \rangle \ll M_{\text{Pl}}$
- Family replication
- Small Yukawa couplings, pattern of masses and mixings
- Gauge group, no anomaly, charge quantization, quantum numbers

• Theoretical problems of the SM:

- Naturalness / unitarity problem
- Cosmological constant problem
- Strong CP problem

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• The unitarity/naturalness argument for new physics @ TeV:  
either a new strong scale  $Q_{\text{NP}} \approx \text{TeV}$  or  
 $G_a \leftrightarrow h$ , with  $\delta m_h^2 / m_h^2 \approx (Q_{\text{NP}} / \text{TeV})^2$

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- Many options; how do they confront the Little Hierarchy problem?

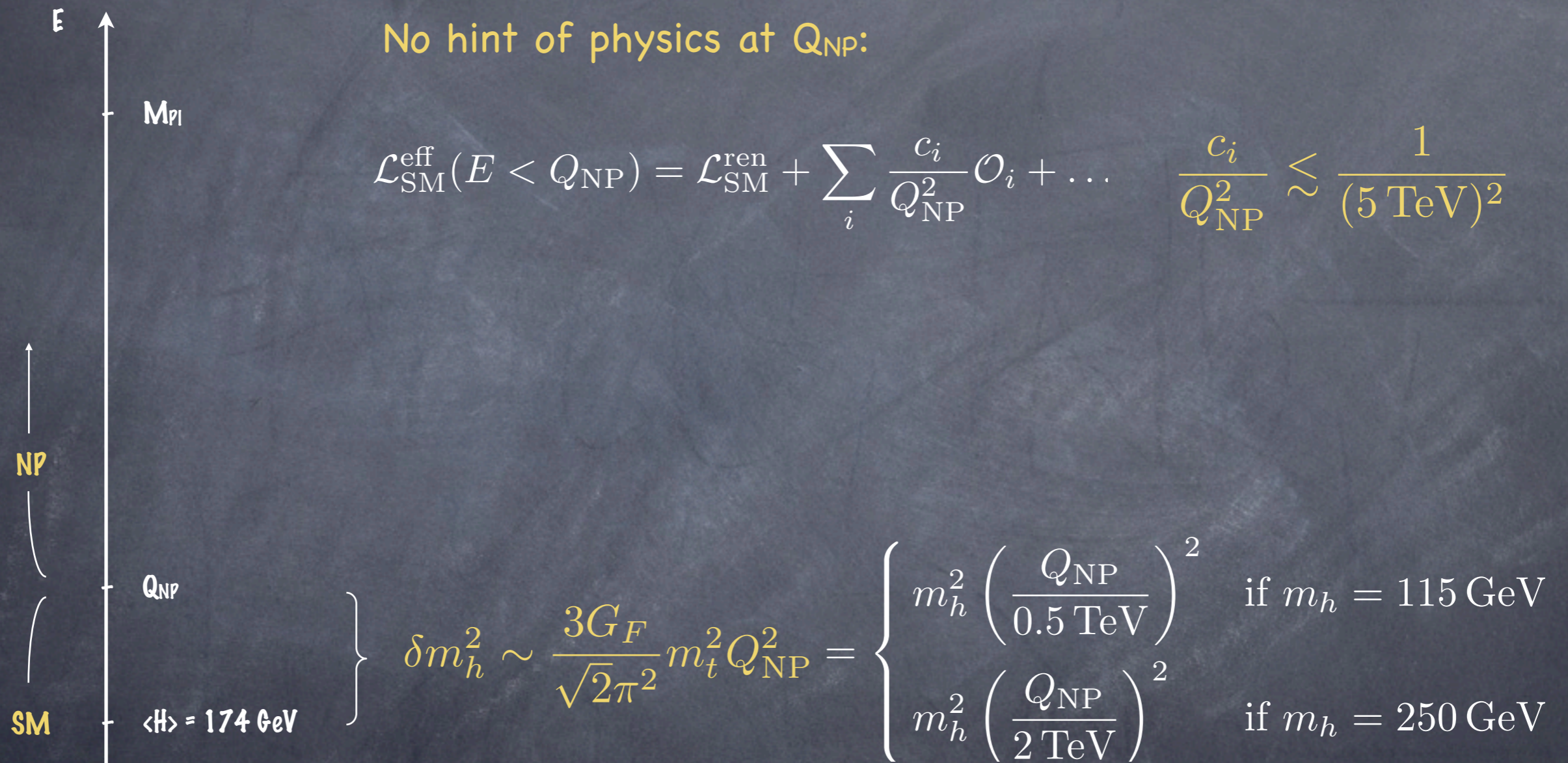
# The little residual hierarchy (e.g. with h)

The diagram illustrates the energy hierarchy. A vertical axis labeled  $E$  has tick marks for  $M_{Pl}$ ,  $Q_{NP}$ , and  $\langle H \rangle = 174 \text{ GeV}$ . The Standard Model (SM) scale is indicated by a bracket from the origin to  $\langle H \rangle$ . The New Physics (NP) scale is indicated by a bracket from  $\langle H \rangle$  to  $Q_{NP}$ . The Planck scale  $M_{Pl}$  is shown at the top.

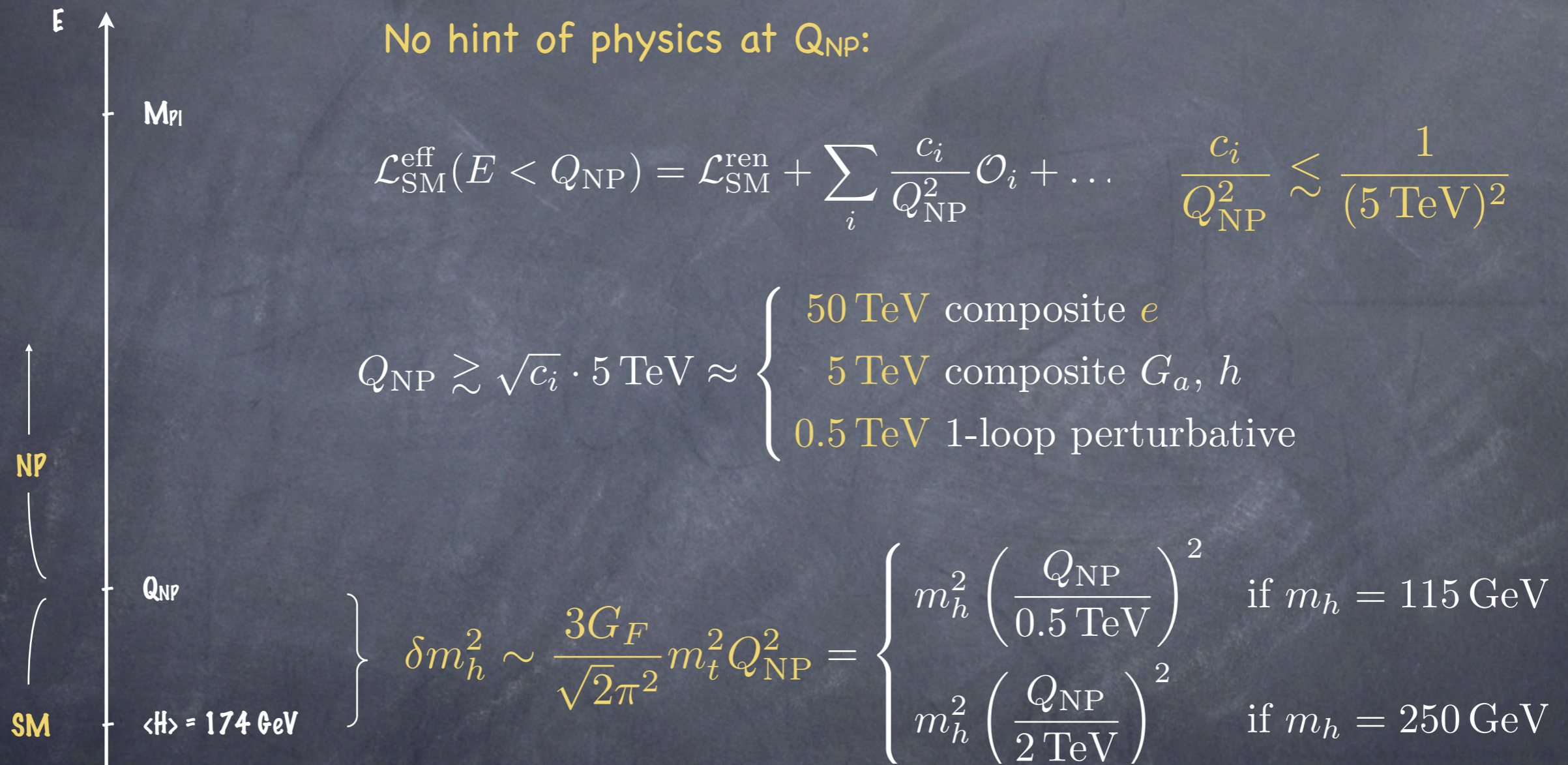
The formula for the shift in the Higgs mass squared is given by:

$$\delta m_h^2 \sim \frac{3G_F}{\sqrt{2}\pi^2} m_t^2 Q_{NP}^2 = \begin{cases} m_h^2 \left( \frac{Q_{NP}}{0.5 \text{ TeV}} \right)^2 & \text{if } m_h = 115 \text{ GeV} \\ m_h^2 \left( \frac{Q_{NP}}{2 \text{ TeV}} \right)^2 & \text{if } m_h = 250 \text{ GeV} \end{cases}$$

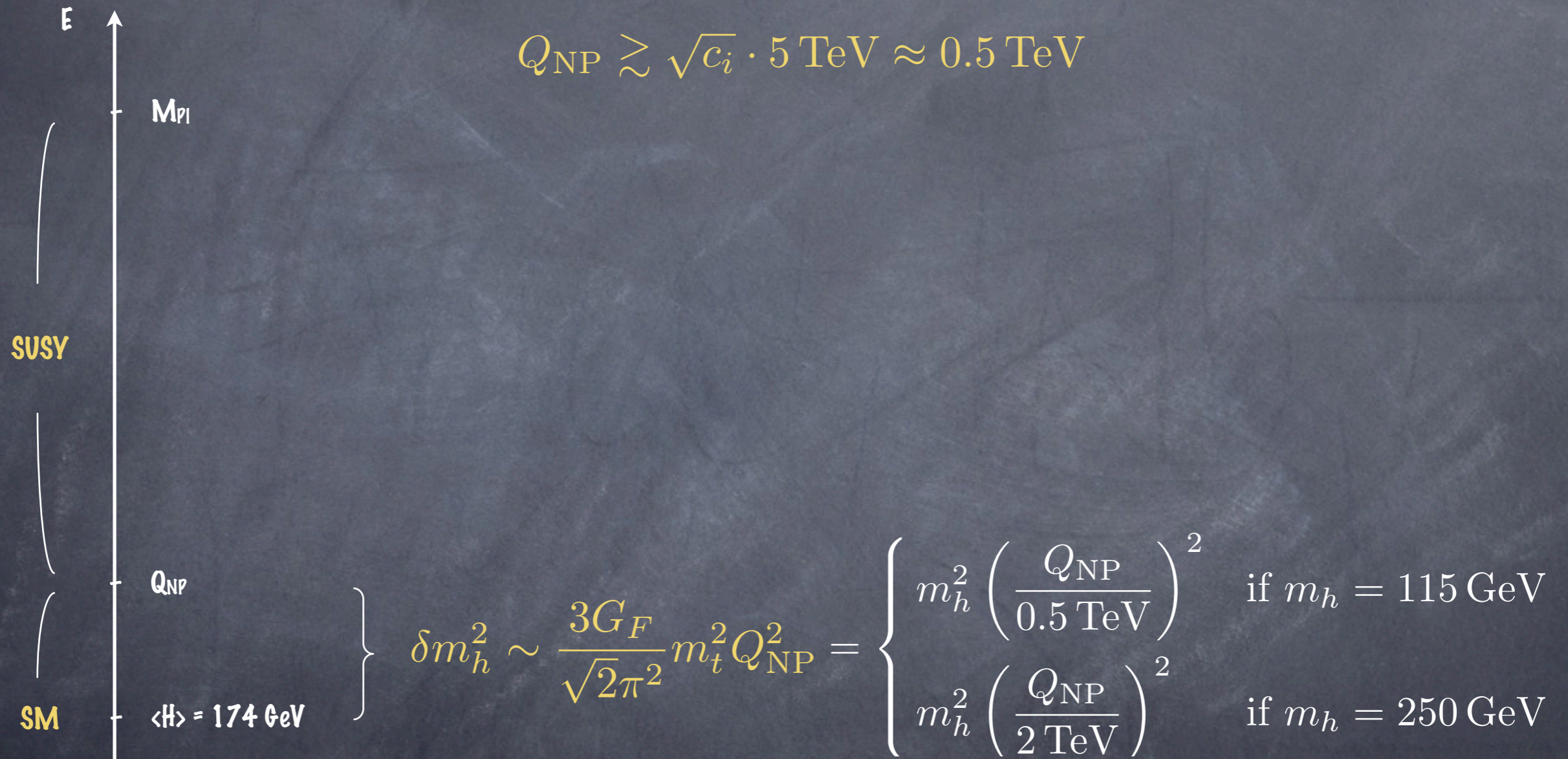
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# The little residual hierarchy (e.g. with $h$ )

