

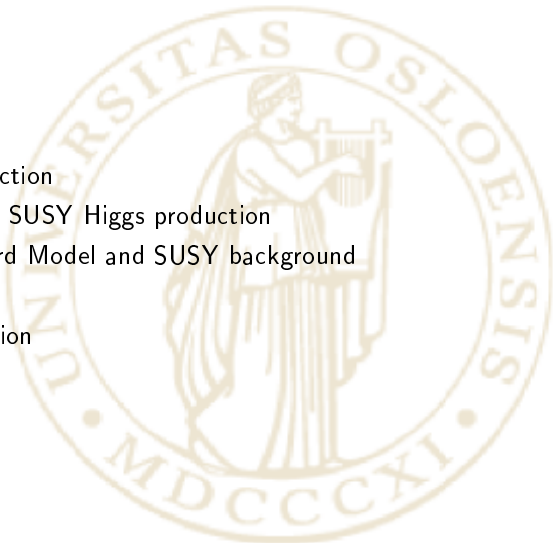
Search for the lightest MSSM Higgs boson in cascades of supersymmetric particles in ATLAS

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- introduction
- lightest SUSY Higgs production
- Standard Model and SUSY background
- results
- conclusion



- Supersymmetry (SUSY) is one of the most believed *new physics* phenomena.
- need two Higgs doublets in MSSM:

$$H_u = \begin{pmatrix} h_u^+ \\ h_u^0 \end{pmatrix}, \quad H_d = \begin{pmatrix} h_d^- \\ h_d^0 \end{pmatrix},$$

to give mass to up and down type fermions

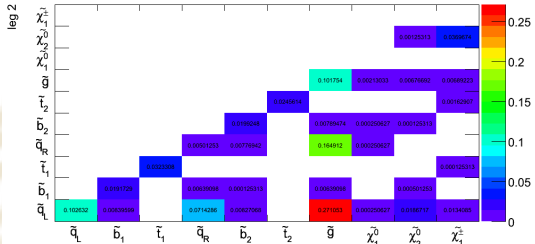
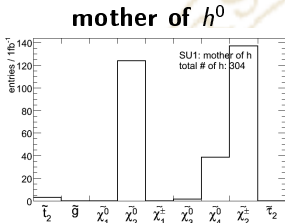
- 8° of freedom, so SUSY predicts 5 Higgs bosons
 - the lightest Higgs, h^0 . $m_h < 135$ GeV
 - heavy CP-even Higgs, H^0
 - charged Higgs, H^\pm
 - CP-odd Higgs, A^0
- Ratio of vacuum expectation values (vev) of up and down field:
 $\tan \beta = v_u/v_d$
- gravity mediated soft SUSY breaking (mSUGRA)
 - described by five parameters at the GUT scale: common scalar (m_0) and gaugino mass ($m_{1/2}$), a common trilinear coupling (A_0) and the two Higgs parameters $\tan \beta$ and the sign of μ

- the lightest Higgs is light (< 135 GeV) so $h^0 \rightarrow b\bar{b}$ dominates ($BR = 0.7 - 0.85$). The remaining is into $\tau^+\tau^-$
 - this motivates us to search for h^0 in SUSY cascade since the direct production (e.g. gluon-gluon fusion) will suffer from large QCD background
- most obvious channel is $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 h$ (if open)
 - only dominant if $\tilde{\chi}_2^0 \rightarrow \tilde{l}l$ is inaccessible
 - $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 h$ dominates $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 Z$ due to the mostly gaugino like $\tilde{\chi}_2^0$ in mSUGRA
- also other Higgs production decays are possible
 - $\tilde{\chi}_{3,4}^0 \rightarrow \tilde{\chi}_{1,2}^0 h$, although $\tilde{\chi}_{3,4}^0$ decays are mostly dominated by decays into $\tilde{\chi}_1^\pm$ and W^\pm
 - $\tilde{\chi}_2^\pm \rightarrow \tilde{\chi}_1^\pm h$, although $\tilde{\chi}_2^\pm$ decay are mostly dominated by decays into $\tilde{\chi}_2^0$ and W^\pm
 - also possible with $\tilde{t}_2 \rightarrow \tilde{t}_1 h$ or $\tilde{\tau}_2 \rightarrow \tilde{\tau}_1 h$

	SU1	SU9
sparticle	mass [GeV]	
h^0	116	114
scalars	$\approx 570 - 760$	$\approx 730 - 950$
gluino	832	990
$\tilde{\chi}_1^0$ (LSP)	137	173
$\tilde{\chi}_2^0$	264	325
$\tilde{\chi}_{3,4}^0$	~ 480	~ 540
$\tilde{\chi}_1^\pm$	262	325
$\tilde{\chi}_2^\pm$	484	545
A_0	512	625
decay	Branching Ratio [%]	
$\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 h$	5	86
$\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 Z$	1	12
$\tilde{\chi}_2^\pm \rightarrow \tilde{\chi}_1^\pm h$	19	27
$\tilde{\chi}_2^\pm \rightarrow \tilde{\chi}_1^\pm Z$	25	22

Two benchmark points
are presented here
($m_0, m_{1/2}, A_0, \tan \beta, \text{sgn}(\mu)$)
SU1:
(70, 350, 0, 10, +)
SU9:
(300, 425, 200, 20, +)

initially produced sparticle pair (when h^0 prod.)



