

# STUDY OF SOLAR CORONAL MASS EJECTIONS IN SPHERICAL GEOMETRY



PIYALI CHATTERJEE  
INDIAN INSTITUTE OF ASTROPHYSICS  
BANGALORE



Pencil-code user meeting, 27-31 Jul, 2020, Glasgow



# INDIA'S FIRST SOLAR MISSION: ADITYA-L1

**Payloads: 7 nos**

**REMOTE instruments-4 nos**

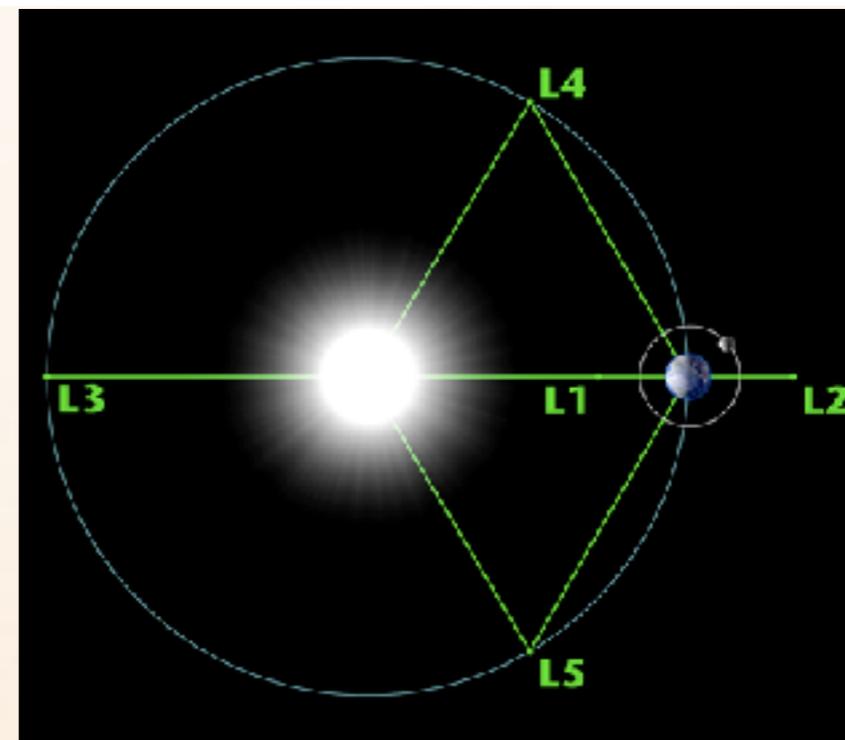
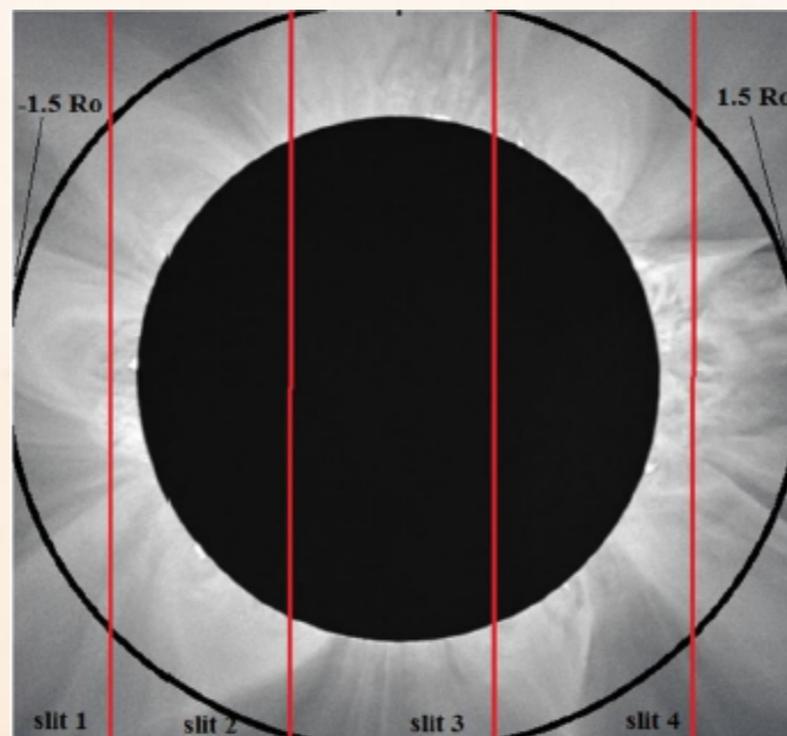
**a. Visible Emission Line Coronagraph (VELC):**

Continuum channel (1.05-1.5  $R_s$ ),  
Spectrograph (coronal red/green lines),  
spectropolarimetry (NIR Fe line, magnetic field measurements)

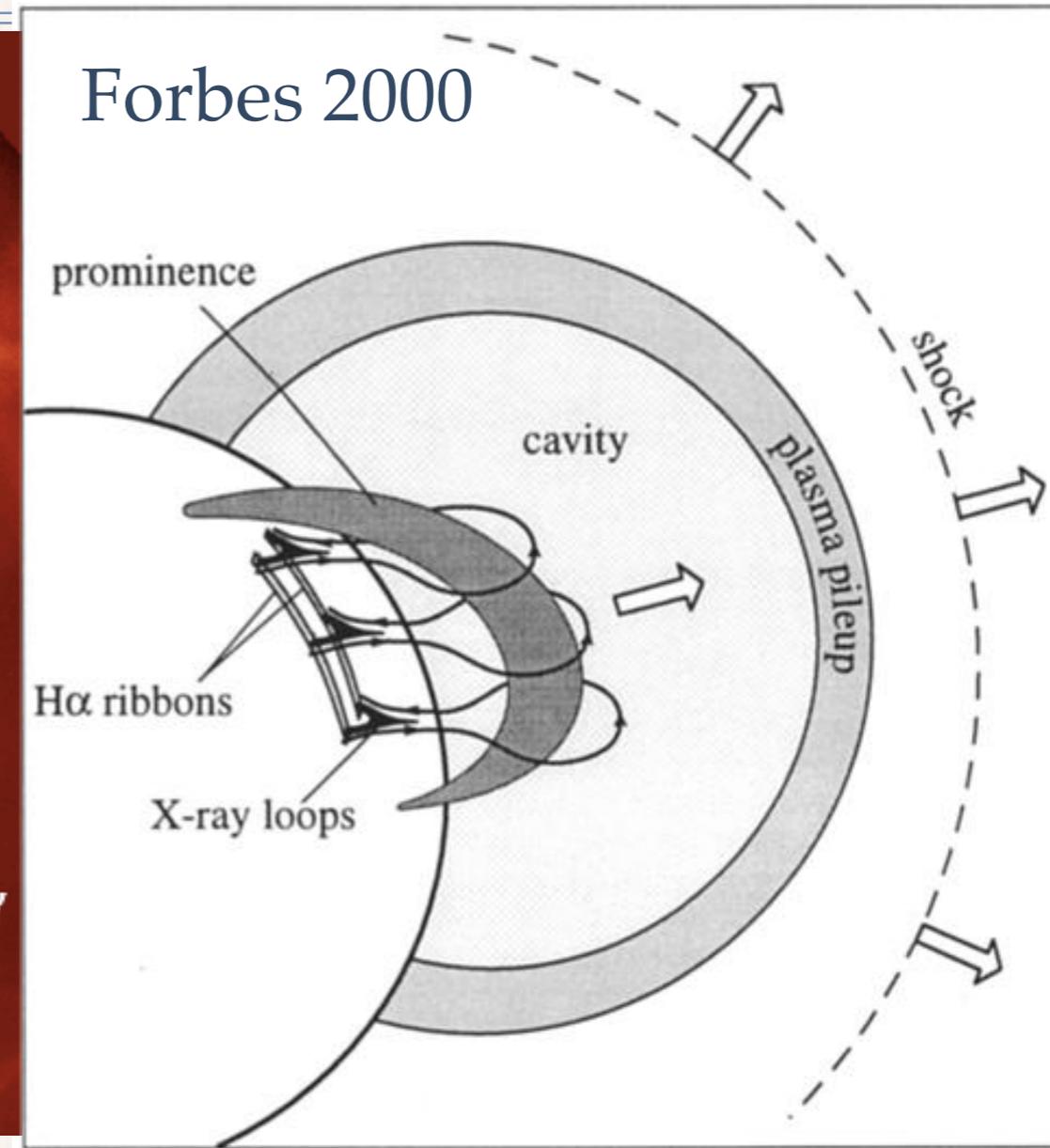
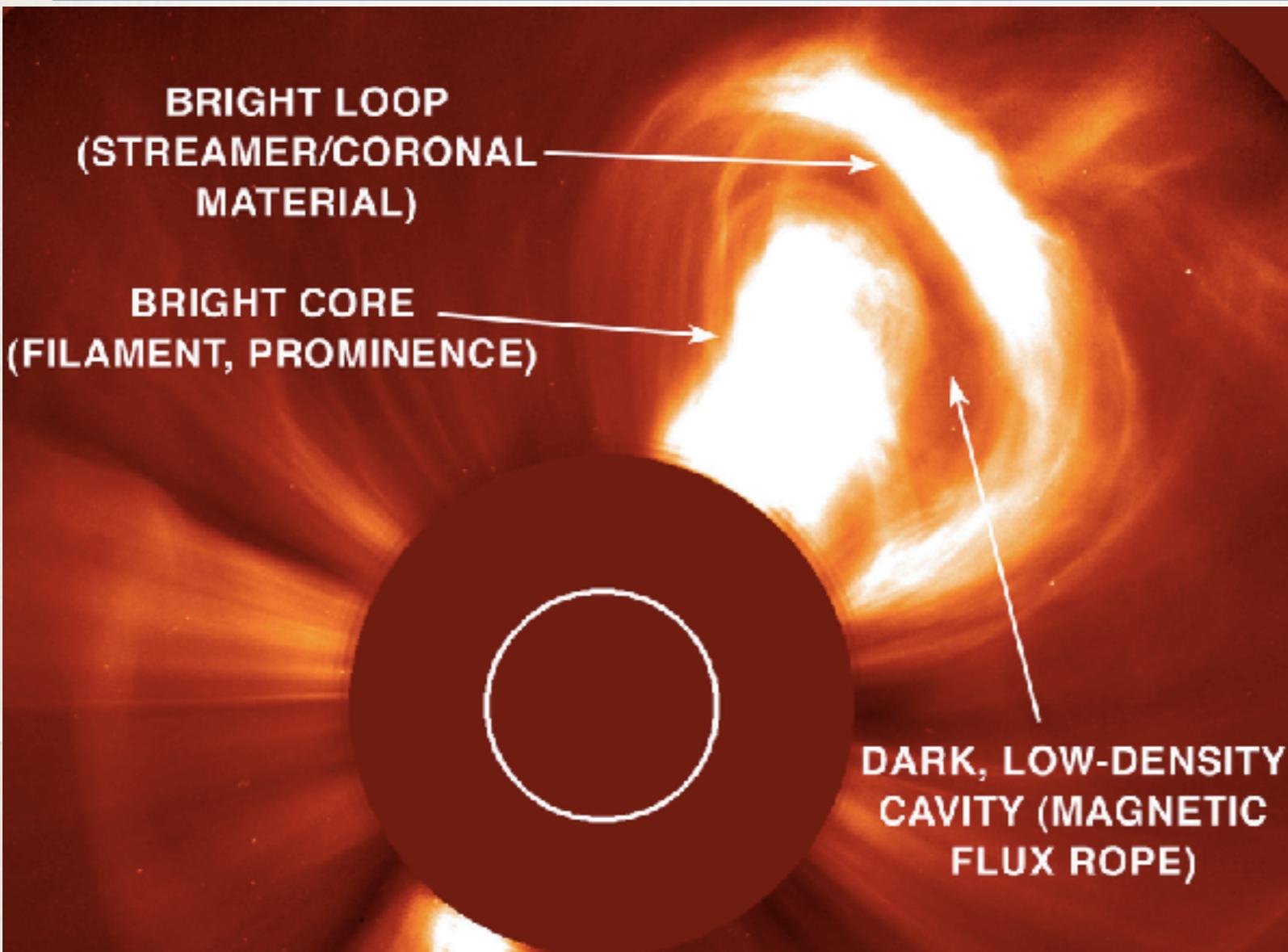
**b. Solar Ultraviolet Imaging Telescope (SUIT):**

Full disk (200-400 nm) Study of prominences,  
solar irradiance

**IN-SITU instruments-3 nos**



# CME: 3-part structure



Considered standard morphology. Although only 30% of observed CMEs show all three parts! (*Webb & Hundhausen 1987*)



# CME precursor: Flux rope formation

---

The magnetic flux rope (MFR) fits beautifully into the 3-part structure of the CME (**Filament-Cavity-Frontal loop**).