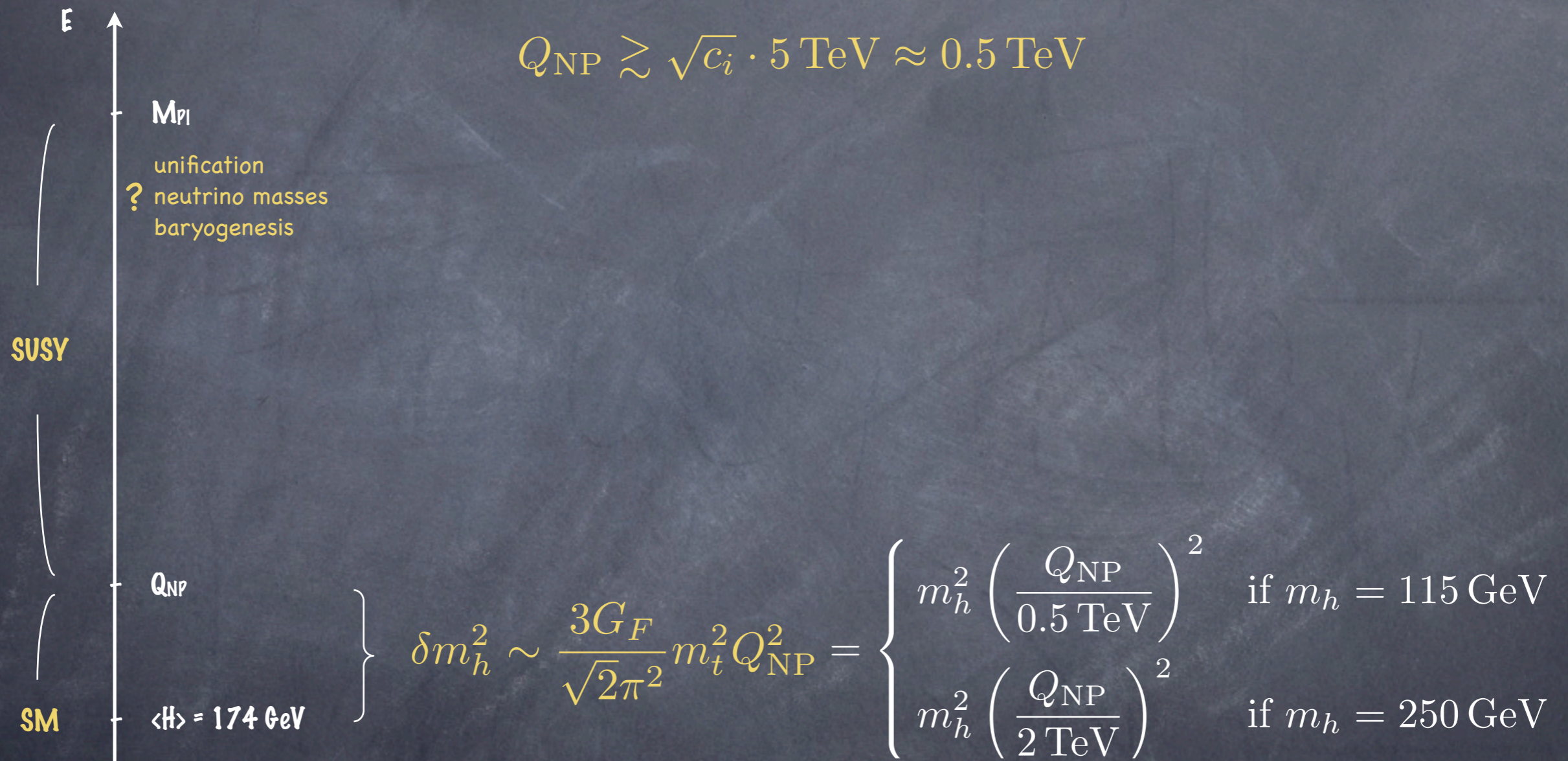


# MSSM

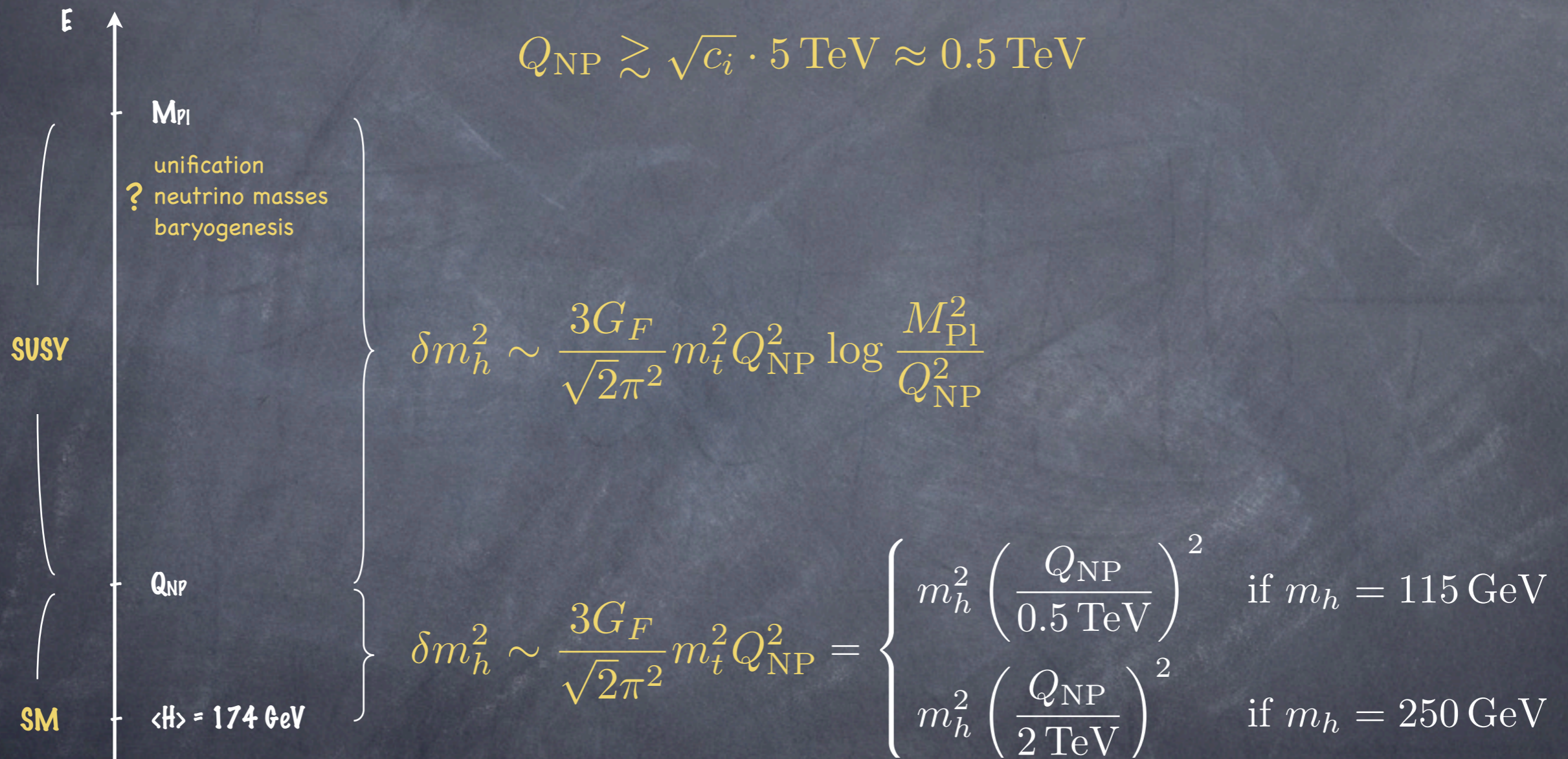


$$Q_{NP} \gtrsim \sqrt{c_i} \cdot 5 \text{ TeV} \approx 0.5 \text{ TeV}$$

# MSSM



# MSSM



# Fine-tuning in the MSSM

- $M_Z^2 = -2 \frac{m_{h_u}^2 \tan^2 \beta - m_{h_d}^2}{\tan^2 \beta - 1} - 2|\mu|^2 \approx -2m_{h_u}^2 - 2|\mu|^2 \quad (\text{large } \tan \beta)$   
 $\approx -2 (m_{h_u}^2(M_0) + |\mu|^2) + 2\delta m_{h_u}^2$
- $\delta m_{h_u}^2 \gg M_Z^2$  (large logs + color + bounds on gluinos and squarks):  
a moderate (up to %) fine-tuning is required to obtain  $M_Z = 91 \text{ GeV}$

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- Indirect bound on stop mass stronger (but direct one is also relevant)
 
$$(114 \text{ GeV})^2 < m_h^2 < M_Z^2 \cos^2 2\beta + \frac{3}{4\pi^2} h_t^2 m_t^2 \log \frac{\tilde{m}_t^2}{m_t^2} \Rightarrow \text{FT} \sim 50 \div 100$$

