Solar surface convection and turbulent magnetism: new insights from SDO

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with Thierry Roudier, Alex Schekochihin, Michel Rieutord, Paul Barrère, Peter Haynes "Our equations for the Sun, for example, as a ball of hydrogen gas, describe a Sun without sunspots, without the rice-grain structure of the surface, without prominences, without coronas. Yet, all of these are really in the equations; we just haven't found the way to get them out."

R. P. Feynman lectures on physics, Chap. 41-12 (1964)

[Williams et al, ApJ 2020] Hi-C coronal imager

# Motivations

- The Sun is the closest astro system in a turbulent fluid MHD state
- Observable regions with different plasma  $\beta$  and B-field strengths
  - Solar convection zone (SCZ), corona
  - Active regions / quiet Sun
- Typical SCZ parameters
  - Re = Linj  $U_{\rm rms} / v \sim 10^{10} 10^{12}$
  - Rm = Linj  $U_{\rm rms} / \eta \sim 10^{6} 10^{10}$
  - $Pm = v / \eta \sim 10^{-6} 10^{-2}$
  - $E = v/(\Omega R^{2}) \sim 10^{-15}$
  - Ro ~  $U_{rms}$  / (L<sub>inj</sub>  $\Omega$ ) down to 10<sup>-1</sup>



- Best available astrophysical fluid dynamics lab
  - Turbulent convection, transport/diffusion, large/small-scale dynamo, rotation

#### La Palma 1m Swedish solar telescope

### Solar supergranulation problem

- Detected since 1954 as a flow pattern, using Doppler imaging
- Light intensity imaging only shows smaller-scale granulation
  - Physical origin of SG long-debated





Rincon & Rieutord, Living Rev. Sol. Phys. 2018

# The Solar Dynamics Observatory (SDO)

- SDO monitors the Sun 365d / 24h with three instruments
  - Helioseismic and Magnetic Imager (HMI) Photosphere
  - AIA, EVE Atmosphere, Corona
- HMI provides full-disk light intensity, Doppler and magnetic maps
  - 4096<sup>2</sup> pix: 45 seconds time-sampling with 350 km<sup>2</sup> resolution



• All high-resolution data is public

# The quiet photosphere

- First goal: study the statistically steady turbulent surface flow
  - Use quiet observation periods with as few active regions as possible
    - Oct. 15th, 2010 (24h), Nov. 26 Dec. 1 2018 (6 days uninterrupted)



#### Raw white-light intensity data



# Raw Doppler data

### Raw magnetic data





- Granulation:  $10^3$  km,  $\tau \sim 5$  min
- Supergranulation:  $3x10^4$  km,  $\tau \sim 24-48$  h

Relevant dynamics well resolved



# Velocity-field inference

- In-plane eulerian velocity field  $(u_x, u_y)$ 
  - Derived from Coherent Structure Tracking (CST) of granule motions
  - $\Delta X = 2500 \text{ km}, \Delta t = 15-30 \text{ min}$
- Out-of-plane Eulerian velocity field  $u_z$ 
  - Derived from Doppler measurements
  - Raw resolution  $\Delta X = 350$  km,  $\Delta t = 45$  s
- Reduction of final data
  - Harmonize resolution  $\Delta X = 2500$  km,  $\Delta t = 15-30$  min
  - Project on Gauss-Legendre-Fourier grid
  - *ℓ*-*v* filter (removal of 5-min *p*-modes)
  - Transform  $(u_x, u_y, u_z)$  to  $(U_r, U_\theta, U_{\varphi})$





### Local velocity field snapshot



Supergranulation:  $3x10^4$  km,  $\tau \sim 24-48$  h,  $u_h \sim 400$  m/s,  $u_v < 30$  m/s

### Spheroidal component (horizontal divergences)



### Toroidal component (vertical vorticity)



#### Velocity spectrum



Rincon et al, A&A 2017

# Main observational conclusions so far

- Surface flows are dominated by horizontal divergences
  - Strong correlation with upflows
- On scales larger than a few 1000 kms, the flow is anisotropic
  - $u_h \sim 400$  m/s,  $u_r < 30$  m/s at supergranulation scale
  - Typical vertical correlation scale is *H* ~ 2000-5000 km
- Spectral break suggests supergranulation is the largest driven scale at the surface

Can we make physical and theoretical sense of this ?

