



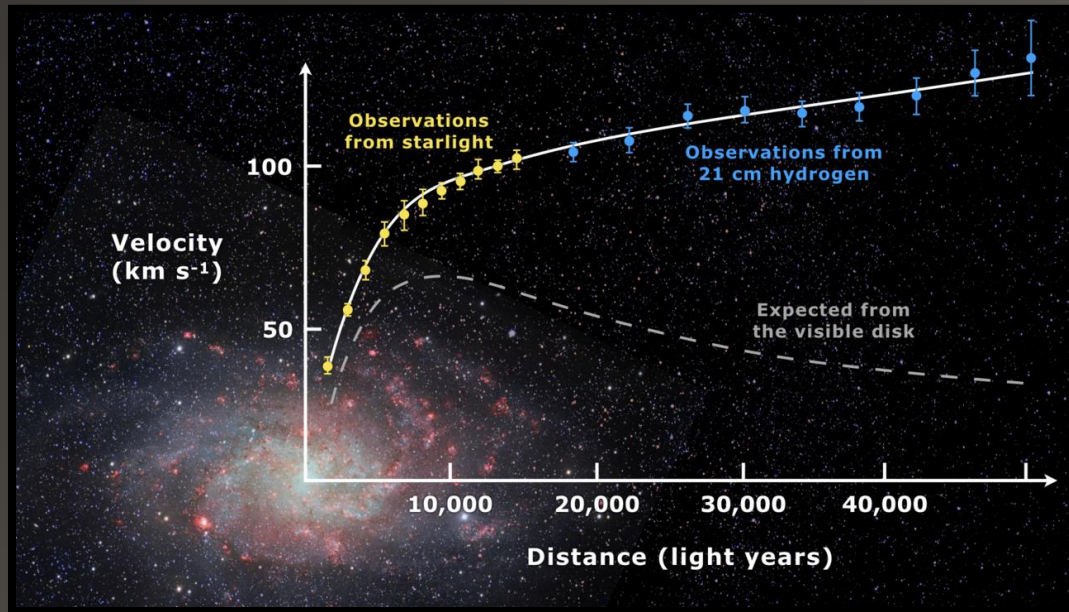
New Constraints on Dark Matter from Organic Targets: What Organic Chemistry can do for Direct Detection

CARLOS BLANCO

Based on work with J. Collar, Y. Kahn, & B. Lillard

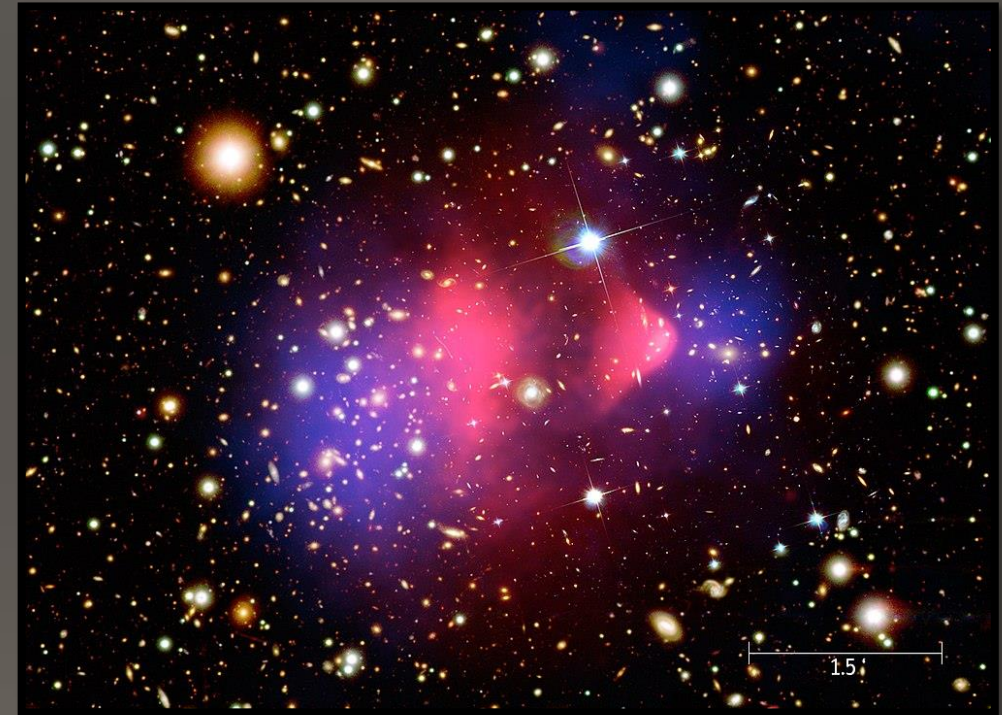
Evidence for Dark Matter

Galactic rotation curves



Solid: Observed
Dashed: Expected from visible disk

Bullet cluster



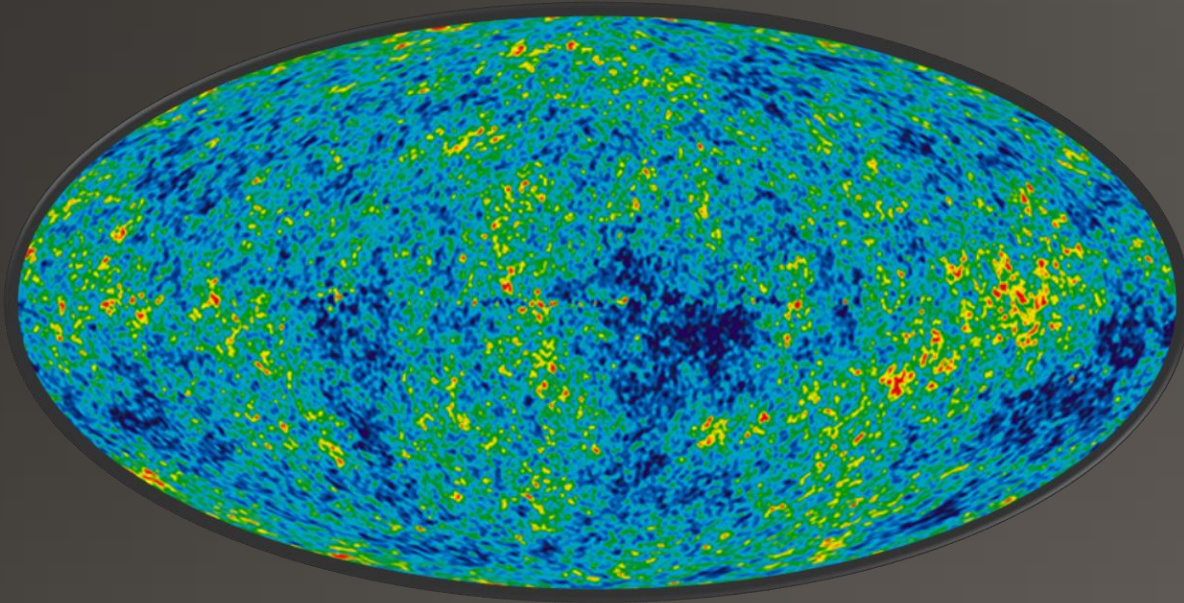
Blue: Mass from grav lensing
Pink: Baryonic mass from x-rays

X-ray: NASA/CXC/CfA/M.Markevitch et al.;

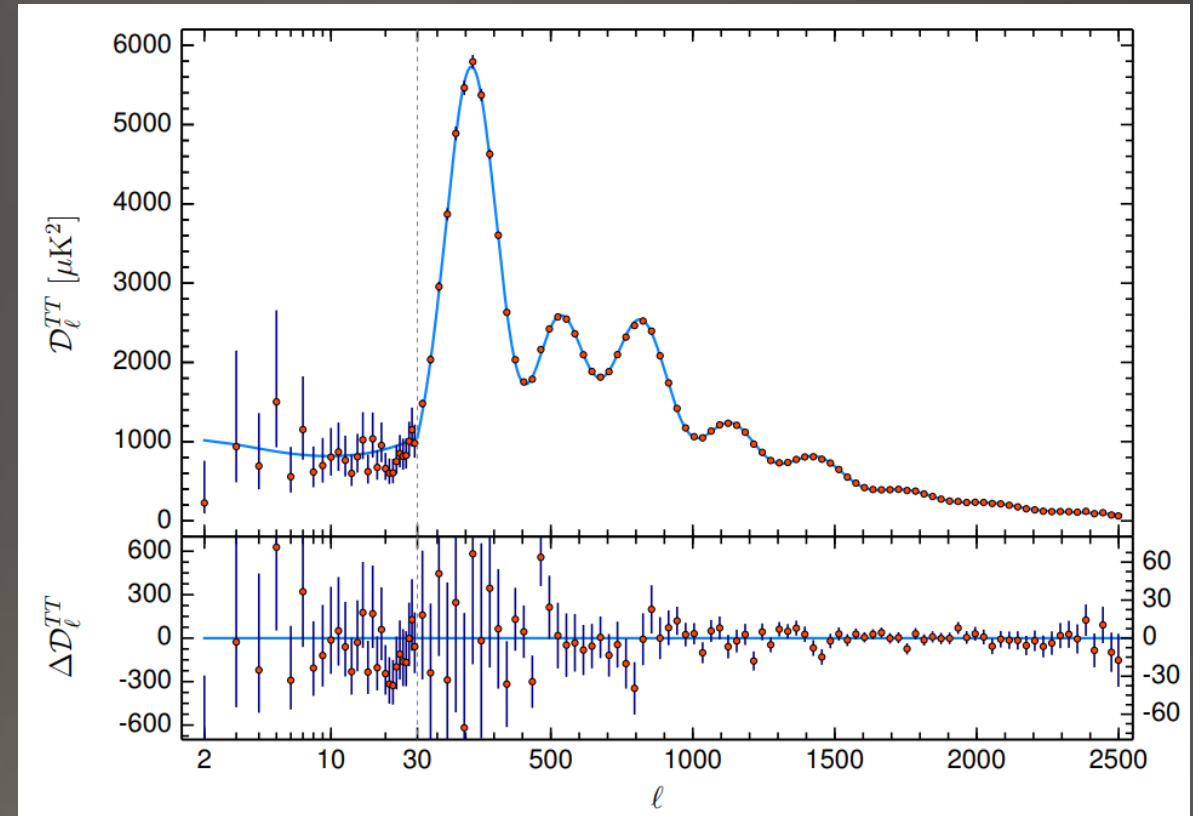
Optical: NASA/STScI; Magellan/U.Arizona/D.Clowe et al.;

Lensing Map: NASA/STScI; ESO WFI; Magellan/U.Arizona/D.Clowe et al.