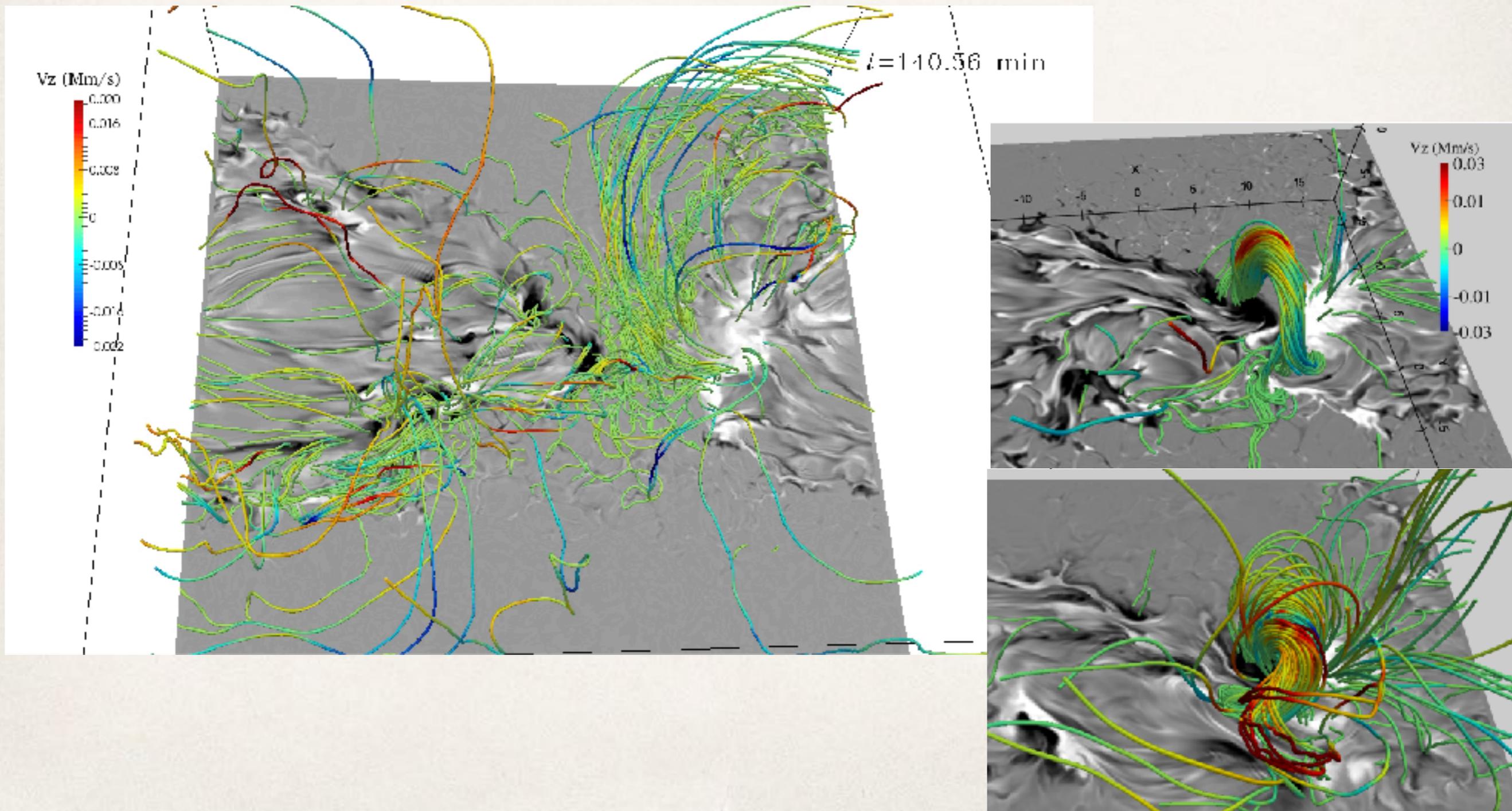
How does a flux rope form?

1. Shearing of foot points and Flux cancellation
2. Emergence of a twisted flux rope

# CME precursor: Flux rope formation

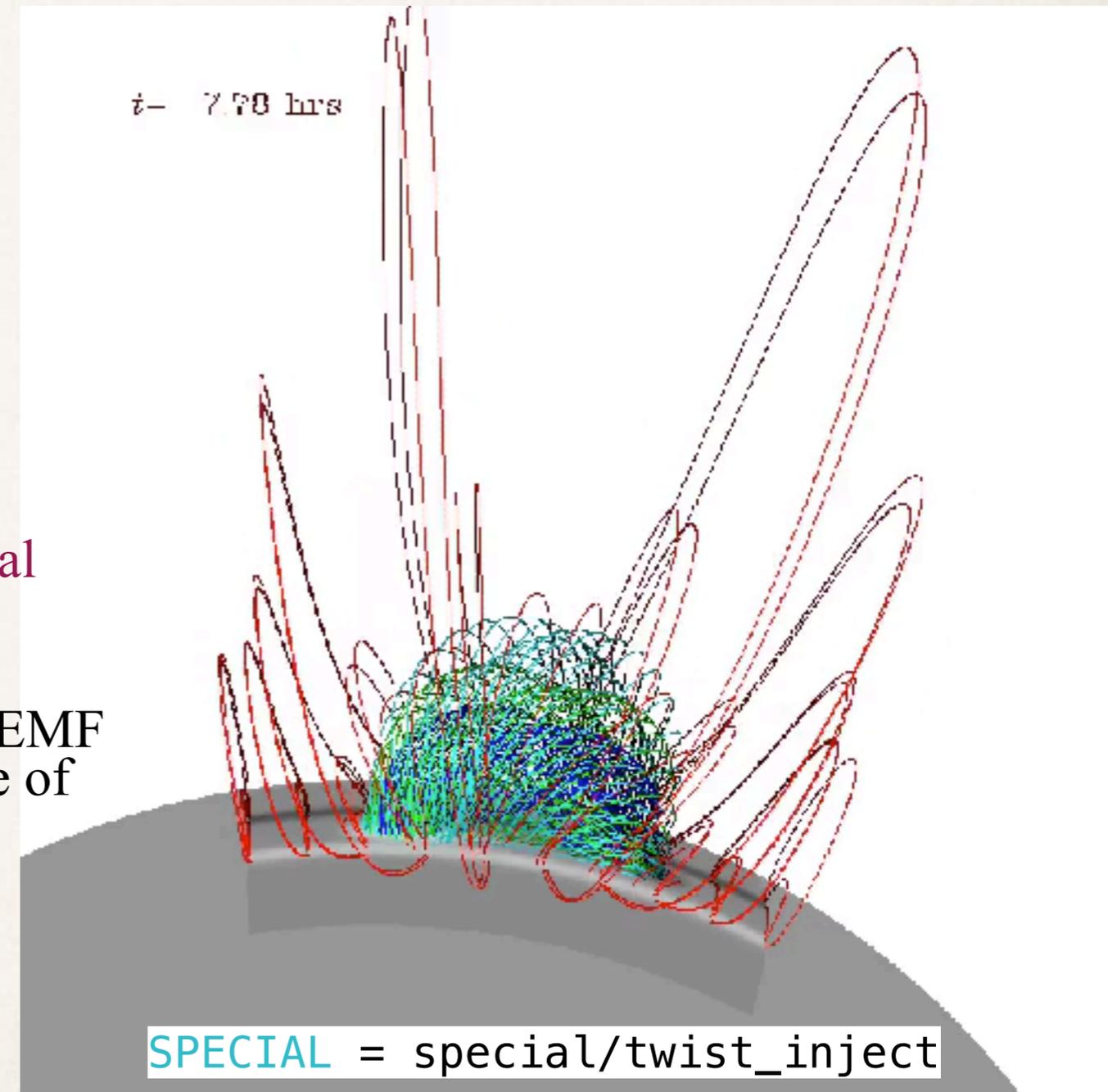
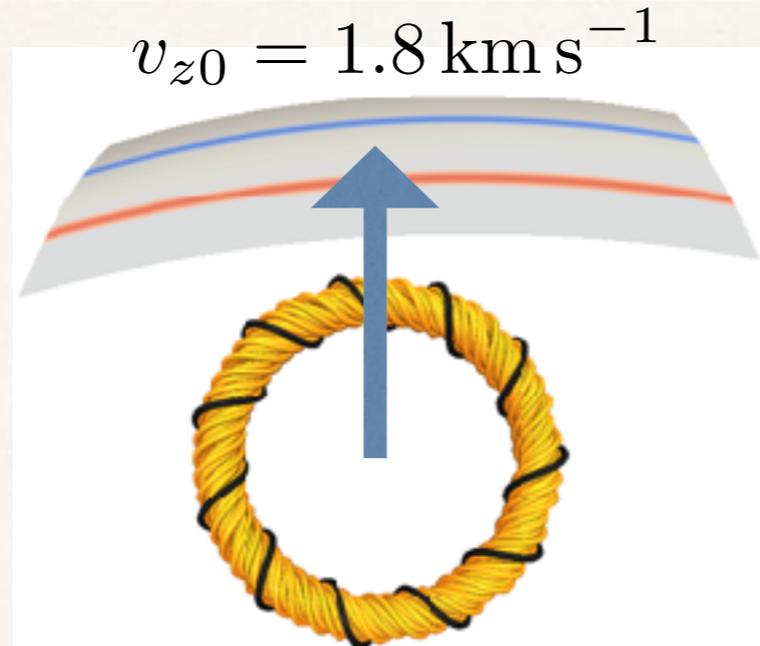


## 1. Shearing and Flux cancellation [Chatterjee et. al. 2016]





# My typical (data inspired) CME initiation model

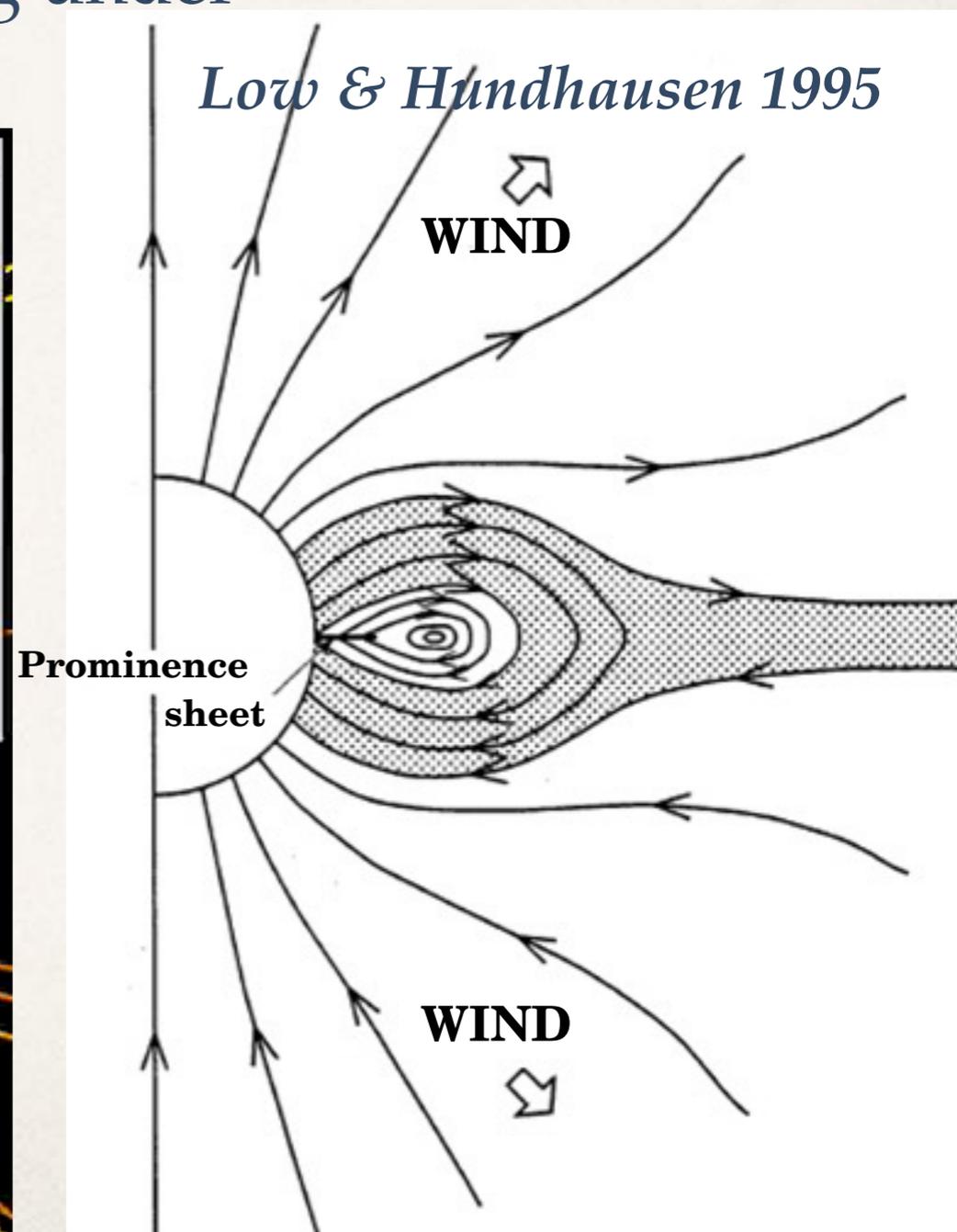
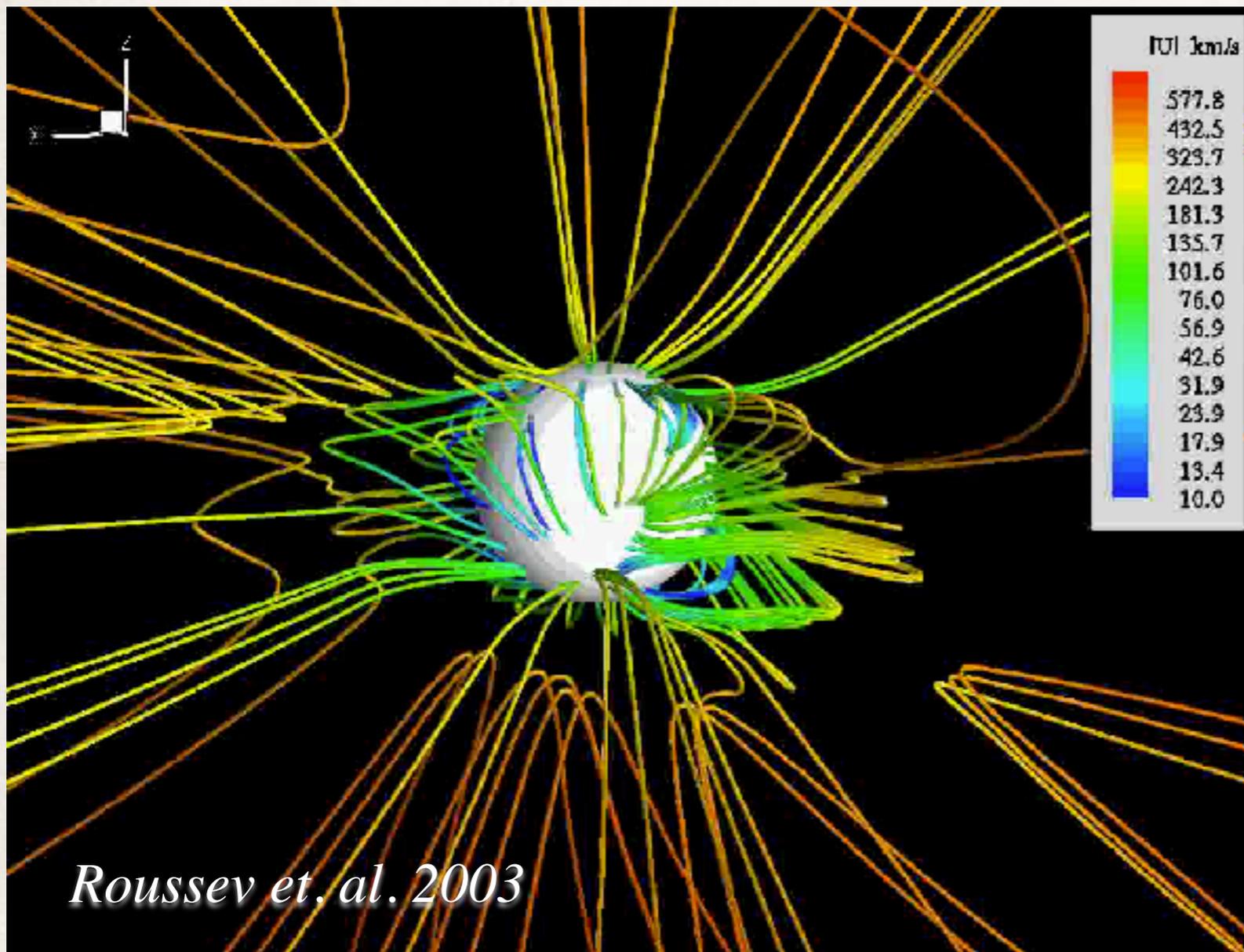


