

Topological spin excitations in non-Hermitian spin chains with a generalized kernel polynomial algorithm

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Physical Background

Dynamics of open quantum systems

$$H = H_s + H_e + H_{se}$$

$$\begin{aligned} & t_e \ll t_r \text{ (Born-Markov)} \\ & t_r \cdot \Delta \gg 1 \text{ (Rotating wave)} \end{aligned}$$

t_e : environment correlation time
 t_r : system relaxation time
 Δ : system energy level spacing

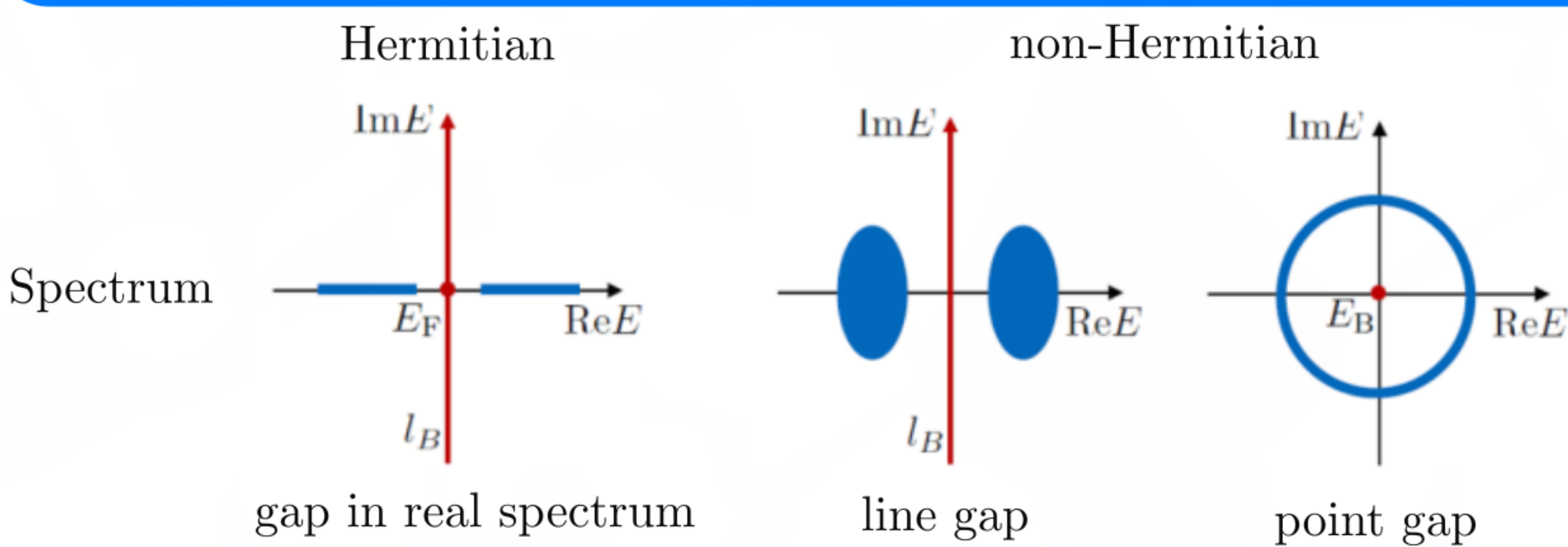
$$\frac{d}{dt}\rho_s = -i[\rho_s, H_s] + \sum_{\alpha} \left(L_{\alpha}\rho_s L_{\alpha}^{\dagger} - \frac{1}{2}\{L_{\alpha}^{\dagger}L_{\alpha}, \rho_s\} \right) := \mathcal{L}[\rho_s]$$

$$\text{vectorize } \rho_s \text{ into } |\tilde{\rho}_s\rangle: \frac{d}{dt}|\tilde{\rho}_s\rangle = \tilde{\mathcal{L}}|\tilde{\rho}_s\rangle$$

$$\text{no quantum jump: } H_{\text{eff}} = H_s - \frac{i}{2} \sum_{\alpha} L_{\alpha}^{\dagger} L_{\alpha}$$

