

Quarter-BPS operators in $\mathcal{N} = 4$ SYM: a case study

30th Nordic Network Meeting on “Strings, Fields and Branes”

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Based on 2111.06857

[Bissi, GF, Manenti]

Introduction: Conformal bootstrap

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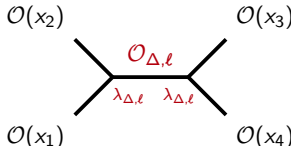
$$\langle \mathcal{O}(x_1) \mathcal{O}(x_2) \mathcal{O}(x_3) \mathcal{O}(x_4) \rangle = \frac{\mathcal{G}(z, \bar{z})}{x_{12}^{2\Delta_{\mathcal{O}}} x_{34}^{2\Delta_{\mathcal{O}}}} = \sum_{\Delta, \ell} \begin{array}{c} \mathcal{O}(x_2) \\ \diagdown \quad \diagup \\ \mathcal{O}_{\Delta, \ell} \\ \diagup \quad \diagdown \\ \mathcal{O}(x_1) \end{array} \begin{array}{c} \mathcal{O}(x_3) \\ \diagdown \quad \diagup \\ \lambda_{\Delta, \ell} \quad \lambda_{\Delta, \ell} \\ \diagup \quad \diagdown \\ \mathcal{O}(x_4) \end{array}$$

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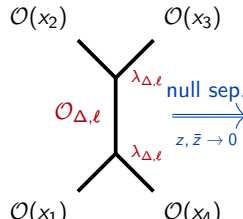
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[Rattazzi et al. (2008); Fitzpatrick et al. (2013)]

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Stronger constraints if we add supersymmetry \rightarrow SCFT

We focus on 4d $\mathcal{N} = 4$ SYM with $SU(N)$

per se: “bootstrap” non protected contributions

AdS/CFT correspondence

Introduction: AdS/CFT correspondence

4d $\mathcal{N} = 4$ SYM: $SU(N)$ gauge group, g_{YM} coupling, $SU(4)_R$



IIB string theory on $\text{AdS}_5 \times S^5$
 $\sqrt{\alpha'}$ string length, g_s coupling

$$g_{\text{YM}}^2 = g_s$$

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$\mathcal{O}_p = \text{tr } \varphi^{\{M_1} \dots \varphi^{M_p\}}$, $\Delta = p$

\mathcal{O}_2 stress tensor

$p \geq 3$



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single particle states

SUGRA excitations, graviton

SUGRA KK modes

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Half-BPS operators

So far, intense study of $\frac{1}{2}$ -BPS operators

- protected dimension and three-point functions
- maximal amount of supersymmetry
- Schur ops in the chiral algebra construction [Beem, Rastelli, Rees (2013)]

Chiral algebra idea: correlators of protected operators reduce to meromorphic functions when after an R-symmetry twist are restricted to lay on a plane.

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Four-point function of \mathcal{O}_2 's

- Superconformal WI are solved [Nirschl, Osborn (2005); Beem et al. (2017)]
 - Completely fixed up to N^{-4} [Alday, Bissi (2017)]
 - **However** as we go to higher orders at large N we run into problems pretty fast
 - mixing among degenerate ops
 - appearance of **higher trace ops**
- approachable
- 4pt of higher trace ops
Richer: ops $\notin \mathcal{O}_2 \times \mathcal{O}_2$
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Quarter-BPS operators: definition

Let us have a closer look to $\frac{1}{4}$ -BPS operators

$$\mathcal{O}_{pq} \sim \text{tr}(\varphi^{M_1} \dots) \text{tr}(\dots \varphi^{M_\Delta}) P_{M_1 \dots M_\Delta} + \frac{1}{N} (\text{single trace})$$

with $P_{M_1 \dots M_\Delta}$ projecting in the (q, p, q) representation of $SU(4)_R$.