# A comment on scanning procedures

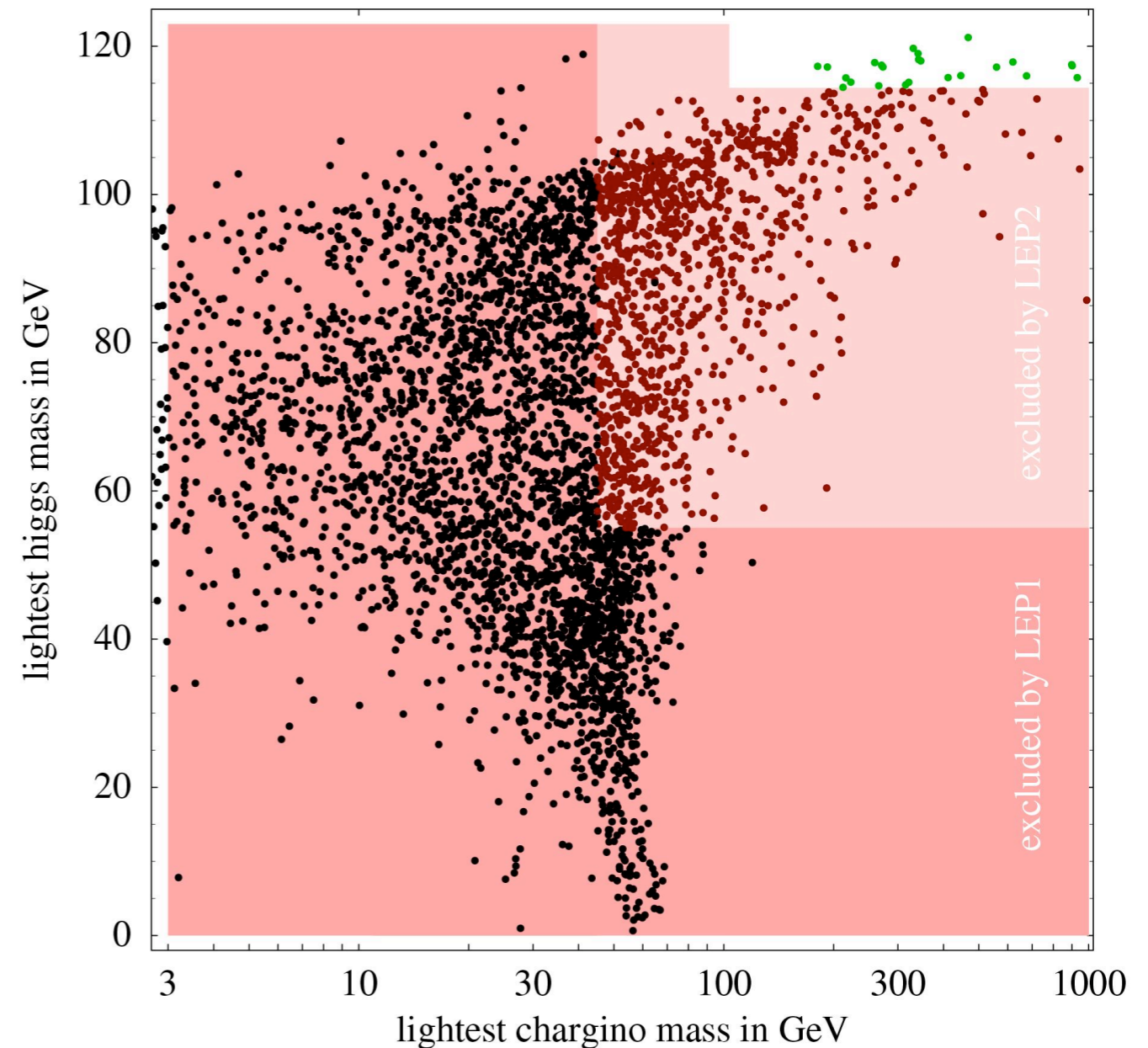
- Numerical procedures typically trade  $|\mu|$  for  $M_Z$ : do not “see” FT

$$M_Z^2 \approx -2 \left( m_{h_u}^2(M_0) + |\mu|^2 \right) + 2 \delta m_{h_u}^2$$

- The FT problem then may introduce a bias in numerical scans of the MSSM parameter space: the (necessary) cancellation is forced to take place between  $\mu^2$  and all the rest
- Example: LSP is rarely an Higgsino (work in progress)

# What is left?

- Quantitative measure of naturalness nicely taking into account and combining all the considerations above
  - Scan the relative sizes of SUSY parameters and the SM parameters in their ranges
  - Set the overall scale of SUSY parameters from  $\langle H \rangle = 174$  GeV
  - Calculate SUSY spectrum and compare with experiment
- Few  $O(1\%)$  of points satisfy all experimental constraints



[Giusti R Strumia]

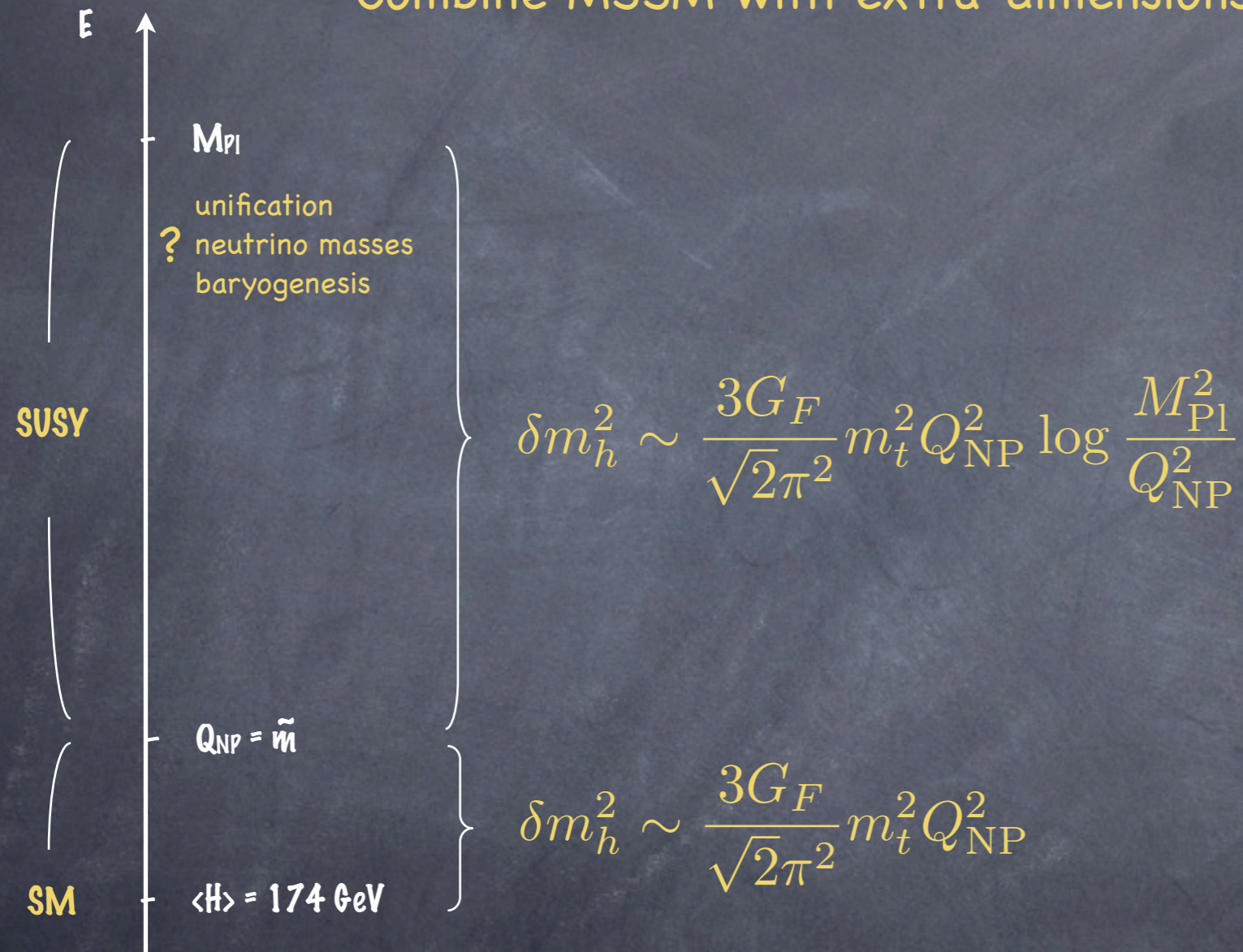
# Beyond MSSM: xMSSM

- Minimal extension:  $\lambda S H_u H_d$  (symmetries forbid  $\mu H_u H_d$ )
  - harmless (unification OK)
  - welcome ( $\mu = \lambda \langle S \rangle \approx \text{susy scale}$ )
- Spectrum:  $h H \rightarrow h_1 h_2 h_3, A \rightarrow a_1 a_2, N_1 \dots N_4 \rightarrow N_0 N_1 \dots N_4$
- Help with FT from  $(114 \text{ GeV})^2 < m_h^2 < M_Z^2 \cos^2 2\beta + \frac{3}{4\pi^2} h_t^2 m_t^2 \log \frac{\tilde{m}_t^2}{m_t^2}$ :
  - $\lambda_h = \frac{g^2 + g'^2}{4} \cos^2 2\beta + \frac{\lambda^2}{2} \sin^2 2\beta + \text{loops}$  gain limited by poles  
 $\lambda(10 \text{ TeV}) < 3$  (EWPTs) best,  $\lambda(M_{\text{GUT}}) < 3$  (unification) OK
  - $m_h^2 < (114 \text{ GeV})^2$  hidden Higgs:  $h \rightarrow aa \rightarrow 4X$  ( $m_a$  protected by PQ, R)
- Persistent FT from
  - direct bounds on SUSY partners
  - arranging the invisible decay [Shuster Toro hep-ph/0512189]

- Invisible Higgs decays:  $h \rightarrow aa \rightarrow 4X$
- 3leptons  $\rightarrow$  multileptons from additional steps in chargino/neutralino decays
  - $C_1 + N_2$  and then
  - $N_2 \rightarrow N_1 + 2l \rightarrow N_0 + 4l$  (if  $N_0$  is lightest and mainly singlino)
  - $C_1 \rightarrow N_0 + l + \nu$  (5l overall) or even  $C_1 \rightarrow N_1 + l + \nu \rightarrow N_0 + 3l + \nu$  (7l overall)
- Deviation from MSSM coupling relations:  $VVh = VhA = \sin^2(\alpha - \beta)$ ,  $VVH = VhA = \cos^2(\alpha - \beta)$  (optimistic)
- $Z'$  if  $\mu$  is protected by a gauge symmetry

# Other variations on the MSSM

Combine MSSM with extra-dimensions not far from TeV



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Combine MSSM with extra-dimensions not far from TeV

A vertical axis with an upward-pointing arrow at the top. On the left side of the axis, there are two labels: 'SUSY' at the top and 'SM' at the bottom, with a curved line connecting them. On the right side of the axis, there are three tick marks. The top tick mark is labeled  $M_f$ . The middle tick mark is labeled  $Q_{NP} = \tilde{m}$ . The bottom tick mark is labeled  $\langle H \rangle = 174 \text{ GeV}$ . A large curly brace on the right side of the axis spans from the middle tick mark to the bottom tick mark. To the right of this brace is the equation  $\delta m_h^2 \sim \frac{3G_F}{\sqrt{2}\pi^2} m_t^2 Q_{NP}^2$ .

SUSY

SM

$M_f$

$Q_{NP} = \tilde{m}$

$\langle H \rangle = 174 \text{ GeV}$

$\delta m_h^2 \sim \frac{3G_F}{\sqrt{2}\pi^2} m_t^2 Q_{NP}^2$

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Combine MSSM with extra-dimensions not far from TeV

Diagram illustrating the relationship between SUSY and SM scales and their effect on the Higgs mass:

Vertical axis (Energy Scale):