- 512X128X192 uniform/non-uniform grid, Initial isothermal stratified atmosphere at 1MK
- Forced kinematically at lower boundary by an EMF ( $= \mathbf{v} \times \mathbf{B}$ ) corresponding to **sustained** emergence of **twisted** torus into an ambient arcade
- Arcade field at lower boundary: 20G;  
Flux rope axis: 50 G
- **Thermal conduction along field lines**
- **Semi relativistic Boris correction (Boris 1970)**



# CMEs under helmet streamers

Use of arcade magnetic field is inspired by observations of CMEs forming under helmet streamers



# Compressible MHD equations



Mass conservation  $\frac{D \ln \rho}{Dt} = -\vec{\nabla} \cdot \mathbf{U}$

Navier-Stokes with semi-relativistic correction  $\frac{D \mathbf{U}}{Dt} = -\frac{\vec{\nabla} p}{\rho} + g_z \hat{z} + \frac{\mathbf{J} \times \mathbf{B}}{\rho} + \vec{F}_L^{\text{corr}} + \rho^{-1} \mathbf{F}_{\text{visc}},$

Induction  $\frac{\partial \mathbf{A}}{\partial t} = \mathbf{U} \times \mathbf{B} - \eta \mu_0 \mathbf{J}$

Temperature `ENTROPY` = temperature\_idealgas

$$\rho c_v T \frac{D \ln T}{Dt} = -(\gamma - 1) \rho c_v T \vec{\nabla} \cdot \mathbf{U} + \vec{\nabla} \cdot \vec{q}_{\text{cond}} + \vec{\nabla} \cdot (\rho T \chi_t \vec{\nabla} \ln T) + \eta \mu_0 J^2 + 2\rho \nu \mathbf{S}_{ij}^2 + \rho \zeta_{\text{shock}} (\nabla \cdot \mathbf{U})^2 - \rho^2 \Lambda(T)$$

Field aligned Spitzer conductivity

`HEATFLUX` = heatflux  $\vec{q}_{\text{cond}} = K_{\text{sp}} T^{5/2} \hat{b} (\hat{b} \cdot \nabla T),$



# ASIDE ON BORIS CORRECTION

The need for having a semi-relativistic correction to the Lorentz force term in the velocity equation comes from the fact that in non-relativistic plasmas, the Alfvén velocity can become comparable to the speed of light,  $c$ . Non-trivial displacement current. Circumvents the time step constraint due to large Alfvén speed.

*Gombosi et. al. 2002*

$$\left[ \mathbf{I} + \frac{v_A^2}{c^2} (\mathbf{I} - \hat{\mathbf{b}}\hat{\mathbf{b}}) \right] \frac{\partial \mathbf{U}}{\partial t} = - (\mathbf{U} \cdot \nabla) \mathbf{U} - \frac{\nabla p}{\rho} + g_z \hat{\mathbf{z}} + \left[ \frac{(\nabla \times \mathbf{B})}{\mu_0 \rho} + \frac{(\nabla \times \mathbf{E}) \times \mathbf{U}}{\mu_0 \rho c^2} \right] \times \mathbf{B}.$$

Enhanced inertia matrix

```
&magnetic_run_pars
  lweyl_gauge=T
  lbb_as_aux=F
  lboris_correction=T,
  cmin=5.0 ! [5.0 Mm/s]
/
```

$$F_L^{\text{corr}} = \frac{\beta_A^2}{1 + \beta_A^2} \left[ \mathbf{I} - \frac{\hat{\mathbf{b}}\hat{\mathbf{b}}}{1 + \beta_A^2} \right] \left[ \mathbf{U} \cdot \nabla \mathbf{U} + \frac{\nabla p}{\rho} - g_z \hat{\mathbf{z}} - \frac{(\nabla \times \mathbf{B}) \times \mathbf{B}}{\mu_0 \rho} \right].$$

```
p%clight2=spread(max(cmin**2, 25*maxval(p%u2), maxval(p%cs2)), 1, nx)
```

Chatterjee, 2019, GAFD, Pencil code special issue

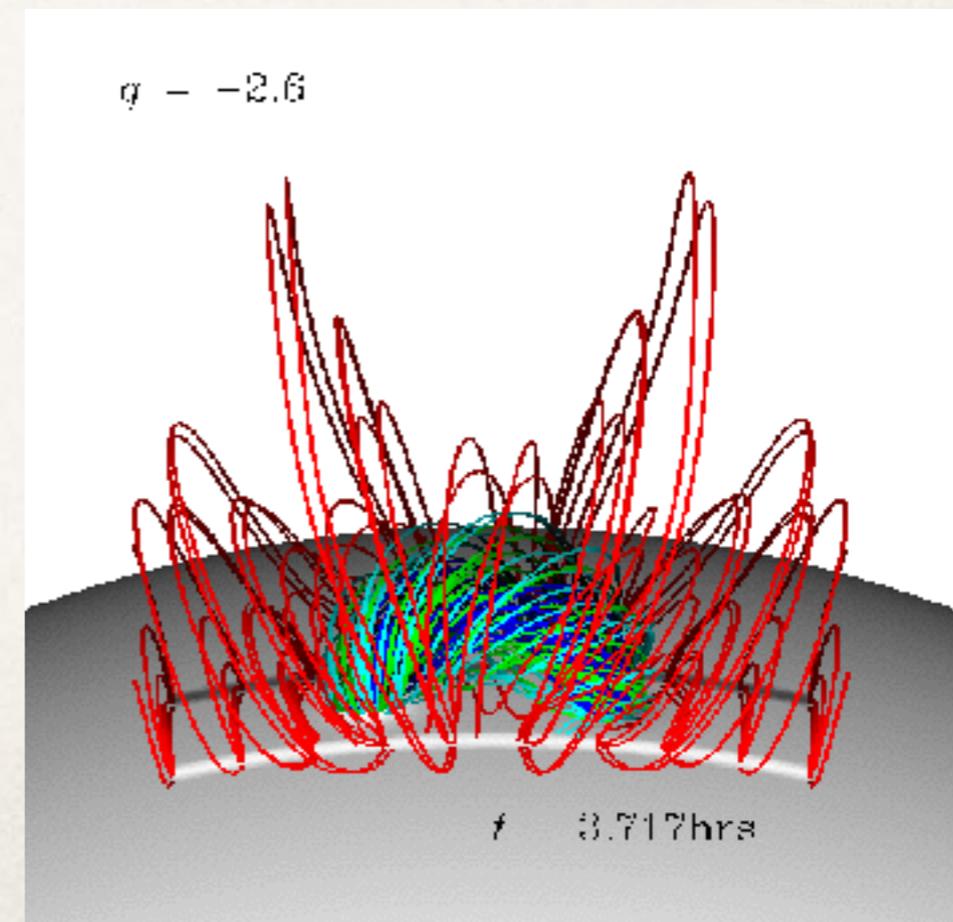


# Storage and release model of CMEs

Slow Emergence of a *twisted* flux rope similar to continually heating a pressure cooker