Evidence for Dark Matter

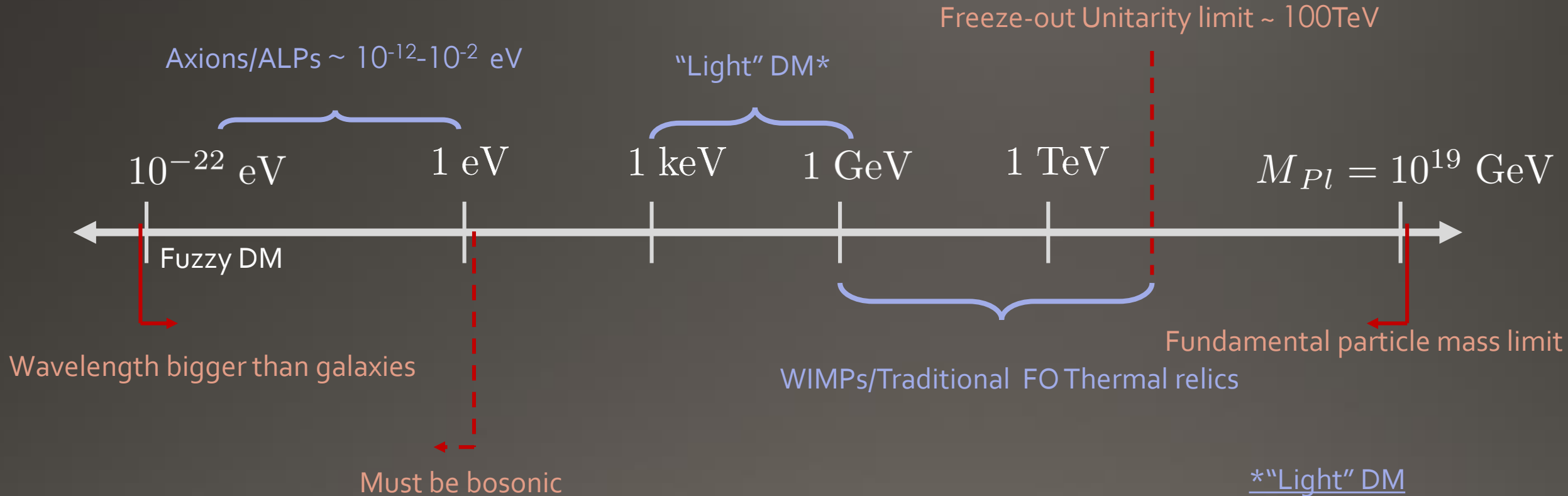


$$\Omega_c = 0.1200 \pm 0.0012$$



Planck results: A&A 2019

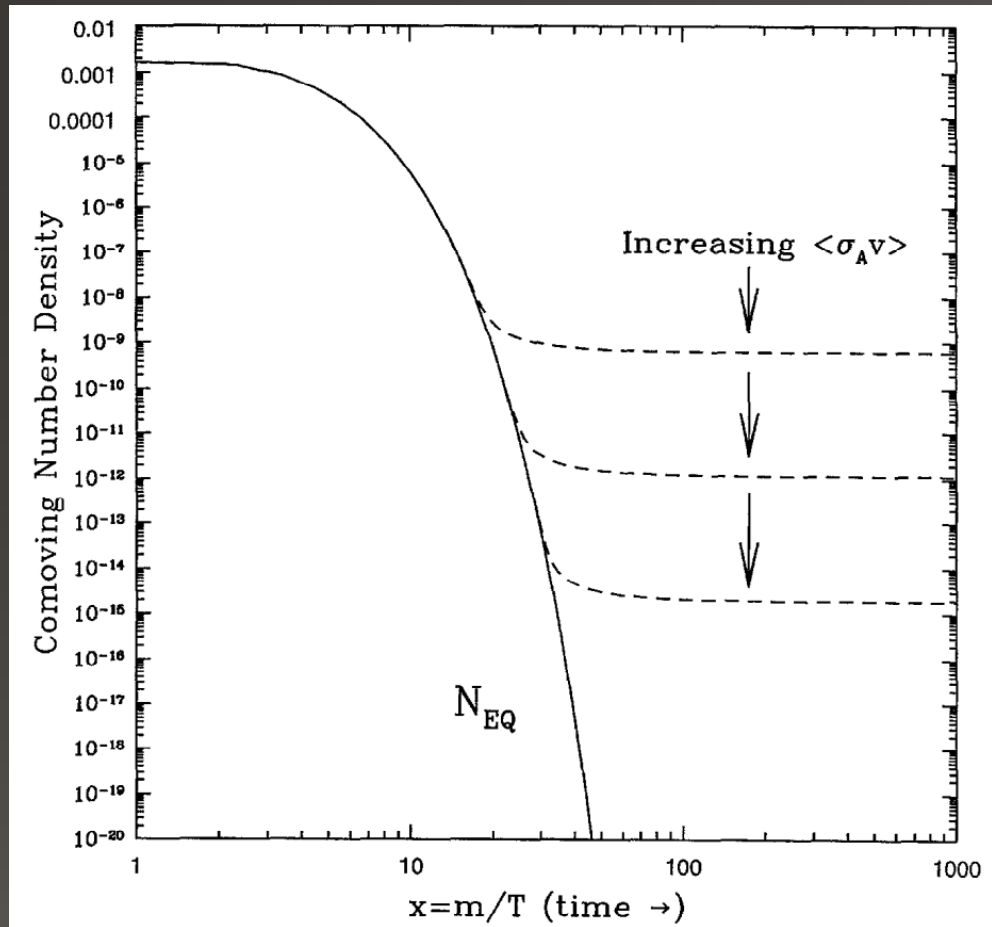
Dark Matter Mass Range



(NOT TO SCALE)

- *"Light" DM
- Assymmetric DM
 - Hidden sector
 - ELDER/SIMPs
 - Freeze-In
 - ...etc

WIMPs: Thermal Freeze-Out



Griest et. al: Phys Rep. 1996

$$\Omega_c h^2 \sim \frac{1}{\langle\sigma v\rangle} \quad \text{Known from Planck}$$

Predicts:

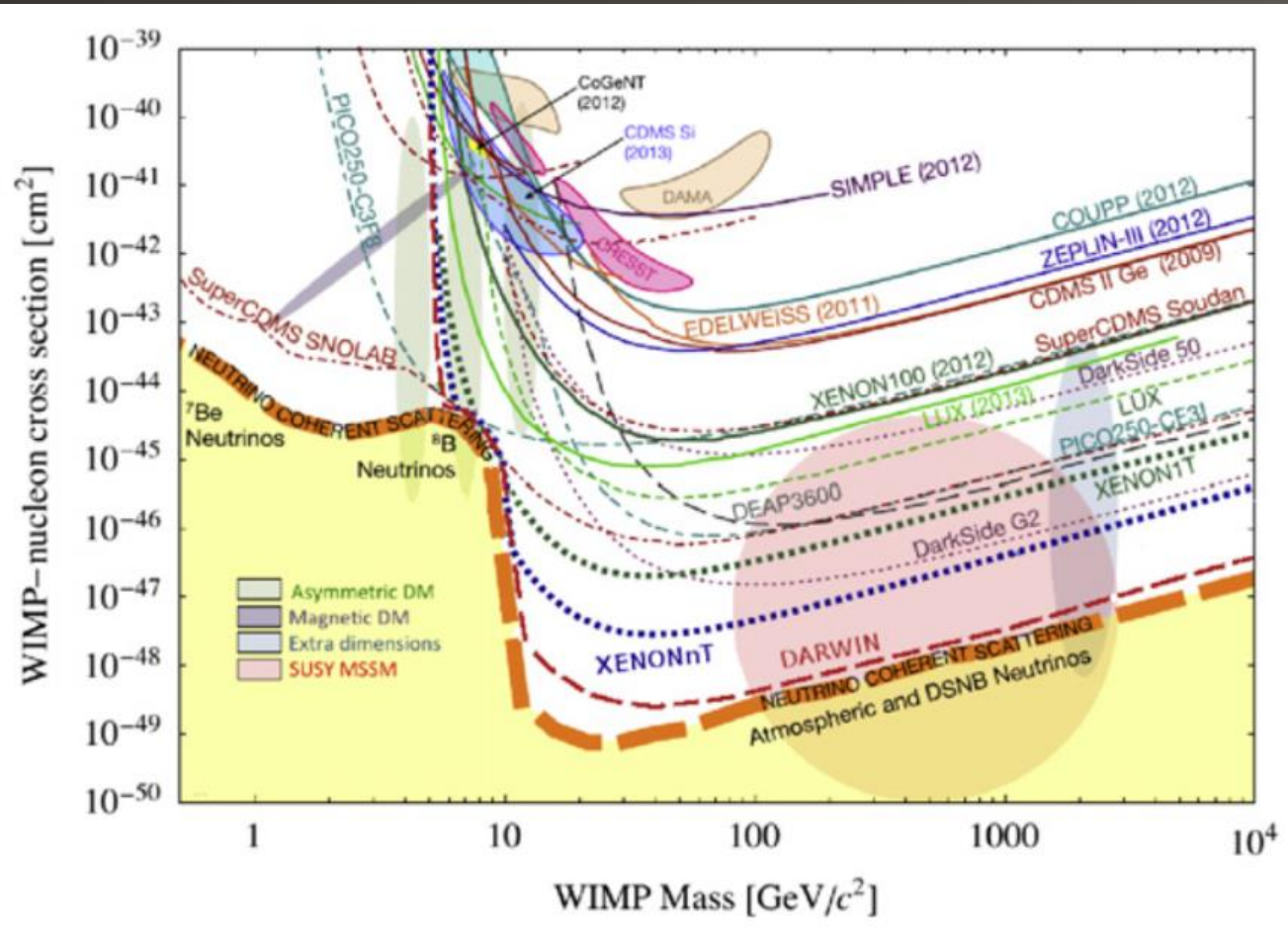
$$\langle\sigma v\rangle_{FO} \approx 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$

Weak scale! For masses between 1 GeV and 100 TeV

Subject to constraints from :

- CMB (recombination)
- Fermi-LAT (Dwarf spheroidal galaxy observations)
- BBN
- LEP, LHC, AMS,...etc

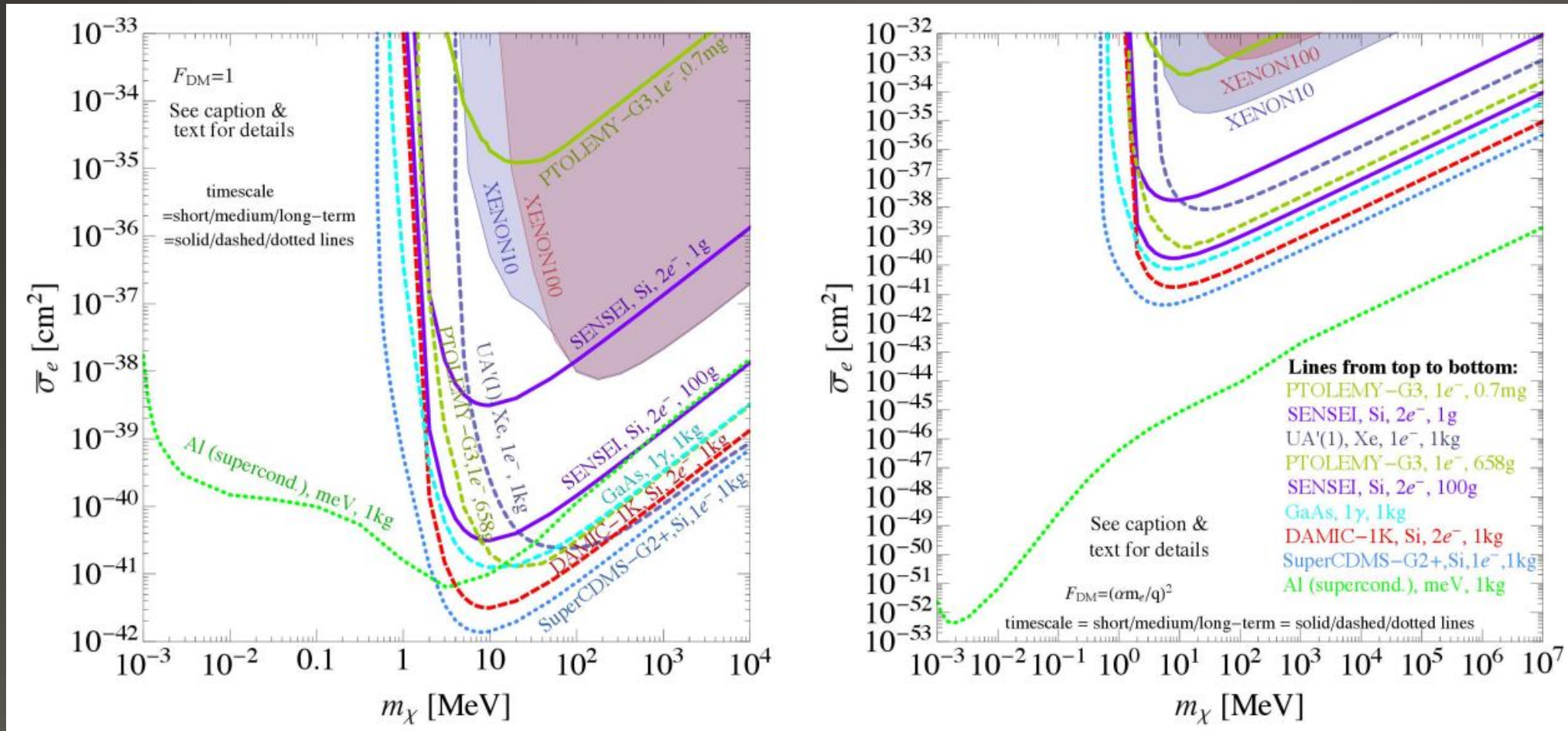
Current Direct Detection



Spin-independent cross section to nucleons

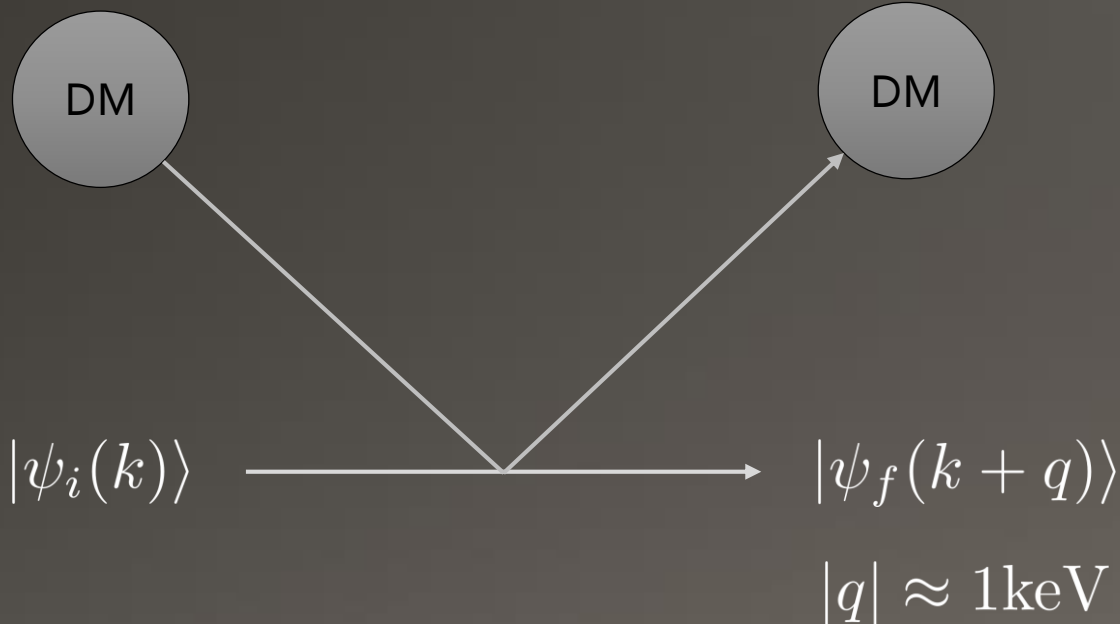
- Noble gases: Very effective at probing masses above 1 GeV
- Will encounter irreducible neutrino floor from
- Linear at high masses due to number density
- Loss of sensitivity due to kinematics at low masses

Sub-GeV Dark Matter (e^- scattering)



Light DM-electron scattering

$v \approx 300\text{km/s}$



Kinematic Requirement

$$\Delta E = \vec{v}_\chi \cdot \vec{q} - \frac{q^2}{2\mu_{\chi m}}$$

$$\Delta E \sim \mathcal{O}(\text{few eV})$$

What has such transition energies?