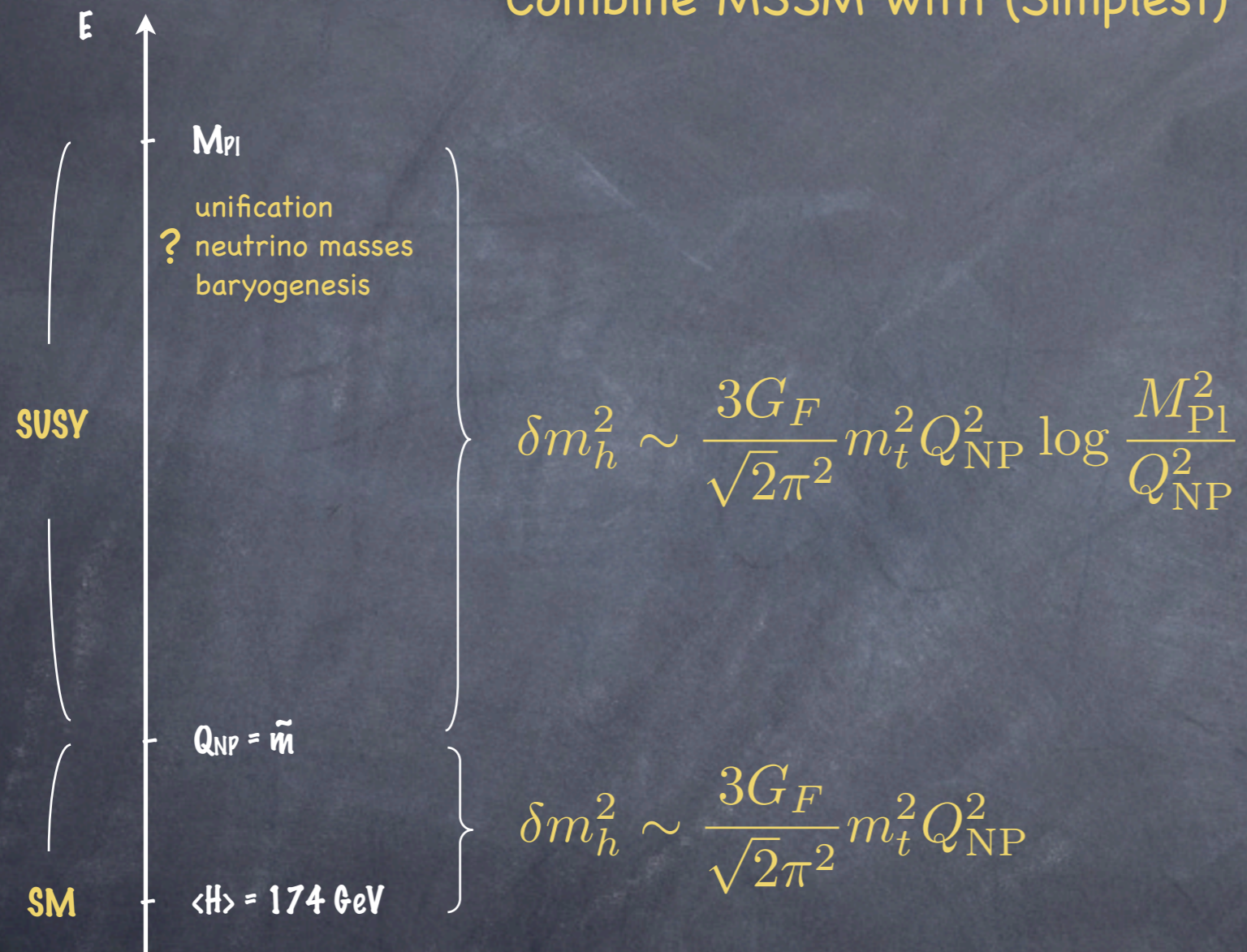
- SUSY** (Supersymmetry) scale:  $M_f$
- SM** (Standard Model) scale:  $\langle H \rangle = 174 \text{ GeV}$
- Intermediate scale:  $Q_{NP} = \tilde{m}$

Mass corrections  $\delta m_h^2$  are shown for two cases:

- For the SUSY scale:  $\delta m_h^2 \sim \frac{3G_F}{\sqrt{2}\pi^2} m_t^2 Q_{NP}^2 \log \frac{M_f^2}{Q_{NP}^2}$
- For the SM scale:  $\delta m_h^2 \sim \frac{3G_F}{\sqrt{2}\pi^2} m_t^2 Q_{NP}^2$

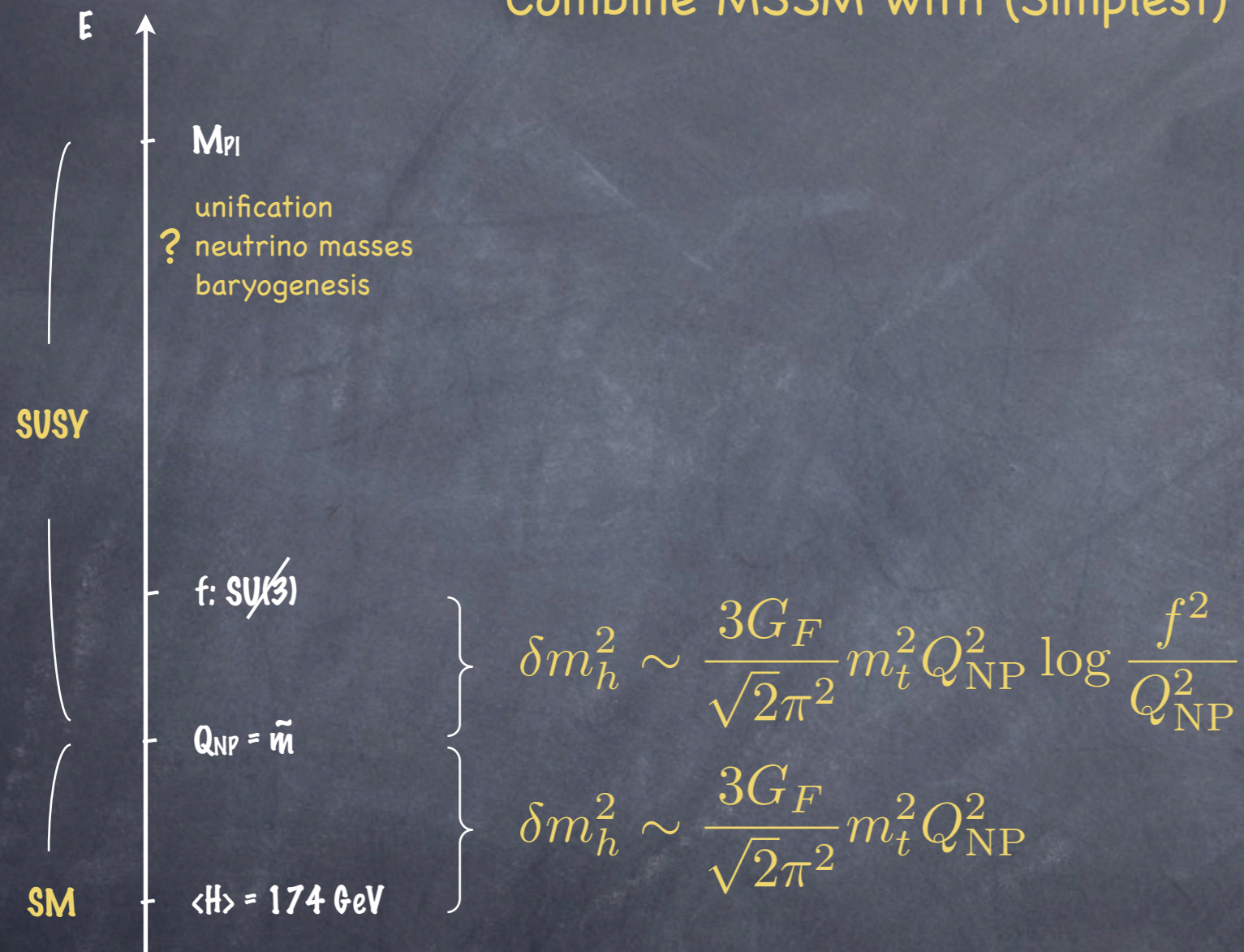
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Combine MSSM with (Simplest) Little Higgs



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# SSM with $Q_3 = (t_L \ b_L) = \text{gaugino}$

- $G = SU(5) \times G'_{SM}$  broken to the diagonal  $G_{SM}$

[Cai Cheng Terning, arXiv:0806.0386]

- Extra vector superfields  $\approx Q + \bar{Q}, g' W' B'$

- $H_5 \approx (5,1) \approx H_d + T^c$

- $g A_i^\dagger T_A^{ij} \lambda_A \psi_j \rightarrow \lambda_t H_d^\dagger Q T^c$

# Strongly coupled models

- **Higgsless (technicolor & C):**  $G_a$  Goldstones of global  $SU(2)_L \times SU(2)_R$   
EWPT not calculable or off; recent progress via duality to 5D

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$$\delta m_h^2 \sim \frac{3G_F}{\sqrt{2}\pi^2} m_t^2 Q_{NP}^2 = m_h^2 \left( \frac{Q_{NP}}{0.5 \text{ TeV}} \right)^2 \text{ for } m_h = 115 \text{ GeV}$$

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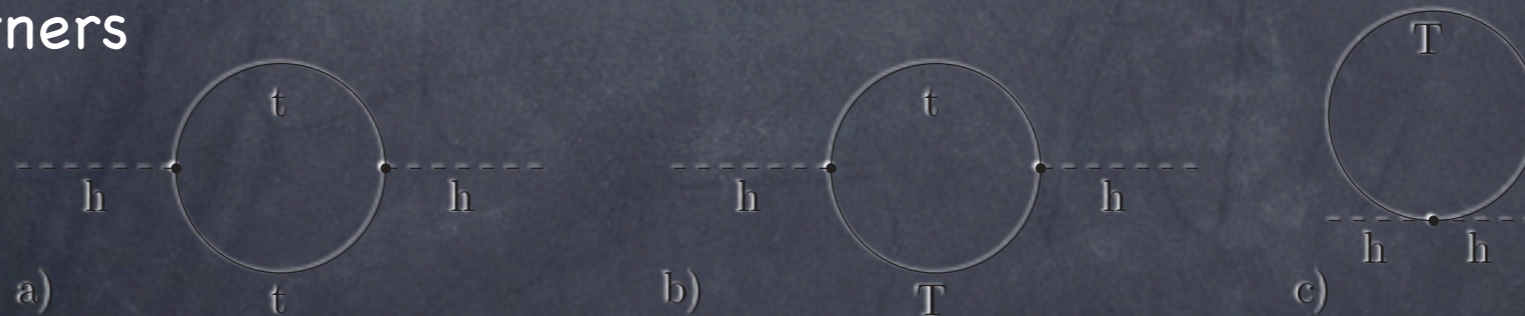
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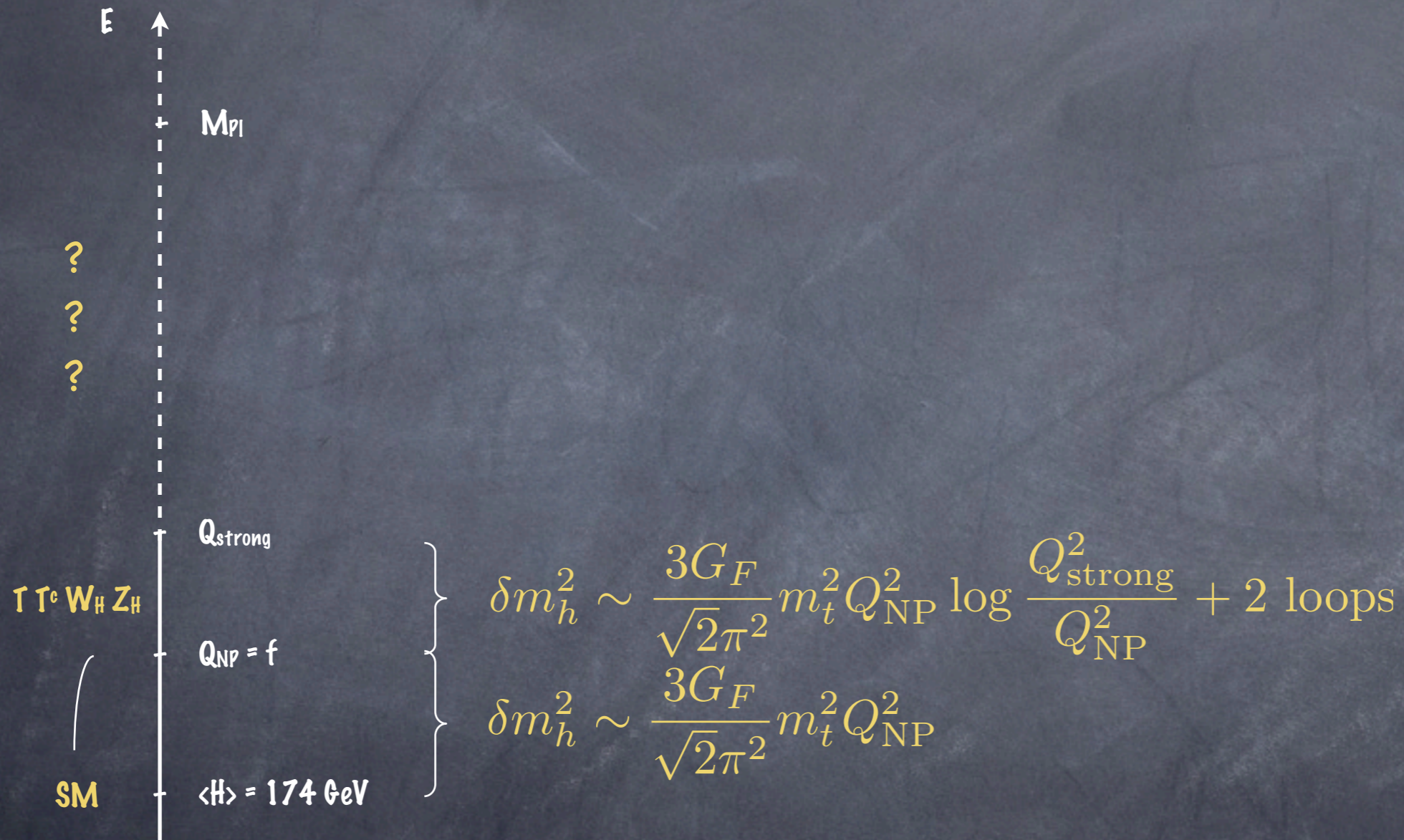
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- More clever explicit breaking ("collective breaking"): **Little Higgs**
  - **no 1-loop  $Q_{NP}^2$  terms** (exact-NGB unless 2+ non-vanishing couplings)
  - the top (gauge, Higgs) loop **must** be cancelled at a lower scale (= global symmetry breaking scale  $f \ll Q_{strong}$ ) by same statistics partners