Dynamics given by non-Hermitian matrices $\tilde{\mathcal{L}}$ and H_{eff}

(non-)Hermitian topology



Topology^[1] topological boundary modes

non-Hermitian skin effect

Challenging to characterize non-Hermitian many-body topology

Technical Background

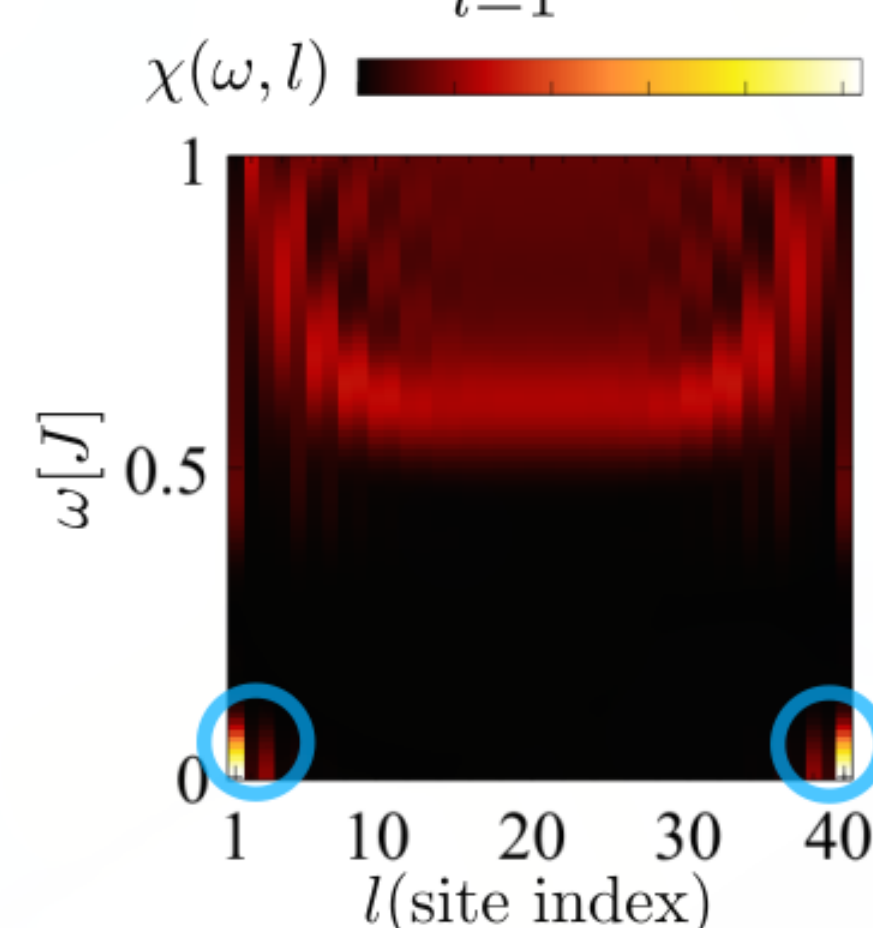
Spectral function for Hermitian many-body topology

$$H = \frac{J}{2} \sum_{l=1}^{20} (S_{2l-1}^x S_{2l}^x + S_{2l-1}^y S_{2l}^y) + J \sum_{l=1}^{19} (S_{2l}^x S_{2l+1}^x + S_{2l}^y S_{2l+1}^y) + J \sum_{l=1}^{39} S_l^z S_{l+1}^z$$

$$\chi(\omega, l) = \langle \text{GS} | S_l^z \delta(\omega - H + E_0) S_l^z | \text{GS} \rangle$$

- computed efficiently with matrix product states and kernel polynomial method

- reveals topological boundary modes



Kernel polynomial method (KPM)^[2]

- single variable function: $\rho(E) = \langle \psi_L | \delta(E - H) | \psi_R \rangle$
- Chebyshev expansion: $\rho(E) = \frac{1}{\pi\sqrt{1-E^2}} (\mu_0 + 2 \sum_{n=1}^{\infty} \mu_n T_n(E))$
- $T_n(E) = \cos(n \arccos(E))$ $T_{n+1}(x) = 2xT_n(x) - T_{n-1}(x)$
- recursively compute μ_n : $\mu_n = \langle \psi_L | T_n(H) | \psi_R \rangle = \langle \psi_L | v_n \rangle$
- $|v_0\rangle = |\psi_R\rangle$ $|v_{n+1}\rangle = 2H|v_n\rangle - |v_{n-1}\rangle$
- can be done with matrix product states representation for $|v_n\rangle$

