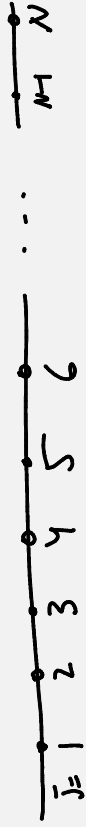


Advantages in Lattice Fermions

For N sites, have 2^N states



N is even so the lattice is bipartite

$$\{c_j^{\dagger}, c_j^{\dagger}\} = \delta_{ij}; \quad \{c_j, c_j^{\dagger}\} = 0,$$

One site $j=1$, $c_1|0\rangle = 0$, $|1\rangle = c_1^{\dagger}|0\rangle$, 2-state system

$$c_1 = \sigma_1^{-}; \quad c_1^{\dagger} = \sigma_1^{+}; \quad \sigma^{\pm} = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} = \frac{1}{2} (\sigma^x \pm i\sigma^y)$$

$$q_1 = \frac{1}{2} [c_1^{\dagger}, c_1] = \frac{1}{2} \sigma_1^z,$$

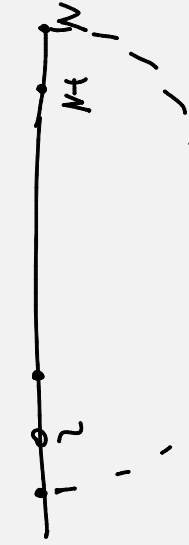
Flipping Hamiltonian: \downarrow PBC $c_{N+1} = c_1$;

$$H_t = -t \sum_{j=1}^{N-1} (c_j^\dagger c_{j+1} + c_{j+1}^\dagger c_j) - t (c_N^\dagger c_1 + c_1^\dagger c_N)$$

$t > 0$; Can flip the sign of t by

$$c_j \rightarrow (-1)^j c_j$$

Either open or closed



Fourier basis

$$c_j = \frac{1}{\sqrt{N}} \sum_k e^{ikj} c_k$$

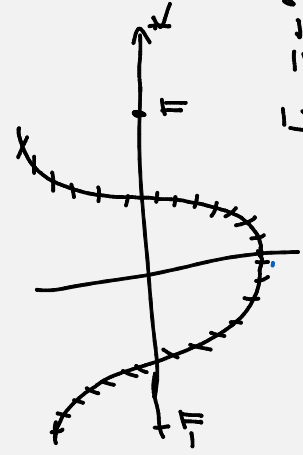
$$k \in \frac{2\pi n}{N}; \quad k \sim k + 2\pi; \quad k = -\frac{N}{2} + 1, -\frac{N}{2} + 2, \dots, 0, \dots, \frac{N}{2}$$

Jordan-Wigner transform.

$$H = -2t \sum_k (\cos k) c_k^\dagger c_k, \quad c_j = \left(\prod_{n=1}^{j-1} \sigma_n^z \right) \sigma_j^-, \quad \{c_j, c_k\} = 0;$$

$$c_j^\dagger = \left(\prod_{n=1}^{j-1} \sigma_n^z \right) \sigma_j^+;$$

Ground State



$$\prod_{k \in [-\frac{\pi}{2}, \frac{\pi}{2}]} c_k^\dagger |0\rangle; \quad \text{As } N \rightarrow \infty$$

$$\Delta k = \frac{2\pi}{N} \Delta n$$

Lieb, Schultz, Mattis

$$E_0 = -2t \frac{N}{2\pi} \int_{-\pi/2}^{\pi/2} \cos k dk = -\frac{2t}{\pi} N;$$

Xy model;

Acts on 2^N states of

$$N \text{ qubits; } [\sigma_j^+, \sigma_{j+2}^-] = 0;$$

$$H = -\frac{t}{2} \sum_j (\sigma_j^x \sigma_{j+1}^x + \sigma_j^y \sigma_{j+1}^y) = -t \sum_j (\sigma_j^+ \sigma_{j+1}^- + \sigma_j^- \sigma_{j+1}^+)$$

Hubbard chain

$$(c_{j\uparrow}, c_{j\downarrow}) \sim c_{j\sigma}$$

$$\epsilon = (\uparrow, \downarrow)$$



$$c_{j,\uparrow} \rightarrow i) c_{j,\sigma}$$

change var to Im hopping parameter. $SU(2)$ spin is manifest, also has hidden $SU(2)$ pseudospin

$$H = i t \sum_j (c_{j+1,\sigma}^\dagger - c_{j,\sigma}^\dagger)(c_{j,\sigma} - c_{j+1,\sigma}) = t \sum_{j,j+1} T_{j,j+1} + U \sum_j n_{j\uparrow} n_{j\downarrow}$$

Free staggered
lattice Dirac
fermion

$$R \begin{pmatrix} c_1 \\ c_2 \\ \vdots \\ c_N \end{pmatrix}; R \in SO(N); R^T R = I_{N \times N}$$

"Scans are $SO(N)$ singlets"

In 1d: Yang

$(-1)^{jx+jy}$ reversed by

$$c_{jxjy} \rightarrow (-1)^{jx+jy} c_{jxjy}$$

$$c_{j\alpha} c_{j\beta} \in \alpha\beta$$

$$H = \sum_j c_{j\uparrow} c_{j\downarrow} = \frac{1}{2} \sum_j (c_{j\uparrow}^2 - c_{j\downarrow}^2)$$

Yang's star-pairing states

$$(h^+)^k |0\rangle; \quad k=0, \dots, N$$

Pseudospin $\frac{N}{2}$ multiplet

Pseudospin generators:

$$SU(2)_{\text{pseudo}} \left\{ \begin{array}{l} h^+ \\ H = \frac{1}{2} \sum_j [c_{j\uparrow}^2 - c_{j\downarrow}^2] \end{array} \right.$$

$$[h^+, H] = 0$$

:

