

# SUSY enhancement on $S_b^3$

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# A puzzle in ABJ(M) on $S_b^3$

ABJ(M) theory from an  $\mathcal{N} = 2$  perspective:

$U(N)_k \times U(N+M)_{-k}$  gauge theory + 4 bi-fundamental chirals  
and  $U(1)_b \times SU(2)_R \times SU(2)_R$  flavor symmetry



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Can be put on the squashed sphere  $S_b^3$  while preserving  $\mathcal{N} = 2$  SUSY.  
Admits 3 mass parameters  $m_1, m_2 + m_3, m_2 - m_3$ .

Partition function can be computed using localization, complicated  $b$  dependence in general.



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Partition function can be computed using localization, complicated  $b$  dependence in general. However [Chester, Kalloor, Sharon '21] found:

$$\begin{aligned} \mathcal{Z} \left( b; m_1 = \frac{bm_+ - b^{-1}m_-}{2}, m_2 = \frac{bm_+ + b^{-1}m_-}{2}, m_3 = i \frac{b - b^{-1}}{2} \right) \\ = \mathcal{Z} \left( 1; m_1 = \frac{m_+ - m_-}{2}, m_2 = \frac{m_+ + m_-}{2}, m_3 = 0 \right) \end{aligned}$$



# Outline

- $\mathcal{N} = 2$  supersymmetry on the squashed three-sphere
- Enhancement to  $\mathcal{N} = 4$  supersymmetry
- Applications



# The squashed three-sphere

In this talk  $S_b^3$  is given by the metric

$$ds^2 = \frac{r_3^2}{4} \left( \frac{b + b^{-1}}{2} \right)^2 (d\psi + \cos\theta d\phi)^2 + \frac{r_3^2}{4} (d\theta^2 + \sin^2\theta d\phi^2)$$

→ we stretch the Hopf fiber by a factor  $\frac{b+b^{-1}}{2} \geq 1$ .



$\mathcal{N} = 2$  supersymmetry on  $S_b^3$ 

To preserve SUSY: embed the metric into a background SUGRA multiplet  
 $\Rightarrow$  Killing spinor equations (KSE)

In 3d  $\mathcal{N} = 2$  (off-shell) conformal SUGRA [Nishimura, Tani '13]:

$$\nabla_\mu \zeta - iA_\mu^{12} \zeta = \gamma_\mu \eta,$$

$$\nabla_\mu \tilde{\zeta} + iA_\mu^{12} \tilde{\zeta} = \gamma_\mu \tilde{\eta}.$$

(In the following I will not write the equations for the  $\tilde{\zeta}$ .)

Correct choice of  $SO(2)_R$ -gauge field  $A^{12} \Rightarrow$  4 Killing spinors preserved.

[Imamura, Yokoyama '11]:  $b$ -dependent partition function from localization.



# Mass-deformed $\mathcal{N} = 2$

Supersymmetric mass: part of a background vector multiplet  $(A_\mu, m, D)$ :

$$\frac{1}{4}\gamma^{\mu\nu}F_{\mu\nu}\zeta - i(D\zeta - m\eta) = 0 \quad \& \text{ corresp. } \tilde{\zeta} \text{ equation.}$$

Fixes  $A, D$  in terms of  $m$  and the metric background.





# Embedding into $\mathcal{N} = 4$ conformal SUGRA

Assume the mass couples to a  $U(1)$  of the R-symmetry:  $A_\mu \rightarrow A_\mu^{34}$ .

Then  $A^{12}, A^{34}$  are components of a background  $SO(4)_R$  gauge field  $A^{IJ}$   
 $\Rightarrow$  all background fields  $g, A^{IJ}, m, D$  part of an  $\mathcal{N} = 4$  Weyl multiplet.



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$\mathcal{N} = 4$  conformal Killing spinor equations [Nishimura, Tanii '13]:

$$\begin{aligned} \nabla_\mu \zeta - iA_\mu^{12} \zeta &= \gamma_\mu \eta, \\ \nabla_\mu \zeta' - iA_\mu^{34} \zeta' &= \gamma_\mu \eta', \\ \frac{1}{4} \gamma^{\mu\nu} F_{\mu\nu}^{34} \zeta - i(D\zeta - m\eta) &= 0, \\ \frac{1}{4} \gamma^{\mu\nu} F_{\mu\nu}^{12} \zeta' - i(D\zeta' - m\eta') &= 0, \end{aligned}$$

& 4 corresponding equations for  $\tilde{\zeta}, \tilde{\zeta}'$ .