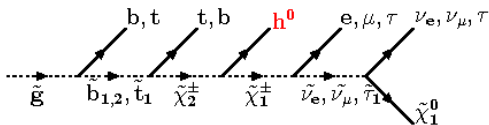
$$BR(\tilde{\chi}_2^\pm \rightarrow \tilde{\chi}_1^\pm h) = 0.19$$

$$BR(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 h) = 0.05$$

$$BR(\tilde{\chi}_3^0 \rightarrow \tilde{\chi}_1^0 h) = 0.02$$

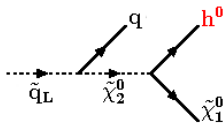
- mostly **gluino** and light **squark** production
- seldom production of **stop** and **sbottom**

main decay through $\tilde{\chi}_2^\pm$:



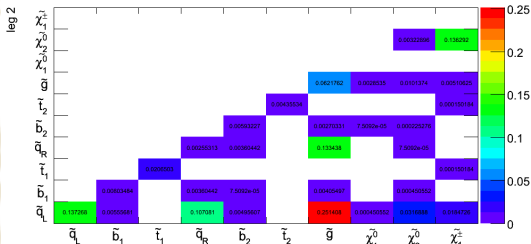
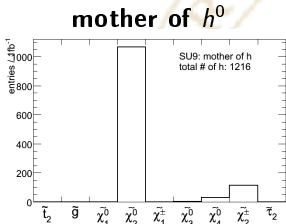
$$BR = 0.014 (\tilde{b}_{1,2}), BR = 0.004 (\tilde{t}_1)$$

main decay through $\tilde{\chi}_2^0$:



$$BR = 0.030 (\tilde{q}_L), BR = 0.017 (\tilde{b}_{1,2}), BR = 0.007 (\tilde{t}_1)$$

initially produced sparticle pair (when h^0 prod.)



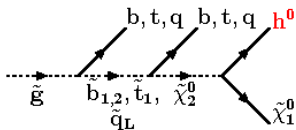
$$BR(\tilde{\chi}_2^\pm \rightarrow \tilde{\chi}_1^\pm h) = 0.22$$

$$BR(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 h) = 0.86$$

- mostly **gluino** and light **squark** production
- seldom production of **stop** and **sbottom**
- also initially production of $\tilde{\chi}_2^0$

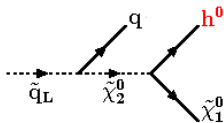
(almost) only h^0 production through $\tilde{\chi}_2^0$

long cascade:



$$BR = 0.022 (\tilde{q}_L), BR = 0.007 (\tilde{b}_{1,2}), BR = 0.036 (\tilde{t}_1)$$

short cascade:

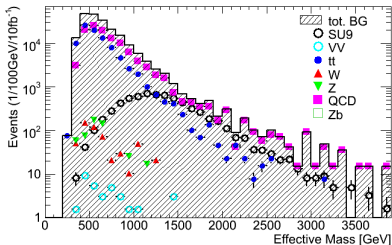
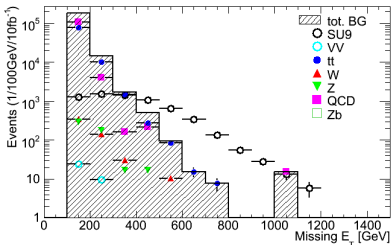


$$BR = 0.82 (\tilde{q}_L), BR = 0.21 (\tilde{b}_{1,2}), BR = 0.12 (\tilde{t}_2)$$

	SU1	SU9
number of Higgs ($\mathcal{L} = 1fb^{-1}$)	304	1216
typical final state	b or light jets, leptons, \cancel{E}_T	light jets, \cancel{E}_T

- Standard Model is not the biggest challenge, can be rejected by using the \cancel{E}_T and effective mass:

$$M_{eff} = \sum_i^{N_{jets}} p_T^i + \cancel{E}_T.$$



..., however the SUSY background is problematic

$$\tilde{g} \rightarrow b \tilde{b}_{1,2} \rightarrow b b \tilde{\chi}_{1,2,3,4}^0$$

$$\tilde{g} \rightarrow b \tilde{b}_{1,2} \rightarrow b t \tilde{\chi}_{1,2}^\pm$$

$$\tilde{g} \rightarrow t \tilde{t}_1 \rightarrow t b \tilde{\chi}_{1,2}^\pm$$

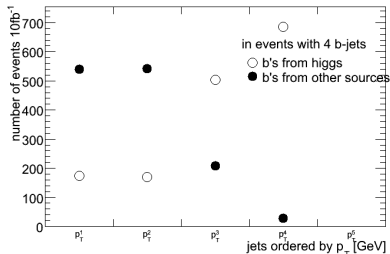
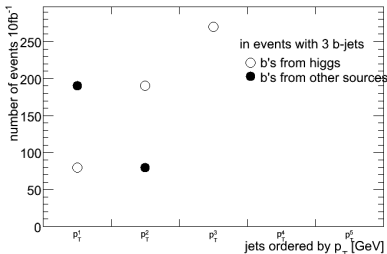
$$\tilde{g} \rightarrow t \tilde{t}_1 \rightarrow t t \tilde{\chi}_{1,2,3,4}^0$$

