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Six-meson amplitude

Bijnens & Husek, "Six-pion amplitude"

Phys.Rev.D, 2107.06291[hep-ph]

Bijnens, Husek & **Sjö**, "*Six-meson amplitude in QCD-like theories*" *Phys.Rev.D*, 2206.14212[hep-ph]



Hans Bijnens, Lund U.



Tomáš Husek, Lund U./Charles U.



Mattias Sjö, Lund U.

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Three-Pion Scattering: From the chiral Lagrangian to the lattice

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Strengths and Weaknesses of the Lattice

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Applicable to low-energy QCD

Non-perturbative

From first principles

Access to experimentally unobservable quantities

Computer-intensive

Euclidean (imaginary) time

Discrete space-time

Finite space-time volume

Often reqires heavier-than-physical masses

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Applicable to low-energy QCD

Non-perturbative

From first principles

Access to experimentally unobservable quantities

Synergy with ChPT

Computer-intensive

Euclidean (imaginary) time

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Often reqires heavier-than-physical masses

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Three-pion K-matrix

B.-B., Bijnens, Husek, R.-L., Sharpe & **Sjö**, *"The isospin-3 three-particle K-matrix at NLO in ChPT"* JHEP, 2303.13206[hep-ph]

B.-B., Bijnens, Husek, R.-L., Sharpe & **Sjö**, *"The three-particle K-matrix at NLO in ChPT for general isospin"* (in preparation, title preliminary)



Stephen Sharpe, U. of Washington



Fernando Romero-López, MIT



Jorge Baeza-Ballesteros, U. de València

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Vertices

X = LO vertex

$$\mathbf{X} = \mathsf{NLO}$$
 vertex

All the LO and NLO diagrams



Simplification & shortening by $\mathcal{O}(6!)$

Flavour structures

$$\begin{aligned} \mathcal{F}_{\{6\}}(a_1,\ldots,a_6) &= \langle t^{a_1}\cdots t^{a_6} \rangle \\ \mathcal{F}_{\{2,4\}}(a_1,\ldots,a_6) &= \langle t^{a_1}t^{a_2} \rangle \langle t^{a_3}\cdots t^{a_6} \rangle \\ \mathcal{F}_{\{3,3\}}(a_1,\ldots,a_6) &= \langle t^{a_1}t^{a_2}t^{a_3} \rangle \langle t^{a_4}t^{a_5}t^{a_6} \rangle \\ \mathcal{F}_{\{2,2,2\}}(a_1,\ldots,a_6) &= \langle t^{a_1}t^{a_2} \rangle \langle t^{a_3}t^{a_4} \rangle \langle t^{a_5}t^{a_6} \rangle \end{aligned}$$

$$\mathcal{M}ig(p_1,a_1;p_2,a_2;\ldots;p_6,a_6ig) = \sum_R \sum_\sigma \mathcal{M}_Rig(\sigma[p_1,\ldots,p_6]ig)\mathcal{F}_Rig(\sigma[a_1,\ldots,a_6]ig)$$

Flavor-ordering

 $\sigma \not\in \mathsf{symmetries of } \mathcal{F}_R \\ \rightarrow \mathsf{well}\mathsf{-}\mathsf{known}, \mathsf{unique}$

Deorbiting

$$\sigma \in \text{symmetries of } \mathcal{F}_R \ o \text{novel, non-unique!}$$

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Finite-volume energy shift

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Energy spectrum: n particles, box of size L

$$\det\left[F_n^{-1}(E, \boldsymbol{P}, L) + \mathcal{K}_n(E^*)\right] = 0$$

n = 2 particles

Lüscher, "Volume Dependence of the Energy Spectrum in Massive Quantum Field Theories" Commun. Math. Phys. (1986)

n = 3 particles

Hansen & Sharpe, "*Relativistic, model-independent, three-particle* quantization condition" Phys. Rev. D, 1408.5933[hep-lat] (and several other approaches)

Properties of $\mathcal{K}_{df,3}^{NLO}$

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3 particles, 2 scattering



Form of \mathcal{M}_2

 $egin{aligned} & o \mathcal{M}_2(oldsymbol{p})_{lm,l'm'} \ oldsymbol{p} &- ext{spectator momentum} \ Y^l_m(heta,\phi) &- ext{pair angular momentum} \end{aligned}$

Anatomy of $\mathcal{K}_{df,3}^{NLO}$

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3 particles, 3 scattering



Form of \mathcal{M}_3

$$ightarrow \mathcal{M}_3({m p},{m p}')_{lm,l'm'}$$

Particle exchange \Leftrightarrow Spectator choice

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On-shell propagator substitute



Properties of G

On-shell only - propagator-like near pole:

$$G(p, p')_{lm, l'm'} \sim rac{1}{(P - p - p')^2 - M_{\pi} + i\epsilon}$$

Smooth cutoff far from pole - non-analytic

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$\mathcal{K}_{df,3}$ at leading order

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One-particle exchange pole



One-particle exchange subtraction



$\mathcal{K}_{df,3}$ at next-to-leading order

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Bull's head cut



Bull's head subtraction



Different approaches



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Threshold expansion

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Expansion parameters

$$egin{aligned} \Delta \propto P^2 - (3M_\pi)^2 \ \Delta_i^{(\prime)} \propto (P - p_i^{(\prime)})^2 - (2M_\pi)^2 \ ilde{t}_{ij} \propto (p_i - p_j')^2 \end{aligned}$$

(system above-thr'ness)

(pair above-thr'ness)
(spectator above-thr'ness)

Compound parameters

$$\Delta_{\rm A} = \sum_i (\Delta_i^2 + \Delta_i'^2) - \Delta^2 \qquad \Delta_{\rm B} = \sum_{ij} \tilde{t}_{ij}^2 - \Delta^2$$

Threshold expansion

$$\mathcal{K}_{\text{df},3} = \mathcal{K}_0 + \mathcal{K}_1 \Delta + \mathcal{K}_2 \Delta^2 + \mathcal{K}_A \Delta_A + \mathcal{K}_B \Delta_B + \mathcal{O}(\Delta^3)$$

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Different approaches



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Different approaches



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Results!