Valve modelled by ambient coronal magnetic field

Vapor pressure of steam build up of magnetic helicity in corona

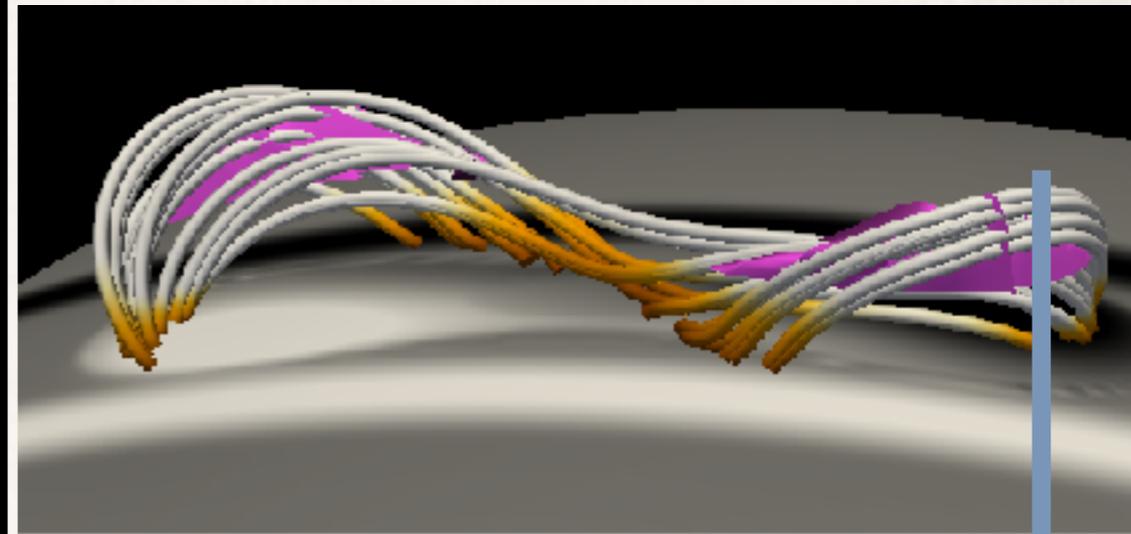


# Tether cutting and formation of soft X-ray sigmoid



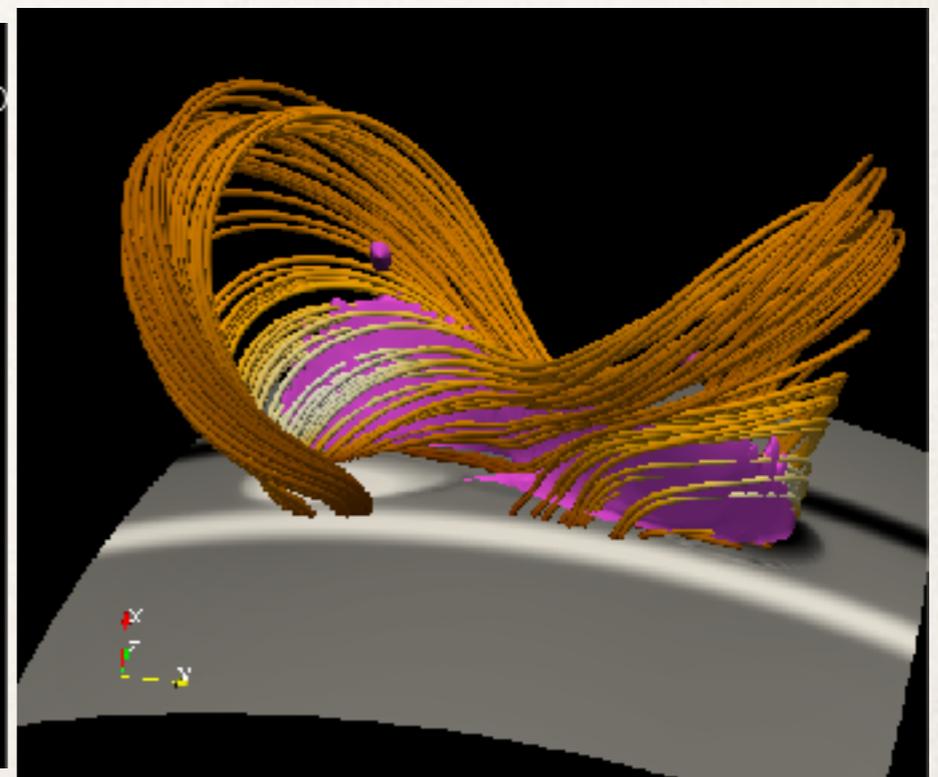
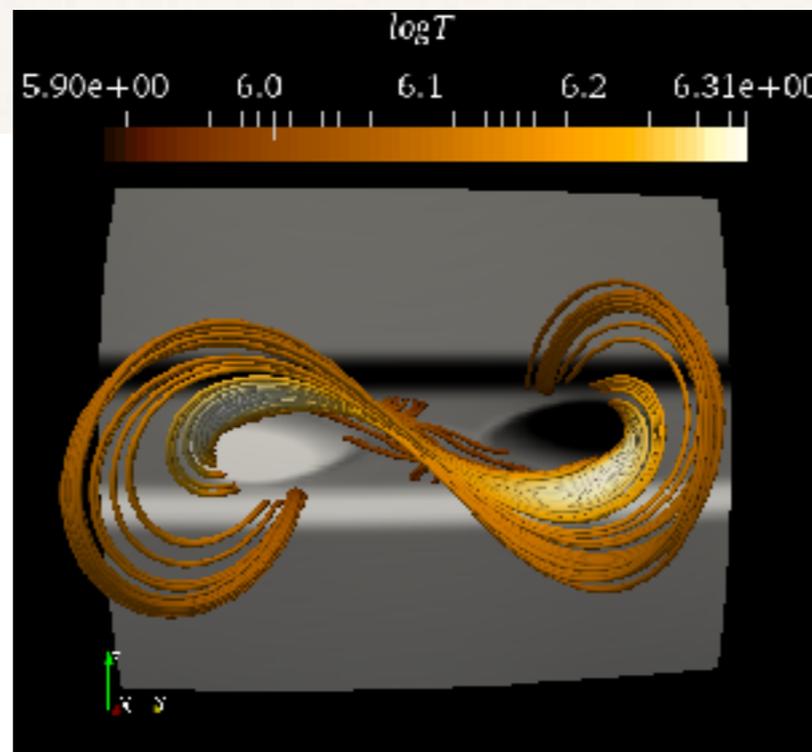
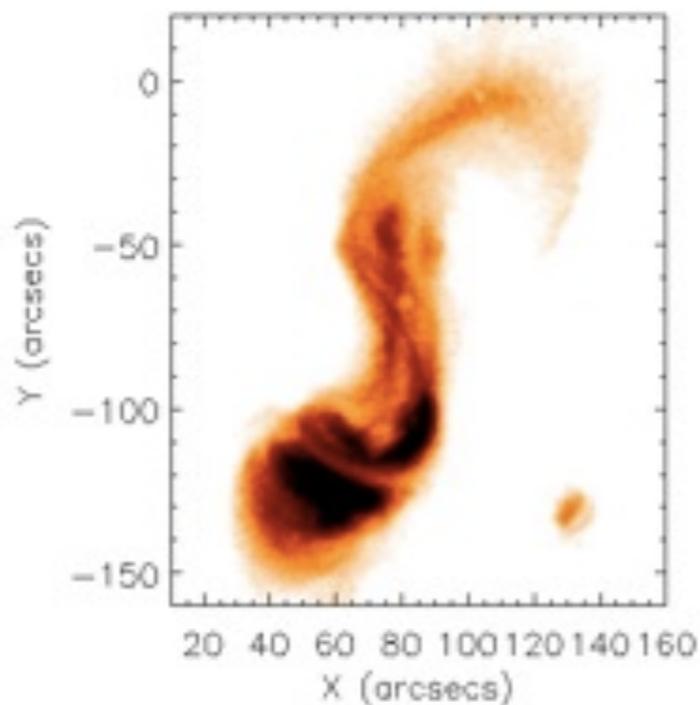
## Slow Build-up phase

Hot current channel forms and field lines passing through the hot channel are sigmoidal in shape and show up in soft X-rays



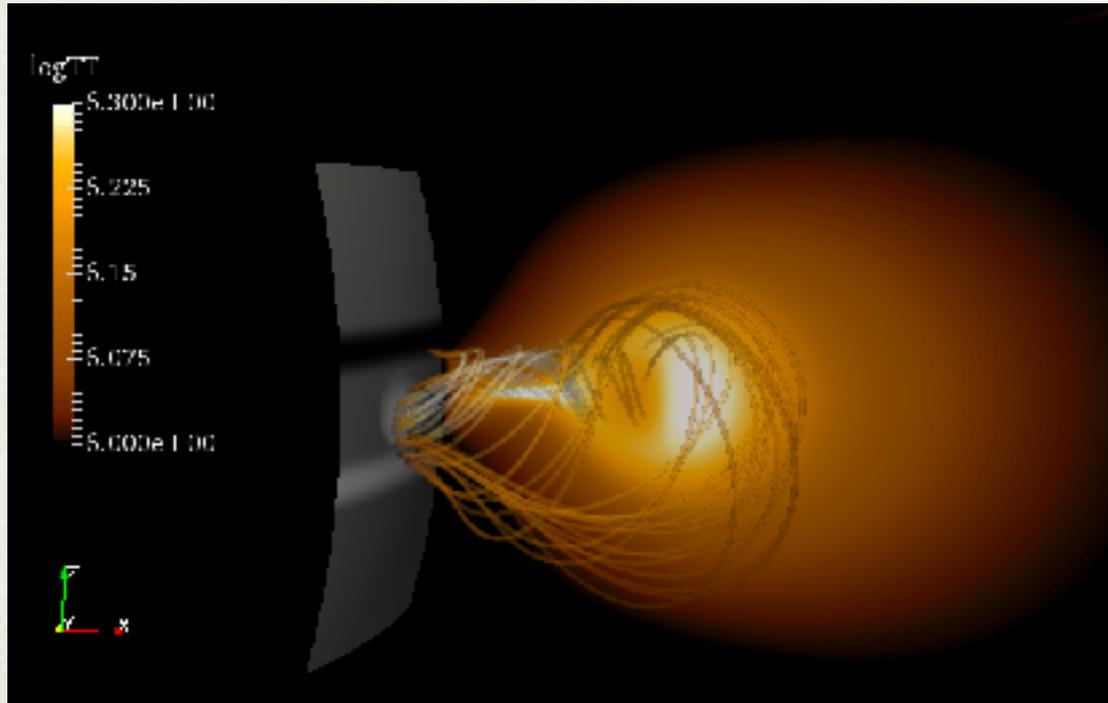
Time

## Hinode

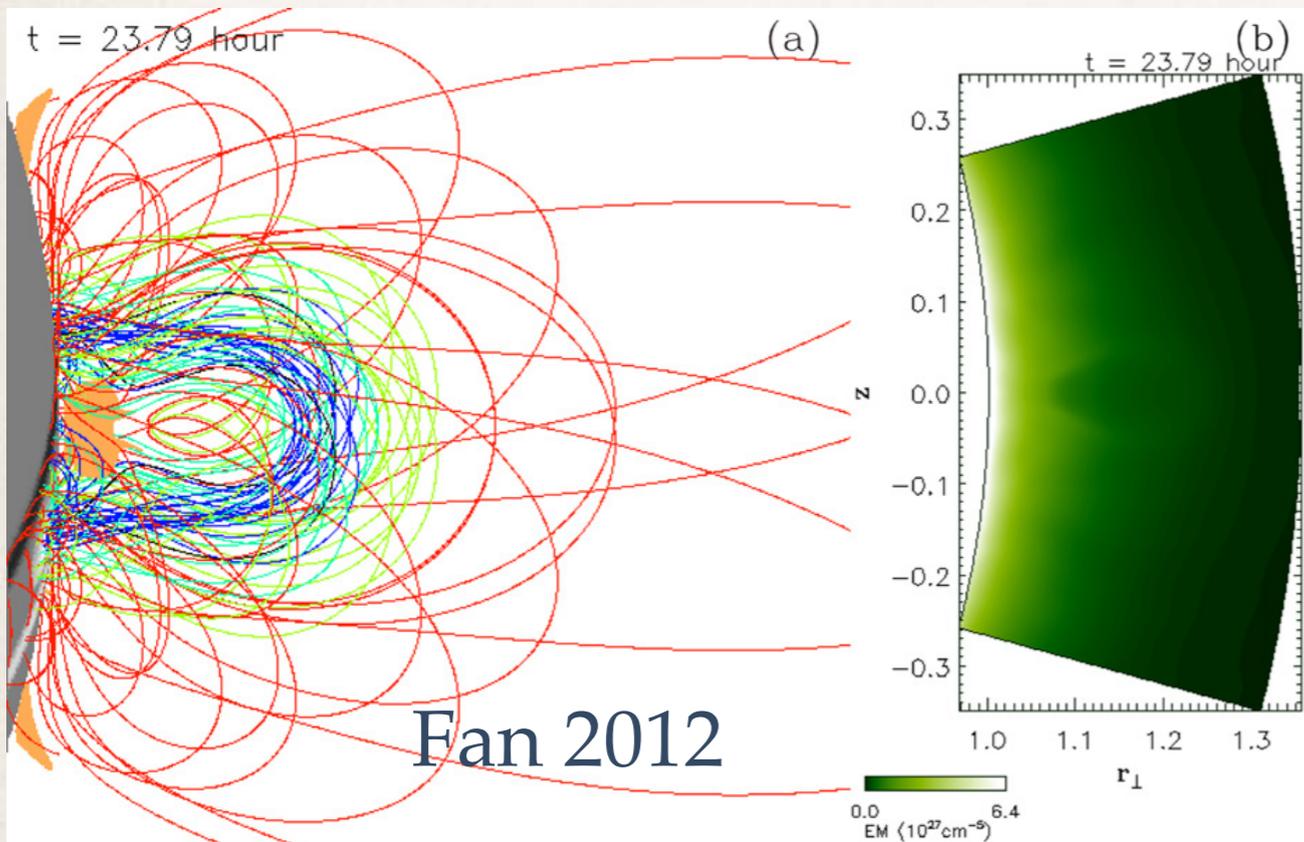
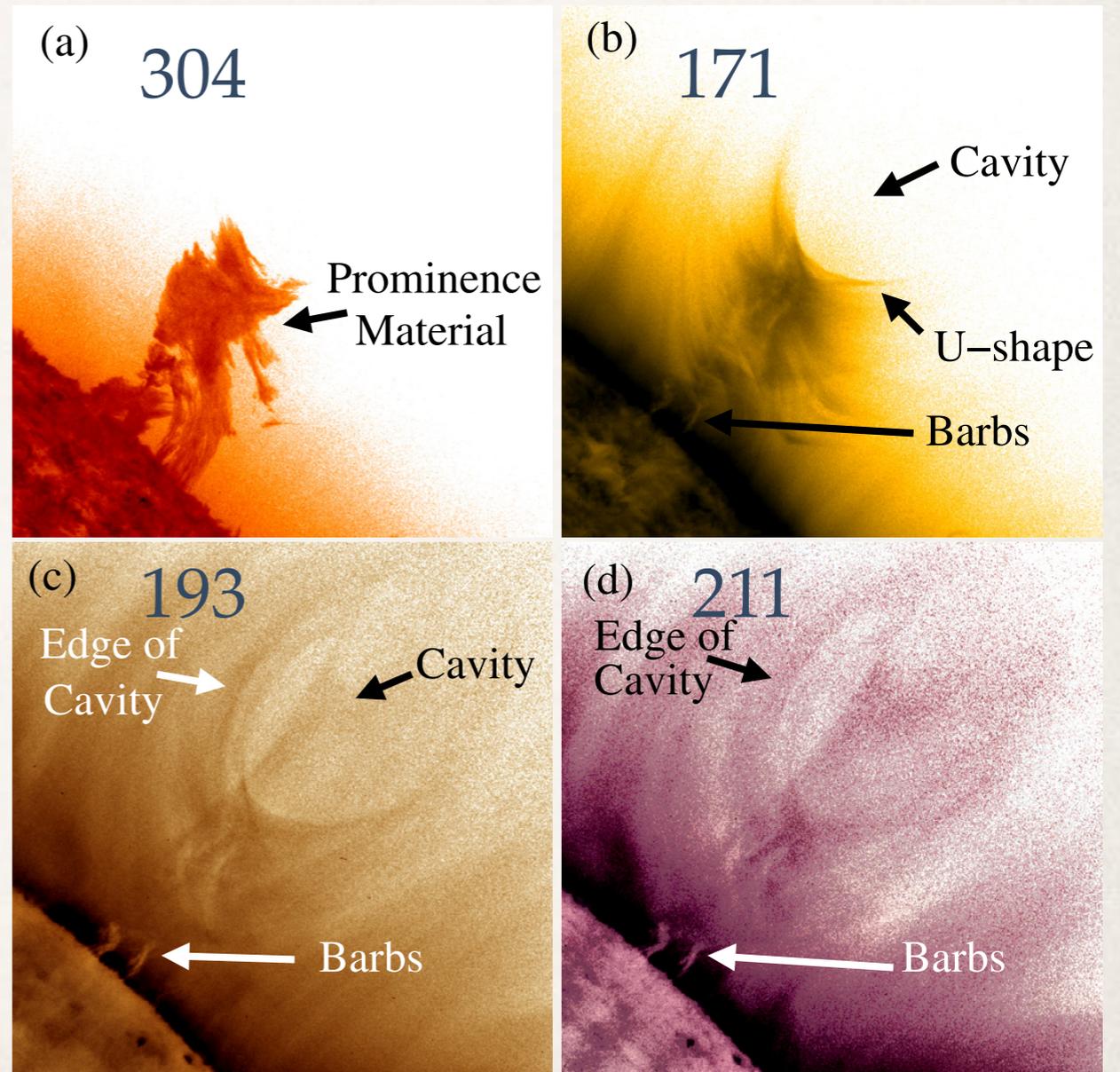