- Semiconductor bandgaps
- Inorganic crystal scintillators
- Organic liquid scintillators

DM-Electron Interaction Rate

Photon Rate

$$R = \xi \frac{2N_B \rho_\chi \bar{\sigma}_e}{8\pi m_\chi \mu_{\chi e}^2} \int \frac{d^3 \mathbf{q}}{q} \eta(v_{\min}) F_{\text{DM}}^2(q) \sum_i |f_i(\mathbf{q})|^2$$

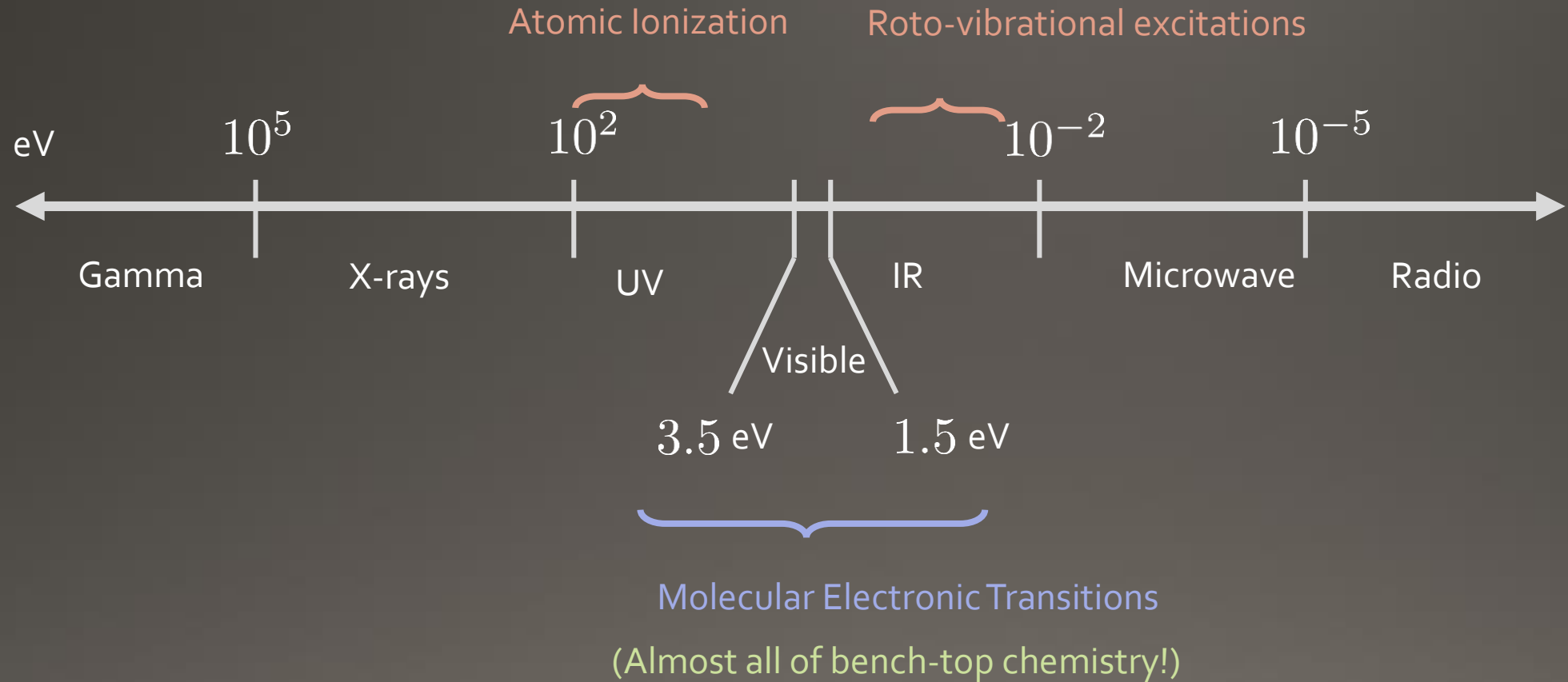
Sensitive when $R > \sim 0.1\text{Hz}$

ρ_χ :DM density ξ :detector efficiency

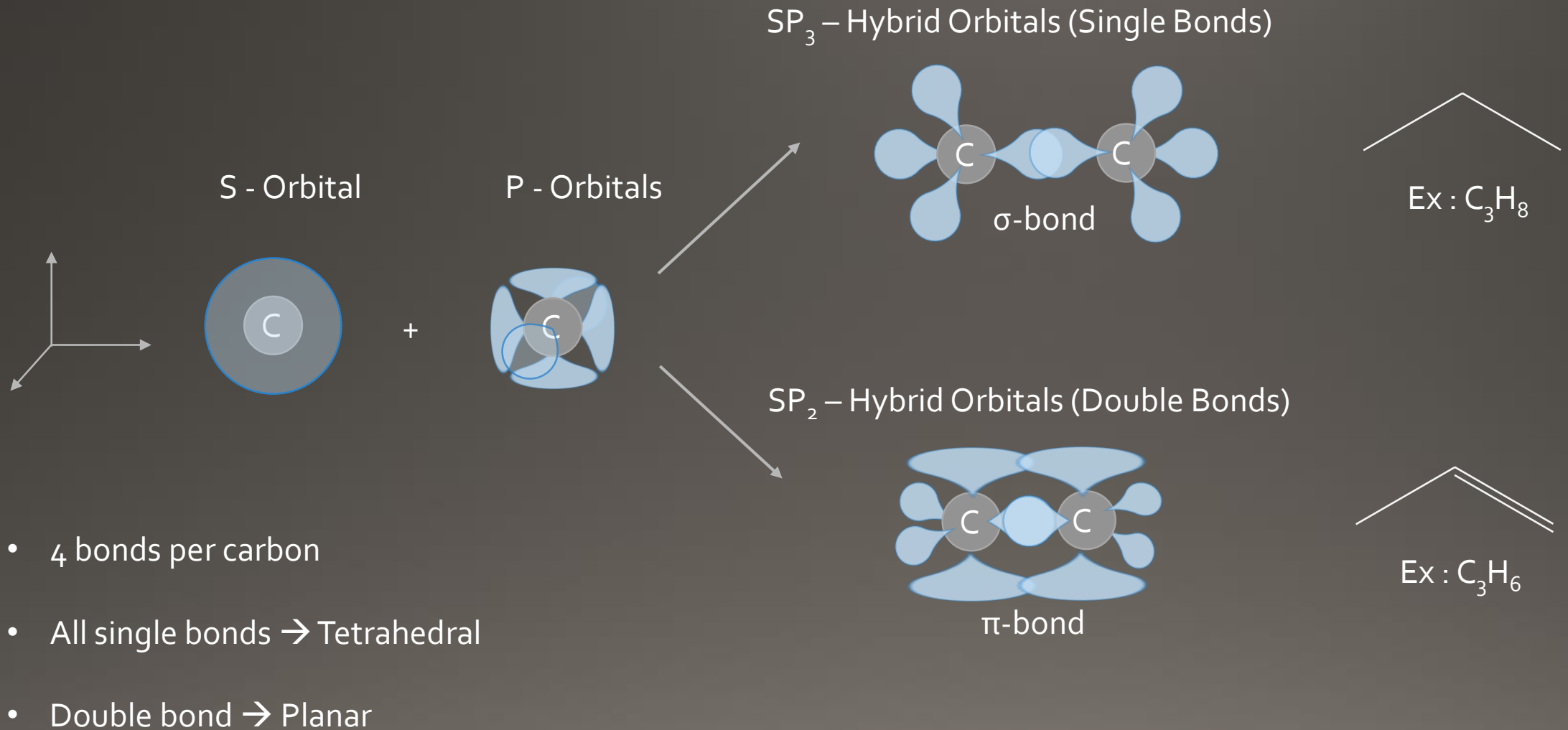
Molecular form factor

$$f_i(\mathbf{q}) = \int d^3 \mathbf{p} \tilde{\psi}_n(\mathbf{p}) \tilde{\psi}_{n+i}^*(\mathbf{p} + \mathbf{q})$$

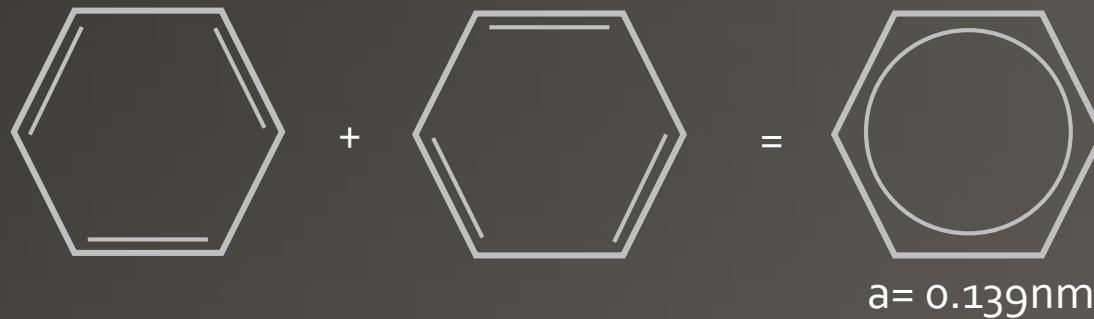
Low Energy Transitions



Organic Chemistry for Physicists: Orbitals



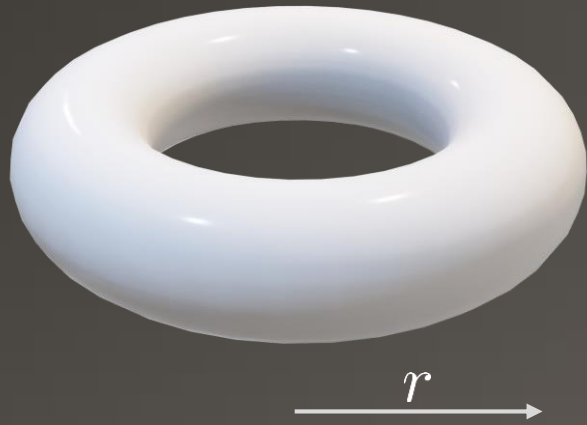
Organic Chemistry for Physicists: Benzene



- Two Kekule resonance forms
- Combined give an “aromatic” ring of delocalized charge
- π -electrons free to move from atom to atom
- σ -electrons still well localized in SP_2 orbital

Free electron model

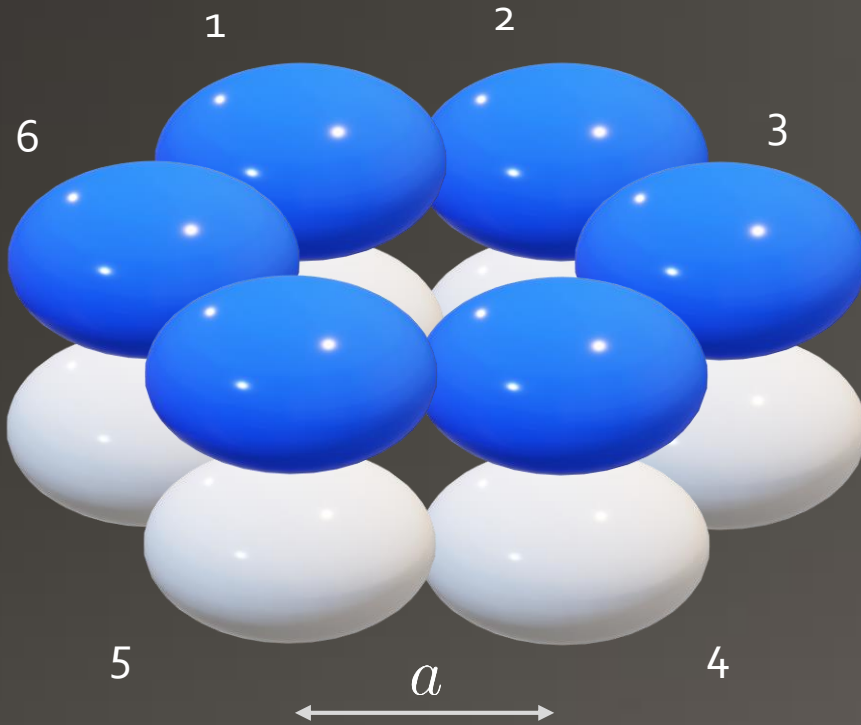
1D ring of charge



$$\Psi_n = \left(\frac{2}{r}\right)^{1/2} \sin \frac{\pi n x}{r} \quad E_n = \frac{h^2 n^2}{8mr^2}$$