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$$\begin{split} \mathcal{K}_{0} &= \left(\frac{M_{\pi}}{F_{\pi}}\right)^{4} 18 + \left(\frac{M_{\pi}}{F_{\pi}}\right)^{6} \left[-3\kappa(35 + 12\log 3) - \mathcal{D}_{0} + 111L + \ell_{(0)}^{r}\right], \\ \mathcal{K}_{1} &= \left(\frac{M_{\pi}}{F_{\pi}}\right)^{4} 27 + \left(\frac{M_{\pi}}{F_{\pi}}\right)^{6} \left[-\frac{\kappa}{20}(1999 + 1920\log 3) - \mathcal{D}_{1} + 384L + \ell_{(1)}^{r}\right], \\ \mathcal{K}_{2} &= \left(\frac{M_{\pi}}{F_{\pi}}\right)^{6} \left[\frac{207\kappa}{1400}(2923 - 420\log 3) - \mathcal{D}_{2} + 360L + \ell_{(2)}^{r}\right], \\ \mathcal{K}_{A} &= \left(\frac{M_{\pi}}{F_{\pi}}\right)^{6} \left[\frac{9\kappa}{560}(21809 - 1050\log 3) - \mathcal{D}_{A} - 9L + \ell_{(A)}^{r}\right], \\ \mathcal{K}_{B} &= \left(\frac{M_{\pi}}{F_{\pi}}\right)^{6} \left[\frac{27\kappa}{1400}(6698 - 245\log 3) - \mathcal{D}_{B} + 54L + \ell_{(B)}^{r}\right]. \end{split}$$

$$\begin{split} \mathcal{D}_0 &\approx -0.0563476589\,, \qquad \mathcal{D}_1 \approx 0.129589681\,, \qquad \mathcal{D}_2 \approx 0.432202370\,, \\ \mathcal{D}_A &\approx 9.07273890 \cdot 10^{-4}\,, \qquad \mathcal{D}_B \approx 1.62394747 \cdot 10^{-4}\,, \end{split}$$

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Reconciliation with the lattice





Phys.Rev.D, 2021.06144[hep-lat]

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* Blanton, Hanlon, Hörz, Morningstar, Romero-López & Sharpe,

"Three-body interactions from the finite-volume QCD spectrum"

** using LECs from FLAG and Colangelo, Gasser & Leutwyler, " $\pi\pi$ scattering"

Phys.Rev.D, 2021.06144[hep-lat] *Nucl.Phys.B*, hep-ph/0103088

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* Blanton, Hanlon, Hörz, Morningstar, Romero-López & Sharpe,

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omplete	Com	plete	
$\pi\pi\pi \to \pi\pi$	xt?	·ππ K -matrix	
ext: Remainin _{	$\pi\pi K$ πKK	ls	
I = 3	Singlet		
I=2	Doublet		
I = 1	Singlet	Doublet	
I = 0	Antisymm	etric singlet	

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Leading order





Osborn (1969) Susskind & Frye (1970)

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Reducible to smaller results





*Bijnens & Lu, "*Meson-meson scattering in QCD-like theories*" JHEP, 1102.0172[hep-ph]

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4-meson but with more legs





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New and difficult



One-Loop Integrals

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One-propagator integral

$$\bigvee^{\mathbf{k}} \sim \int \frac{\mathrm{d}^d \mathbf{k}}{(2\pi)^d} \frac{1}{\mathbf{k}^2 - M^2} \sim \frac{1}{4 - d} + \text{(finite)}$$

Two-propagator integral

$$q \xrightarrow{k} \sim \int \frac{\mathrm{d}^d \mathbf{k}}{(2\pi)^d} \frac{\{1, \mathbf{k}^\mu, \mathbf{k}^\mu \mathbf{k}^\nu\}}{(\mathbf{k}^2 - M^2) \left[(\mathbf{k} - q)^2 - M^2\right]}$$
$$\sim \frac{1}{4 - d} + \bar{\mathcal{I}}(q^2) + \text{(finite)}$$

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One-Loop Integrals

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Three-propagator integral

$$\int \frac{\mathrm{d}^{d} \mathbf{k}}{(2\pi)^{d}} \frac{\{1, \mathbf{k}^{\mu}, \mathbf{k}^{\mu} \mathbf{k}^{\nu}, \mathbf{k}^{\mu} \mathbf{k}^{\nu} \mathbf{k}^{\rho}\}}{(\mathbf{k}^{2} - M^{2}) \left[(\mathbf{k} - q_{1})^{2} - M^{2}\right] \left[(\mathbf{k} + q_{2})^{2} - M^{2}\right]}$$

In principle reducible to \overline{f} but impractical. Instead: elegant but redundant basis

$$\{C_0, C_1, C_2, C_3\}(p_1, \ldots, p_6)$$

where
$$C_3 \sim rac{1}{4-d} + ar{C}_3$$
 (others are finite)

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The bull's head

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How to integrate?