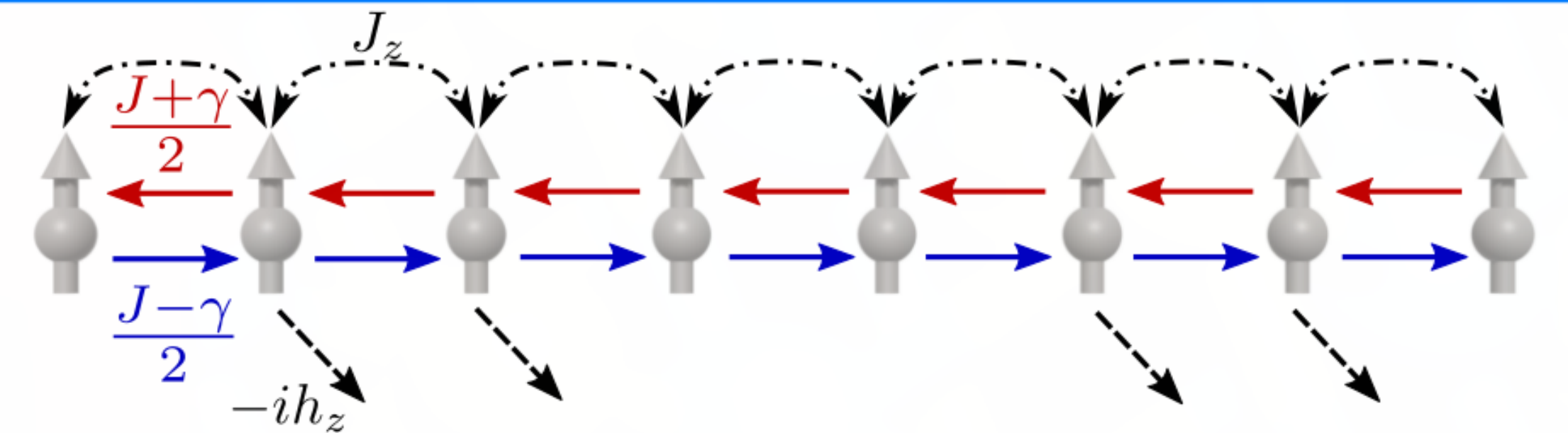
Spectral function for non-Hermitian many-body topology?

Bibliography

- [1] Y. Ashida, Z. Gong, and M. Ueda, Advances in Physics 69, 249 (2020).
[2] A. Weiße, G. Wellein, A. Alvermann, and H. Fehske, Rev. Mod. Phys. 78, 275 (2006)
[3] <https://github.com/GUANGZECHEN/NHKPM.jl>
[4] G. Chen, F. Song and J. L. Lado, arXiv:2208.06425

Our findings

Model



$$H = \sum_{l=1}^{L-1} \left(\frac{J+\gamma}{2} S_l^+ S_{l+1}^- + \frac{J-\gamma}{2} S_l^- S_{l+1}^+ + J_z S_l^z S_{l+1}^z \right) + \sum_{l=1}^L i h_l^z S_l^z$$

$$h_l^z = -h_z \text{ if } l \bmod 4 = 2, 3$$

$J_z = \gamma = 0$: topological for $h_z > 0$ Topology survive for $J_z > 0$?

non-Hermitian Kernel polynomial method (NHKPM)

- double variable function: $f(\omega) = \langle \psi_L | \delta^2(\omega - H) | \psi_R \rangle$

$$\delta^2(\omega - H) = \sum_n \delta(\Re(\omega - E_n)) \delta(\Im(\omega - E_n)) |\Psi_{R,n}\rangle \langle \Psi_{L,n}|$$

- no direct Chebyshev expansion, encounter with Hermitization:

$$\mathcal{H} = \begin{pmatrix} \omega - H & \\ \omega^* - H^\dagger & \end{pmatrix} \text{ s.t. } f(\omega) = \frac{1}{\pi} \partial_{\omega^*} G(E=0)$$

$$G(E) = \langle \text{L} | (E - \mathcal{H})^{-1} | \text{R} \rangle \quad | \text{L} \rangle = \begin{pmatrix} 0 \\ | \psi_L \rangle \end{pmatrix} \quad | \text{R} \rangle = \begin{pmatrix} | \psi_R \rangle \\ 0 \end{pmatrix}$$