# Pre-eruption coronal cavity: *Lollipop on a stick!*

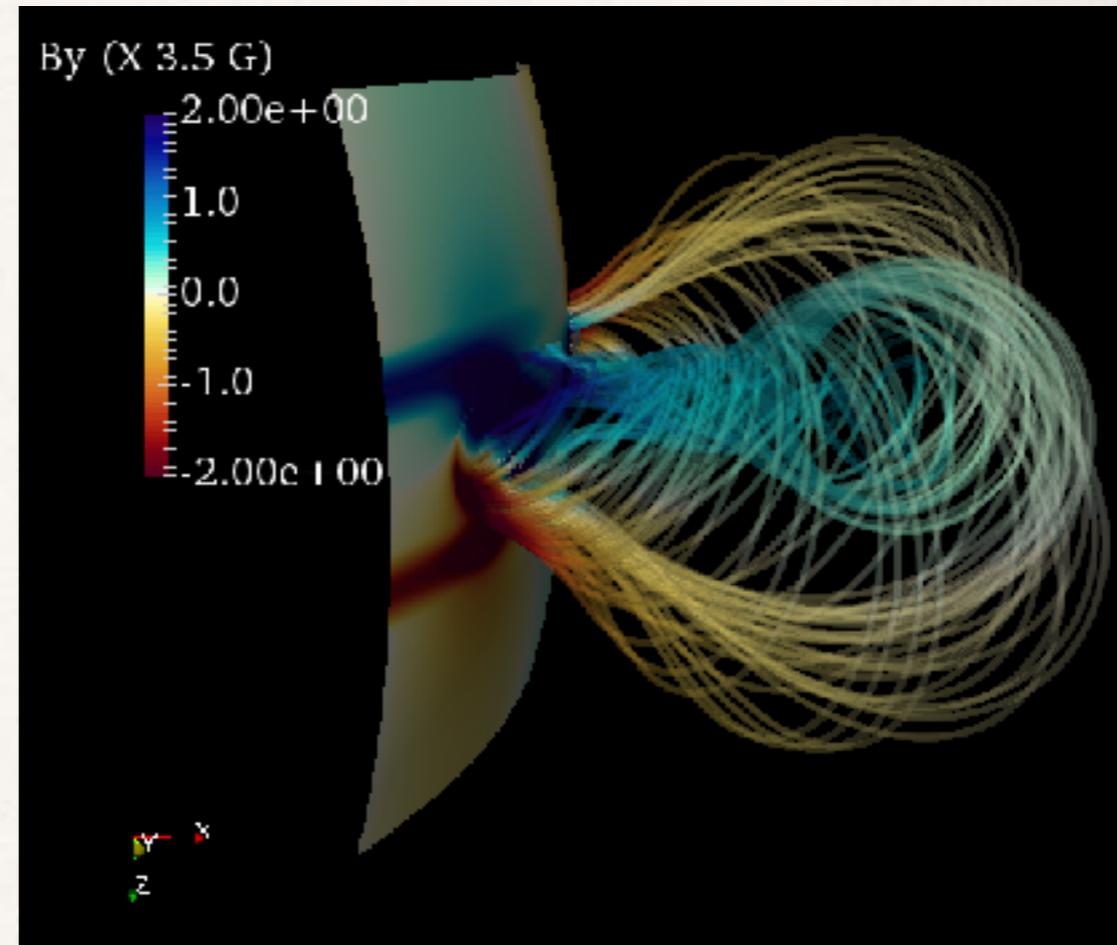
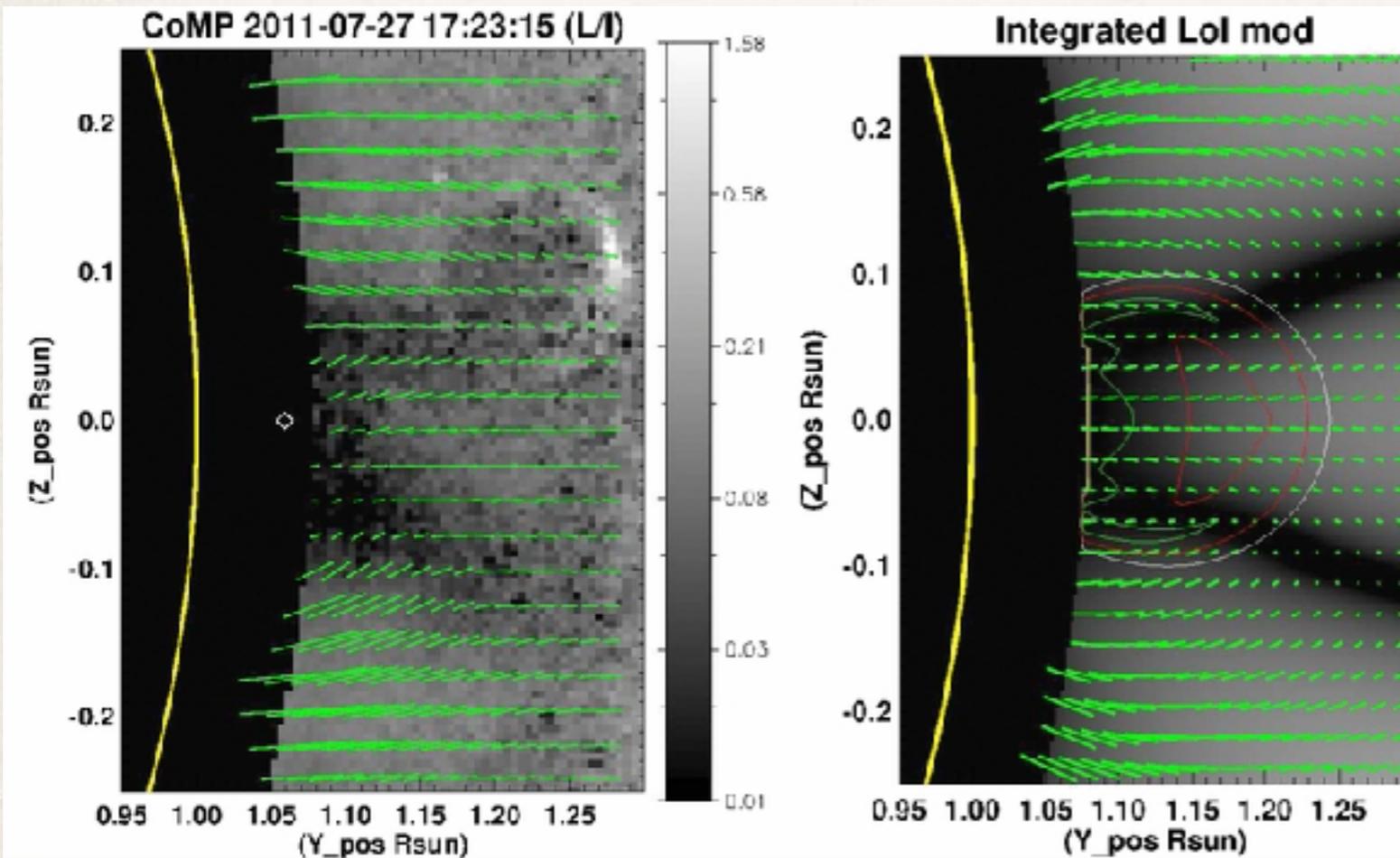


Prominence cavity SDO/AIA [negative images] on 13 June 2010 on NW limb (Regnier et. al. 2011)





# Pre-eruption coronal cavity: *Bunny ears*



Null patterns in Stokes L/I in a Polar crown cavity  
(Bak-Stešlicka et. al. 2013)

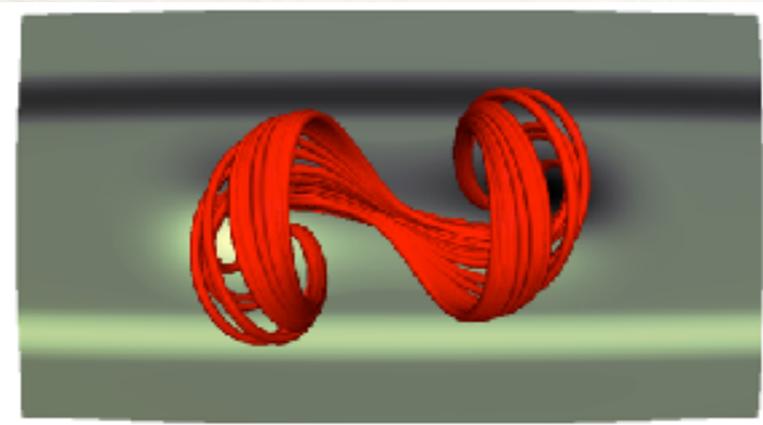
$$L = \sqrt{Q^2 + U^2} \propto \sin^2 \theta$$



# Supporting a prominence

1. **Mass supply:** pressure gradient due to formation of hot current channel

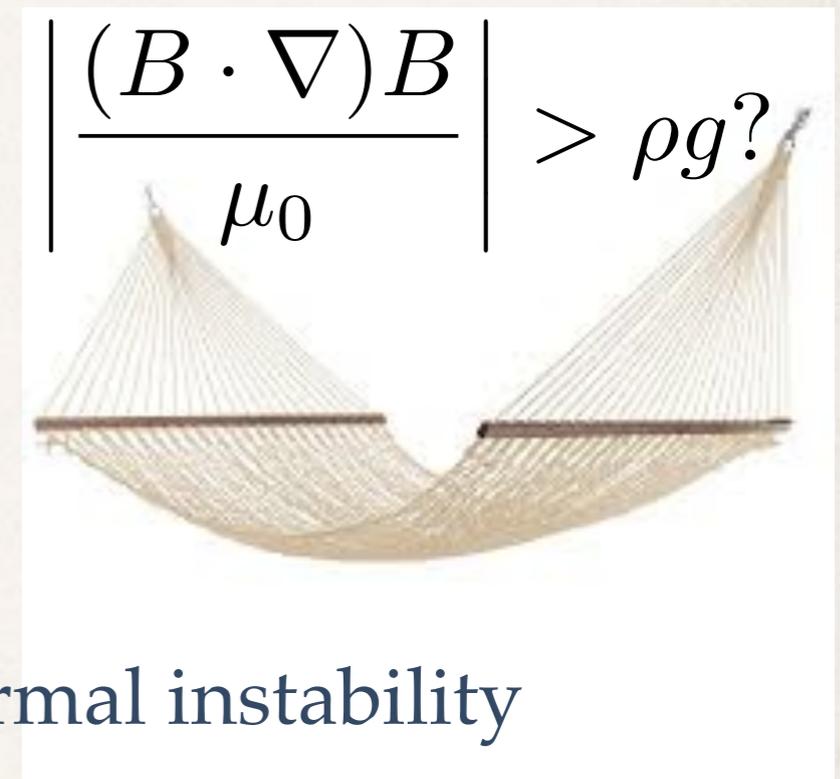
$$\mathbf{q}_s = -\kappa_0 T^{5/2} \hat{\mathbf{b}} \hat{\mathbf{b}} \cdot \nabla T,$$



