

Celestial Soft Algebras

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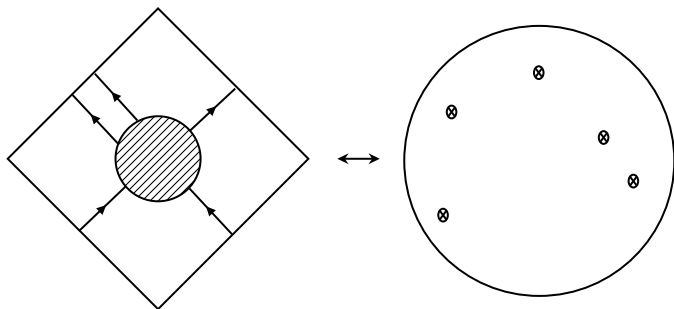
Outline

- Celestial Holography
- Celestial Soft Algebras for Gravity and Gauge Theory
- Dynamic and Central Extensions of Celestial Soft Algebras

Celestial Holography

Celestial Holography

Four-dimensional gravitational S -matrices have a holographic description as correlation functions on a 2D Euclidean sphere.



Celestial Holography

- We want examples where there is an independently defined 2D theory that computes 4D scattering amplitudes.
- In the absence of an independent boundary description, we would like to identify universal aspects of these dual theories starting from properties of gauge and gravitational amplitudes.
- We focus on a set of higher-spin currents that are present in the celestial dual due to universal low-energy features of gauge and gravitational amplitudes.

Celestial Holography

Conformal Primary Wavefunction

A conformal primary wavefunction is a solution to the Klein-Gordon equation

$$(\partial^\mu \partial_\mu - m^2) \phi_{\Delta, m}(x^\mu | z, \bar{z}) \quad (1)$$

such that under Lorentz transformations

$$\phi_{\Delta, m}(\Lambda^\mu{}_\nu x^\nu | \frac{az + b}{cz + d}, \frac{\bar{a}\bar{z} + \bar{b}}{\bar{c}\bar{z} + \bar{d}}) = |cw + d|^{2\Delta} \phi_{\Delta, m}(x^\mu | z, \bar{z}) \quad (2)$$

^a

^aPasterski, S., Shao, S.H., Strominger, A. Flat space amplitudes and conformal symmetry of the celestial sphere.

Celestial Holography

Massless Conformal Primary Wavefunction

$$\begin{aligned}\phi_{\Delta}^{\pm}(x|z, \bar{z}) &= \int d\omega \omega^{\Delta-1} e^{\pm i\omega \hat{q}(z, \bar{z}) \cdot x} \\ \hat{q}^{\mu}(z, \bar{z}) &= (1 + z\bar{z}, z + \bar{z}, -i(z - \bar{z}), 1 - z\bar{z})\end{aligned}\tag{3}$$

Celestial Amplitude

$$\begin{aligned}\mathcal{A}(1_{\Delta_1, J_1, z, \bar{z}}^{\epsilon_1, a_1} \cdots n_{\Delta_n, J_n, z_n, \bar{z}_n}^{\epsilon_n, a_n}) &= \int \prod_j d\omega_j \omega_j^{\Delta_j-1} A^{a_1 \cdots a_n}(\epsilon_1 \omega_1 q(z_1, \bar{z}_1), \dots) \\ &= \langle \mathcal{O}_{\Delta_1, J_1}^{\epsilon_1, a_1}(z_1, \bar{z}_1) \cdots \mathcal{O}_{\Delta_n, J_n}^{\epsilon_n, a_n}(z_n, \bar{z}_n) \rangle\end{aligned}\tag{4}$$

Conformally Soft Theorems¹

Energetically Soft Theorems

$$\lim_{\omega \rightarrow 0} \omega \langle \text{out} | a_+(\omega \hat{q}) \mathcal{S} | \text{in} \rangle = e \sum_k \frac{Q_k \cdot p_k \cdot \epsilon^+}{p_k \cdot \hat{q}} \langle \text{out} | \mathcal{S} | \text{in} \rangle + O(\omega) \quad (5)$$

Conformally Soft Theorems

$$\lim_{\Delta_1 \rightarrow 1} (\Delta_1 - 1) \mathcal{A}(1_{\Delta_1}^+, 2_{\Delta_2}^{\epsilon_2}, \dots) = e \sum_k \frac{Q_k}{z_1 - z_k} \mathcal{A}(2_{\Delta_2}^{\epsilon_2}, \dots) \quad (6)$$

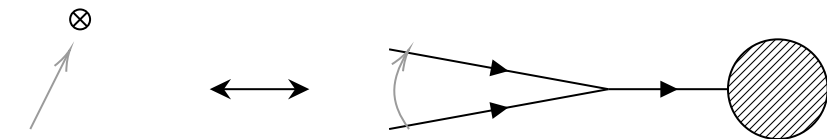
¹Pate, M., Raclariu, A.M., Strominger, A. Conformally soft theorem in gauge theory. Phys. Rev. D 100, (2019).