- Let the π -electron travel on a one dimensional ring
- Well known solutions
- Quadratic spacing, two electrons per state
- No degenerate energy states
- Radius becomes an empirical parameter
- Not great for anything except benzene's first excitation

Linear Combination of Atomic Orbitals



Take linear combination of Slater atomic orbitals

$$\phi_{2p_z}(\mathbf{r}) = \mathcal{N} a_0^{-3/2} \frac{r \cos \theta}{a_0} \exp\left(\frac{-Z_{\text{eff}} r}{2a_0}\right)$$

One centered at each nuclei

$$\Psi_i = \sum_{i=1}^6 c_i \phi_{2P_z}(\mathbf{r} - \mathbf{R}_i)$$

Solutions come from diagonalizing the Hamiltonian matrix

$$\sum_{i=1}^6 [(H_{ji} - E_i \delta_{ji}) c_i] = 0, \text{ for } j = 1, 2, \dots, 6$$

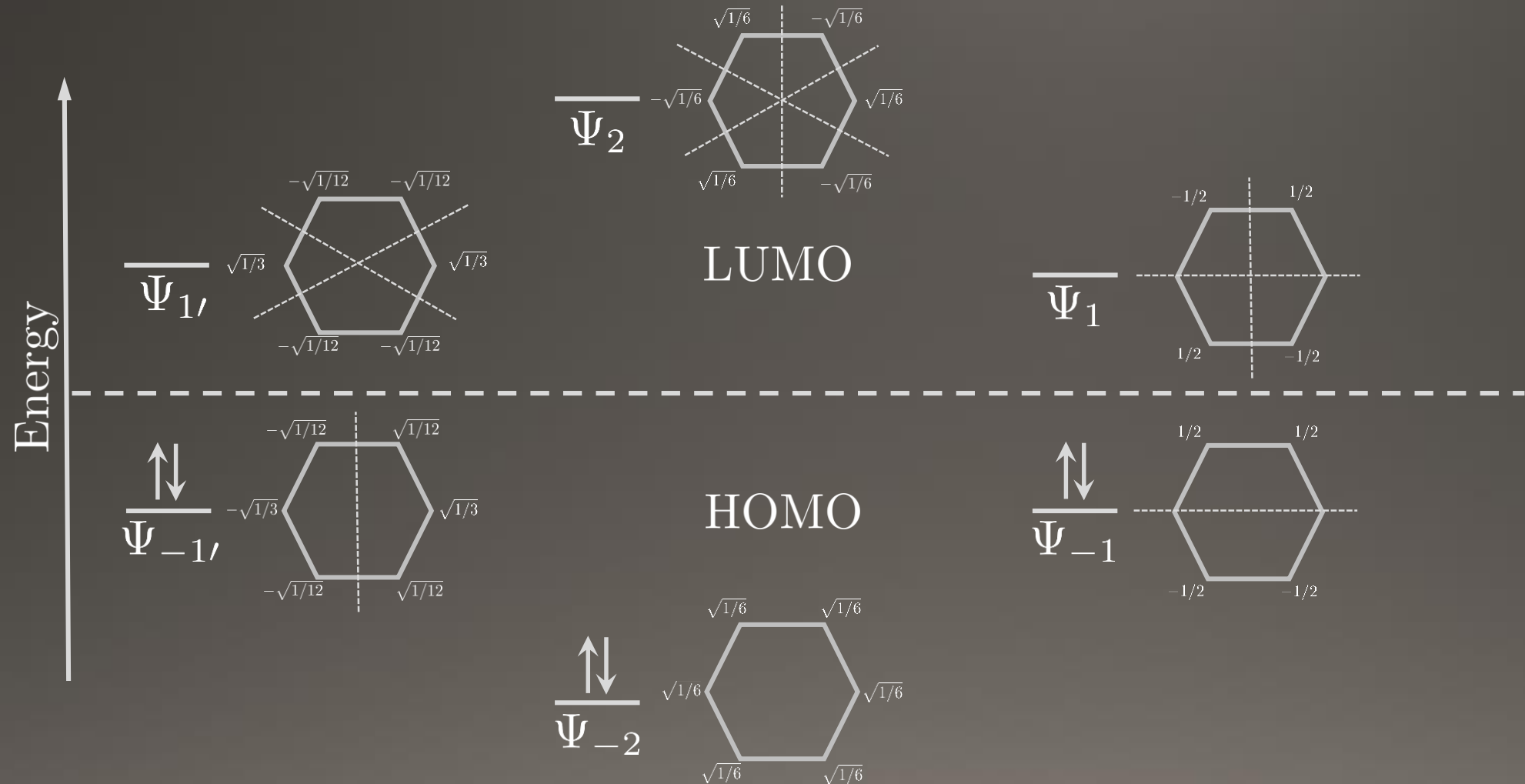
$$H_{ji} = \langle \Psi_j | \mathcal{H} | \Psi_i \rangle$$

Energies depend only on on-site, H_{ii} , energy and nearest neighbor, H_{ij} , resonance integrals

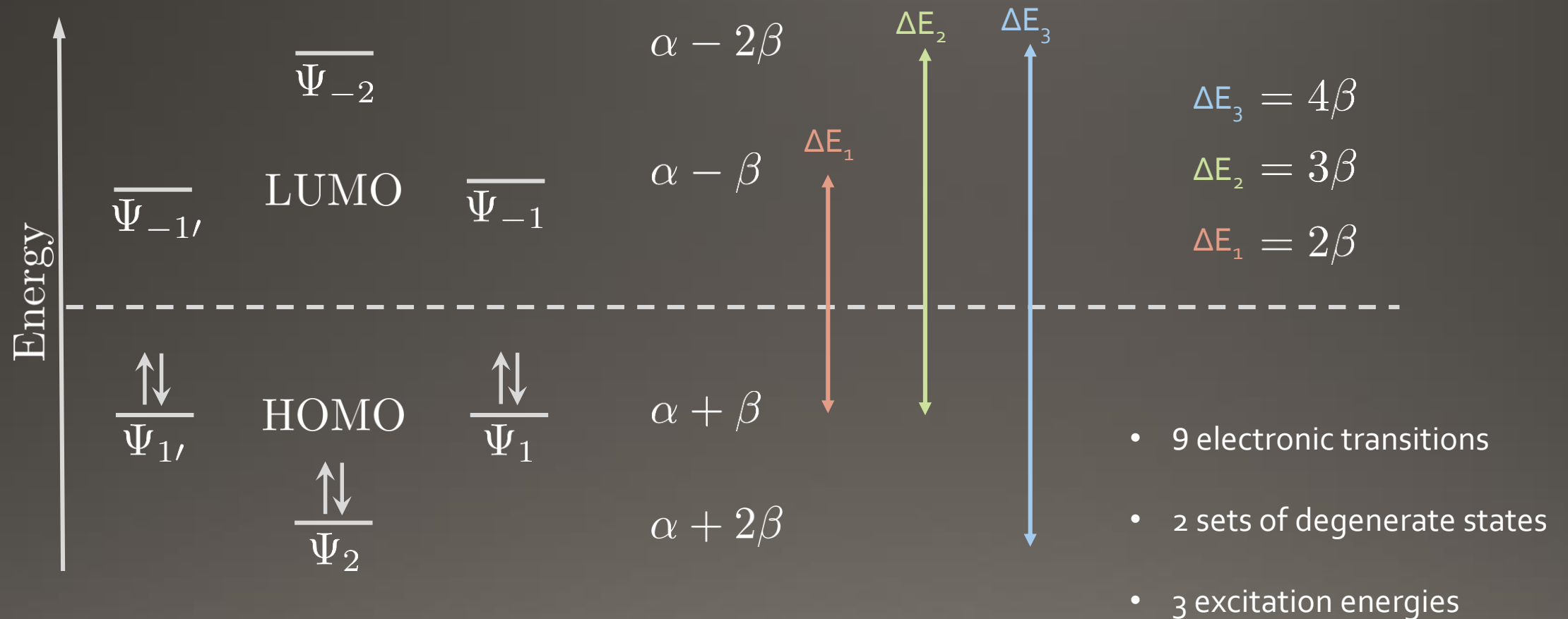
$$\alpha \equiv H_{ii}$$

$$\beta \equiv H_{ij}$$

LCAO: Benzene Hückel solutions



Hückel Transitions



Anti-symmetry & Slater Determinants

Electrons need to live in anti-symmetric states! (exclusion principle)

$$\begin{aligned} \Psi(\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_N) &= \frac{1}{\sqrt{N!}} \begin{vmatrix} \chi_1(\mathbf{x}_1) & \chi_2(\mathbf{x}_1) & \cdots & \chi_N(\mathbf{x}_1) \\ \chi_1(\mathbf{x}_2) & \chi_2(\mathbf{x}_2) & \cdots & \chi_N(\mathbf{x}_2) \\ \vdots & \vdots & \ddots & \vdots \\ \chi_1(\mathbf{x}_N) & \chi_2(\mathbf{x}_N) & \cdots & \chi_N(\mathbf{x}_N) \end{vmatrix} \\ &\equiv |\chi_1, \chi_2, \dots, \chi_N\rangle \\ &\equiv |1, 2, \dots, N\rangle \end{aligned}$$

Example, a simple 2 particle state:

$$\begin{aligned} \Psi(\mathbf{x}_1, \mathbf{x}_2) &= \frac{1}{\sqrt{2}} \{ \chi_1(\mathbf{x}_1) \chi_2(\mathbf{x}_2) - \chi_1(\mathbf{x}_2) \chi_2(\mathbf{x}_1) \} \\ &= \frac{1}{\sqrt{2}} \begin{vmatrix} \chi_1(\mathbf{x}_1) & \chi_2(\mathbf{x}_1) \\ \chi_1(\mathbf{x}_2) & \chi_2(\mathbf{x}_2) \end{vmatrix} \end{aligned}$$

Pariser, Pople, & Parr Method

$$\mathcal{H}_{ppp} = \mathcal{H}_{core} + \sum_{ij} \frac{e^2}{r_{ij}}$$

Hückel "core"
Hamiltonian

Electron repulsion term

1st singlet excitations

$$\Psi_1^{s_0} = 1/\sqrt{2} (\Psi_1^{-1} + \Psi_{1'}^{-1'})$$

$$\Psi_2^{s_0} = 1/\sqrt{-2} (\Psi_1^{-1} - \Psi_{1'}^{-1'})$$

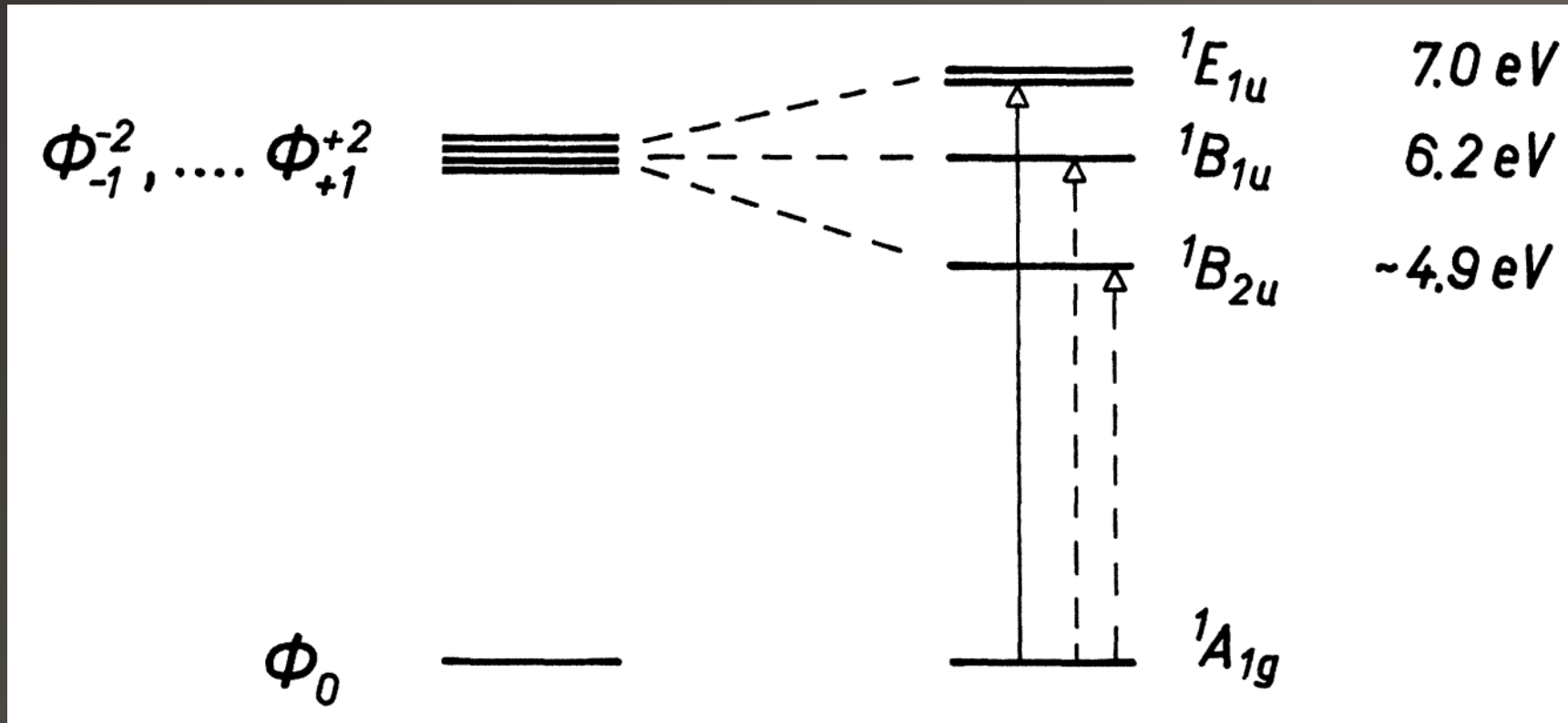
$$\Psi_3^{s_0} = 1/\sqrt{2} (\Psi_1^{-1'} + \Psi_{1'}^{-1})$$

$$\Psi_4^{s_0} = 1/\sqrt{-2} (\Psi_1^{-1'} - \Psi_{1'}^{-1})$$

$$\Psi_G = |\psi_2 \overline{\psi_2} \psi_1 \overline{\psi_1} \psi_{1'} \overline{\psi_{1'}}|$$

$$\Psi_i^j = \frac{1}{\sqrt{2}} (|\psi_1 \overline{\psi_1} \dots \psi_i \overline{\psi_j} \dots \psi_N \overline{\psi_N}| - |\psi_1 \overline{\psi_1} \dots \psi_j \overline{\psi_i} \dots \psi_N \overline{\psi_N}|)$$