Increase in Gas pressure at Chromosphere



Plasma sucked out of chromosphere



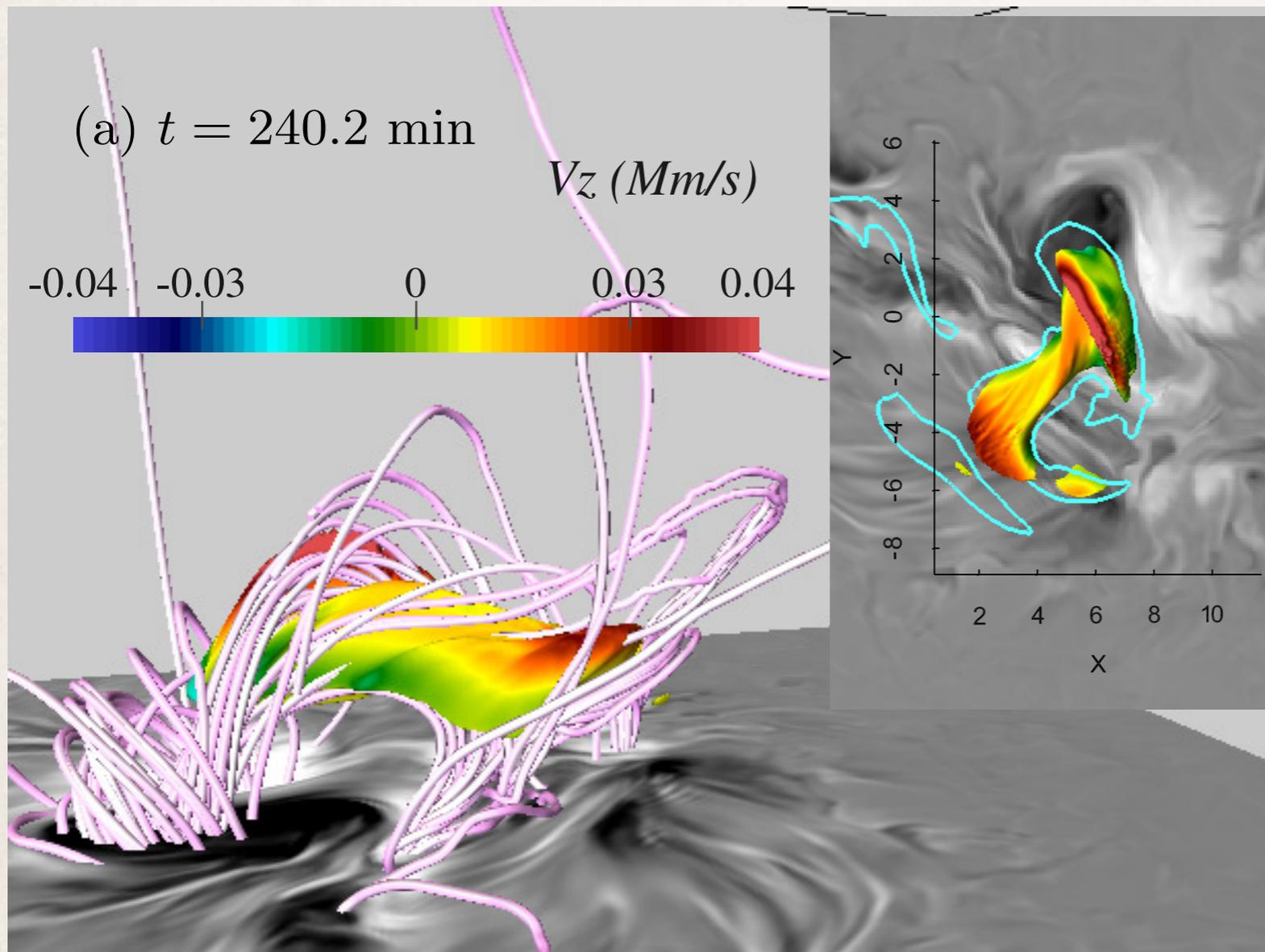
2. **Formation:** Prominence condensation due to thermal instability

Parker (1953), Field (1965)

$$\frac{\partial e}{\partial t} + \nabla \cdot (e\vec{u}) = \text{Heating} - \rho^2 \Lambda(T) + \text{Others}$$

Condensed plasma falls into preexisting dips

# Supporting a prominence



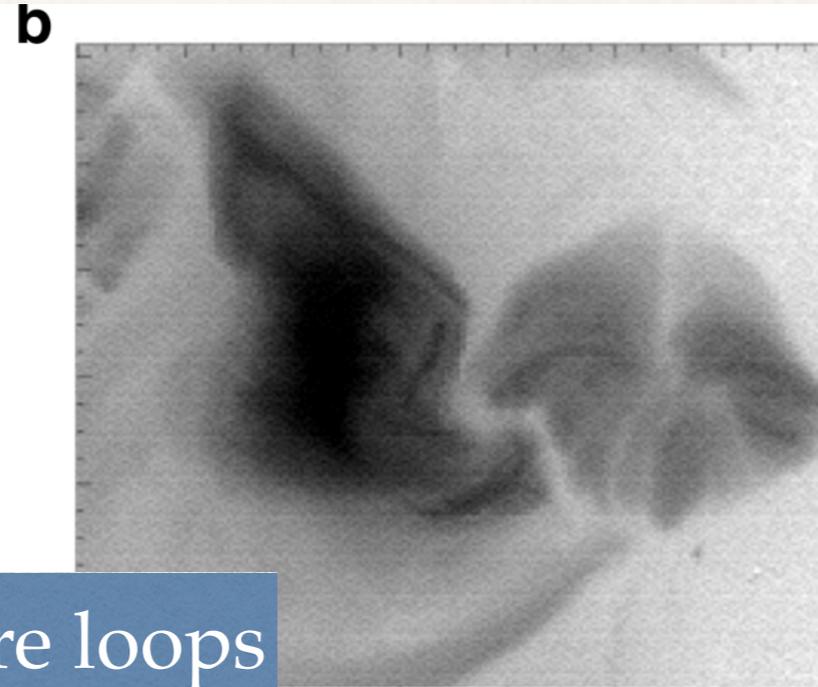
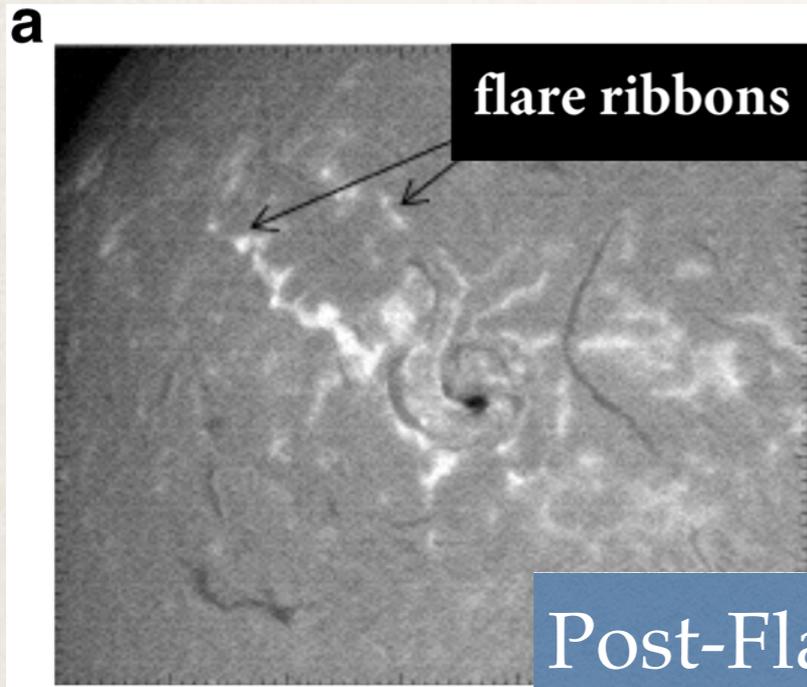
**2. Formation:** Field lines dip under the weight of the plasma. Depends on magnetic tension vs gravity

Estimated mass inside the iso-surface  $\sim 1.2 \times 10^{13}$  kg  
*PC et. al. 2016*

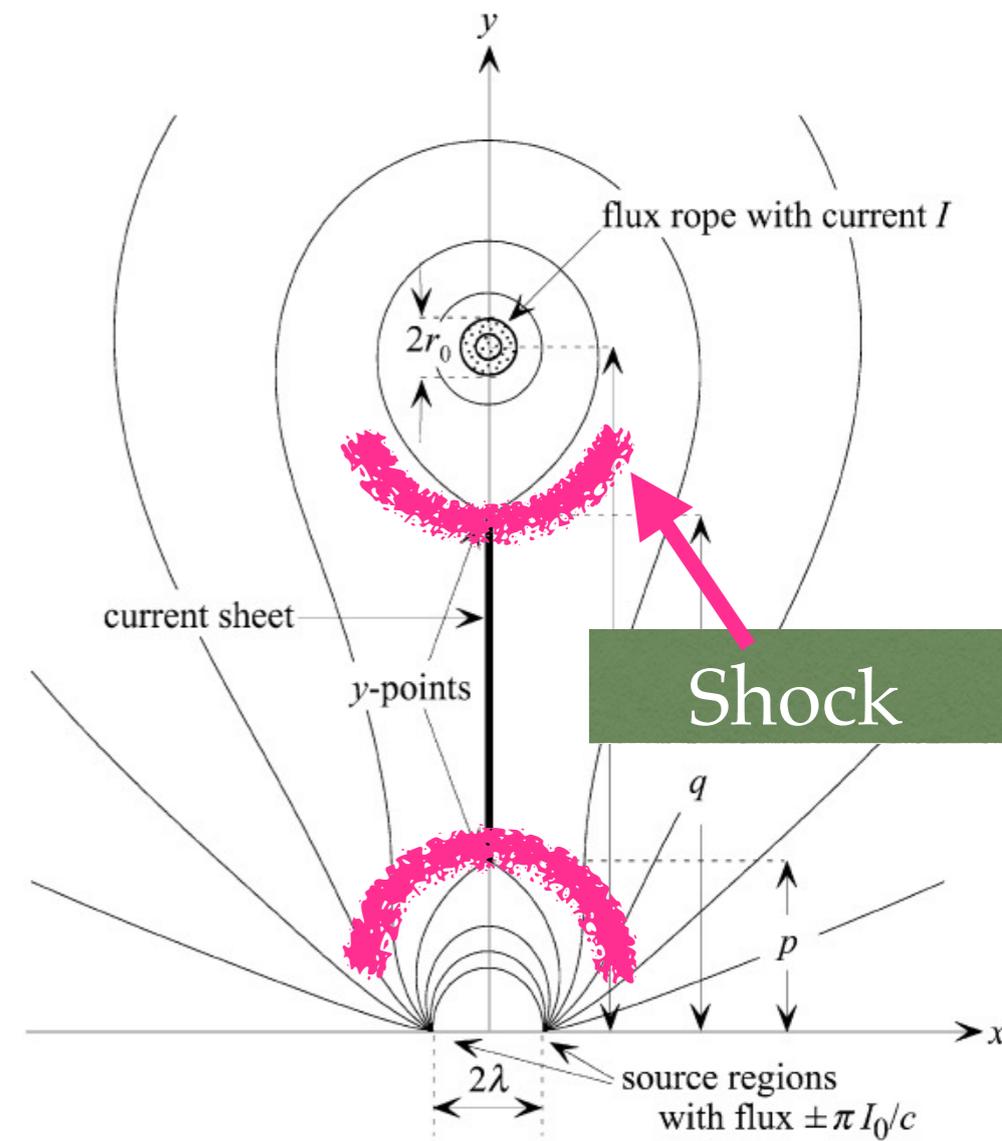
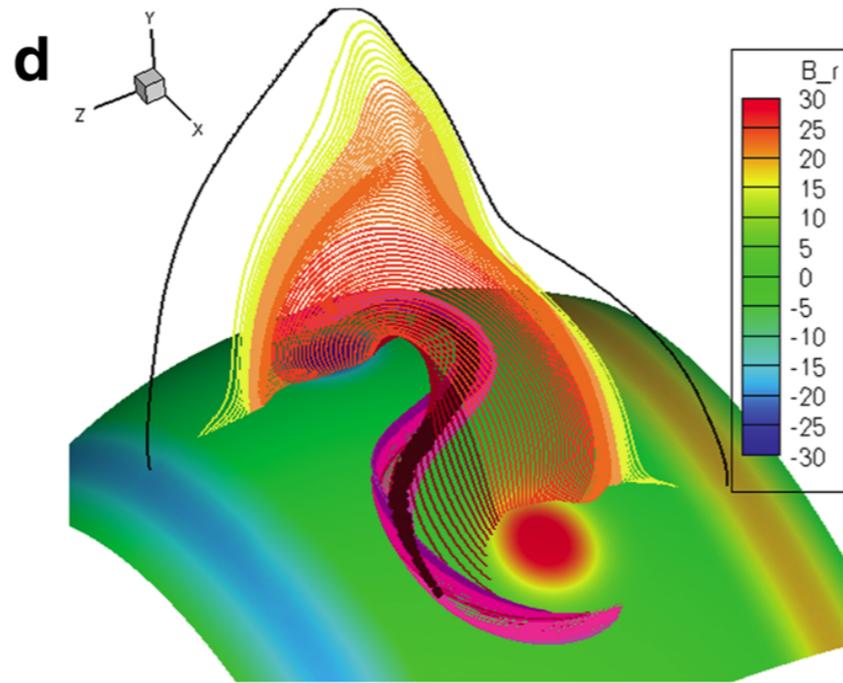
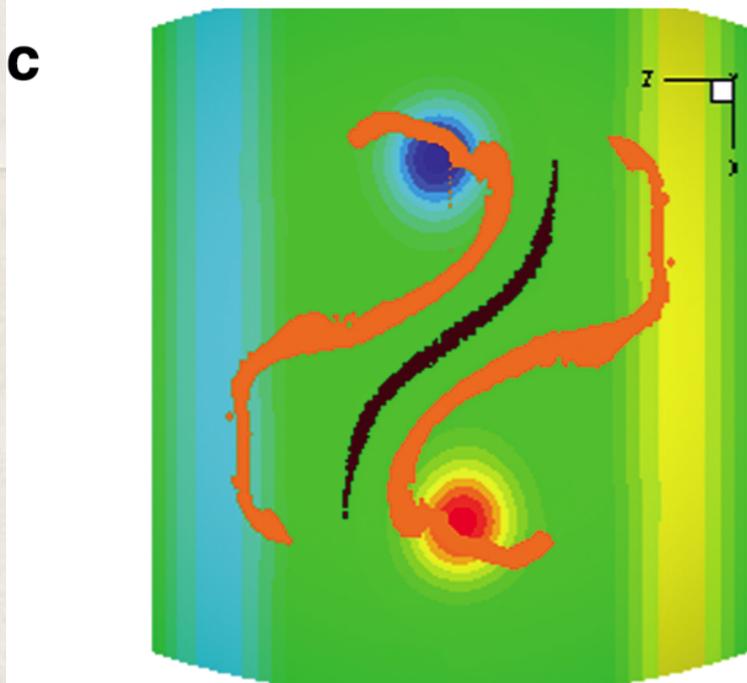
Typical Observed CME mass  $\sim 10^{11}$  to  $4 \times 10^{13}$  kg