# Nonlinear evolution of large-aspect ratio convection simulations



# Energy budget/transfer

• Lin's equation





Rincon et al., A&A 2005



# Dependence of injection scale on size of surface entropy jump



Cossette & Rast, ApJ 2016

### Latest generation of global simulations





Hotta et al., ApJ 2014

# Supergranulation: emerging picture

- Supergranulation is the largest buoyancy-driven scale at the photospheric level
  - Supported by both observational and numerical analysis
- It appears to be the outcome of a nonlinear self-organization of turbulent thermal convection
  - Much more complex than thought for decades: a lesson for AFD ?
- Not entirely understood
  - Scale-dependence on convective-driving intensity (stellar luminosity)
  - Lack of strong thermal signature (radiative granulation boundary layer blanket, weak thermal flux at SG scales ?)

Rincon & Rieutord, Living Rev. Sol. Phys. 2018

# Turbulent convection phenomenology revisited

Dynamical equations for fluctuations

$$\begin{aligned} \frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} &= -\nabla p + \frac{\theta}{\Theta_0} g \, \mathbf{e}_z + \nu \Delta \mathbf{u} \\ \frac{\partial \theta}{\partial t} + \mathbf{u} \cdot \nabla \theta &= -u_z \nabla_z \, \Theta_0 + \kappa \Delta \theta \\ \nabla \cdot (\rho_0 \mathbf{u}) &= 0 \qquad \qquad \Theta = \Theta_0 + \theta \\ \mathbf{g} &= -q \, \mathbf{e}_z \end{aligned}$$

Derive evolution laws for statistics of fluctuation increments

$$\delta f(\mathbf{x}, \mathbf{r}) = f(\mathbf{x} + \mathbf{r}) - f(\mathbf{x})$$

### Isotropic, homogeneous theory

• Generalised Kolmogorov-Yaglom relations [Yakhot, PRL 1992]

$$\left\langle \delta u_r^3 \right\rangle = -\frac{4}{5} \varepsilon_u r + \frac{6}{r^4} \frac{g}{\Theta_0} \int_0^r y^4 \left\langle \delta \theta \delta u_z \right\rangle \mathrm{d}y + 6\nu \frac{\partial}{\partial r} \left\langle (\delta u_r)^2 \right\rangle$$
$$\left\langle (\delta \theta)^2 \delta u_r \right\rangle = -\frac{4}{3} \varepsilon_\theta r + \frac{2}{r^2} \int_0^r y^2 \left\langle \delta \theta \delta u_z \right\rangle \mathrm{d}y \frac{\partial \Theta_0}{\partial z} + 2\kappa \frac{\partial}{\partial r} \left\langle (\delta \theta)^2 \right\rangle$$

- Kolmogorov 41, passive scalar (constant fluxes):
  - $\delta u \sim r^{1/3}, \delta \theta \sim r^{1/3}$   $E(k) \sim k^{-5/3}, E_{\theta}(k) \sim k^{-5/3}$
- Bolgiano-Obukhov 59 (inertia/buoyancy balance + constant thermal flux)

$$\delta u \sim r^{3/5}, \delta \theta \sim r^{1/5}$$
  $E(k) \sim k^{-11/5}, E_{\theta}(k) \sim k^{-7/5}$ 

### Isotropic theory

• Bolgiano scale

$$L_B = \frac{\varepsilon_u^{5/4} \Theta_0^{3/2}}{\varepsilon_\theta^{3/4} g^{3/2}} \sim \frac{N u^{1/2} H}{(Ra \, Pr)^{1/4}}$$

• BO59 for *r* > *L*<sub>B</sub>

• K41 for *r* < *L*<sub>*B*</sub>

(within isotropic framework)

- BO59 never observed in aspect ratio O(1) situations because  $L_B$  is always O(H) [Rincon et al., JFM 2006, Kumar et al., PRE 2014]
- At the solar surface,  $L_B \sim H_{\rho} \sim 2000-5000 \text{ km}$ 
  - Transition takes place around granulation scale
- What happens for  $k_h H < 1$ ? (i.e. in our observational scale-range)
  - Anisotropic generalization of BO59 needed