PPP Transitions



Wagniere, 1976

Hückel excitations

PPP excitations

1st singlet excitations

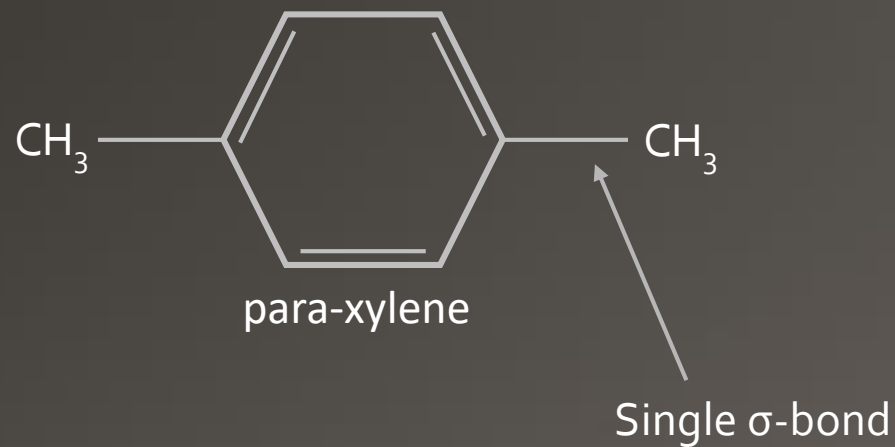
- Transform as symmetry subgroups of the symmetry group of Benzene
- Splitting due to electron repulsion, "Configurational Interactions"

Extension to p-xylene (EJ-301)

EJ-301

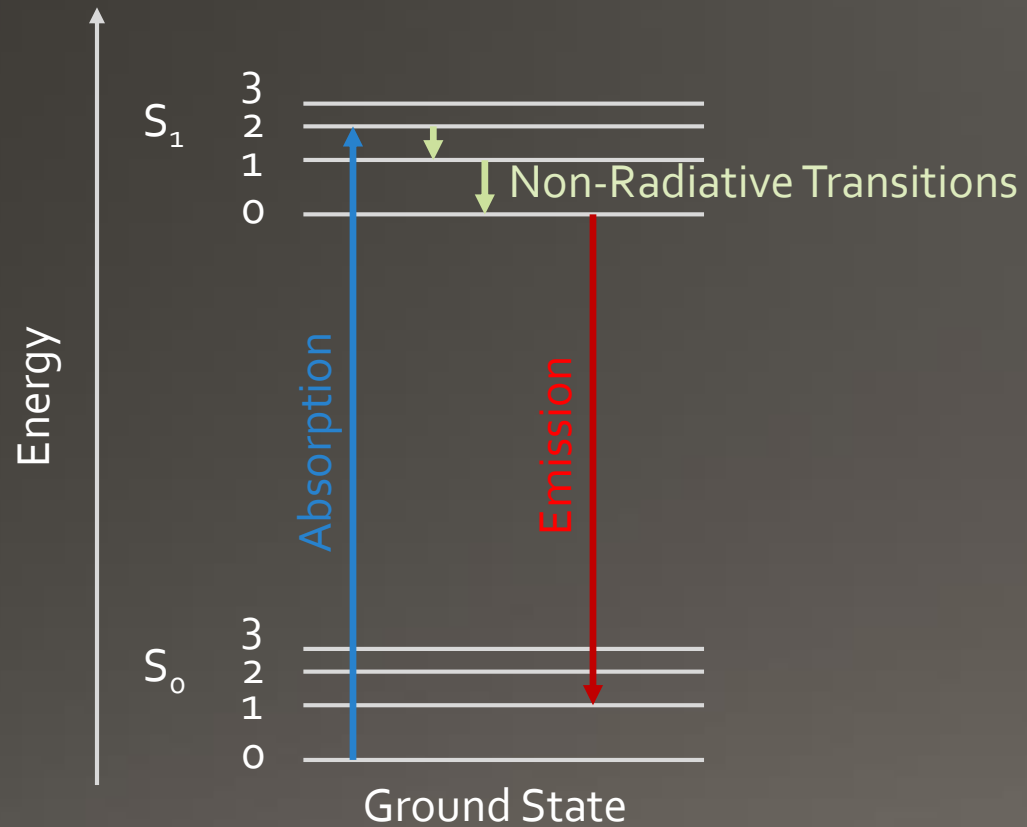
Solvent: para-xylene

Fluor: 5% by mass



- Absorption spectra of p-xylene is well described by slightly perturbed aromatic peaks
- Expected since methyl groups don't affect LCAO at leading order
- Produce perturbation at ~5% level in energy,
 - Benzene HOMO/LUMO gap = 4.9 eV
 - p-xylene HOMO/LUMO gap = 4.7 eV

Fluorescence: Quick review



- **Absorption:** Blue photon (E_A) promotes electronic transition: $S_0 \rightarrow S_1$ $\Delta J \neq 0$
- **Non-Radiative Transition:** Internal conversion or vibrational deexcitation
- **Emission:** Red Photon (E_E) emitted as electronic state relaxes back to S_0
- $E_E < E_A$

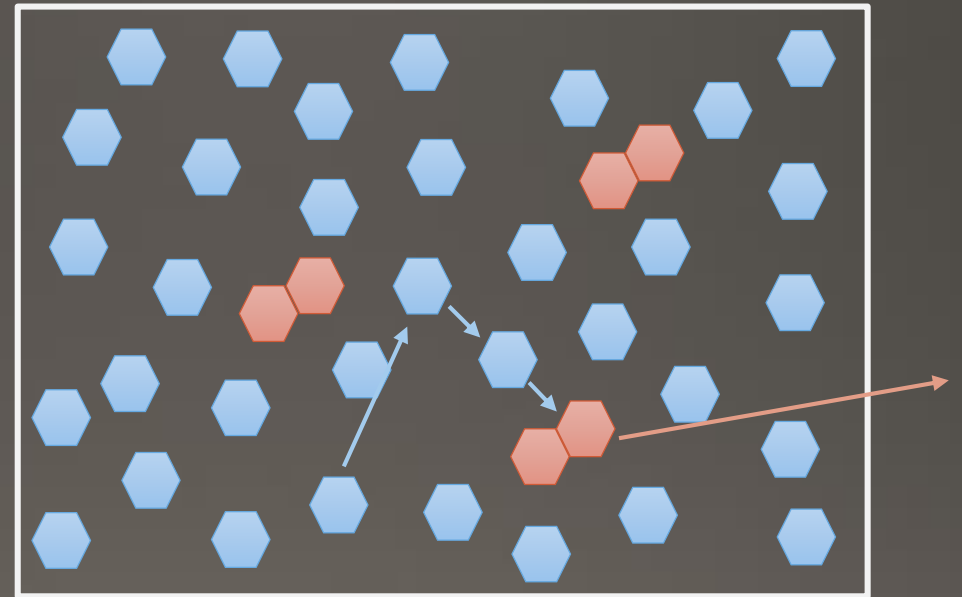
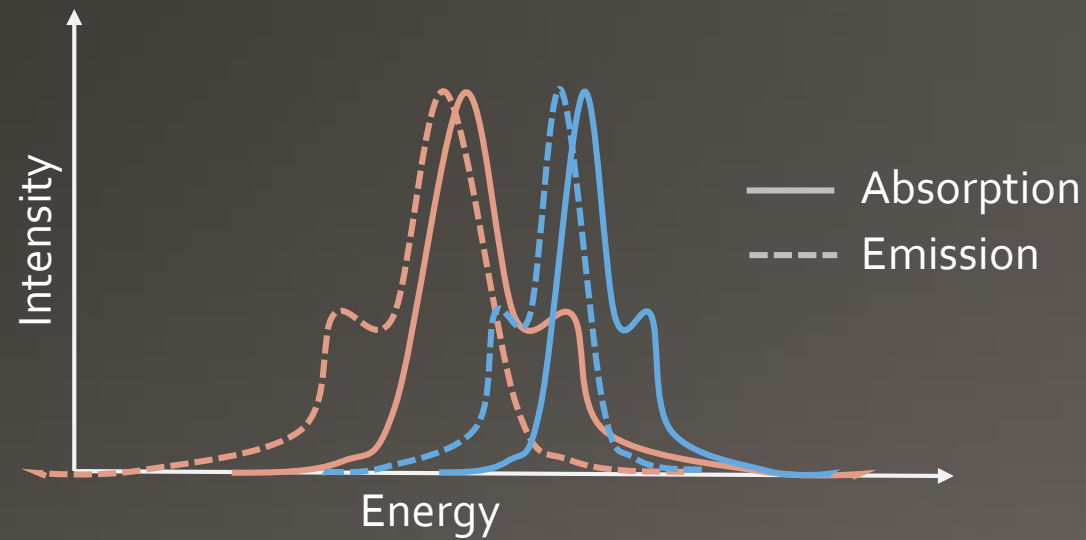
Fluorescence: Binary Scintillators



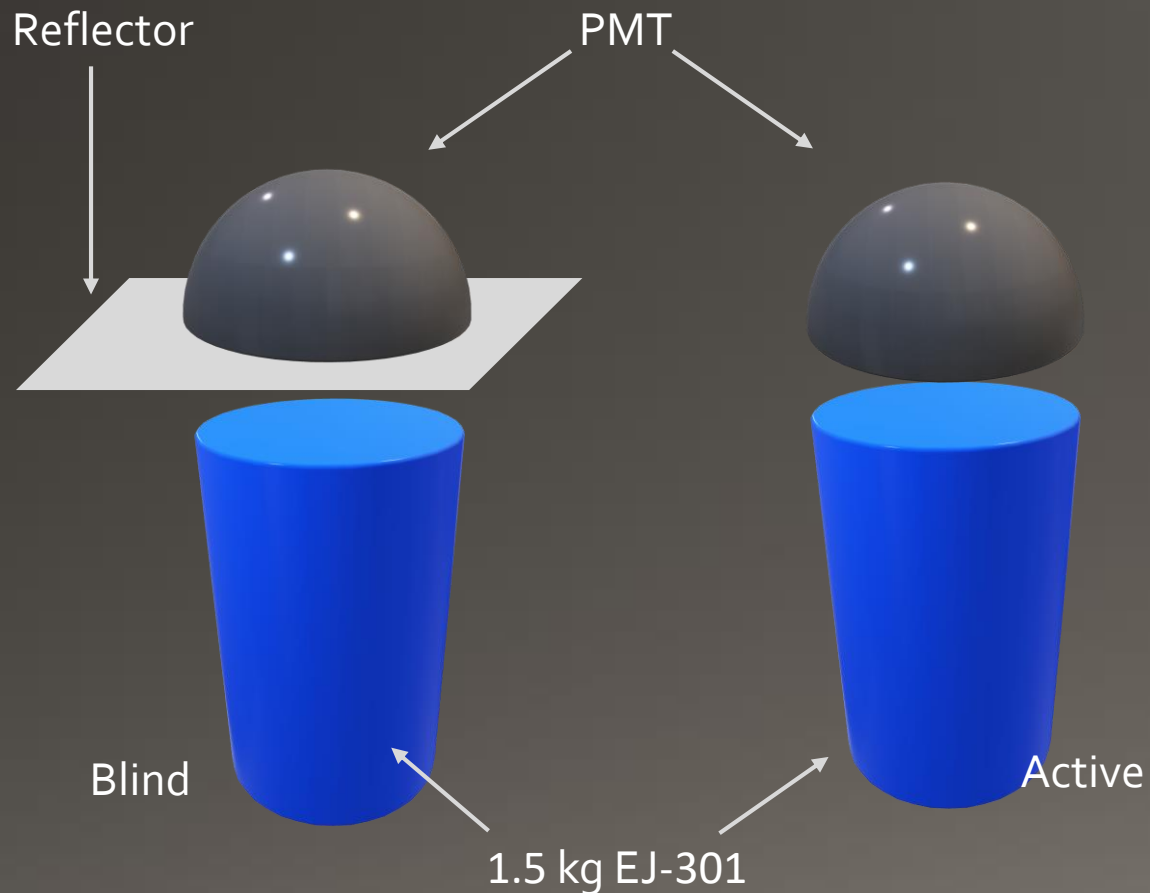
Solvent: Primary target starts the signal