- do Chebyshev expansion on $G(E)$ to efficiently compute $f(\omega)$ ^[3]

Spectral function for non-Hermitian many-body topology

$$\mathcal{S}(\omega, l) = \langle \text{GS}_L | S_l^- \delta^2(\omega + E_{\text{GS}} - H) S_l^+ | \text{GS}_R \rangle + \langle \text{GS}_L | S_l^+ \delta^2(\omega - E_{\text{GS}} - H) S_l^- | \text{GS}_R \rangle$$

- Faithfully revealing topological boundary states

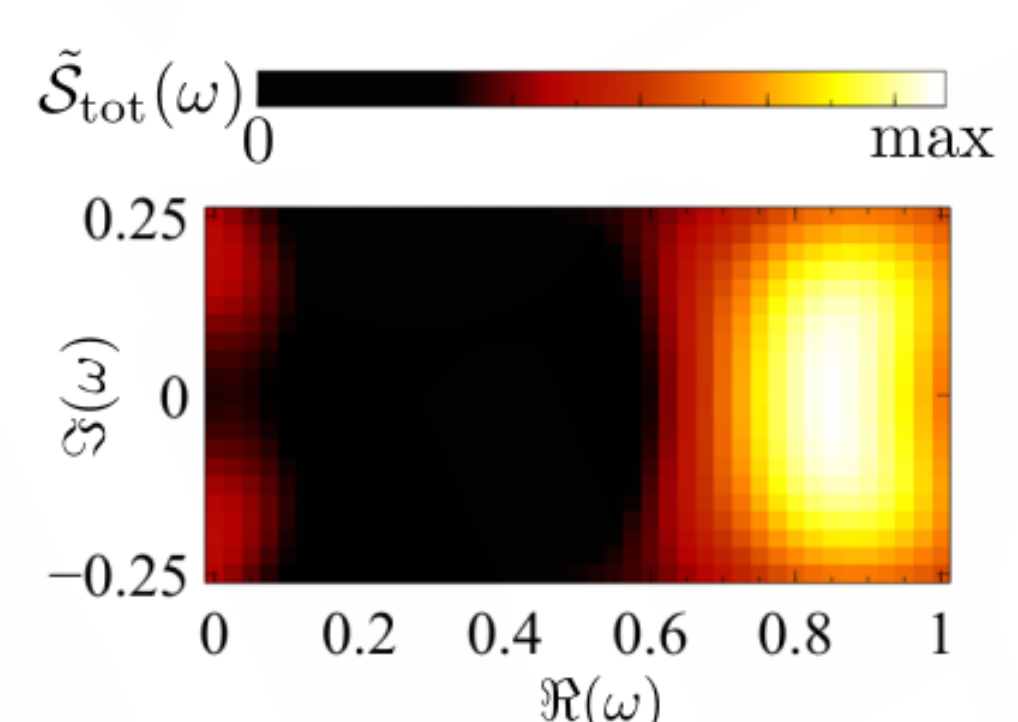
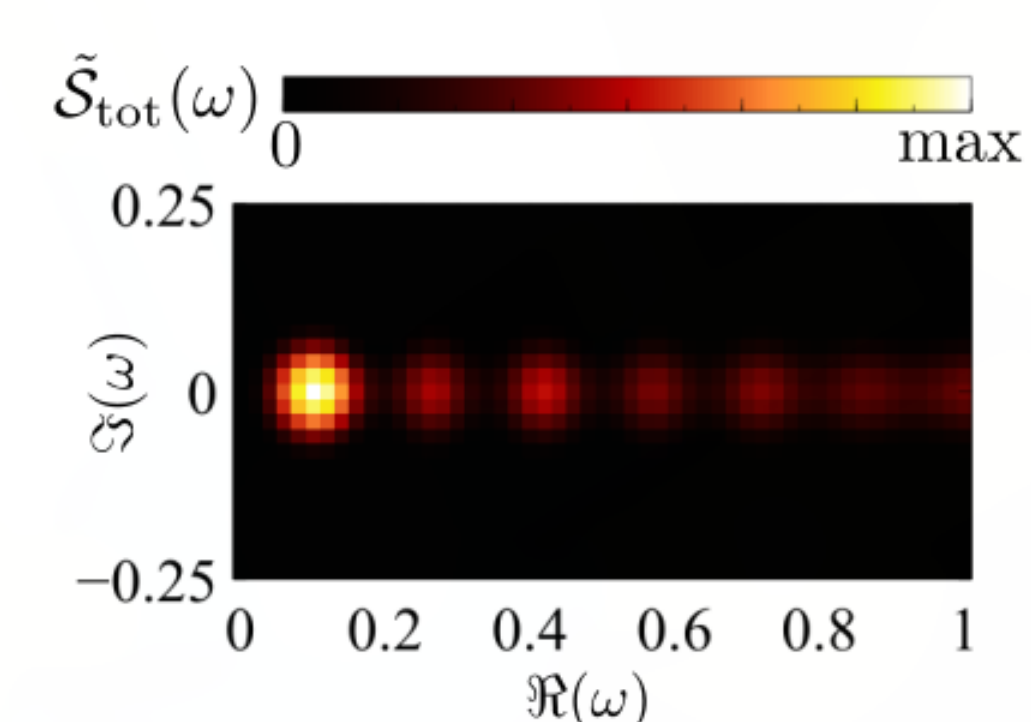
$$L = 24 \\ J_z = J/2, \gamma = 0$$