Twice as many equations as before: At general  $m$  no non-trivial  $\zeta', \tilde{\zeta}'$ .



$\mathcal{N} = 4$  on the squashed sphere

For  $m = i\frac{b-b^{-1}}{2}$  the  $\mathcal{N} = 4$  KSE admit additionally non-trivial Killing spinors

$$\zeta' = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \tilde{\zeta}' = \begin{pmatrix} 0 \\ 1 \end{pmatrix},$$

i.e. we now have a total of 6 Killing spinors.



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Localization with only these two Killing spinors  $\zeta', \tilde{\zeta}'$

[Hama, Hosomichi, Lee '11]: squashing independent partition function.

$\Rightarrow$  Coupling  $R$ -symmetry to a mass  $m = i\frac{b-b^{-1}}{2}$  allows for SUSY enhancement, the squashing deformation is  $\mathcal{Q}$ -exact under these additional Killing spinors [Closset, Dumitrescu, Festuccia, Komargodski '12, '13]



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Note: Same happens for  $m = -i\frac{b-b^{-1}}{2}$



# Back to ABJ(M) and beyond

For ABJ(M):  $m_3$  coupled to a  $U(1) \subset SU(2)_R \times SU(2)_R \subset SO(6)_R$   
 $\Rightarrow$  ABJ(M) features this SUSY enhancement for  $m_3 = i\frac{b-b^{-1}}{2}$   
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Our conclusions hold for other theories as well.

Example: theories constructed of  $\mathcal{N} = 4$  vector and hypers.

Fine-tuned mass couples to axial  $U(1)$  of the  $SU(2) \times SU(2) = SO(4)_R$ .

We find then

$$\mathcal{Z}_{\mathcal{N}=4}(b; b^{\mp 1}\eta; b^{\pm 1}\vec{m}, m = \pm i \frac{b-b^{-1}}{2}) = \mathcal{Z}_{\mathcal{N}=4}(1; \eta; \vec{m}, 0)$$



Testing  $\mathcal{N} = 4$  dualities on  $S_b^3$ 

Wide variety of IR dualities for  $3d$   $\mathcal{N} = 4$  theories, usually on conformally flat manifolds. E.g. [Kapustin, Willett, Yaakov '10]:

$$\mathcal{Z}_{\text{ABJM}}(1; \zeta; 2\xi, 0) = \mathcal{Z}_{\mathcal{N}=8}(1; \xi + 2\zeta; \xi - 2\zeta, 0).$$

Here  $\mathcal{N} = 8$  SYM:  $\mathcal{N} = 4$  vector + 1 fund. hyper + 1 massive adj. hyper





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Putting things together with our results:

$$\mathcal{Z}_{\text{ABJM}}\left(b; \zeta; 2\xi, i\frac{b - b^{-1}}{2}\right) = \mathcal{Z}_{\mathcal{N}=8}\left(b; \xi + 2\zeta; \xi - 2\zeta, -i\frac{b - b^{-1}}{2}\right)$$



# Conclusion

What I have shown:

- Introducing a fine tuned mass in  $\mathcal{N} = 2$  on  $S_b^3$  leads to supersymmetry enhancement.
- This SUSY enhancement leads to squashing independence of partition functions.
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What I had no time to show:

- We also find an enhancement at  $m = i\frac{b+b^{-1}}{2}$ . Coincides with a simplification of the partition function.
- 4d index perspective.
- THF and the ellipsoid.
- A similar story in 4d [Naseer, Thull '21].



# Outlook

Ongoing work and future directions:

- BPS-Wilson lines.
- Mirror symmetry on  $S_b^3$ .
- Brane constructions for the dualities on  $S_b^3$ .
- Seiberg-like dualities on  $S_b^3$  ?
- What happens if we go away from this special point? Do the dualities break?
- Understand the  $m = i \frac{b+b^{-1}}{2}$  theories.