Solute: Dilute fluor gets the signal out of the bulk

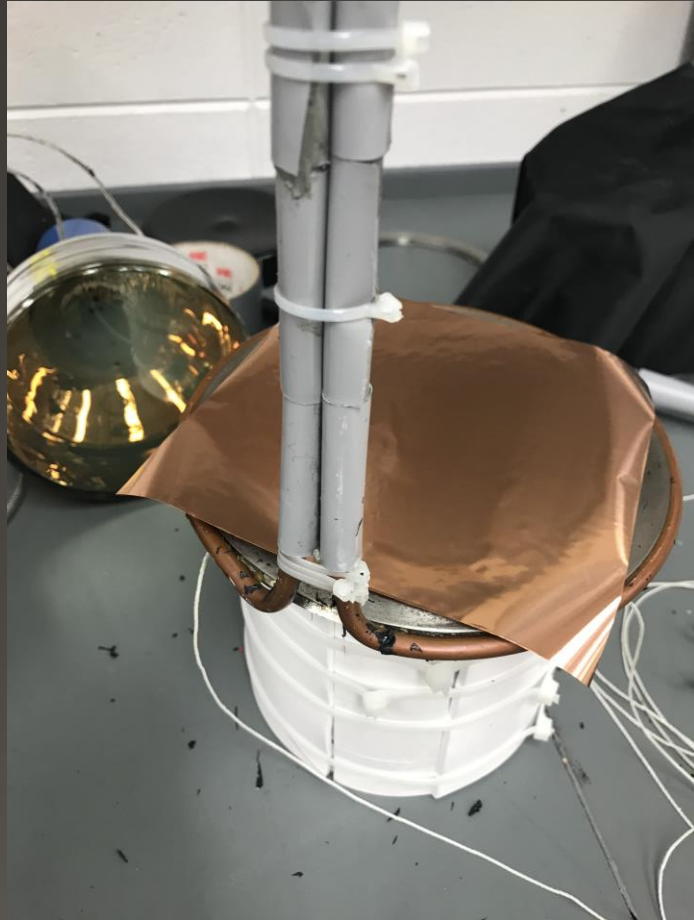


Experiment

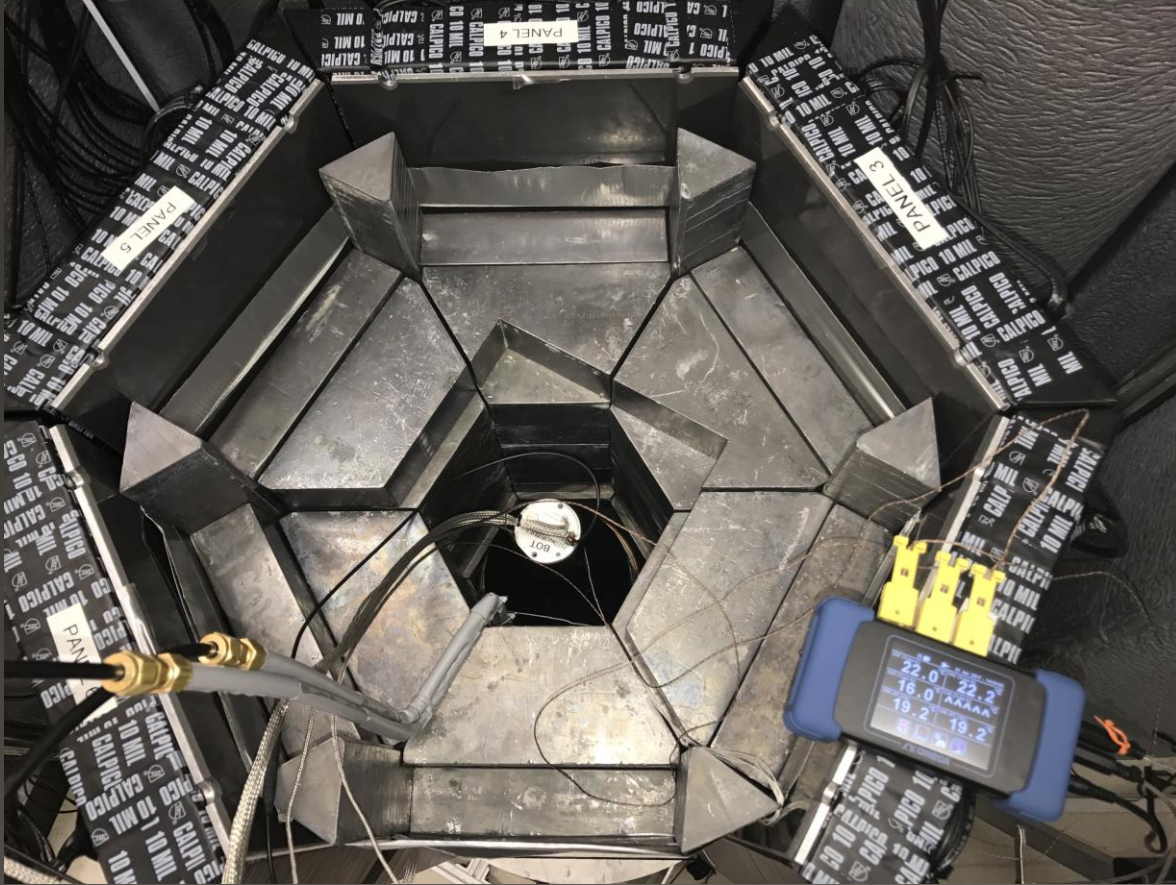


- 1.5 kg of EJ-301
- 0.1° Temperature control
- Need only about a week in each mode
- Record PMT rate in two modes,
 - Blind: Reflector shielding from the scintillator
 - Active: PMT observing the scintillator

Experimental Setup



Experimental Setup



Interaction Rate

Photon Rate

$$R = \xi \frac{2N_B \rho_\chi \bar{\sigma}_e}{8\pi m_\chi \mu_{\chi e}^2} \int \frac{d^3 \mathbf{q}}{q} \eta(v_{\min}) F_{\text{DM}}^2(q) \sum_i |f_i(\mathbf{q})|^2$$

Sensitive when $R > \sim 0.1\text{Hz}$

ρ_χ :DM density ξ :detector efficiency

Molecular form factor

$$f_i(\mathbf{q}) = \int d^3 \mathbf{p} \tilde{\psi}_n(\mathbf{p}) \tilde{\psi}_{n+i}^*(\mathbf{p} + \mathbf{q})$$

$$\xi \approx (0.77)(0.28) = 0.166$$

PMT efficiency

EJ-301 quantum efficiency

SAO Fourier transform

$$\tilde{\Psi}(\mathbf{k}) = \left(\sum_{i=1}^6 c_i^\psi e^{-i\mathbf{k} \cdot \mathbf{R}_i} \right) \tilde{\phi}(\mathbf{k}) \equiv \mathcal{B}_\psi(\mathbf{k}) \tilde{\phi}(\mathbf{k})$$

\mathcal{B}_ψ : k_x & k_y dependent phases

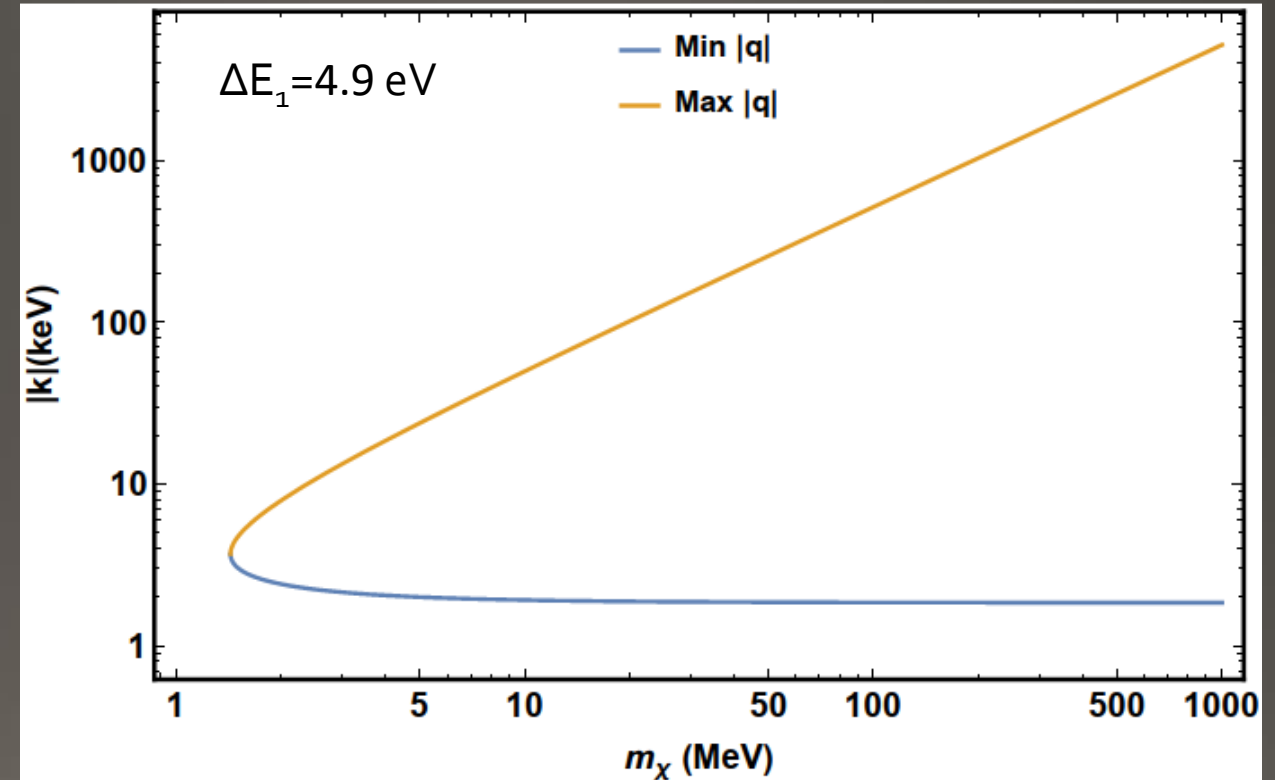
Kinematically Allowed Momenta

- Only has support on a range of momenta.
 - If q is too large, the DM would have escaped
 - If q is too small, there's not enough energy to excite the transition

$$\eta(v_{\min}) = \int \frac{4\pi v^2 dv}{v} g_{\chi}(v) \Theta(v - v_{\min})$$

$$v_{\min}(\mathbf{q}) = \frac{\Delta E_{ij}}{q} + \frac{q}{2m_{\chi}}$$

$$g_{\chi}(v) = \frac{1}{K} \exp\left(-\frac{|\mathbf{v} + \mathbf{v}_{\mathbf{E}}|^2}{v_0^2}\right) \Theta(v_{\text{esc}} - |\mathbf{v} + \mathbf{v}_{\mathbf{E}}|)$$