$h_z = 0$, trivial

$h_z = J$, topological

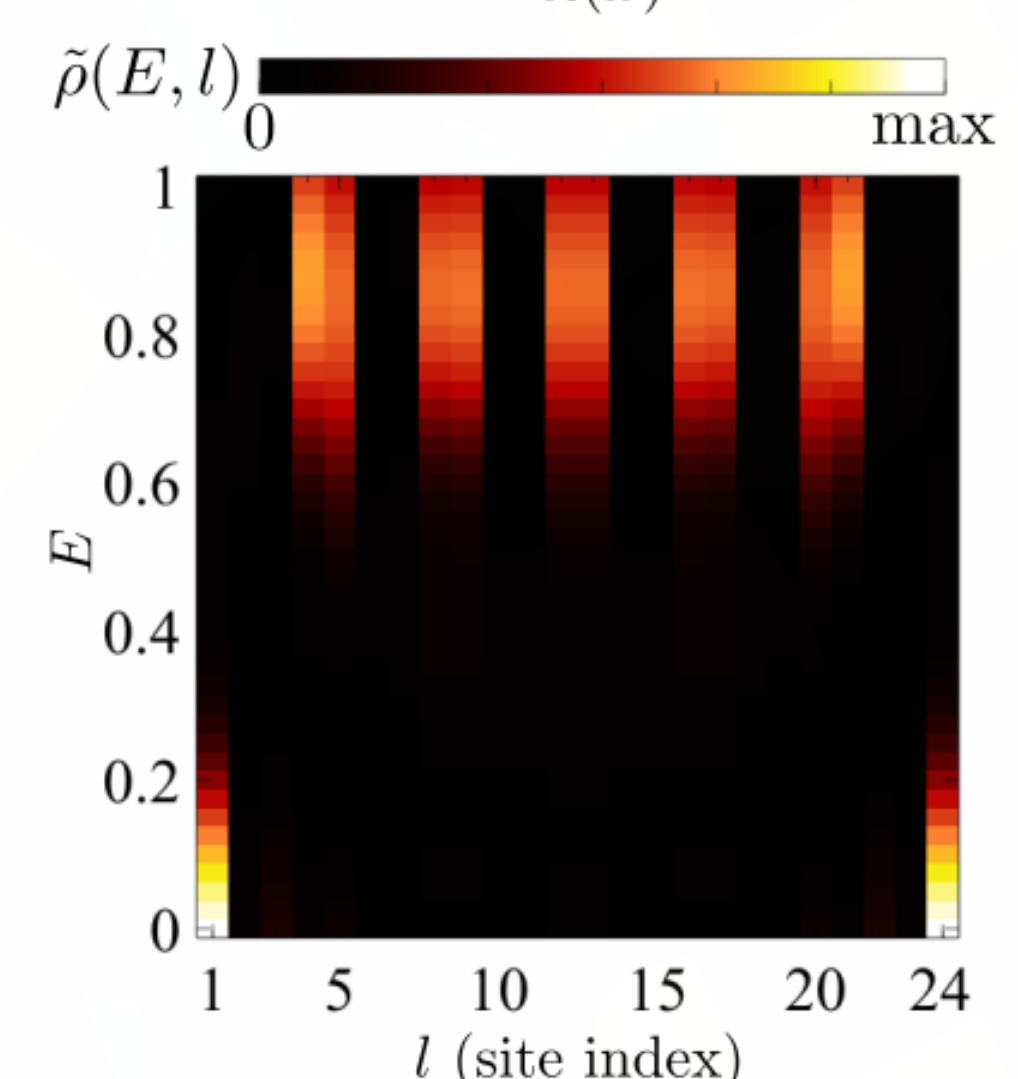
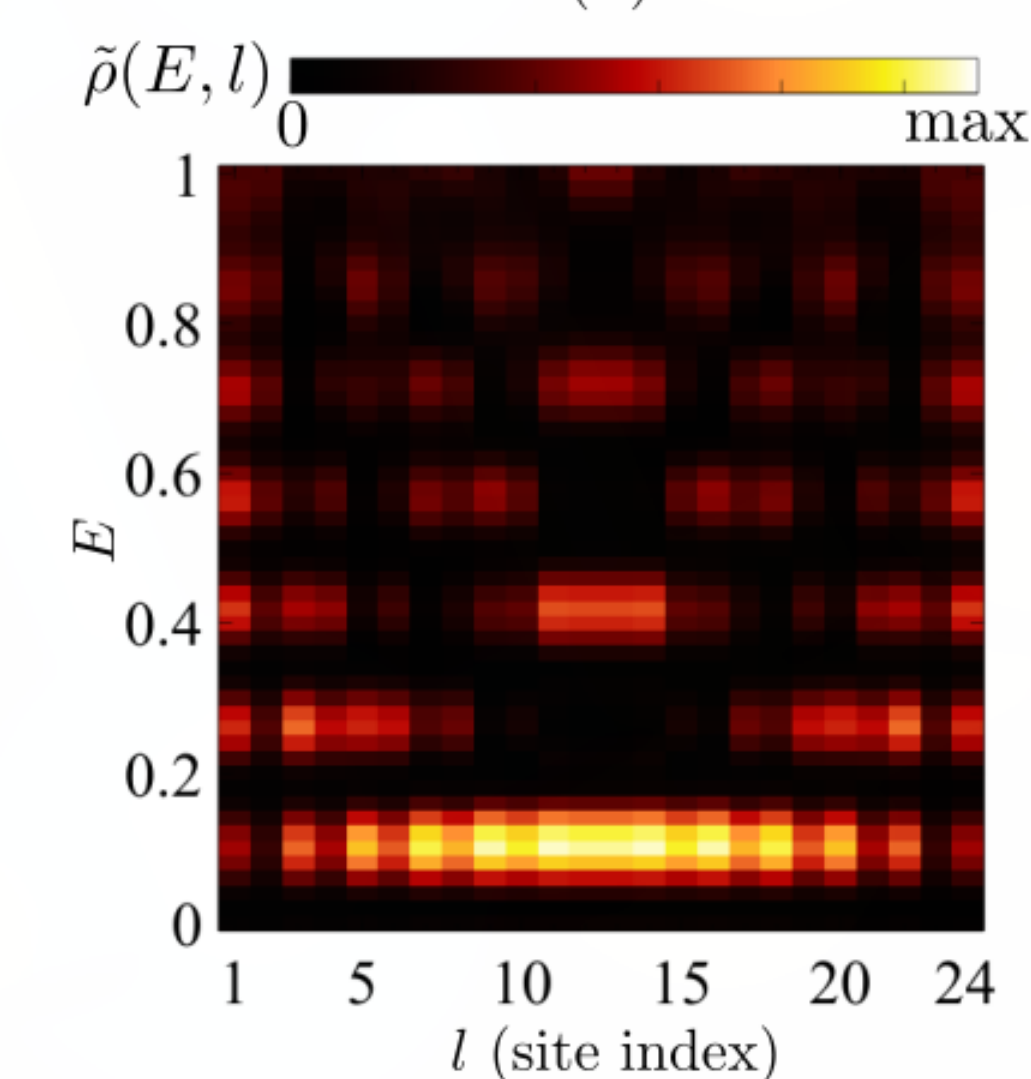
total spectrum:

$$\tilde{\mathcal{S}}_{\text{tot}}(\omega) = \left| \sum_{l=1}^L \mathcal{S}(\omega, l) \right|$$