Conformally Soft Theorems

Conformally Soft Gluons and Gravitons

$$\begin{aligned} R^{k,a} &= \lim_{\epsilon \rightarrow 0} \epsilon \mathcal{O}_{k+\epsilon}^{+a}(z, \bar{z}), \quad k = 1, 0, \dots \\ H^\ell &= \lim_{\epsilon \rightarrow 0} \epsilon \mathcal{G}_{\ell+\epsilon}^+(z, \bar{z}), \quad k = 2, 1, 0, -1, \dots \end{aligned} \tag{7}$$

Celestial OPE



Operators becoming coincident \leftrightarrow Particles becoming collinear

Celestial OPE²

Hard OPE

$$\begin{aligned} \mathcal{O}_{\Delta_1}^{+a}(z_1, \bar{z}_1) \mathcal{O}_{\Delta_2}^{+b}(z_2, \bar{z}_2) &\sim -\frac{if^{ab}}{z_{12}} \sum_{m=0}^{\infty} B(\Delta_1 + m - 1, \Delta_2 - 1) \frac{\bar{z}_{12}^m}{m!} \\ &\quad \bar{\partial}^m \mathcal{O}_{\Delta_1 + \Delta_2 - 1}^{+c}(z_2, \bar{z}_2) + \mathcal{O}(z_{12}^0) \\ \mathcal{G}_{\Delta_1}^+(z_1, \bar{z}_1) \mathcal{G}_{\Delta_2}^+(z_2, \bar{z}_2) &\sim -\frac{\kappa \bar{z}_{12}}{2 z_{12}} \sum_{m=0}^{\infty} B(\Delta_1 + m - 1, \Delta_2 - 1) \frac{\bar{z}_{12}^m}{m!} \\ &\quad \times \bar{\partial}^m \mathcal{G}_{\Delta_1 + \Delta_2}^+(z_2, \bar{z}_2) + \mathcal{O}(z_{12}^0) \end{aligned} \tag{8}$$

²Pate, M., Raclariu, A.M., Strominger, A., Yuan, E.Y. Celestial Operator Products of Gluons and Gravitons. *Reviews in Mathematical Physics*, 33, no 9 (2021).

Celestial OPE

Soft OPE

$$R^{k,a} R^{\ell,b} \sim -\frac{if^{ab}_c}{z_{12}} \sum_{j=0}^{1-k} \frac{(2-k-\ell-j)!}{(1-k-j)!(1-\ell)!} \frac{\bar{z}_{12}^j}{j!} \bar{\partial}^j R^{k+\ell-1,c}$$
$$H^k H^\ell \sim -\frac{\kappa}{2} \frac{\bar{z}_{12}}{z_{12}} \sum_{j=0}^{1-k} \frac{(2-k-\ell-j)!}{(1-k-j)!(1-\ell)!} \frac{\bar{z}_{12}^j}{j!} \bar{\partial}^j H^{k+\ell}$$
(9)

Celestial Soft Algebras³

Modes of Soft Operators

$$R^{3-2q,a} = \sum_n \sum_{m=1-q}^{q-1} \frac{S_{n,m}^{q,a}}{(q+m-1)!(q-m-1)!z^{n+2-q}\bar{z}^{m+1-q}} \quad (10)$$
$$H^{4-2q} = \sum_n \sum_{m=1-q}^{q-1} \frac{\kappa W_{n,m}^q}{(q+m-1)!(q-m-1)!z^{n+3-q}\bar{z}^{m+1-q}}$$

Gauge and Gravitational Soft Algebras

$$[S_{n,m}^{q,a}, S_{n',m'}^{p,b}] = -if_c^{ab} S_{n+n',m+m'}^{p+q-1,c} \quad (11)$$
$$[W_{n,m}^p, W_{n',m'}^q] = [m(q-1) - m'(p-1)] W_{n+n',m+m'}^{p+q-2}$$

³Melton, W., Narayanan, S.A., Strominger, A. Deforming soft algebras for gauge theory. J. High Energy. Phys. 2023, 233 (2023).