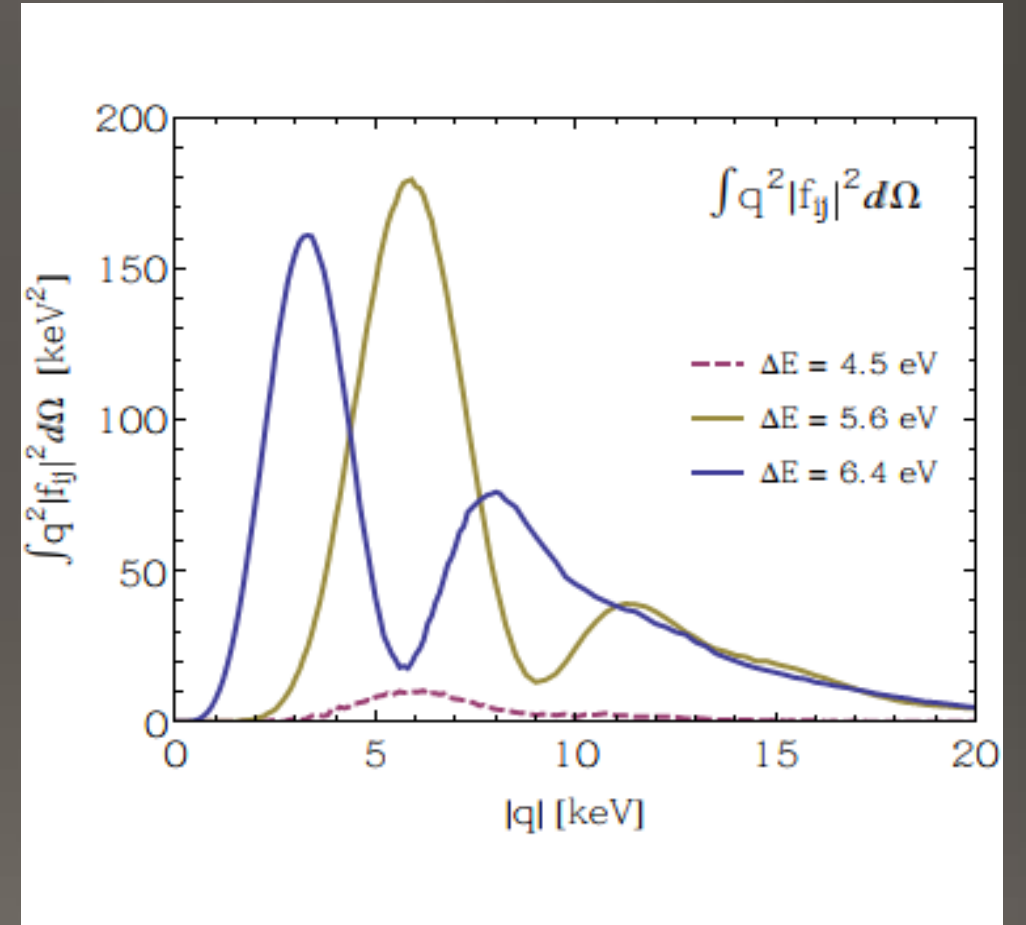
Molecular Form Factor

3 Dimensional integral in k-space

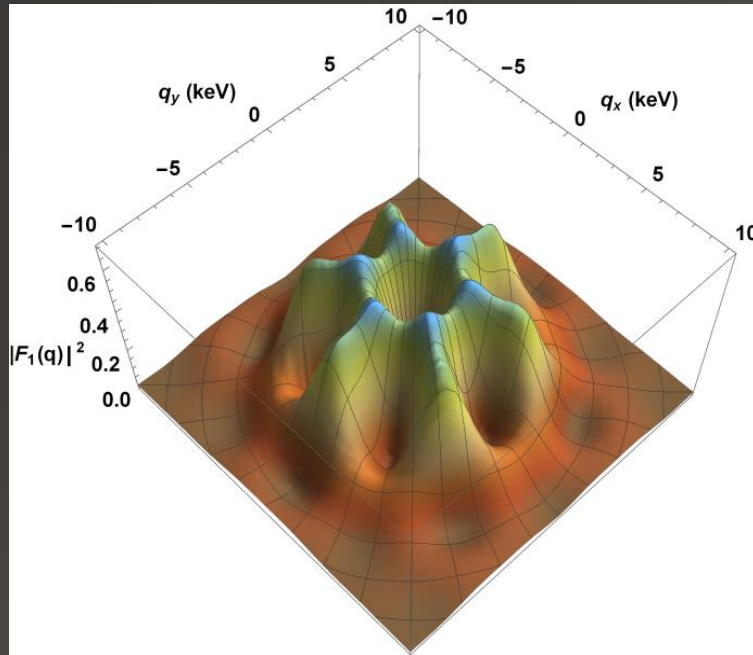
$$f_i(\mathbf{q}) = \int d^3\mathbf{p} \tilde{\psi}_n(\mathbf{p}) \tilde{\psi}_{n+i}^*(\mathbf{p} + \mathbf{q})$$

$$= \langle \psi_f(k + q) | e^{i\mathbf{q} \cdot \mathbf{x}} | \psi_i(k) \rangle$$

$$\tilde{\Psi}(\mathbf{k}) = \left(\sum_{i=1}^6 c_i^\psi e^{-i\mathbf{k} \cdot \mathbf{R}_i} \right) \tilde{\phi}(\mathbf{k}) \equiv \mathcal{B}_\psi(\mathbf{k}) \tilde{\phi}(\mathbf{k})$$