- also b -jets from W^\pm and Z -decays
- since all SUSY events contain two LSPs, \cancel{E}_T can not be used to reduce SUSY background
- in case of $SU9$ a lepton veto could be efficient, remove some of the background, both SM and SUSY
- ... however; not good for $SU1$

lepton veto after cuts: $\cancel{E}_T > 100$ GeV and ≥ 2 jets, $p_T^{1/2} > 100/50$:

$\mathcal{L} = 10fb^{-1}$	$SU1$		$SU9$	
	S	B	S	B
no lep. veto	165	116875	1413	115030
lep. veto	101	95618	1121	94855
efficiency(ε), rejection(R)	0.6, 0.18		0.80, 0.18	

- there are often produced b -jets in SUSY cascades
- might have high p_T since they come from decays of heavy squarks or gluinos



From this:

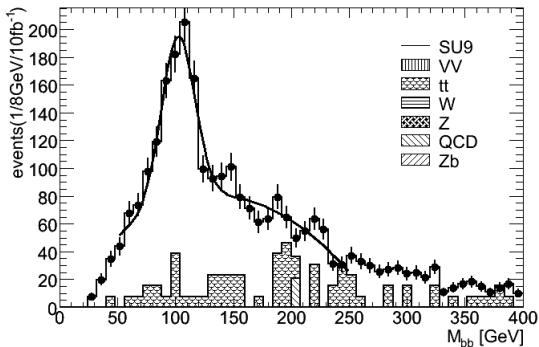
- omitting the hardest (and second hardest) b -jet(s) in events with 3 (> 3) b -jets we remove some of the wrong b -jet combinations in the M_{bb} calculation

Final cut used:

CUT #	VALUE
CUT 1	2 b -jets with $p_T^{1,2} > 50, 20$ GeV
CUT 2	omit hardest (and sec. hardest) b -jet in events with 3 (> 3) b -jets from M_{bb} calculation
CUT 3	$\cancel{E}_T > 300$ GeV
CUT 4	$M_{\text{eff}} > 800$ GeV
CUT 5	≥ 2 additional jets $p_T^{1,2} > 100, 100$ GeV

b -tagging efficiency: 60% \Rightarrow efficiency of 36% per b -jet pair.

$$\mathcal{L} = 10\text{fb}^{-1}$$

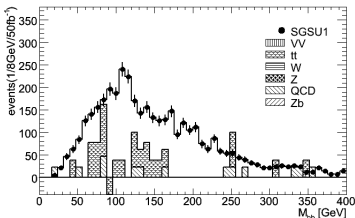


- have required the two additional jets to be light flavored
- see a clear signal above SUSY + SM background
- fitted with a Gaussian superimposed to a second degree polynomial

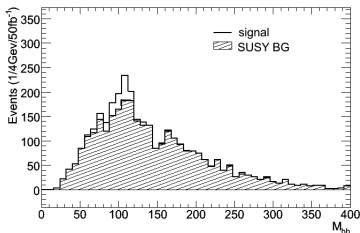
Number of events in $\pm 1\sigma$:

	signal	SUSY BG	SM BG	significance
events	681	777	426	15.1

$\mathcal{L} = 50\text{fb}^{-1}$



- have required one of the two additional jets to be *b*-jet
- more difficult to extract a resonance over the SUSY background
 - however in *SU1* we see a signal with $\mathcal{L} = 50\text{fb}^{-1}$



- lightest MSSM Higgs boson is light (< 135 GeV) so $h \rightarrow b\bar{b}$ dominates
- when $\tilde{\chi}_2^0 \rightarrow \tilde{l}l$ is not kinematically allowed, $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 h$ dominates
 - if not: h^0 is produced mainly through chargino and heavy neutralino decays
 - these decays are never dominated by the decay into h^0 , leading to few Higgses
- Standard Model background not the main challenge (E_T and M_{eff})
- the SUSY events contain many b -jets and is the main background, some is removed by:
 - omit the hardest (and second hardest) b -jet in the M_{bb} calculation in events with 3 (> 3) b -jets
 - require more jets with specific flavors
- in all SUX samples except $SU9$ the $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 h$ is suppressed and to extract a Higgs resonance one will need high luminosity and/or more sophisticated methods