**3. Destruction:** Evaporation, Rayleigh Taylor Instability, Coronal Rain

# Tether cutting: Fast Release phase



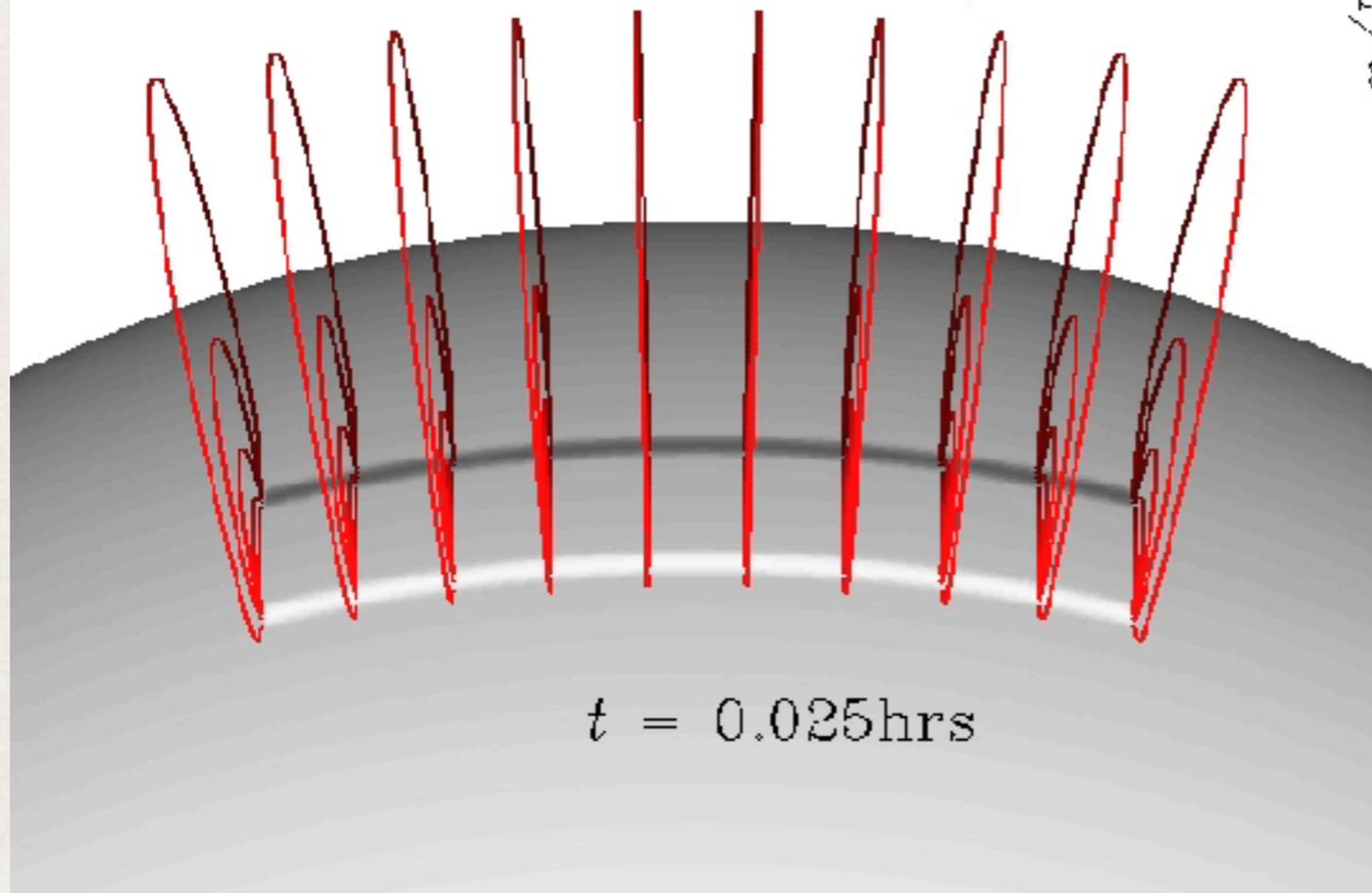
Two-Ribbon flare with foot points moving apart with time  
*Gibson, 2018 (LRSP)*



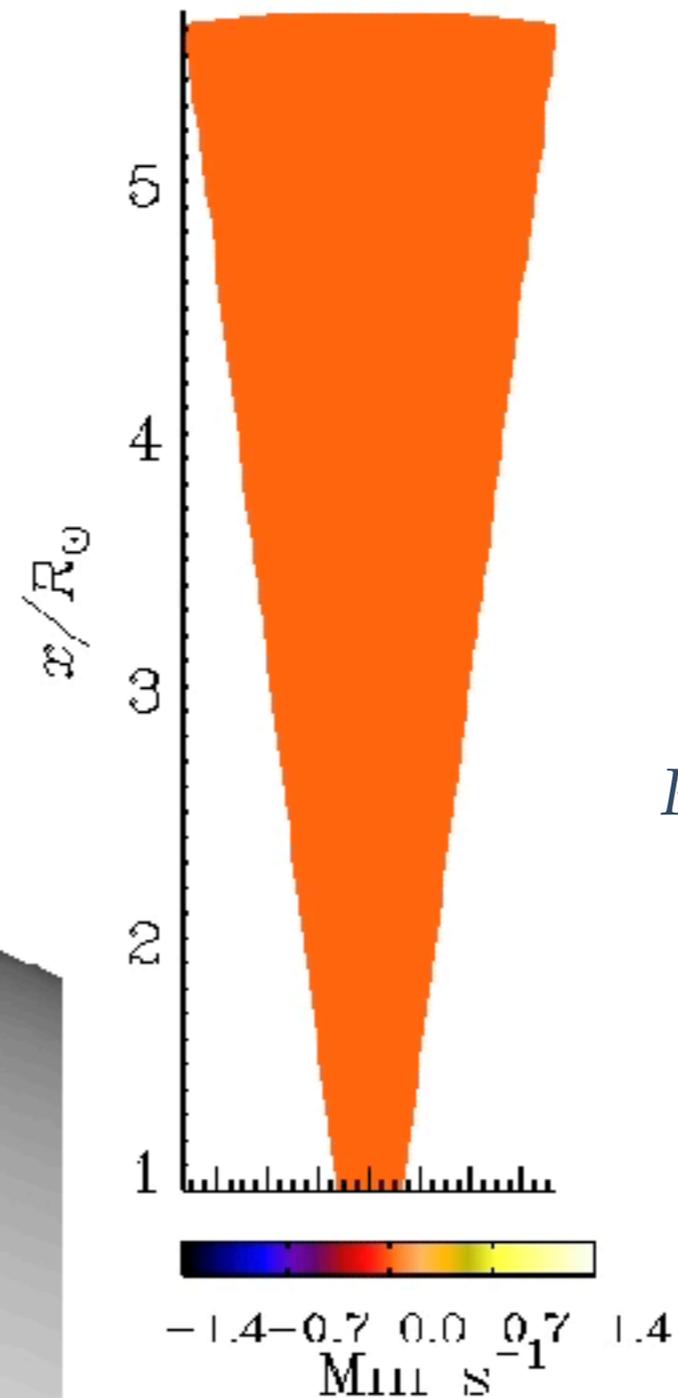
# Homologous eruptions: Full MHD simulation



1. Sequence of **four homologous eruptions initiated** due to helical kink instability.
2. Flux rope **reforms** after every eruption
3. **Second** in the series is **cannibalistic**



Radial Velocity



*PC & Fan 2013*



# Height time curves for erupting flux ropes

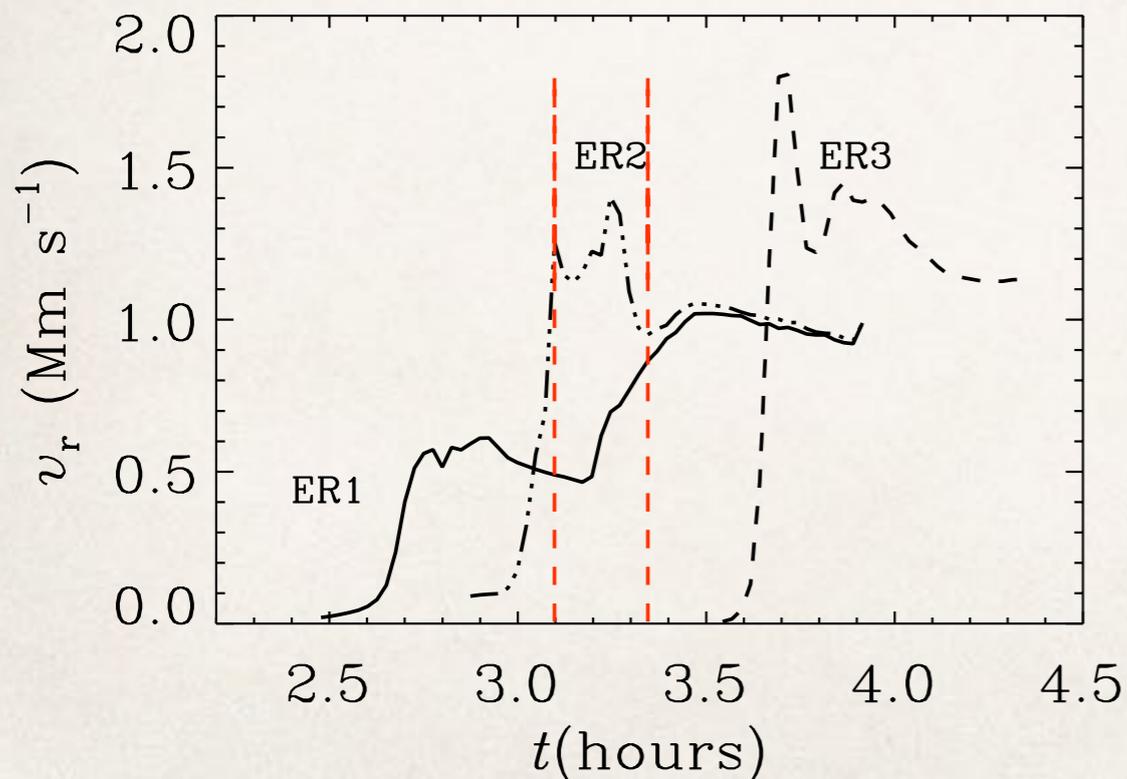
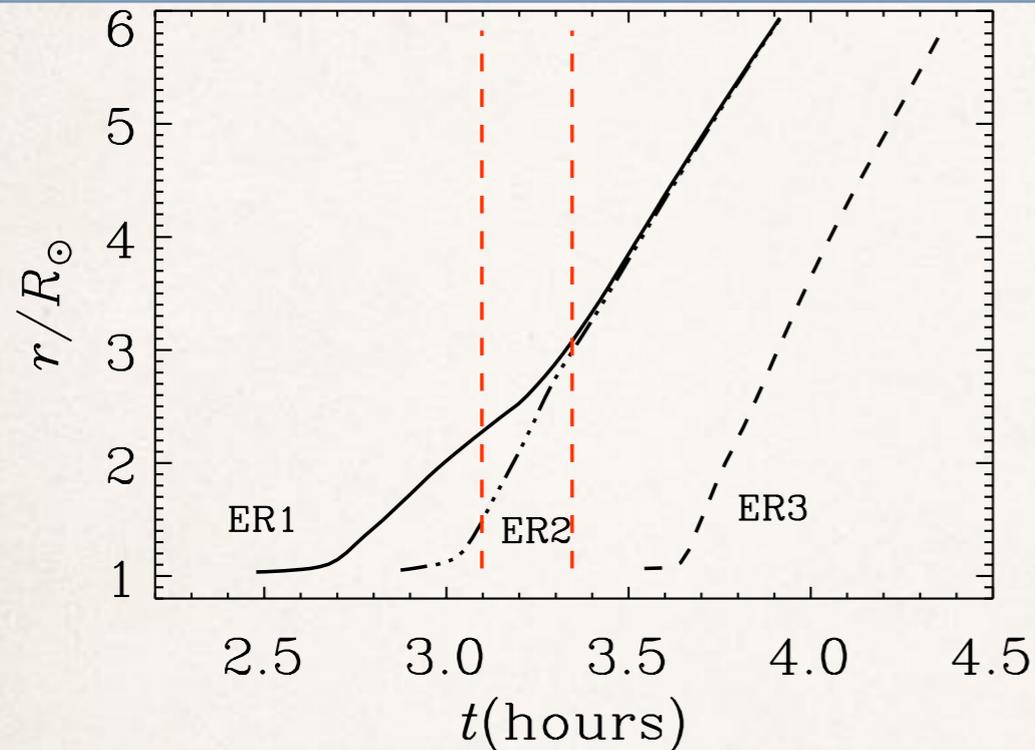
PC & Fan 2013

CME speeds :