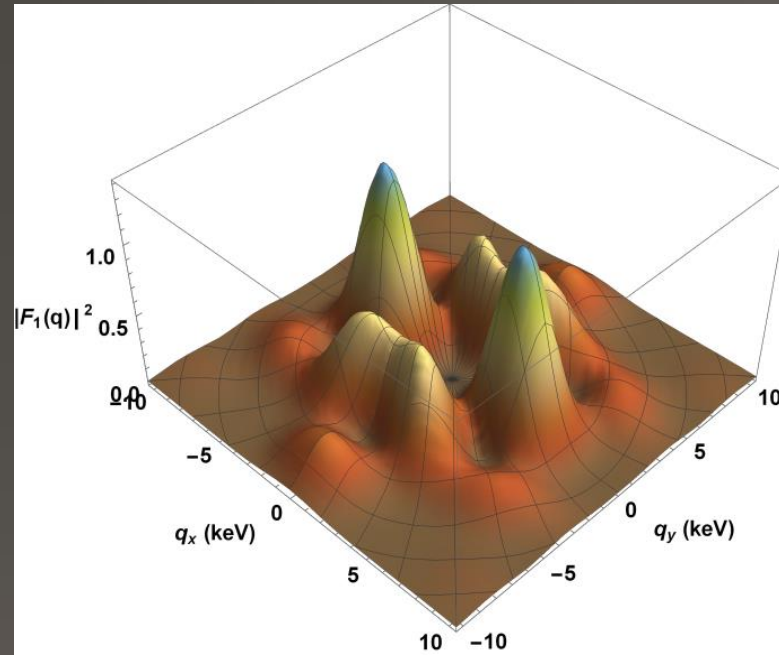


Form Factors Illustrated



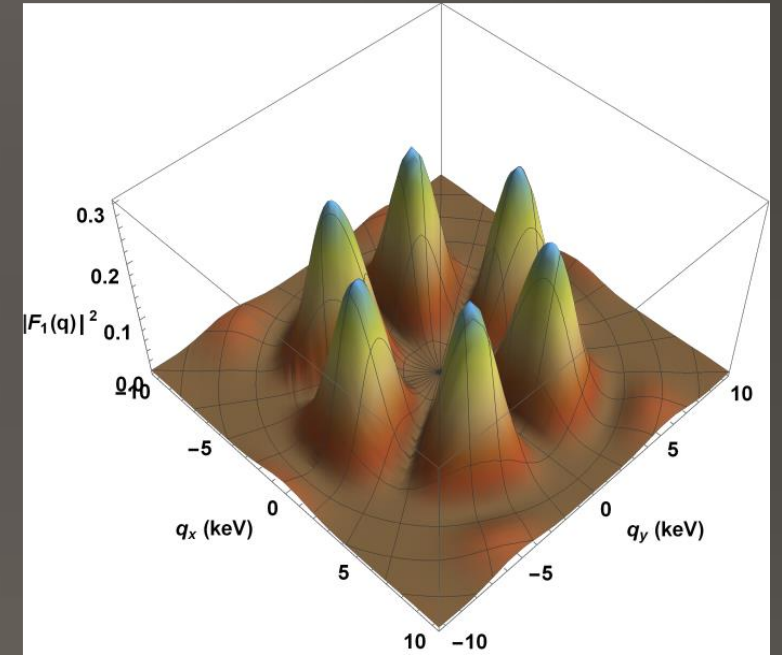
$$|F_1(q)|^2$$

$$\theta = 90^\circ$$



$$|F_2(q)|^2$$

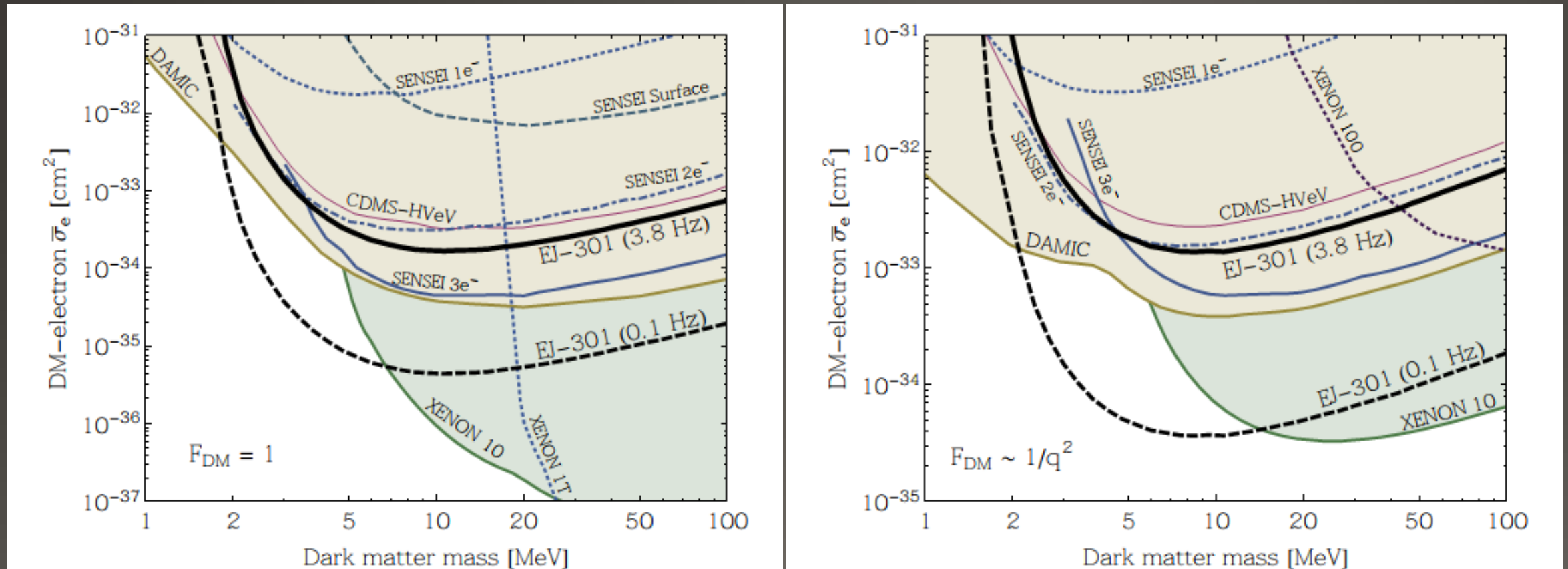
$$\theta = 90^\circ$$



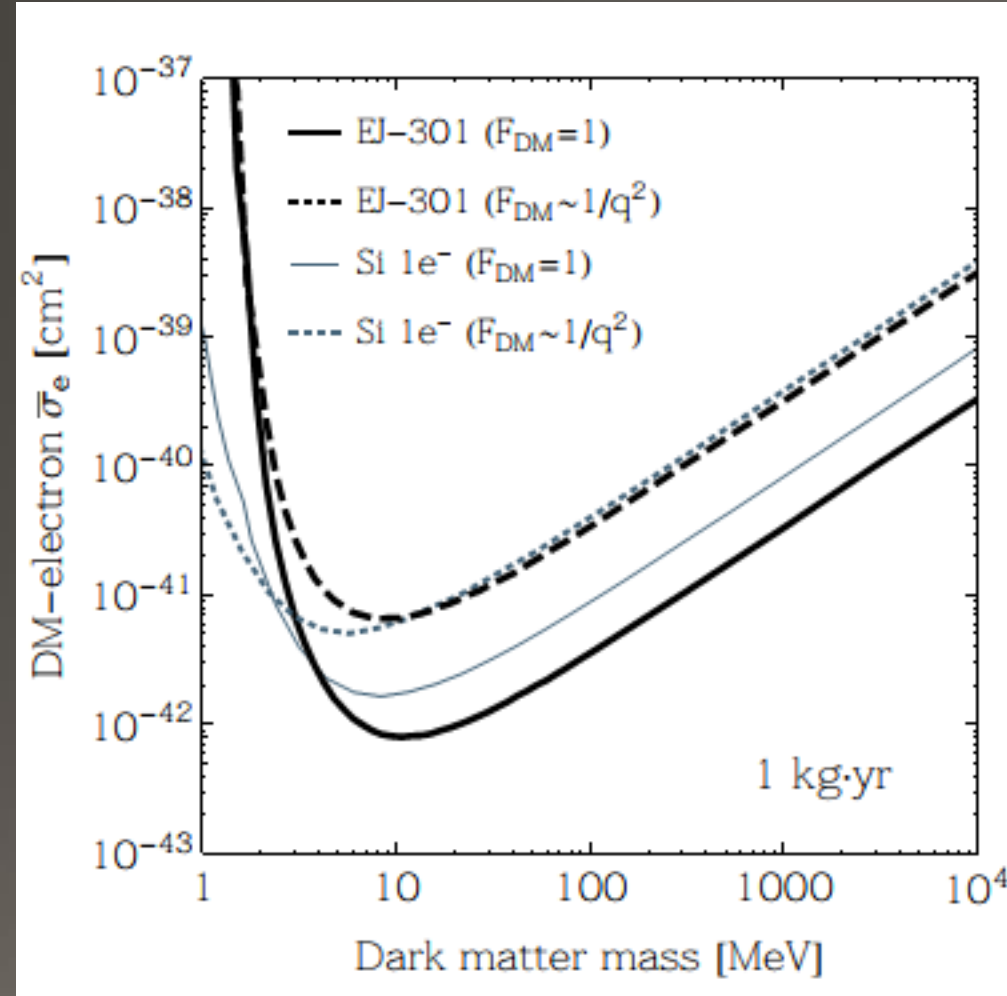
$$|F_3(q)|^2$$

$$\theta = 90^\circ$$

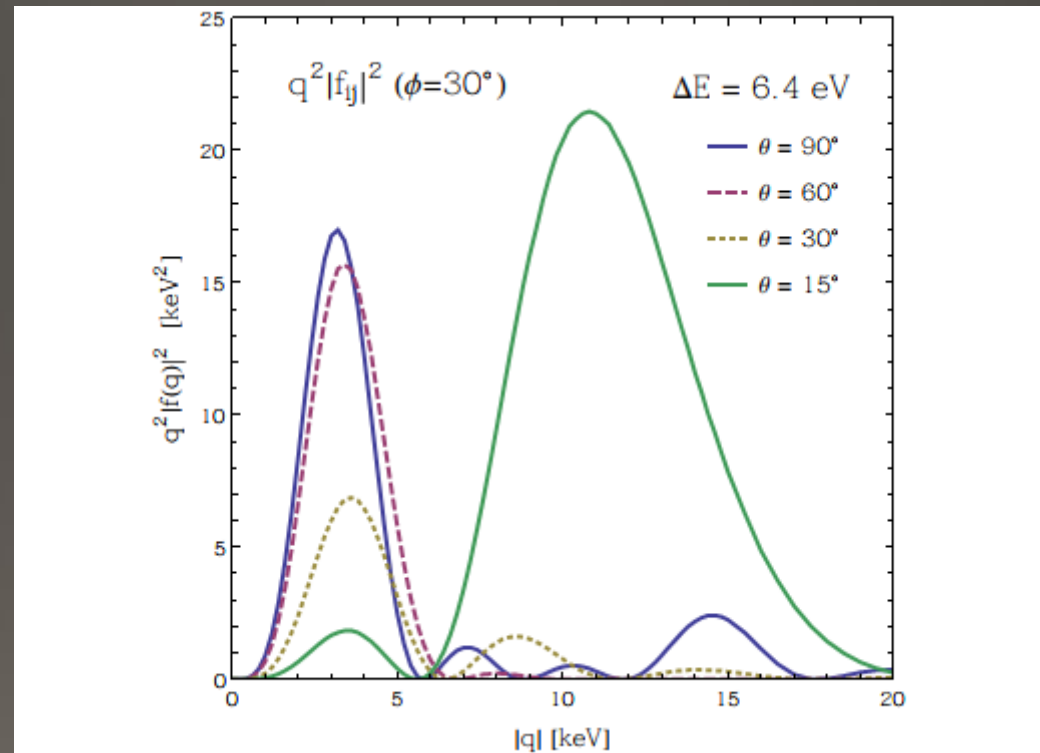
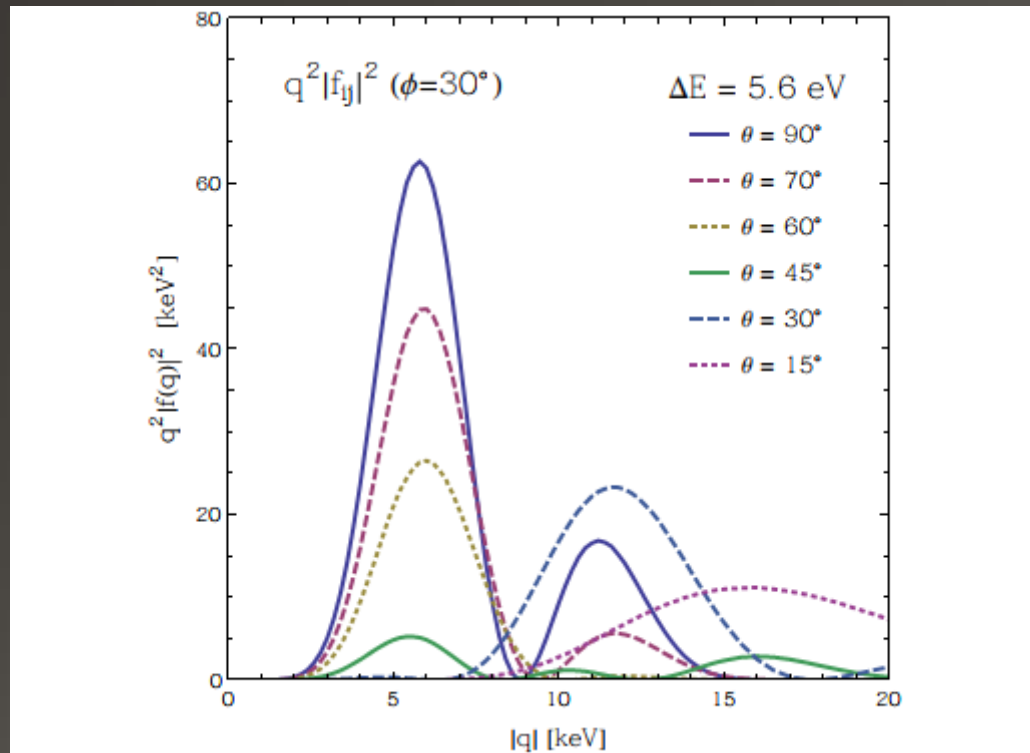
Predicted Sensitivity



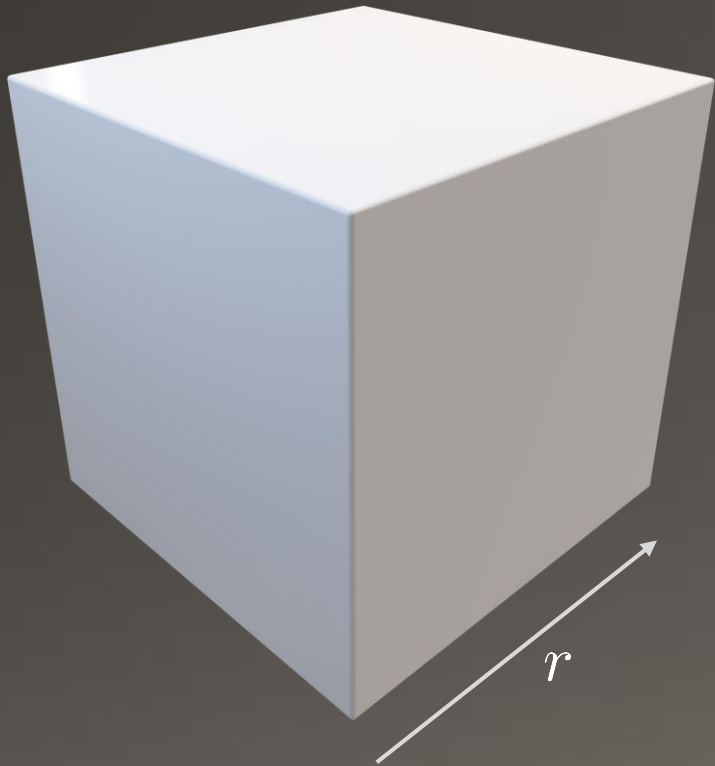
Sensitivity Comparison: Silicon



Directional Detection



Future Targets: e⁻s in Bigger Boxes



$$E_n \sim \frac{n^2 \hbar^2 \pi^2}{2mr^2}$$

$$E_1 \sim 0.1\text{eV} \left(\frac{10^{-7}\text{m}}{r} \right)^2$$

Need a box that is nano- to micro-scale!

i.e. Need systems with tens to thousands of atoms

Future Targets: Larger Polyacenes



Naphthalene

$$E_1 = 3.9\text{eV}$$

Phenanthrene

$$E_1 = 3.5\text{eV}$$

Natural extension, add more rings!

- Anthracene is known to be one of the best organic scintillators
- More complex ring structures are possible
- Long poly-ene bridges form the backbone of many dyes
- Quantum yield near unity
- Asymmetric ground state!

Future Targets: Quantum Dots

Semiconducting nano-spheres



$$E_{\text{confinement}} = \frac{\hbar^2 \pi^2}{2a^2} \left(\frac{1}{m_e} + \frac{1}{m_h} \right) = \frac{\hbar^2 \pi^2}{2\mu a^2}$$

$$E_{\text{exciton}} = -\frac{1}{\epsilon_r^2} \frac{\mu}{m_e} R_y = -R_y^*$$

$$\begin{aligned} E &= E_{\text{bandgap}} + E_{\text{confinement}} + E_{\text{exciton}} \\ &= E_{\text{bandgap}} + \frac{\hbar^2 \pi^2}{2\mu a^2} - R_y^* \end{aligned}$$

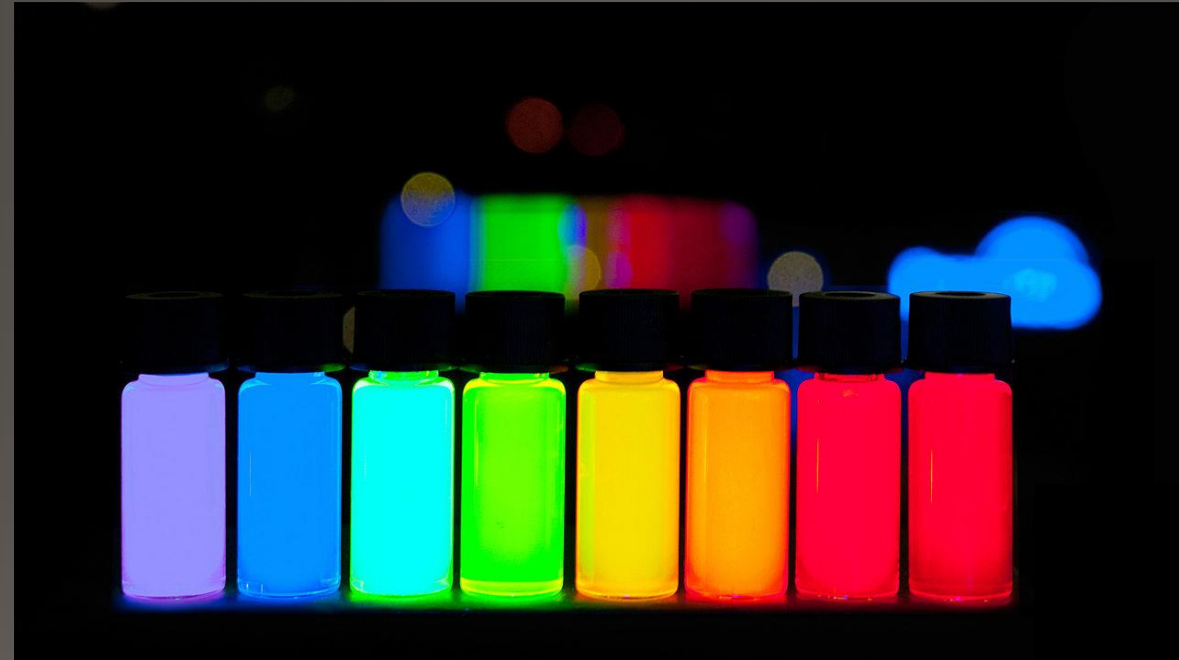
ϵ_r Dielectric constant

Future Targets: Quantum Dots

PlasmaChem ZnCdSeS C.C. 3.0



$a \sim 1 - 10 \text{ nm}$



Scalable, Cheap, tunable!

E.x. 60% by mass scintillators

Future Targets: Solvated Electrons



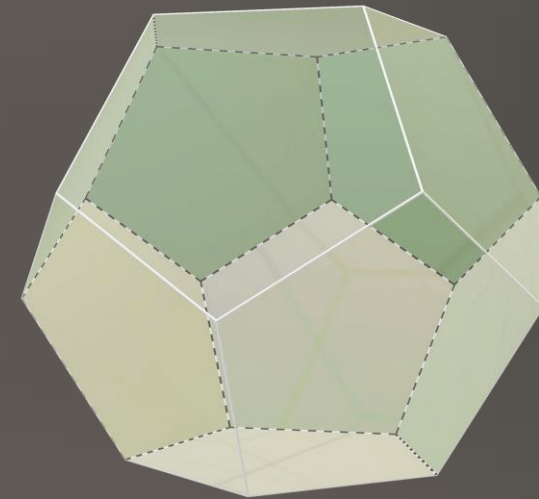
Purple color comes from eV-scale electronic transition at low conc.



Metallic color comes from quasi-conduction band at high conc.

LiNH₃ CC 3.0

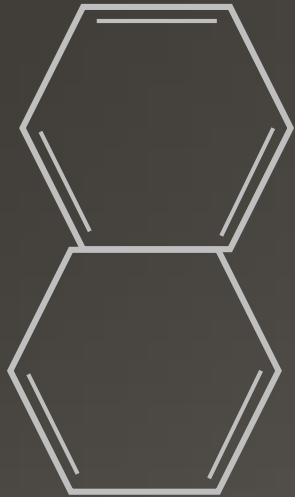
Solvation shell encapsulates a quasi-free electron in alkali metal solutions



$a \sim 0.1 \text{ nm}$

Electron acts like a particle in a cavity with electronic properties which are concentration dependent.

Future Considerations: Nuclear Scattering



Nuclear recoils can induce electronic excitations

- Molecular analog to Migdal effect
- Can be used to problem DM-Nucleon cross section
- Simple molecules and solvation shell are particularly well-suited
- Again, for aromatics, quantum yield near unity!

Future Prospects

- Explore scintillators with lower thresholds
- Crystal organic scintillators for directionality
- Expand formalism to more complicated molecules
- Expand formalism to include molecular deformations

Acknowledgements

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- Also thanks to Rouven Essig and Oren Slone for helpful conversations
- CB is supported by NSF-GRF under grants No. DGE-1144082 and DGE-1746045 .

Summary

- Covalent molecules have electronic transitions with $\Delta E \sim$ few eV
- Well characterized excitation spectra
- Experiment limited by threshold not statistics
- Short duration experiments
- Comparatively inexpensive
- Formalism being developed for DM interactions w/ covalent molecules
- Table-top experiments probing “high” energy BSM physics.