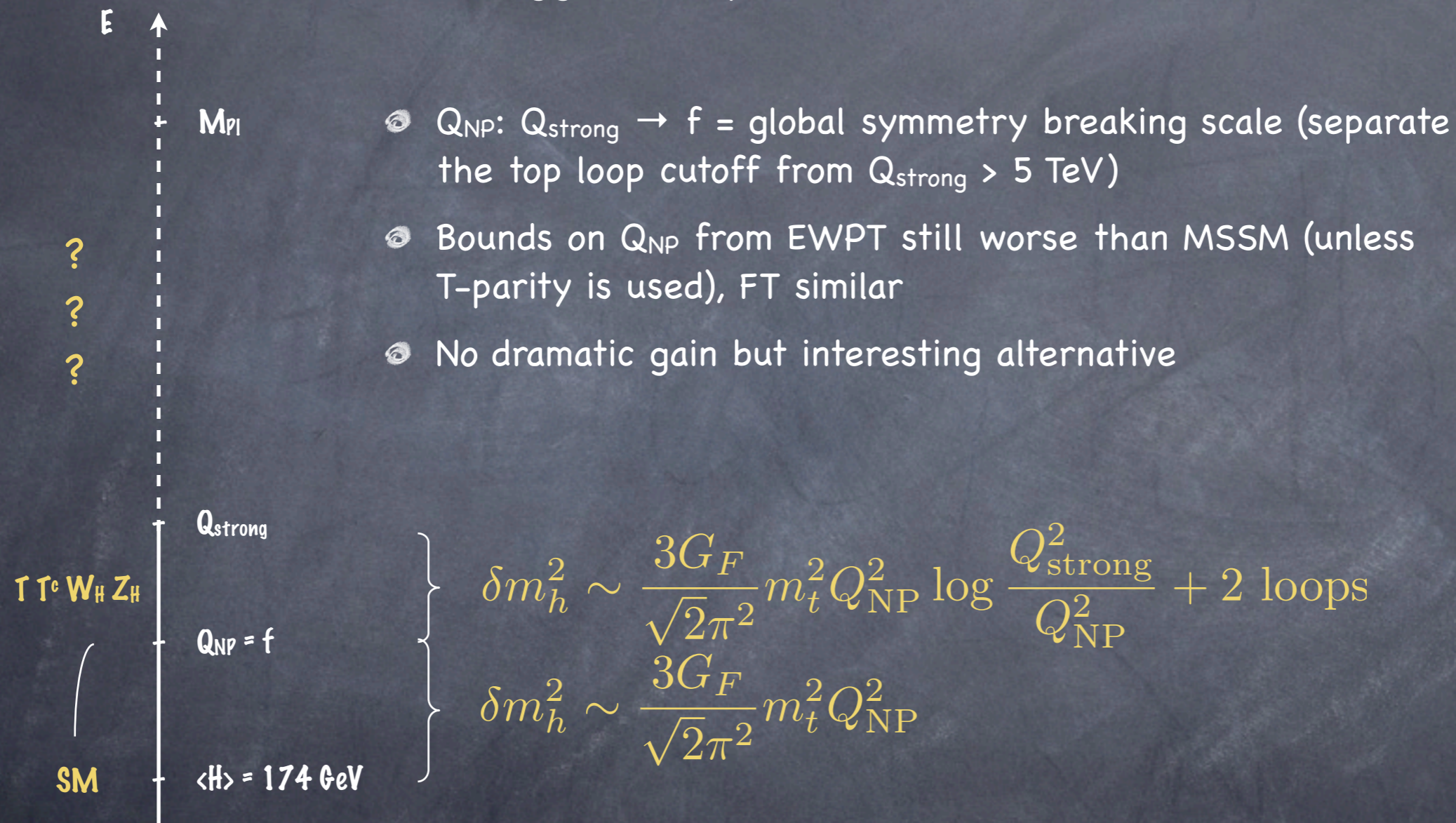
# Little Higgs

Higgs mass protected by  $H(x) \rightarrow H(x) + c$



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# LH @ LHC

- Observe the partners responsible for the divergence cancellation

- $qq \rightarrow Z_H \rightarrow l^+l^-$  up to few TeV (standard); also  $\rightarrow ff, VV, Vh$

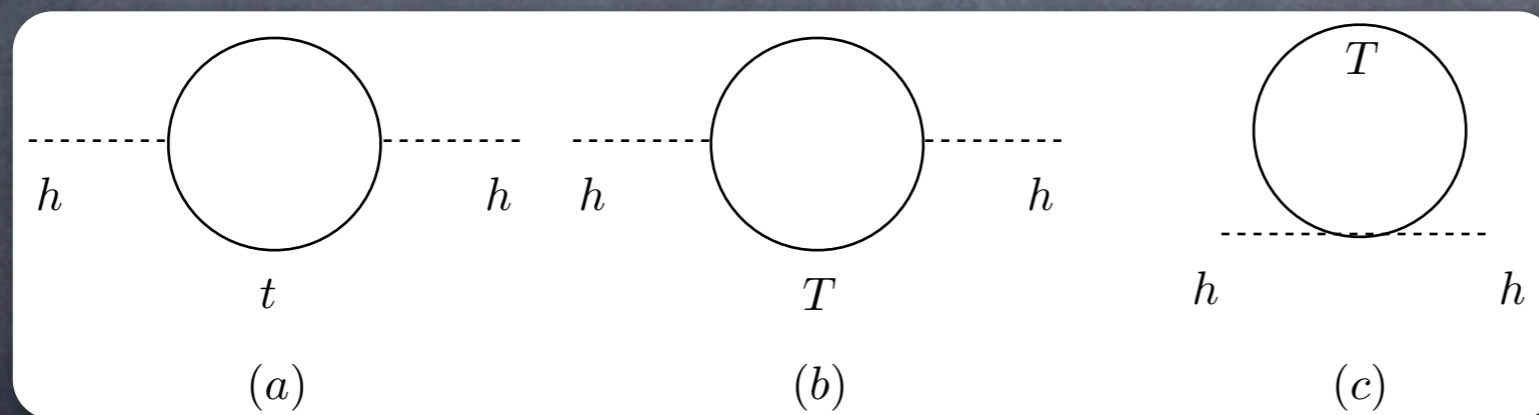
- $T, T^c$ : single production (bWT) dominates (b pdf up to  $\approx 0.2$ )

- $\Gamma(T \rightarrow th) = \Gamma(T \rightarrow tZ) = \Gamma(T \rightarrow bW)/2$  all identifiable:

$tZ \rightarrow bWl^+l^- (m_T), th \rightarrow bWbb (m_h, m_T), bW \rightarrow bl\nu$

- additional (++) Higgs states

- Observe the divergence cancellation



$$-\lambda_t^2$$

$$-\lambda_T^2$$

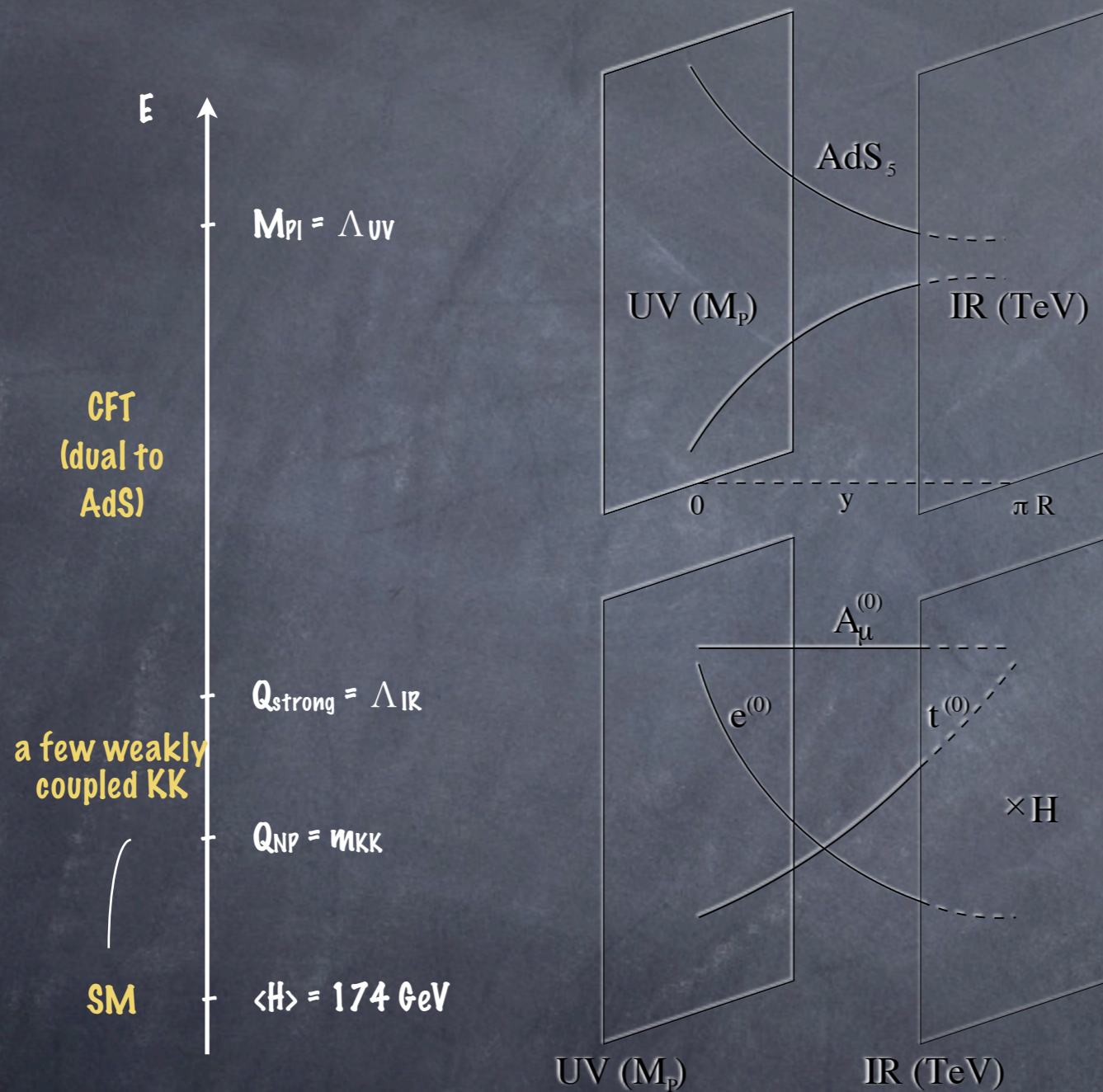
$$+\lambda_T \frac{m_T}{f} = 0$$

$$(a) = -6\lambda_t^2 \int \frac{d^4k}{(2\pi)^4} \frac{1}{k^2},$$

$$(b) = -6\lambda_T^2 \int \frac{d^4k}{(2\pi)^4} \frac{1}{k^2 - m_T^2},$$

$$(c) = +6\frac{\lambda_T}{f} \int \frac{d^4k}{(2\pi)^4} \frac{m_T}{k^2 - m_T^2}$$