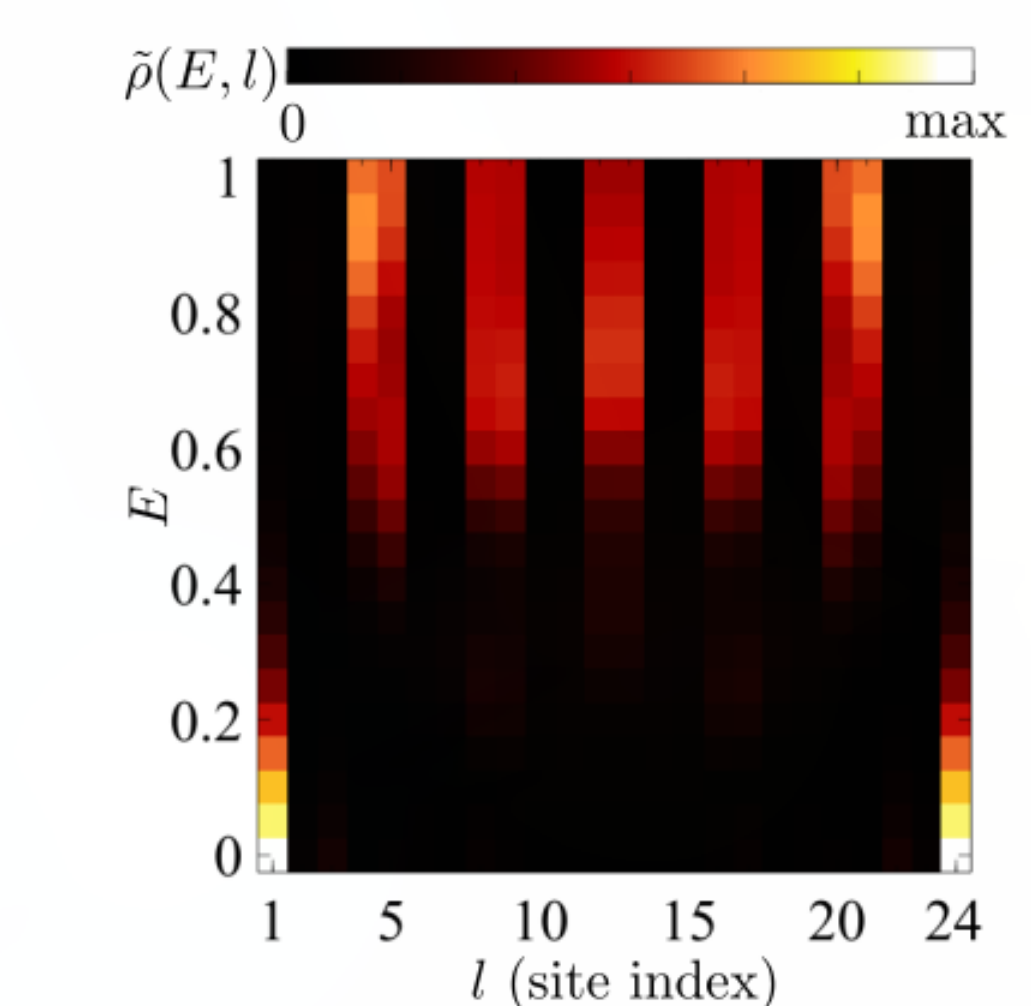


local spectrum:

$$\tilde{\rho}(E \in \mathbb{R}, l) = \left| \int \mathcal{S}(E + iy, l) dy \right|$$



- Stable in the presence of non-Hermitian skin effect



Take home

We developed a generalized kernel polynomial algorithm to efficiently compute spectral functions in non-Hermitian many-body systems, which is particularly useful for characterizing non-Hermitian many-body topology^[4].