Extensions of Gauge Soft Algebras

Yang-Mills and Neutral Scalar Theory

$$S = \int d^4x \left(-\frac{1}{4} \text{Tr} F^2 - \partial^\mu \bar{\phi} \partial_\mu \phi - \frac{\mu}{4} (\phi \text{Tr} F_+^2 + \bar{\phi} \text{Tr} F_-^2) \right) \quad (12)$$

Extended Soft Algebra

$$\begin{aligned} [S_{n,m}^{q,a}, S_{n',m'}^{r,b}] &= -if^{ab}{}_c S_{n+n',m+m'}^{q+r-1,c} \\ &\quad - \mu \delta^{ab} (m'(q-1) - m(r-1)) \bar{\sigma}_{n+n',m+m'}^{q+r-2} \\ [S_{n,m}^{q,a}, \sigma_{n',m'}^r] &= -\mu (m'(q-1) - m(r-1)) \bar{s}_{n+n',m+m'}^{q+r-2,a} \end{aligned} \quad (13)$$
$$\lim_{\epsilon \rightarrow 0} \epsilon \bar{\Phi}_{2-2q+\epsilon}^+ = \sum_{m=1-q}^{q-1} \frac{\bar{\sigma}_{n,m}^q}{(q+m-1)!(q-m-1)! z^{n+1-q} \bar{z}^{m+1-q}}$$

Central Terms From Sourced Backgrounds

Soft Amplitudes in a Background

$$\begin{aligned}\delta S &= -\frac{\mu}{4} \int d^4x \phi_0(x) \text{Tr } F^2 \\ \phi_0(p^2) &= \sum_j a_j (p^2)^j \\ \langle R^{k,a}(z_1, \bar{z}_1) R^{j,b}(z_2, \bar{z}_2) \rangle &= \frac{2\mu \delta^{kj} \delta^{ab} a_{-1-k}}{z_{12}^2 |z_{12}|^{2(k-1)}}\end{aligned}\tag{14}$$

Central Extensions to Gauge Algebra

$$\begin{aligned}[S_{n,m}^{q,a}, S_{n',m'}^{r,b}] &= -if^{ab} c S_{n+n',m+m'}^{q+r-1,c} \\ &\quad + (c_1 \delta^{q1} n + c_{3/2} \delta^{q3/2} m) \delta^{pq} \delta_{n+n',0} \delta_{m+m',0} \\ c_1 &\propto a_{-2}, \quad c_{3/2} \propto a_{-1}\end{aligned}\tag{15}$$

Soft Algebras and Form Factors

- For certain local theories on twistor space, it is known that *conformal blocks* of the soft algebra are in a one-to-one correspondence with local operators of the theory.
- A conformal block is a set of correlation functions consistent with the OPE of the soft algebra.
- The conformal block gives bulk scattering amplitudes in the presence of a local operator in the bulk.⁴

⁴Costello, K., Paquette, N.M. Celestial holography meets twisted holography: 4d amplitudes from chiral correlators. J. High Energ. Phys. 2022, 193 (2022).

Sourced Backgrounds as Conformal Blocks

$$\int [D\phi][DA] e^{-S} (1 + \alpha\phi(0)) \mathcal{O}(\phi, A) \approx \int [D\phi][DA] e^{-S + \alpha\phi(0)} \mathcal{O}(\phi, A) \quad (16)$$

$$\square\phi_0(x) = \alpha\delta^{(4)}(x) \implies \phi_0(p) = \frac{\alpha}{p^2} \quad (17)$$

$$\langle R^{k,a} R^{\ell,b} \rangle = \frac{2\mu\alpha\delta^{k0}\delta^{\ell 0}\bar{z}_{12}}{z_{12}} \quad (18)$$

Sourced Backgrounds as Conformal Blocks

- For YM coupled to a canonical scalar, no operator deforms leading soft theorem.
- For YM coupled to scalar with higher-order kinetic term, ϕF^2 can deform the leading soft theorem.
- $\phi_0(p) \propto 1/p^4$ gives conformal block for $\phi(0)$:

$$\langle R^{k,a} R^{\ell,b} \rangle = \frac{2\mu\alpha\delta^{k1}\delta^{\ell1}}{z_{12}^2} \quad (19)$$

Conclusions

- Celestial amplitudes organize a tower of soft theorems into a simple algebra.
- Near universality of conformally soft theorems implies that these algebras should be present in explicit examples of celestial holography.
- Scattering amplitudes in the presence of sources or local operators inserted in the bulk are naturally described by conformal blocks of the soft algebra.