650 km/s, 1400 km/s, 1800 km/s

Slowest CME acceleration  $\sim 1.6 \text{ km s}^{-2}$

Height at which CME acceleration initiated  $\sim 1.035 R_s$

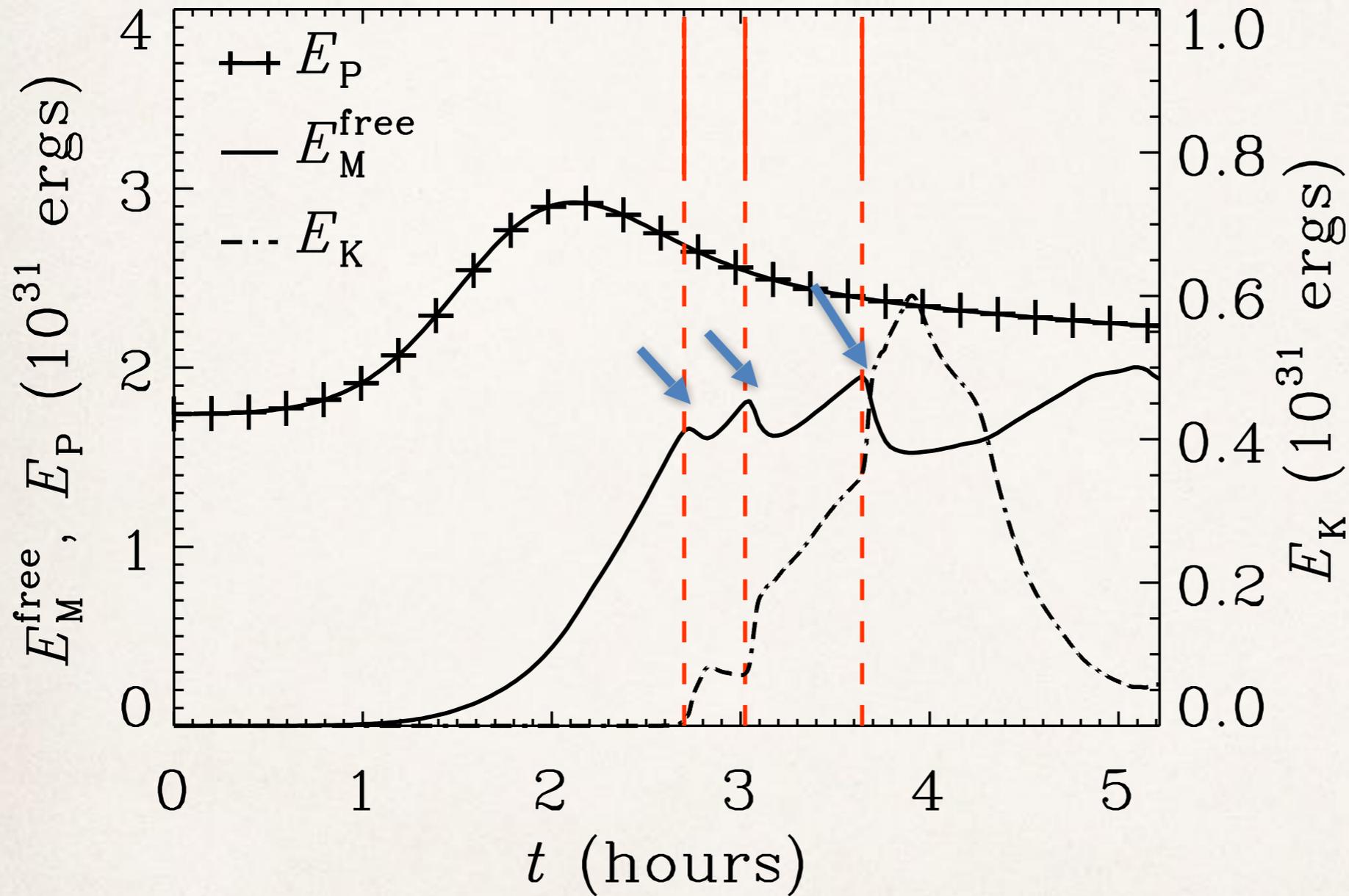


$$-\frac{d \ln B_{\text{pot}}}{d \ln r} \approx 1.0, 1.05, 1.1 < n_{\text{crit}} = 1.5$$

At the point of initiation all CMEs stable to TI

*The twist of the field lines near the axis  
 $\sim 2.1$  winds between the anchored footpoints.*

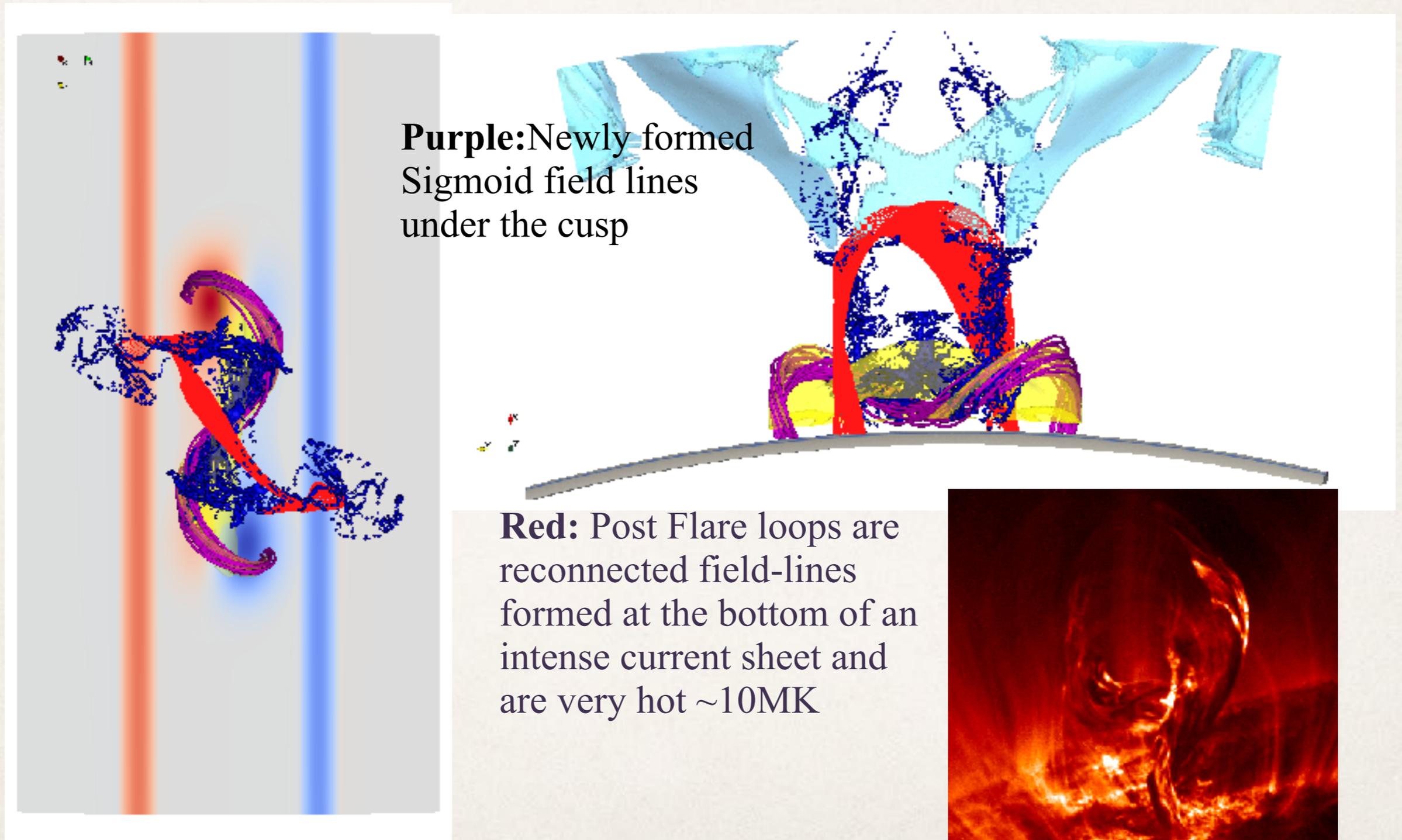
# Energetics of homologous eruptions



1. More magnetic free energy (MFE) available for conversion to kinetic energy (KE) for successive CMEs

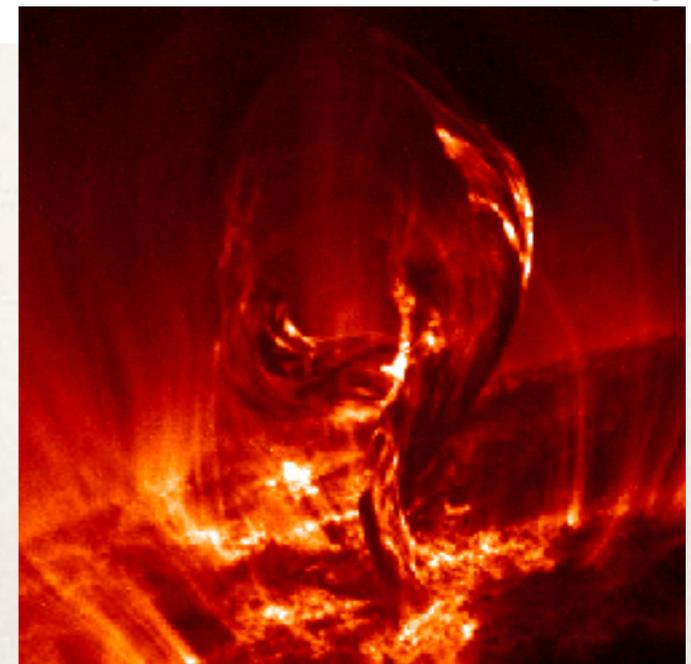
2. Each CME releases the free energy only partially

# Partial eruption: sigmoid-under-cusp

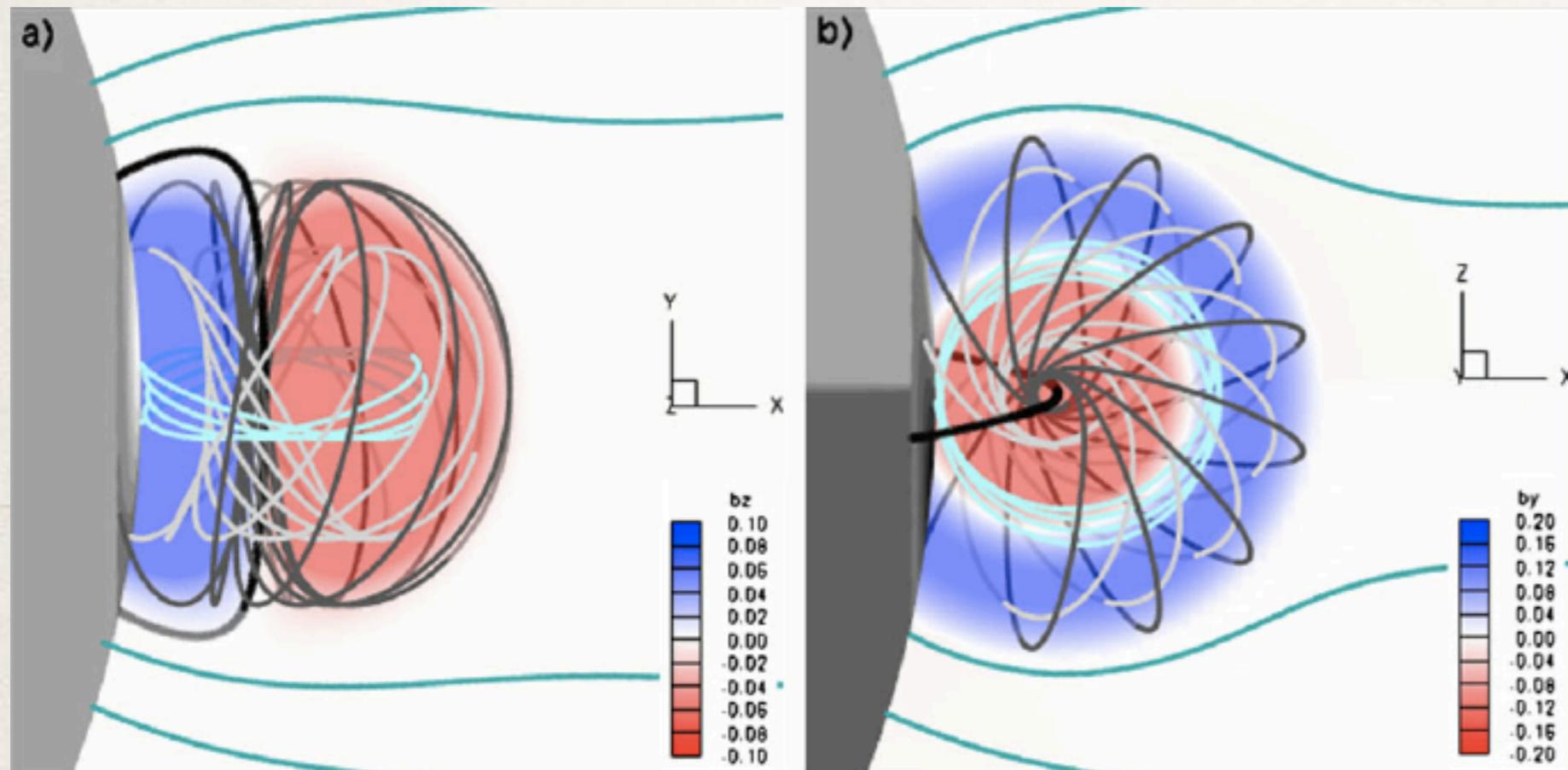


**Purple:** Newly formed Sigmoid field lines under the cusp

**Red:** Post Flare loops are reconnected field-lines formed at the bottom of an intense current sheet and are very hot  $\sim 10\text{MK}$



# Structure of CME far from the sun



Gibson & Low (1998)  
Flux rope

Spheromak

Chatterjee & Fan 2013 model