# Warping and composite Higgs



- RS + bulk fermions + H as  $(A_5)_0$  + deconstruction = Little Higgs + UV completion
- $m_H$  protection: collective breaking = bc breaking of 5D gauge symmetry
- **4D dual:** UV brane: elementary  
IR brane: composite (H,  $t_R$ )
- $Q_{strong} > 5 \text{ TeV}$  as usual  
 $m_{KK} > \text{TeV}$ , watch  $Z \rightarrow b\bar{b}$
- Gauge coupling unification in a novel way (but limited calculability)

$$m_h \sim M_5 e^{-\pi k R}$$

$k$  = curvature

# @LHC (a first look)

- Keep only first excitation:

- $|SM\rangle = \cos\varphi |elem\rangle + \sin\varphi |comp\rangle$

- $|KK\rangle = -\sin\varphi |elem\rangle + \cos\varphi |comp\rangle$

- Production:

- $A(SM_1 SM_2 \rightarrow KK_3) \propto -g_w \cos\varphi_1 \cos\varphi_2 \sin\varphi_3 + g_s \sin\varphi_1 \sin\varphi_2 \cos\varphi_3$

- $SM_3$  needs to be substantially composite:

- $t_R$  (bW fusion) or  $V_{long}$  (DY) (analogous to LH)

- Decay:

- into  $V_{long}$  and heavier particles ( $t_R$   $b_R$ ,  $\tau$  if non negligible) dominates

- also:  $(gluon)_{KK} \rightarrow t_R \bar{t}_R$

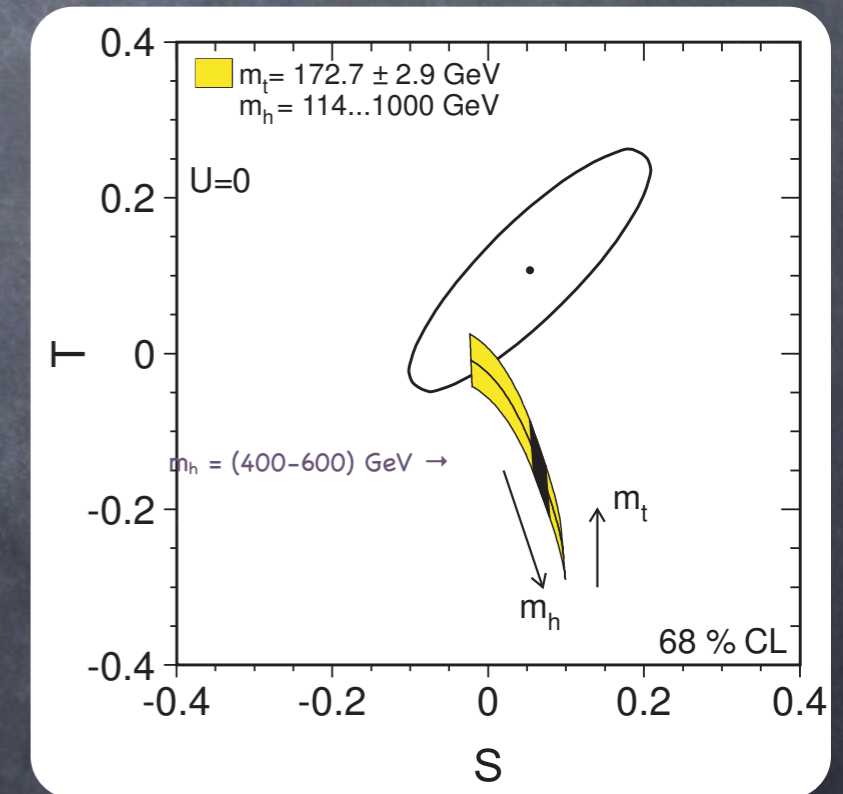
- possibly lepton excitations (if open)

# Back to the residual hierarchy

$$\delta m_h^2 \sim \frac{3G_F}{\sqrt{2}\pi^2} m_t^2 Q_{\text{NP}}^2 = \begin{cases} m_h^2 \left( \frac{Q_{\text{NP}}}{0.5 \text{ TeV}} \right)^2 & \text{if } m_h = 115 \text{ GeV} \\ m_h^2 \left( \frac{Q_{\text{NP}}}{2 \text{ TeV}} \right)^2 & \text{if } m_h = 250 \text{ GeV} \end{cases}$$

$$Q_{\text{NP}} \gtrsim \sqrt{c_i} \cdot 5 \text{ TeV} \approx \begin{cases} 50 \text{ TeV composite SM fermions} \\ 5 \text{ TeV composite Higgs} \\ 0.5 \text{ TeV 1-loop perturbative} \end{cases}$$

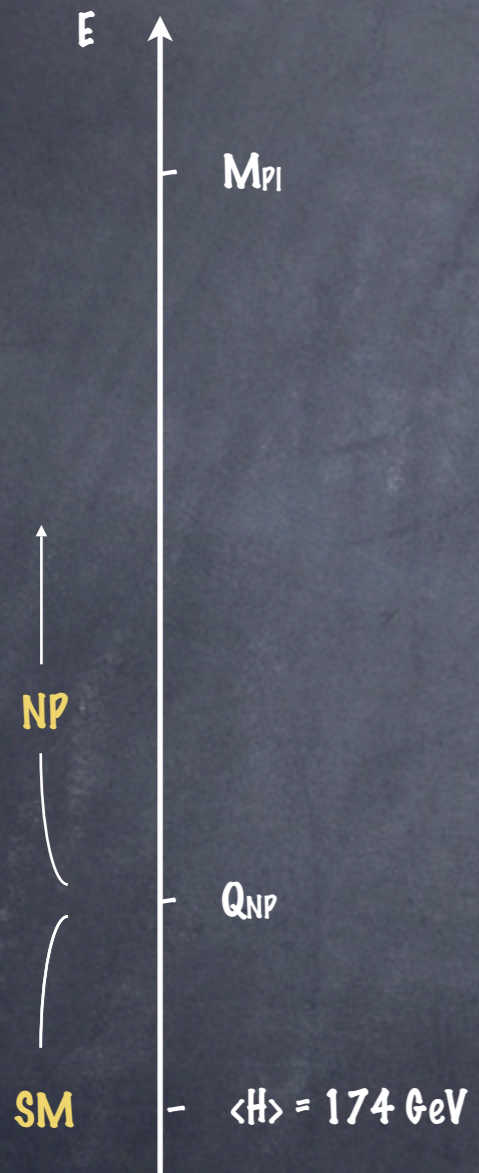
- $m_h = 500 \text{ GeV}$  would help; disfavoured by EWPTs only within the SM
- Cancel SM heavy Higgs contributions to EWPT with NP (goodness of SM + light H fit accidental but not too much fine-tuned)
- Generic prediction of NP giving  $\Delta T = 0.25 \pm 0.1$



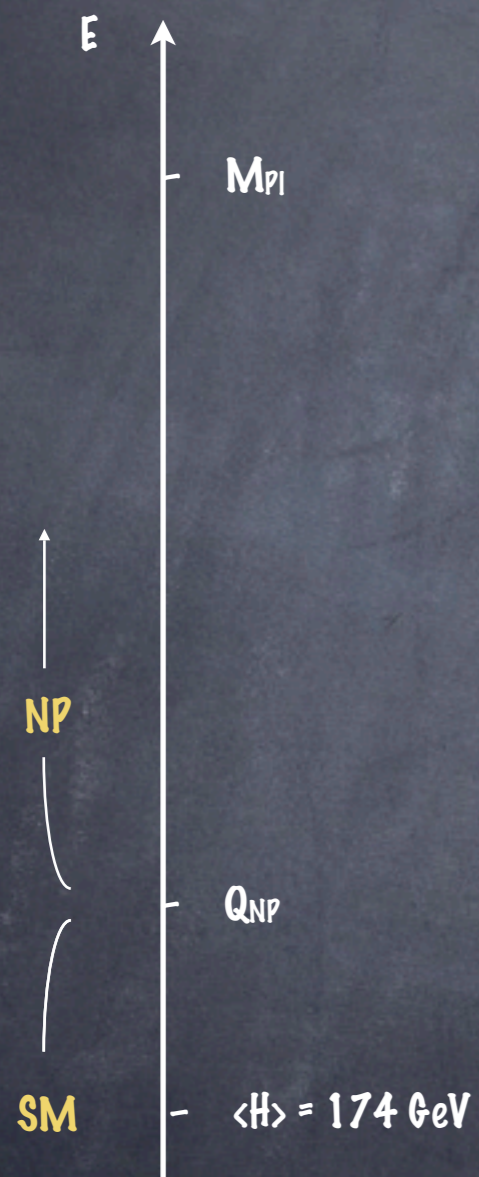
# An inert Higgs

- $H_1$  ( $h$ ): usual Higgs (but heavier): EWSB,  $M_W$   $M_Z$ ,  $m_f$
- $H_2$  ( $H$ ,  $A$ ,  $H^\pm$ ): inert Higgs (60 GeV–1TeV): no vev, no coupling to fermions ( $H_2 \rightarrow -H_2$ ), gives  $\Delta T = 0.25 \pm 0.1$
- DM candidate for  $m_H \approx 70$  GeV (LEP?)
- Pair production:  $pp \rightarrow W^* \rightarrow H^+H$ ,  $H^+A$  or  $pp \rightarrow Z^* \rightarrow H^+H^-$ ,  $HA$
- Decay into the lightest + gauge bosons (no fermions)  $\rightarrow$  charged leptons in the final states
- UV completion?