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- Proliferation of $SU(4)_R$ tensor structures in their OPEs.
Ex: 10 structures in $\mathcal{O}_{02} \times \mathcal{O}_2$, 42 in $\mathcal{O}_{02} \times \mathcal{O}_{02}$ vs 6 in $\mathcal{O}_2 \times \mathcal{O}_2$.
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 \Rightarrow We introduce null polarization vectors and we build invariants out of them similar to [Cuomo, Karateev, Kravchuk (2018)]
- Contain $\frac{1}{2}$ -BPS Schur operators when expanded in $\mathcal{N} = 2$ supermultiplet \Rightarrow Ward Identities

Quarter-BPS operators: four-point function

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$$\langle \mathcal{O}_{p_1 q_1} \mathcal{O}_{p_2 q_2} \mathcal{O}_{p_3 q_3} \mathcal{O}_{p_4 q_4} \rangle \sim \sum_k \mathbb{T}_k \mathcal{G}_k(z, \bar{z}),$$

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$$\mathcal{G}_k(z, \bar{z}) = w_k(z, \bar{z}) + \sum_{m=1}^{\dim(\ker \chi)} \mathcal{H}_k(z, \bar{z}) v_k^{(m)}(z, \bar{z})$$



protected
 χ algebra



non
protected

$$\begin{aligned} \chi \left[\sum \mathbb{T}_k w_k \right] &= f(z) \\ \chi \left[\sum \mathbb{T}_k v_k^{(m)} \right] &= 0 \end{aligned}$$

Extract new
OPE data



[Alday, Caron-Huot (2018)]

OPE data

In the regime under analysis, the OPE data of a non protected exchanged operator get expanded as

$$a_{\Delta,\ell} = a_{\Delta,\ell}^{(0)} + \sum_{\kappa} \frac{a_{\Delta,\ell}^{(\kappa)}}{N^{2\kappa}}, \quad \Delta = \Delta^{(0)} + \sum_{\kappa} \frac{\gamma_{\Delta,\ell}^{(\kappa)}}{N^{2\kappa}}.$$

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[Caron-Huot (2017)]

We have studied

- $\langle \mathcal{O}_{02} \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \rangle \rightarrow$ protected (see next)
- $\langle \mathcal{O}_2 \mathcal{O}_{02} \mathcal{O}_{02} \mathcal{O}_2 \rangle \rightarrow$ OPE data in $\mathcal{O}_{02} \times \mathcal{O}_2$
- $\langle \mathcal{O}_{02} \mathcal{O}_{02} \mathcal{O}_{02} \mathcal{O}_{02} \rangle \rightarrow$ OPE data in $\mathcal{O}_{02} \times \mathcal{O}_{02}$

Higher-point functions

Why higher-points? Quarter-BPS operators appear as leading singularities in the OPE of $\mathcal{O}_2 \times \mathcal{O}_2$. For instance

$$\lim_{x_1 \rightarrow x_2} \langle \mathcal{O}_2(x_1) \mathcal{O}_2(x_2) \cdots \rangle |_{(2,0,2)} \supset \lambda_{\mathcal{O}_{02}} \langle \mathcal{O}_{02} \cdots \rangle$$

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So to extract the contribution of a specific operator, we just have to take limits of higher-point functions. Consider the 5-point function of [Gonçalves, Pereira, Zhou (2019)] in supergravity

$$\langle \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \mathcal{O}_2 \rangle = \frac{1}{N} \mathcal{G}_{\text{disc}}^{(5)} + \frac{1}{N^3} \mathcal{G}_{\text{tree}}^{(5)}$$

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Conclusions

Summary

We have started a systematic study of quarter-BPS operators in $\mathcal{N} = 4$ SYM. In particular

- How to derive Ward Identities
- How to deal with R-symmetry structures
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Outlook

Several directions can be pursued

- Find the explicit form of **superconformal blocks**, similarly to what has been done for $\frac{1}{2}$ -BPS
- Understand the **basis of function** of the SUGRA correlator (not only \bar{D} ?)
- **Higher-point functions**
- Determine the OPE data of **triple trace** operators appearing in $\mathcal{O}_2 \times \mathcal{O}_2$