# Tentative theory of large-scale, anisotropic turbulent convection

- Mass conservation  $\frac{\delta u_h}{\lambda_h} \sim \frac{\delta u_z}{H}$   $\lambda_z \sim H$
- Constant flux of thermal fluctuations
- Dominant inertia/buoyancy balance

$$\frac{\delta u_h \, \delta \theta^2}{\lambda_h} \sim \varepsilon_\theta = \text{const},$$

$$\frac{\delta u_h \, \delta u_z}{\lambda_h} \sim \left(\frac{H}{\lambda_h}\right)^2 g \frac{\delta \theta}{\Theta_0}$$

$$\delta u_{z} \sim \left(\frac{\varepsilon_{\theta}}{\Theta_{0}^{2}}\right)^{1/5} g^{2/5} H^{7/5} \lambda_{h}^{-4/5}$$

$$\delta u_{h} \sim \left(\frac{\varepsilon_{\theta}}{\Theta_{0}^{2}}\right)^{1/5} g^{2/5} H^{2/5} \lambda_{h}^{1/5}$$

$$E_{z}(k_{h}) \sim \left(\frac{\varepsilon_{\theta}}{\Theta_{0}^{2}}\right)^{2/5} g^{4/5} H^{14/5} k_{h}^{3/5}$$

$$E_{h}(k_{h}) \sim \left(\frac{\varepsilon_{\theta}}{\Theta_{0}^{2}}\right)^{2/5} g^{4/5} H^{4/5} k_{h}^{-7/5}$$

$$\delta \theta / \Theta_{0} \sim \left(\frac{\varepsilon_{\theta}}{\Theta_{0}^{2}}\right)^{2/5} g^{-1/5} H^{-1/5} \lambda_{h}^{2/5}$$

$$E_{\theta}(k_{h}) \sim \left(\frac{\varepsilon_{\theta}}{\Theta_{0}^{2}}\right)^{4/5} g^{2/5} H^{-2/5} k_{h}^{-9/5}$$

#### Theory vs. observations $E_z(k_h) \sim \left(\frac{\varepsilon_\theta}{\Theta_0^2}\right)^{2/5} g^{4/5} H^{14/5} k_h^{3/5} \qquad E_h(k_h) \sim \left(\frac{\varepsilon_\theta}{\Theta_0^2}\right)^{2/5} g^{4/5} H^{4/5} k_h^{-7/5}$ Doppler $E_r(\ell)$ $E_S(\ell)$ $E_T(\ell)$ $10^{4}$ Spectral energy $E(\ell)$ (m<sup>2</sup> s<sup>-2</sup>) $\propto \ell^2$ $\propto \ell^{-7/5} \\ \propto \ell^{-5/3}$ 10<sup>3</sup> $\propto \ell$ $\propto \ell^{3/5}$ $\propto \ell^{1/3}$ $10^{2}$ $10^{1}$ $\lambda_r \,(\mathrm{Mm})$ $10^{1}$ $10^{\circ}$ $10^{2} \ell$ $10^{3}$ $10^{2}$ $10^{1}$ $10^{3}$

Angular degree  $\ell$ 

Rincon et al., A&A 2017

# Signatures of self-similar buoyant dynamics: trees of fragmenting granules



Roudier et al., A&A 2003



# Lagrangian Coherent Structures

 Finite-Time Lyapunov Exponent field derived from separation of grid of evolving passive tracers



24h positive and negative-time FTLEs

Pros sea

Mordita, April 2020.

# Conclusions

- Unprecedented characterisation of strongly driven astrophysical fluid turbulence
  - Determination of full 3D velocity field in a plane, over almost two scale decades
  - High-resolution in time and space, up to global scales
  - Followed over several typical turnover times
- Observations, numerics paint a complicated nonlinear dynamical picture
  - Significant implications for the understanding of solar convection (supergranulation)
  - Motivates the development of new turbulence phenomenology
  - Promising preliminary results on turbulent transport
- More discovery/understanding potential
  - Dynamo/MHD: *α*-effect, MHD turbulence, low Pm small-scale dynamo?
  - Implications for stellar physics [and exoplanet detection, spectral noise problem]
  - Relevance to other astrophysical transport/turbulence problems (galaxies, disks, ICM)