# Is a natural $m_H$ unavoidable?

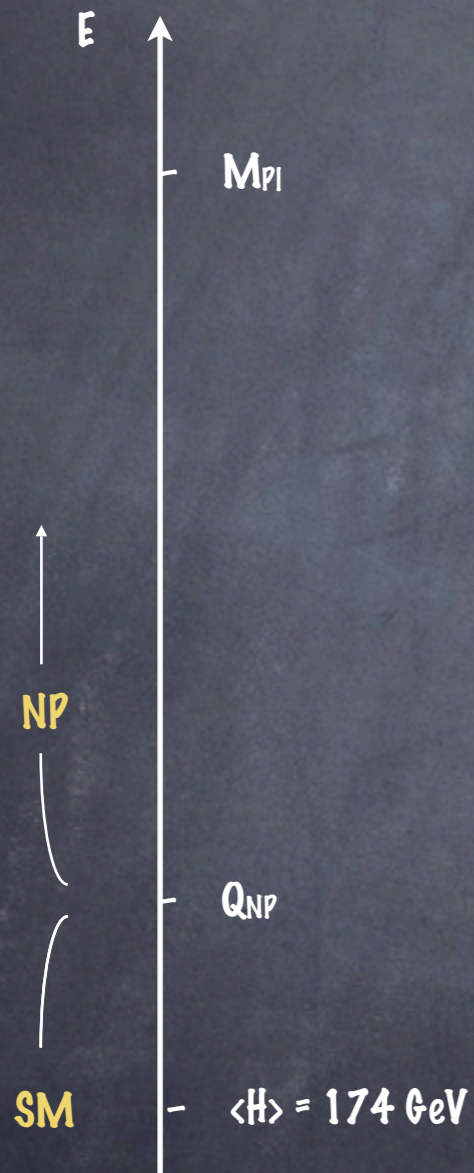


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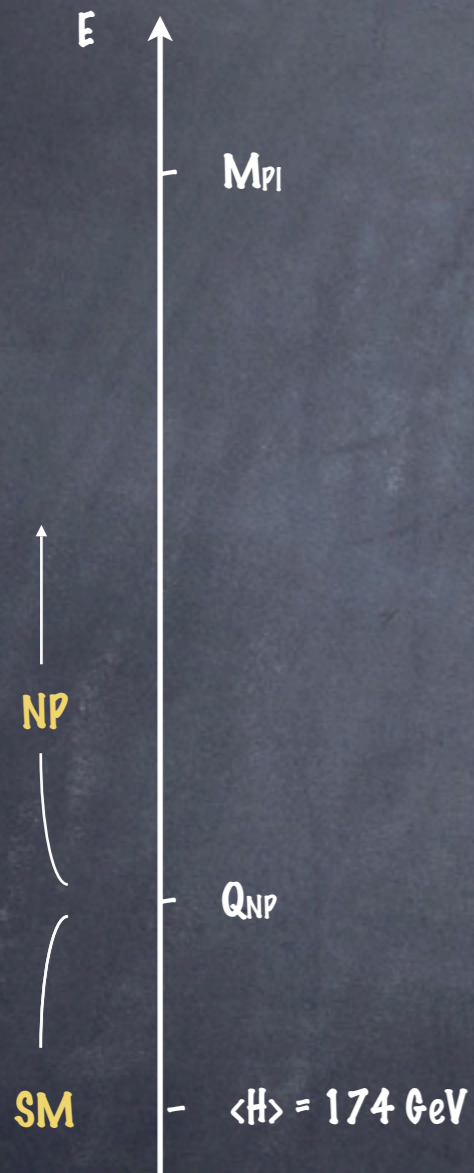
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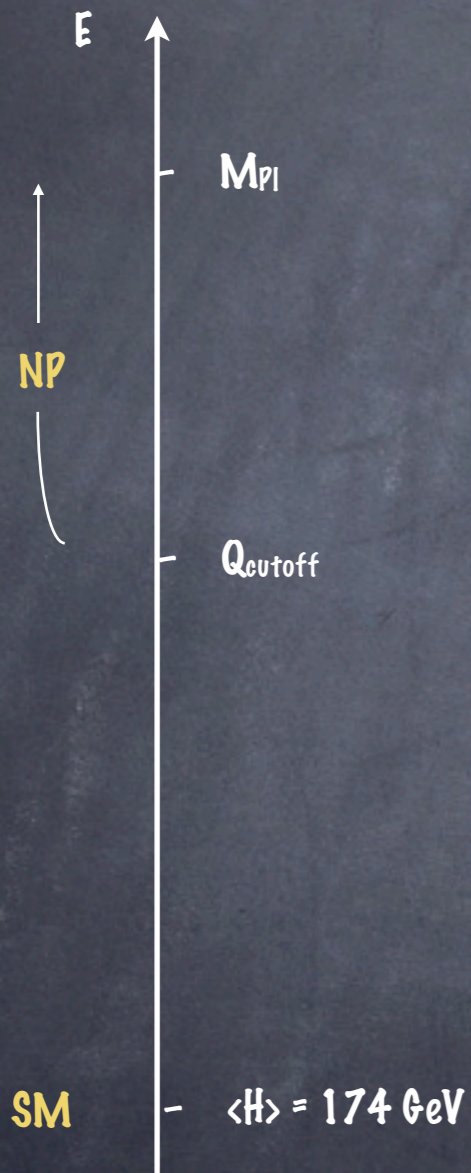
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- If the  $m_h$  naturalness criterium is irrelevant, what are the observable consequences?

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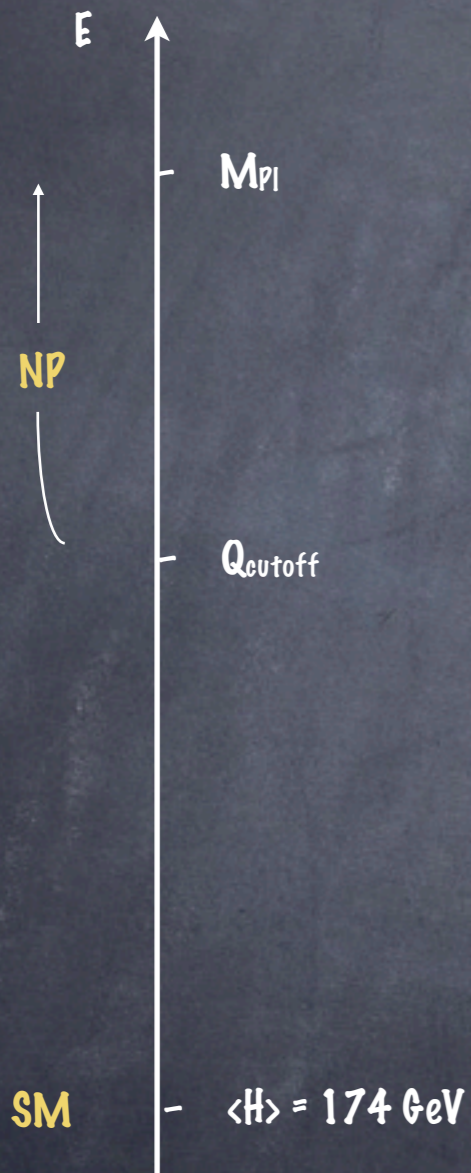
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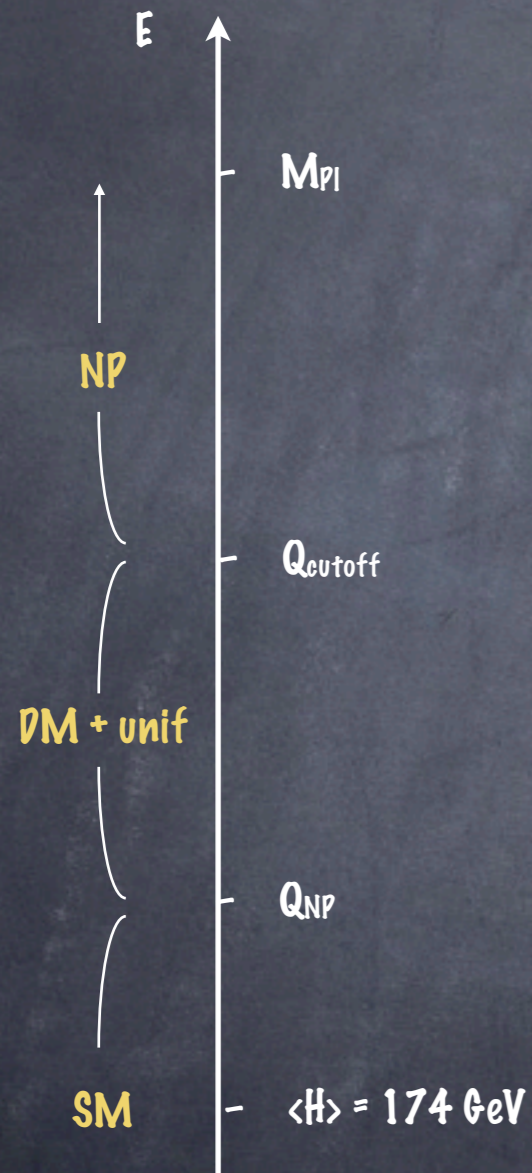
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- If the  $m_h$  naturalness criterium is irrelevant, what are the observable consequences?
- LHC..?
- Dark matter still motivates NP at the TeV scale

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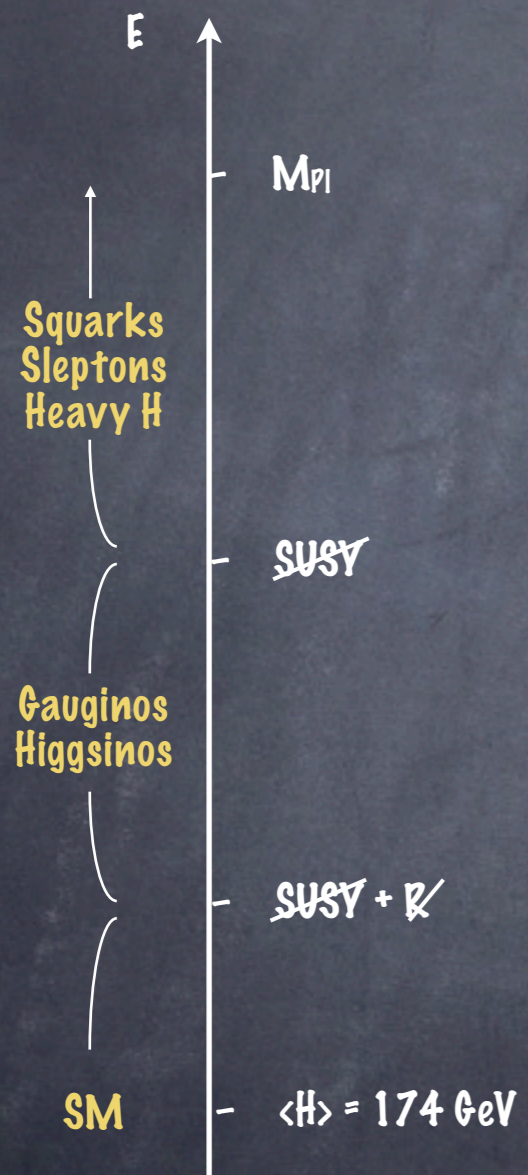
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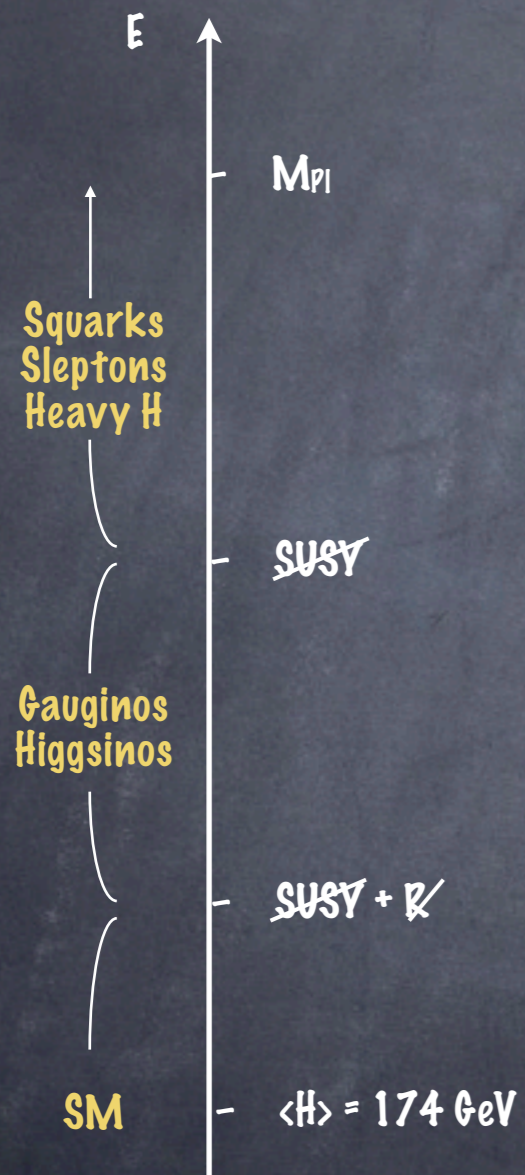
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# Split Supersymmetry



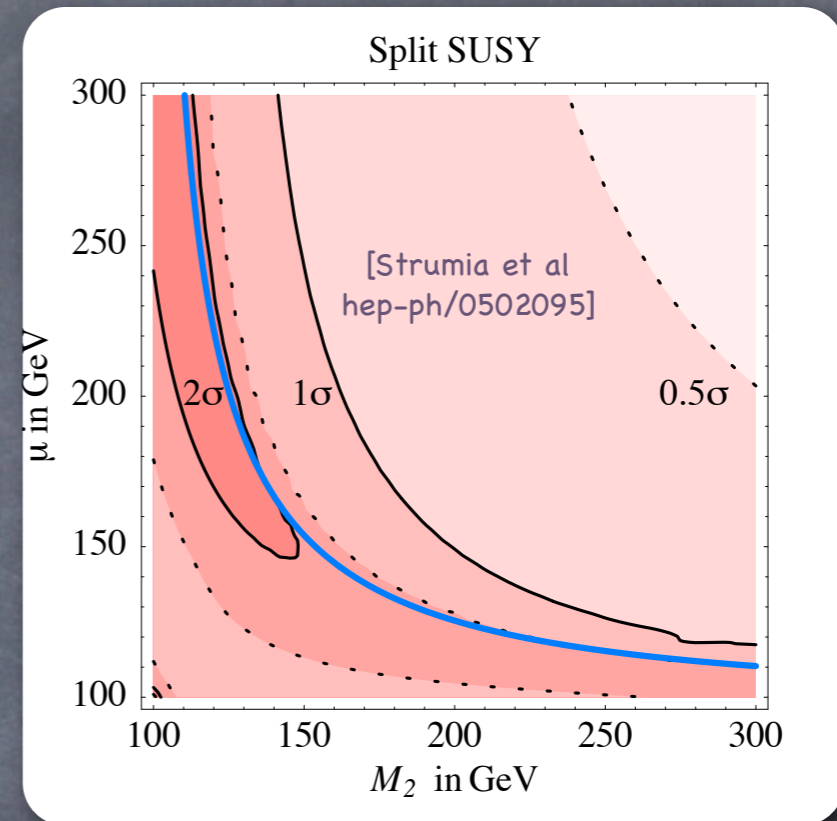
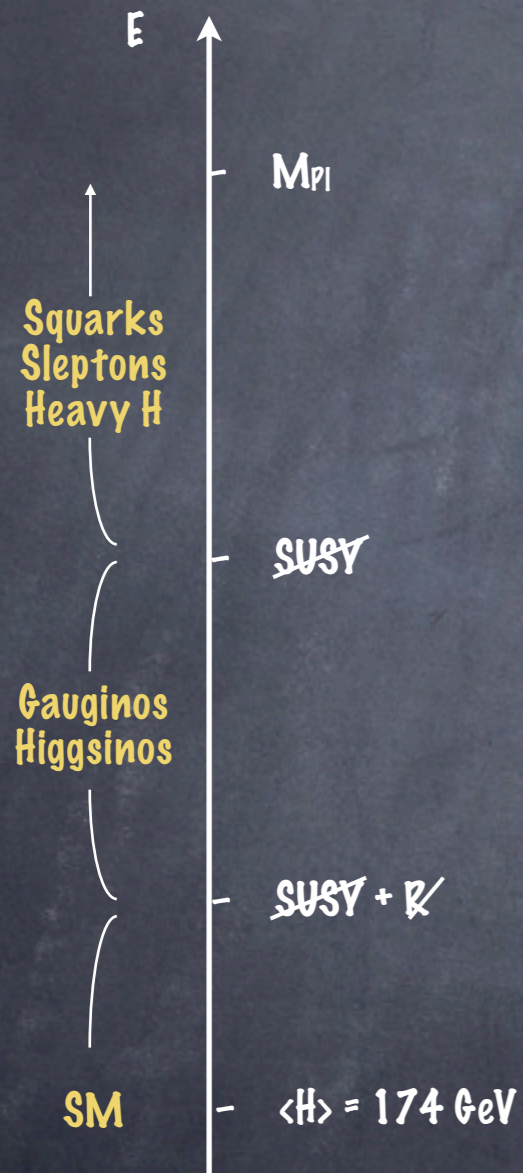
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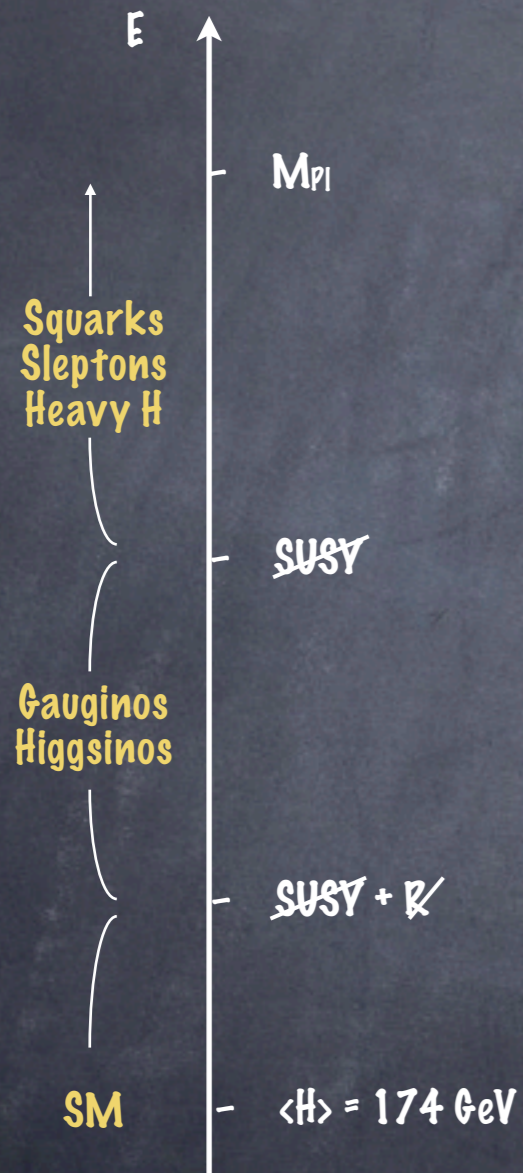
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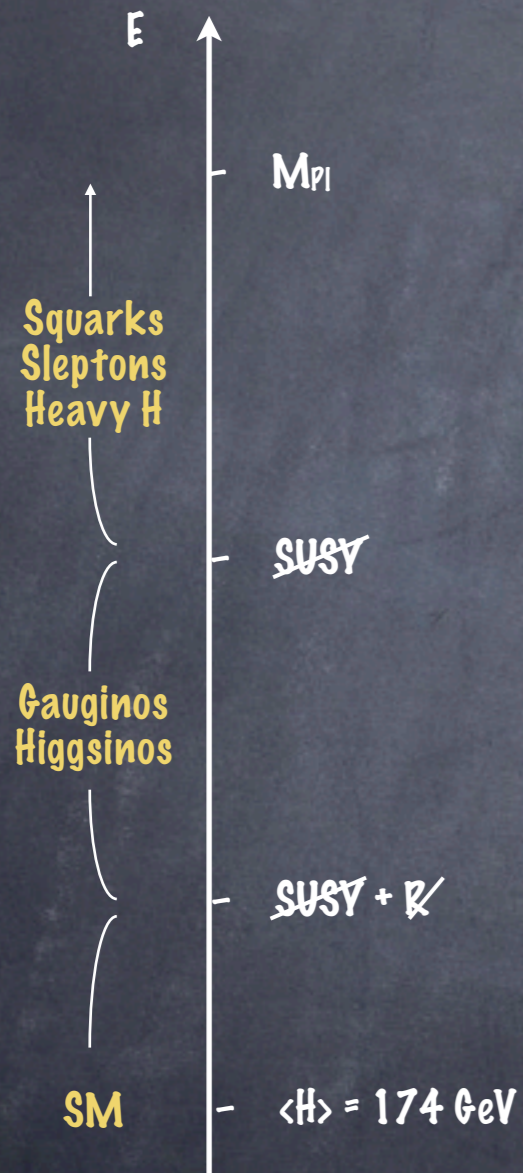
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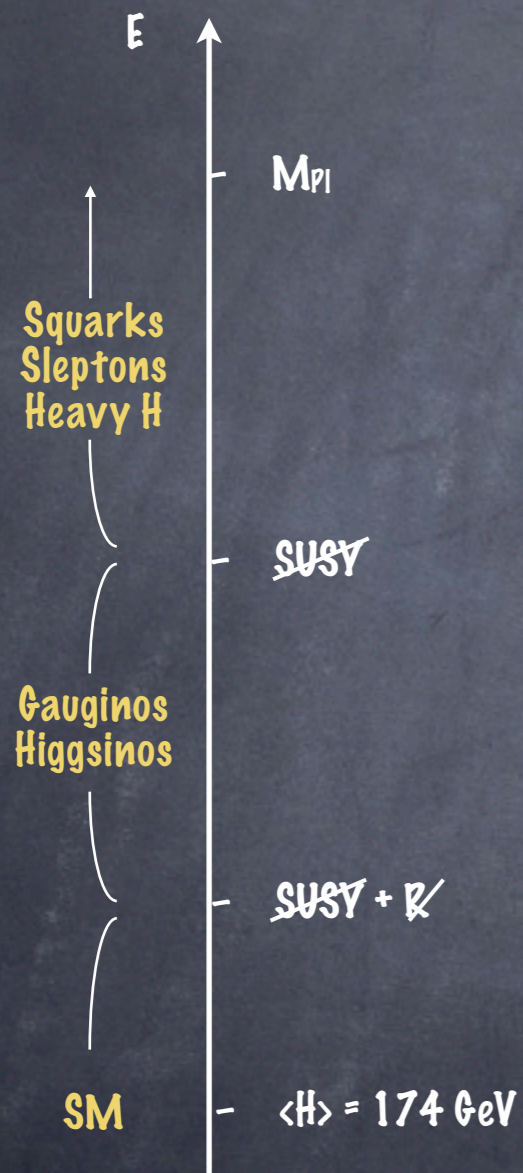


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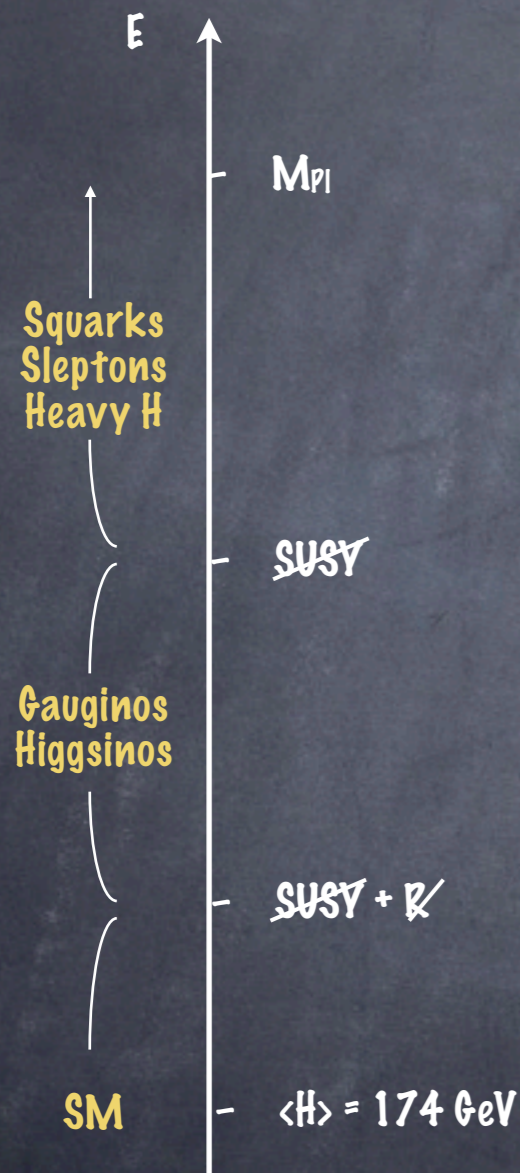


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- **Wilder**: stopping gluinos (1-2 jets in any direction from denser parts of the detector + m.e.), displaced vertexes (low m), charge flips

# Summary

- Is a % tuning really worth worrying?
- If not, NP could as well be out of reach of the LHC
- Barring independent arguments (e.g. DM)
- Useful and fruitful guideline within models addressing the naturalness issue
- Surprises are not unlikely

- Experimental “problems” of the SM:
  - Gravity
  - Dark matter
  - Baryon asymmetry
- Experimental “hints” of physics beyond the SM
  - Neutrino masses
  - Quantum number unification
- Theoretical puzzles of the SM:
  - $\langle H \rangle \ll M_{Pl}$
  - Family replication
  - Small Yukawa couplings, pattern of masses and mixings
  - Gauge group, no anomaly, charge quantization, quantum numbers
- Theoretical problems of the SM:
  - Naturalness / unitarity problem
  - Cosmological constant problem
  - Strong CP problem